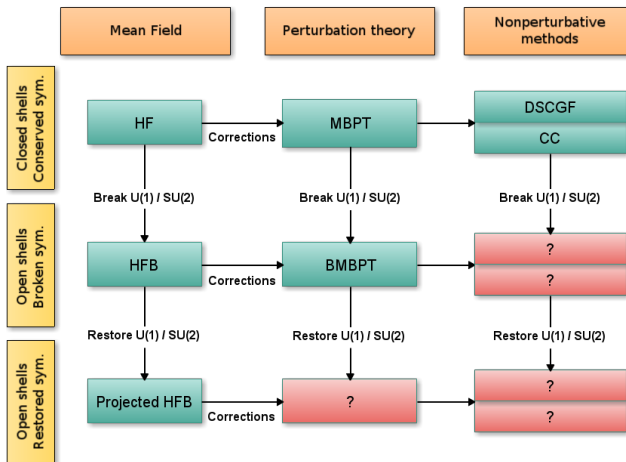


# Bogoliubov Many-Body Perturbation Theory for Open-Shell Nuclei

Pierre Arthuis  
IRFU, CEA, Université Paris - Saclay

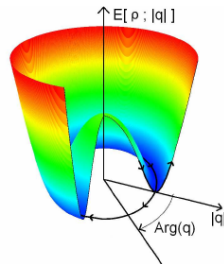
with T. Duguet (CEA Saclay), J.-P. Ebran (CEA DAM), H. Hergert (MSU),  
R. Roth (TU Darmstadt) & A. Tichai (ESNT, CEA Saclay)

Workshop MBPT in modern quantum chemistry and nuclear physics  
CEA Saclay - March 29th 2018

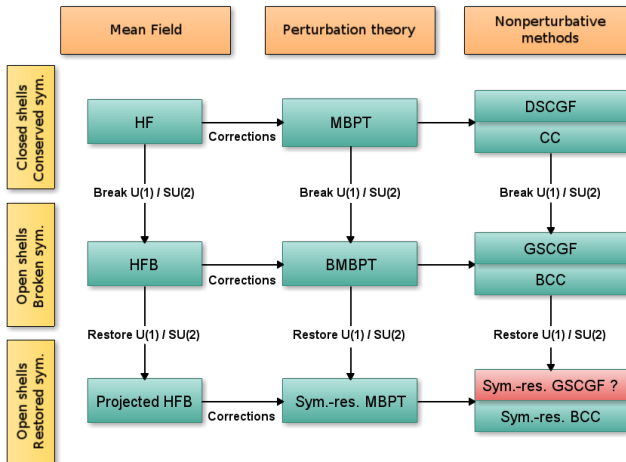


Expansion methods around unperturbed product state

- Symmetry breaking helps incorporating non-dynamical correlations:
  - ◇ Superfluid character:  $U(1)$  (particle number)
  - ◇ Deformations:  $SU(2)$  (angular momentum)
- But nuclei carry good quantum numbers (e.g. number of particles)  
 $\Rightarrow$  Symmetries must eventually be restored

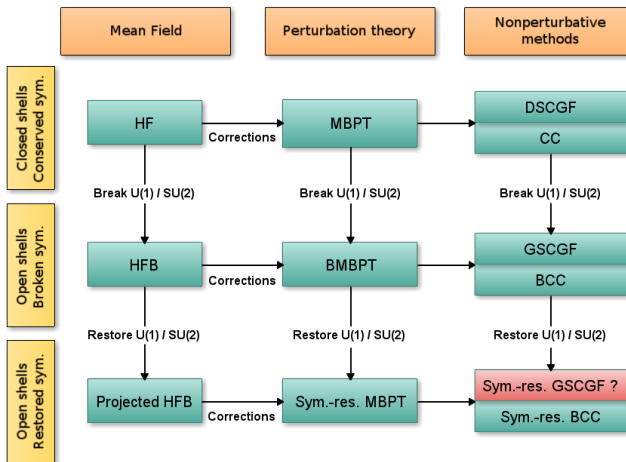


- See Thomas' talk on Monday



New methods recently proposed and implemented

- GSCGF, BCC [Somà *et al.* 2011, Signoracci *et al.* 2014]
- Sym.-res. BCC & sym.-res. BMBPT [Duguet 2015, Duguet & Signoracci 2017, Qiu *et al.* 2017]



MBPT reimplemented using SRG-evolved H in closed shell [Tichai *et al.* 2016] → Robert's talk

➡ MBPT competes with non-perturbative methods

Current objective: extend to (symmetry-projected) BMBPT for open shell

## Particle-number-restored BMBPT formalism

Exact diagrammatic expansion with symmetry breaking *and* restoration

[Duguet and Signoracci, *J. Phys. G* **44**, 2017] → Thomas' talk on Monday



## Formalism actualization

Expand off-diagonal kernels

$$\langle \Psi_0^A | H | \Phi(\phi) \rangle$$

$$\langle \Psi_0^A | \Phi(\phi) \rangle$$

Symmetry restoration

Diagonal reduction

$$\langle \Psi_0^A | H | \Phi \rangle$$

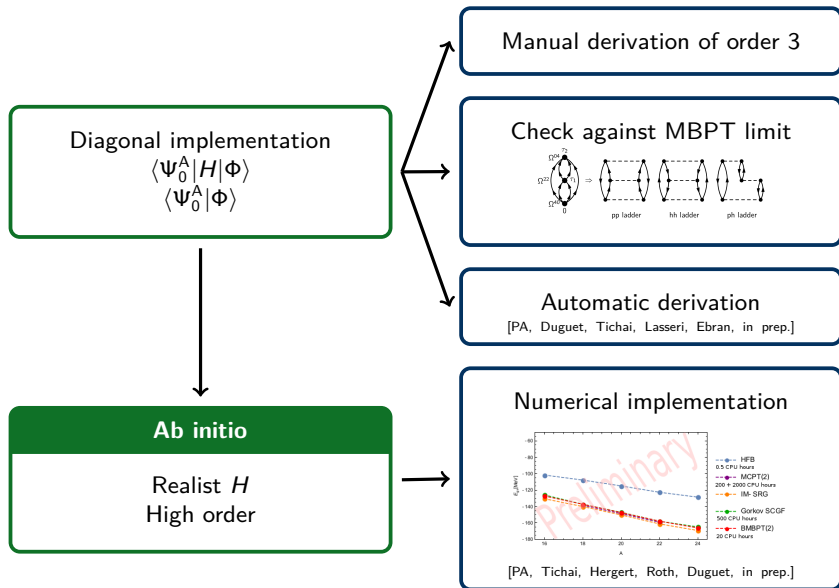
$$\langle \Psi_0^A | \Phi \rangle$$

No symmetry restoration



Ab initio

Realist  $H$   
High order



- Bogoliubov vacuum  $|\Phi\rangle$ ,  $\beta_k|\Phi\rangle = 0 \forall k$  with

$$\beta_k = \sum_p U_{pk}^* c_p + V_{pk}^* c_p^\dagger$$

$$\beta_k^\dagger = \sum_p U_{pk} c_p^\dagger + V_{pk} c_p$$

- Particle number symmetry broken:  $A|\Phi\rangle \neq A|\Phi\rangle$
- Grand potential  $\Omega \equiv H - \lambda A$  in qp basis, normal-ordered w.r.t.  $|\Phi\rangle$

$$\begin{aligned} \Omega = & \Omega^{00} + \frac{1}{1!} \sum_{k_1 k_2} \Omega_{k_1 k_2}^{11} \beta_{k_1}^\dagger \beta_{k_2} + \frac{1}{2!} \sum_{k_1 k_2} \left\{ \Omega_{k_1 k_2}^{20} \beta_{k_1}^\dagger \beta_{k_2}^\dagger + \Omega_{k_1 k_2}^{02} \beta_{k_2} \beta_{k_1} \right\} \\ & + \frac{1}{(2!)^2} \sum_{k_1 k_2 k_3 k_4} \Omega_{k_1 k_2 k_3 k_4}^{22} \beta_{k_1}^\dagger \beta_{k_2}^\dagger \beta_{k_4} \beta_{k_3} \\ & + \frac{1}{3!} \sum_{k_1 k_2 k_3 k_4} \left\{ \Omega_{k_1 k_2 k_3 k_4}^{31} \beta_{k_1}^\dagger \beta_{k_2}^\dagger \beta_{k_3}^\dagger \beta_{k_4} + \Omega_{k_1 k_2 k_3 k_4}^{13} \beta_{k_1}^\dagger \beta_{k_4} \beta_{k_3} \beta_{k_2} \right\} \\ & + \frac{1}{4!} \sum_{k_1 k_2 k_3 k_4} \left\{ \Omega_{k_1 k_2 k_3 k_4}^{40} \beta_{k_1}^\dagger \beta_{k_2}^\dagger \beta_{k_3}^\dagger \beta_{k_4}^\dagger + \Omega_{k_1 k_2 k_3 k_4}^{04} \beta_{k_4} \beta_{k_3} \beta_{k_2} \beta_{k_1} \right\} + \dots \end{aligned}$$



- Fully correlated state obtained via the evolution operator

$$\begin{aligned} |\Psi_0^A(\tau)\rangle &\equiv \mathcal{U}(\tau)|\Phi\rangle \\ &= e^{-\tau\Omega_0} \mathcal{T} e^{-\int_0^\tau d\tau \Omega_1(\tau)} |\Phi\rangle \end{aligned}$$

## Ground state energy of an open-shell nucleus

$$E_0^A = \lim_{\tau \rightarrow \infty} \langle \Psi_0^A(\tau) | \Omega | \Phi \rangle_c$$

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## Ground state energy of an open-shell nucleus

$$E_0^A = \lim_{\tau \rightarrow \infty} \langle \Psi_0^A(\tau) | \Omega | \Phi \rangle_c$$

- Diagonal propagators (no anomalous)

$$G_{k_1 k_2}^{+- (0)}(\tau_1, \tau_2) \equiv \frac{\langle \Phi | \mathcal{T} [\beta_{k_1}^\dagger(\tau_1) \beta_{k_2}(\tau_2)] | \Phi \rangle}{\langle \Phi | \Phi \rangle}$$

$$G_{k_1 k_2}^{-+ (0)}(\tau_1, \tau_2) \equiv \frac{\langle \Phi | \mathcal{T} [\beta_{k_1}(\tau_1) \beta_{k_2}^\dagger(\tau_2)] | \Phi \rangle}{\langle \Phi | \Phi \rangle}$$

with antisymmetry relation

$$G_{k_1 k_2}^{+- (0)}(\tau_1, \tau_2) = -G_{k_2 k_1}^{-+ (0)}(\tau_2, \tau_1)$$

- Perturbative expansion of ground-state energy ( $\Omega = \Omega_0 + \Omega_1$ )

$$\begin{aligned}
 E_0 &= \langle \Phi | \left\{ \Omega(0) - \int_0^\infty d\tau_1 T [\Omega_1(\tau_1) \Omega(0)] + \frac{1}{2!} \int_0^\infty d\tau_1 d\tau_2 T [\Omega_1(\tau_1) \Omega_1(\tau_2) \Omega(0)] + \dots \right\} | \Phi \rangle_c \\
 &= \sum_{p=0}^{\infty} \frac{(-1)^p}{p!} \sum_{\substack{i_0+j_0=2,4 \\ \vdots \\ i_p+j_p=2,4}} \int_0^\infty d\tau_1 \dots d\tau_p \\
 &\quad \times \sum_{\substack{k_1 \dots k_{i_1}, k_{i_1+1} \dots k_{i_1+j_1} \\ \vdots \\ l_1 \dots l_{i_p}, l_{i_p+1} \dots l_{i_p+j_p}}} \frac{\Omega_{k_1 \dots k_{i_1} k_{i_1+1} \dots k_{i_1+j_1}}^{i_1 j_1}}{(i_1)! (j_1)!} \dots \frac{\Omega_{l_1 \dots l_{i_p} l_{i_p+1} \dots l_{i_p+j_p}}^{i_p j_p}}{(i_p)! (j_p)!} \frac{\Omega_{m_1 \dots m_{i_0} m_{i_0+1} \dots m_{i_0+j_0}}^{i_0 j_0}}{(i_0)! (j_0)!} \\
 &\quad \times \langle \Phi | T \left[ \beta_{k_1}^\dagger(\tau_1) \dots \beta_{k_{i_1}}^\dagger(\tau_1) \beta_{k_{i_1+j_1}}(\tau_1) \dots \beta_{k_{i_1+1}}(\tau_1) \dots \beta_{l_1}^\dagger(\tau_p) \dots \beta_{l_{i_p}}^\dagger(\tau_p) \dots \right. \\
 &\quad \left. \times \beta_{l_{i_p+j_p}}(\tau_p) \dots \beta_{l_{i_p+1}}(\tau_p) \beta_{m_1}^\dagger(0) \dots \beta_{m_{i_0}}^\dagger(0) \beta_{m_{i_0+j_0}}(0) \dots \beta_{m_{i_0+1}}(0) \right] | \Phi \rangle_c
 \end{aligned}$$

Diagonal case:  $\varphi = 0$

- No anomalous propagator, no self-contraction
- Standard Wick's theorem with respect to  $|\Phi\rangle$

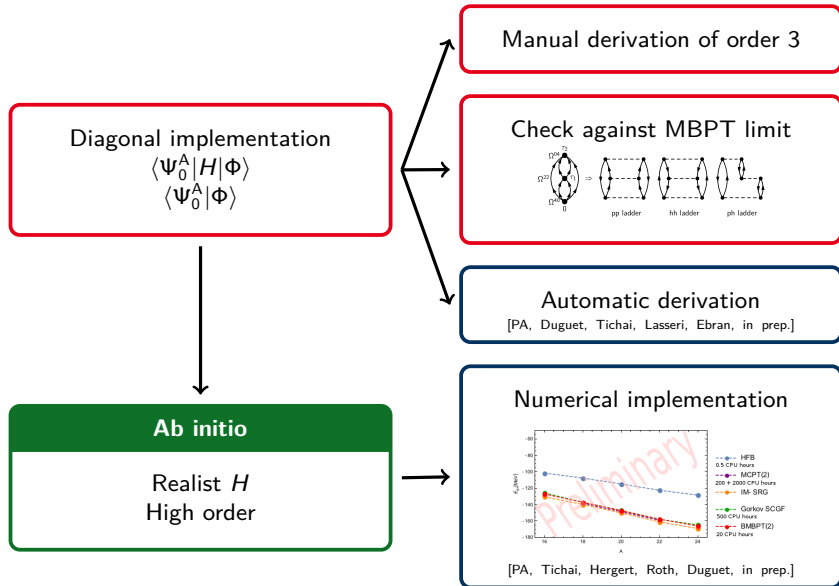
- Normal-ordered form of  $\Omega$  with respect to  $|\Phi\rangle$

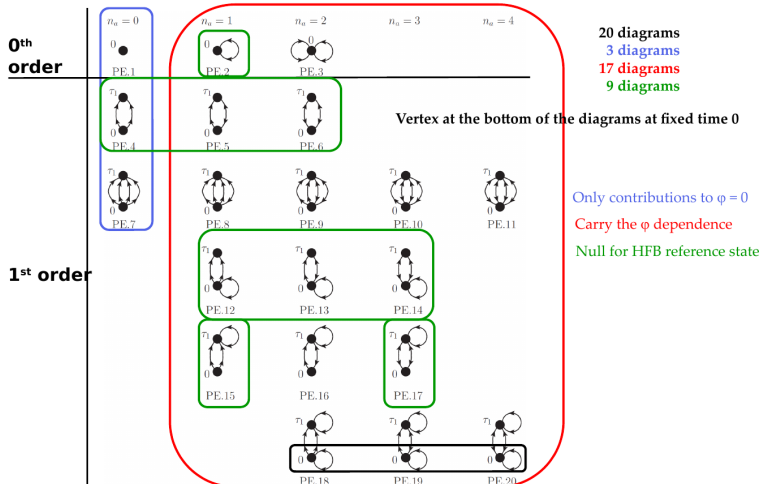
$$\Omega = \underbrace{\bullet}_{\Omega^{00}} + \underbrace{\begin{array}{c} \uparrow \\ \bullet \\ \uparrow \end{array}}_{\Omega^{11}} + \underbrace{\begin{array}{c} \swarrow \quad \nearrow \\ \bullet \end{array}}_{\Omega^{20}} + \underbrace{\begin{array}{c} \nearrow \quad \swarrow \\ \bullet \end{array}}_{\Omega^{02}} + \dots$$

- Diagonal propagators

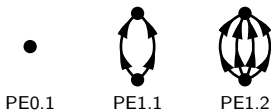
$$G_{k_1 k_2}^{+- (0)}(\tau_1, \tau_2) \quad \begin{array}{c} k_2 \tau_2 \\ \uparrow \\ \uparrow \\ k_1 \tau_1 \end{array} \quad G_{k_1 k_2}^{-+ (0)}(\tau_1, \tau_2) \quad \begin{array}{c} k_2 \tau_2 \\ \downarrow \\ \downarrow \\ k_1 \tau_1 \end{array}$$

- Main diagrammatic rules from Wick theorem
  - ◇ No external legs
  - ◇ No oriented loop between vertices
  - ◇ No self-contraction
  - ◇ Propagators go out of the  $\Omega$  vertex at time 0
  - ◇ Equivalent lines
  - ◇ Discard topologically equivalent diagrams

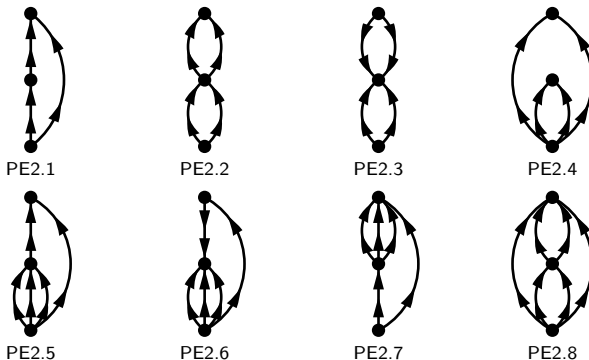




- First- and second-order diagrams [Duguet and Signoracci, *J. Phys. G* 44, 2017]

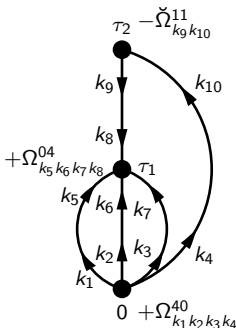


- Third-order diagrams



Validation of the manual derivation by checking the MBPT limit

# Derivation of a third-order diagram



Feynman (time-dependent) and Goldstone (time-integrated) expressions:

$$\begin{aligned}
 \text{PE2.6} &= -\frac{1}{3!} \sum_{k_1 k_2 k_3 k_4 k_8} \Omega_{k_1 k_2 k_3 k_4}^{40} \Omega_{k_1 k_2 k_3 k_8}^{04} \check{\Omega}_{k_8 k_4}^{11} \int_0^\infty d\tau_1 d\tau_2 \theta(\tau_1 - \tau_2) e^{-\tau_1 (E_{k_1} + E_{k_2} + E_{k_3} + E_{k_8})} e^{\tau_2 (E_{k_8} - E_{k_4})} \\
 &= -\frac{1}{3!} \sum_{k_1 k_2 k_3 k_4 k_8} \frac{\Omega_{k_1 k_2 k_3 k_4}^{40} \Omega_{k_1 k_2 k_3 k_8}^{04} \check{\Omega}_{k_8 k_4}^{11}}{(E_{k_1} + E_{k_2} + E_{k_3} + E_{k_4}) (E_{k_1} + E_{k_2} + E_{k_3} + E_{k_8})}
 \end{aligned}$$

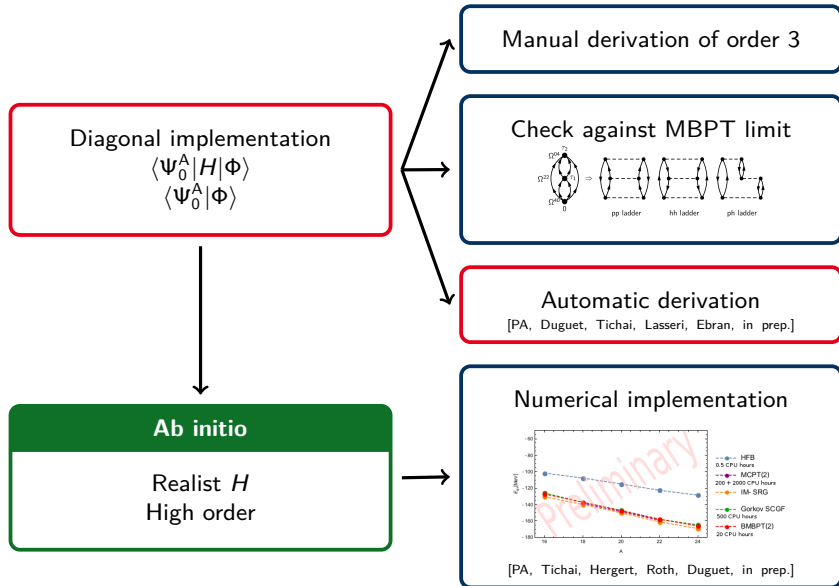


- All diagrams from  $2N$  vertices derived and implemented up to order 3  
[PA, Tichai, Ebran, Duguet]
- Want to go to higher orders
  - ◇ At least up to order 4
    - ▶ Check convergence pattern
    - ▶ Grasp effect from quadruples  $\leftrightarrow$  8 qp excitations
  - ◇ Derivation time-consuming
  - ◇ Derivation error-prone

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## Develop automatic tool

- ◇ To generate all possible connected diagrams at order  $n$
- ◇ To extract associated time-integrated expressions
- ◇ To be both quick and safe



## Our goal

An automatic and systematic way of producing diagrams

## Our tool

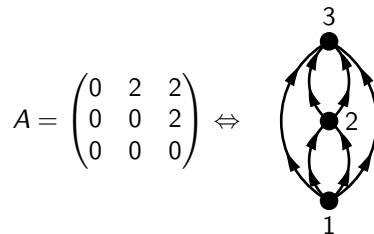
Adjacency matrices in graph theory

## Our challenge

From BMBPT diagrammatic rules to constraints on matrices

Each Feynman diagram to be represented by an adjacency matrix

- $a_{ij}$  indicate the number of edges going from node  $i$  to node  $j$



- ◇ Carry detailed information for directed graphs
- ◇ Symmetry properties and connectivity properties directly readable
- Only two propagators, readable as one once reading direction is fixed
  - ◇ Perfectly adapted for diagonal BMBPT
  - ◇ Extension needed for off-diagonal diagrams with anomalous propagator

**Each vertex belongs to  $\Omega^{[2]}$ ,  $\Omega^{[4]}$  or  $\Omega^{[6]}$**

For each vertex  $i$ ,  $\sum_j (a_{ij} + a_{ji})$  is 2, 4 or 6

**No self-contraction**

Every diagonal element is zero

**Every propagator coming out of the vertex at time 0 goes upward**

First column of the matrix is zero

**No loop between vertices**

Can restrict to upper triangular matrices

- Generate all upper triangular  $n \times n$  matrices for  $n$ -th order BMBPT diagrams
  - ◇ Fill the matrices "vertex-wise" with all allowed integers
  - ◇ Check the degree of each vertex before moving on

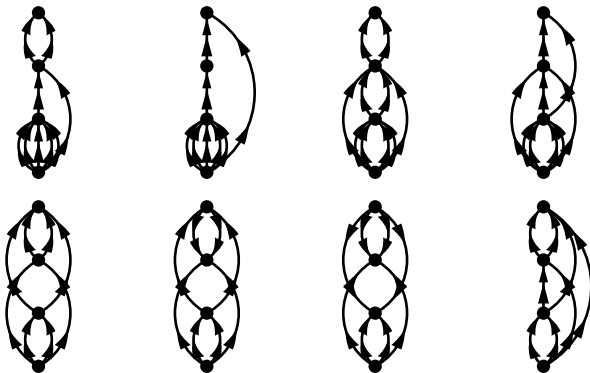
$$\begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \rightarrow \begin{pmatrix} 0 & a_{12} & a_{13} \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \rightarrow \begin{pmatrix} 0 & a_{12} & a_{13} \\ 0 & 0 & a_{23} \\ 0 & 0 & 0 \end{pmatrix}$$

- Discard matrices leading to topologically identical diagrams
- Read the matrix and translate it into drawing instructions

```
\begin{fmfgraph*}(60,60)
\fmftop{v2}\fmfbottom{v0}
\fmf{phantom}{v0,v1}
\fmfv{d.shape=circle,d.filled=full,d.size=3thick}{v0}
\fmf{phantom}{v1,v2}
\fmfv{d.shape=circle,d.filled=full,d.size=3thick}{v1}
\fmfv{d.shape=circle,d.filled=full,d.size=3thick}{v2}
\fmffreeze
\fmf{prop_pm}{v0,v1}
\fmf{prop_pm,right=0.6}{v0,v2}
\fmf{prop_pm}{v1,v2}
\fmf{prop_pm,left=0.5}{v1,v2}
\fmf{prop_pm,right=0.5}{v1,v2}
\end{fmfgraph*}
```



Run the code at order 4 with 2N and 3N interactions, obtain...



...and 388 others!



- Number of diagrams with 2N interactions (using an HFB vacuum)
  - ◇ 8 (1) diagrams at order 3
  - ◇ 59 (10) diagrams at order 4
  - ◇ 568 (82) diagrams at order 5
  - ◇ 6 805 (938) diagrams at order 6
- Number of diagrams with 2N and 3N interactions (using an HFB vacuum)
  - ◇ 23 (8) diagrams at order 3
  - ◇ 396 (177) diagrams at order 4
  - ◇ 10 716 (5 055) diagrams at order 5
  - ◇ 100 000+ diagrams at order 6?
- Obtained in only a few minutes...

All BMBPT diagrams produced automatically at a given order

➡ Need to derive automatically the diagrams' expressions

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- Feynman diagrams recast different time-orderings
  - ✓ Less diagrams to set up
  - ✗ But time-integrated (Goldstone) expressions are to be coded

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- Goldstone diagrams capture each time ordering separately
  - ✓ Time-integrated expressions obtained directly from diagrammatic rules
  - ✗ Many more diagrams to consider

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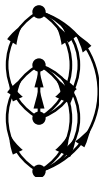
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## Challenge: Extract Goldstone expressions from Feynman diagrams

- ◇ Capture all time ordering at once
- ◇ Challenging because of structure of corresponding time integrals
- ◇ Undone task to our knowledge (even for standard diagrammatic)

- Extract graph structure info as well
  - ◇ Associate labels with vertices, propagators, etc.
  - ◇ In- / out-degree of vertices associated with annihilators / creators
  - ◇ Run routines for symmetry factors
- Have your code write the corresponding equations in your .tex file

$$\frac{-(-1)^3}{(3!)^2} \sum_{k_i} \Omega_{k_1 k_2 k_3 k_4}^{40} \Omega_{k_5 k_6 k_7 k_8}^{40} \Omega_{k_5 k_1 k_2 k_3}^{04} \Omega_{k_6 k_7 k_8 k_4}^{04} \int_0^\tau d\tau_1 d\tau_2 d\tau_3 \theta(\tau_2 - \tau_1) \theta(\tau_3 - \tau_1) \\ \times e^{-\tau_1(-E_{k_5} - E_{k_6} - E_{k_7} - E_{k_8})} e^{-\tau_2(E_{k_1} + E_{k_2} + E_{k_3} + E_{k_5})} e^{-\tau_3(E_{k_4} + E_{k_6} + E_{k_7} + E_{k_8})}$$



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Sign, prefactor and operators left unchanged in Goldstone expression

➡ Only need to extract the denominator

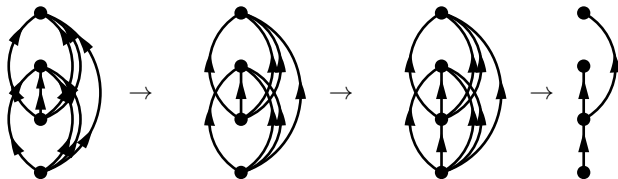
- Introduce time-structure diagrams (TSDs)
  - ◇ Links carry time-ordering relations, moving towards higher times
  - ◇ Contain only the minimal set of links to describe all the time relations



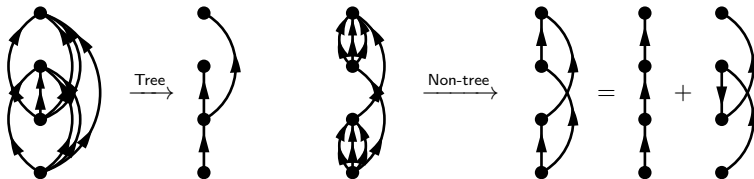
- Determine the time-structure diagram (TSD) associated to BMBPT one
  - ◇ Propagators carry time-ordering relations
  - ◇  $\Omega$  vertex at time 0 is a lower limit for time
  - ◇ One TSD recast several Feynman, even more Goldstone



- Each TSD produced from the BMBPT diagram
  - ◇ Replace propagators by links
  - ◇ Add links between vertex at time 0 and other vertices
  - ◇ Remove links carrying unnecessary information

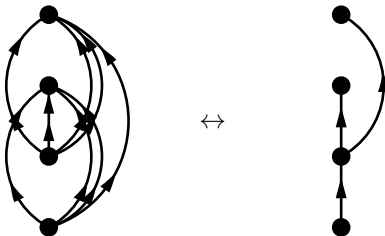


- Extraction of time-integrated expression depends on tree / non-tree



For each perturbation vertex in the diagram with an associated tree TSD

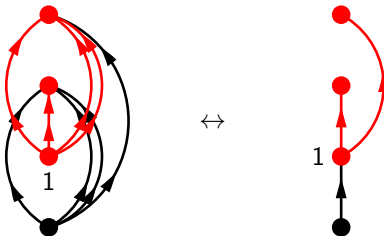
- ① Determine all its descendants using the TSD diagram
- ② Form a subgraph using the vertex and its descendants
- ③ For all propagators entering the subgraph, add the associated qpe



$$\frac{-(-1)^3}{(3!)^2} \sum_{k_i} \frac{\Omega_{k_1 k_2 k_3 k_4}^{40} \Omega_{k_5 k_6 k_7 k_8}^{40} \Omega_{k_5 k_1 k_2 k_3}^{04} \Omega_{k_6 k_7 k_8 k_4}^{04}}{(E_{k_1} + E_{k_2} + E_{k_3} + E_{k_4})(E_{k_1} + E_{k_2} + E_{k_3} + E_{k_5})(E_{k_4} + E_{k_6} + E_{k_7} + E_{k_8})}$$

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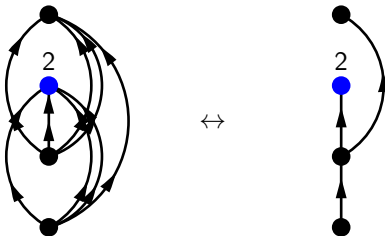
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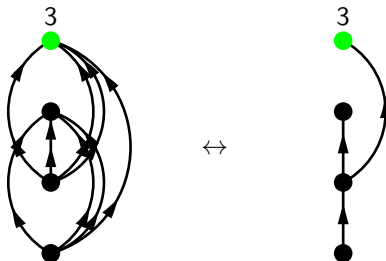
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$$\frac{-(-1)^3}{(3!)^2} \sum_{k_i} \frac{\Omega_{k_1 k_2 k_3 k_4}^{40} \Omega_{k_5 k_6 k_7 k_8}^{40} \Omega_{k_5 k_1 k_2 k_3}^{04} \Omega_{k_6 k_7 k_8 k_4}^{04}}{(E_{k_1} + E_{k_2} + E_{k_3} + E_{k_4})(E_{k_1} + E_{k_2} + E_{k_3} + E_{k_5})(E_{k_4} + E_{k_6} + E_{k_7} + E_{k_8})}$$

## Why a so simple denominator algorithm for all trees?

Link between tree TSD structure and time integrals structure

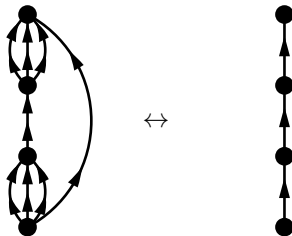


$$\begin{aligned} D &= \lim_{\tau \rightarrow \infty} \int_0^\tau d\tau_1 d\tau_2 d\tau_3 \theta(\tau_3 - \tau_1) \theta(\tau_2 - \tau_1) e^{a\tau_1} e^{b\tau_2} e^{c\tau_3} \\ &= \lim_{\tau \rightarrow \infty} \int_0^\tau d\tau_1 e^{a\tau_1} \int_0^{\tau_1} d\tau_2 e^{b\tau_2} \int_0^{\tau_1} d\tau_3 e^{c\tau_3} \\ &= \lim_{\tau \rightarrow \infty} \frac{1}{bc} \int_0^\tau d\tau_1 e^{a\tau_1} (e^{b\tau} - e^{b\tau_1}) (e^{c\tau} - e^{c\tau_1}) \\ &= \frac{1}{bc(a + b + c)} \end{aligned}$$

- Integrate from the leaves first
- Go down each branch
- Each vertex depends on the vertices above it

## Same algorithm applied on linear tree

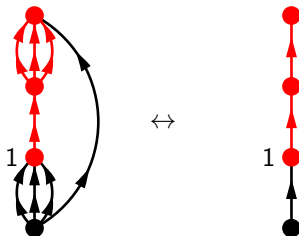
Classic Goldstone rule recovered on a Feynman graph



$$\frac{(-1)^3}{(3!)^2} \sum_{k_i} \frac{\Omega_{k_1 k_2 k_3 k_4}^{40} \Omega_{k_5 k_1 k_2 k_3}^{13} \Omega_{k_6 k_7 k_8 k_5}^{31} \Omega_{k_6 k_7 k_8 k_4}^{04}}{(E_{k_1} + E_{k_2} + E_{k_3} + E_{k_4})(E_{k_5} + E_{k_4})(E_{k_4} + E_{k_6} + E_{k_7} + E_{k_8})}$$

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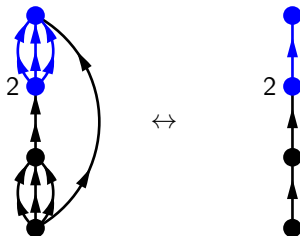


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## Same algorithm applied on linear tree

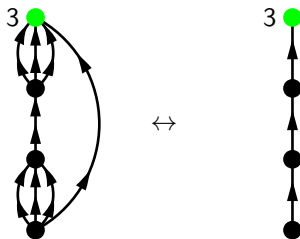
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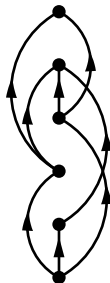
## Same algorithm applied on linear tree

Classic Goldstone rule recovered on a Feynman graph

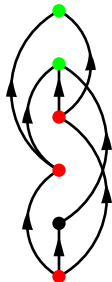


$$\frac{(-1)^3}{(3!)^2} \sum_{k_i} \frac{\Omega_{k_1 k_2 k_3 k_4}^{40} \Omega_{k_5 k_1 k_2 k_3}^{13} \Omega_{k_6 k_7 k_8 k_5}^{31} \Omega_{k_6 k_7 k_8 k_4}^{04}}{(E_{k_1} + E_{k_2} + E_{k_3} + E_{k_4})(E_{k_5} + E_{k_6})(E_{k_4} + E_{k_6} + E_{k_7} + E_{k_8})}$$

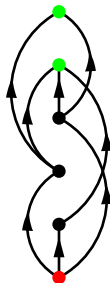
- For each node\_a with out\_degree  $\geq 2$ :
  - ◇ For each node\_b different from node\_a:
    - ▶ List all paths going from node\_a to node\_b
    - ▶ If in\_degree(node\_b)  $\geq 2$  and nb\_paths  $\geq 2$ :  
node\_a and node\_b are end nodes of the cycle
    - ▶ Check that the two paths share only their end ones



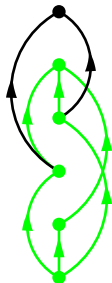
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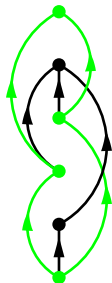


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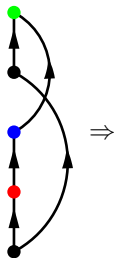
➡ Two pairs of end nodes producing cycles to be addressed

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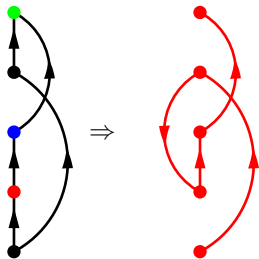
➡ Two pairs of end nodes producing cycles to be addressed

- Set `node_to_insert` as the first node of `path_1` after start node
- For each `daughter_node` in `path_2` but the starting node:
  - ◇ Make a copy of the graph
  - ◇ Add an edge from `node_to_insert` to `daughter_node`
  - ◇ Set `mother_node` as the node preceding `daughter_node` in `path_2`
  - ◇ Add an edge from `mother_node` to `daughter_node`
  - ◇ Remove the edges carrying unnecessary information

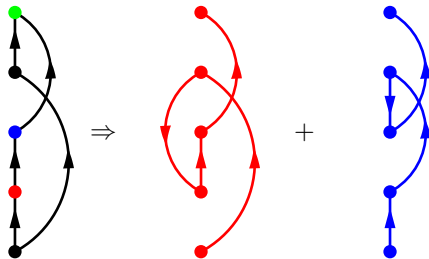




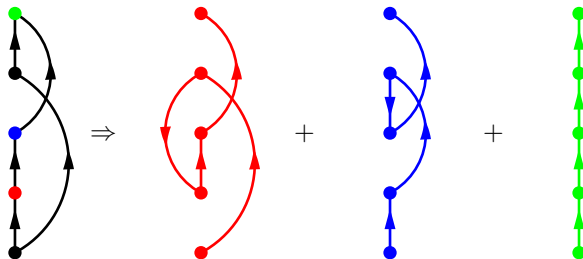
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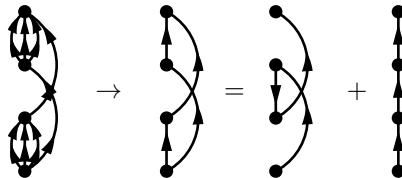


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If the associated time-structure diagram is not a tree:

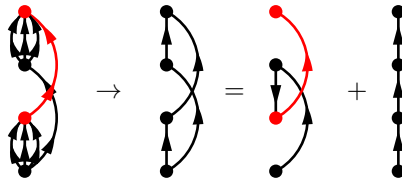
- Separate the TSD in a sum of tree TSDs
- Apply the tree denominator algorithm, sum the results



$$\frac{-(-1)^3}{(3!)^2} \sum_{k_i} \Omega_{k_1 k_2 k_3 k_4}^{40} \Omega_{k_5 k_1 k_2 k_3}^{13} \Omega_{k_6 k_7 k_8 k_4}^{31} \Omega_{k_6 k_7 k_8 k_5}^{04} \\ \times \left[ \frac{1}{(E_{k_1} + E_{k_2} + E_{k_3} + E_{k_6} + E_{k_7} + E_{k_8})(E_{k_1} + E_{k_2} + E_{k_3} + E_{k_4})(E_{k_5} + E_{k_6} + E_{k_7} + E_{k_8})} \right. \\ \left. + \frac{1}{(E_{k_1} + E_{k_2} + E_{k_3} + E_{k_4})(E_{k_4} + E_{k_5})(E_{k_5} + E_{k_6} + E_{k_7} + E_{k_8})} \right]$$

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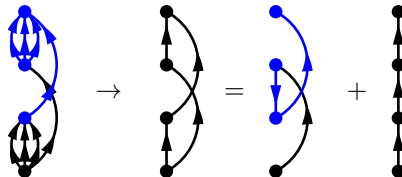
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 \end{aligned}$$

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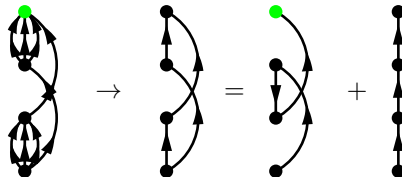
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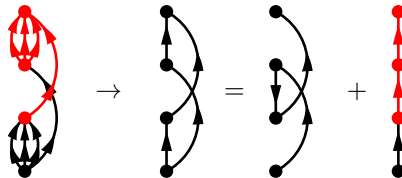
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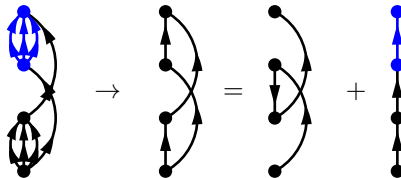


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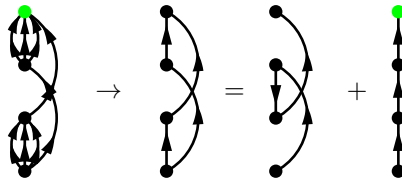
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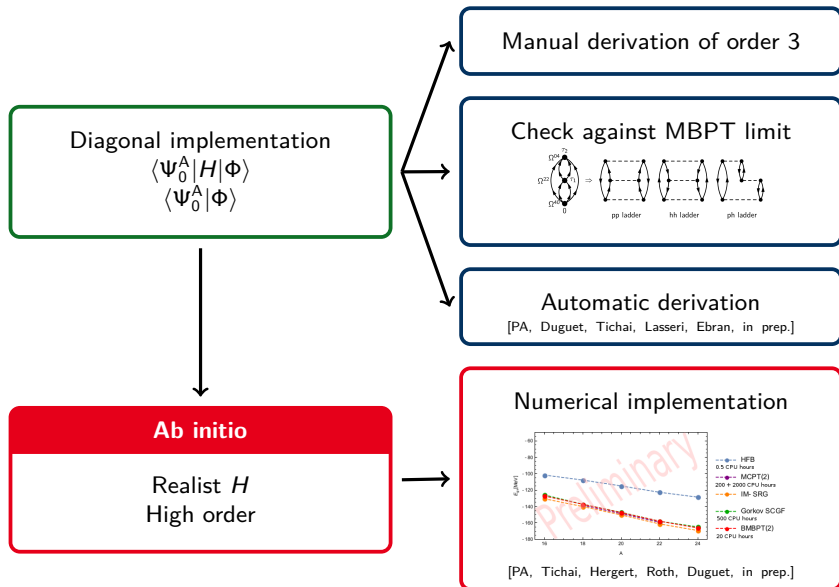
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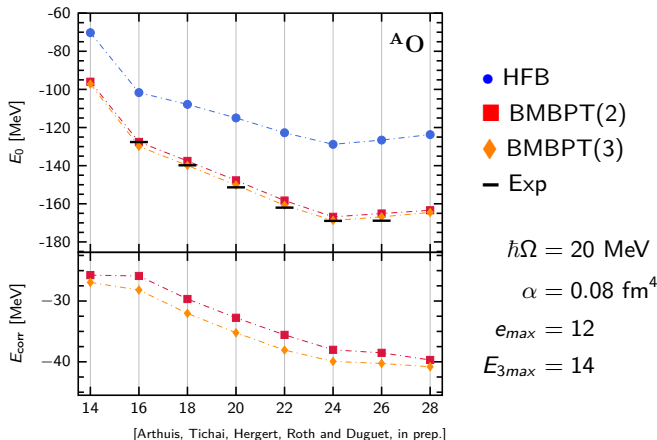


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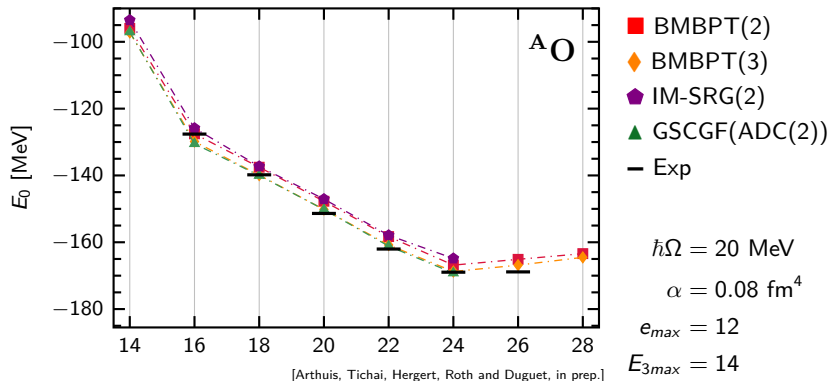
All BMBPT expressions produced automatically at a given order

➡ Need to implement them numerically

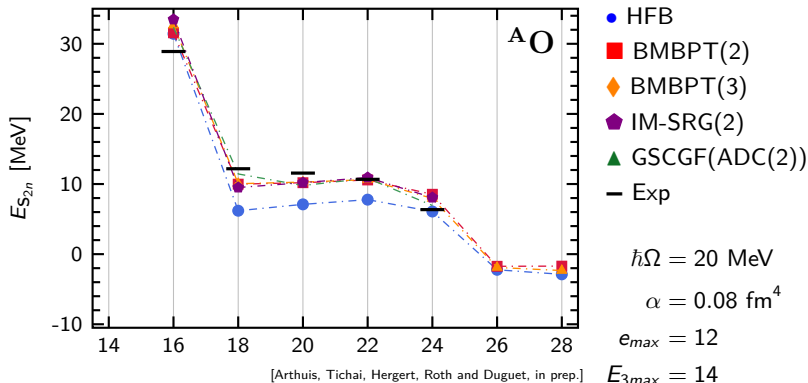




- $E^{(3)}$  one order of magnitude smaller than  $E^{(2)}$
- Computer resources independent of system size: **10-20 CPU hours**
- Error estimate on 3rd order correction:  $\Delta E = \Delta A^{(3)} \cdot 8\text{MeV}/A \approx 5\text{MeV}$
- **Calculations of Ca, Ni and Sn chains coming soon**



- Consistent with different non-perturbative methods
- **Comparable accuracy within 1-5 % of computing time**
- Computational scaling independent of system size



- **Very good agreement with state-of-the-art approaches**
- Reproduction of experimentally observed shell gaps
- Little overall effect of particle-number breaking (similar to GGF)
- **Particle-number restoration could impact near magic numbers**

- BMBPT diagrams now generated automatically
  - ✓ Fast and error-safe
  - ✓ No intrinsic upper limit on the order
- BMBPT analytical expressions automatically derived to all order as well
  - ✓ Feynman and Goldstone expressions for all diagrams
  - ✓ Order 4 to be implemented in BMBPT code in near future
- Project still moving on
  - ◇ Code to be published
  - ◇ Open to collaborations regarding other diagrammatic methods
- Numerical implementation of BMBPT(2) and BMBPT(3)
  - ✓ Very low-cost correlated method
  - ✓ Competes with state-of-the-art *ab initio* methods

- Extend the scope of ADG
  - ◇ Gorkov SCGF
  - ◇ Off-diagonal BMBPT
- Extend the scope of diagonal BMBPT
  - ◇ Excited states and new observables
  - ◇ Developments used in parallel in future BCC implementation
- Move towards symmetry-restored BMBPT
  - ◇ Extensive work on the theory
  - ◇ Automated diagram generation and derivation
  - ◇ Implementation in the BMBPT numerical code



## BMBPT Project



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TECHNISCHE  
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DARMSTADT

R. Roth



MICHIGAN STATE  
UNIVERSITY

H. Hergert

## On broader aspects



M. Drissi  
J. Ripoché



R.-D. Lasserri