



Shape coexistence in $N \approx Z$ nuclei around $A \approx 70$



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Shape Coexistence and Electric Monopole Transitions in Atomic Nuclei
Saclay
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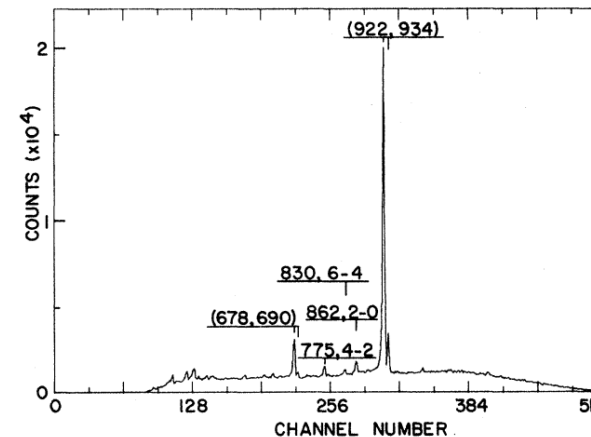
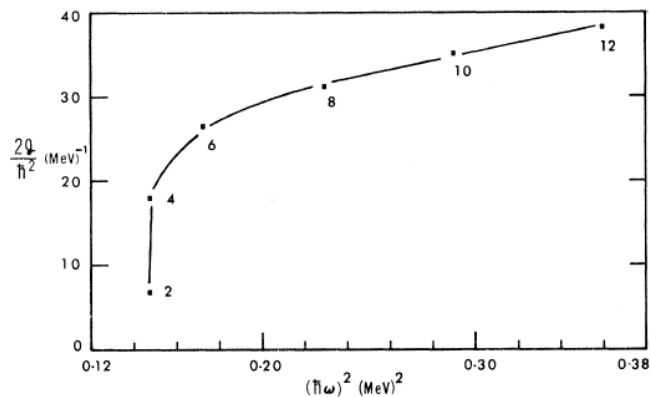
Some history...

- low-lying 0_2^+ states in ^{70}Ge and ^{72}Ge interpreted as shape coexistence between spherical vibrational and deformed rotational states

K. W. C. Stewart, B. Castel, *Lett. Nuovo Cimento* 4, 589 (1970)

- Observation of a low-lying 0_2^+ state in ^{72}Se by in-beam conversion electron spectroscopy

W.G. Wyckoff, J.E. Draper, *PRC* 8, 796 (1973)

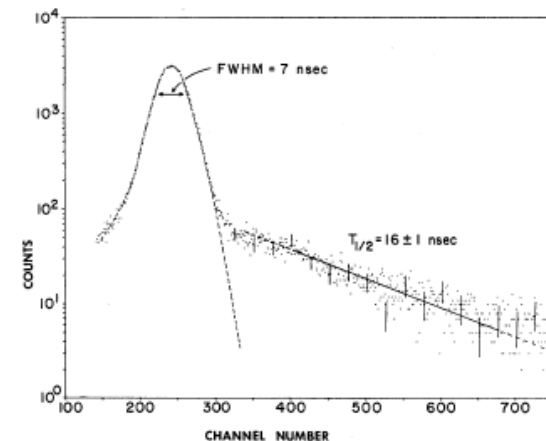


- distortion of ground-state band at low spin in ^{72}Se explained by shape coexistence and mixing

J.H. Hamilton et al., *PRL* 32, 239 (1974)

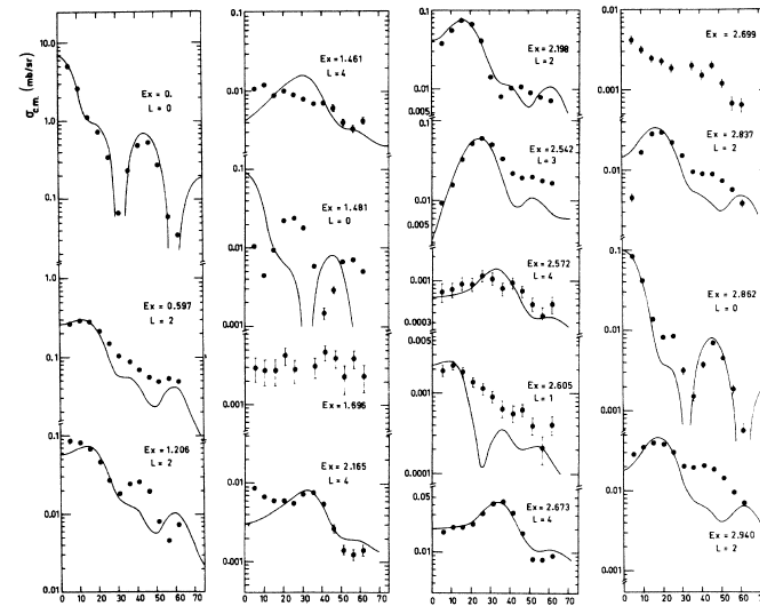
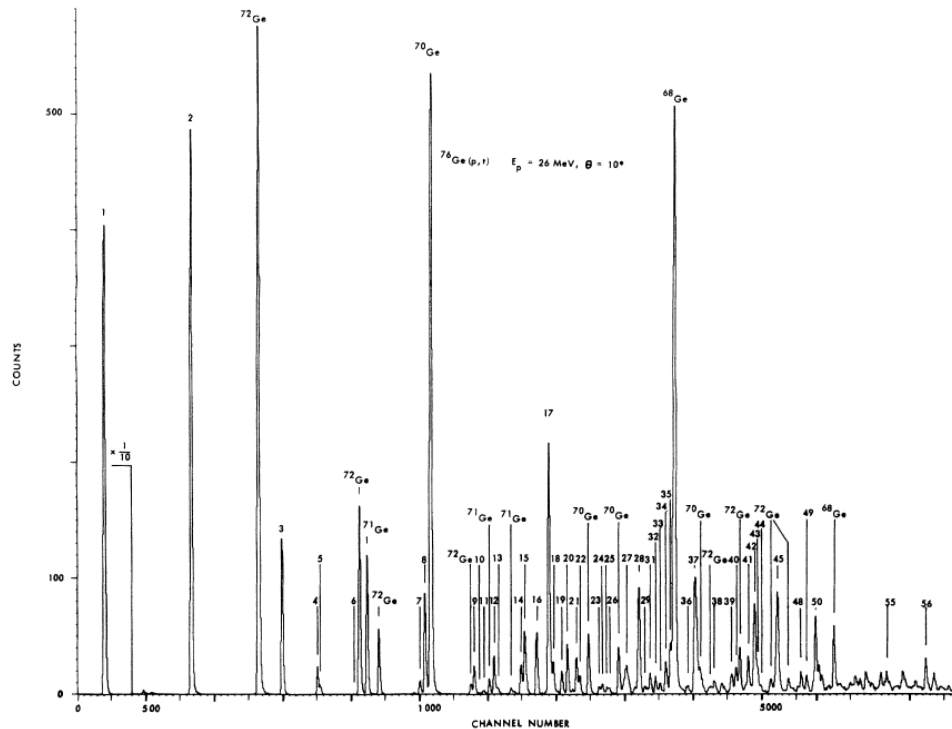
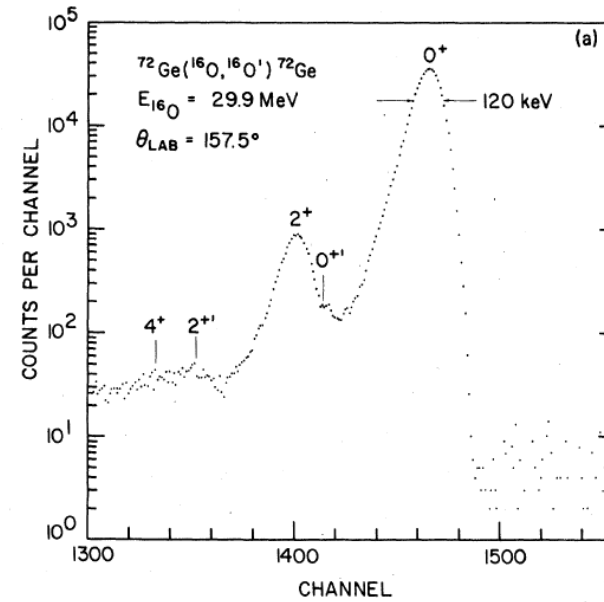
- lifetime measurement for 0_2^+ state in $^{72}\text{Se} \Rightarrow \rho^2(E0)$

A.V. Ramayya et al, *PRC* 12, 1360 (1975)



Experimental results for Ge isotopes

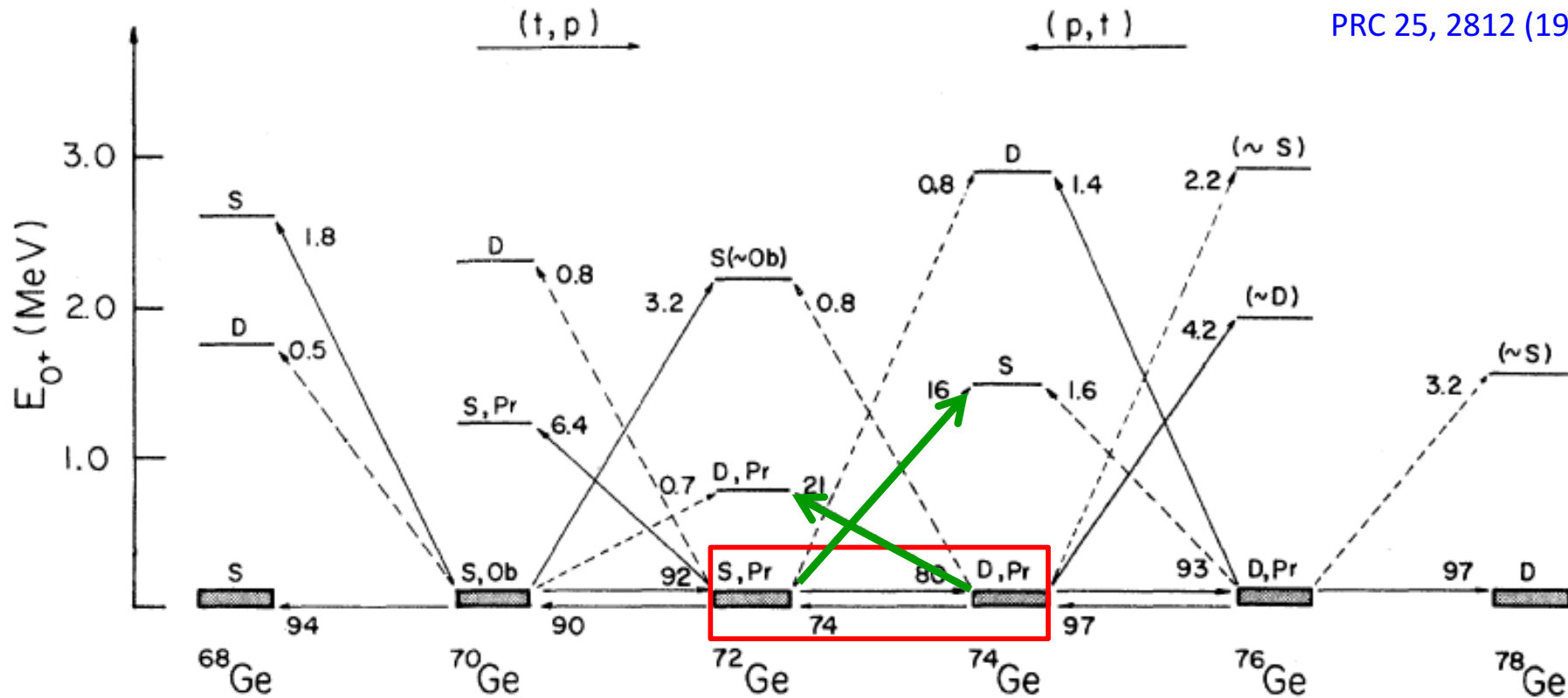
- Coulomb excitation of stable Ge isotopes
 $B(E2)$ values and $Q_s(2_1^+)$ (reorientation)
[R.Lecomte et al., PRC 22, 1530 \(1980\)](#)
- (p, t) and (t, p) pair transfer reactions
 cross sections and angular distributions
[F. Guilbault et al., PRC 16, 1840 \(1977\)](#)



Systematics for Ge isotopes

- combining quadrupole moments and transfer cross sections

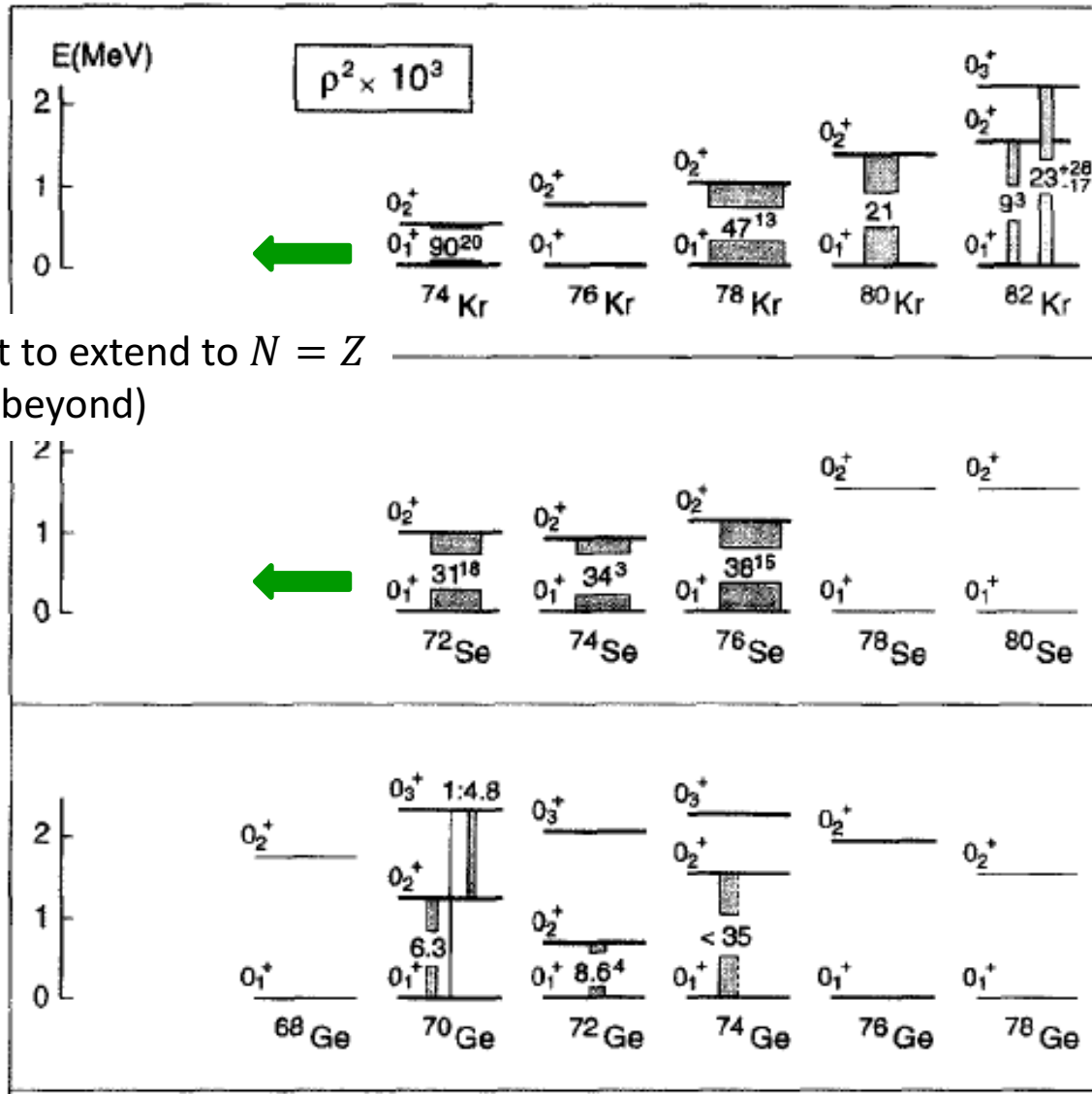
R. Lecomte et al.,
PRC 25, 2812 (1982)



- transition from spherical to deformed ground state with increasing N
- reduced cross section for $0_1^+ \rightarrow 0_1^+$ for $^{72}\text{Ge} \leftrightarrow ^{74}\text{Ge}$
- enhanced cross section for $0_1^+ \rightarrow 0_2^+$ where configurations cross

Electric monopole strengths

- low-lying 0_2^+ states and strong $E0$ transitions observed throughout the region



effort to extend to $N = Z$
(and beyond)

largest $\rho^2(E0)$
for ^{74}Kr

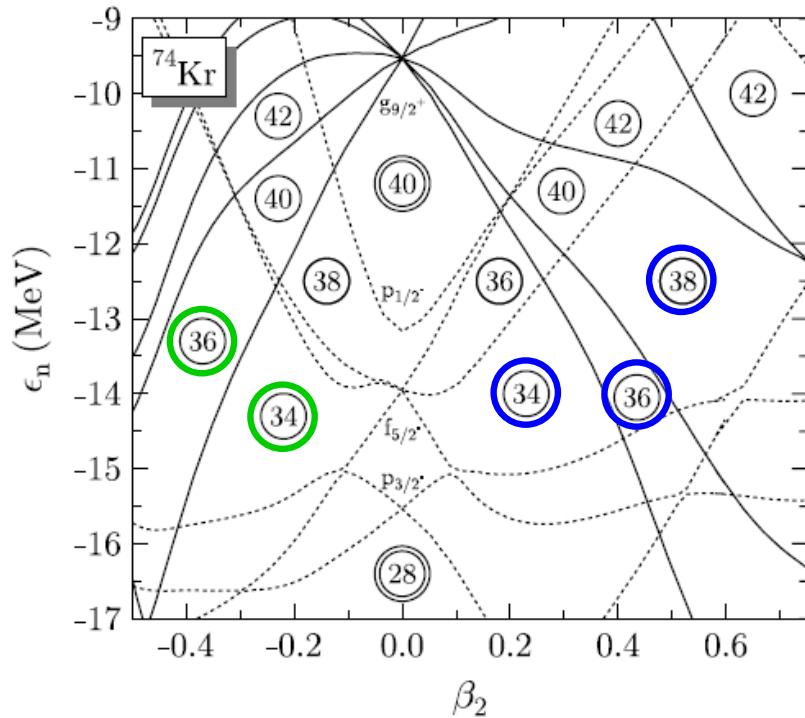
Ge and Se:
min $E(0_2^+)$ and
max $\rho^2(E0)$ for
 $Z \approx 40$.

J.L. Wood et al.
Nucl. Phys. A 651, 323 (1999)

Deformed shell gaps and potential energies

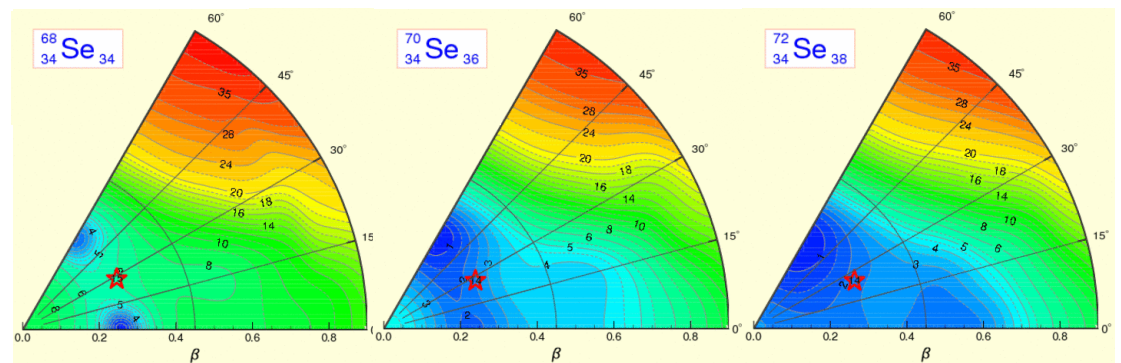
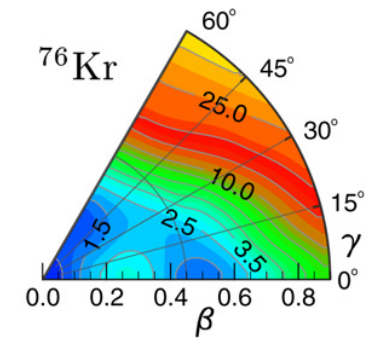
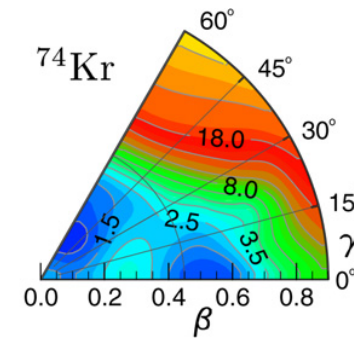
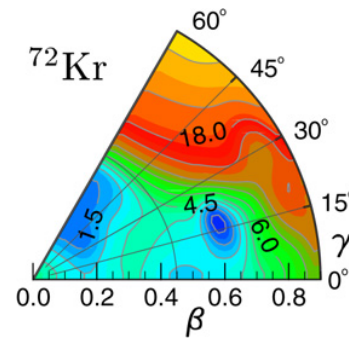
➤ neutron single-particle energies (SLy6)

➤ potential energy surfaces (Gogny D1S)



M. Bender et al., PRC 74, 024312 (2006)

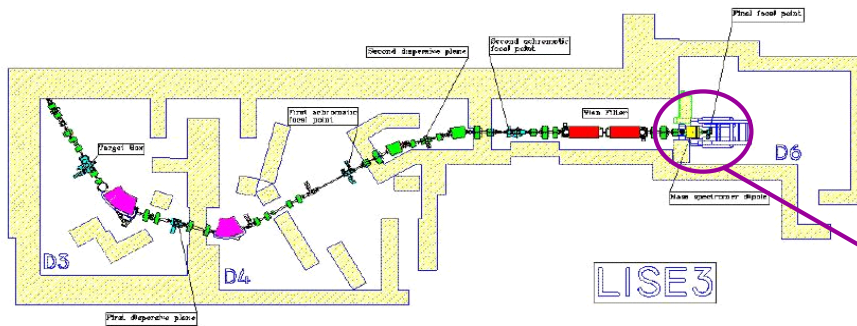
➤ competition between prolate and oblate shapes



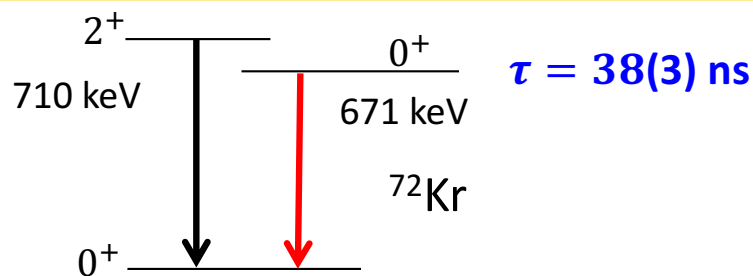
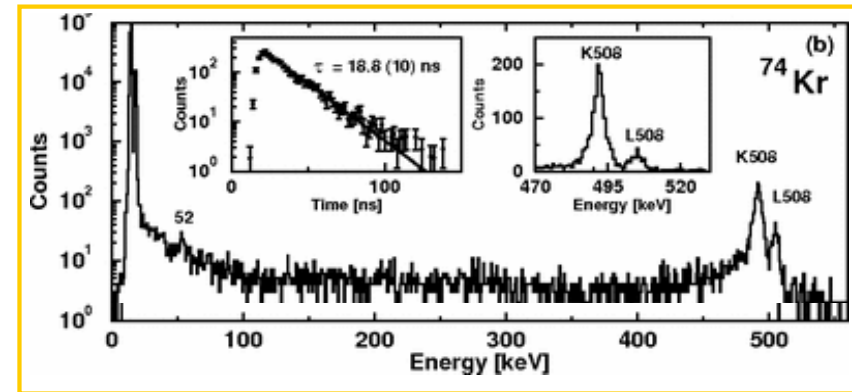
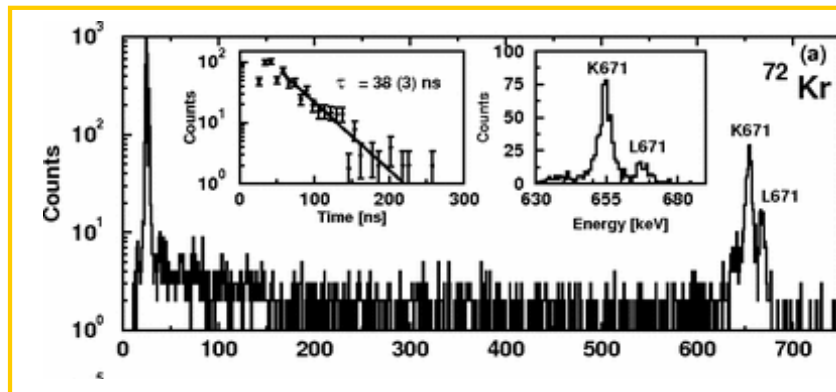
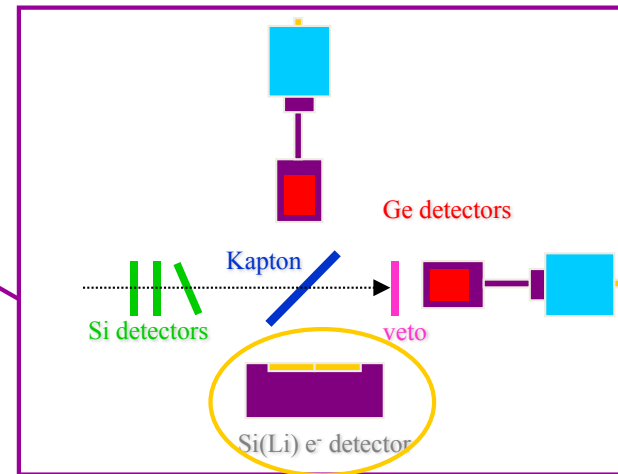
M. Girod et al., Phys. Lett. B 676, 39 (2009)
<http://www-phynu.cea.fr/>

Shape isomers in ^{72}Kr and ^{74}Kr

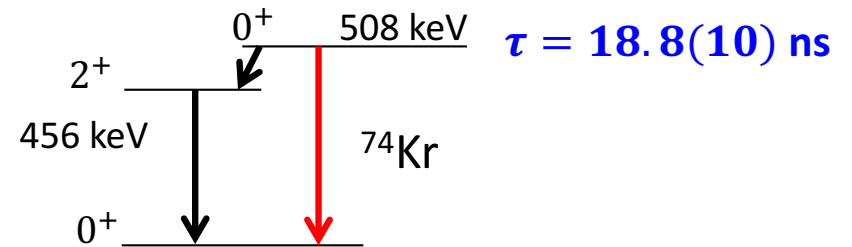
➤ isomer spectroscopy at GANIL / LISE



E. Bouchez et al.,
PRL 90, 082502 (2003)

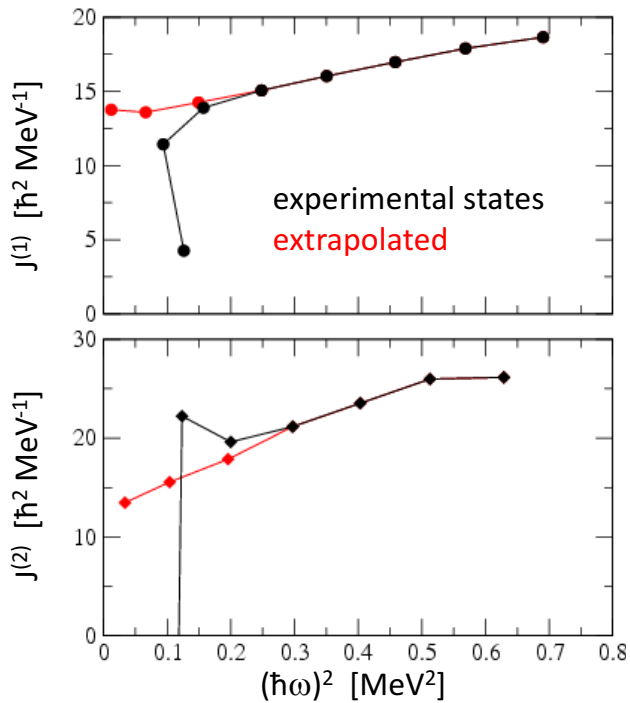


$$\rho^2(E0; 0_2^+ \rightarrow 0_1^+) = 72(6) \cdot 10^{-3}$$



$$\rho^2(E0; 0_2^+ \rightarrow 0_1^+) = 85(19) \cdot 10^{-3}$$

Band mixing

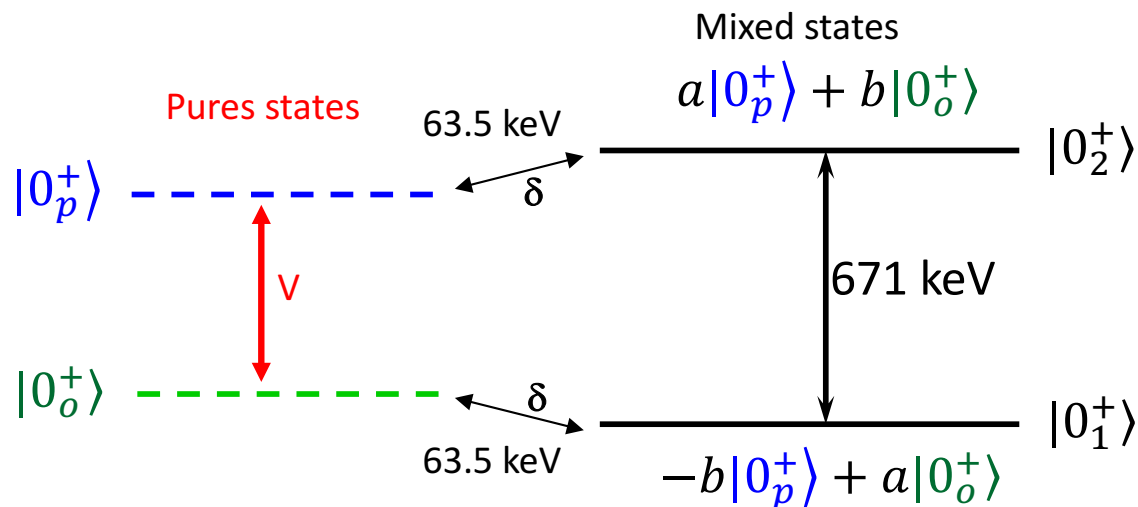


Regular rotational cascade at high spin:

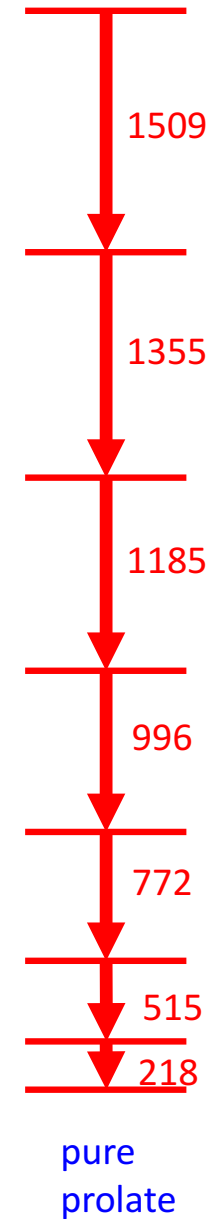
$$E(I) = \frac{\hbar^2}{2\mathfrak{I}} I(I + 1)$$

Rotational band is distorted at low spin.
 ⇒ influence of mixing

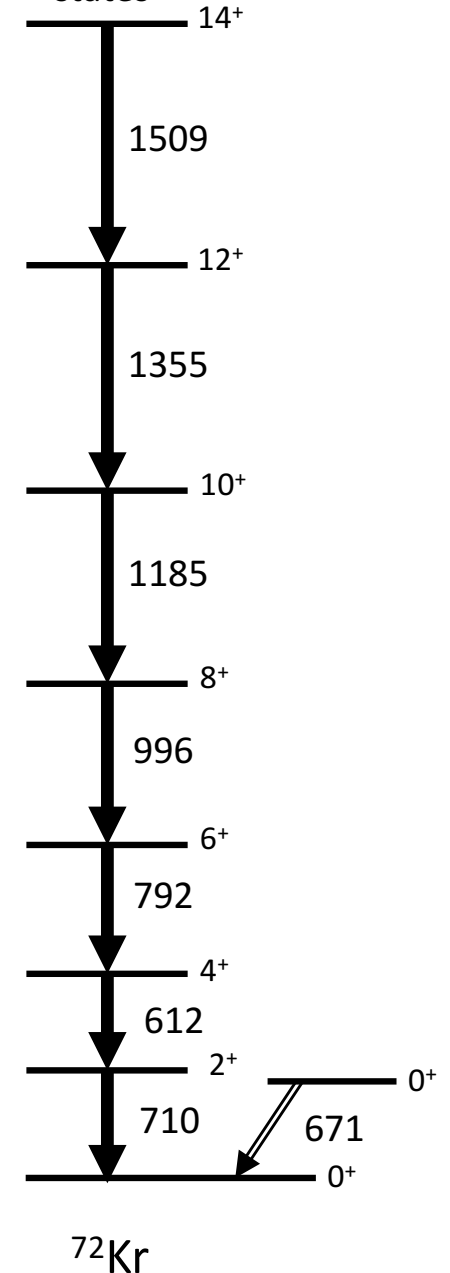
- Interaction V
- mixing amplitudes a, b



extrapolated states



experimental states

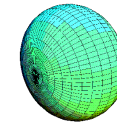
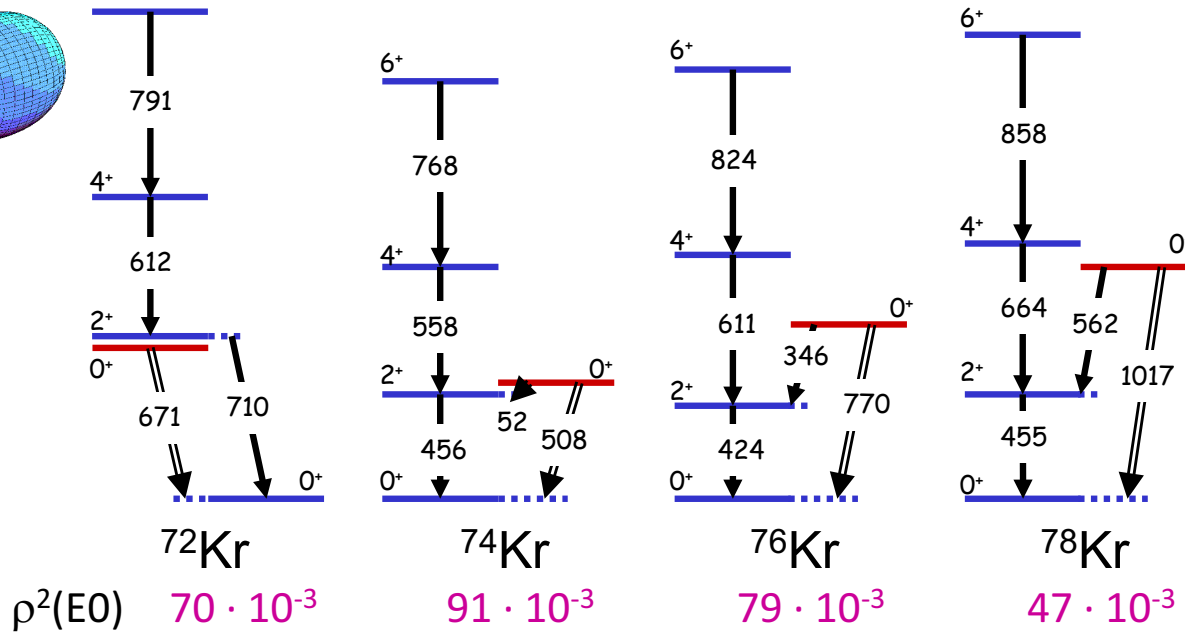
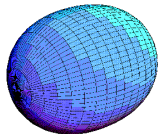


Systematics of the light krypton isotopes

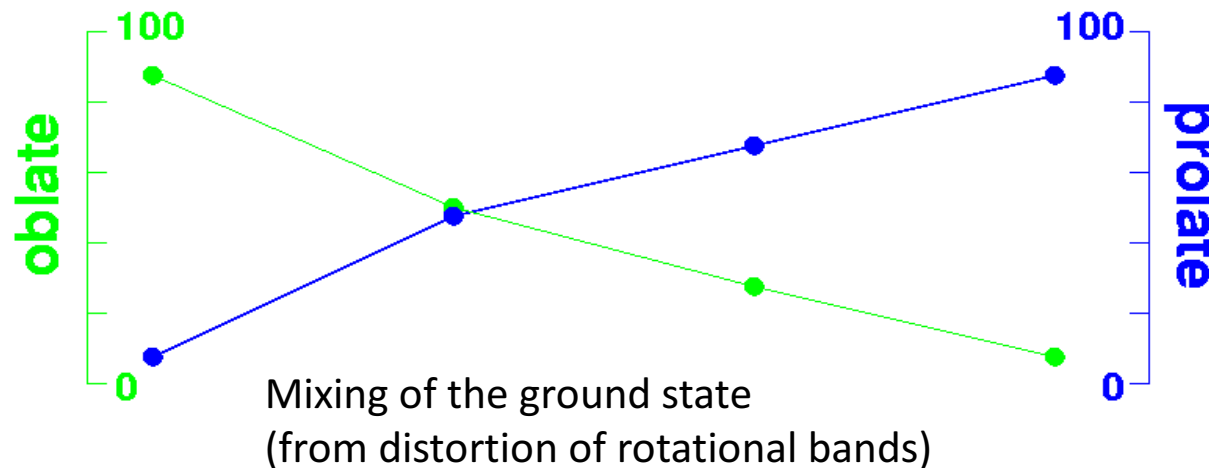
- energy of excited 0^+
- E0 strengths $\rho^2(E0)$
- configuration mixing

- Inversion of ground state shape for ^{72}Kr

prolate



oblate

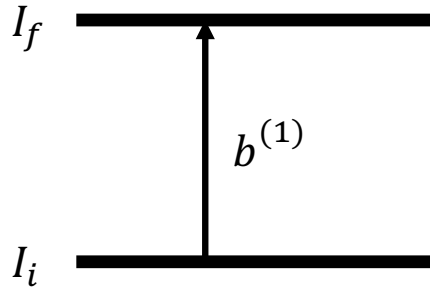


- next step:
Coulomb excitation

E. Bouchez et. al.,
Phys. Rev. Lett. 90, 082502 (2003)

Second-order effects in Coulomb excitation

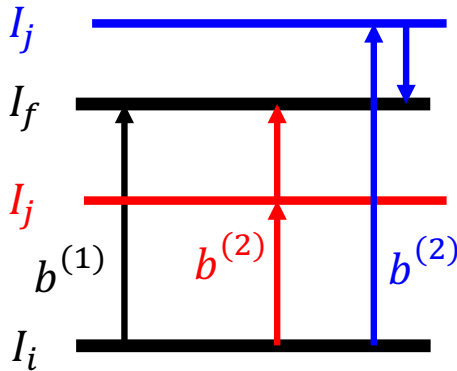
$$P_{i \rightarrow f} = \frac{1}{2I_i + 1} |b_{i \rightarrow f}^{(1)}|^2$$



1st order:

$$b_{i \rightarrow f}^{(1)} \propto \langle I_f | M(\sigma\lambda) | I_i \rangle$$

$$P_{i \rightarrow f} = \frac{1}{2I_i + 1} |b_{i \rightarrow f}^{(1)} + b_{i \rightarrow f}^{(2)}|^2$$



2nd order:

$$b_{i \rightarrow f}^{(2)} \propto \sum_j \langle I_f | M(\sigma\lambda) | I_j \rangle \langle I_j | M(\sigma\lambda) | I_i \rangle$$

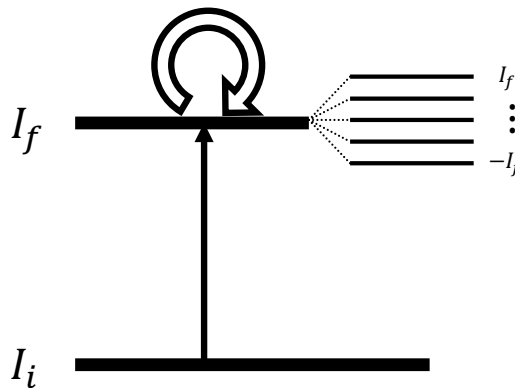
$P_{i \rightarrow f}$ depends on:

- transitional matrix element $B(\sigma\lambda; I_i \rightarrow I_f) \propto |\langle I_f | M(\sigma\lambda) | I_i \rangle|^2$
- diagonal matrix element $\langle I_f | M(\sigma\lambda) | I_f \rangle \propto Q_s(I_f)$

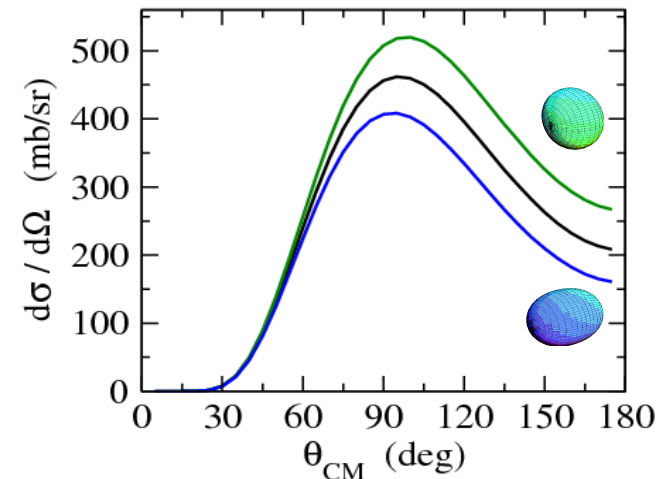
special case:
intermediate state
another magnetic substate

“reorientation effect”

$$b_{i \rightarrow f}^{(2)} \propto \langle I_f | M(\sigma\lambda) | I_f \rangle \langle I_f | M(\sigma\lambda) | I_i \rangle$$

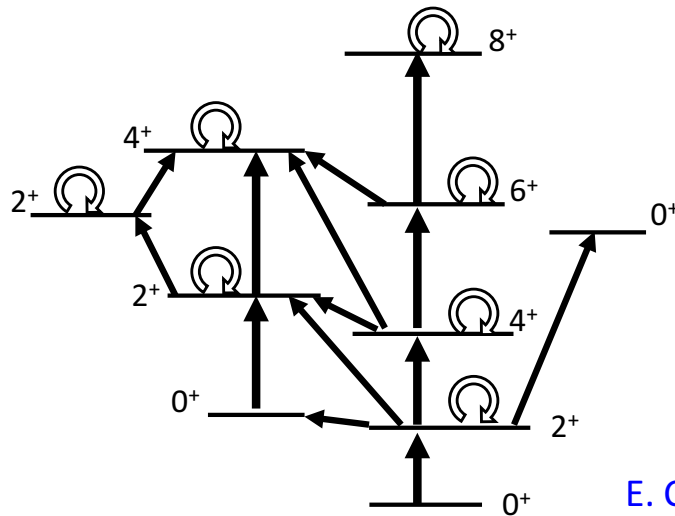


Coulomb excitation sensitive
to spectroscopic quadrupole moments

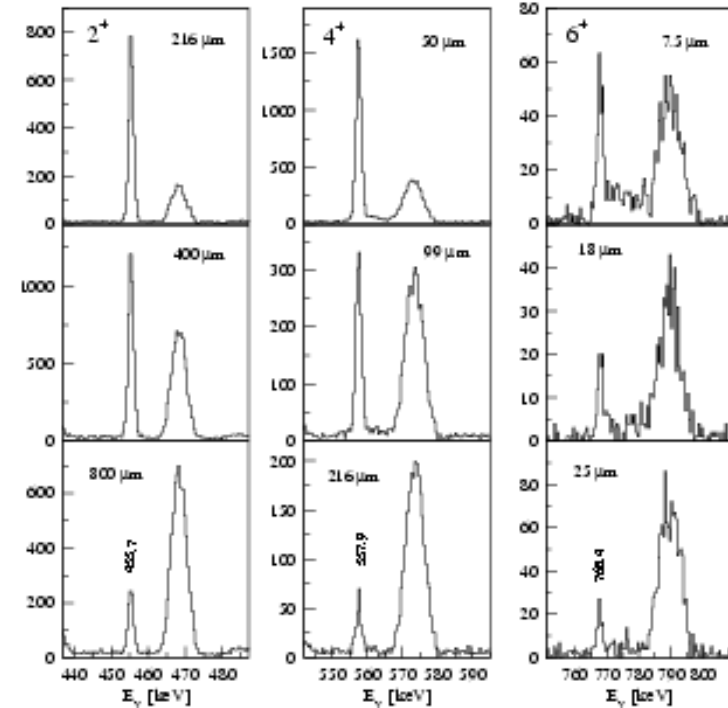


Multi-step Coulomb excitation of ^{74}Kr and ^{76}Kr at GANIL/SPIRAL

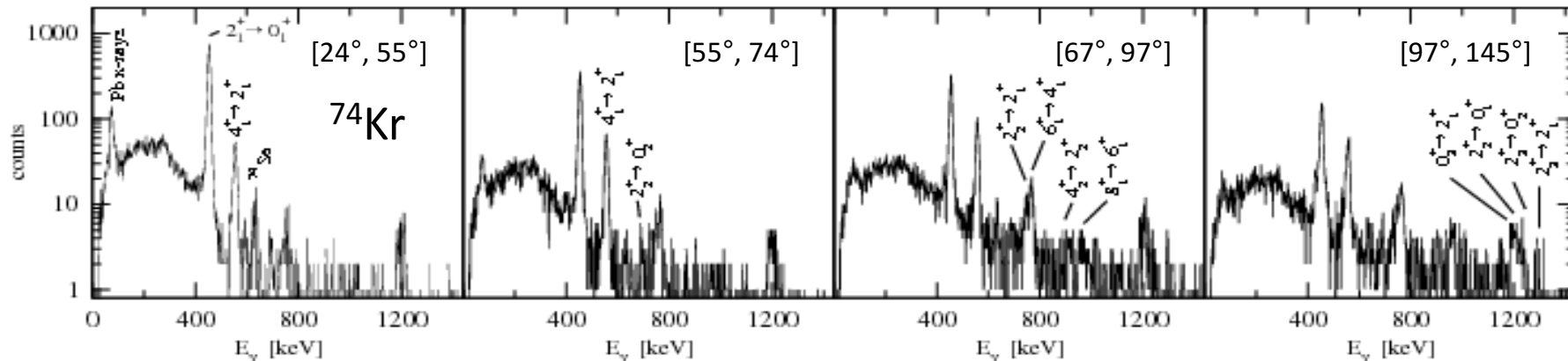
- γ yields as function of scattering angle: differential cross section
- least squares fit of ~ 30 matrix elements (transitional and diagonal)
- experimental spectroscopic data
 - lifetimes, branching and mixing ratios



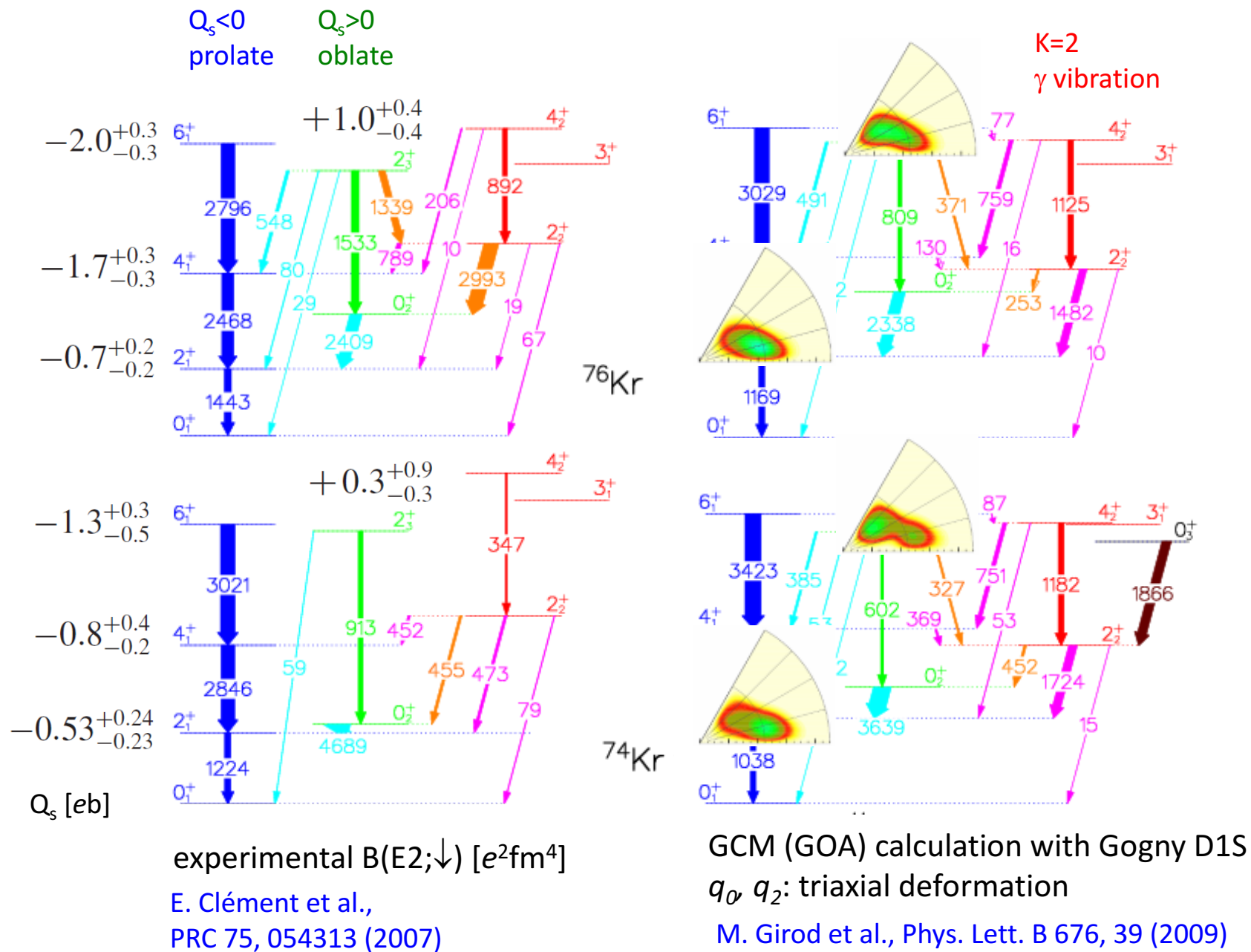
E. Clément et al., PRC 75, 054313 (2007)



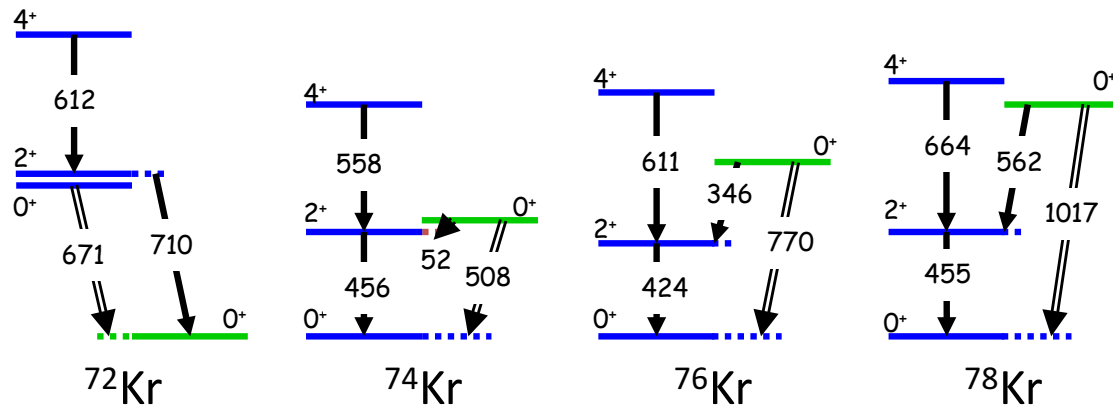
A. G3rgen, Eur. Phys. J. A 26, 153 (2005)



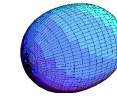
Prolate – oblate shape coexistence



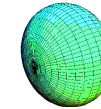
Shape transition in the light krypton isotopes



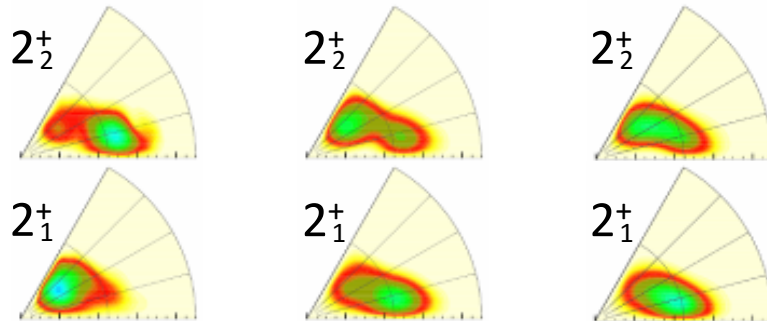
Experimental
(direct and indirect evidence)



prolate



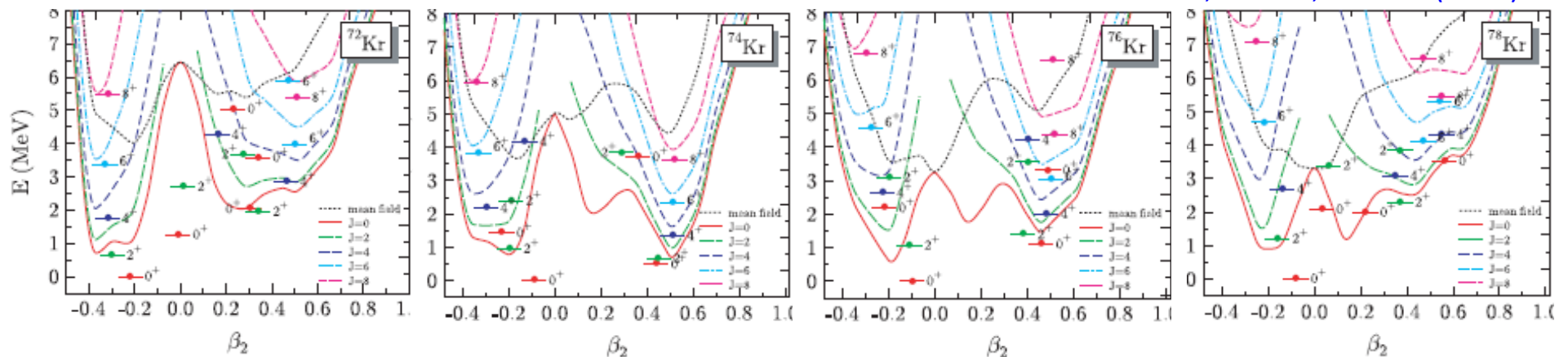
oblate



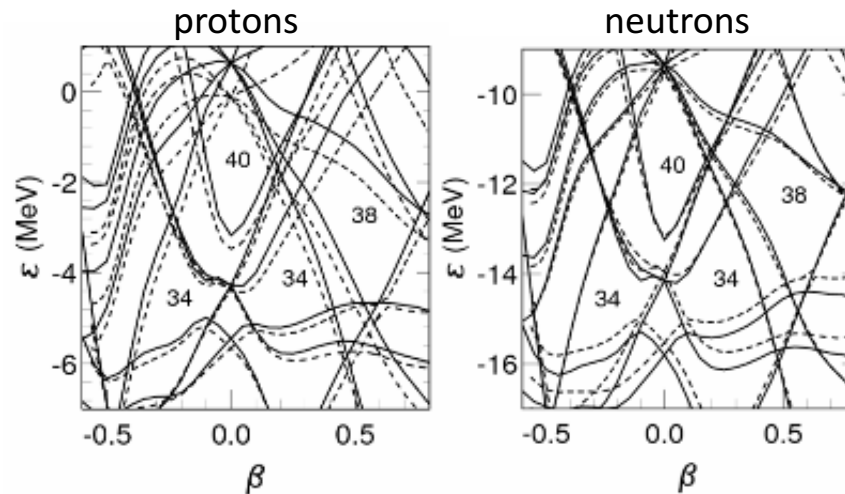
also mixing of
K=0 and K=2
states

Gogny D1S 5D GCM (GOA):
M. Girod et al., PLB 676, 39 (2009)

Skyrme SLy6 GCM:
M. Bender et al., PRC 74, 024312 (2006)



Configuration mixing calculations



Difference #1: effective interaction

Skyrme SLy6 \Leftrightarrow **Gogny D1S**
 Bender et al. Girod et al.

very similar single-particle energies
 \Rightarrow no big differences on the mean-field level

M. Girod et al., PLB 676, 39 (2009)

Difference #2: generator coordinates

axial quadrupole deformation q_0 \Leftrightarrow
 (exact GCM formalism)

- good agreement for in-band B(E2) and quadrupole moments
- wrong ordering of states: oblate ground-state shape for $^{72}\text{Kr} \rightarrow ^{78}\text{Kr}$
- K=2 states outside model

triaxial quadrupole deformation q_0, q_2
 Euler angles $\Omega=(\theta_1, \theta_2, \theta_3)$
 \Rightarrow 5-dimensional collective Hamiltonian
 (Gaussian overlap approximation)

- excellent agreement for excitation energies, B(E2), and quadrupole moments
- inversion of ground-state shape from prolate in ^{76}Kr to oblate in ^{72}Kr reproduced
- assignment of prolate, oblate, and K=2 states

\Rightarrow triaxiality seems to be important to describe prolate-oblate shape coexistence in this region

Quadrupole sum rules

method to determine the **intrinsic shape** (also for 0^+)

rotate electric quadrupole tensor into principal axis frame:

⇒ only two non-zero quadrupole moments

⇒ two parameters (Q, δ) in analogy with Bohr's parameters (β, γ)

$$\mathcal{M}(E2, \mu = 0) = Q \cos \delta$$

$$\mathcal{M}(E2, \mu = \pm 1) = 0$$

$$\mathcal{M}(E2, \mu = \pm 2) = \frac{1}{\sqrt{2}} Q \sin \delta$$

zero-coupled products of the E2 operators are **rotationally invariant**:

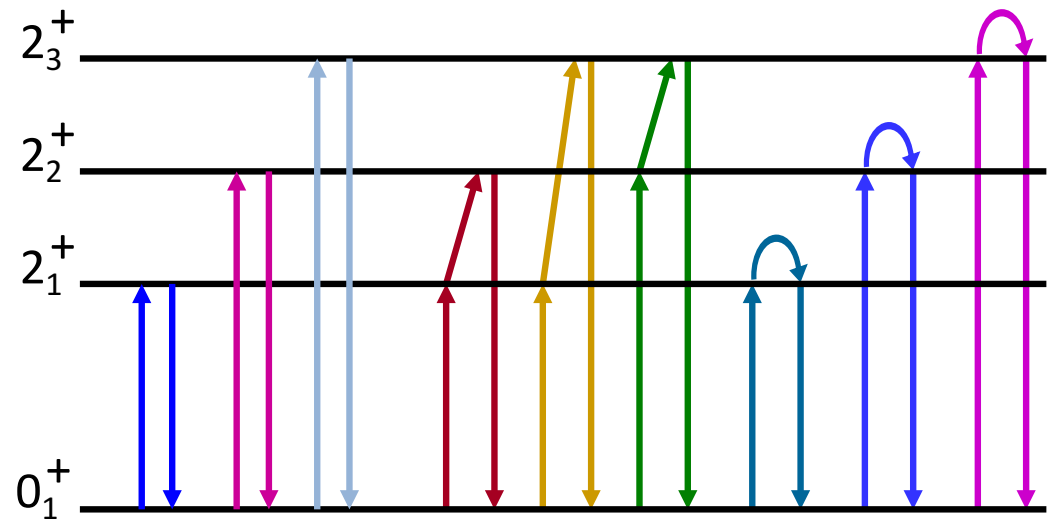
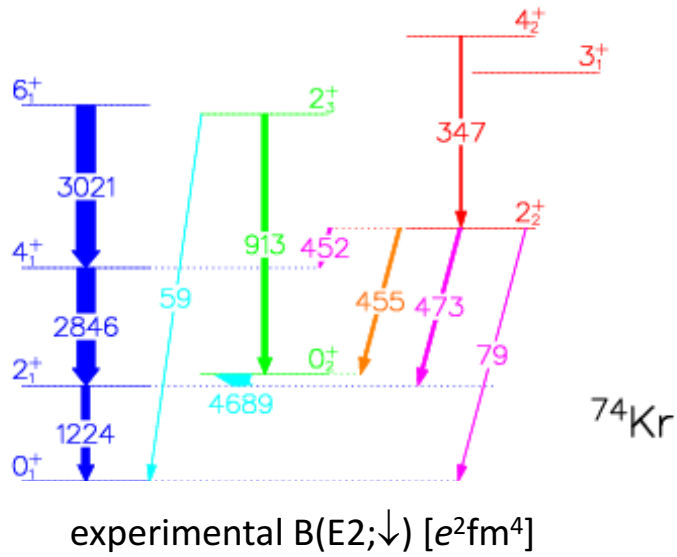
$$\langle s | [E2 \times E2]_0 | s \rangle = \frac{1}{\sqrt{5}} Q^2 = \frac{(-1)^{2s}}{\sqrt{2s+1}} \sum_t \langle s || E2 || t \rangle \langle t || E2 || s \rangle \left\{ \begin{matrix} 2 & 2 & 0 \\ s & s & t \end{matrix} \right\}$$

$$\langle s | [[E2 \times E2]_2 \times E2]_0 | s \rangle = -\sqrt{\frac{2}{35}} Q^3 \cos(3\delta) = \frac{1}{2s+1} \sum_{tu} \langle s || E2 || t \rangle \langle t || E2 || u \rangle \langle u || E2 || s \rangle \left\{ \begin{matrix} 2 & 2 & 2 \\ s & t & u \end{matrix} \right\}$$

D. Cline, *Ann. Rev. Nucl. Part. Sci.* 36, 683 (1986)

- model-independent method to determine the intrinsic quadrupole shape from a set of E2 matrix elements

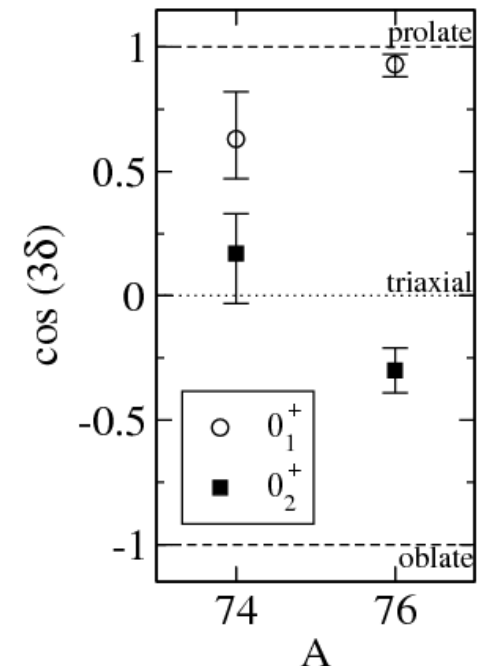
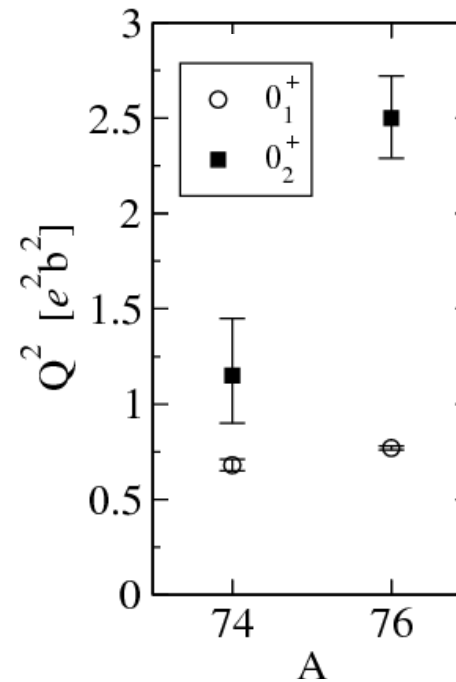
Example for ^{74}Kr , ^{76}Kr



$$\sum_t \langle s || E2 || t \rangle \langle t || E2 || s \rangle \begin{Bmatrix} 2 & 2 & 0 \\ s & s & t \end{Bmatrix}$$

$$\sum_{tu} \langle s || E2 || t \rangle \langle t || E2 || u \rangle \langle u || E2 || s \rangle \begin{Bmatrix} 2 & 2 & 2 \\ s & t & u \end{Bmatrix}$$

needs complete set of matrix elements
 \Rightarrow usually only feasible for 0^+ states



Matrix elements from Coulomb excitation

$$|0_1^+\rangle = a_0|0_p^+\rangle + b_0|0_o^+\rangle$$

$$|2_1^+\rangle = a_2|0_p^+\rangle + b_2|0_o^+\rangle$$

$$|0_2^+\rangle = -b_0|0_p^+\rangle + a_0|0_o^+\rangle$$

$$|2_2^+\rangle = -b_2|0_p^+\rangle + a_2|0_o^+\rangle$$

$$a_0^2 + b_0^2 = 1$$

$$a_2^2 + b_2^2 = 1$$

4 equations:

$$\langle 2_1^+ | E2 | 0_1^+ \rangle = b_0 b_2 \langle 2_o^+ | E2 | 0_o^+ \rangle + a_0 a_2 \langle 2_p^+ | E2 | 0_p^+ \rangle$$

$$\langle 2_1^+ | E2 | 0_2^+ \rangle = a_0 b_2 \langle 2_o^+ | E2 | 0_o^+ \rangle - a_2 b_0 \langle 2_p^+ | E2 | 0_p^+ \rangle$$

$$\langle 2_2^+ | E2 | 0_1^+ \rangle = a_2 b_0 \langle 2_o^+ | E2 | 0_o^+ \rangle - a_0 b_2 \langle 2_p^+ | E2 | 0_p^+ \rangle$$

$$\langle 2_2^+ | E2 | 0_2^+ \rangle = a_0 a_2 \langle 2_o^+ | E2 | 0_o^+ \rangle + b_0 b_2 \langle 2_p^+ | E2 | 0_p^+ \rangle$$

4 unknowns:

2 mixing amplitudes: a_0, a_2

$$Q_{0,p} = \sqrt{\frac{16\pi}{5}} \langle 2_p^+ | E2 | 0_p^+ \rangle$$

$$Q_{0,o} = \sqrt{\frac{16\pi}{5}} \langle 2_o^+ | E2 | 0_o^+ \rangle$$

no transitions between intrinsic prolate and oblate states:

$$\langle I_p^+ | E2 | J_o^+ \rangle = 0$$

using experimental coulex matrix elements for ^{74}Kr :

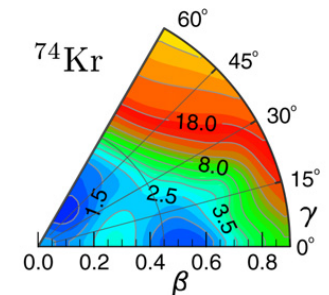
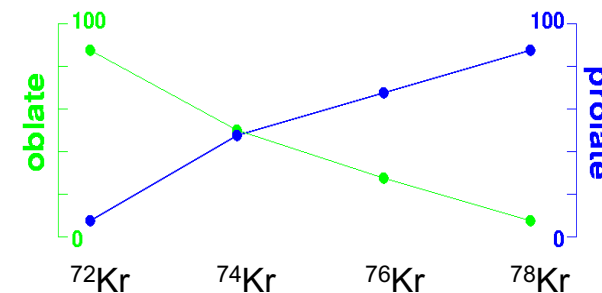
$$a_0^2 = 0.48(2)$$

$$a_2^2 = 0.82(20)$$

$$Q_{0,p} = 3.62(48) \text{ eb}$$

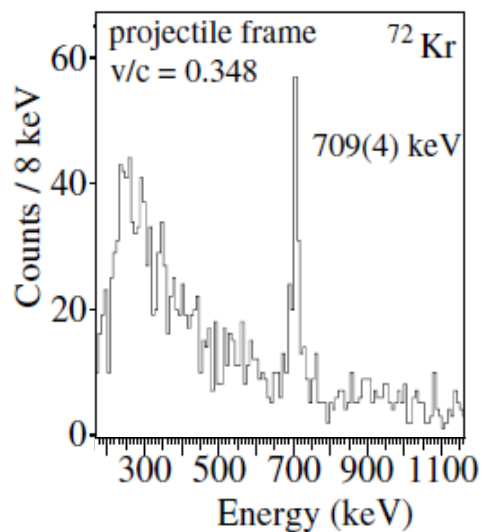
$$Q_{0,o} = -0.66(86) \text{ eb}$$

E. Clément et al.,
PRC 75, 054313 (2007)

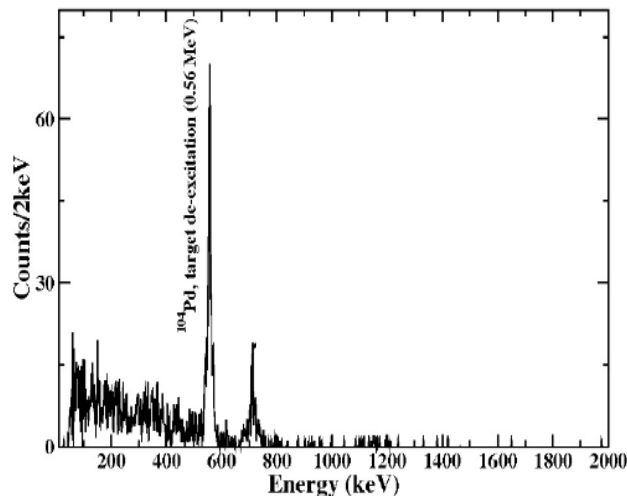


consistent with mixing obtained from E0 transitions

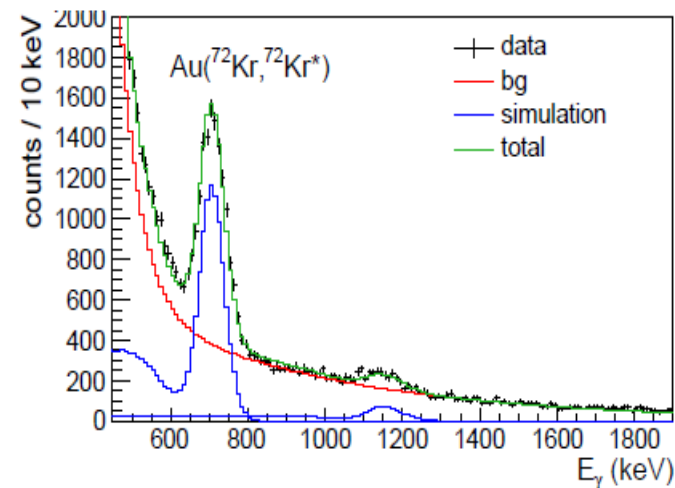
$B(E2; 0^+ \rightarrow 2^+)$ measurements for ^{72}Kr



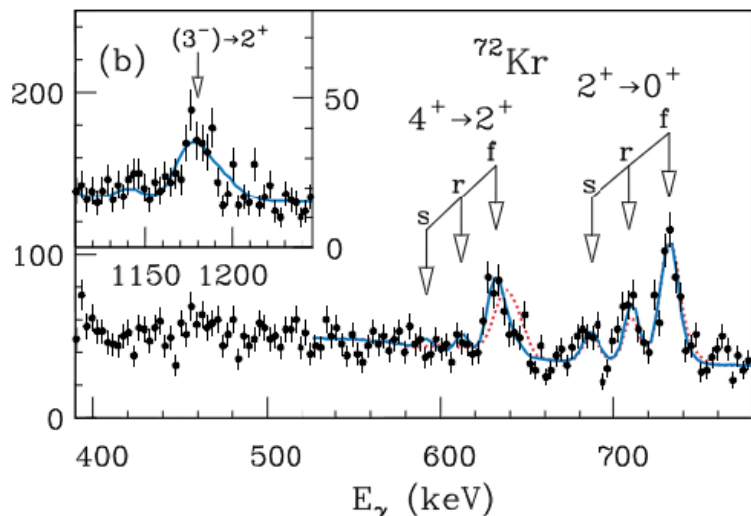
$\text{Au}(^{72}\text{Kr}, ^{72}\text{Kr}')$, 70 MeV/u, MSU
A.Gade et al., PRL 95, 022502 (2005)



$^{104}\text{Pd}(^{72}\text{Kr}, ^{72}\text{Kr}')$, 2.8 MeV/u
ISOLDE IS478 (2012)
B.S.Nara Singh et al. (unpublished)



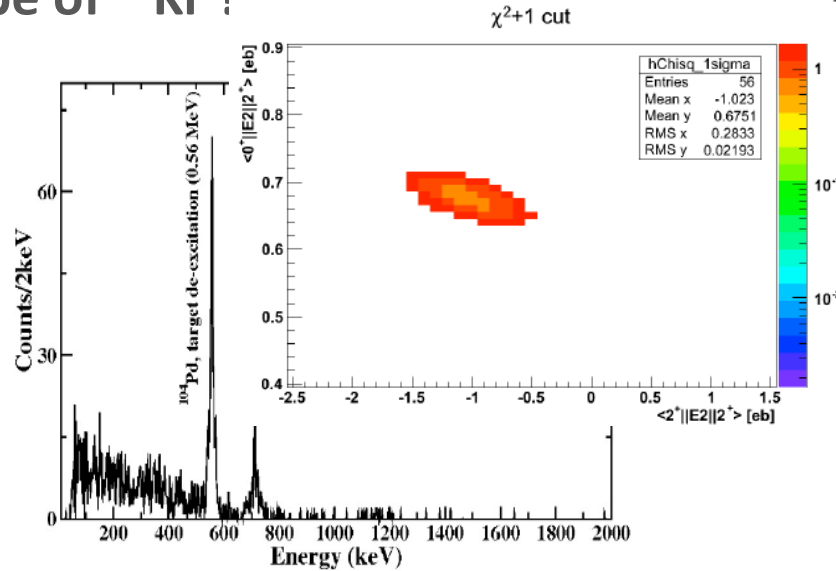
$\text{Au}(^{72}\text{Kr}, ^{72}\text{Kr}')$, 150 MeV/u, RIKEN
T.Arici, PhD Univ. Giessen (2017)



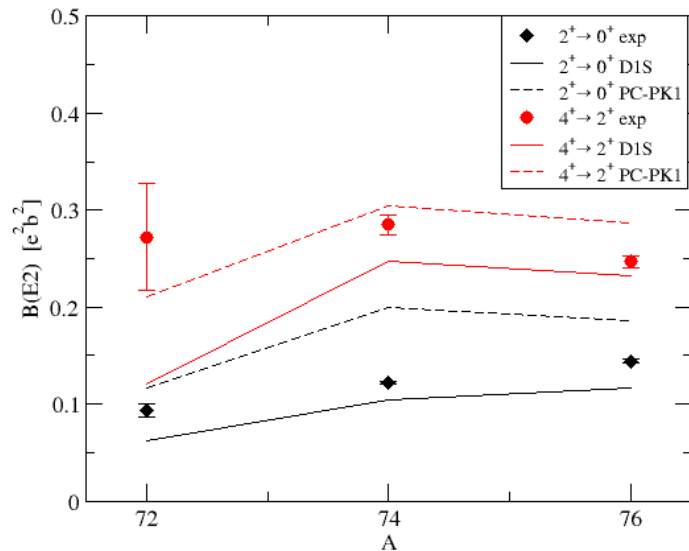
$^9\text{Be}(^{74}\text{Kr}, ^{72}\text{Kr})$, RDDS lifetime, MSU
H.Iwasaki et al., PRL 112, 142502 (2014)

$B(E2; 0^+ \rightarrow 2^+)$ [$e^2\text{fm}^4$]	Author
4997(647)	Gade
4600(600)	Nara Singh
4050(750)	Iwasaki
4910(500)	Arici

Shape of ^{72}Kr ?



B.S.Nara Singh, HIE-ISOLDE Workshop 2016



$^{74,76}\text{Kr}$: E.Clément et al., PRC 75, 054313 (2007)

^{72}Kr : H.Iwasaki et al., PRL 112, 142502 (2014)

high-energy Coulomb excitation (MSU, RIKEN)

➤ not sensitive to Q_s

low-energy Coulomb excitation (ISOLDE)

➤ sensitive to $Q_s(2^+)$

➤ very low statistics

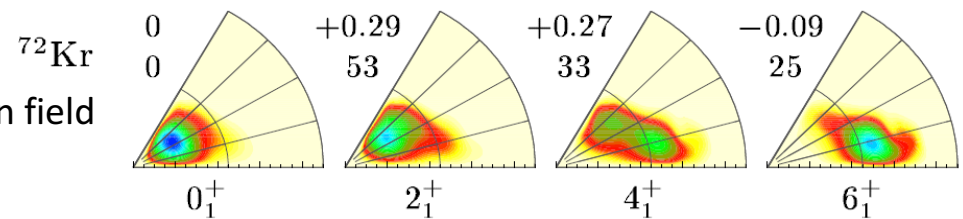
➤ favors prolate shape for 2^+ state

lifetime measurement (MSU)

➤ not sensitive to $Q_s(2^+)$

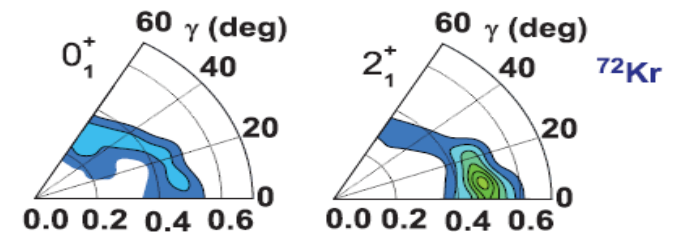
➤ large $B(E2; 4^+ \rightarrow 2^+)/B(E2; 2^+ \rightarrow 0^+)$

beyond mean field
Gogny D1S
5DCH



M.Girod et al., PLB 676, 39 (2009)

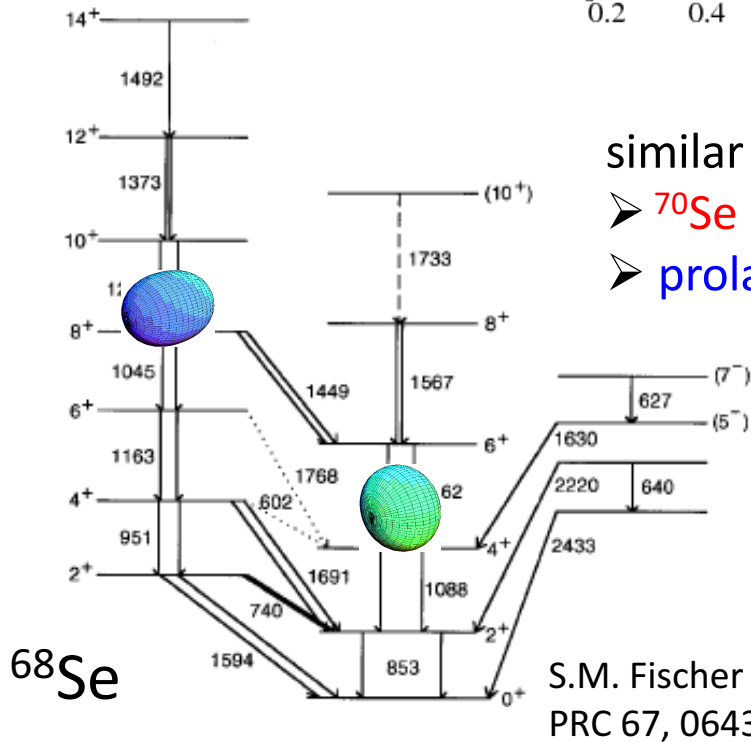
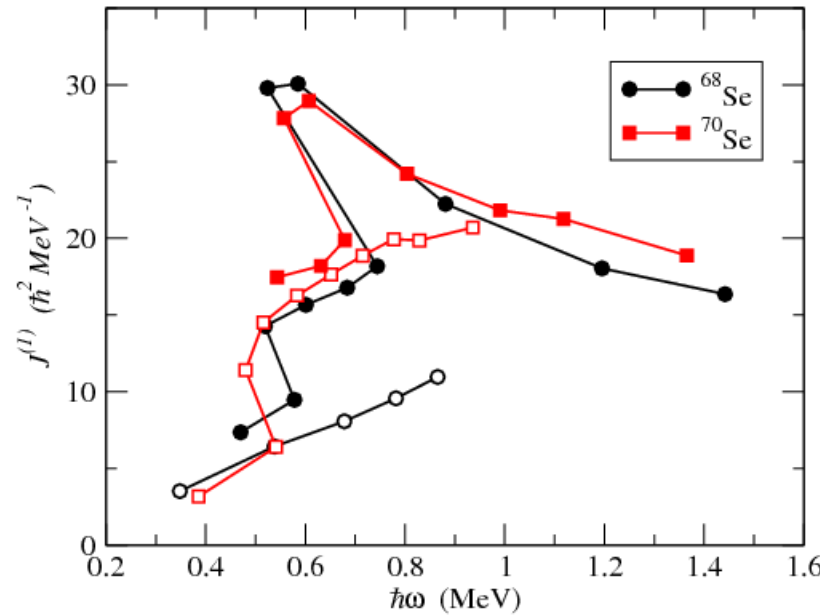
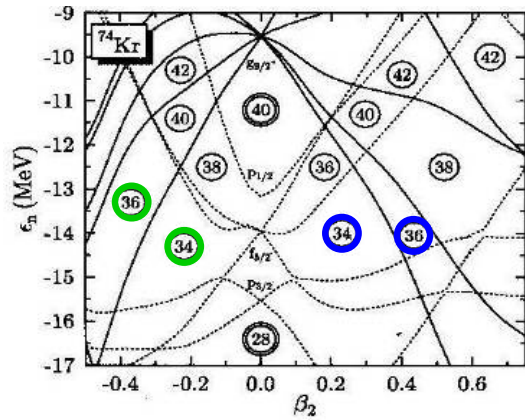
beyond mean field
RMF PC-PK1
5DCH



Y.Fu et al., PRC 87, 054305 (2013)

➤ rapid transition from oblate ground state to prolate 2^+ in ^{72}Kr ?

Shape coexistence in light Selenium isotopes

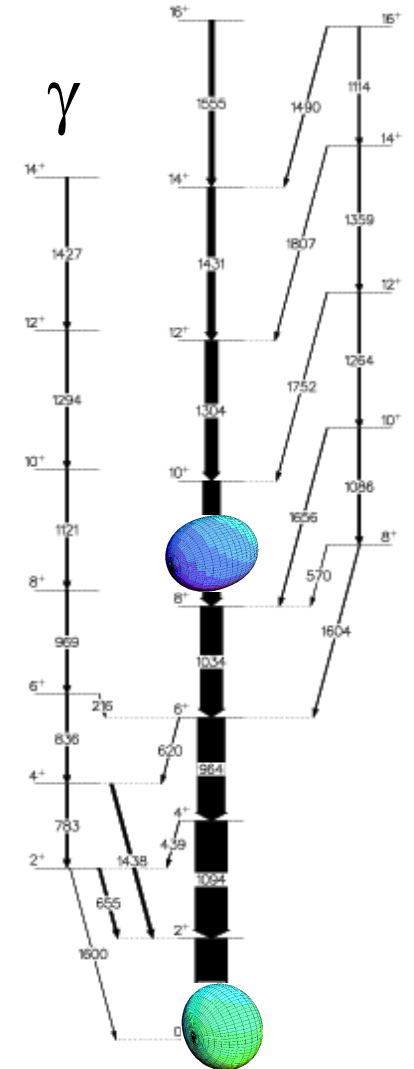


similar $J^{(1)}$ in ^{68}Se and ^{70}Se :

- ^{70}Se oblate near ground state
- prolate at higher spin

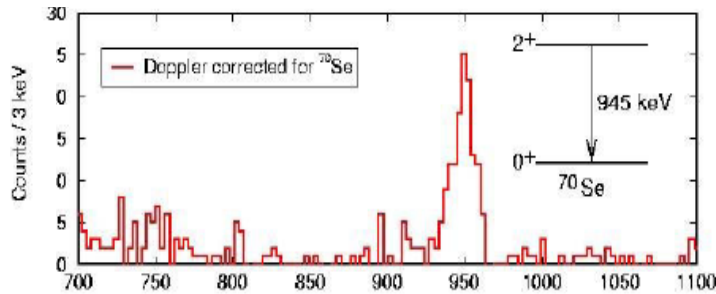
S.M. Fischer et al.,
PRC 67, 064318 (2003)

^{70}Se $\pi(g_{9/2})^2$
 $\otimes \nu(g_{9/2})^2$

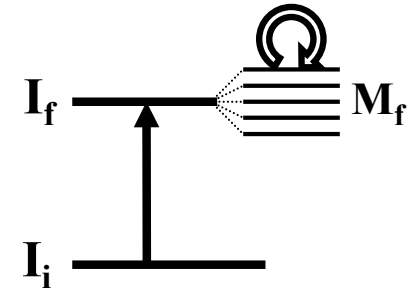


G. Rainovski et al.,
J.Phys.G 28, 2617 (2002)

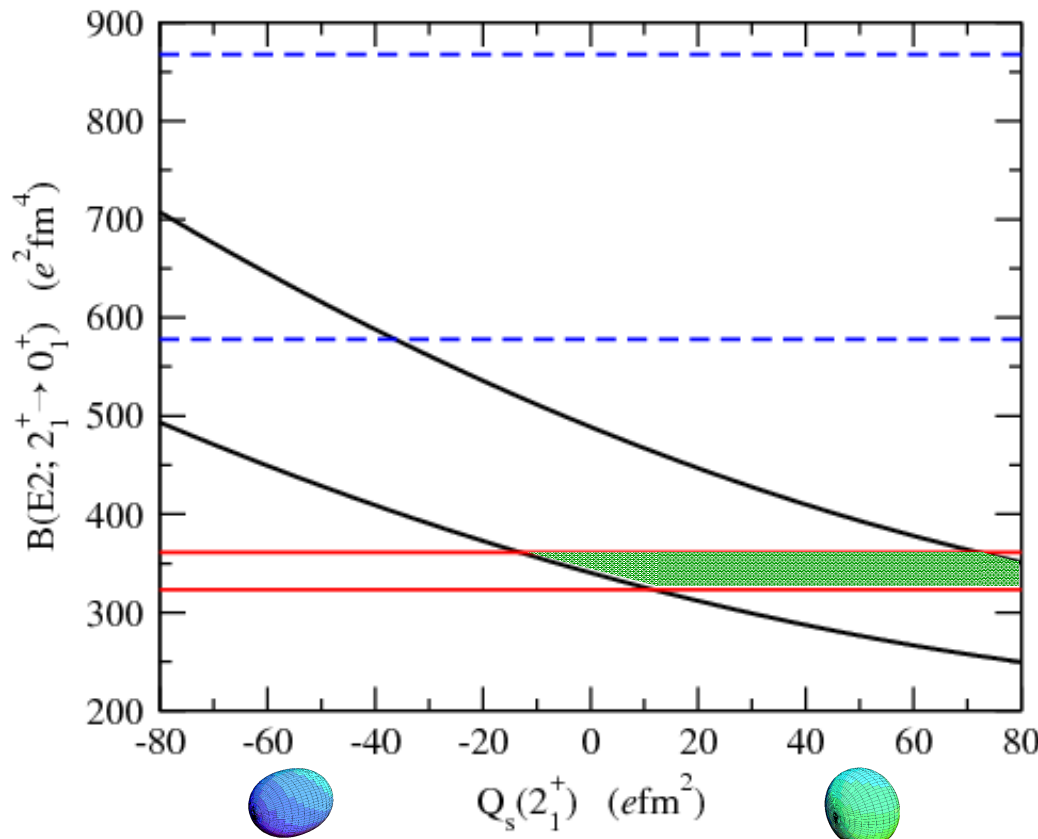
The shape of ^{70}Se



- Coulomb excitation at ISOLDE
- ^{70}Se on ^{104}Pd at 2.94 MeV/u
- integral measurement
- depends on $B(E2)$ and Q_s
- one measurement, two unknowns !



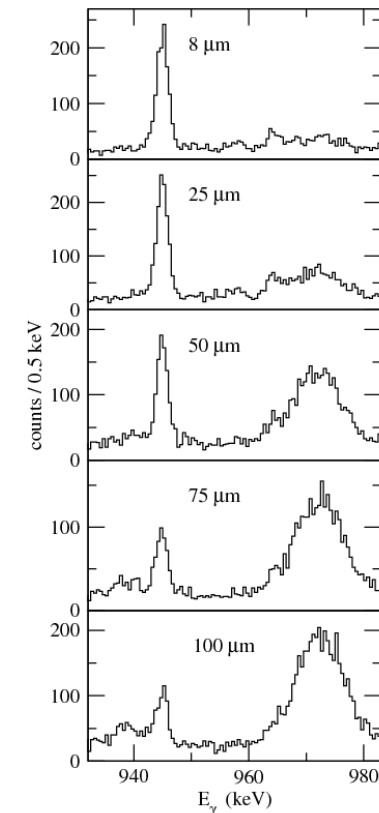
A.M. Hurst et al.,
PRL 98, 072501 (2007)



$\tau = 1.5(3)$ ps
J. Heese et al.,
Z. Phys. A 325, 45 (1986)

$\tau = 3.2(2)$ ps
J. Ljungvall et al.,
PRL 100, 102502 (2008)

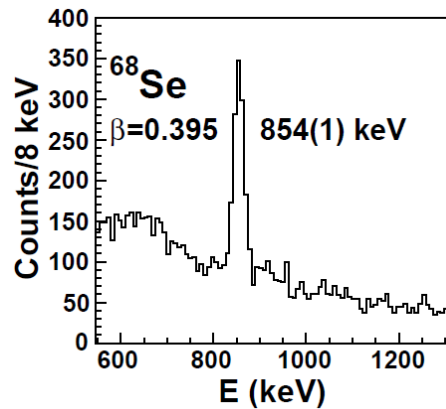
- new lifetime measurement



Shape evolution in the light Selenium isotopes

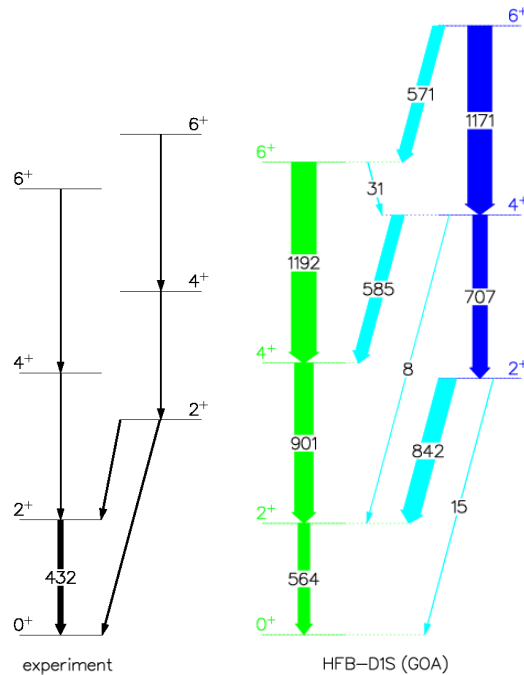
- **oblate** rotation prevails only in ^{68}Se
 ⇒ best example for shape coexistence in A=70 region?
- need experimental quadrupole moments
- **where is the 0_2^+ state?**

intermediate-energy Coulomb excitation MSU
 $B(E2; 2_1^+ \rightarrow 0_1^+) = 432(58) e^2\text{fm}^4$

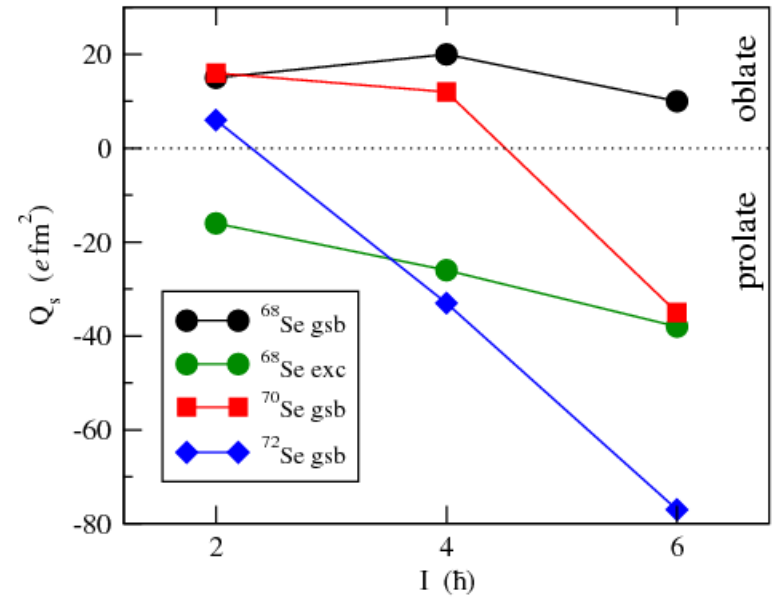


A. Obertelli et al.
 PRC 80 (2009) 031304(R)

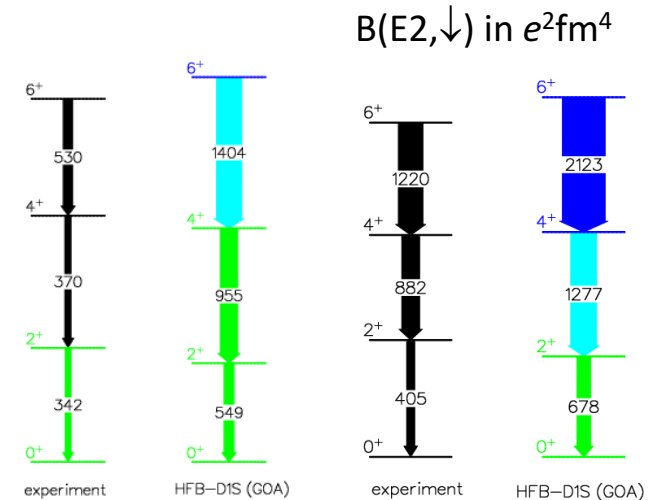
lifetime measurement at MSU
 $B(E2; \downarrow) = 392(70) e^2\text{fm}^4$
 A.J. Nichols et al., PLB 733, 52 (2014)



^{68}Se



Q_s from Gogny D1S calculation



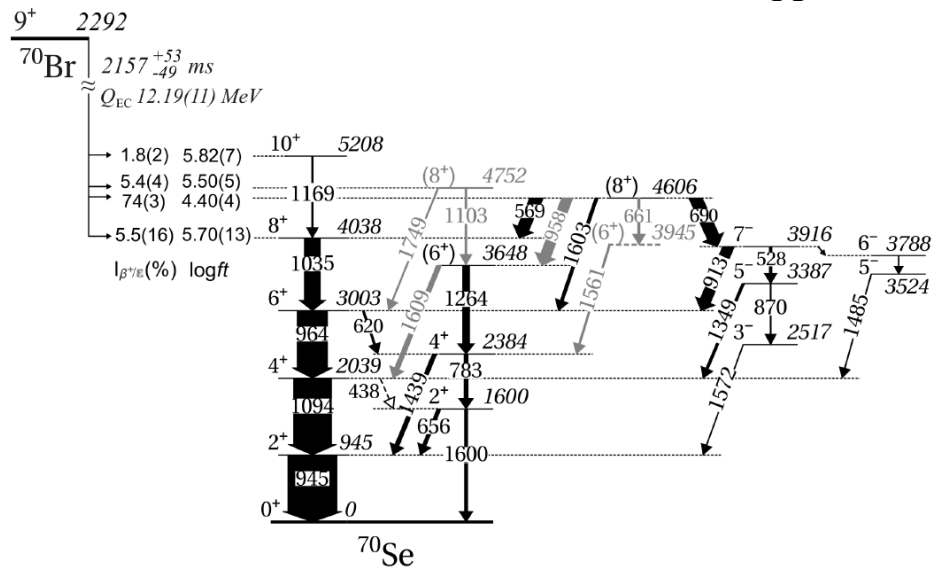
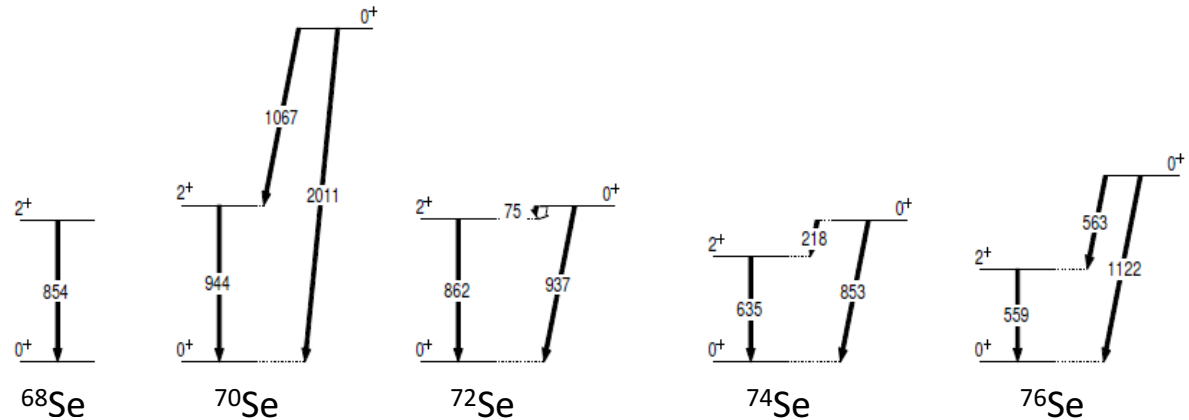
^{70}Se

^{72}Se

$B(E2, \downarrow)$ in $e^2\text{fm}^4$

0_2^+ states in Se isotopes

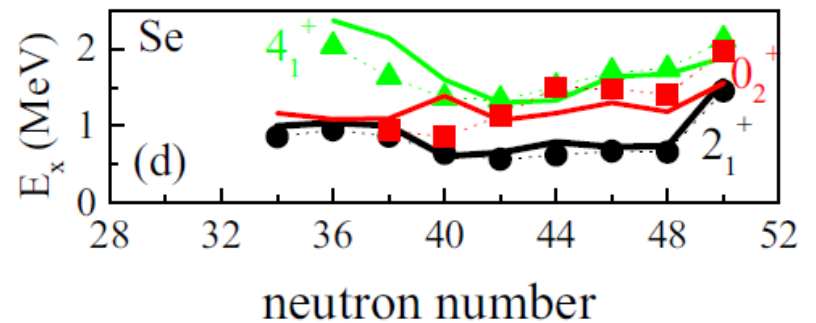
- spectroscopy of ^{70}Se following β decay of ^{70}Br at RIKEN
- no sign of low-lying 0_2^+ state



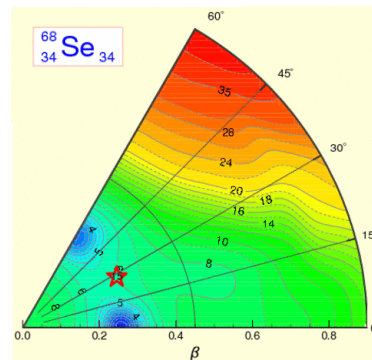
A.I. Morales et al.
 Phys.Rev. C 95, 064327 (2017)

- isomer spectroscopy at GANIL / LISE
- 0_2^+ state near or below 2_1^+ state excluded
 E. Clément et al., NIM A 587, 292 (2008)

- shell model predicts 0_2^+ states near 2_1^+ for ^{68}Se and ^{70}Se

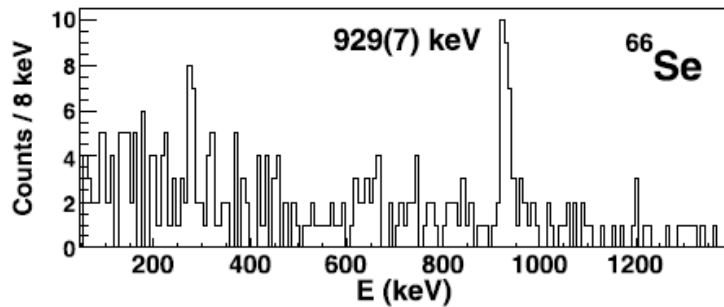


K. Kaneko et al., PRC 92, 044331 (2015)

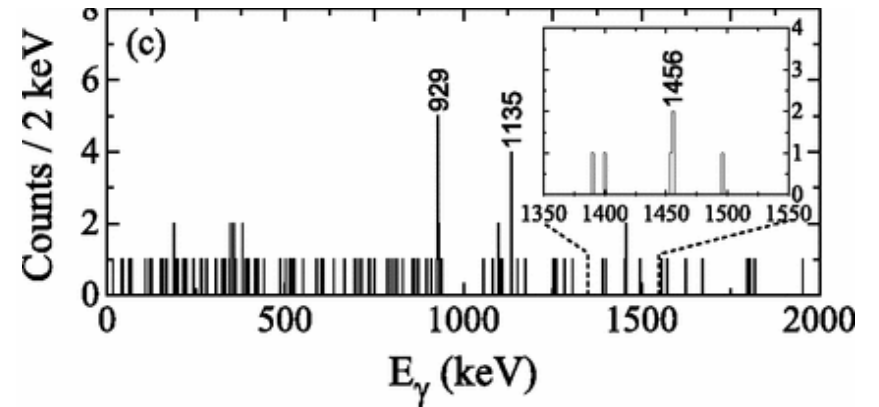
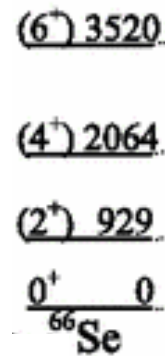


Beyond $N = Z$: Se

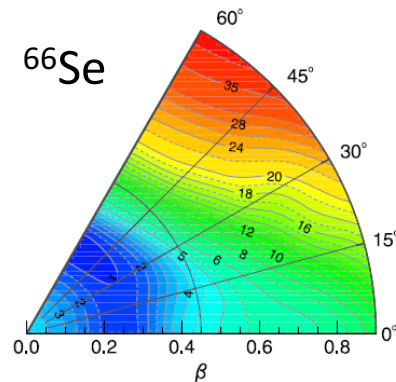
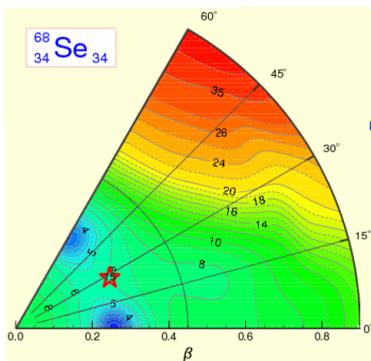
- $^9\text{Be}(^{68}\text{Se}, ^{66}\text{Se})$ @ 78 MeV
- two-nucleon knockout @ MSU
- 2_1^+ identified
- β -decay tagging @ Jyväskylä
- candidates for ground-state band up to 6^+



A. Obertelli et al.,
Phys. Lett. B 701, 417 (2011)



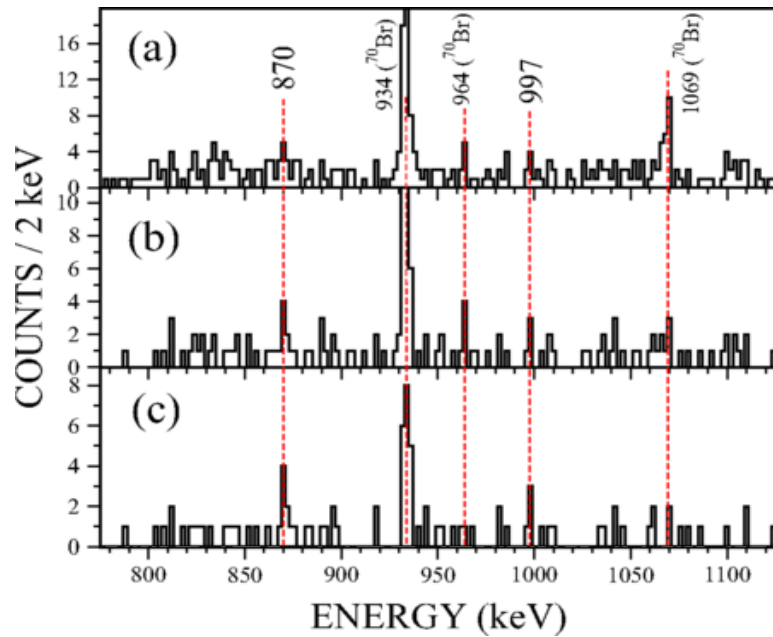
P. Ruotsalainen et al.,
Phys.Rev. C 88, 041308 (2013)



- PES (Gogny D1S) suggests vanishing barrier between oblate and prolate minima
- γ softness
- more experimental data needed

Beyond $N = Z$: Kr

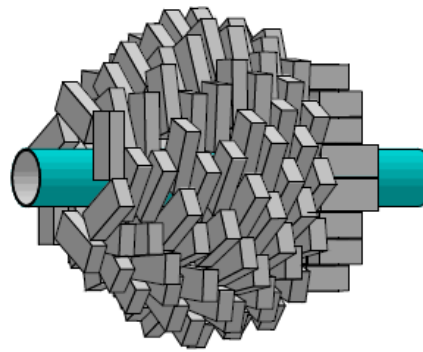
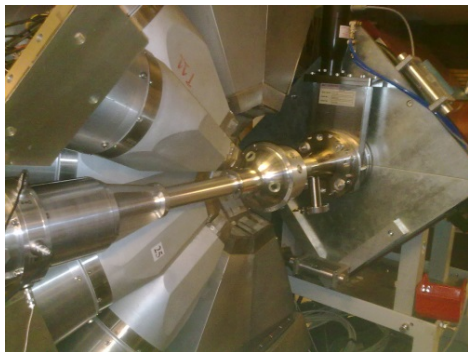
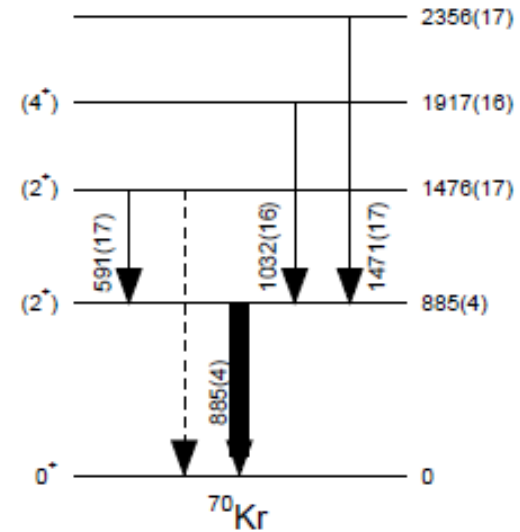
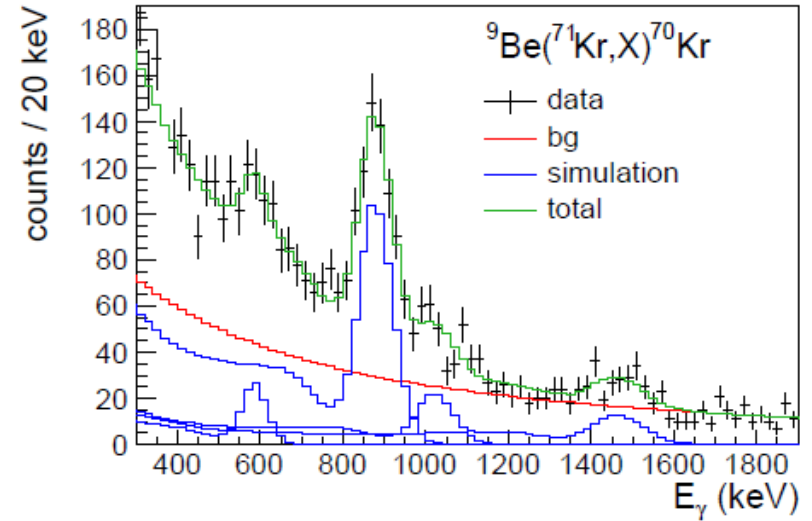
- $^{40}\text{Ca}(^{32}\text{S}, 2n)^{70}\text{Kr}$
- β -decay tagging Jyväskylä
- candidates for 2^+ and 4^+ states



D.M. Debenham *et al.*, PRC 94, 054311 (2016)

- spectroscopy following knockout @ RIKEN

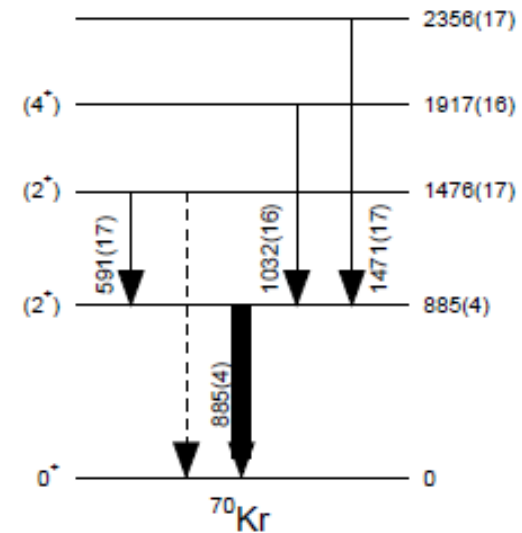
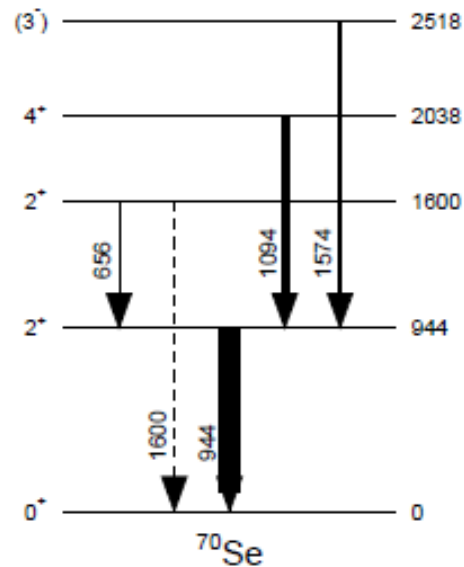
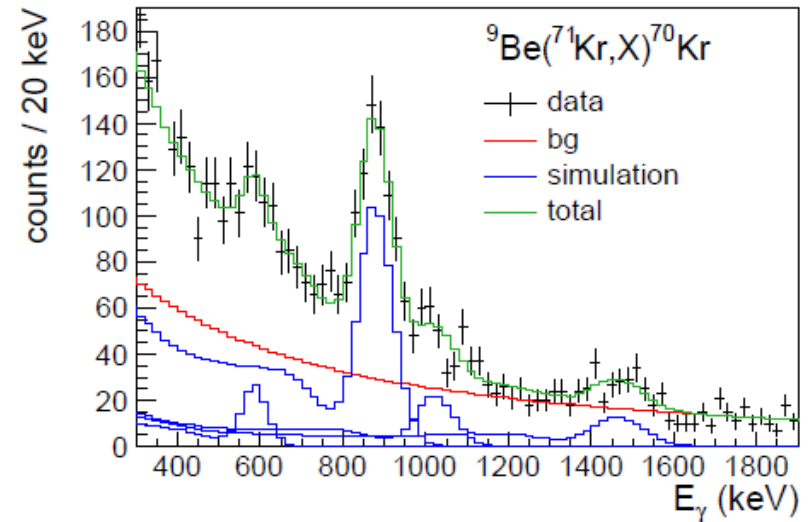
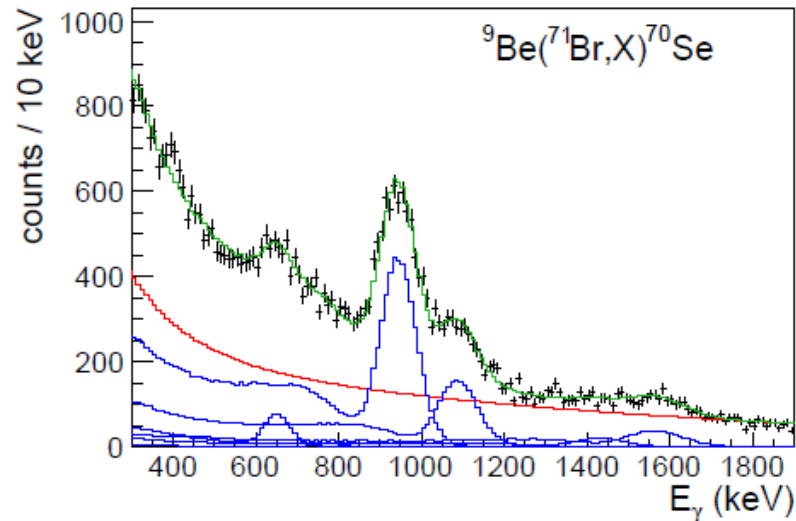
T. Arici, PhD Univ. Giessen (2017)



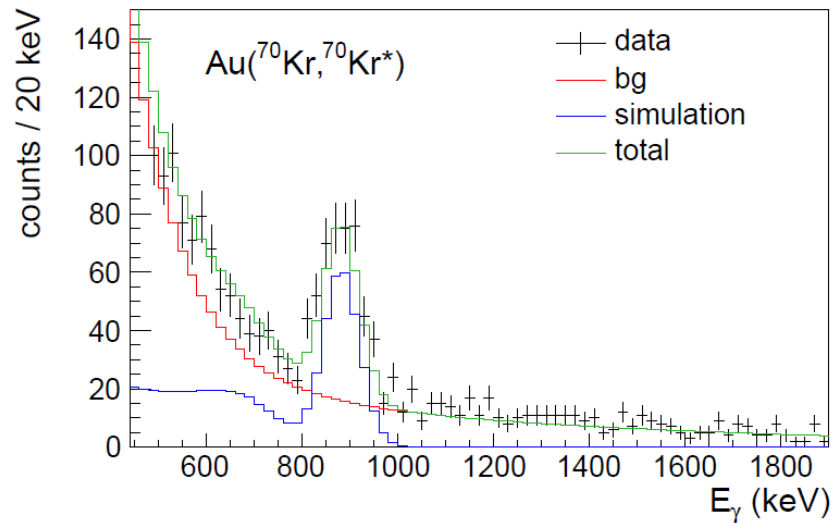
Mirror symmetry $^{70}\text{Se} - ^{70}\text{Kr}$

➤ spectroscopy following knockout @ RIKEN

T. Arici, PhD Univ. Giessen (2017)



Coulomb excitation of ^{70}Kr

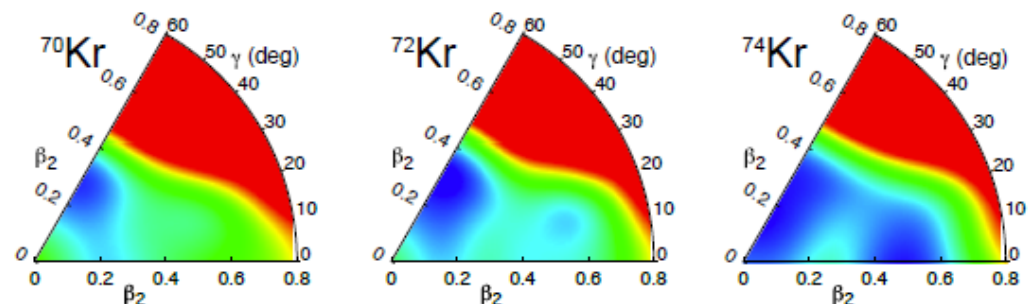
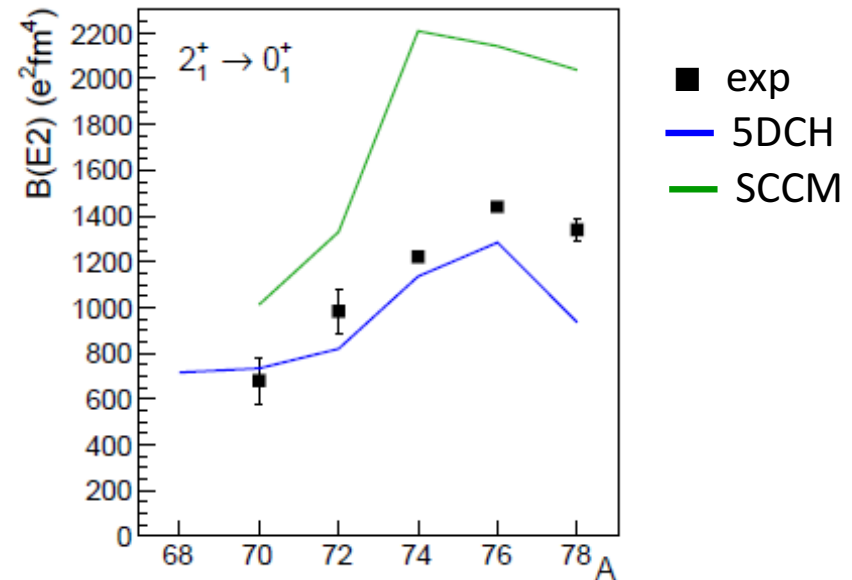


^{70}Kr : $B(E2; \downarrow) = 680(100)e^2\text{fm}^4$
 T. Arici, PhD Univ. Giessen (2017)

comparison with theory:

- 5DCH: Gogny D1S, $Q_{20} + Q_{22}$, GCM + GOA
- SCCM (symmetry-conserving configuration mixing)
 Gogny D1S, $Q_{20} + Q_{22}$, exact spin and
 particle-number restoration

T.R. Rodriguez, PRC 90, 034306 (2014)



Questions / Discussion

- Do we understand ^{72}Kr ?
 - oblate ground state shape that turns prolate already for 2^+ ?
 - remarkable agreement for Gogny 5DCH for ^{76}Kr , ^{74}Kr , (^{70}Kr) but less for ^{72}Kr ?

- Is ^{68}Se the best case for shape coexistence in this region (two rotational bands)?
- Why are there no low-lying 0_2^+ states in ^{68}Se and ^{70}Se ?

- Do we understand the large $B(E2)$ value in ^{70}Kr ?
(or the low $B(E2)$ value in ^{70}Se)

- Which experimental measurements should we focus on?