

# Density functional theory with spatial-symmetry breaking and configuration mixing

Thomas Lesinski

Espace de structure Nucléaire Théorique,  
DSM/IRFU/SPhN, CEA Saclay

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# Outline

- 1 Introduction
  - Density-functional models for nuclei
- 2 Hohenberg-Kohn scheme
  - Hohenberg-Kohn scheme
- 3 Form of the functional
  - Form of the functional
- 4 Kohn-Sham scheme
  - Kohn-Sham scheme

# Nuclear structure: methods

- Unified, microscopic theory ?
  - Perturbation theory: fail (repulsive core)
  - Ladder resummation (Brueckner): fail (involved, 3N force needed)
  - **Effective** NN+3N interactions: fail (pairing, INM EoS, ferromagnetism, parameters)
  - ➔ Effective density-dependent “interactions”: **density functionals**

$$E_{\text{mf}}(q) = \min_{\Phi_0(q)} \langle \Phi_0(q) | \hat{T} + \hat{V}_{\text{eff}}[\rho_{\Phi_0}] - \lambda \hat{Q} | \Phi_0(q) \rangle \rightarrow |\Phi_0(q)\rangle$$

- Deformed rotor/vibrator but  $[\hat{H}, \hat{J}^2] = 0$ : break/restore sym.  
M. Bender, P.-H. Heenen, P.-G. Reinhard, RMP 75, 121 (2003)
- Beyond-Mean-field scheme (Hill-Wheeler-Griffin)  
for mixing (quasi-)degenerate minima
- ...but theory with functionals is **ill defined**:
  - ① Back to interactions (or “pseudopotentials”) ?
  - ② Approximate schemes (Lipkin, GOA+collective Hamiltonian)

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# Hohenberg-Kohn-Sham scheme (quick version)

- Consider a system with Hamiltonian  $\hat{H} = \hat{T} + \hat{U} + \hat{V}$  with  $\hat{T}$  kinetic,  $\hat{U}$  interaction (NN+3N+...),  $\hat{V} \sim v(\vec{r})$  ext. potential

$$F[v] = \min_{\Psi} \langle \Psi | \hat{T} + \hat{U} + \hat{V} | \Psi \rangle$$

- Legendre transform with  $\rho(\vec{r}) = \partial E / \partial v(\vec{r})$

$$E[\rho] = \min_v \left[ F[v] - \int d^3\vec{r} v(\vec{r}) \rho(\vec{r}) \right] = \min_{\Psi \rightarrow \rho} \langle \Psi | \hat{T} + \hat{U} | \Psi \rangle$$

- $E[\rho]$  “universal” w.r.t choice of  $v$
- Kohn-Sham scheme: write

$$E[\rho] = T_s[\rho] + E_H[\rho] + E_{xc}[\rho]$$

$$T_s[\rho] = \min_{\{\phi_i\} \rightarrow \rho} \left[ -\frac{1}{2} \int d^3\vec{r} \sum_{i=1}^N \phi_i^*(\vec{r}) \Delta \phi_i(\vec{r}) \right]$$

- G.s.: minimum of  $E[\rho]$

# Definitions

- Consider a system of  $N$  particles with Hamiltonian  $\hat{H} = \hat{T} + \hat{U} + \hat{V}$  (kinetic, interaction, ext.)

$$\mathbf{R} \equiv (\vec{r}_1, \dots, \vec{r}_N), \quad d^{3N} \mathbf{R} \equiv d^3 \vec{r}_1 \dots d^3 \vec{r}_N$$

- Now consider real functions  $Q_\mu(\vec{r})$ ,  $\mu = 1 \dots n$ ,  $\underline{q} = (q_1, \dots, q_n)$

$$\hat{Q}_\mu(\mathbf{R}) \equiv \sum_i Q_\mu(\vec{r}_i)$$

$$\hat{P}(\underline{q}, \mathbf{R}) \equiv \prod_\mu \delta(\hat{Q}_\mu(\mathbf{R}) - q_\mu)$$

- $\hat{P}$  projects on an eigenspace of  $\hat{Q}$ , and

$$\hat{P}(\underline{q}, \mathbf{R}) \hat{P}(\underline{q}', \mathbf{R}) = \delta^{(n)}(\underline{q} - \underline{q}') \hat{P}(\underline{q}, \mathbf{R})$$

- Define the **generalized** density

$$D(\underline{q}, \vec{r}) \equiv N \int d^{3N} \mathbf{R} \delta^{(3)}(\vec{r} - \vec{r}_1) \hat{P}(\underline{q}, \mathbf{R}) \Psi^*(\mathbf{R}) \Psi(\mathbf{R})$$

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# Collective coordinates

- $(Q_1, Q_2, Q_3)(\vec{r}) = \frac{1}{N}(x, y, z)$ , call  $(q_1, q_2, q_3) \equiv \vec{R}$
- Now, add to  $(q_\mu)$  the inertia tensor

$$\mathbb{J} \equiv \int d^3\vec{r} \begin{pmatrix} y^2 + z^2 & -xy & -xz \\ -yx & x^2 + z^2 & -yz \\ -zx & -zy & x^2 + y^2 \end{pmatrix} d(\vec{R}, \mathbb{J}, \vec{r})$$

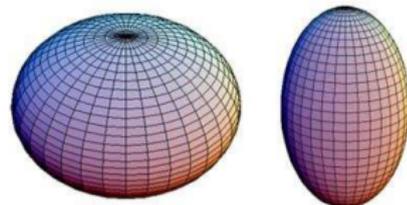
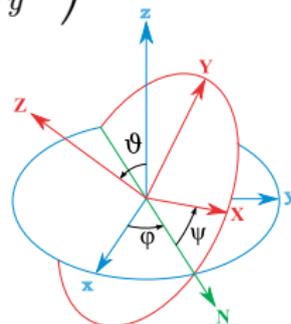
- Huygens-Steiner theorem for intrinsic inertia tensor

$$\mathbb{J}_0 = \mathbb{J} - N(R^2\mathbb{I} - \vec{R} \otimes \vec{R})$$

- $(\vec{R}, \mathbb{J}) \rightarrow (\vec{R}, r_{\text{rms}}, \beta, \gamma, \varphi, \vartheta, \psi)$ :  
Rotation + quadrupole vibrations

$$\beta \cos \gamma = \sqrt{\frac{\pi}{5}} \frac{\langle 2z^2 - x^2 - y^2 \rangle_{\text{int}}}{N r_{\text{rms}}^2}$$

$$\beta \sin \gamma = \sqrt{\frac{3\pi}{5}} \frac{\langle x^2 - y^2 \rangle_{\text{int}}}{N r_{\text{rms}}^2}$$



oblate spheroid

prolate spheroid

## External potential

- Let  $w(\underline{q}, \vec{r})$  be a real (bounded) function and

$$w(\underline{q}, \mathbf{R}) = \sum_i w(\underline{q}, \vec{r}_i) \quad \hat{W}(\mathbf{R}) = \int d^n \underline{q} w(\underline{q}, \mathbf{R}) \hat{P}(\underline{q}, \mathbf{R})$$

- We have

$$\langle \Psi | \hat{W} | \Psi \rangle = \int d^n \underline{q} \int d^3 \vec{r} w(\underline{q}, \vec{r}) D(\underline{q}, \vec{r})$$

- First, define the functional (assume non-degenerate)

$$F[w] = \min_{\Psi} \langle \Psi | \hat{T} + \hat{U} + \hat{W} | \Psi \rangle$$

- then

$$\begin{aligned} E[D] &= \min_w \left[ F[w] - \int d^n \underline{q} \int d^3 \vec{r} w(\underline{q}, \vec{r}) D(\underline{q}, \vec{r}) \right] \\ &= \min_{\Psi \rightarrow D} \langle \Psi | \hat{T} + \hat{U} | \Psi \rangle \end{aligned}$$

- Universality:  $\hat{V}$  is a special case of  $\hat{W}$  ( $w(\underline{q}, \vec{r}) = v(\vec{r})$ )

# Energy functional

- Define the **collective w.f.**, and  $\underline{q}$ -dependent density

$$f(\underline{q}) \equiv e^{i\theta(\underline{q})} \left[ \frac{1}{N} \int d^3\vec{r} D(\underline{q}, \vec{r}) \right]^{1/2}$$

$$d(\underline{q}, \vec{r}) \equiv |f(\underline{q})|^{-2} D(\underline{q}, \vec{r})$$

➔  $E[D] = E[f, d]$

- $\underline{q}$ -dependent wave function (“slice”)

$$\Psi(\underline{q}, \mathbf{R}) = f^{-1}(\underline{q}) \hat{P}(\underline{q}, \mathbf{R}) \Psi(\mathbf{R})$$

$$\int d^{3N}\mathbf{R} \Psi^*(\underline{q}, \mathbf{R}) \Psi(\underline{q}', \mathbf{R}) = \delta^{(n)}(\underline{q} - \underline{q}')$$

- $d(\underline{q}, \vec{r})$  is the density of  $\Psi(\underline{q}, \mathbf{R})$

$$\int d^{3N}\mathbf{R} \delta^{(3)}(\vec{r} - \vec{r}_1) \Psi^*(\underline{q}, \mathbf{R}) \Psi(\underline{q}', \mathbf{R}) = \delta^{(n)}(\underline{q} - \underline{q}') d(\underline{q}, \vec{r})$$

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# Properties

## ■ Normalisation

$$\int d^n \underline{q} f^*(\underline{q}) f(\underline{q}) = 1$$

$$\forall \underline{q}, \int d^3 \vec{r} d(\underline{q}, \vec{r}) = N$$

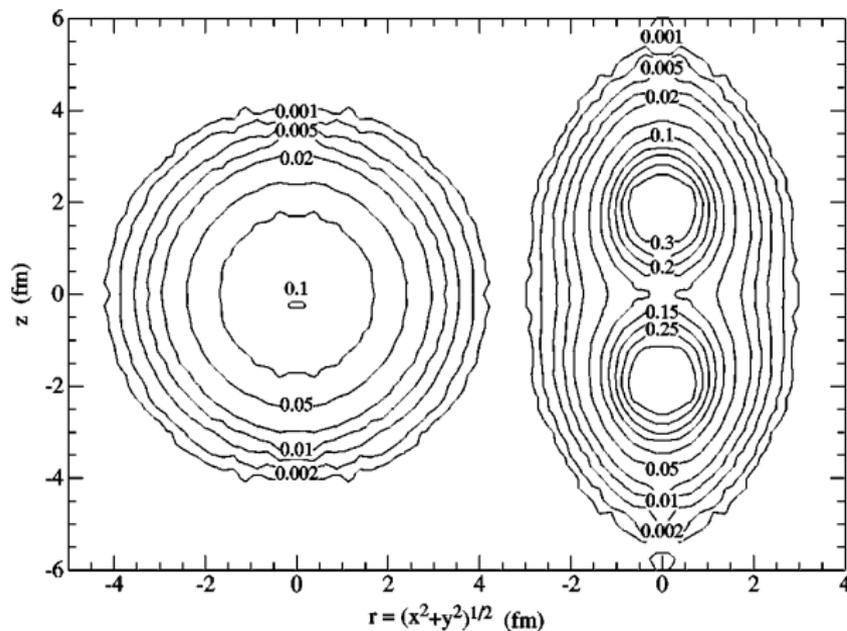
$$\forall \vec{r}, \int d^n \underline{q} D(\underline{q}, \vec{r}) = \rho(\vec{r})$$

$$\int d^n \underline{q} \int d^3 \vec{r} D(\underline{q}, \vec{r}) = N$$

## ■ Verify that

$$\int d^3 \vec{r} Q_\mu(\vec{r}) d(\underline{q}, \vec{r}) = q_\mu$$

## Internal/intrinsic density



- $^8\text{Be}$ , AV18+UIX  
GFMC
  - Left:  $\vec{R}$
  - Right:  $\vec{R}, \varphi, \vartheta$
- Wiringa, Pieper,  
Carlson,  
Pandharipande,  
Phys. Rev. C 62,  
014001 (2000)

FIG. 15. Contours of constant density, plotted in cylindrical coordinates, for  $^8\text{Be}(0^+)$ . The left side is in the “laboratory” frame while the right side is in the intrinsic frame.

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# Collective coordinates and symmetries

- Assume...
  - $f$  given by symmetry
  - $d(\underline{q}, \vec{r})$  and  $d(\underline{q}', \vec{r})$  related by sym. transformation
  
- ... then  $E[f, d] = E[\rho_{\text{int}}]$ , with  $\rho_{\text{int}}(\vec{r}) = d(\underline{Q}, \vec{r})$ 
  - Functional of the density of “pinned down” w.f.
  - $(Q_1, Q_2, Q_3)(\vec{r}) = \frac{1}{N}(x, y, z)$ : internal-frame DFT
  - ➔ Messud, Bender, Suraud, Phys.Rev.C 80 054314 (2009)
  
- Note:  $D(\underline{q}, \vec{r})$  conserves symmetries !

$$D(\underline{\mathcal{S}}(\underline{q}), \vec{\mathcal{S}}(\vec{r})) = D(\underline{q}, \vec{r})$$

# Collective equation ?

- Rewrite w.f.  $\Psi(\mathbf{R})$  as

$$\Psi(\mathbf{R}) = \int d^n \underline{q} f(\underline{q}) \Psi(\underline{q}, \mathbf{R})$$

- Assume  $\hat{U}$  local

$$E[f, d] = \langle \Psi[f, d] | \hat{T} + \hat{U} + \hat{V} | \Psi[f, d] \rangle$$

$$E[f, d] = \int d^n \underline{q} f^*(\underline{q}) \left[ -\frac{1}{2} \sum_{\mu\nu} \partial_\mu \mathcal{A}_{\mu\nu}(\underline{q}) \partial_\nu + \mathcal{U}(\underline{q}) - \frac{i}{2} \sum_{\mu} (\partial_\mu \mathcal{V}_\mu(\underline{q}) + \mathcal{V}_\mu(\underline{q}) \partial_\mu) \right] f(\underline{q})$$

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## Collective equation ?

### ■ Collective mass

$$\mathcal{A}_{\mu\nu}(\underline{q}) = \frac{1}{\langle \hat{P}(\underline{q}) \rangle} F_{\mu\nu}(\underline{q})$$

$$F_{\mu\nu}(\underline{q}) = \int d^{3N} \mathbf{R} \hat{P}(\underline{q}, \mathbf{R}) [\vec{\nabla} \hat{Q}_\mu(\mathbf{R})] \cdot [\vec{\nabla} \hat{Q}_\nu(\mathbf{R})] \Psi^*(\mathbf{R}) \Psi(\mathbf{R})$$

### ■ Current term ( $\langle \hat{P}(\underline{q}) \rangle = |f(\underline{q})|^2$ )

$$\mathcal{V}_\mu(\underline{q}) \equiv \sum_\nu F_{\mu\nu}(\underline{q}) \frac{\partial_\nu \theta(\underline{q})}{\langle \hat{P}(\underline{q}) \rangle} + \frac{J_\mu(\underline{q})}{\langle \hat{P}(\underline{q}) \rangle}$$

$$J_\mu(\underline{q}) \equiv \frac{i}{2} \int d^{3N} \mathbf{R} \hat{P}(\underline{q}, \mathbf{R}) \vec{\nabla} \hat{Q}_\mu(\mathbf{R}) \cdot [\vec{\nabla} \Psi^*(\mathbf{R}) \Psi(\mathbf{R}) - \Psi^*(\mathbf{R}) \vec{\nabla} \Psi(\mathbf{R})]$$

### ■ Choose $\theta(\underline{q})$ to make $f$ continuous, and/or cancel $\mathcal{V}$ with

$$\sum_\nu F_{\mu\nu}(\underline{q}) \partial_\nu \theta(\underline{q}) = -J_\mu(\underline{q})$$

# Collective equation ?

## ■ Collective potential

$$\begin{aligned}
 \mathcal{U}(\underline{q}) \equiv & \sum_{\mu\nu} F_{\mu\nu}(\underline{q}) \left[ \frac{1}{2} \frac{\partial_\mu \theta(\underline{q}) \partial_\nu \theta(\underline{q})}{\langle \hat{P}(\underline{q}) \rangle} + \frac{1}{8} \frac{\partial_\mu \langle \hat{P}(\underline{q}) \rangle \partial_\nu \langle \hat{P}(\underline{q}) \rangle}{\langle \hat{P}(\underline{q}) \rangle^3} \right] \\
 & + \frac{1}{4} \sum_{\mu\nu} \partial_\mu \left[ \frac{F_{\mu\nu}(\underline{q})}{\langle \hat{P}(\underline{q}) \rangle^2} \partial_\nu \langle \hat{P}(\underline{q}) \rangle \right] - \frac{1}{2} \sum_{\mu} \frac{J_{\mu}(\underline{q}) \partial_\mu \theta(\underline{q})}{\langle \hat{P}(\underline{q}) \rangle} \\
 & + \frac{1}{\langle \hat{P}(\underline{q}) \rangle} \int d^{3N} \mathbf{R} \hat{P}(\underline{q}, \mathbf{R}) \Psi^*(\mathbf{R}) \left[ -\frac{1}{2} \hat{\Delta} + \hat{U}(\mathbf{R}) \right] \Psi(\mathbf{R}) \\
 & + \int d^3 \vec{r} v_{\text{ext}}(\vec{r}) d(\underline{q}, \vec{r})
 \end{aligned}$$

➔ Minimize energy: collective Schrödinger equation

$$\left[ -\frac{1}{2} \sum_{\mu\nu} \partial_\mu \mathcal{A}_{\mu\nu}[f, d](\underline{q}) \partial_\nu + \mathcal{U}[f, d](\underline{q}) + \mathcal{U}_{\text{tra}}[f, d](\underline{q}) - E' \right] f(\underline{q}) = 0$$

■ For ex. if  $\underline{q}$  is  $\vec{R}$ ,  $\mathcal{A}_{\mu\nu}[f, d](\underline{q}) = \frac{1}{N}$  and  $\mathcal{U}[f, d](\underline{q}) = E_{\text{int}}$

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# Collective equation ?

## ■ Collective potential

$$\begin{aligned}
 \mathcal{U}(\underline{q}) \equiv & \sum_{\mu\nu} F_{\mu\nu}(\underline{q}) \left[ \frac{1}{2} \frac{\partial_{\mu}\theta(\underline{q}) \partial_{\nu}\theta(\underline{q})}{\langle \hat{P}(\underline{q}) \rangle} + \frac{1}{8} \frac{\partial_{\mu}\langle \hat{P}(\underline{q}) \rangle \partial_{\nu}\langle \hat{P}(\underline{q}) \rangle}{\langle \hat{P}(\underline{q}) \rangle^3} \right] \\
 & + \frac{1}{4} \sum_{\mu\nu} \partial_{\mu} \left[ \frac{F_{\mu\nu}(\underline{q})}{\langle \hat{P}(\underline{q}) \rangle^2} \partial_{\nu}\langle \hat{P}(\underline{q}) \rangle \right] - \frac{1}{2} \sum_{\mu} \frac{J_{\mu}(\underline{q}) \partial_{\mu}\theta(\underline{q})}{\langle \hat{P}(\underline{q}) \rangle} \\
 & + \frac{1}{\langle \hat{P}(\underline{q}) \rangle} \int d^{3N} \mathbf{R} \hat{P}(\underline{q}, \mathbf{R}) \Psi^*(\mathbf{R}) \left[ -\frac{1}{2} \hat{\Delta} + \hat{U}(\mathbf{R}) \right] \Psi(\mathbf{R}) \\
 & + \int d^3 \vec{r} v_{\text{ext}}(\vec{r}) d(\underline{q}, \vec{r})
 \end{aligned}$$

➔ Minimize energy: collective Schrödinger equation

$$\left[ -\frac{1}{2} \sum_{\mu\nu} \partial_{\mu} \mathcal{A}_{\mu\nu}[f, d](\underline{q}) \partial_{\nu} + \mathcal{U}[f, d](\underline{q}) + \mathcal{U}_{\text{tra}}[f, d](\underline{q}) - E' \right] f(\underline{q}) = 0$$

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# Outline

- 1 Introduction
  - Density-functional models for nuclei
- 2 Hohenberg-Kohn scheme
  - Hohenberg-Kohn scheme
- 3 Form of the functional
  - Form of the functional
- 4 Kohn-Sham scheme
  - Kohn-Sham scheme

## Kohn-Sham scheme

- Write the collective potential  $\mathcal{U}$  as (def  $\rho_{\underline{q}}(\vec{r}) \equiv d(\underline{q}, \vec{r})$ )

$$\mathcal{U}[f, d](\underline{q}) = T_s[\rho_{\underline{q}}] + \mathcal{U}^{\text{ext}}[f, d](\underline{q}) + \mathcal{U}^{\text{ic}}[f, d](\underline{q})$$

$$T_s[\rho_{\underline{q}}] = \min_{\{\phi_i\} \rightarrow \rho_{\underline{q}}} \left[ -\frac{1}{2} \int d^3\vec{r} \sum_{i=1}^N \phi_i^*(\underline{q}; \vec{r}) \Delta \phi_i(\underline{q}; \vec{r}) \right]$$

- Kohn-Sham equation

$$\frac{\delta \left[ E - |f(\underline{q})|^2 \varepsilon_k(\underline{q})(\underline{q}k|\underline{q}k) - |f(\underline{q})|^2 \sum_{\mu} \lambda_{\mu}(Q_{\mu}|\rho_{\underline{q}}) \right]}{\delta \phi_k^*(\underline{q}; \vec{r})} =$$

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- $v(\underline{q}; \vec{r})$  optimizes energy: functional derivatives of  $\mathcal{U}$ ,  $\mathcal{A}$

$$v_s(\underline{q}; \vec{r}) = |f(\underline{q})|^{-2} \int d^n \underline{q}' f^*(\underline{q}') \left[ -\partial'_{\mu} \frac{1}{2} \frac{\delta \mathcal{A}_{\mu\nu}(\underline{q}')}{\delta d(\underline{q}, \vec{r})} \partial'_{\nu} + \frac{\delta \mathcal{U}^{\text{ic+ext}}(\underline{q}')}{\delta d(\underline{q}, \vec{r})} \right] f(\underline{q}')$$

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## ■ Nothing new ! Collective Hamiltonian

Review: Próchniak, Rohoziński, J. Phys. G 36, 123101 (2009)

■ Equation for quantum vibrating droplet (Bohr)

■  $\mathcal{A}_{\mu\nu}$  ? Inglis(-Belyaev), ATDHFB,...

■ Approximation to Hill-Wheeler equations (GCM-GOA)

## ■ $\delta\mathcal{A}_{\mu\nu}(\underline{q}')/\delta d(\underline{q}, \vec{r})$ term: use OEP

## ■ Assume...

■ Weak feedback from  $f$  (neglect above)

■  $\mathcal{U}[f, d](\underline{q})$  depends on  $d(\underline{q}', \vec{r})$  only, for  $\underline{q}' = \underline{q}$  only

⇒  $\mathcal{U}[f, d](\underline{q}) = \mathcal{U}[\rho_{\underline{q}}]$

⇒  $v_s(\underline{q}, \vec{r}) = \delta\mathcal{U}/\delta\rho_{\underline{q}}(\vec{r})$

⇒  $\mathcal{U}^{\text{ic}}[\rho_{\underline{q}}]$ : Skyrme/Gogny/relativistic/... functional

⇒ Collective equation can be solved a posteriori

## ■ But in general, coupling between slow (collective) and fast (s.p.) d.o.f.

Beyond adiabatic approximation

■ NB(1): Any prescription for  $\mathcal{A}$ ,  $\mathcal{U}$  is part of the functional definition

■ NB(2): Minimum of  $\mathcal{U}$  of no physical significance.

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# Summary and outlook

- Generalized DFT: we can obtain a model with
  - A collective Hamiltonian
  - ...coupled to single particle degrees of freedom
  - ...potentially exact
  - ...from first principles
  - ...that looks just like a Bohr Hamiltonian
  - ➔ TL, arXiv:1301.0807 (2013)
  
- Determine the functional  $(\mathcal{A}, \mathcal{U})[f, d] ? \dots$