

Optimization and theoretical uncertainties

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Acknowledgements

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S. M. Wild	Argonne
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J. A. Maruhn	Frankfurt
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Outline

1 Models and observables

2 Optimization of model parameters and estimate of uncertainties

- Error propagation – covariance analysis
- Strategies to estimate errors (systematic and statistical)
- 3 short examples: trends of residuals, compare with data, far extrapolations

3 More detailed results for 3 strategies: comparison of models, trends, covariances

- Comparison of models
- Dedicated variation of NMP: example symmetry energy
- Covariances

4 Conclusions

Models and observables

The four models in the contest

(T = isospin: $T=0 \leftrightarrow$ isoscalar, $T=1 \leftrightarrow$ isovector)

Skyrme Hartree Fock (SHF):

point couplings $(T=0 \text{ density})^2$, $(T=0 \text{ kinetic dns.})^2$, $(T=0 \text{ l*s dns.})^2$, $(T=0 \nabla \text{ dns.})^2$
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simple density dependence for $(\text{density})^2$ terms

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simple density dependence for (density)² terms

RMF – density-dependent point-coupling (DD-PC):

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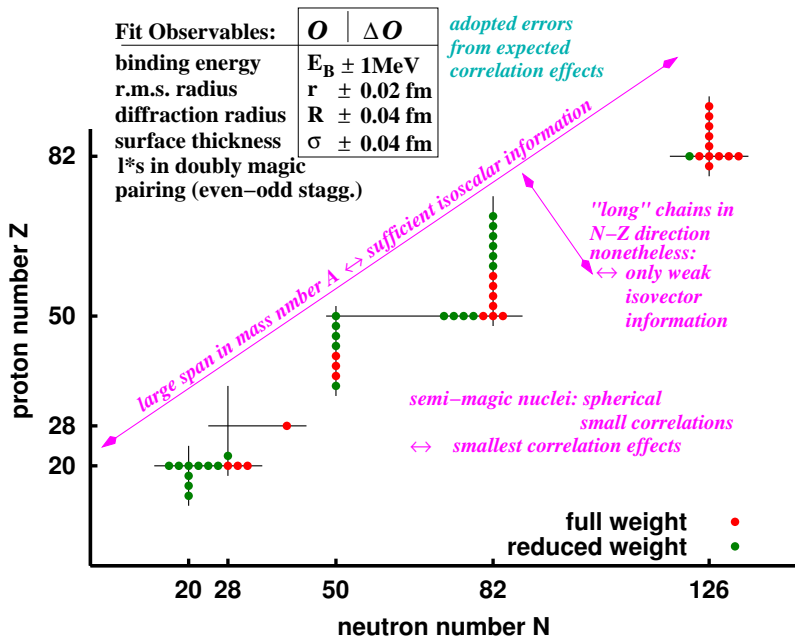
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all models: 6–11 free parameters, to be determined empirically by fit to data

Observables in the pool of fit data



The nuclear matter parameters (NMP)

given $E/A(\rho)$ = energy per particle in symmetric nuclear matter (ρ = total density)
this allows to define basic properties near equilibrium:

E/A_{eq}	binding energy per particle at equilibrium point
ρ_{eq}	equilibrium density
$K = 9\rho_0^2 \partial_\rho^2 \frac{E}{A}$	incompressibility (isoscalar static response)
$\frac{m^*}{m}$	effective mass (isoscalar dynamic response)
$J = a_{\text{sym}}$	symmetry energy (isovector static response)
κ_{TRK}	TRK sum rule enhancement \leftrightarrow isovector $\frac{m_1^*}{m}$ (dynamic response)
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2 interpretations for NMP: extrapolation to huge systems (matter, neutron stars)
equivalent to part of the model parameters (volume params.)

“Predicted” observables in the forthcoming test cases

- neutron r.m.s. radius r_n in ^{208}Pb
- neutron skin $r_n - r_p$ in ^{208}Pb
- weak-charge formfactor $F_W(q_{\text{PREX}} = 0.475/\text{fm})$ in ^{208}Pb (related to PREX)
- binding energy E_B for extremely exotic nuclei (super heavy, very neutron rich)
- Q_α value (energy gain in α emission) for super-heavy element
- S_{2n} two-neutron separation energy for Sn chain
- B_f = fission barrier in ^{266}Hs
- dipole polarizability $\alpha_D = \int_0^\infty d\omega S_D(\omega) \omega^{-1}$ in ^{208}Pb
- dipole strength distribution $S_D(\omega) = B(E1, \omega)$ (note: $\sigma_{\text{photoabs}}(\omega) = \omega S_D(\omega)$)

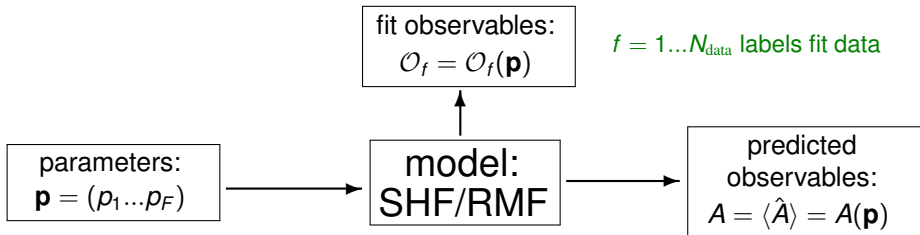
Optimization of model parameters and subsequent variation

model:
SHF/RMF

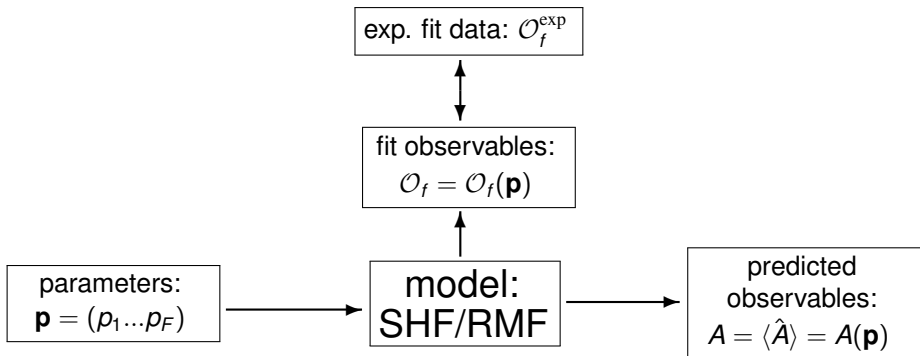
Optimization of model parameters and exploring its variation



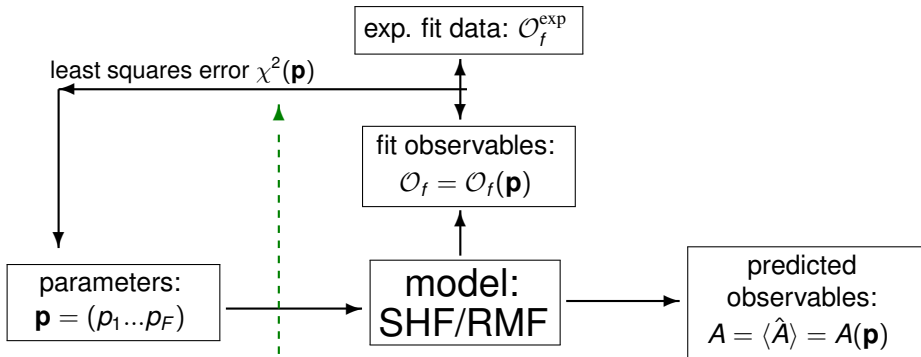
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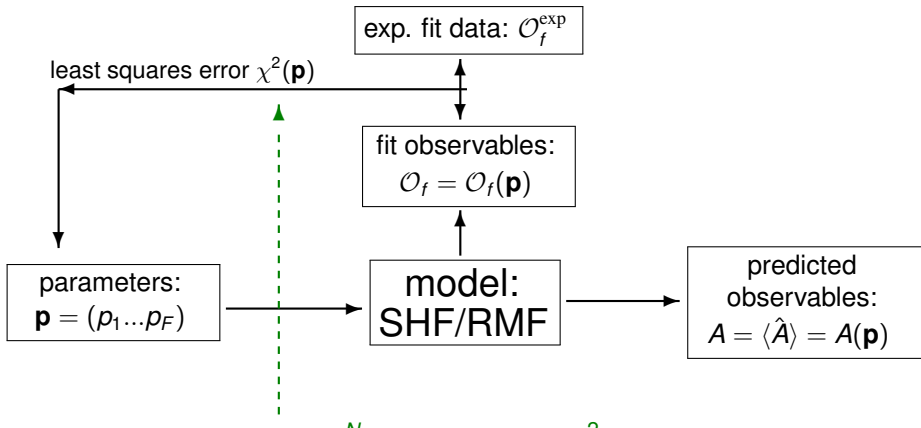


Optimization of model parameters and exploring its variation



$$\chi^2(\mathbf{p}) = \sum_{f=1}^{N_{\text{data}}} \frac{(O_f(\mathbf{p}) - O_f^{\text{exp}})^2}{\Delta O_f^2}$$

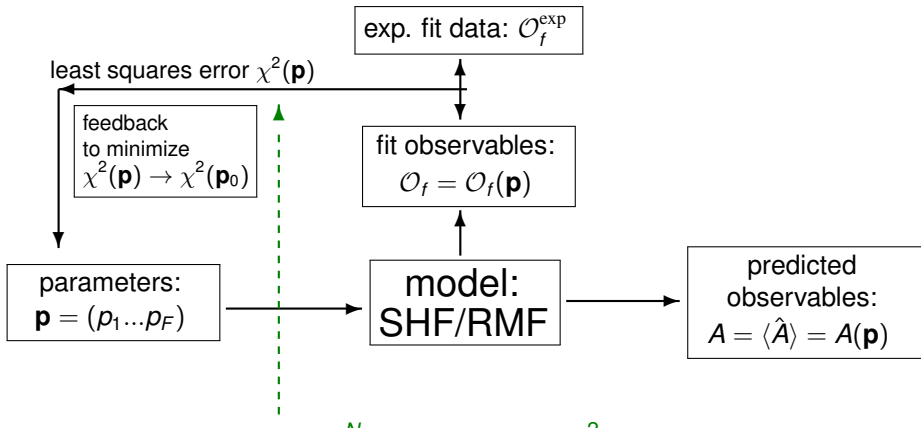
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adopted error $\Delta \mathcal{O}_f$: $\chi^2(\mathbf{p}_0) = N_{\text{data}} - N_{\text{params}}$

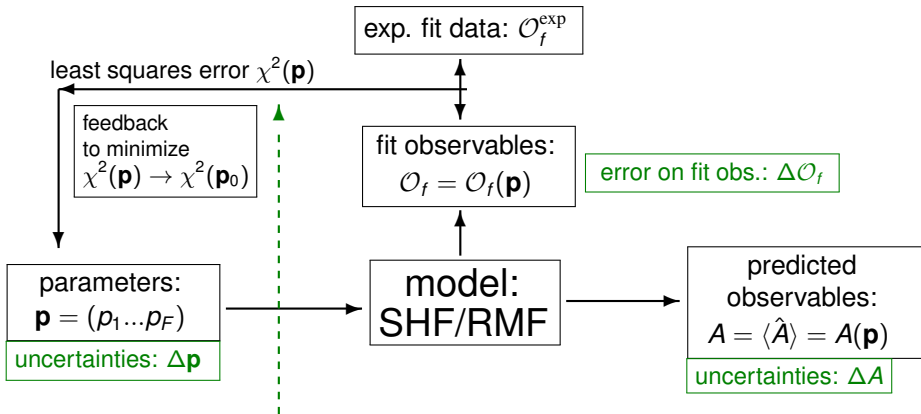
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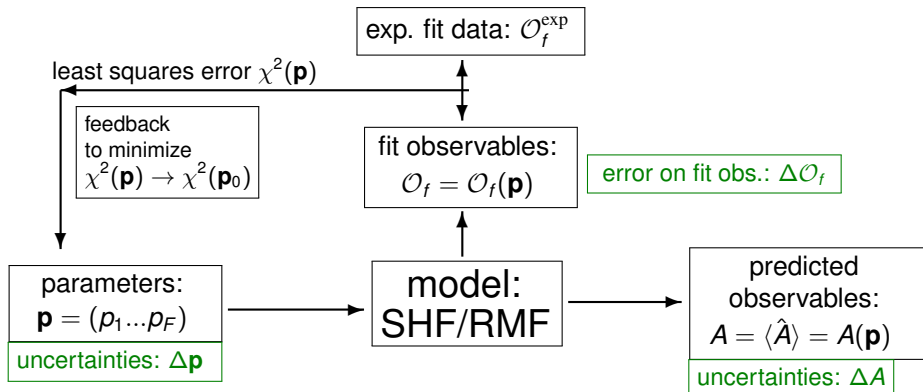
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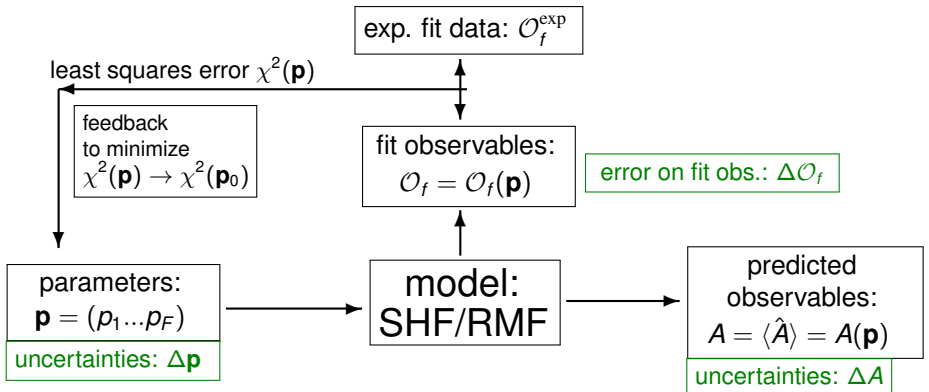
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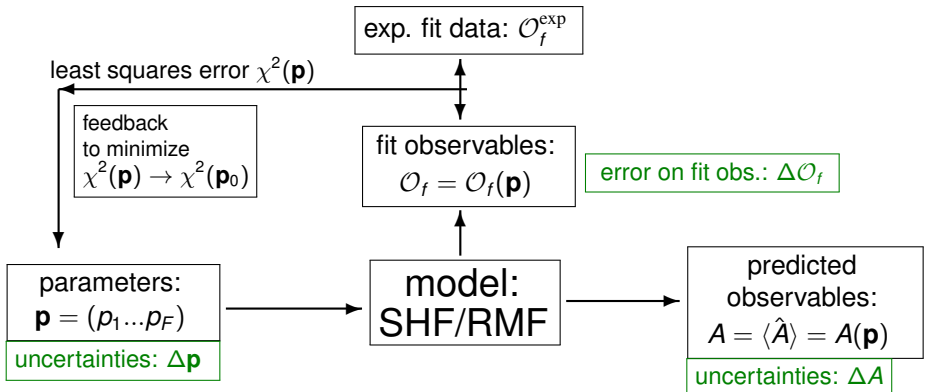


Optimization of model parameters and exploring its variation



← reasonable range of \mathbf{p} : $\chi^2(\mathbf{p}) = \chi^2(\mathbf{p}_0) + 1$

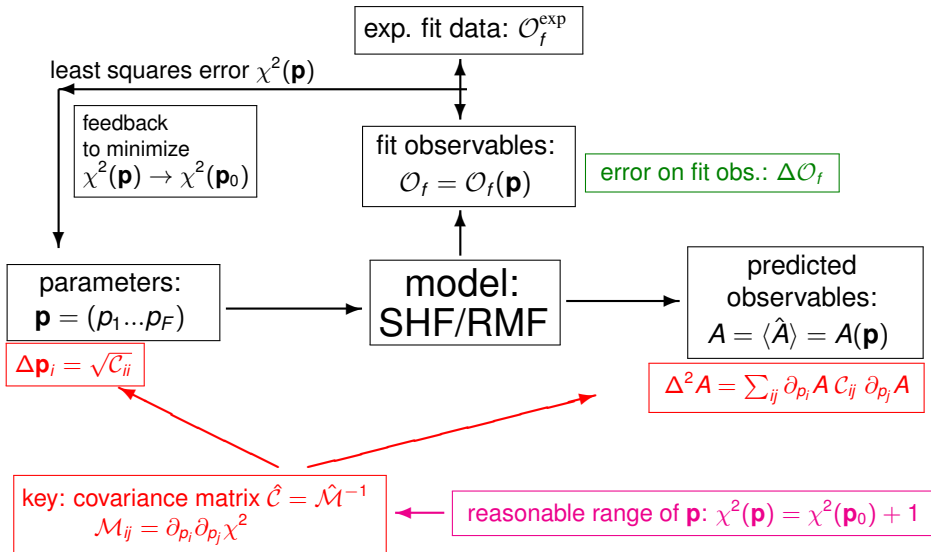
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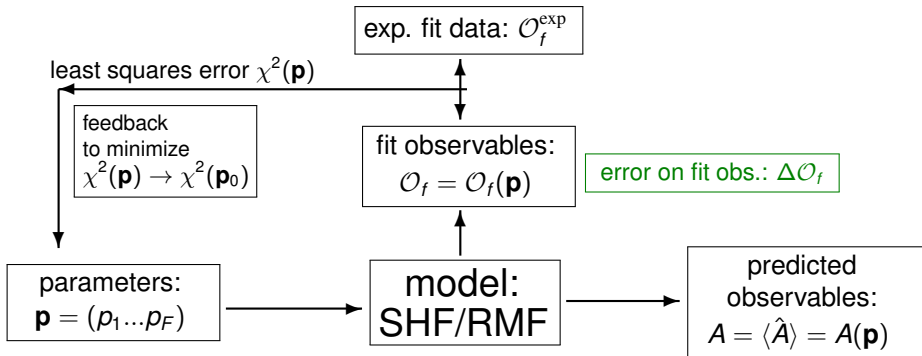
key: covariance matrix $\hat{C} = \hat{M}^{-1}$
 $M_{ij} = \partial_{p_i} \partial_{p_j} \chi^2$

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mixed variance:

$$\Delta^2(AB) = \sum_{ij} \partial_{p_i} A C_{ij} \partial_{p_j} B$$

covariance: $C_{AB} = \frac{\Delta^2(AB)}{\sqrt{\Delta^2(AA)\Delta^2(BB)}}$

$C_{AB} = 1 \leftrightarrow$ highly correlated

$C_{AB} = 0 \leftrightarrow$ uncorrelated

Error propagation – “reasonable parametrizations” and uncertainties

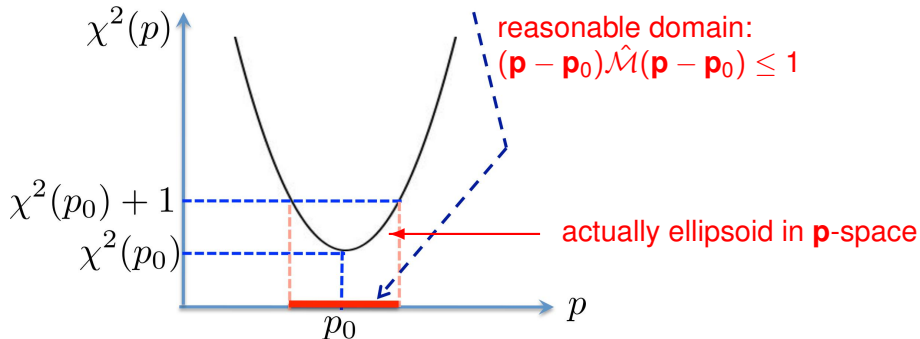
probability distribution of “reasonable” parameters:

$$W(\mathbf{p}) \propto e^{-\chi^2(\mathbf{p})} \approx \exp\left(-\frac{1}{2}(\mathbf{p} - \mathbf{p}_0)\hat{\mathcal{M}}(\mathbf{p} - \mathbf{p}_0)\right), \quad \mathcal{M}_{ij} = \partial_{p_i}\partial_{p_j}\chi^2$$

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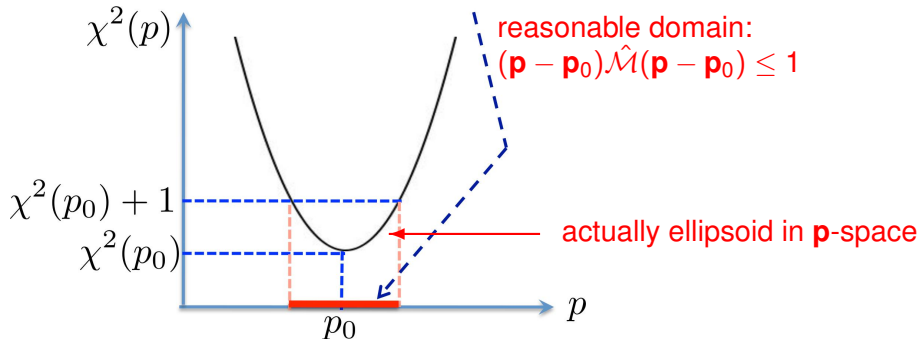
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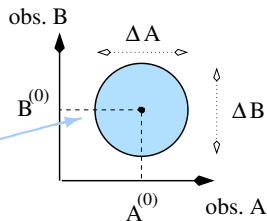
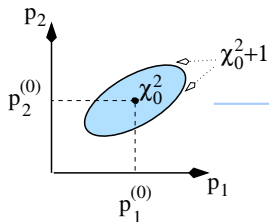
mean values and variances as probability averages (statistical):

$$\bar{A} = \int d\mathbf{p} W(\mathbf{p}) A(\mathbf{p}) \approx \int d\mathbf{p} W(\mathbf{p}) (A(\mathbf{p}_0) + (\mathbf{p} - \mathbf{p}_0) \cdot \nabla A) = A(\mathbf{p}_0)$$

$$\overline{AB} = \int d\mathbf{p} W(\mathbf{p}) A(\mathbf{p})B(\mathbf{p}) \approx \nabla A \cdot \hat{\mathcal{M}}^{-1} \cdot \nabla B$$

Error propagation – covariance analysis

ellipsoid of “reasonable” parameters:



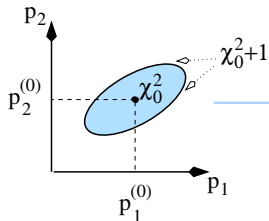
given by: $\Delta \mathbf{p} \cdot (\nabla \otimes \nabla \chi^2) \cdot \Delta \mathbf{p} = 1$

observables: $A = A(\mathbf{p})$, $B = B(\mathbf{p})$

$$\Rightarrow \boxed{\overline{\Delta A \Delta B} = \nabla A \cdot (\nabla \otimes \nabla \chi^2)^{-1} \cdot \nabla B}$$

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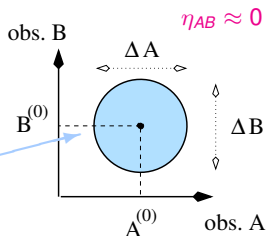
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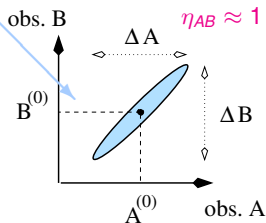
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correlation: $\eta_{AB} = \frac{\overline{\Delta A \Delta B}}{\sqrt{\overline{\Delta^2 A} \overline{\Delta^2 B}}}$

uncorrelated observables:

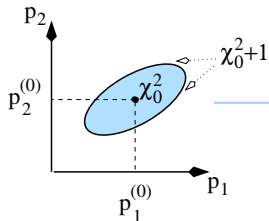


highly correlated observables:



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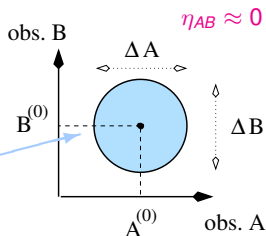
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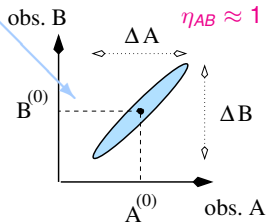
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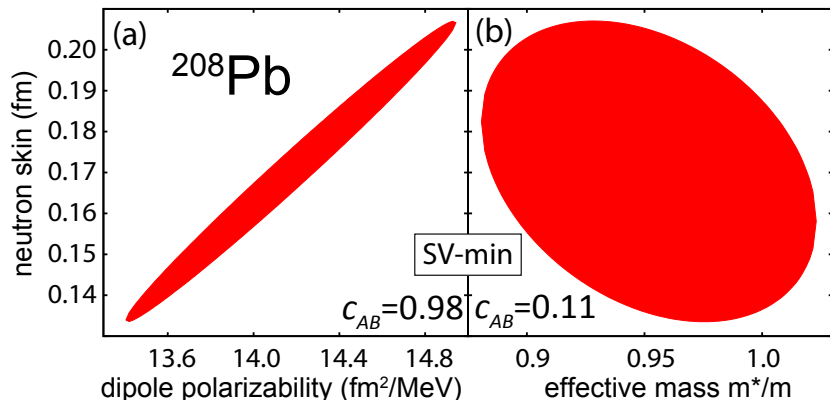


highly correlated observables:



correlations \longleftrightarrow underlying nuclear model \Rightarrow “reasonable model” required

Covariance analysis: example r_n , α_D , and m^*/m in ^{208}Pb



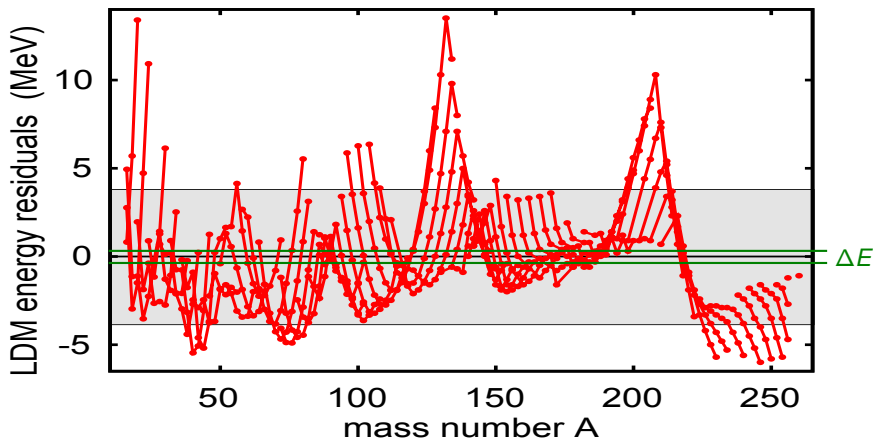
highly correlated: polarizability $\alpha_D \leftrightarrow$ neutron skin $r_n - r_p$

almost uncorrelated: isoscalar effective mass $m^*/m \leftrightarrow$ neutron skin $r_n - r_p$

Strategies to estimate errors (systematic and statistical)

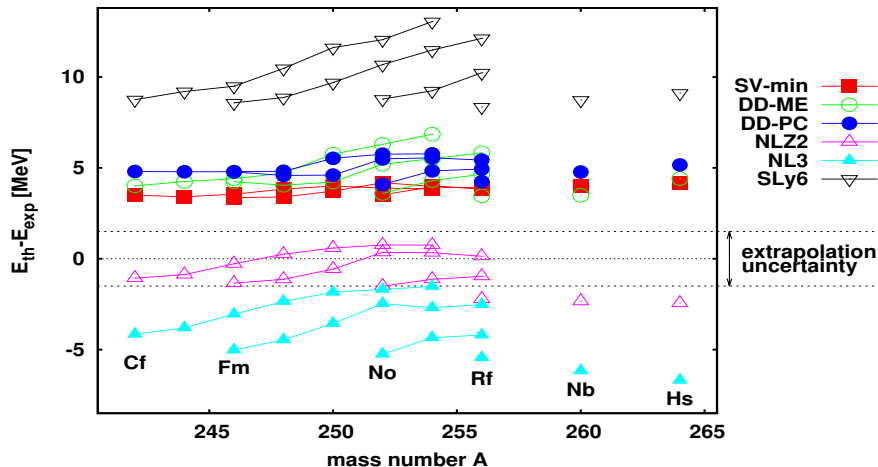
- 1 Dedicated variations of parameters (NMP): (statistical)**
perform χ^2 with fixed \mathcal{O}_{NMP} , watch the trend of $A(\mathbf{p})$ when varying \mathcal{O}_{NMP}
- 2 Extrapolation uncertainties from statistical analysis (statistical):**
 $\chi^2(\mathbf{p})$ & observable $A(\mathbf{p}) \implies$ uncertainty ΔA
- 3 Correlations between observables from statistical analysis (statistical):**
 $\chi^2(\mathbf{p})$ & two observables $A(\mathbf{p}), B(\mathbf{p}) \implies$ covariance c_{AB}
- 4 Sensitivity analysis for the model parameters \mathbf{p} (statistical):**
Take a model parameter p_α as “observable” ($p_\alpha \equiv B$), look at covariance c_{Ap_α}
- 5 Variations of fit data: (statistical/systematic)**
switch on/off groups of fit observables $\hat{\mathcal{O}}_i$, energy, radius, ...
- 6 Trends of residual errors: (systematic)**
is distribution of residuals $\mathcal{O}_i(\mathbf{p}) - \mathcal{O}_i^{\text{exp}}$ stochastic? \leftrightarrow trends signal insufficiencies
- 7 Comparison with predicted data: (systematic)**
actual deviation should fulfill $A(\mathbf{p}_0) - A^{\text{(exp)}} \leq \Delta A$, larger deviations indicate problems
- 8 Variations of the model: (systematic)**
compare different models (RMF \leftrightarrow SHF), extend a model by new terms

Trends of residual errors: example LDM and energy systematics



LDM misses quantum shell effects \leftrightarrow can describe only average trends of $E_B(N, Z)$
residual errors on E_B not stochastically distributed \leftrightarrow trends reflect shell effects
actual average error (grey area) \gg extrapolation error from χ^2 analysis (ΔE)

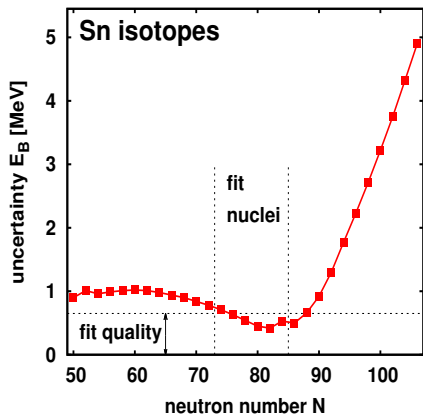
Comparison with predicted data: example known super-heavy elements



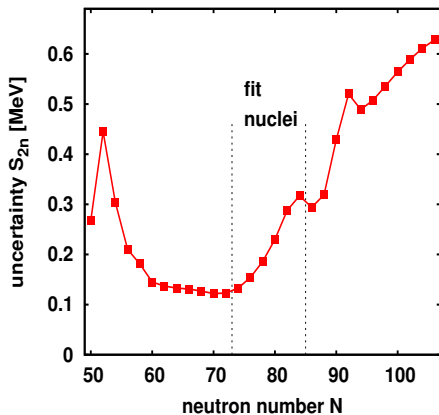
residual errors on E_B stay $\leq \pm 1$ MeV for $A < 220$ for all parametrizations
 deviations of E_B grow systematically with $A \Leftrightarrow$ something missing in all models
 deviations have different sign for RMF vs. SHF \Leftrightarrow something different is missing
 good performance of (RMF) NL-Z irrelevant \leftrightarrow NL-ME has poor quality else wise

Extrapolation uncertainties: example chain of Sn isotopes

binding energies



2-neutron separation energies



uncertainty E_B grows quickly with $N \nearrow$, far beyond ΔE_B in fit

\leftrightarrow extrapolations deteriorate with distance to fit data

but uncertainty S_{2n} generally smaller & not dramatically growing (factor < 2)

\rightarrow differences E_B often more reliable than E_B as such

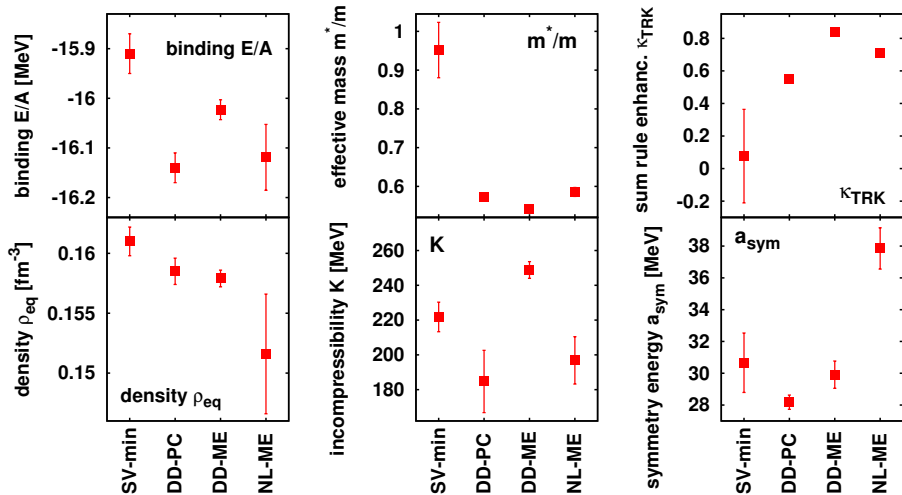
\implies : uncertainties ΔA from statistical analysis very useful indicators

Comparison SHF \leftrightarrow $3 \times$ RMF

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Goal: large differences in predictions indicate intrinsic differences of models
and give an idea about possible systematic errors

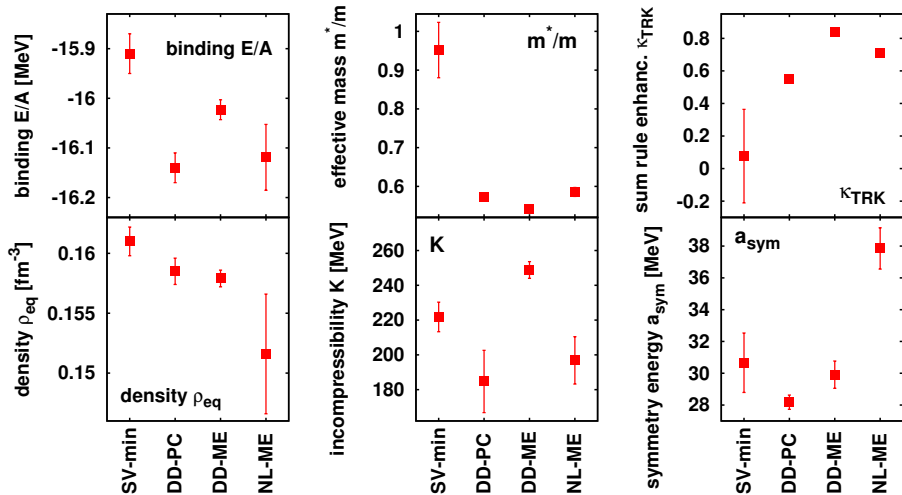
SHF \leftrightarrow 3 \times RMF: predictions for NMP



uncertainties $\Delta\mathcal{O}$: very different \leftrightarrow specific structure of models

m^*/m , κ_{TRK} fixed in RMF; a_{sym} larger in SHF; NL-ME generally weak

SHF \leftrightarrow 3 \times RMF: predictions for NMP

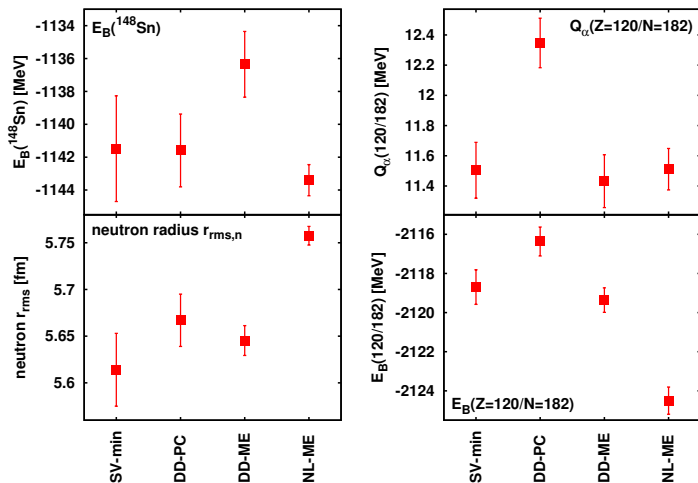


uncertainties ΔO : very different \leftrightarrow specific structure of models

m^*/m , κ_{TRK} fixed in RMF; a_{sym} larger in SHF; NL-ME generally weak

deviations between models: often larger than ΔO \leftrightarrow systematic differences of models

SHF \leftrightarrow 3 \times RMF: predictions for observables in (super-)heavy nuclei



T=1 obs. (left panels): SHF larger $\Delta\mathcal{O} \leftrightarrow$ isovector freedom
 NL-ME: smallest $\Delta\mathcal{O} \leftrightarrow$ predictive power? – too rigid model!
 deviations $> \Delta\mathcal{O}$: not as dramatic as for NMP

Dedicated variation: example trends with symmetry energy a_{sym} (J)

strategy: produce parametrizations with constraint on a_{sym}

⇒ a set of parametrizations with systematically varied a_{sym}
and else wise similar properties

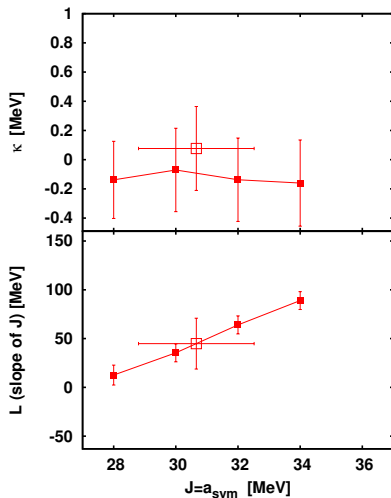
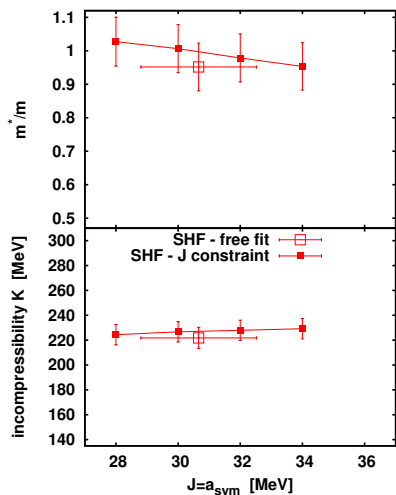
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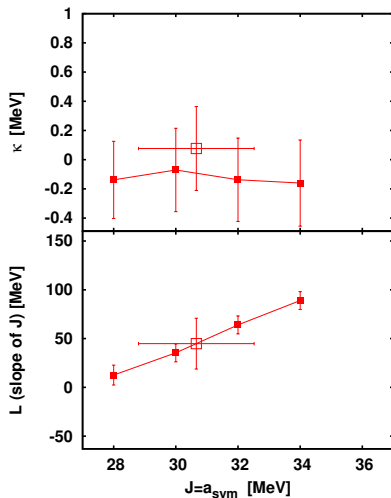
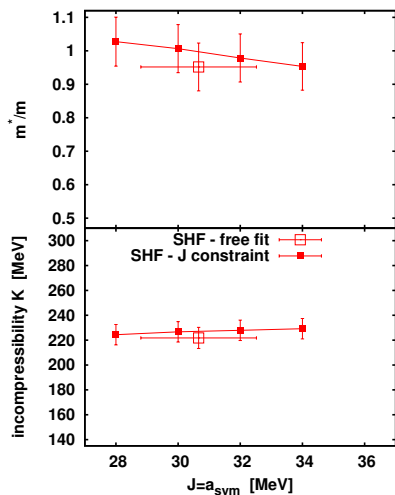
combined with comparison of models – SHF & 3×RMF

Trends of nuclear matter properties (NMP) with a_{sym} for the case of SHF



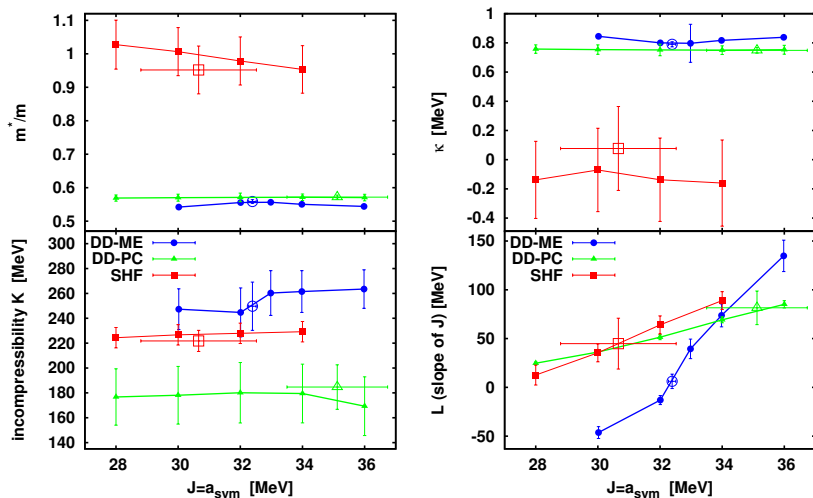
ΔJ of free fit = reasonable range $J \leftrightarrow \Delta \mathcal{O}$ of free fit complies with slope of constr. fits

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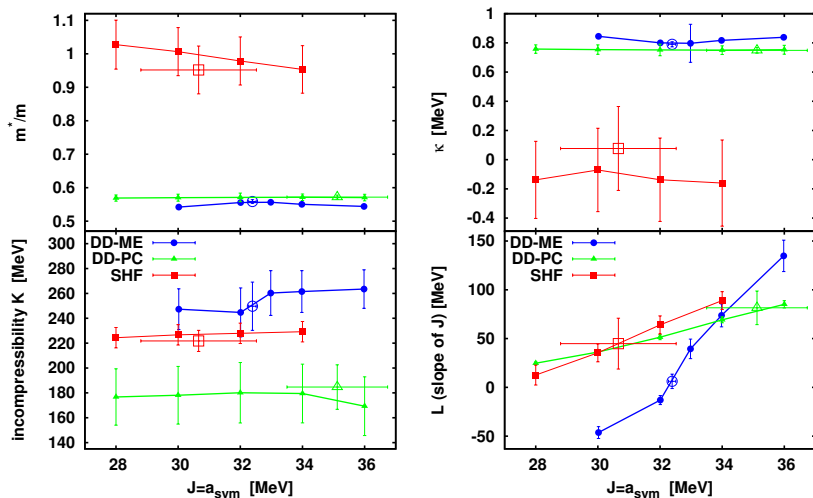
ΔJ of free fit = reasonable range $J \leftrightarrow \Delta \mathcal{O}$ of free fit complies with slope of constr. fits
 $K, m^*/m, \kappa$ independent of $J (= a_{\text{sym}}) \implies$ four independent model parameters
 $L = \text{slope of } J \text{ strongly linked with } J \longleftrightarrow$ hidden correlation in data (and model)

Trends of nuclear matter properties (NMP) with a_{sym} – SHF & RMF



for all cases: K , m^*/m , κ independent of a_{sym} , and L strongly linked with J
 but $L(J)$ differs between models \leftrightarrow density dependent $T=1$ terms?

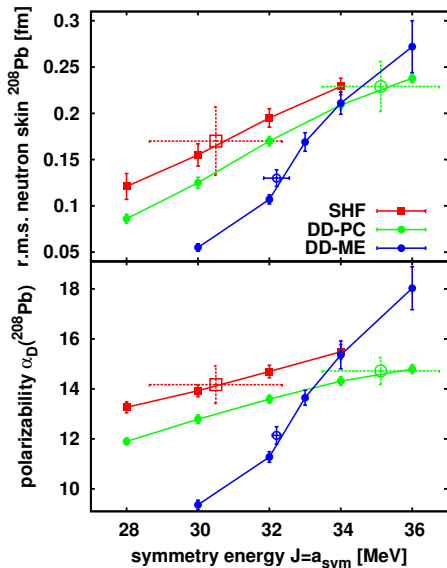
Trends of nuclear matter properties (NMP) with a_{sym} – SHF & RMF



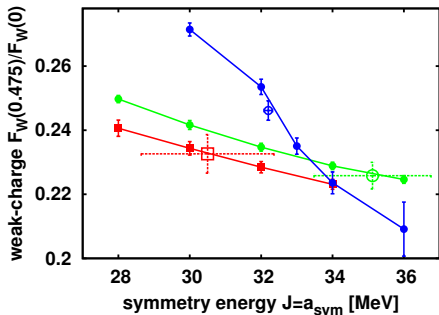
for all cases: K , m^*/m , κ independent of a_{sym} , and L strongly linked with J
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large differences SHF \leftrightarrow RMF: mass parameters m^*/m and κ (already seen before)

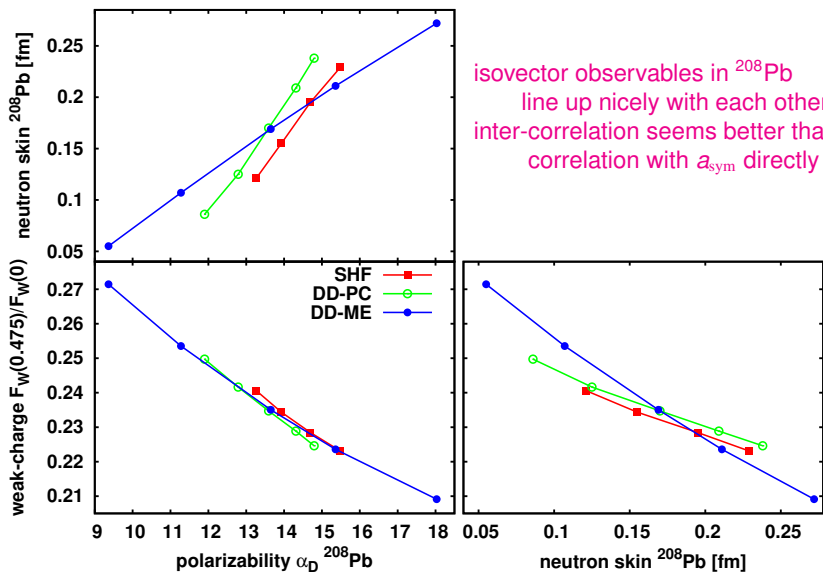
Trends of isovector observables in ^{208}Pb with a_{sym}



for each model: clear trends with a_{sym}
 SHF \leftrightarrow DD-PC: same slope,
 different offset
 DD-ME: even different slope
 \leftrightarrow different ρ -dependence
 (but small variance of a_{sym})

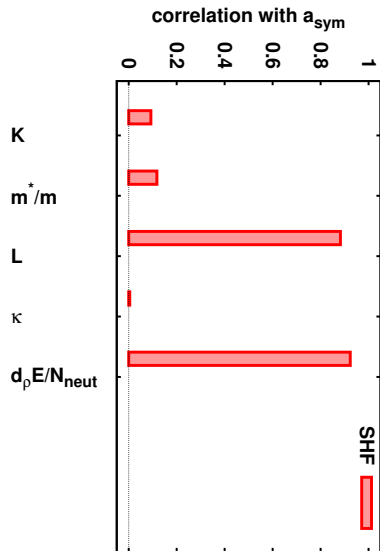


Trends of isovector observables in ^{208}Pb – drawn against each other



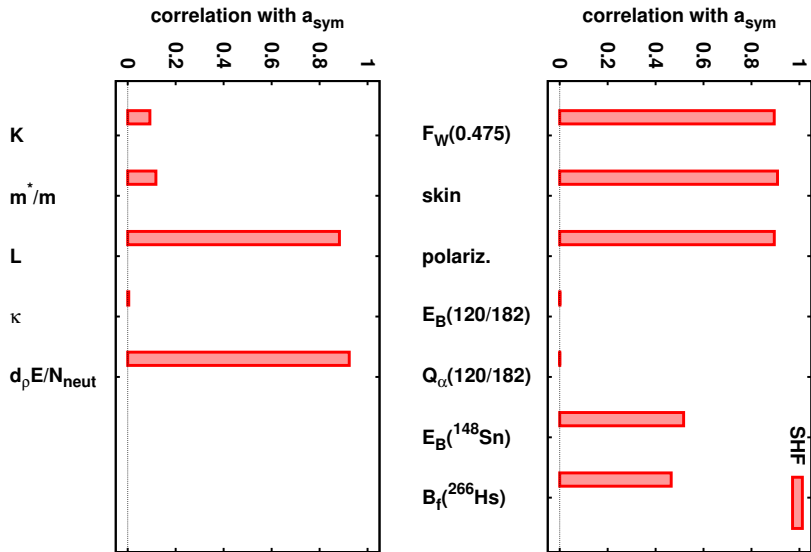
Examples from covariance analysis

Correlations of a_{sym} with NMP, for the case of SHF



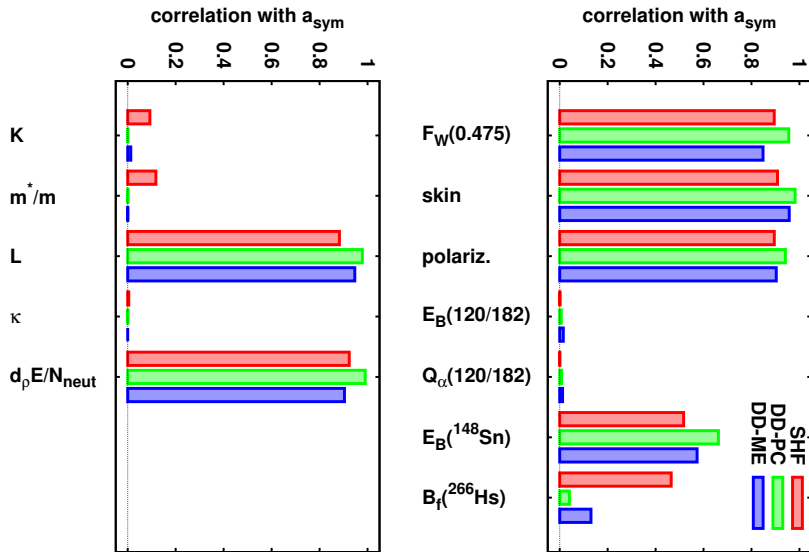
Again: K , m^*/m , κ perfectly independent; but strong correlation with L , $\partial_\rho E_{\text{neut}}/N$
surface properties a_{surf} , $a_{\text{surf,sym}}$ show a mixed picture

Correlations of a_{sym} with NMP and observables in finite nuclei – SHF



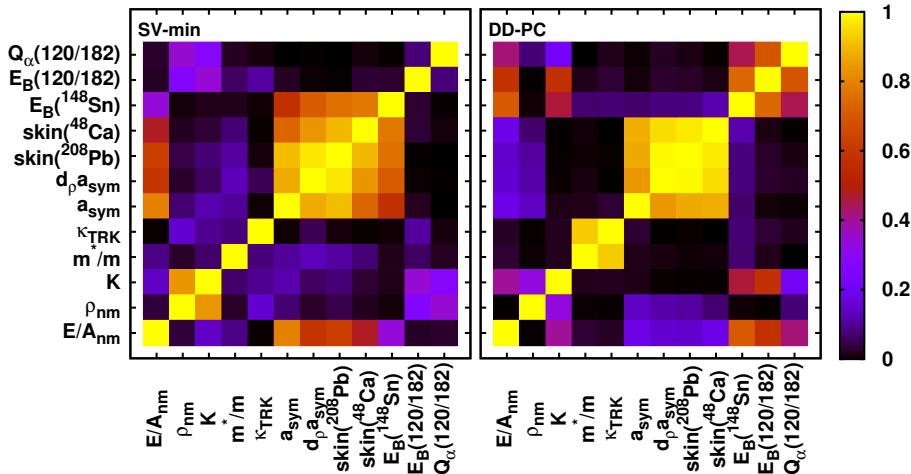
Good isovector observables (in ^{208}Pb): weak-charge formfactor F_W , neutron skin $r_n - r_p$, polarizability α , superheavy elements uncorrelated, neutron rich nuclei & fission somewhat correlated

Correlations of a_{sym} with NMP and obs. in finite nuclei – SHF & RMF



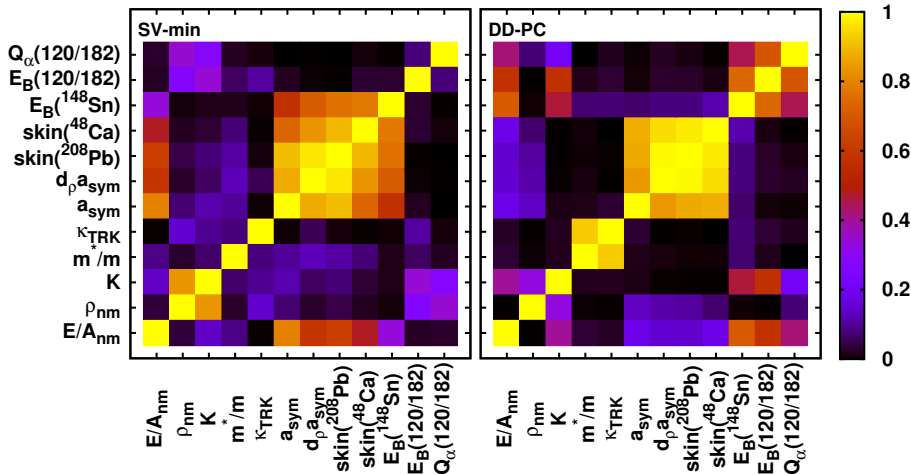
RMF models shows nearly the same correlations as SHF except for surface properties and fission

Correlations of a_{sym} with NMP and obs. in finite nuclei – SHF & RMF



DD-PC \leftrightarrow SHF show at first glance similar correlations

Correlations of a_{sym} with NMP and obs. in finite nuclei – SHF & RMF



DD-PC \leftrightarrow SHF show at first glance similar correlations

differences: $\kappa_{\text{TRK}}^* \leftrightarrow m^*/m$ – due to “frozen mass” in RMF

$K \leftrightarrow \rho_{\text{nm}}$ – due to different modeling of density dependence

exotic nuclei – density dep. & isovector terms & shell structure (m^*/m)

Least-squares fits and statistical analysis:

powerful tool: extrapolation errors, covariances, trends \leftrightarrow insight into model

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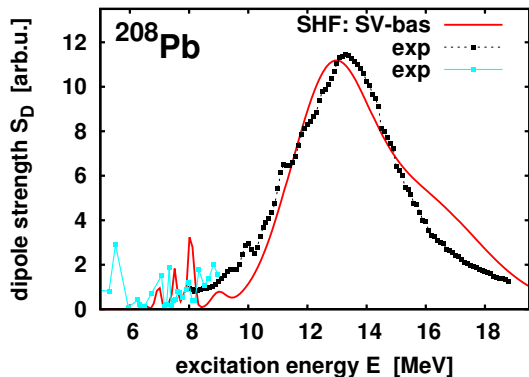
Systematic errors:

several different tools: comparison models, trend residuals, data, ...

no ultimate solution: there is no systematic way to evaluate systematic errors

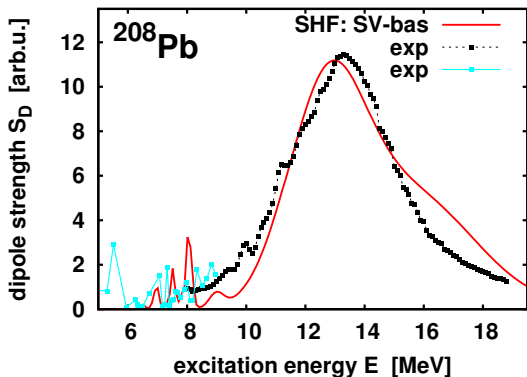
Correlations with dipole strength

Dipole strength and related observables



integrated strength: $\mathcal{I}_D^{(n)}(E) = \int_0^E dE' S_D(E') E'^n$

Dipole strength and related observables



polarizability:

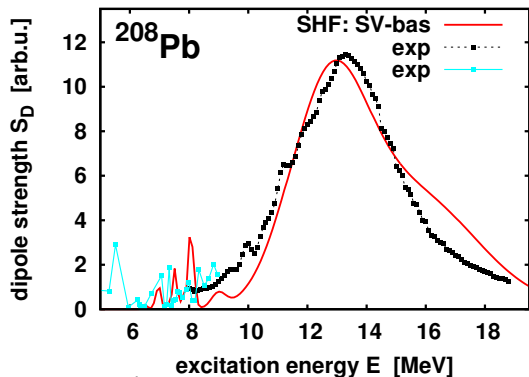
$$\begin{aligned}\alpha_D &= \mathcal{I}_D^{(-1)}(\infty) \\ &= \int_0^\infty dE S_D(E) E^{-1}\end{aligned}$$

TRK sum-rule enhancement:

$$\begin{aligned}\kappa_{\text{TRK}} &= \mathcal{I}_D^{(1)}(\infty) \\ &= \int_0^\infty dE S_D(E) E^{+1}\end{aligned}$$

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Dipole strength and related observables



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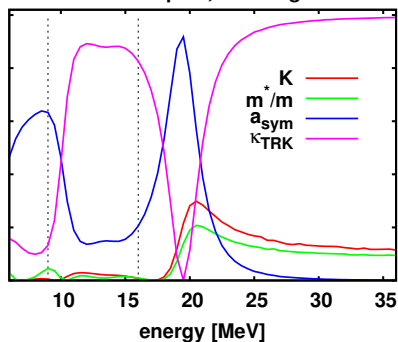
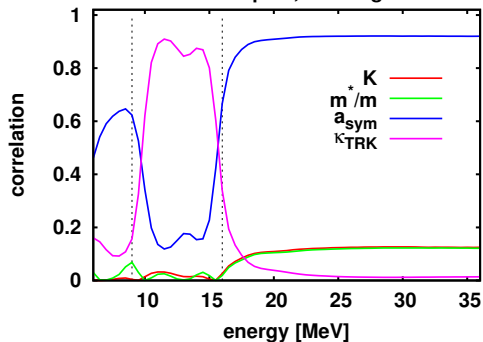
$$= \int_0^\infty dE S_D(E) E^{+1}$$

integrated strength: $\mathcal{I}_D^{(n)}(E) = \int_0^E dE' S_D(E') E'^n$

low-lying dip. strength ("pygmy"): $\mathcal{I}_D^{(n)}(9 \text{ MeV}) \leftrightarrow$ signal for $\alpha_D, \kappa_{\text{TRK}}$?

Correlation of $\int_0^E dE' \sigma(E') E'^n$ with nuclear matter parameters

correlation: integrated dip.strength with NMP, ^{208}Pb , SV-min
 with T=1 dipole, E^{-1} weighted with T=1 dipole, E^{+1} weighted

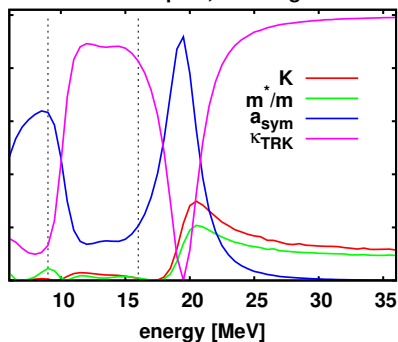
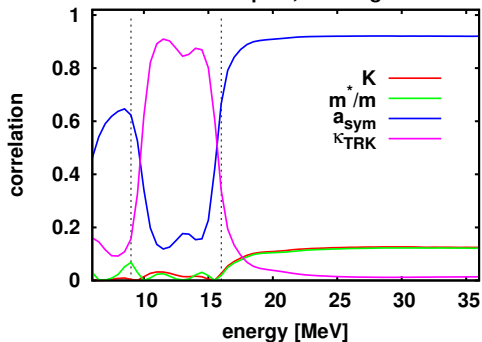


E^{-1} weight favors a_{sym} for $E \rightarrow \infty$

E^{+1} weight favors κ for $E \rightarrow \infty$

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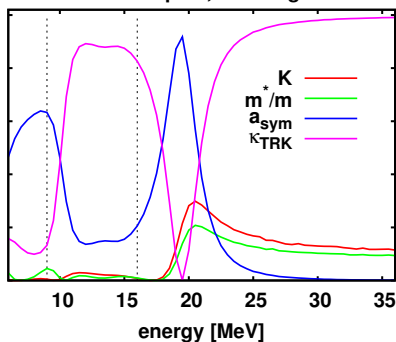
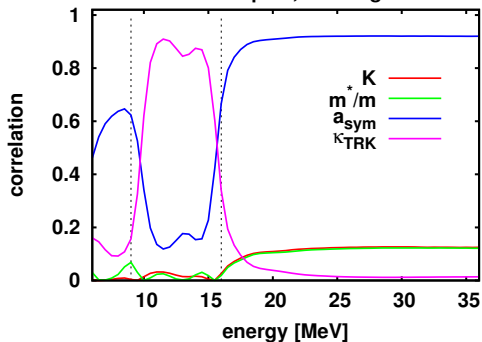
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GDR region favors κ for all E weights

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E^{-1} weight favors a_{sym} for $E \rightarrow \infty$

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GDR region favors κ for all E weights

pygmy region yields mixed info about a_{sym} and $\kappa \implies$ useful data in combined analysis
 (open problem: robust choice of cutoff energy E_{cut})

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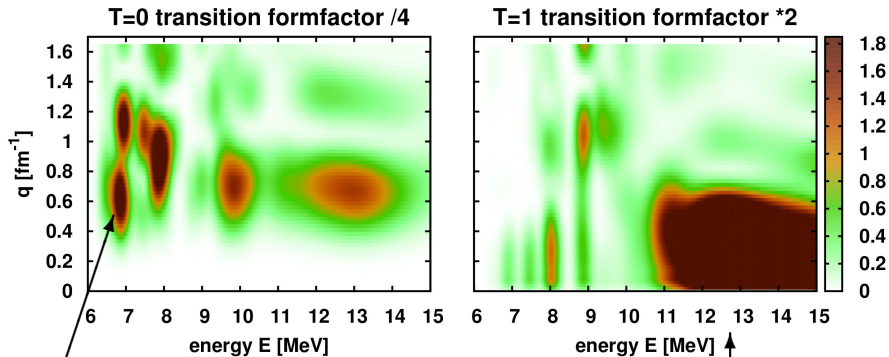
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Transition formfactor $f_{\text{trans}}(q, E)$ for $L = 1$ modes in ^{208}Pb

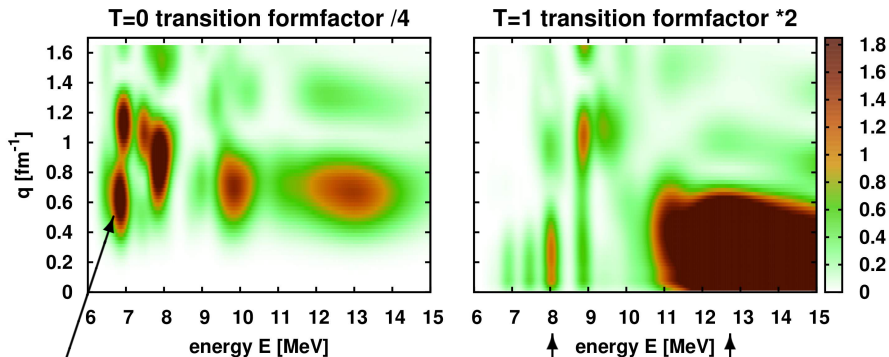
^{208}Pb , SV-bas, L=1 modes



strong $T = 0$ modes, $q \approx 0.6$
($q = 0$ occupied by c.m. mode)

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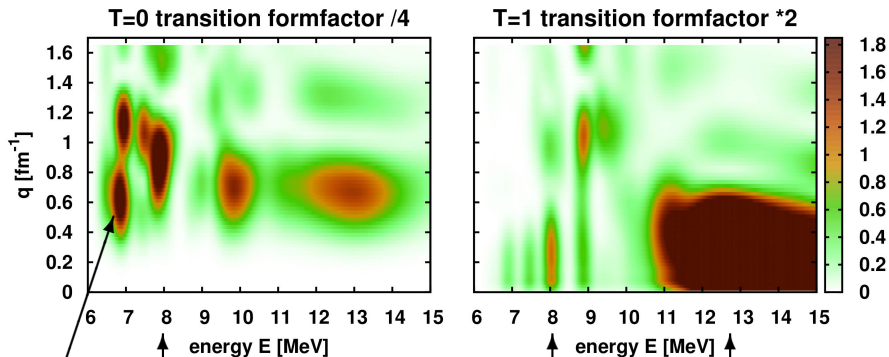
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^{208}Pb , SV-bas, L=1 modes



related to strong $T = 0$ mode ← low E dip.str.

GDR

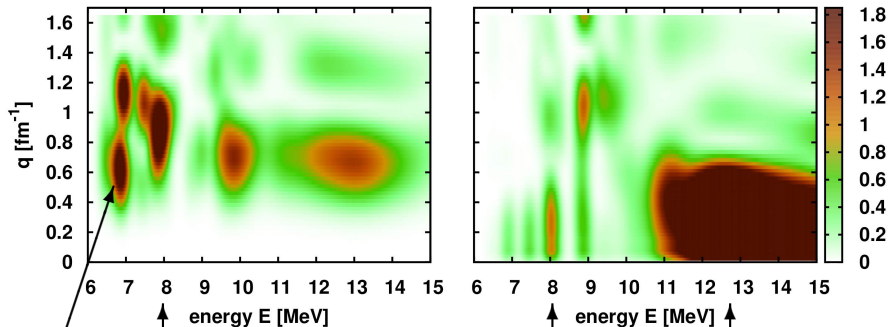
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^{208}Pb , SV-bas, $L=1$ modes

T=0 transition formfactor /4

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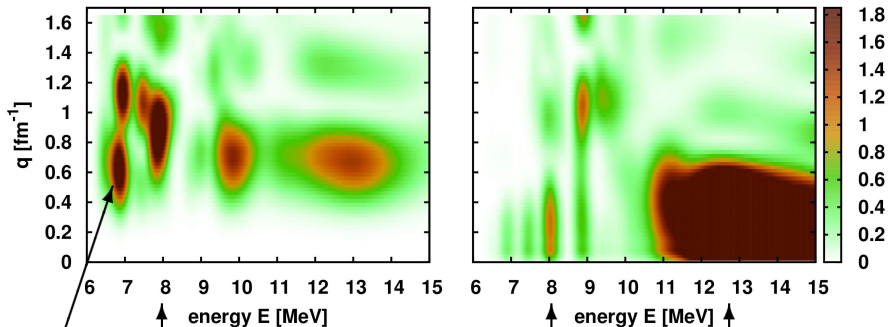
pure $L=1$, $q \approx 0.6$

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 ($q = 0$ occupied by c.m. mode)

pure $L=1$, $q \approx 0.6$

low E dipoles: mixed modes, predominantly isoscalar, nonetheless useful info on NMP