

Impact of correlations on the phase diagram of nuclear matter

Drissi, Rios & Barbieri, Ann. Phys. **469** 169729 (2024)
Drissi, Rios & Barbieri, Ann. Phys. **469** 169730 (2024)
Drissi & Rios, EPJA **58** 90 (2022)

M Drissi



TECHNISCHE
UNIVERSITÄT
DARMSTADT

C Barbieri



UNIVERSITÀ
DEGLI STUDI
DI MILANO

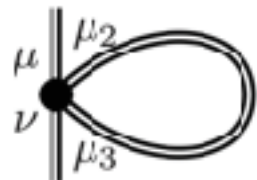
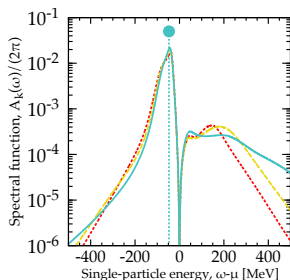
Dr Arnau Rios Huguet

Institute of Cosmos Sciences
Universitat de Barcelona

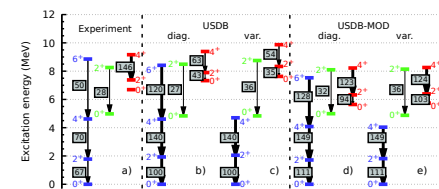
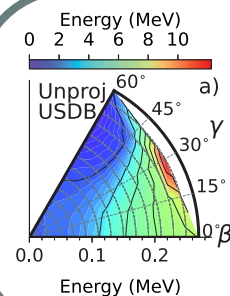
ESNT Pairing Workshop
21 May 2025

Nuclear Physics in Barcelona

Dense matter

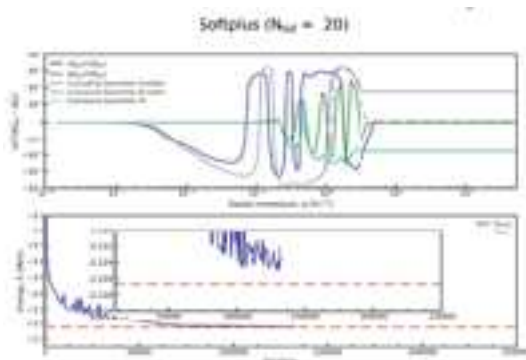


Nuclear structure



Frycz et al PRC 110, 054326 (2024)

Machine learning

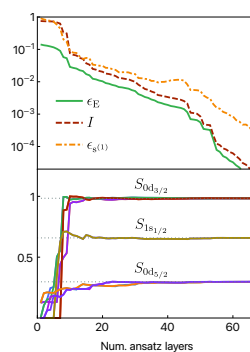
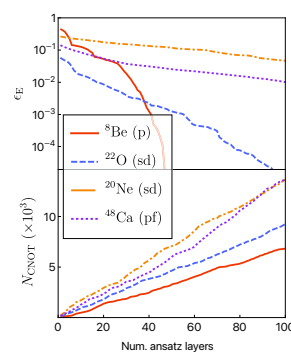


JWT Keeble



J Rozalén

Quantum computing



AM Romero



J Menéndez

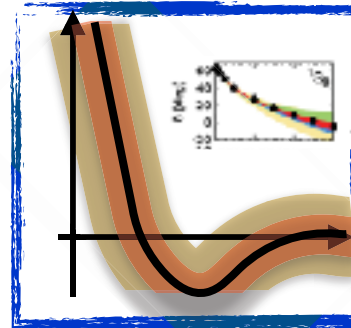


Fellowships

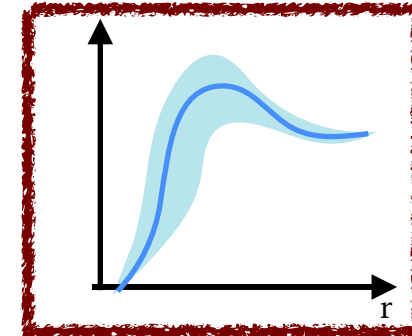


• Motivation

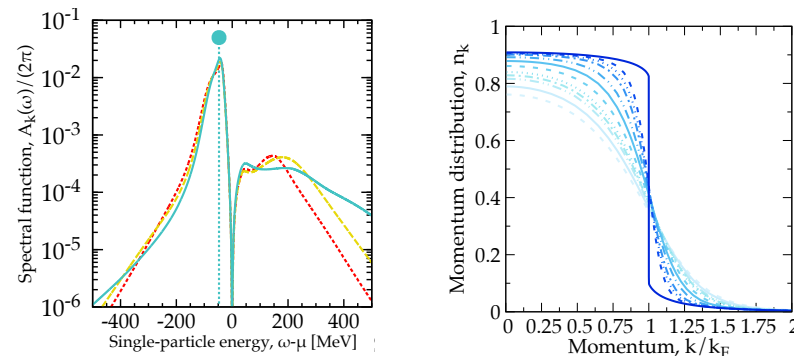
Hamiltonian



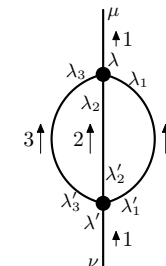
Many-body method



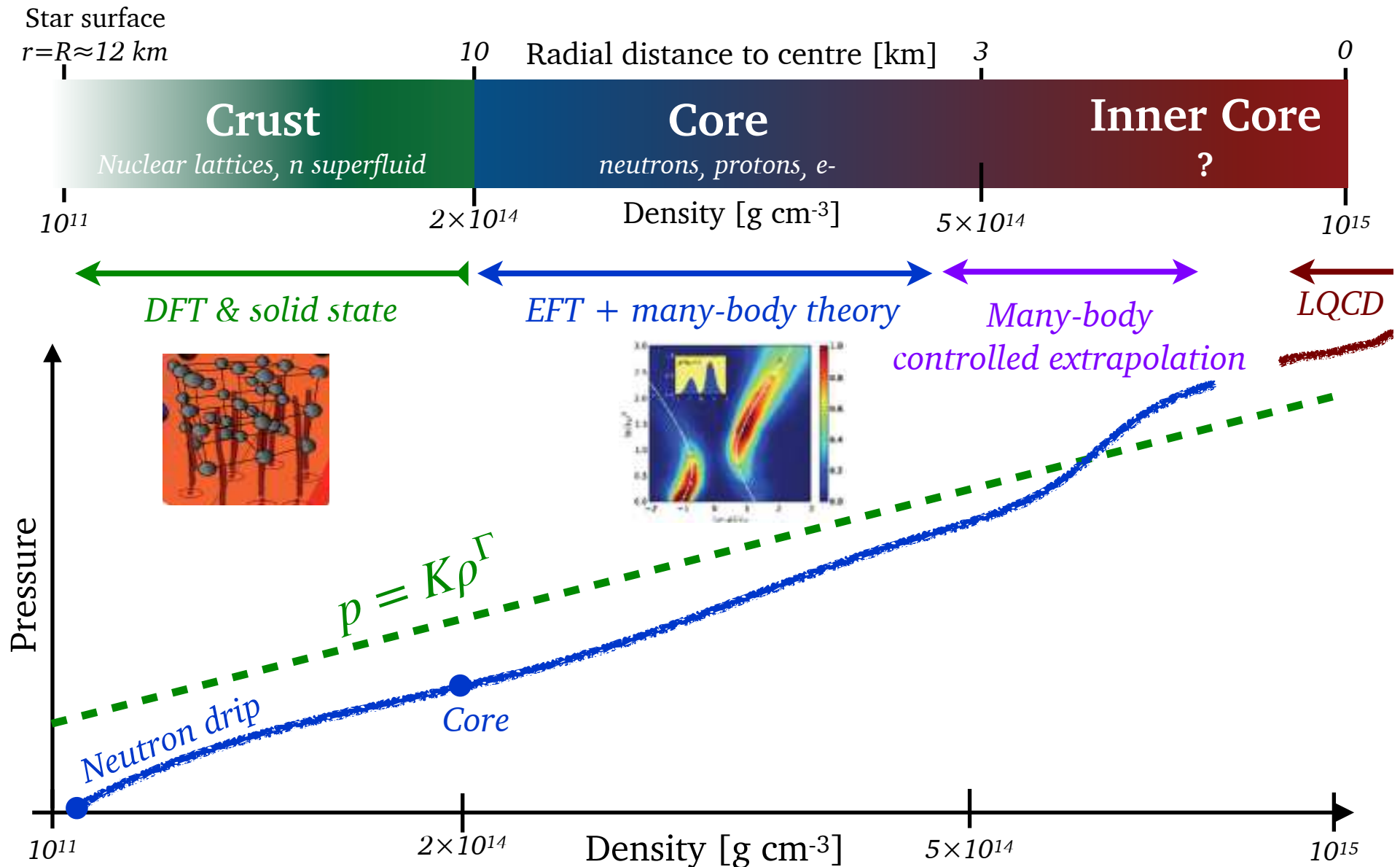
• “Normal” self-consistent Green’s functions



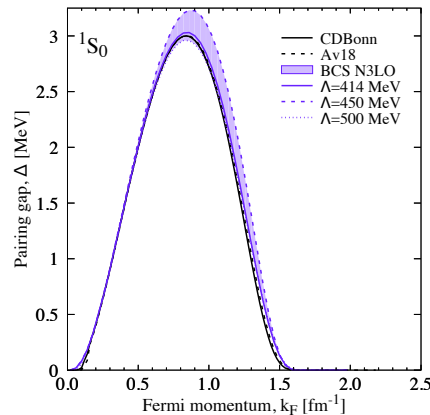
• Nambu-Covariant Green’s functions



Neutron-star modelling

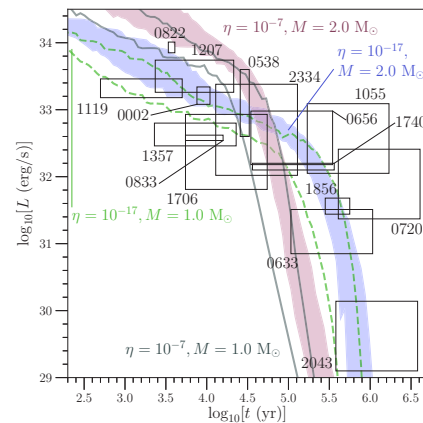


Why are pairing gaps necessary?



• Glitches in period

- Pinning of neutron superfluid, crustal physics
- Proton superconductor?

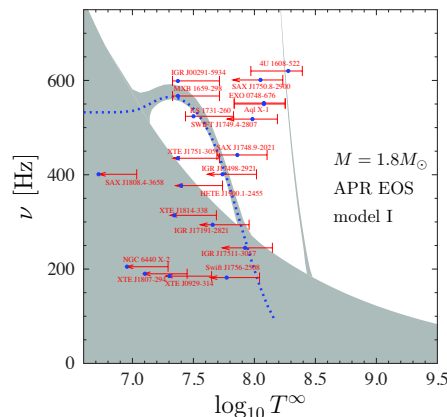


• Neutron star cooling

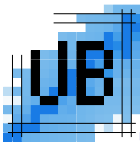
- Neutrino rates through pair-breaking

• Astereoseismology

- Superfluid modes affected
- Potential observable signal in GW in BNS

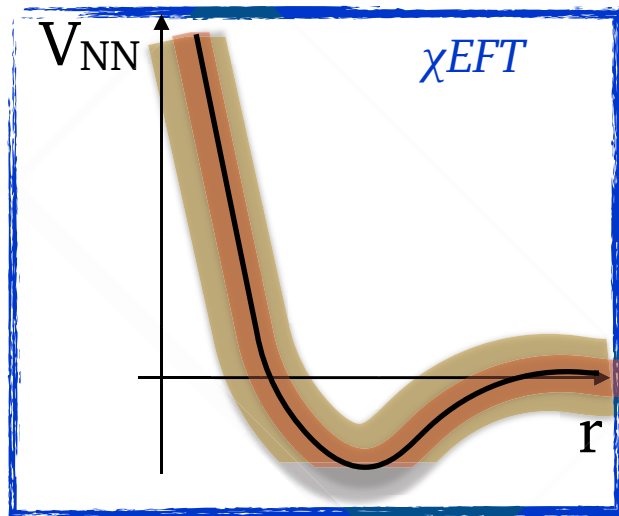


Kantor et al. Phys. Rev. Lett. **125** 151101 (2020),
Rau & Wasserman, MNRAS **481**, 4427-4444 (2018)

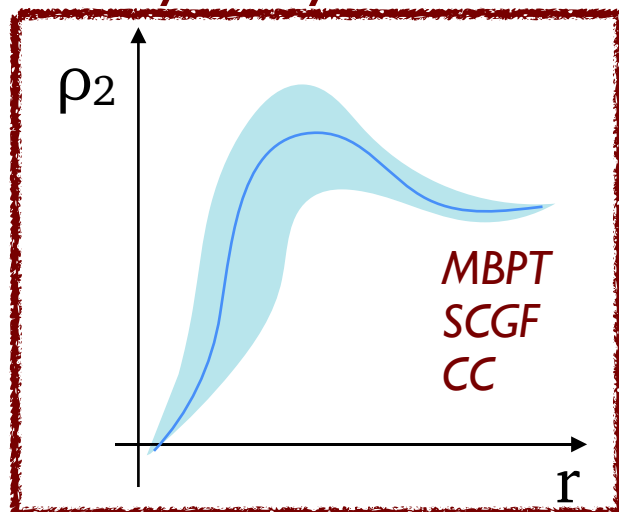


Nuclear uncertainty quantification

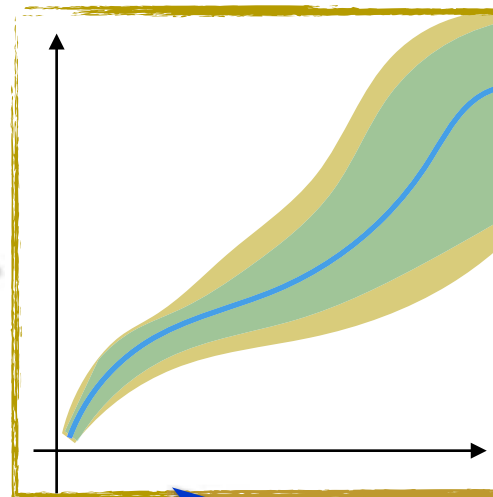
Hamiltonian



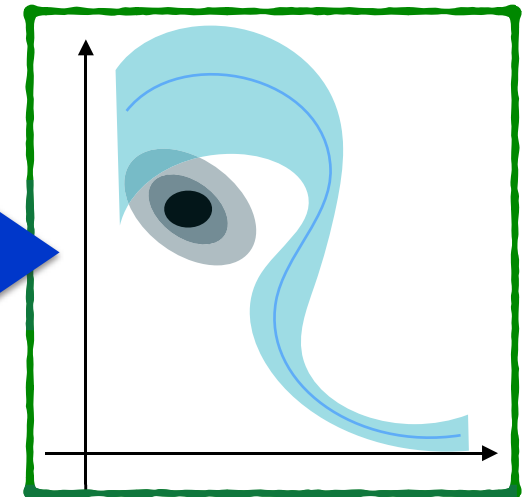
Many-body method



Astronuclear
property

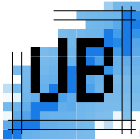


Neutron star
observations



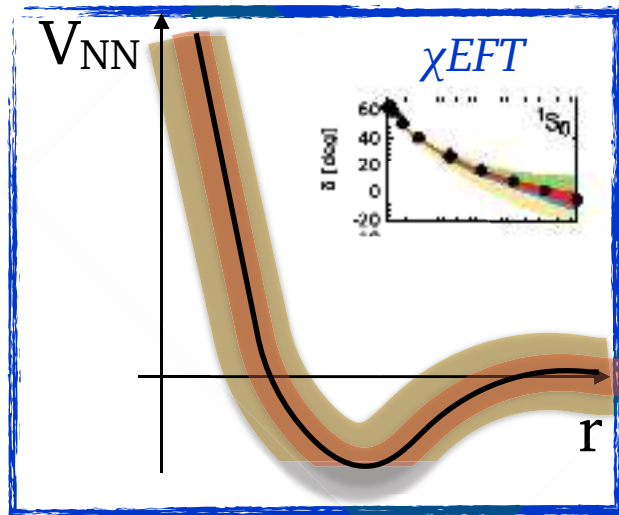
Forwards modelling

- Statistical propagation
- Bayesian analysis
- Emulators

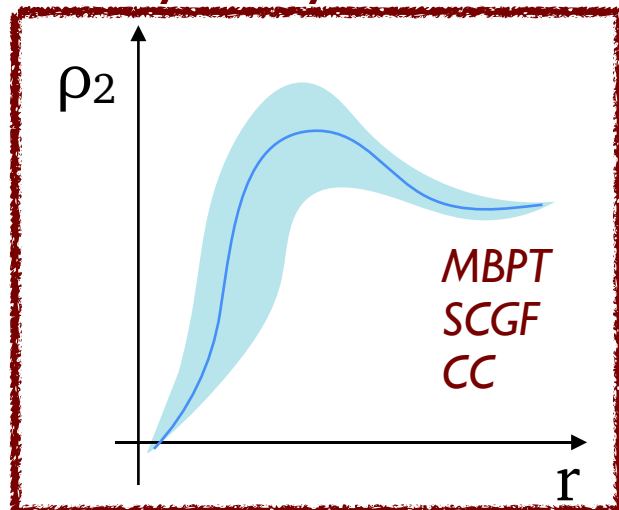


Nuclear uncertainty quantification

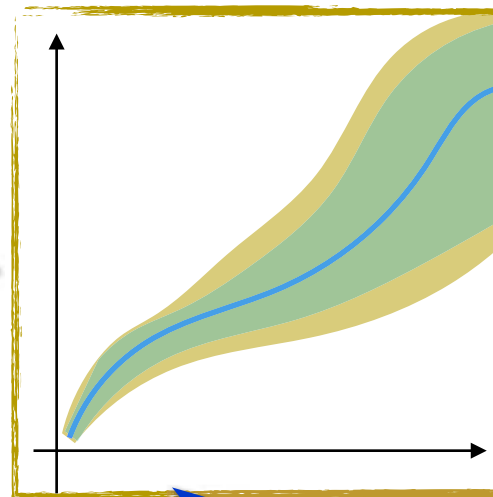
Hamiltonian



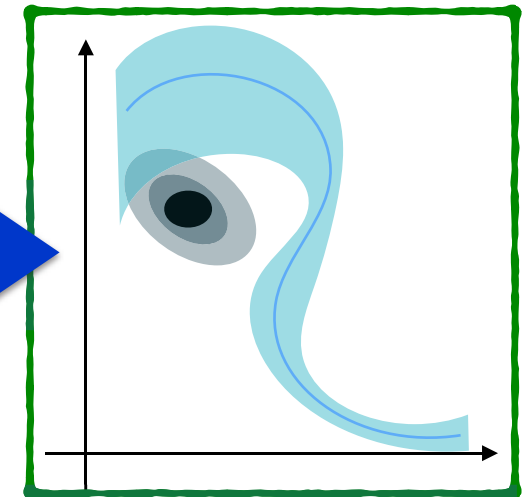
Many-body method



Astronuclear
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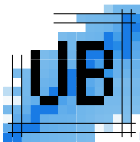


Neutron star
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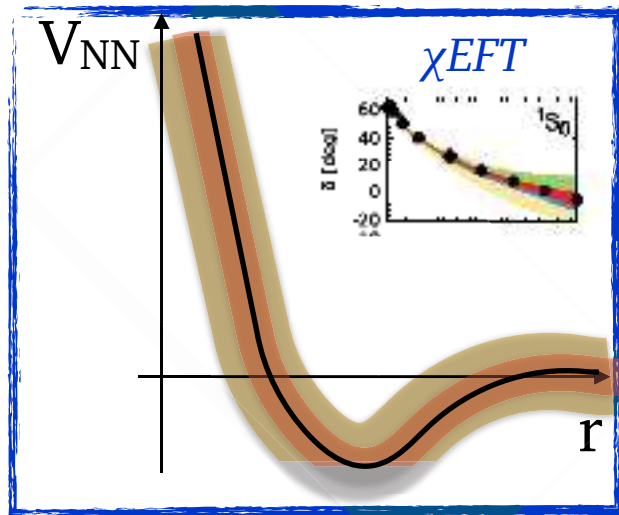
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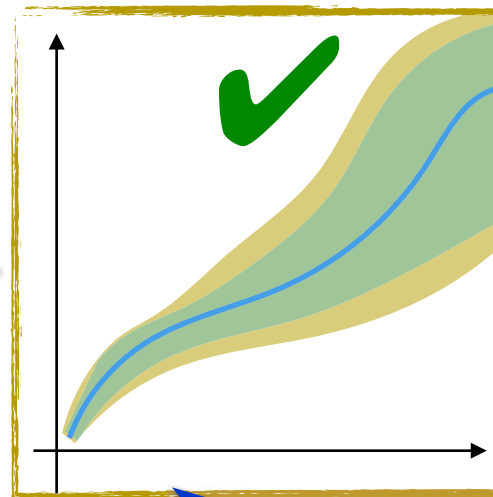


Nuclear uncertainty quantification

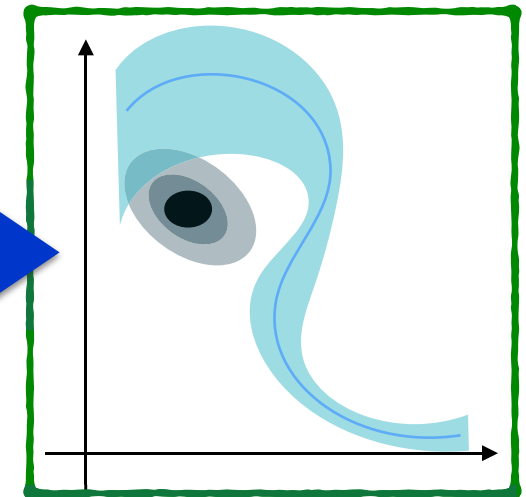
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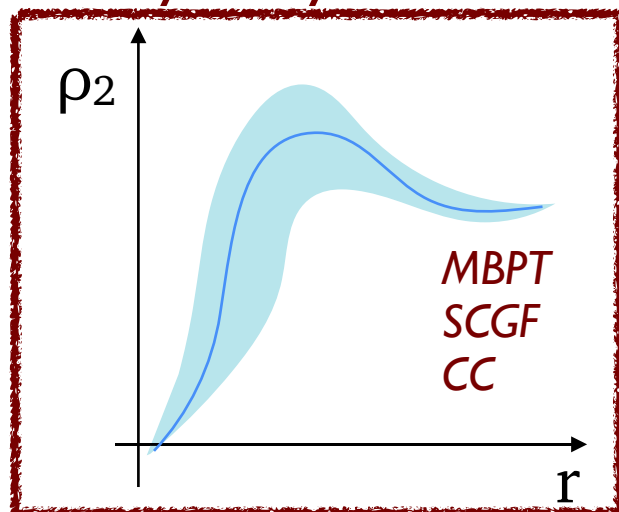
Astronuclear
property



Neutron star
observations

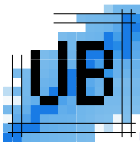


Many-body method



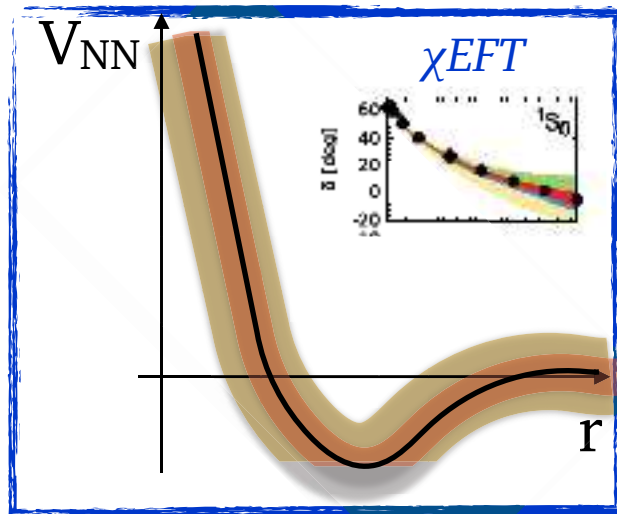
Forwards modelling

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- Bayesian analysis
- Emulators

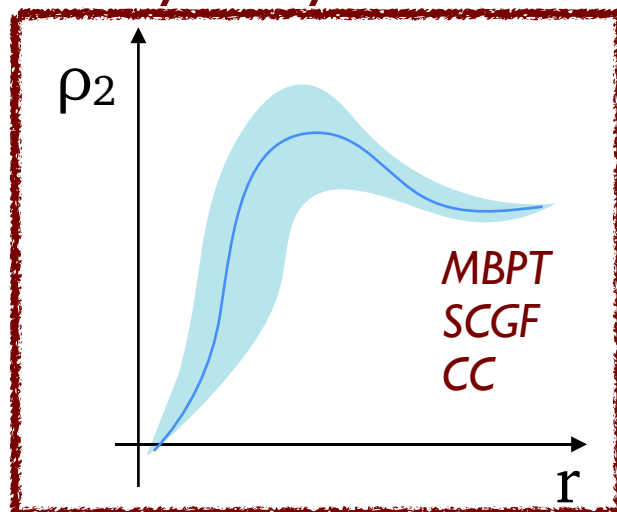


Nuclear uncertainty quantification

Hamiltonian

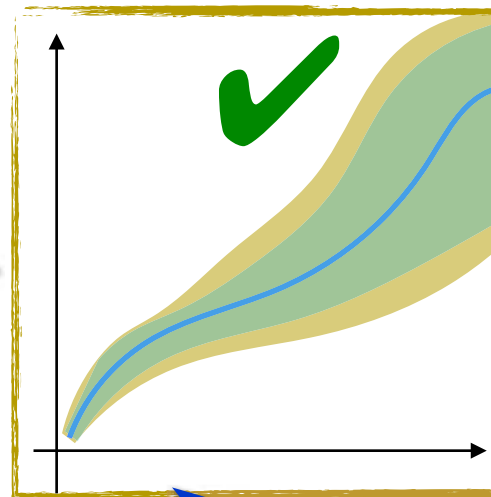


Many-body method

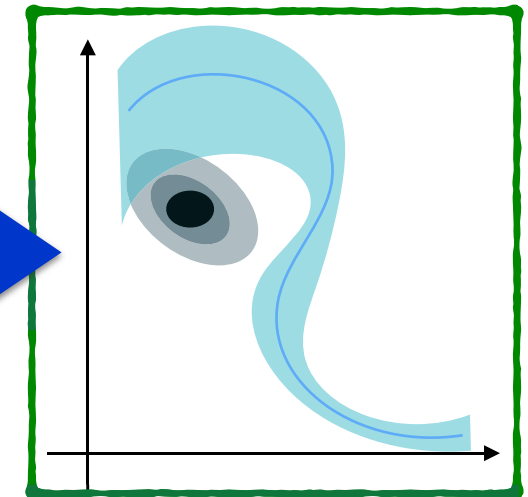


Forwards

Astronuclear
property

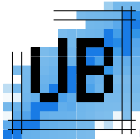


Neutron star
observations

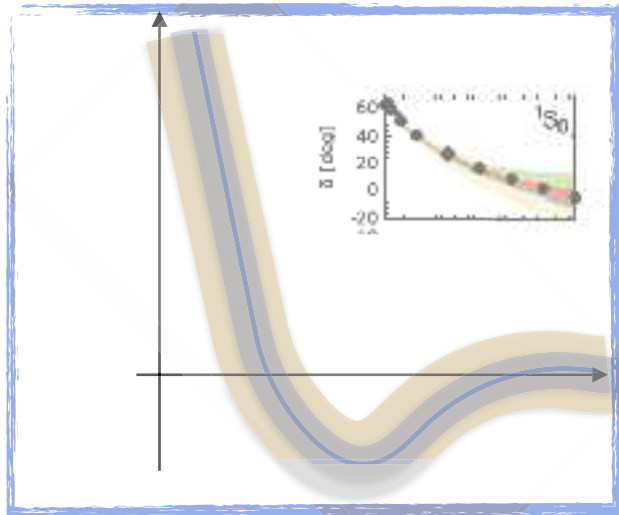


Forwards modelling

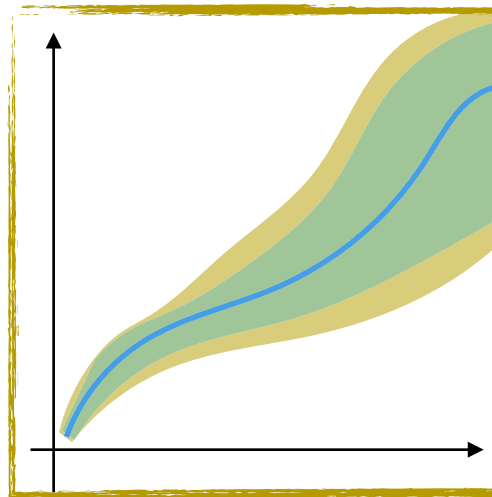
- Statistical propagation
- Bayesian analysis
- Emulators



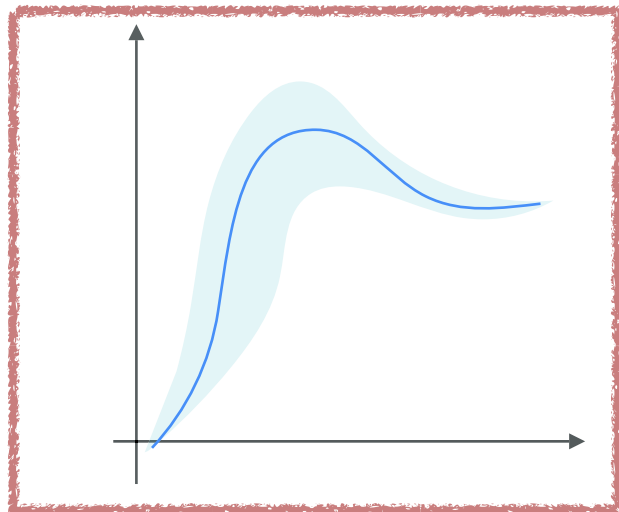
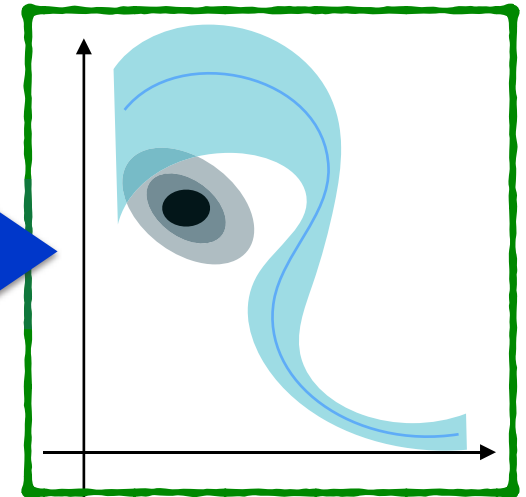
Nuclear uncertainty quantification

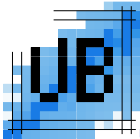


Astronuclear
property

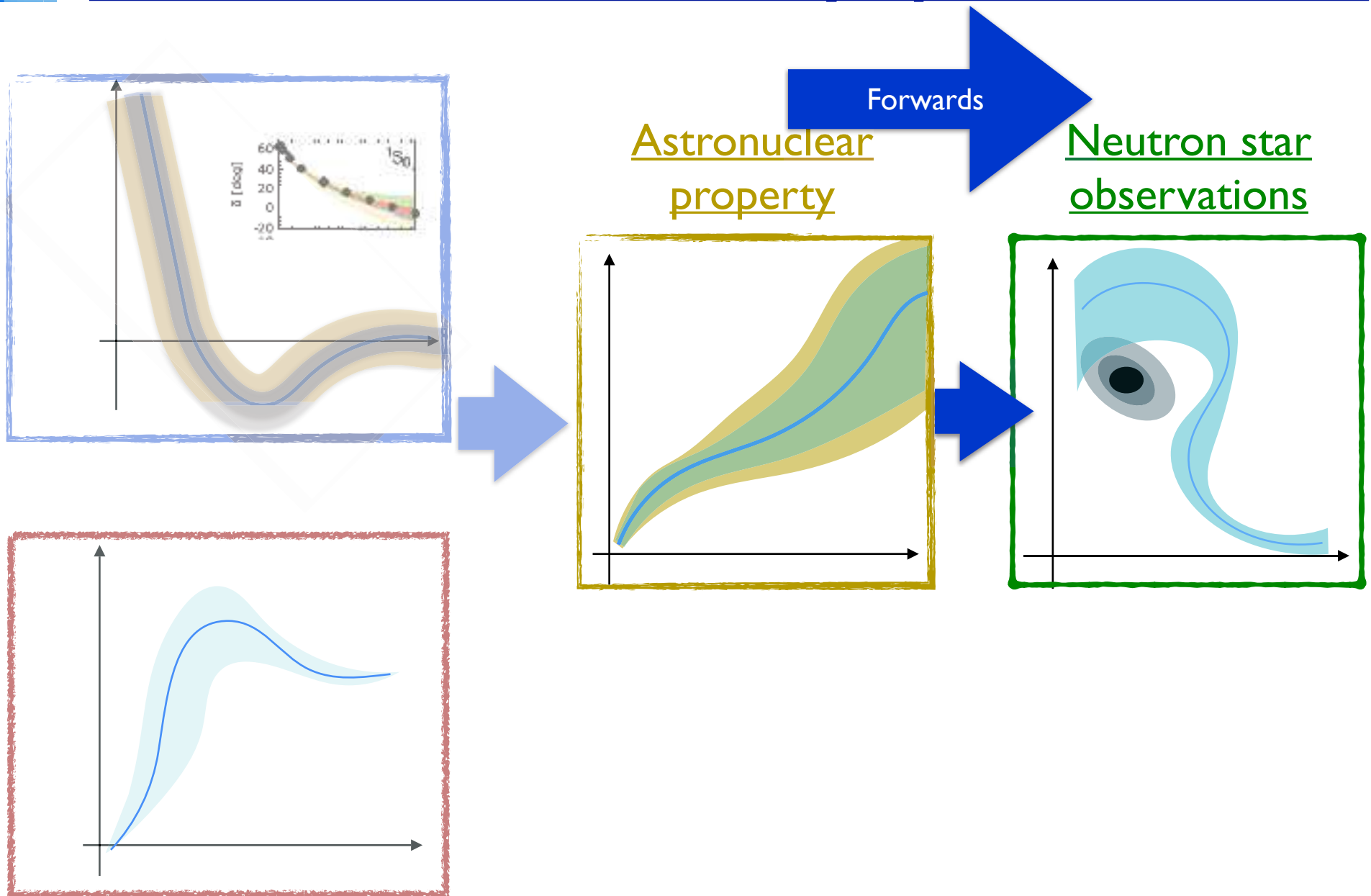


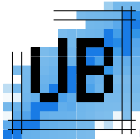
Neutron star
observations



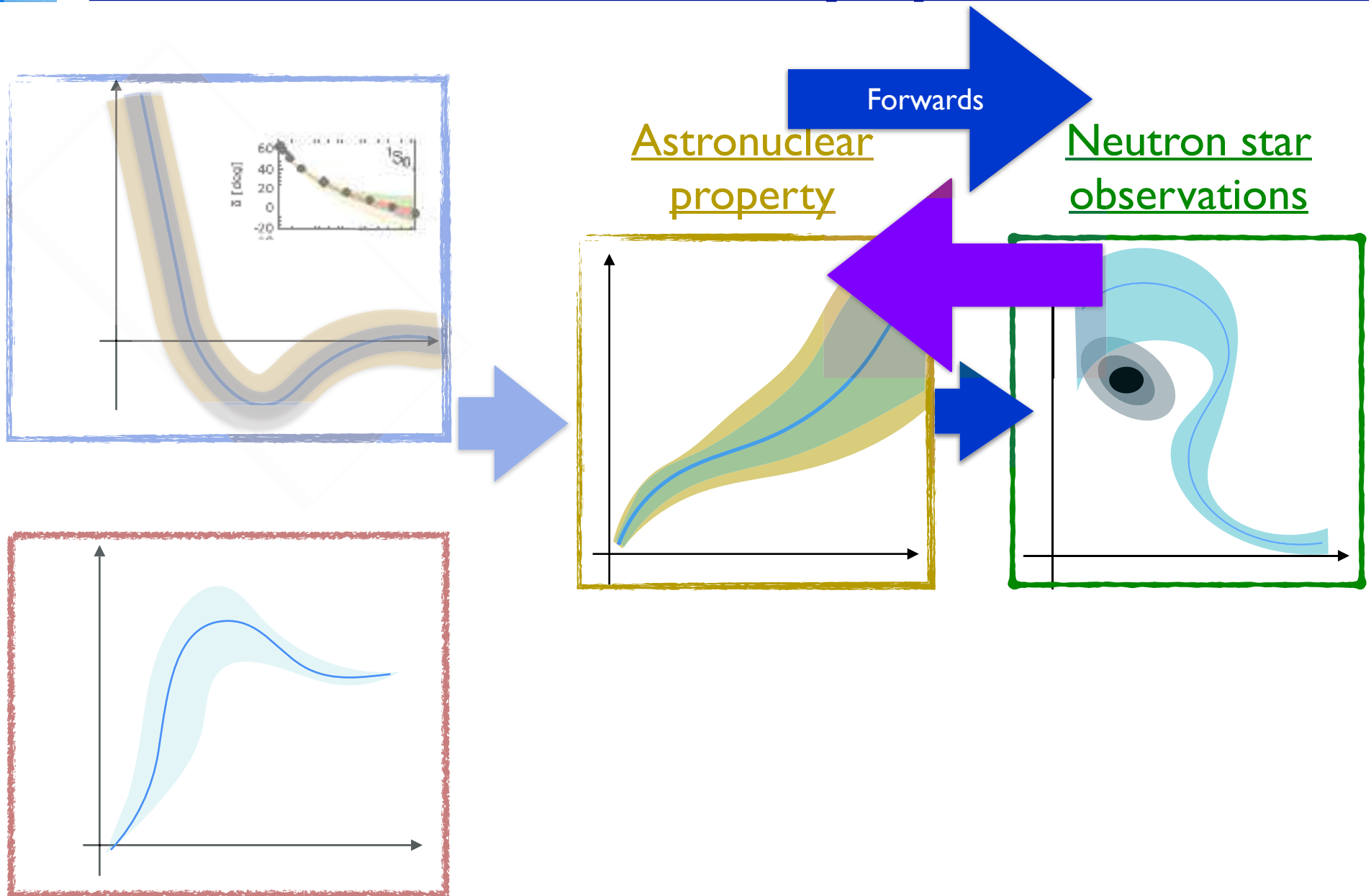


Nuclear uncertainty quantification

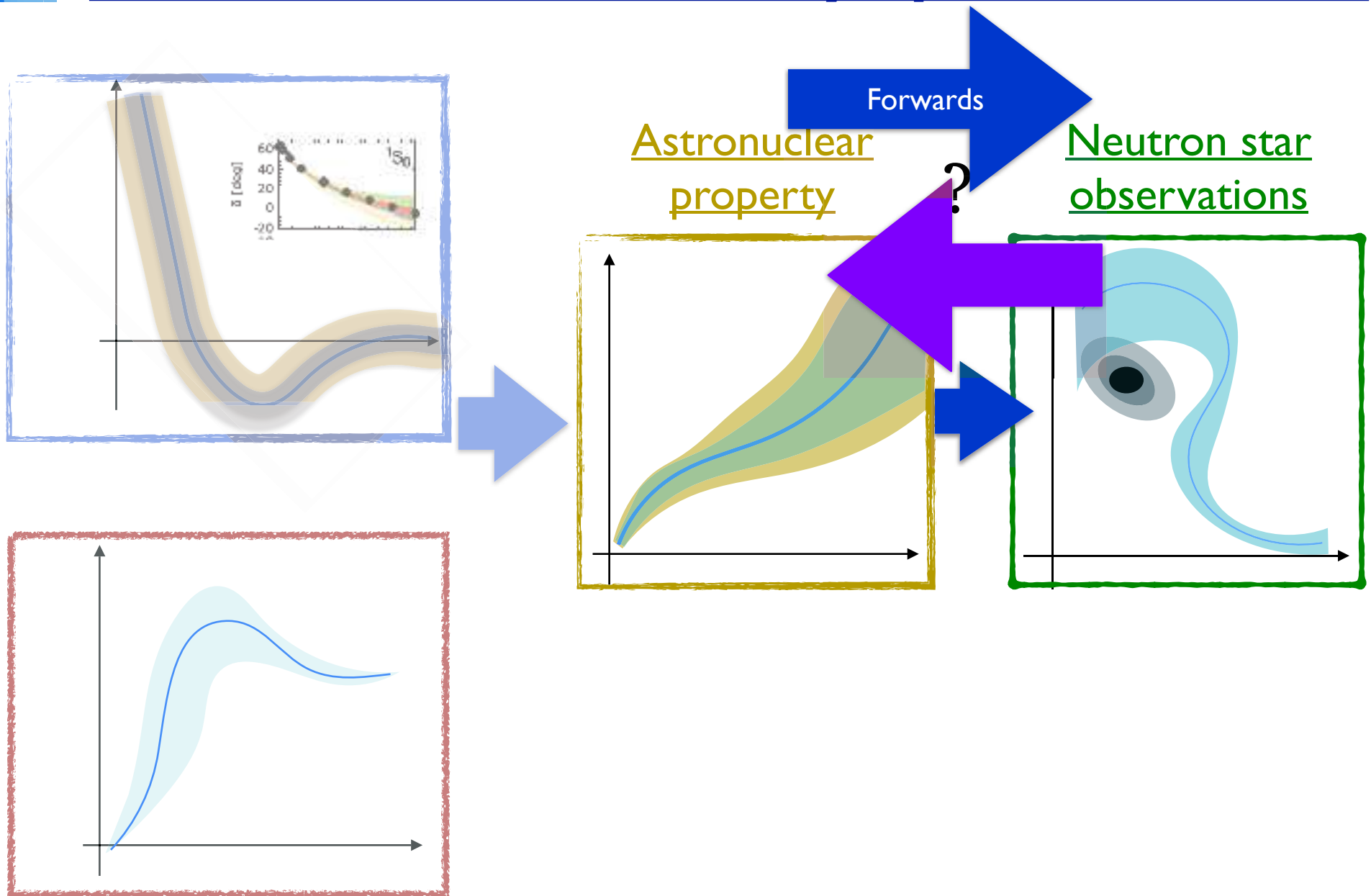




Nuclear uncertainty quantification



Nuclear uncertainty quantification



Tolman-Oppenheimer-Volkov equations

$$\frac{dP(r)}{dr} = -\frac{G}{c^2} \frac{[m_{<}(r) + 4\pi P(r)r^3] [\epsilon(r) + P(r)]}{1 - \frac{2Gm_{<}(r)}{rc^2}}$$

$$\frac{dm_{<}(r)}{dr} = 4\pi r^2 \rho(r)$$

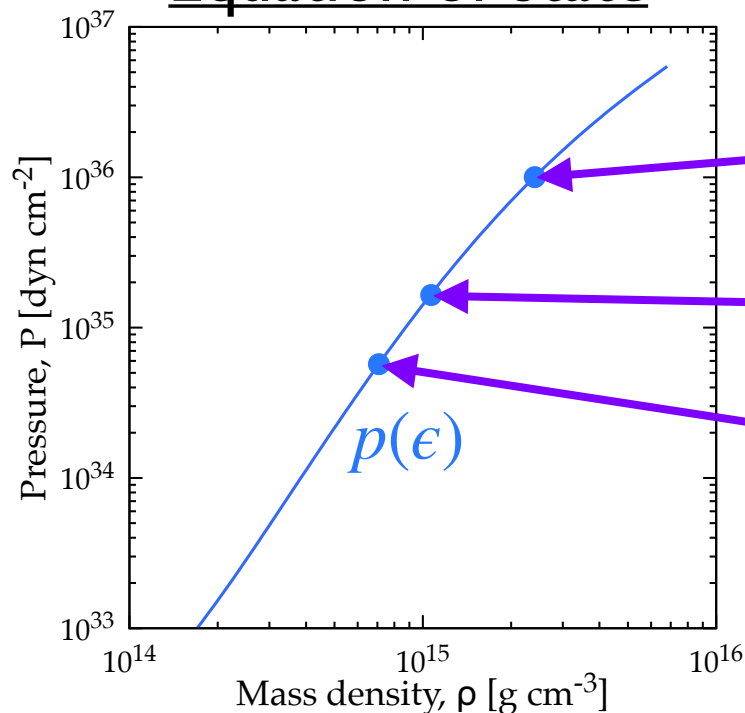
+

$$P \equiv P(\rho)$$

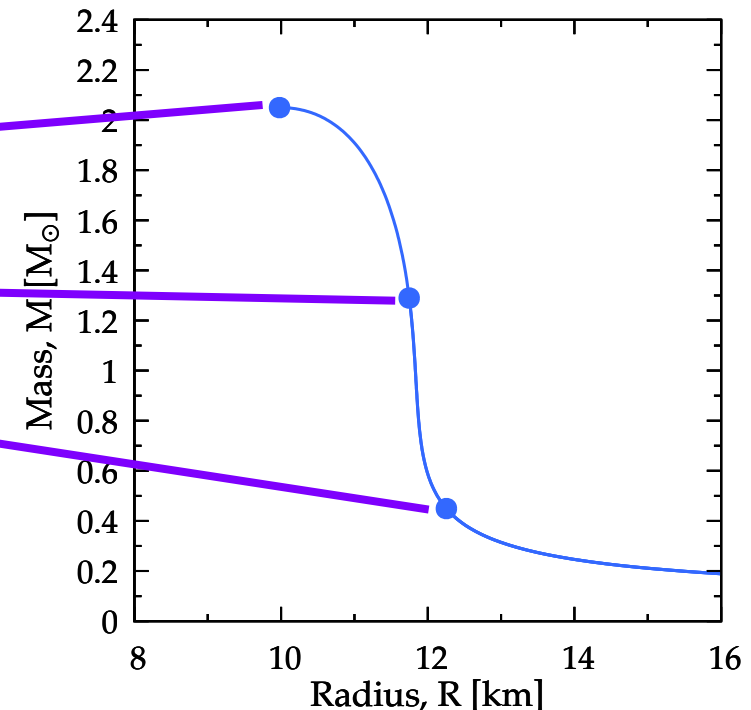
$$\epsilon = \rho c^2$$

Backwards

Equation of State



Mass-Radius



From M-R to EoS

Tolman-Oppenheimer-Volkov equations

$$\frac{dP(r)}{dr} = -\frac{G}{c^2} \frac{[m_{<}(r) + 4\pi P(r)r^3] [\epsilon(r) + P(r)]}{1 - \frac{2Gm_{<}(r)}{rc^2}}$$

$$\frac{dm_{<}(r)}{dr} = 4\pi r^2 \rho(r)$$

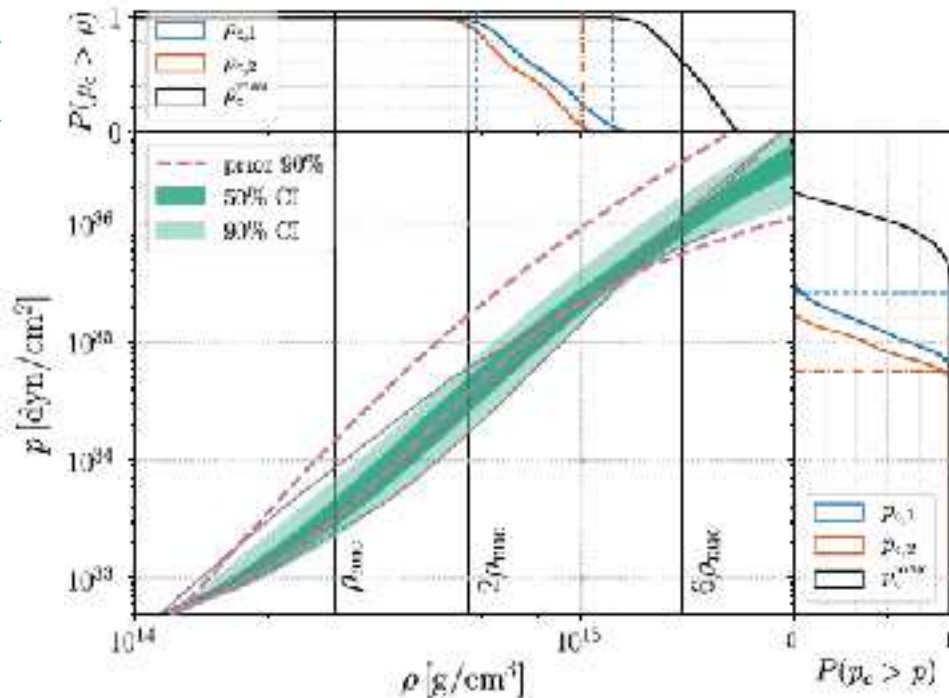
+

$$P \equiv P(\rho)$$

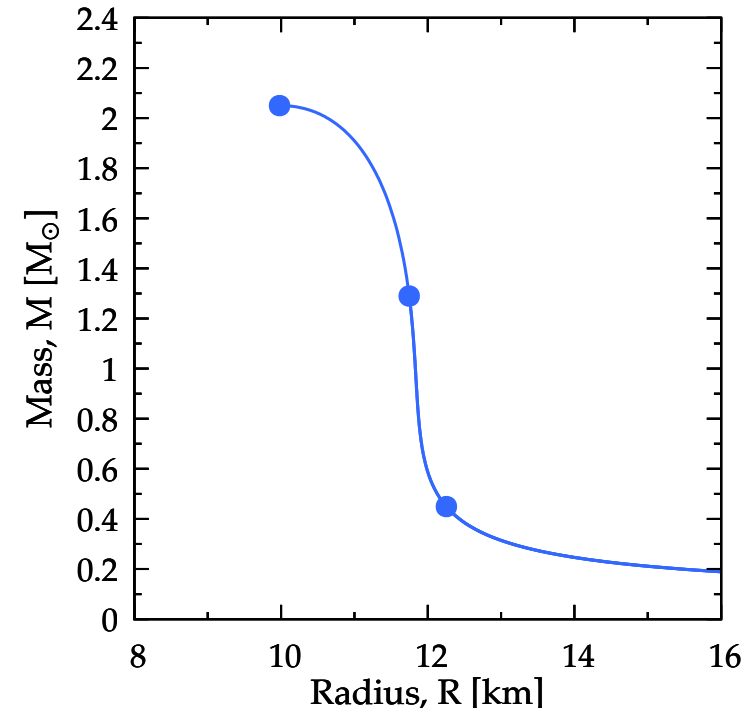
$$\epsilon = \rho c^2$$

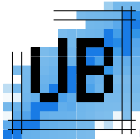
Backwards

Equation of State



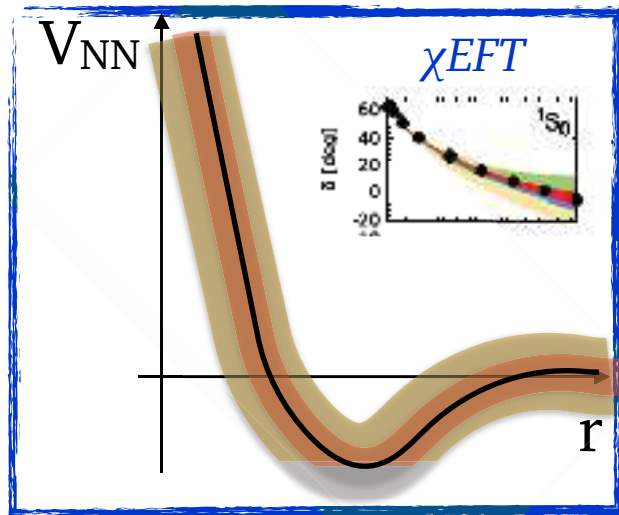
Mass-Radius



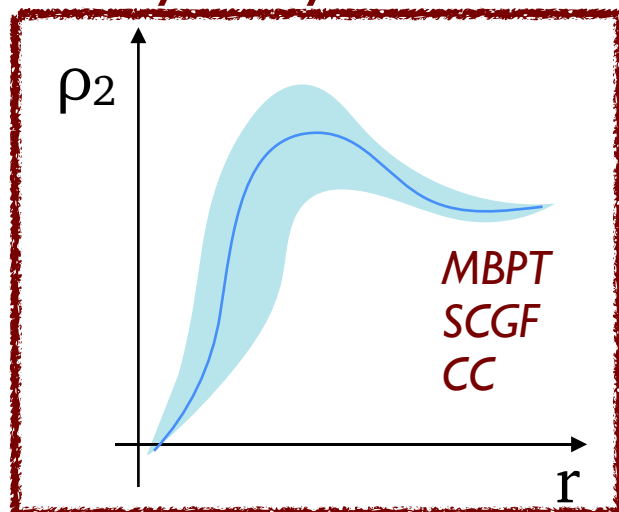


Nuclear uncertainty quantification

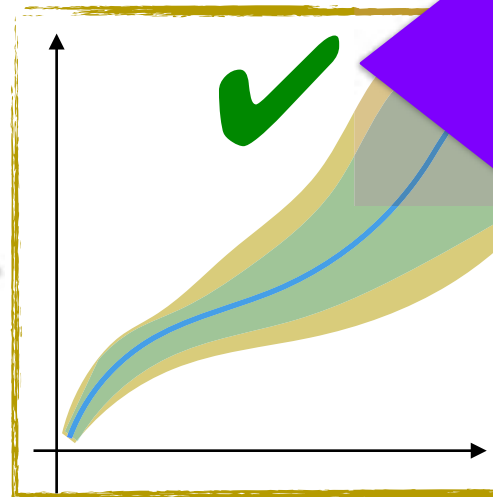
Hamiltonian



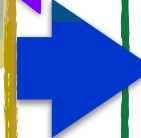
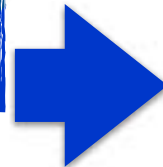
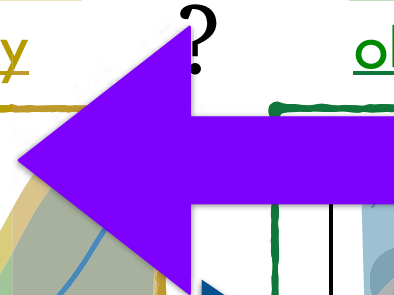
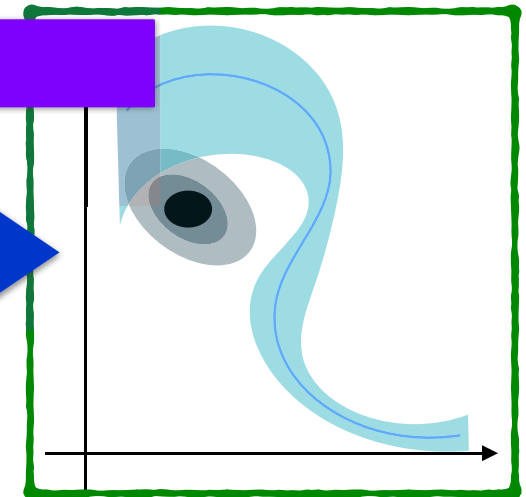
Many-body method

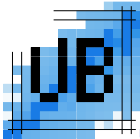


Astronuclear
property



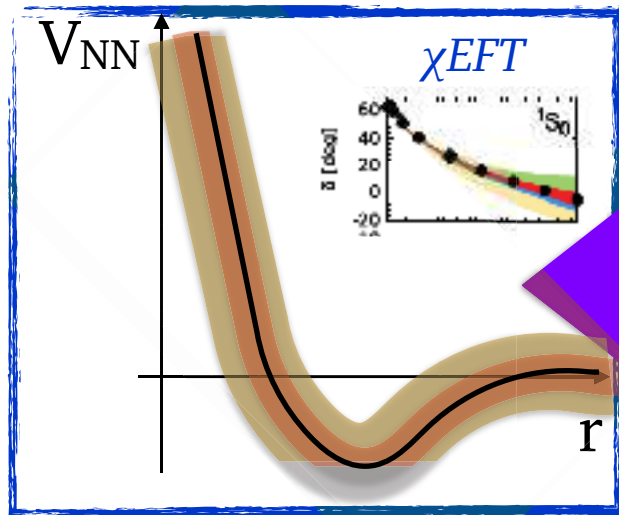
Neutron star
observations





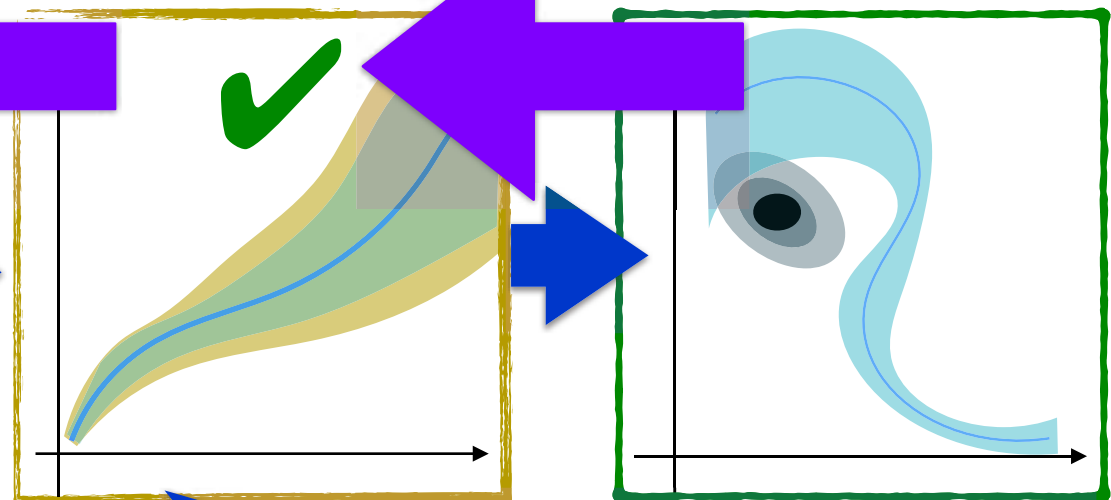
Nuclear uncertainty quantification

Hamiltonian

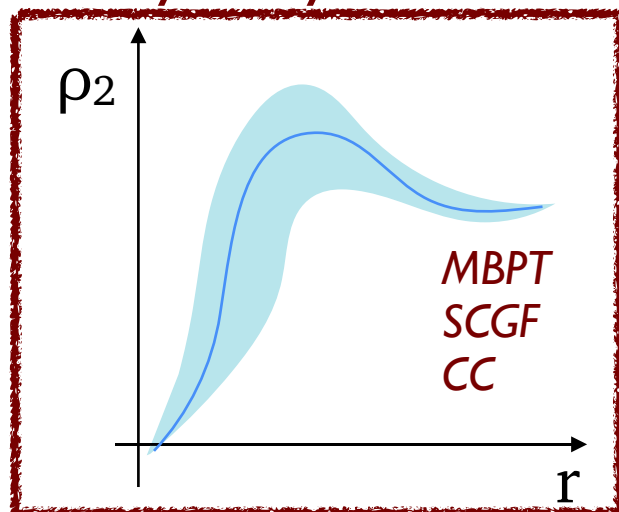


Astronuclear property

Neutron star observations



Many-body method



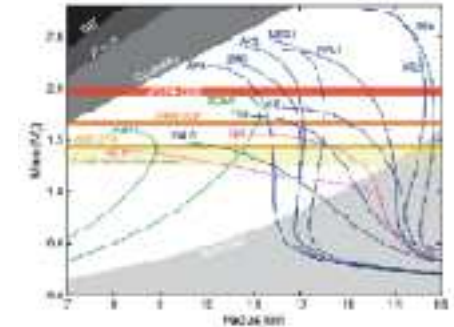
Requires

- Matter composition
- Regulators & density reach
- Degeneracies?

Neutron star modelling

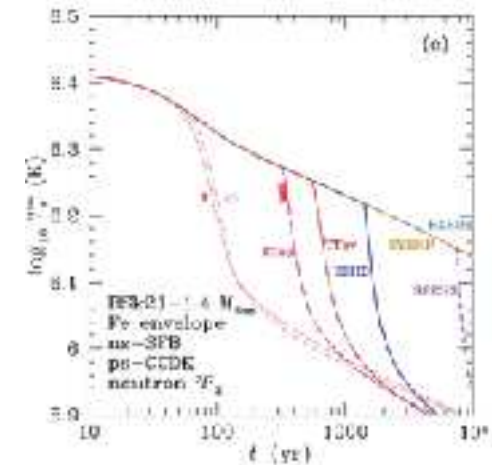
Input #1
EoS

Observable #1
Mass-Radius relation



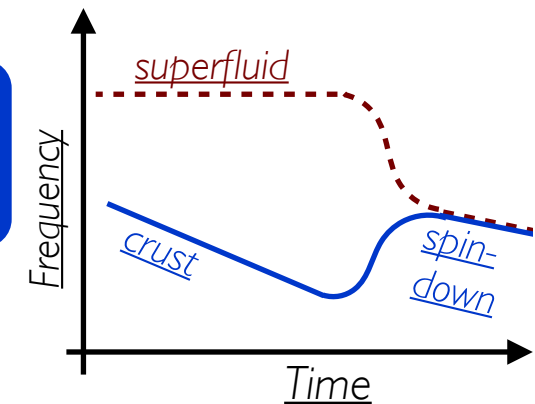
Input #2
Pairing gap

Observable #2
Cooling curve

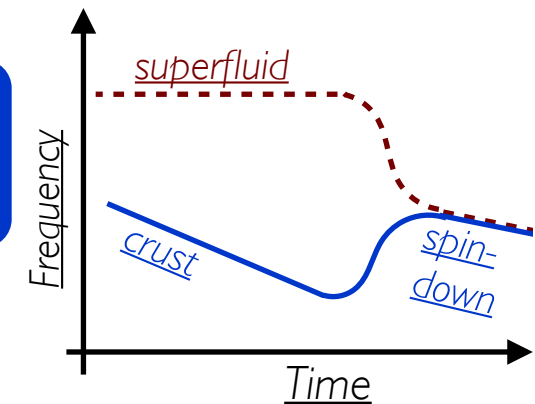
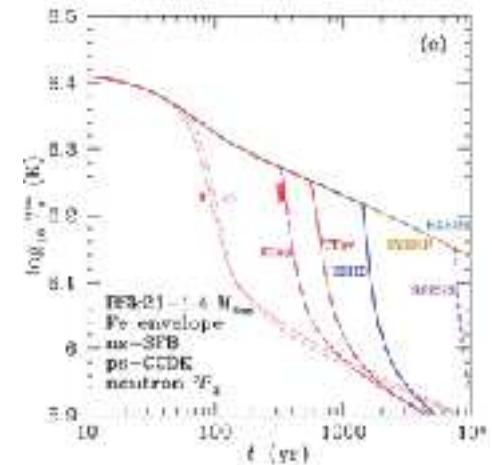
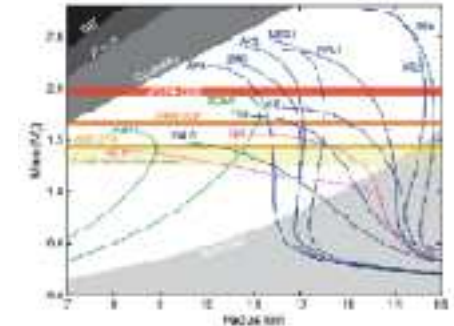
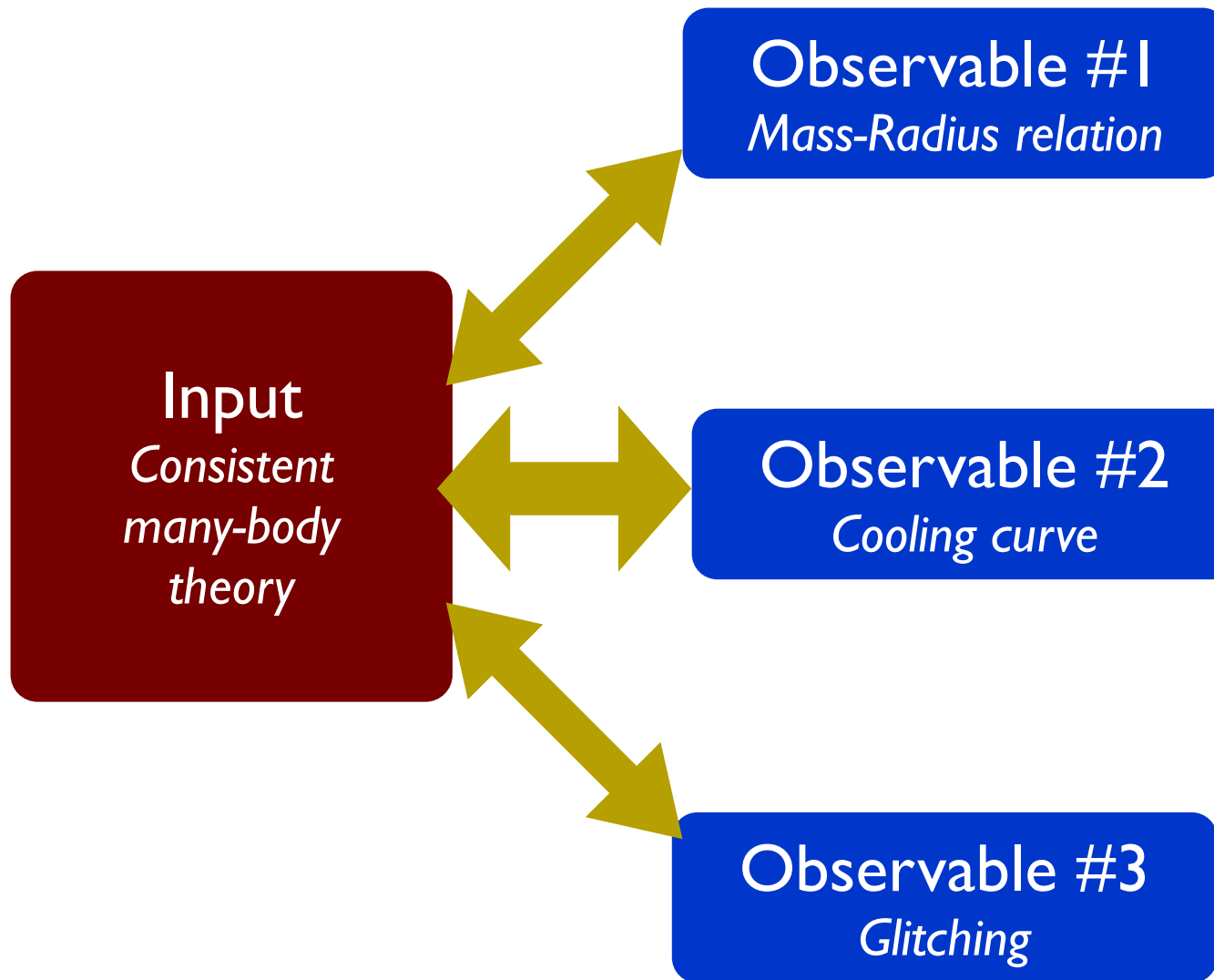


Input #3
Crust-core

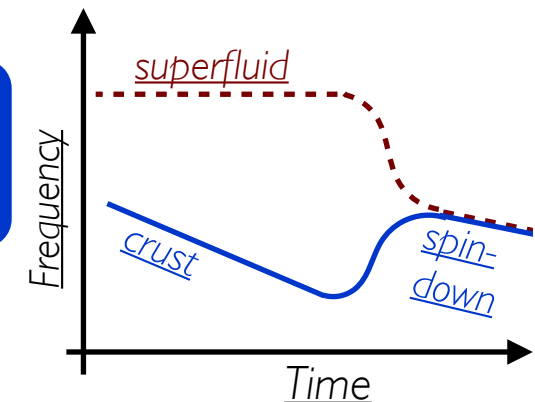
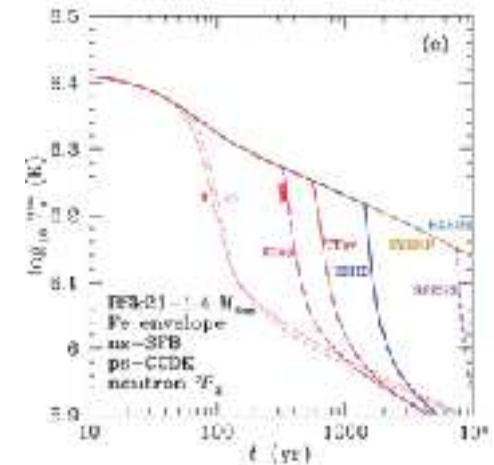
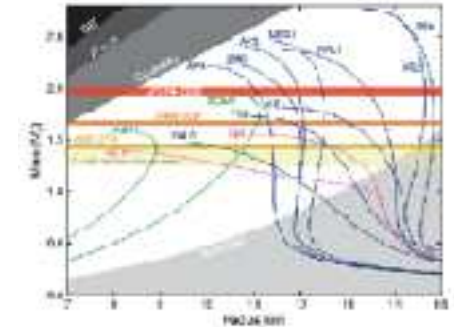
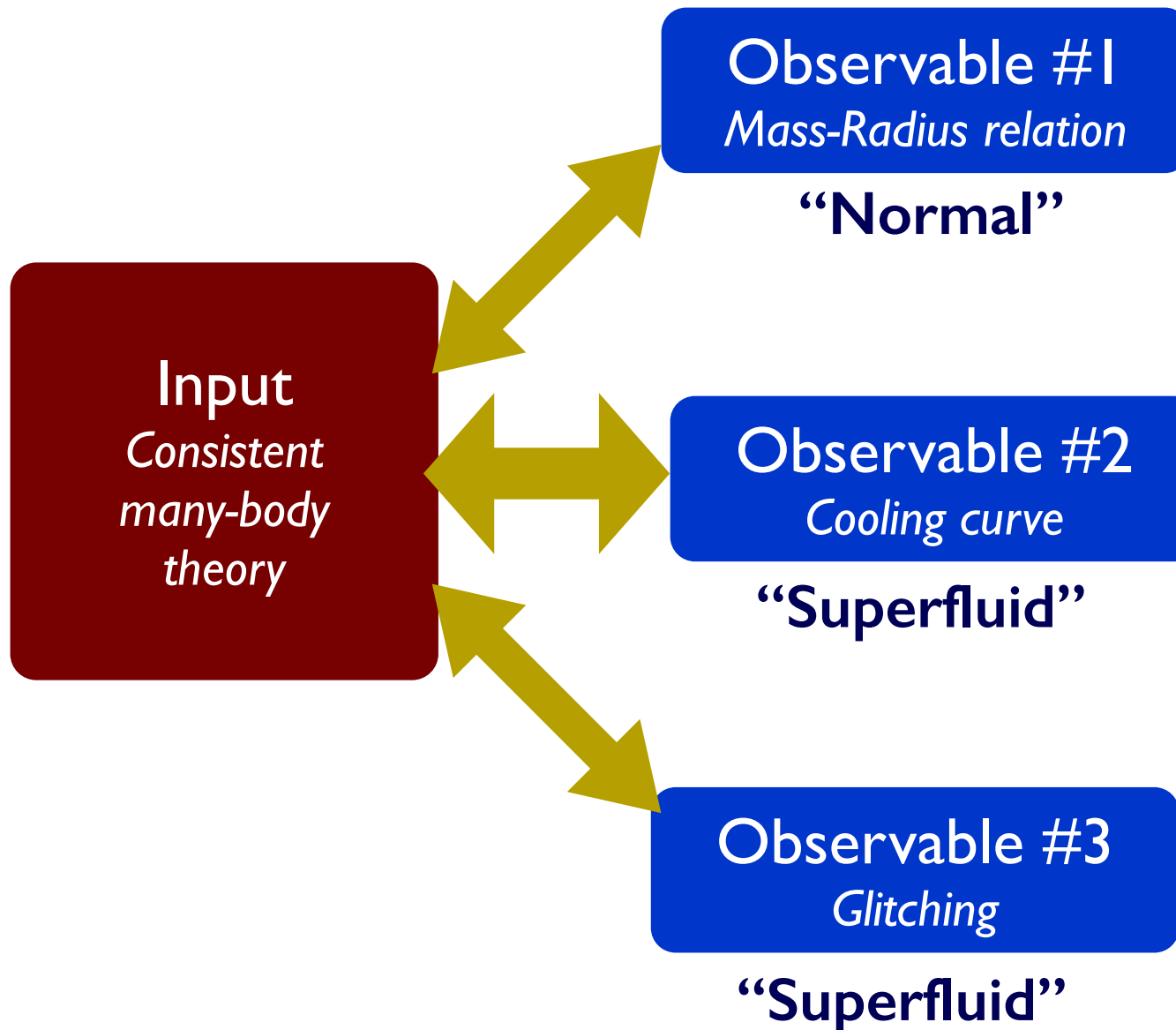
Observable #3
Glitching



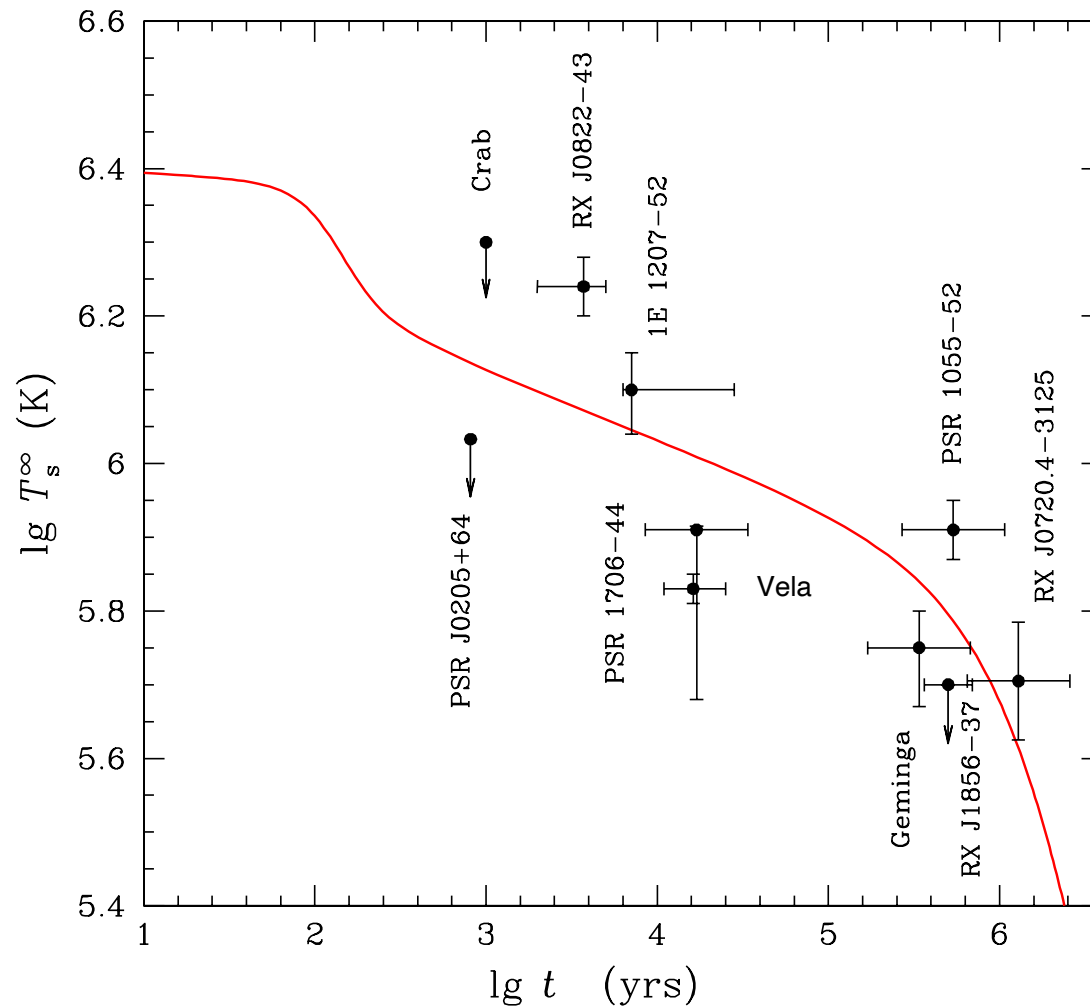
Neutron star modelling



Neutron star modelling

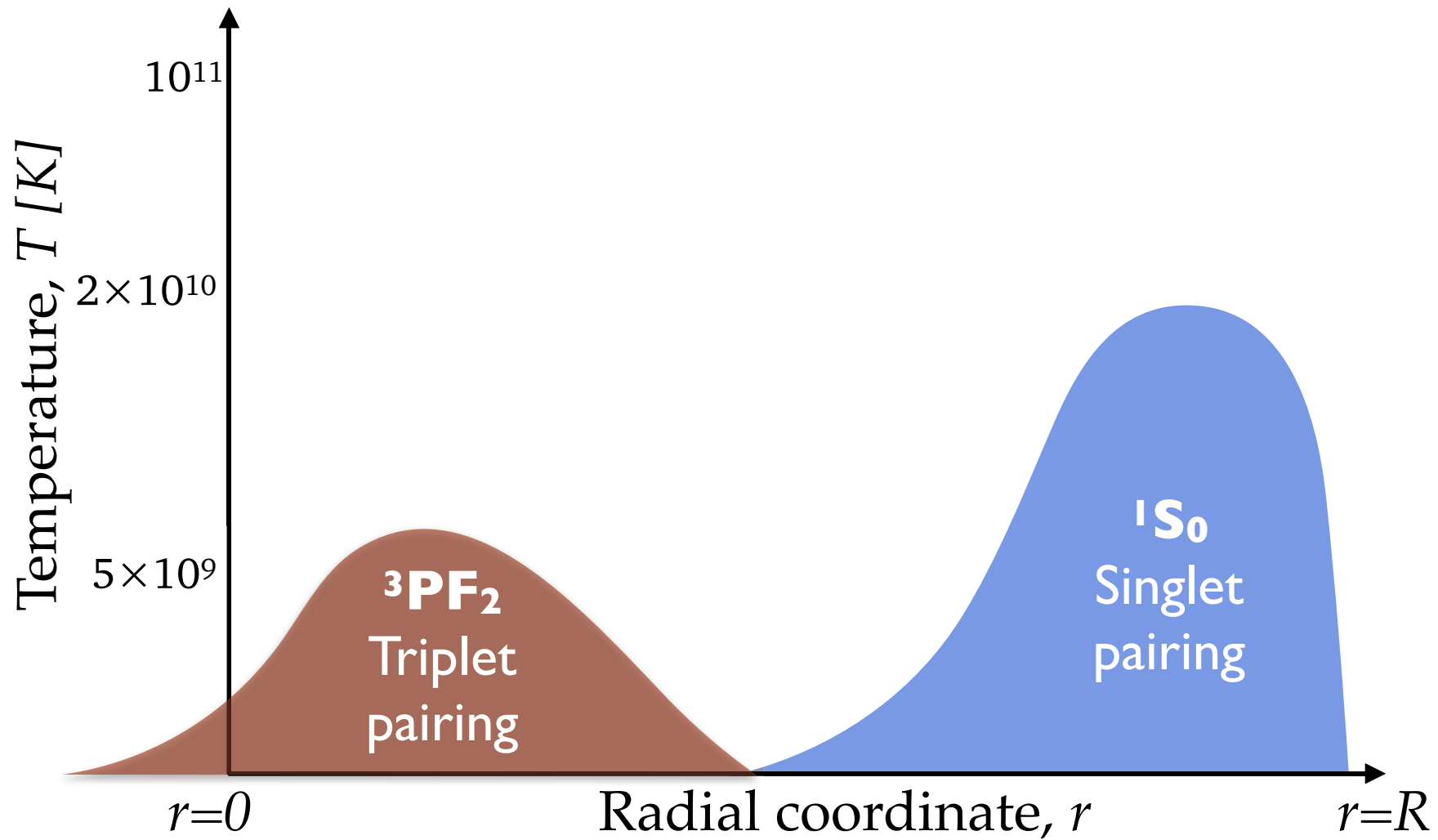


Cooling curve of neutron stars

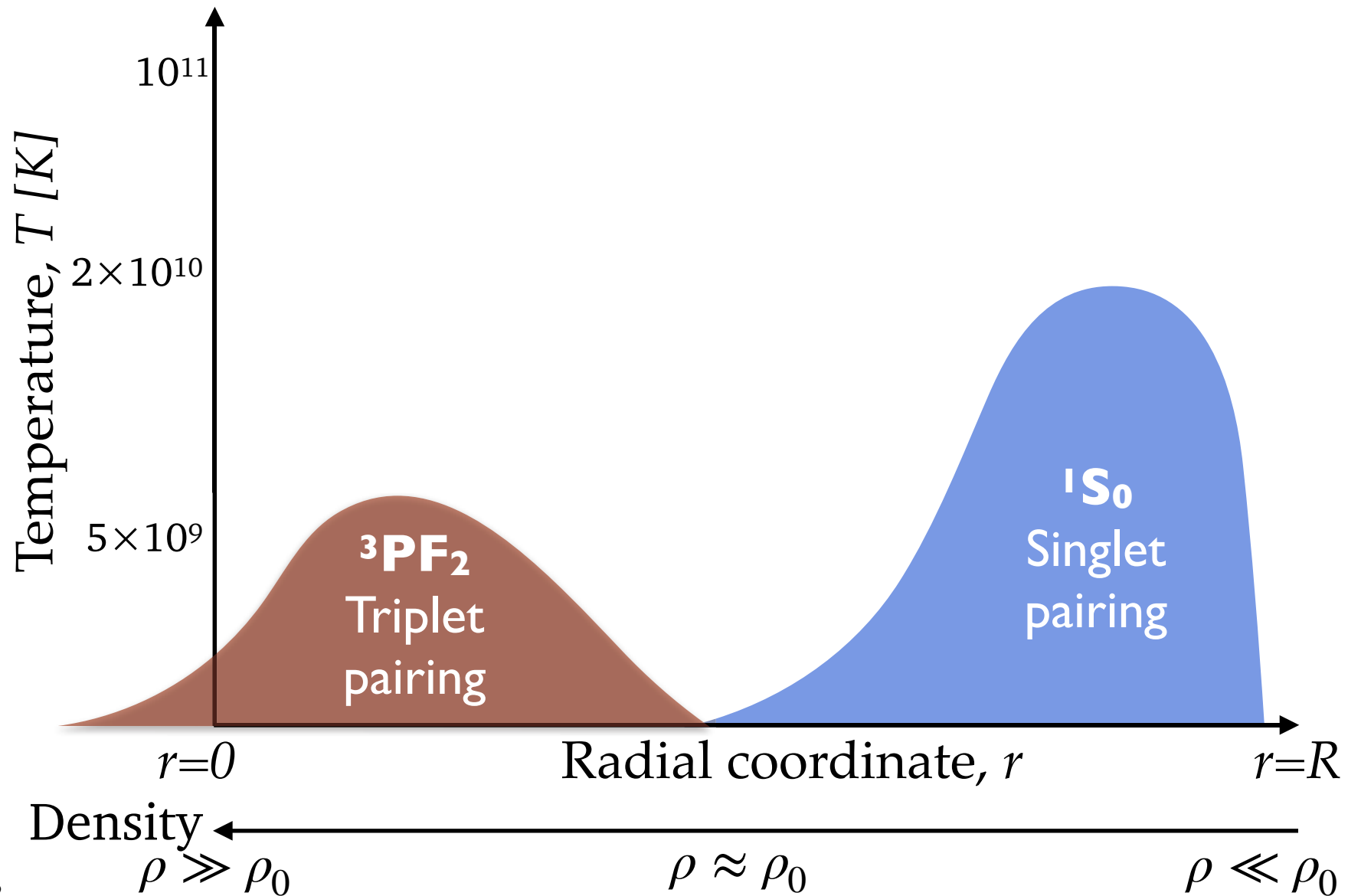


- Observational data available for a handful of NS
- Sensitive to **interior** physics (mostly **pairing**)

Pairing gaps & cooling

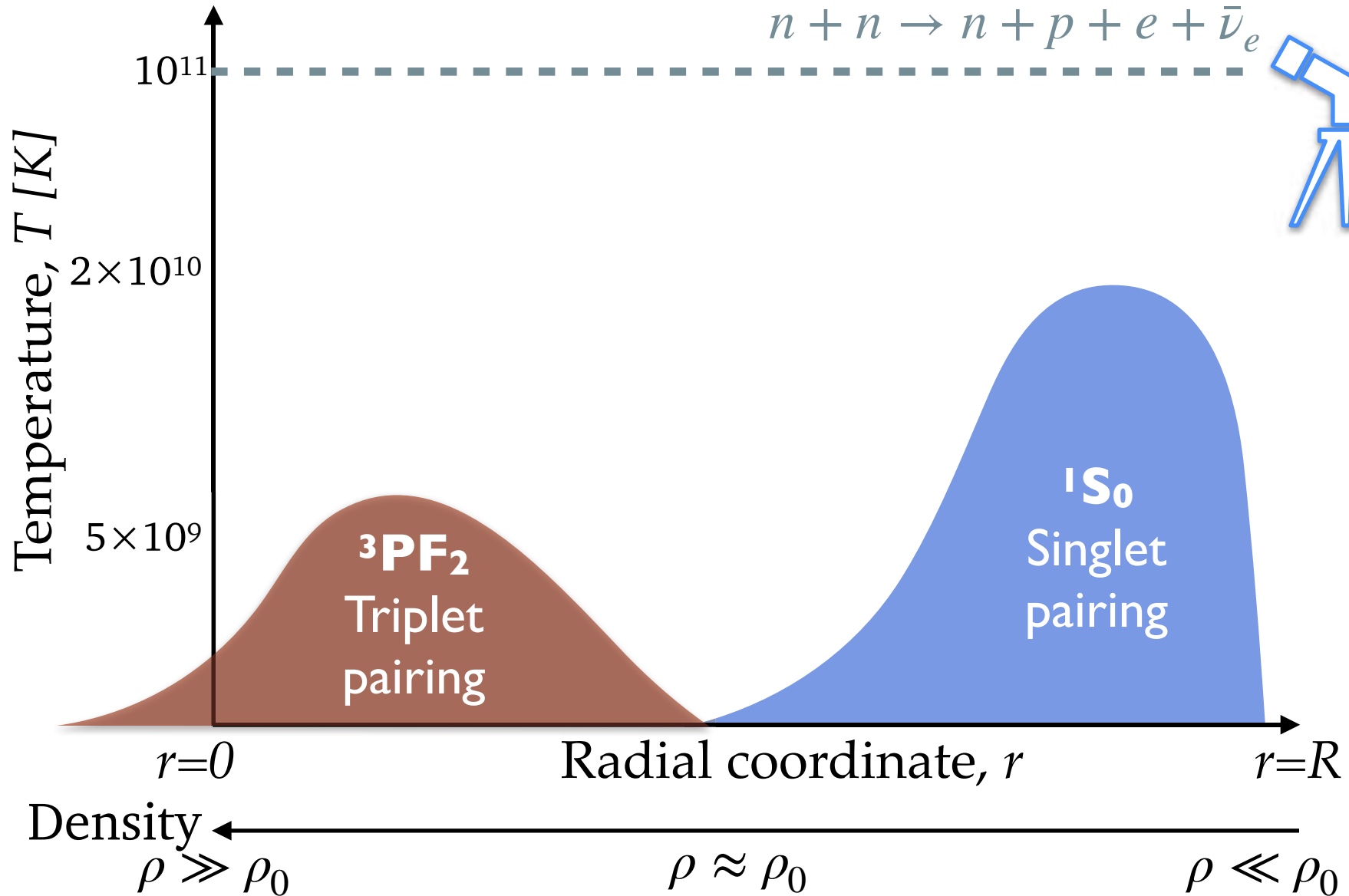


Pairing gaps & cooling

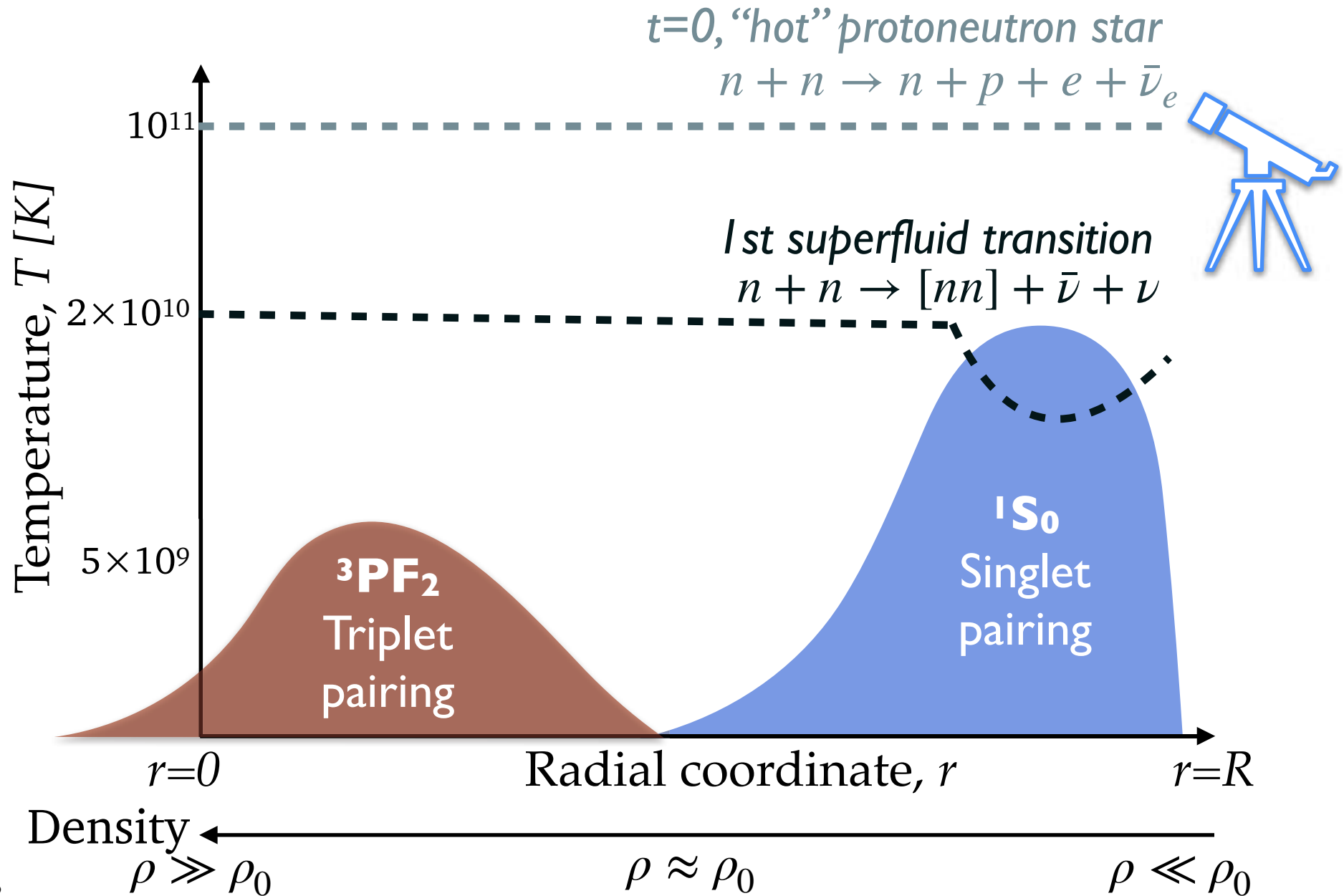


Pairing gaps & cooling

$t=0$, "hot" protoneutron star



Pairing gaps & cooling

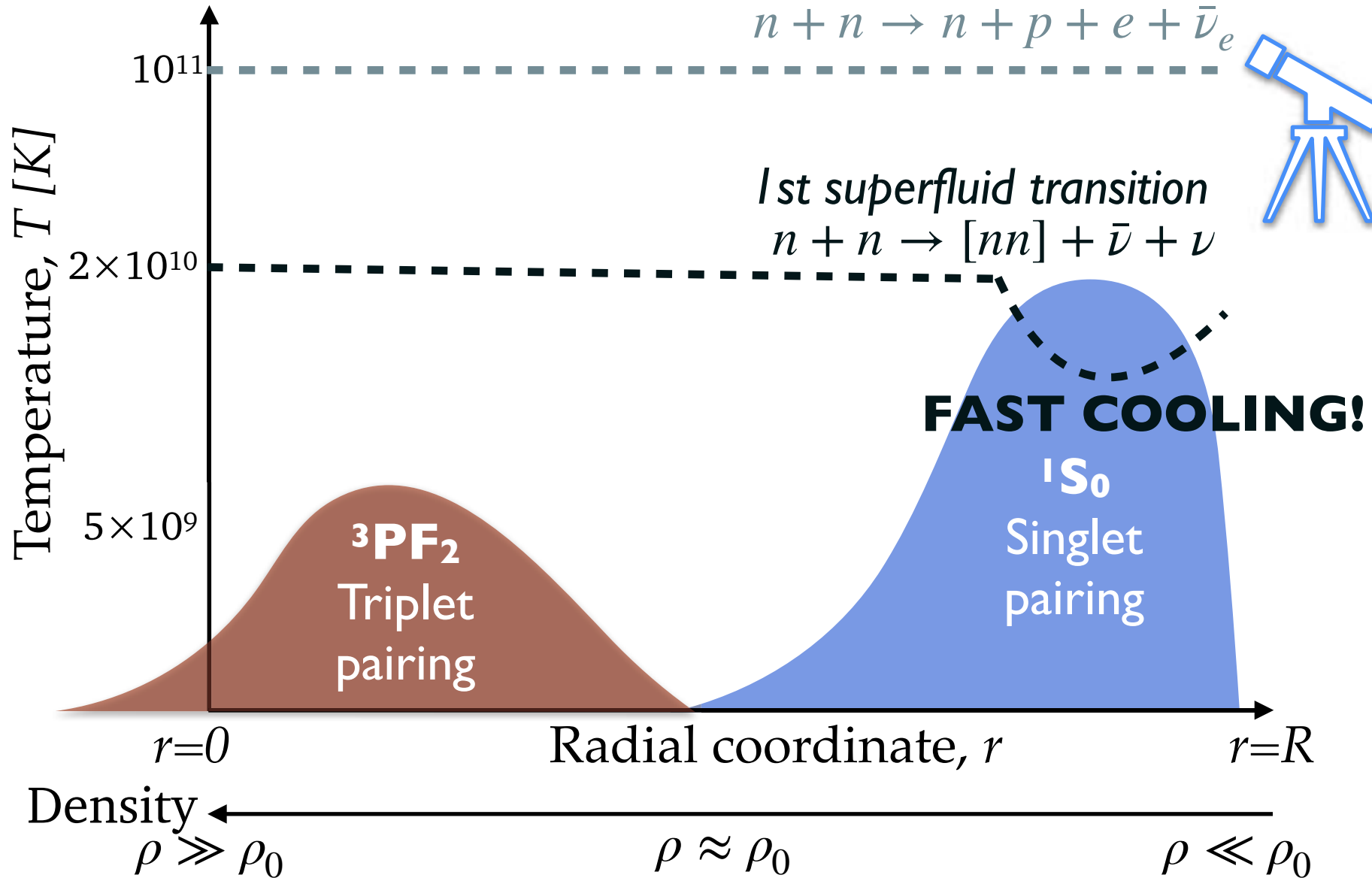


Pairing gaps & cooling

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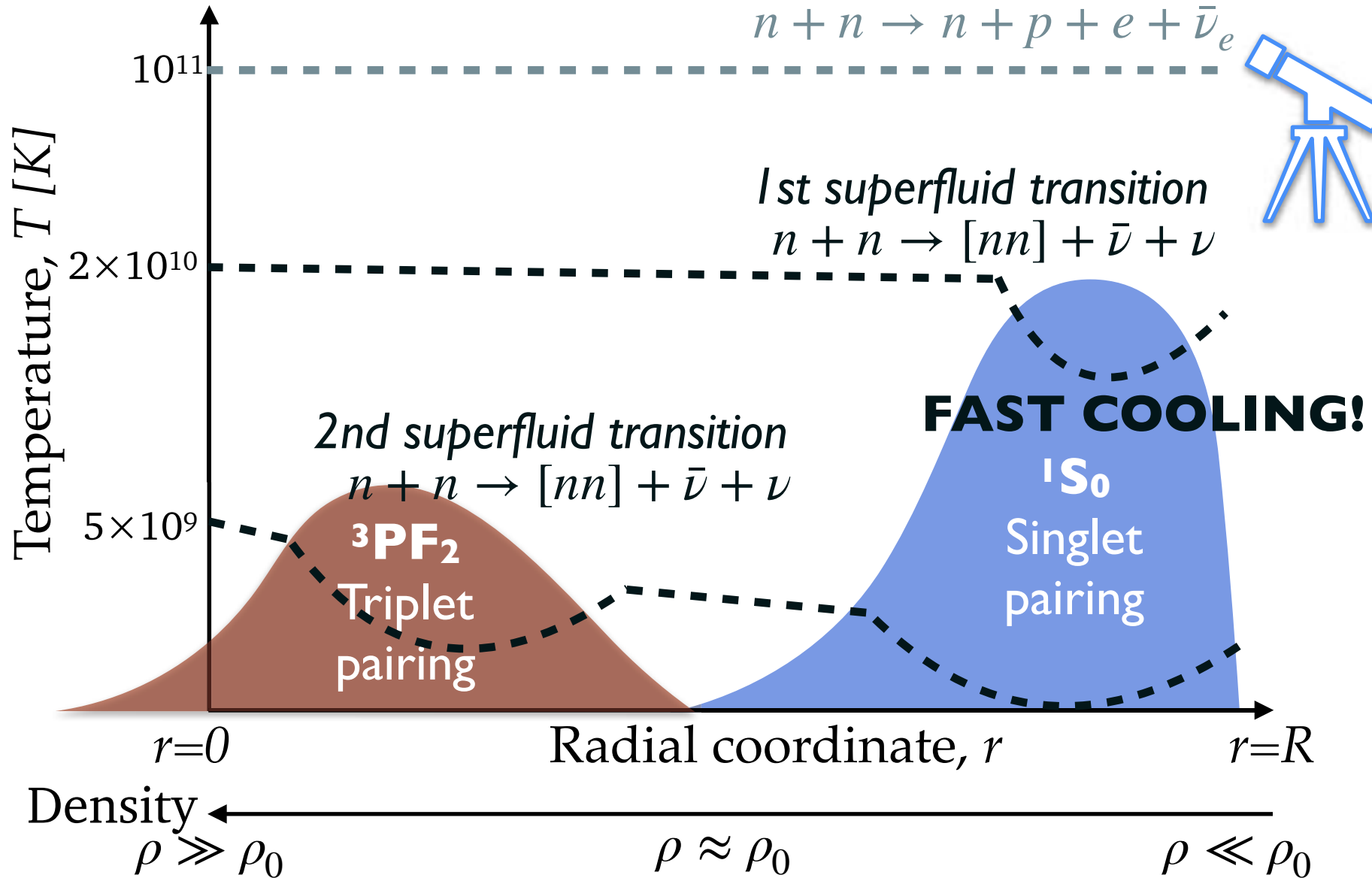


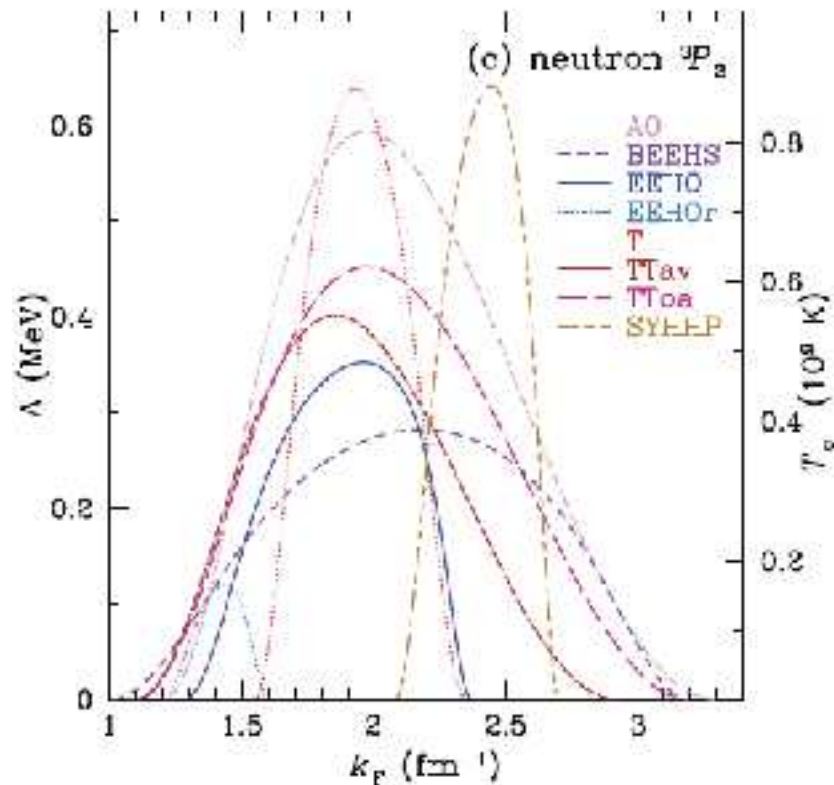
1st superfluid transition



Pairing gaps & cooling

$t=0$, "hot" protoneutron star



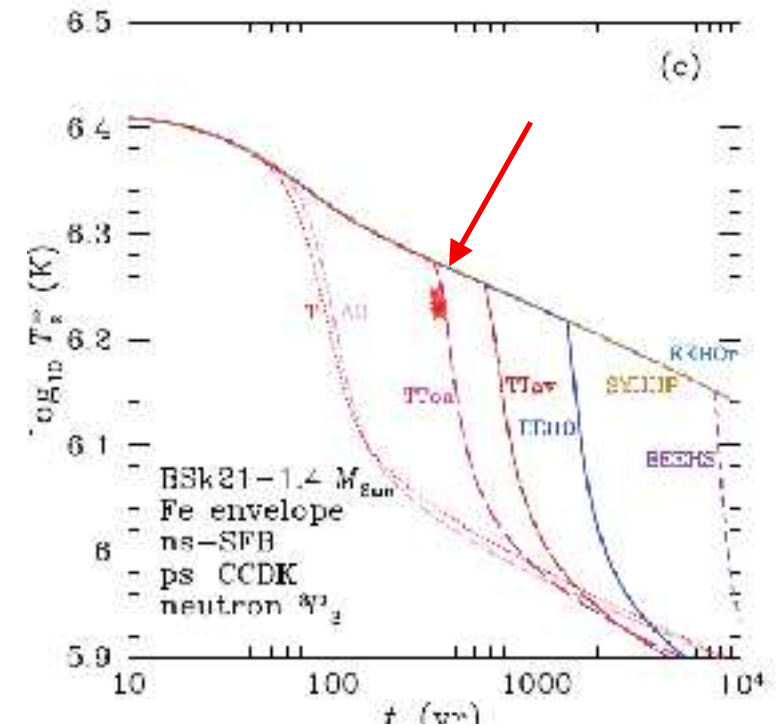


Ho, et al., PRC **91** 015806 (2015)

Page, et al., PRL **106** 081101 (2011)

Ingredients

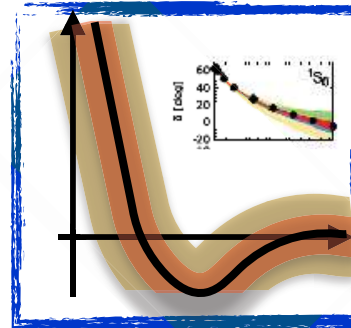
- (a) Mass of pulsar
- (b) EoS (determines radius)
- (c) Internal composition
- (d) **Pairing gaps** (1S_0 & 3P_F channels)
- (e) Atmosphere composition



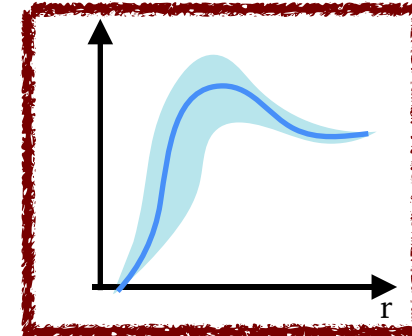
Name	Process	Emissivity (erg cm ⁻³ s ⁻¹)
Modified Urca (neutron branch)	$n + n \rightarrow n + p + e^- + \bar{\nu}_e$ $n + p + e^- \rightarrow n + n + \nu_e$	$\sim 2 \times 10^{21} RT_9^8$
Modified Urca (proton branch)	$p + n \rightarrow p + p + e^- + \bar{\nu}_e$ $p + p + e^- \rightarrow p + n + \nu_e$	$\sim 10^{21} RT_9^8$
Bremsstrahlungs	$n + n \rightarrow n + n + \nu + \bar{\nu}$ $n + p \rightarrow n + p + \nu + \bar{\nu}$ $p + p \rightarrow p + p + \nu + \bar{\nu}$	$\sim 10^{19} RT_9^8$
Cooper pair	$n + n \rightarrow [nn] + \nu + \bar{\nu}$ $p + p \rightarrow [pp] + \nu + \bar{\nu}$	$\sim 5 \times 10^{21} RT_9^7$ $\sim 5 \times 10^{19} RT_9^7$
Direct Urca (nucleons)	$n \rightarrow p + e^- + \bar{\nu}_e$ $p + e^- \rightarrow n + \nu_e$	$\sim 10^{27} RT_9^6$

- Motivation

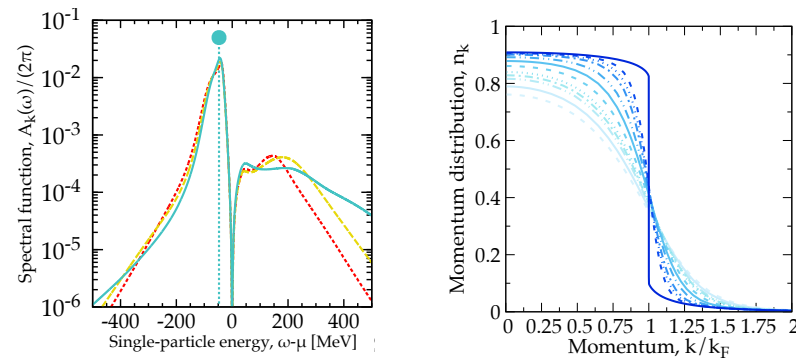
Hamiltonian



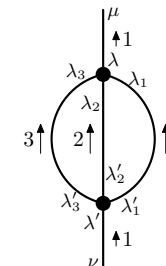
Many-body method

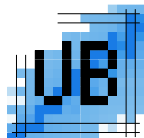


- “Normal” self-consistent Green’s functions



- Nambu-Covariant Green’s functions





Self-Consistent Green's Functions

(ρ, T)



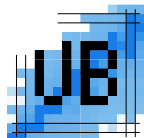
2N & 3N forces

$$\text{wavy line} = \text{dashed line} + \text{dotted line} \circlearrowleft$$

$$\text{wavy line with cross} = \text{dashed line} \circlearrowleft + \frac{1}{2} \text{dotted line} \circlearrowleft \circlearrowleft$$

Carbone, Rios & Polls PRC **88** 044302 (2013);
PRC **90**, 054322 (2014);
Carbone PhD Thesis

Forward



Self-Consistent Green's Functions

(ρ, T)



2N & 3N forces

$$\begin{aligned} \text{wavy line} &= \text{dashed line} + \text{dotted line} \\ \text{wavy line with cross} &= \text{dashed line} + \frac{1}{2} \text{dotted line} \end{aligned}$$



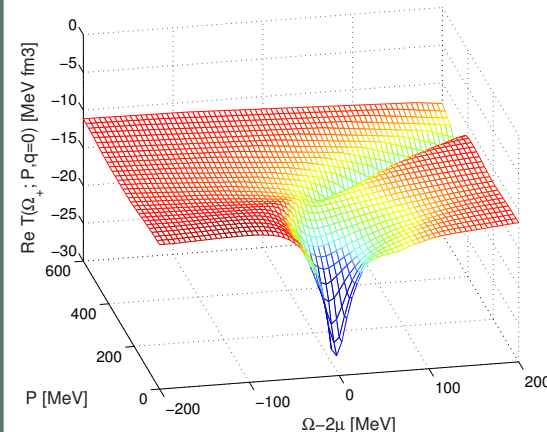
In-medium interaction

$$T = \text{wavy line} + \text{loop diagram with T}$$

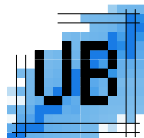
Carbone, Rios & Polls PRC **88** 044302 (2013);
PRC **90**, 054322 (2014);
Carbone PhD Thesis

Forward

T-matrix at $T=5$ MeV

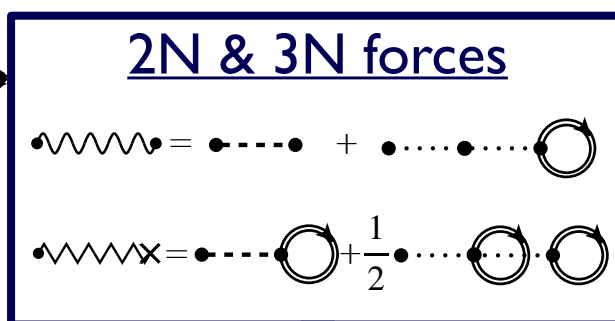


Ramos, Polls & Dickhoff, NPA **503** 1 (1989)
Alm *et al.*, PRC **53** 2181 (1996)
Dewulf *et al.*, PRL **90** 152501 (2003)
Frick & Muther, PRC **68** 034310 (2003)
Rios, PhD Thesis, U. Barcelona (2007)
Soma & Bozek, PRC **78** 054003 (2008)
Rios & Soma PRL **108** 012501 (2012)



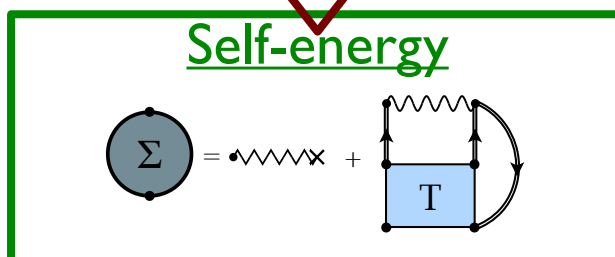
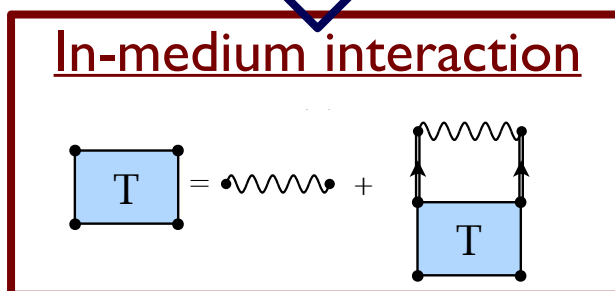
Self-Consistent Green's Functions

(ρ, T)

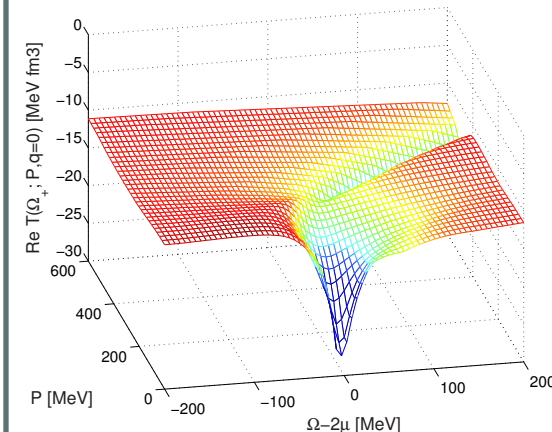


Carbone, Rios & Polls PRC **88** 044302 (2013);
PRC **90**, 054322 (2014);
Carbone PhD Thesis

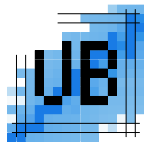
Forward



T-matrix at $T=5$ MeV

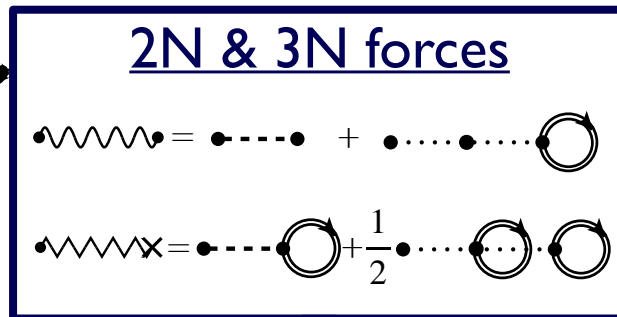


Ramos, Polls & Dickhoff, NPA **503** 1 (1989)
Alm *et al.*, PRC **53** 2181 (1996)
Dewulf *et al.*, PRL **90** 152501 (2003)
Frick & Muther, PRC **68** 034310 (2003)
Rios, PhD Thesis, U. Barcelona (2007)
Soma & Bozek, PRC **78** 054003 (2008)
Rios & Soma PRL **108** 012501 (2012)



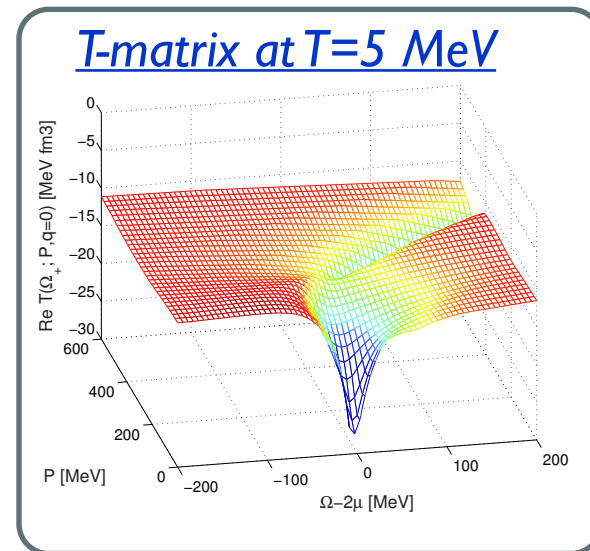
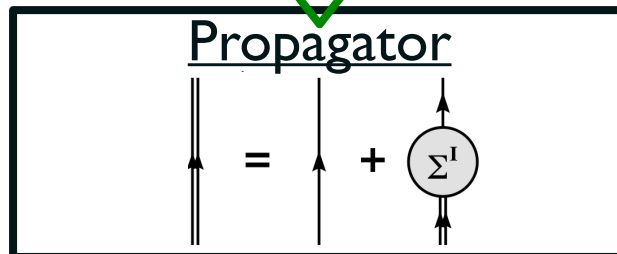
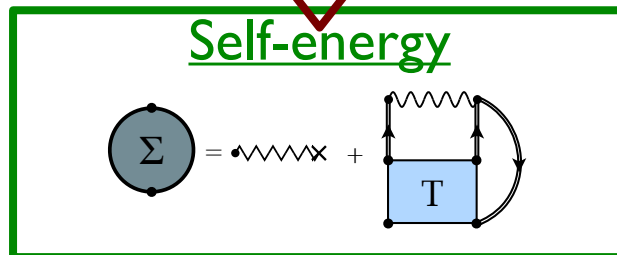
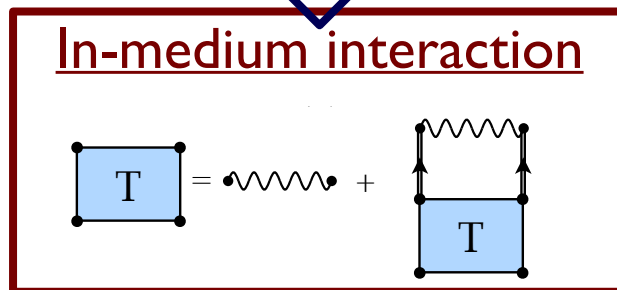
Self-Consistent Green's Functions

(ρ, T)

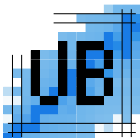


Carbone, Rios & Polls PRC **88** 044302 (2013);
PRC **90**, 054322 (2014);
Carbone PhD Thesis

Forward



Ramos, Polls & Dickhoff, NPA **503** 1 (1989)
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Frick & Muther, PRC **68** 034310 (2003)
Rios, PhD Thesis, U. Barcelona (2007)
Soma & Bozek, PRC **78** 054003 (2008)
Rios & Soma PRL **108** 012501 (2012)



Self-Consistent Green's Functions

(ρ, T)

2N & 3N forces

$$\begin{aligned} \text{wavy line} &= \text{dashed line} + \text{dotted line} \\ \text{wavy line with cross} &= \text{dashed line} + \frac{1}{2} \text{dotted line} \end{aligned}$$

In-medium interaction

$$T = \text{wavy line} + \text{loop diagram with } T \text{ in the middle}$$

Self-energy

$$\Sigma = \text{wavy line with cross} + \text{loop diagram with } T \text{ in the middle}$$

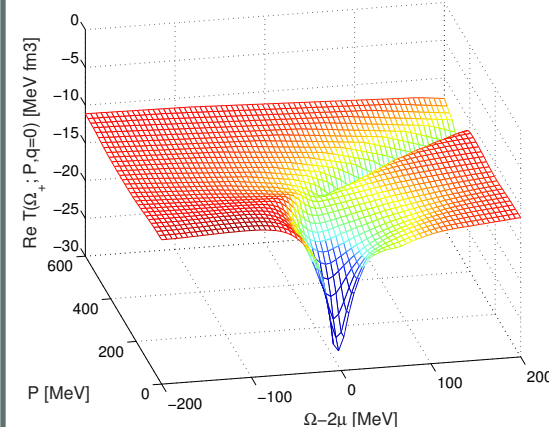
Propagator

$$\text{double line} = \text{single line} + \text{loop diagram with } \Sigma^I \text{ in the middle}$$

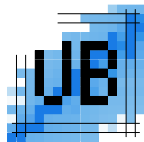
Carbone, Rios & Polls PRC **88** 044302 (2013);
PRC **90**, 054322 (2014);
Carbone PhD Thesis

Forward

T-matrix at $T=5$ MeV



Ramos, Polls & Dickhoff, NPA **503** 1 (1989)
Alm *et al.*, PRC **53** 2181 (1996)
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Self-Consistent Green's Functions

(ρ, T)

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In-medium interaction

$$T = \text{wavy line} + \text{wavy line} \text{ with } T \text{ box}$$

Self-energy

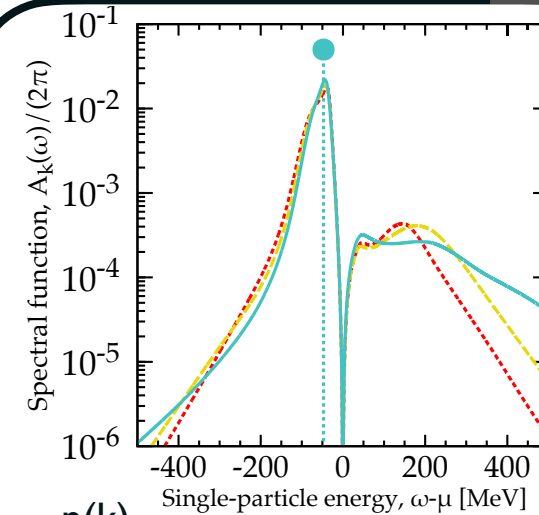
$$\Sigma = \text{wavy line with cross} + \text{wavy line with } T \text{ box}$$

Propagator

$$\text{double line} = \text{single line} + \text{single line with } \Sigma^I \text{ box}$$

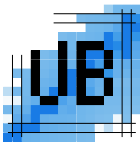
Carbone, Rios & Polls PRC **88** 044302 (2013);
PRC **90**, 054322 (2014);
Carbone PhD Thesis

Forward



$n(k)$
Thermodynamics & EoS
Transport

Ramos, Polls & Dickhoff, NPA **503** 1 (1989)
Alm et al., PRC **53** 2181 (1996)
Dewulf et al., PRL **90** 152501 (2003)
Frick & Muther, PRC **68** 034310 (2003)
Rios, PhD Thesis, U. Barcelona (2007)
Soma & Bozek, PRC **78** 054003 (2008)
Rios & Soma PRL **108** 012501 (2012)



Self-Consistent Green's Functions

(ρ, T)

2N & 3N forces

$$\begin{aligned} \text{wavy line} &= \text{dashed line} + \text{dotted line} \\ \text{wavy line with cross} &= \text{dashed line} + \frac{1}{2} \text{dotted line} + \text{dotted line} \end{aligned}$$

In-medium interaction

$$T = \text{wavy line} + \text{wavy line} \text{ with } T \text{ box}$$

Self-energy

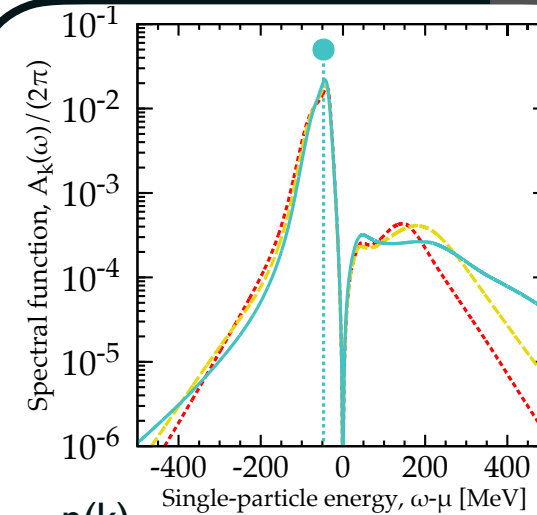
$$\Sigma = \text{wavy line with cross} + \text{wavy line with } T \text{ box}$$

Propagator

$$\text{double line} = \text{single line} + \text{single line with } \Sigma^I \text{ box}$$

Carbone, Rios & Polls PRC **88** 044302 (2013);
PRC **90**, 054322 (2014);
Carbone PhD Thesis

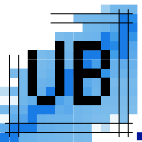
Forward



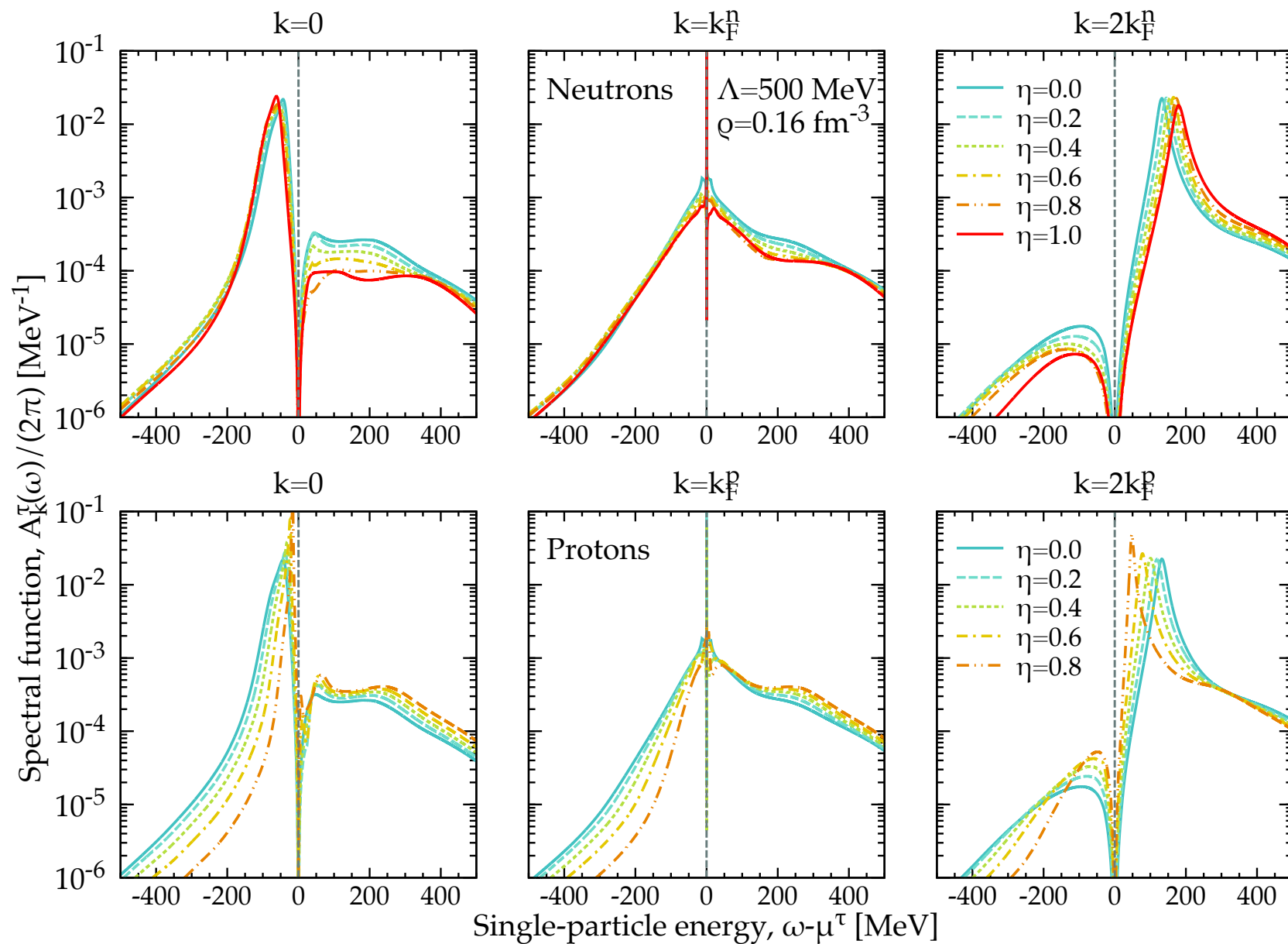
$n(k)$
Thermodynamics & EoS
Transport

- Off-shell ✓
- Matsubara formalism ✓
- Φ -derivable ✓

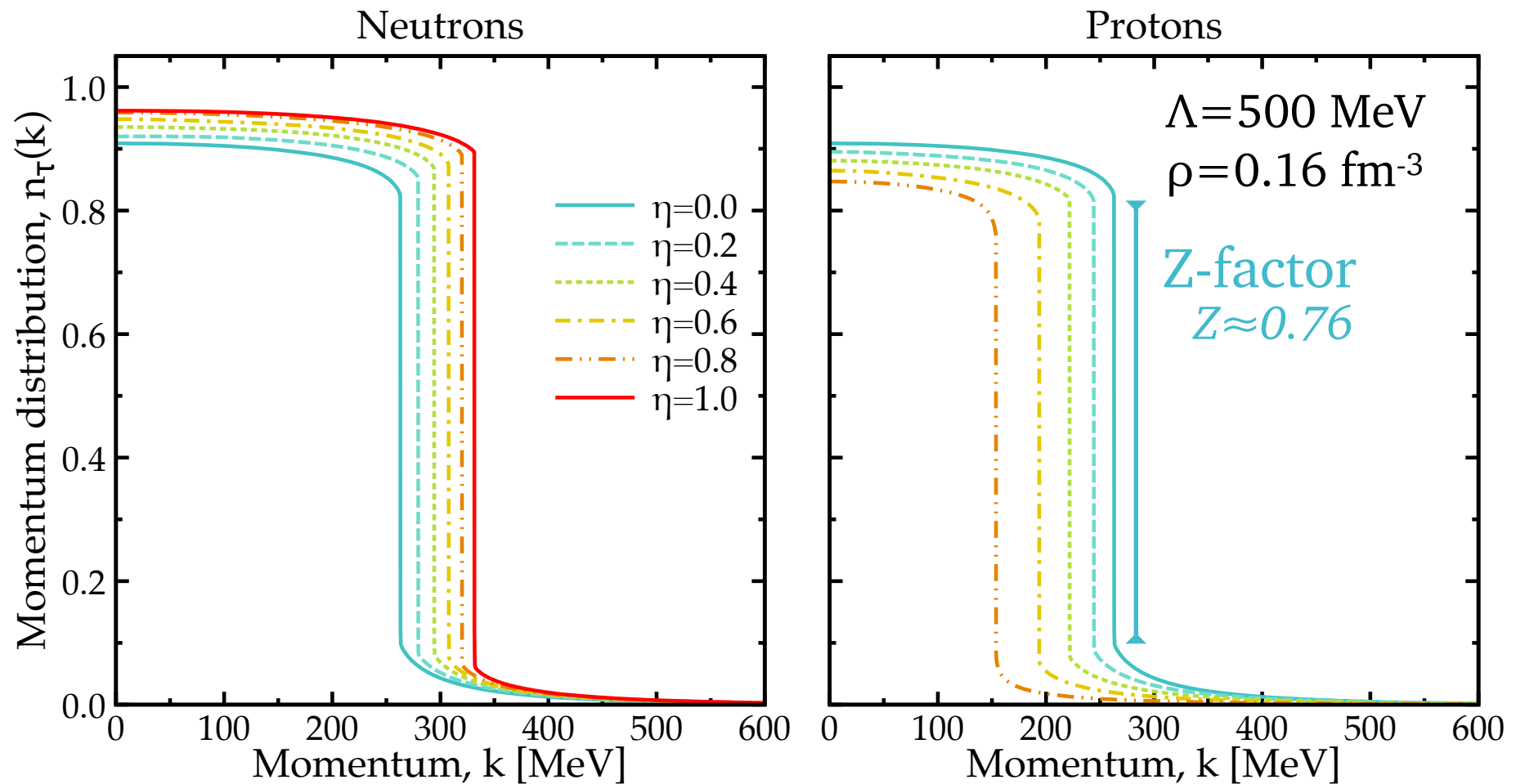
Ramos, Polls & Dickhoff, NPA **503** 1 (1989)
Alm et al., PRC **53** 2181 (1996)
Dewulf et al., PRL **90** 152501 (2003)
Frick & Muther, PRC **68** 034310 (2003)
Rios, PhD Thesis, U. Barcelona (2007)
Soma & Bozek, PRC **78** 054003 (2008)
Rios & Soma PRL **108** 012501 (2012)



Isospin-asymmetric spectral functions

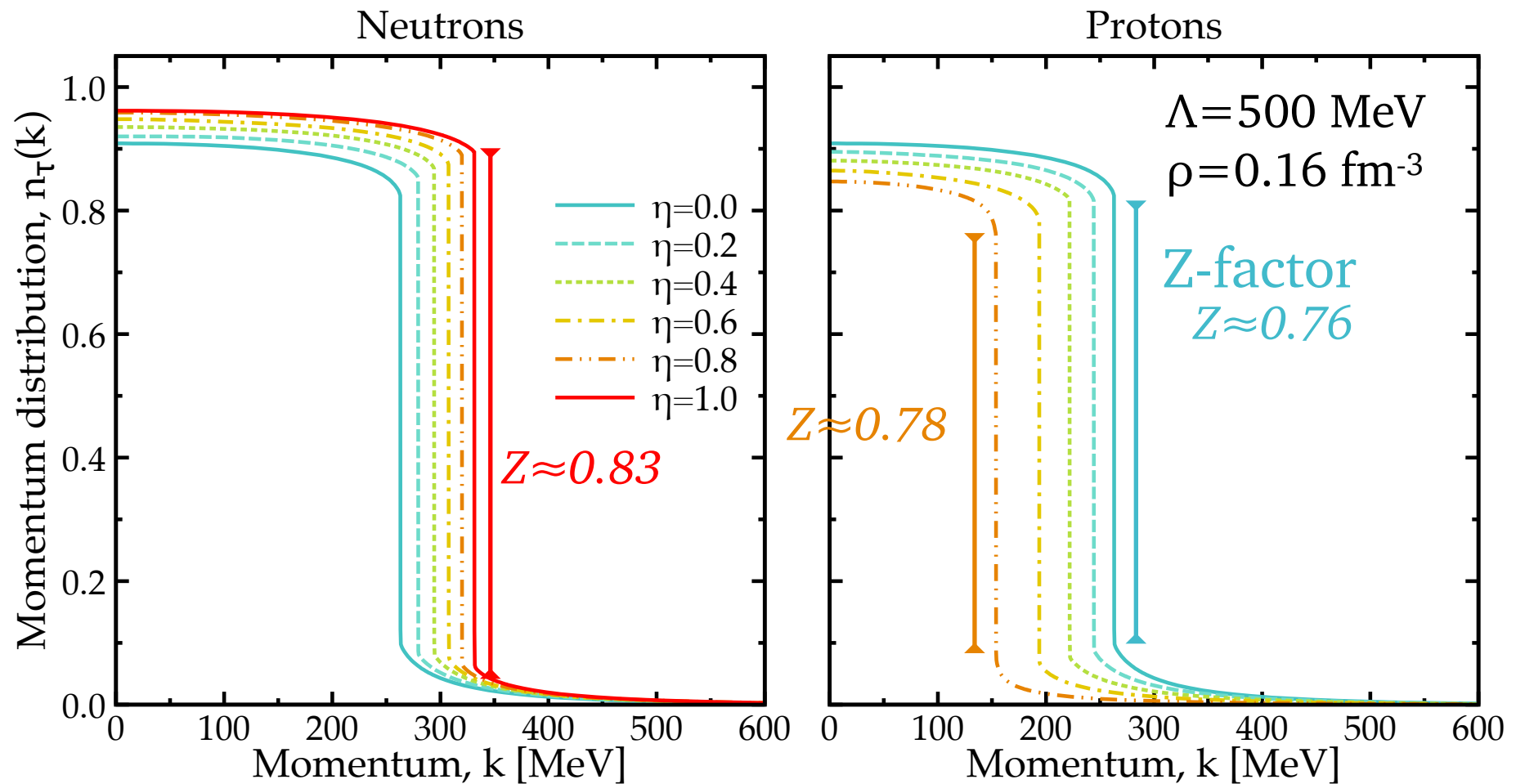


Momentum distribution



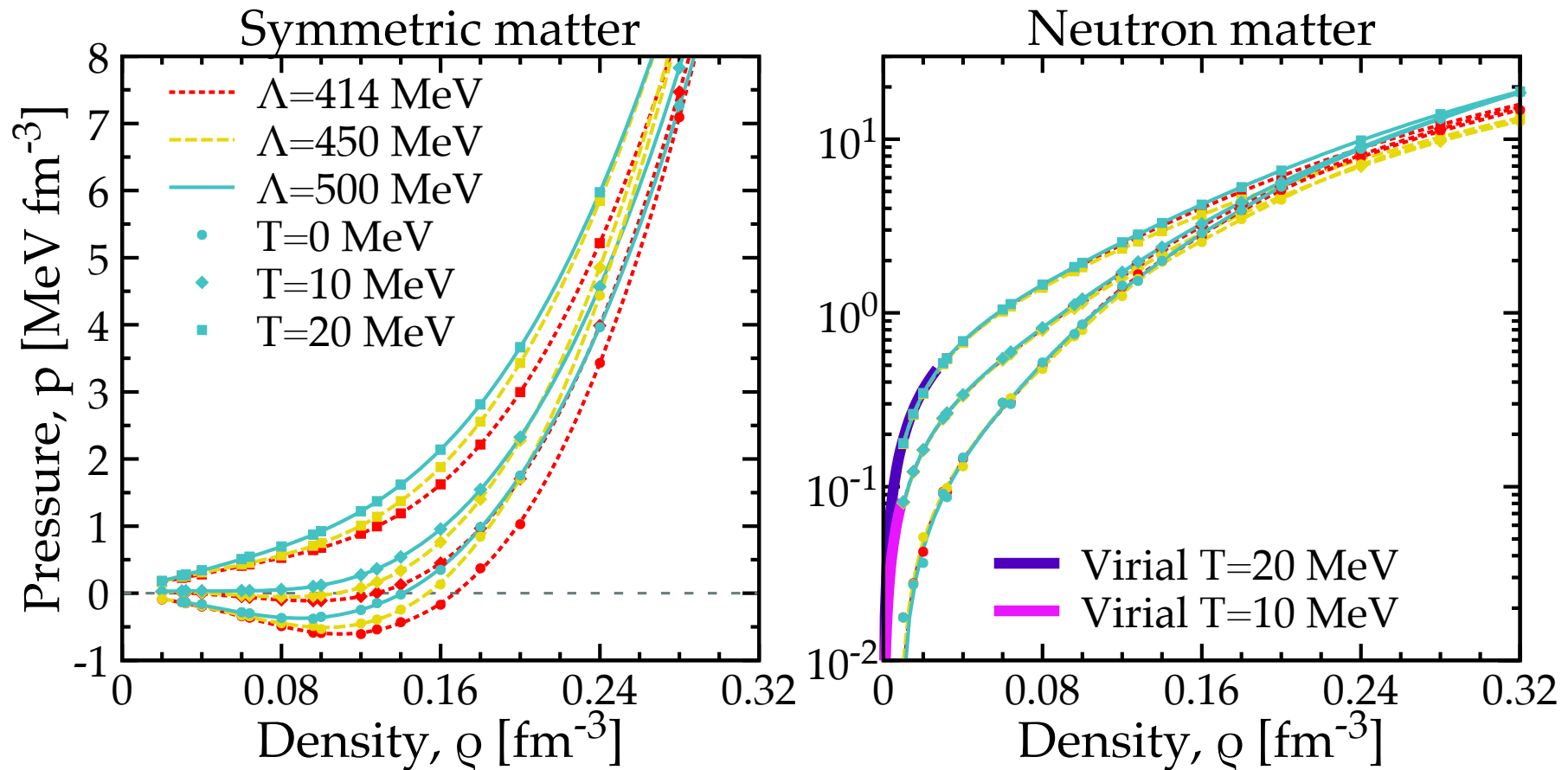
Forward

Momentum distribution



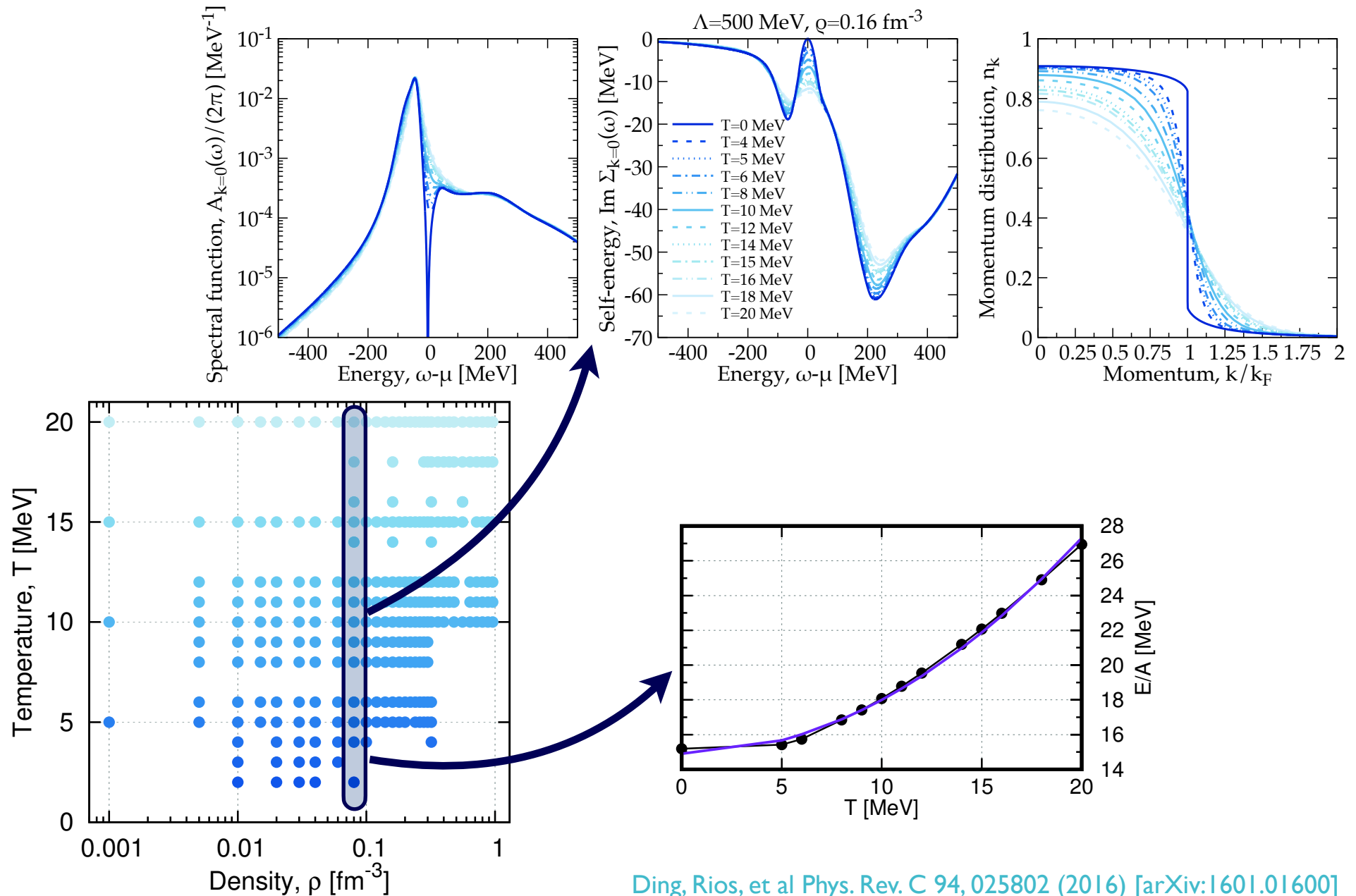
Forward

EoS at finite temperature

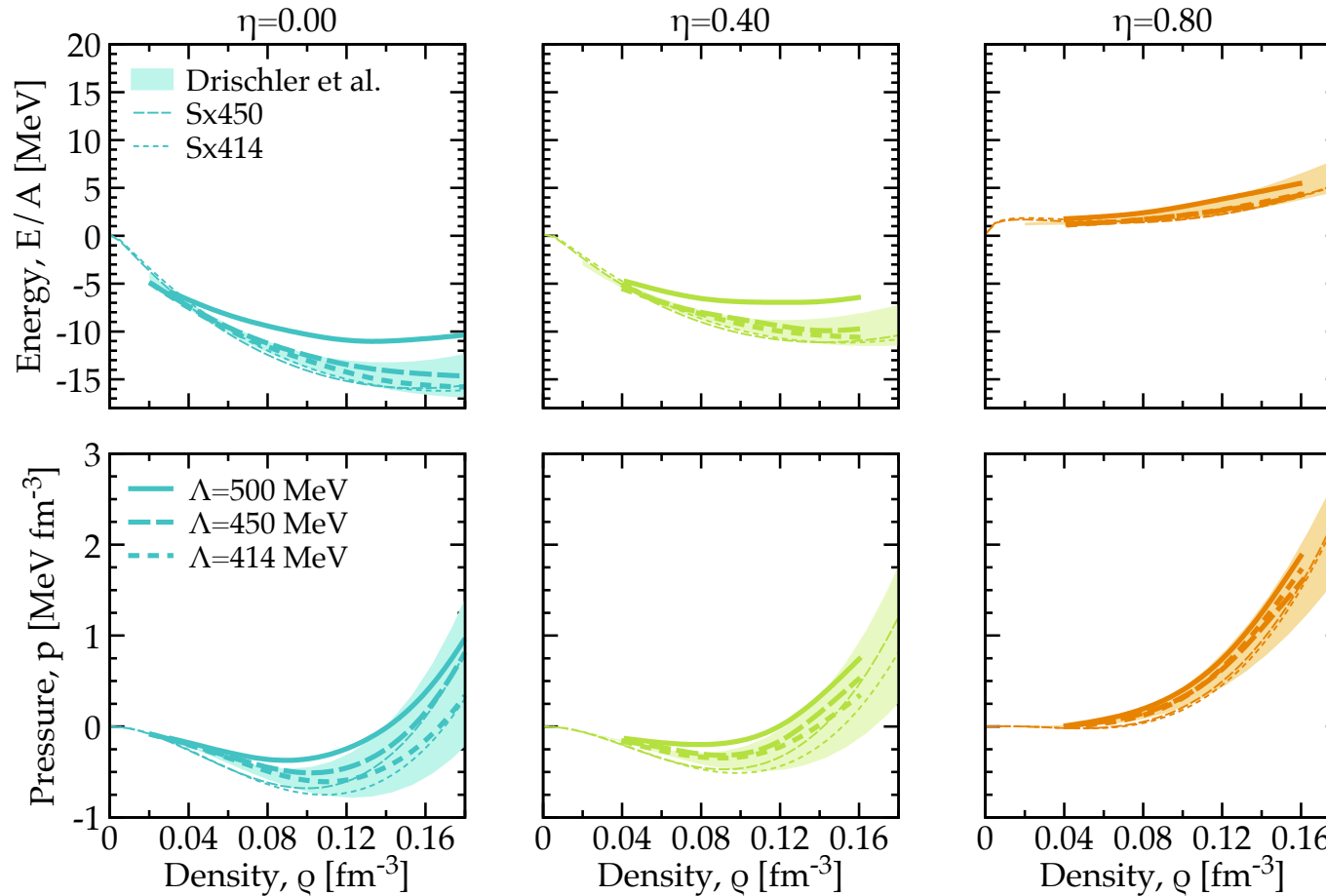


- Relevant & **necessary** for binary NS simulations
- **Parameter-free** first principles calculation
- Reproduces **virial** at low density

Zero temperature extrapolation



Self-Consistent Green's Functions



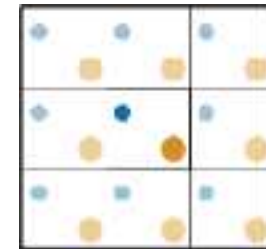
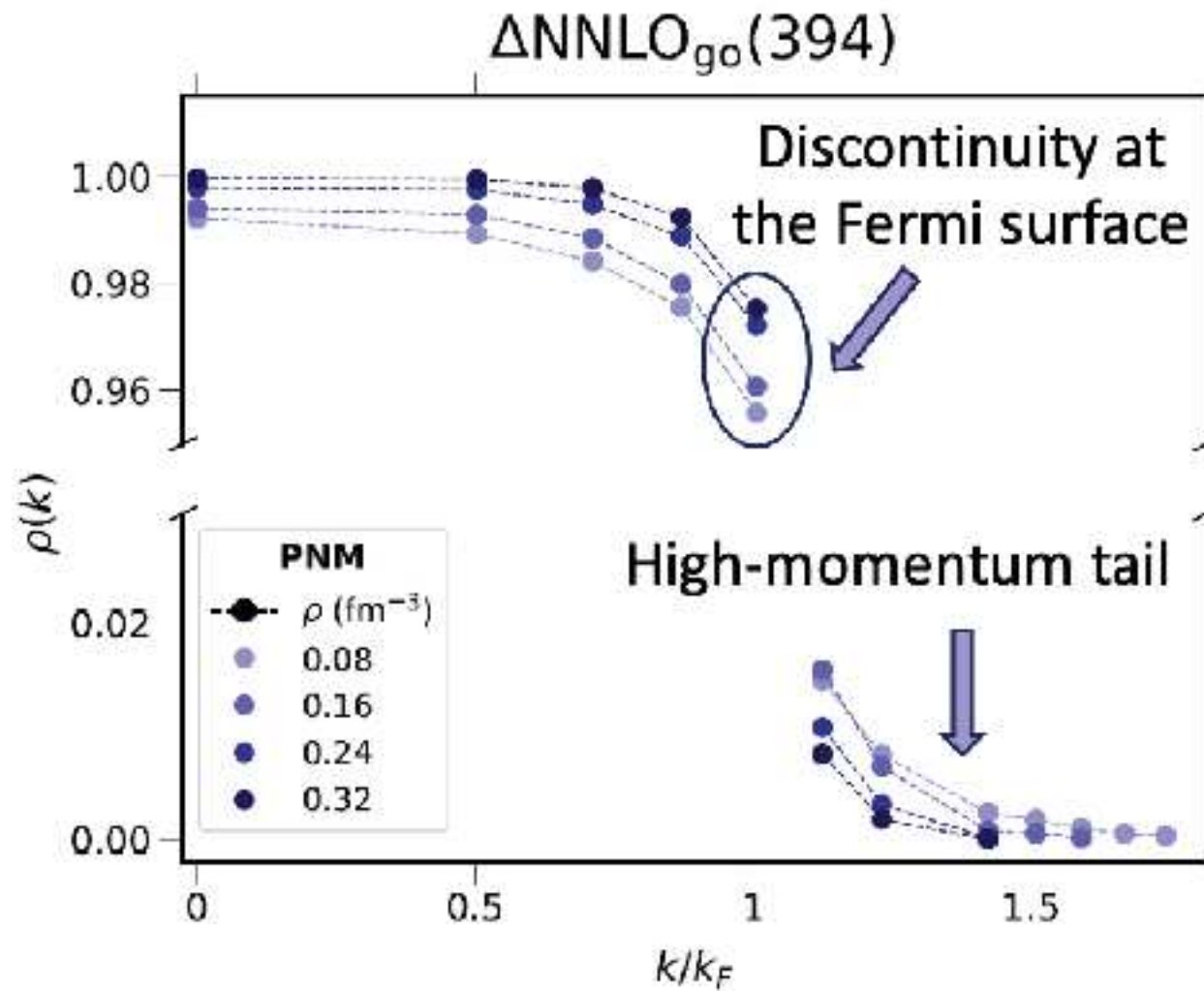
SCGF can treat:

- Explicitly asymmetric matter ✓
- Finite temperature ✓
- Systematic expansion ✓
- 3 nucleon forces ✓

[Rios, Frontiers Physics fphy.2020.00387 \(2020\)](#)
[\[arXiv:2006.10610\]](#)

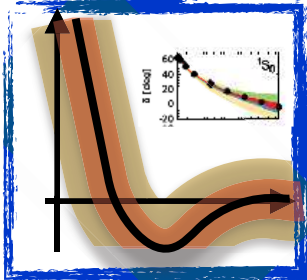
but:

- Numerically intensive
- Pairing?

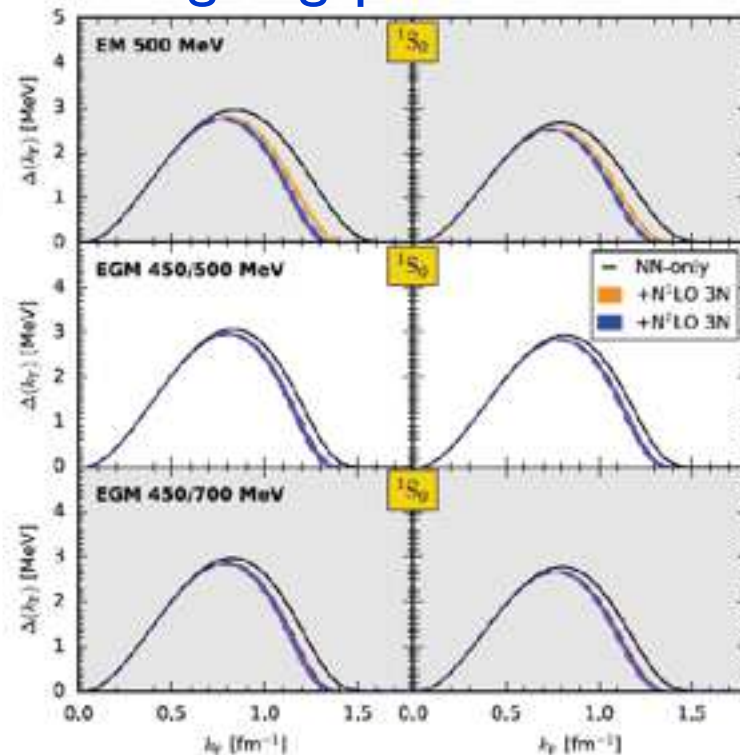


BCS+HF gaps in neutron matter

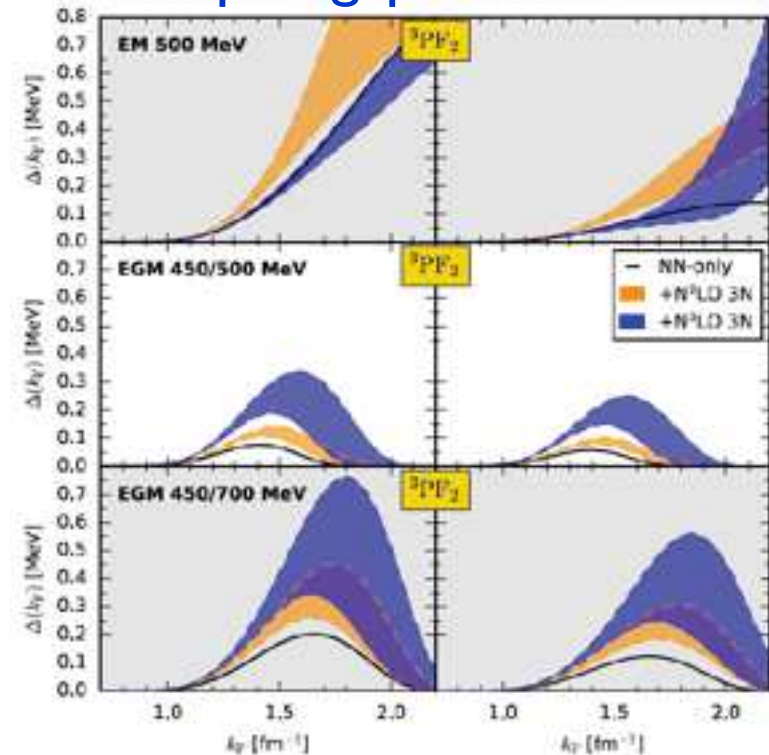
Hamiltonian



Singlet gaps with 3NF



Triplet gaps with 3NF



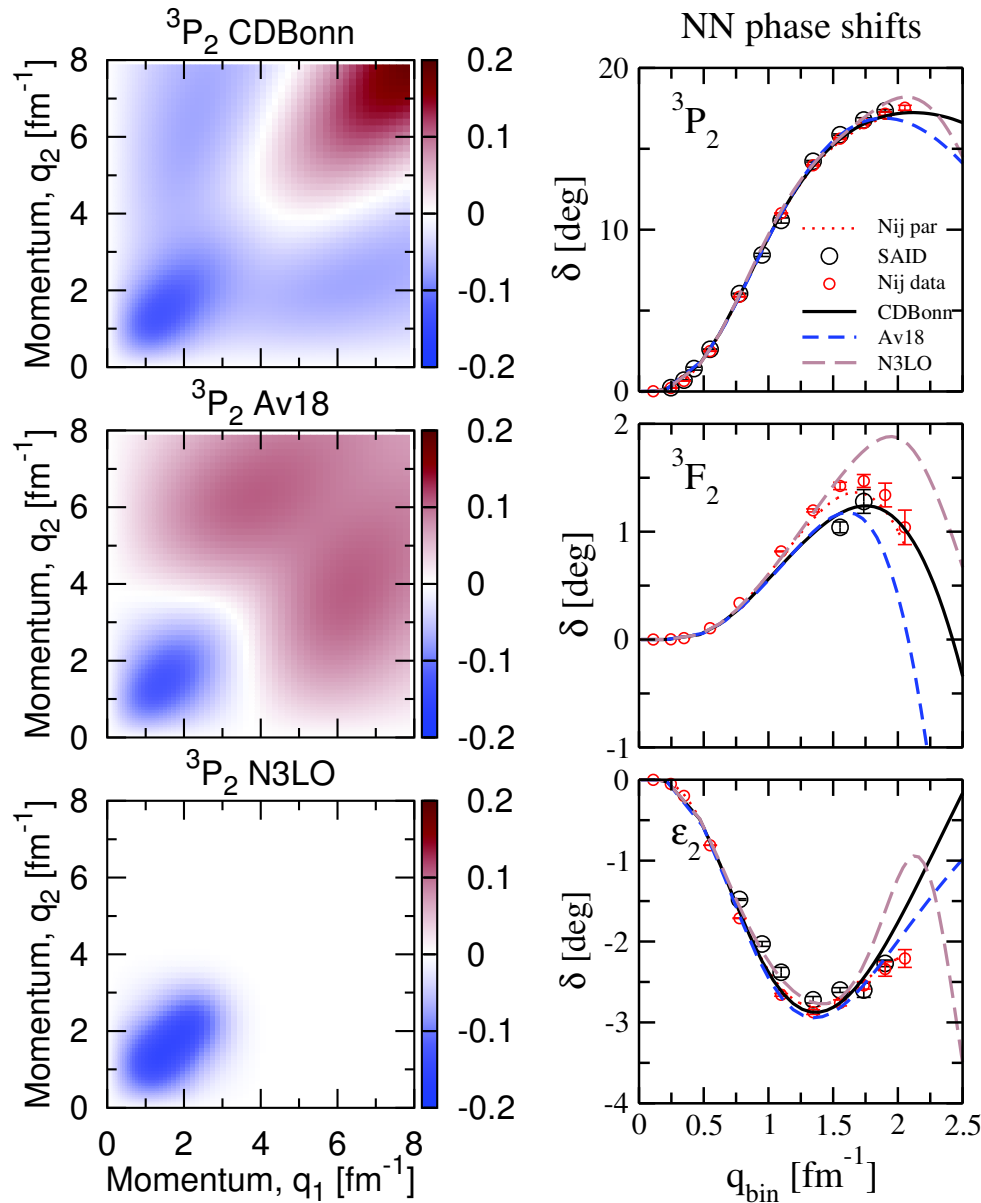
Drischler, Kruger, Hebeler, Schwenk, *PRC* 95 024302 (2017) [arXiv:1610.05213]

BCS

$$\Delta_k^L = - \sum_{L'} \int_{k'} \frac{\langle k | V_{nn}^{LL'} | k' \rangle}{2\sqrt{\chi_{k'}^2 + |\Delta_{k'}|}} \Delta_{k'}^{L'} + \quad \begin{aligned} \chi_k &= \varepsilon_k - \mu \\ \varepsilon_k &= \frac{k^2}{2m} + U(k) - \mu \end{aligned}$$

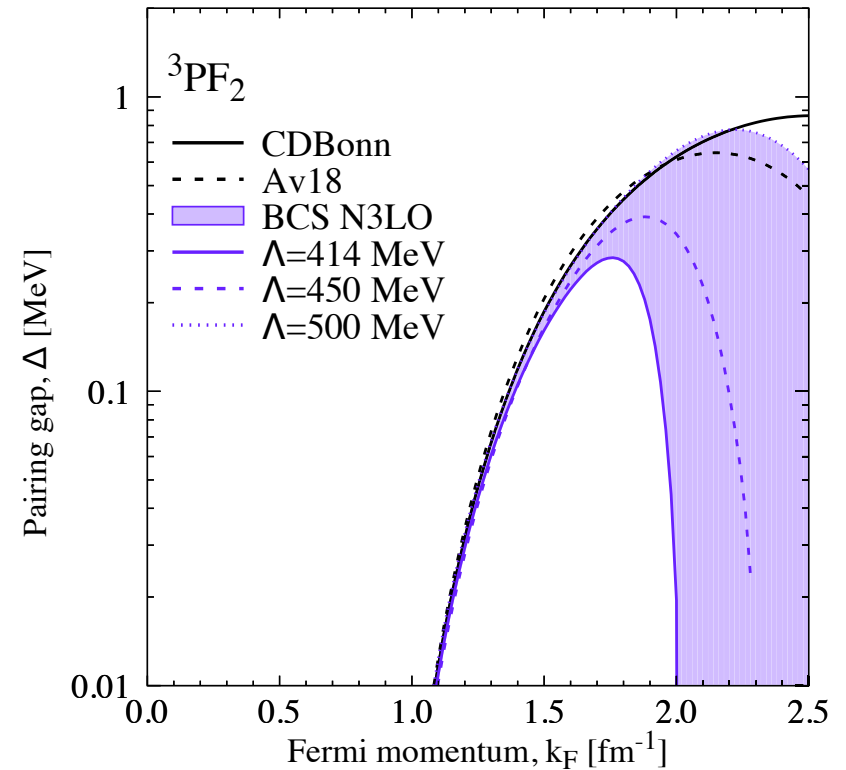
- **Error estimates** from nuclear force (chiral expansion) ✓
- **Many-body uncertainty?** ✗

Triplet channel: limits of EFT



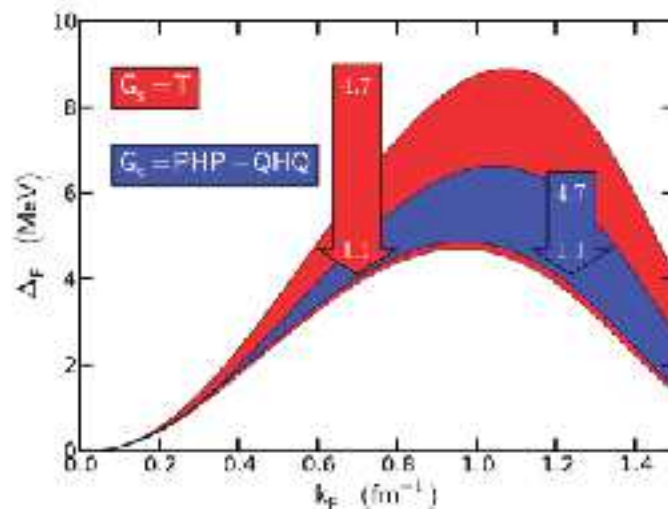
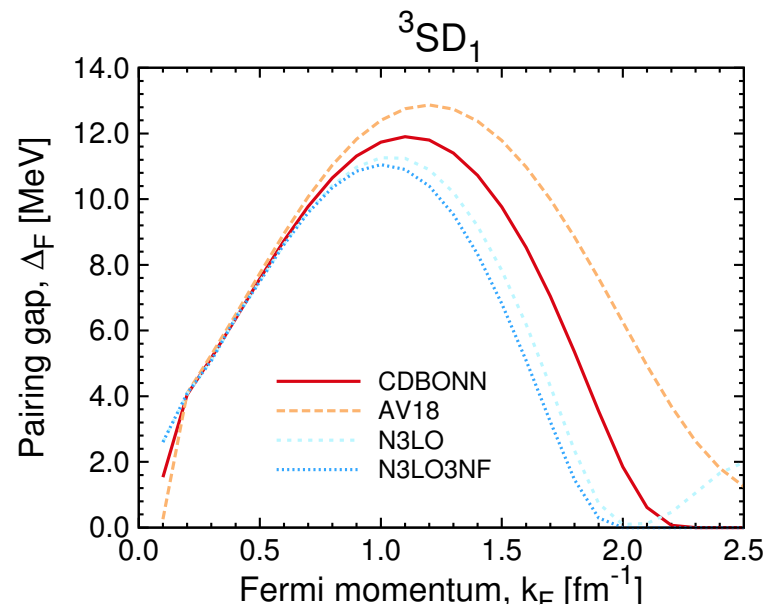
BCS equation

$$\Delta_k^L = - \sum_{L'} \int_{k'} \frac{\langle k | V_{nn}^{LL'} | k' \rangle}{2\sqrt{\chi_{k'}^2 + |\Delta_{k'}|^2}} \Delta_{k'}^{L'}$$



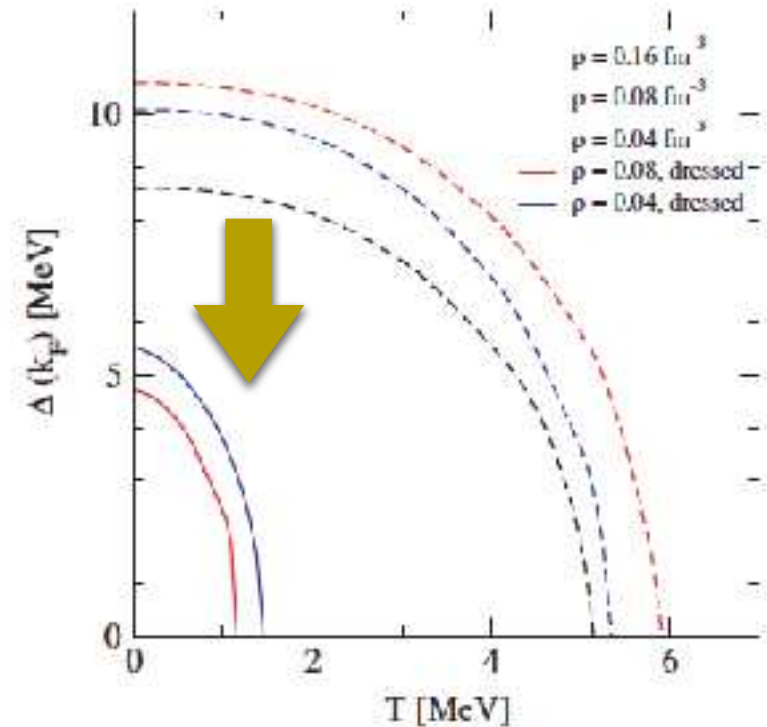
Triplet pairing: symmetric matter

$^3\text{SD}_1$ nuclear matter BCS gaps



Maurizio, Holt & Finelli, PRC **90**, 044003 (2014)

SRC-depleted $^3\text{SD}_1$ gaps

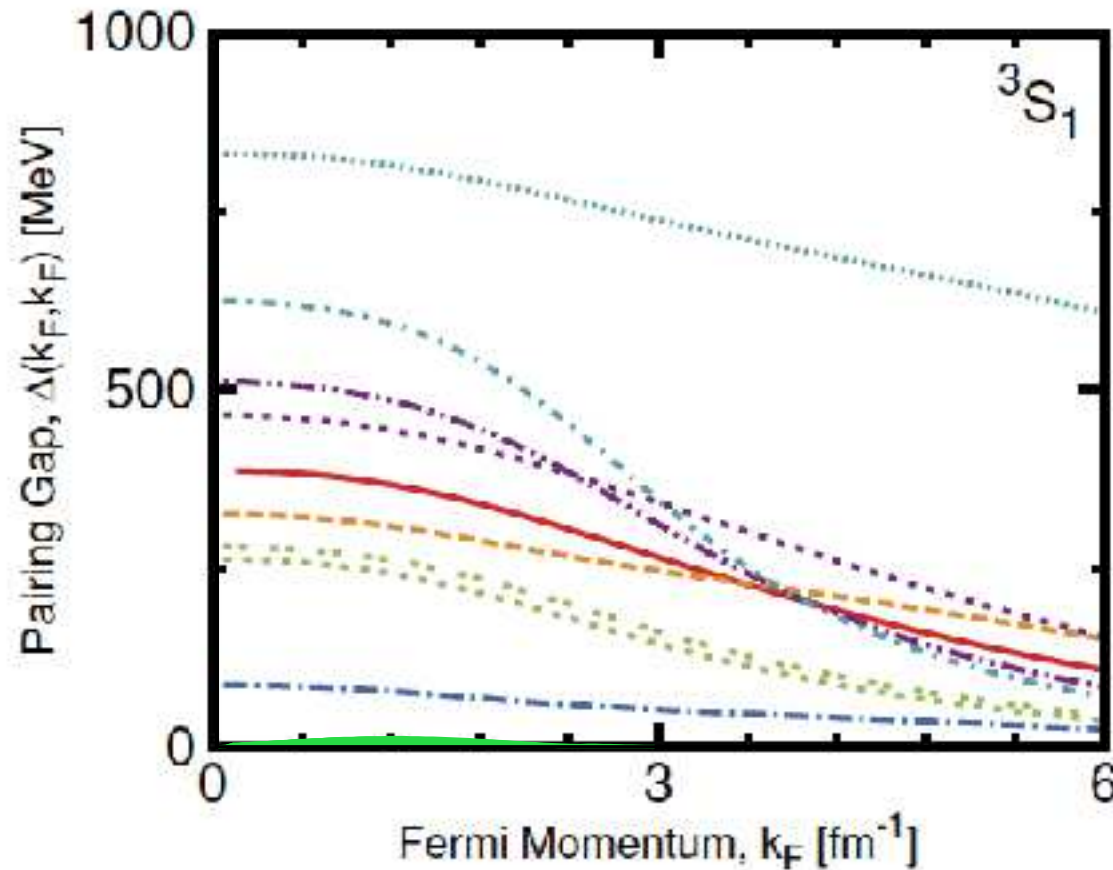


Muether & Dickhoff, PRC **72** 054313 (2005)

- **Massive** gaps $^3\text{SD}_1$ channel but...
- No **evidence** of strong np nuclear pairing
- Short-range correlations **deplete** gap
- 3BF effect? Short-range effects? Deformation?

Gezerli's talk

Triplet pairing: Gogny?



— D1 - - - D250 ... D1N
 - - - D1S - - - D260 ... NR71
 ... D1P - - - D280 - - - D1AS
 - . . D1M . . . D300

S=1

$$V_{L=0}^S \approx t_0^i \rho^{\alpha_i} \left[1 - (-)^S x_0^i \right]$$

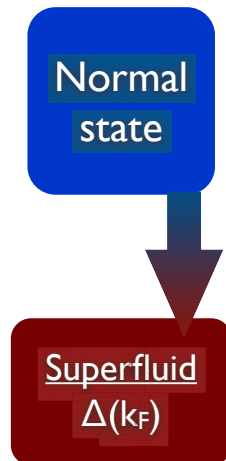
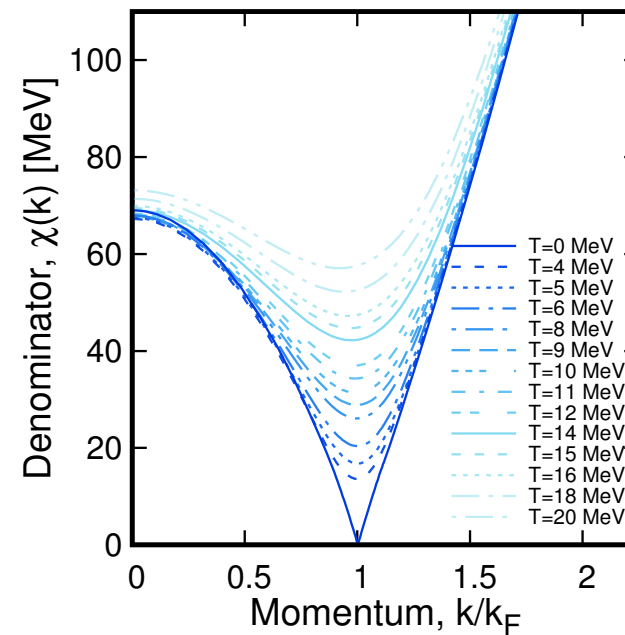
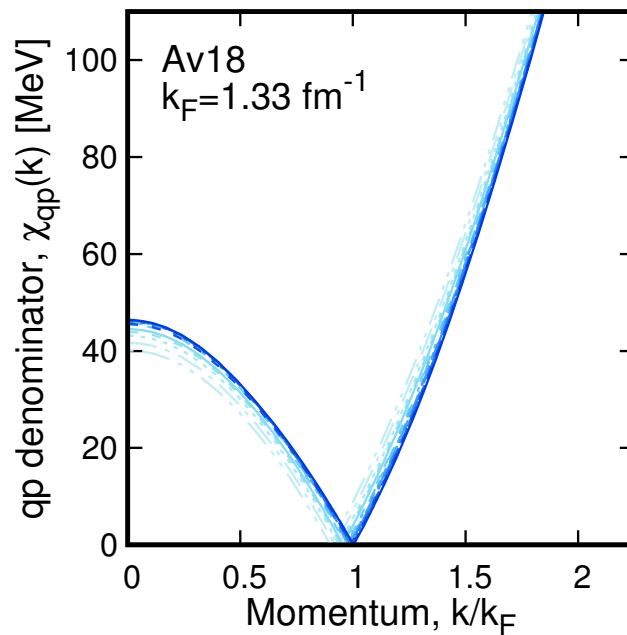
a) we neglect it

b) will be there in nuclei!

BCS equation

$$\Delta_k^L = - \sum_{L'} \int_{k'} \frac{\langle k | V_{nn}^{LL'} | k' \rangle}{2\sqrt{\chi_{k'}^2 + |\Delta_{k'}|^2}} \Delta_{k'}^{L'} + \quad \chi_k = \varepsilon_k - \mu$$

Beyond-BCS: SRCs



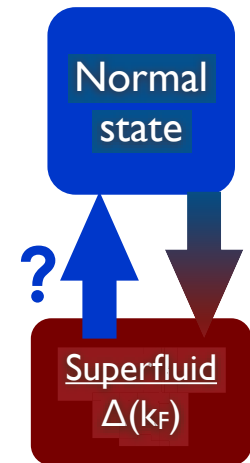
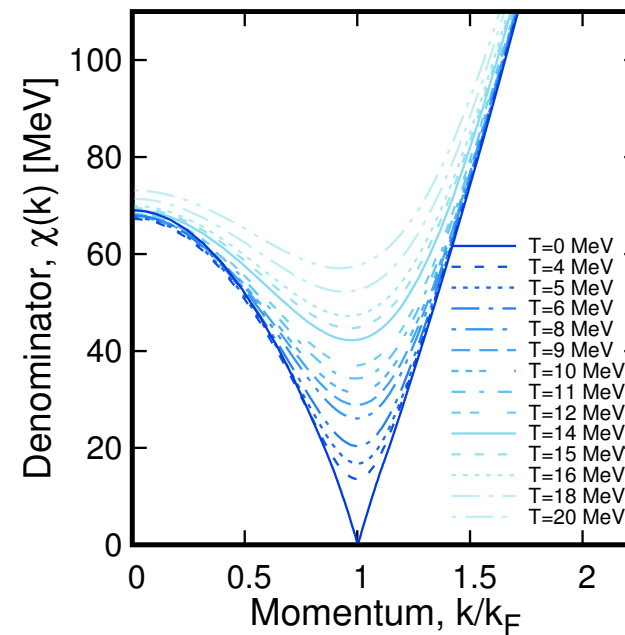
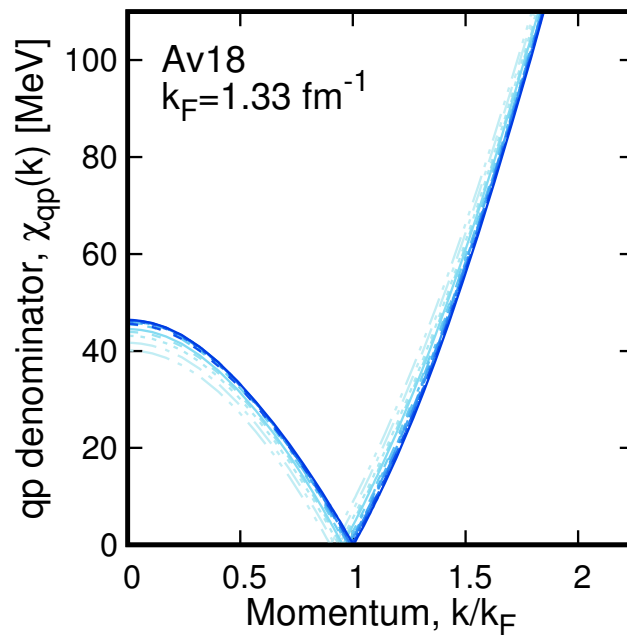
BCS gap equation

$$\Delta_k^L = - \sum_{L'} \int_{k'} \frac{\langle k | V_{nn}^{LL'} | k' \rangle}{2\sqrt{\chi_{k'}^2 + |\Delta_{k'}|^2}} \Delta_{k'}^{L'} + \frac{1}{2\bar{\chi}_k} = \frac{1}{2|\varepsilon_k - \mu|}$$

BCS+SRC gap equation

$$\Delta_k^L = - \sum_{L'} \int_{k'} \frac{\langle k | V_{nn}^{LL'} | k' \rangle}{2\sqrt{\bar{\chi}_{k'}^2 + |\Delta_{k'}|^2}} \Delta_{k'}^{L'} + \frac{1}{2\bar{\chi}_k} = \int_{\omega} \int_{\omega'} \frac{1 - f(\omega) - f(\omega')}{\omega + \omega'} A(k, \omega) A_s(k, \omega')$$

Beyond-BCS: SRCs



BCS gap equation

$$\Delta_k^L = - \sum_{L'} \int_{k'} \frac{\langle k | V_{nn}^{LL'} | k' \rangle}{2\sqrt{\chi_{k'}^2 + |\Delta_{k'}|^2}} \Delta_{k'}^{L'} + \frac{1}{2\bar{\chi}_k} = \frac{1}{2|\varepsilon_k - \mu|}$$

BCS+SRC gap equation

$$\Delta_k^L = - \sum_{L'} \int_{k'} \frac{\langle k | V_{nn}^{LL'} | k' \rangle}{2\sqrt{\bar{\chi}_{k'}^2 + |\Delta_{k'}|^2}} \Delta_{k'}^{L'} + \frac{1}{2\bar{\chi}_k} = \int_{\omega} \int_{\omega'} \frac{1 - f(\omega) - f(\omega')}{\omega + \omega'} A(k, \omega) A_s(k, \omega')$$

Bozek, PRC **62** 054316 (2000)

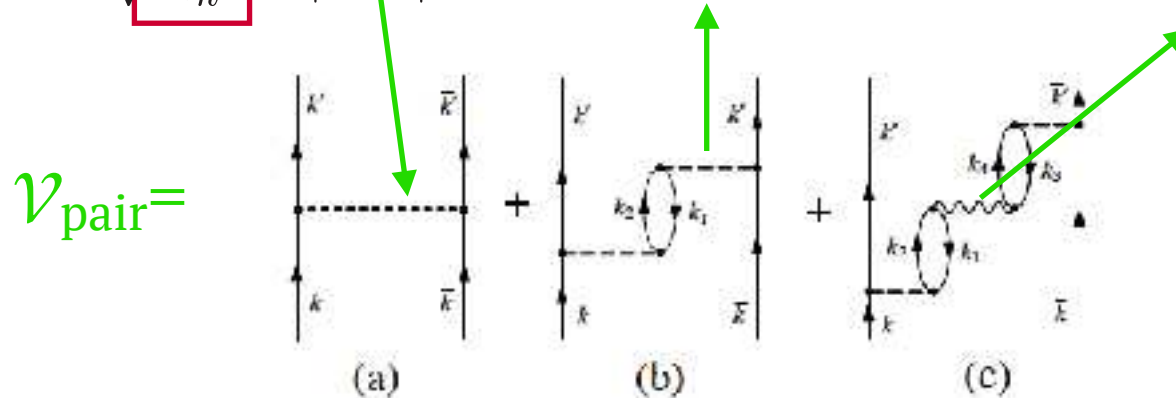
Muther & Dickhoff, PRC **72** 054313 (2005)

Rios, Polls & Dickhoff, J Low T Phys **189** 234 (2017)

$$\Delta_k^L = - \sum_{L'} \int_{k'} \frac{\langle k | V_{nn}^{LL'} | k' \rangle}{2\sqrt{\chi_{k'}^2 + |\Delta_{k'}|^2}} \Delta_{k'}^{L'}$$

ph recoupled G-matrix

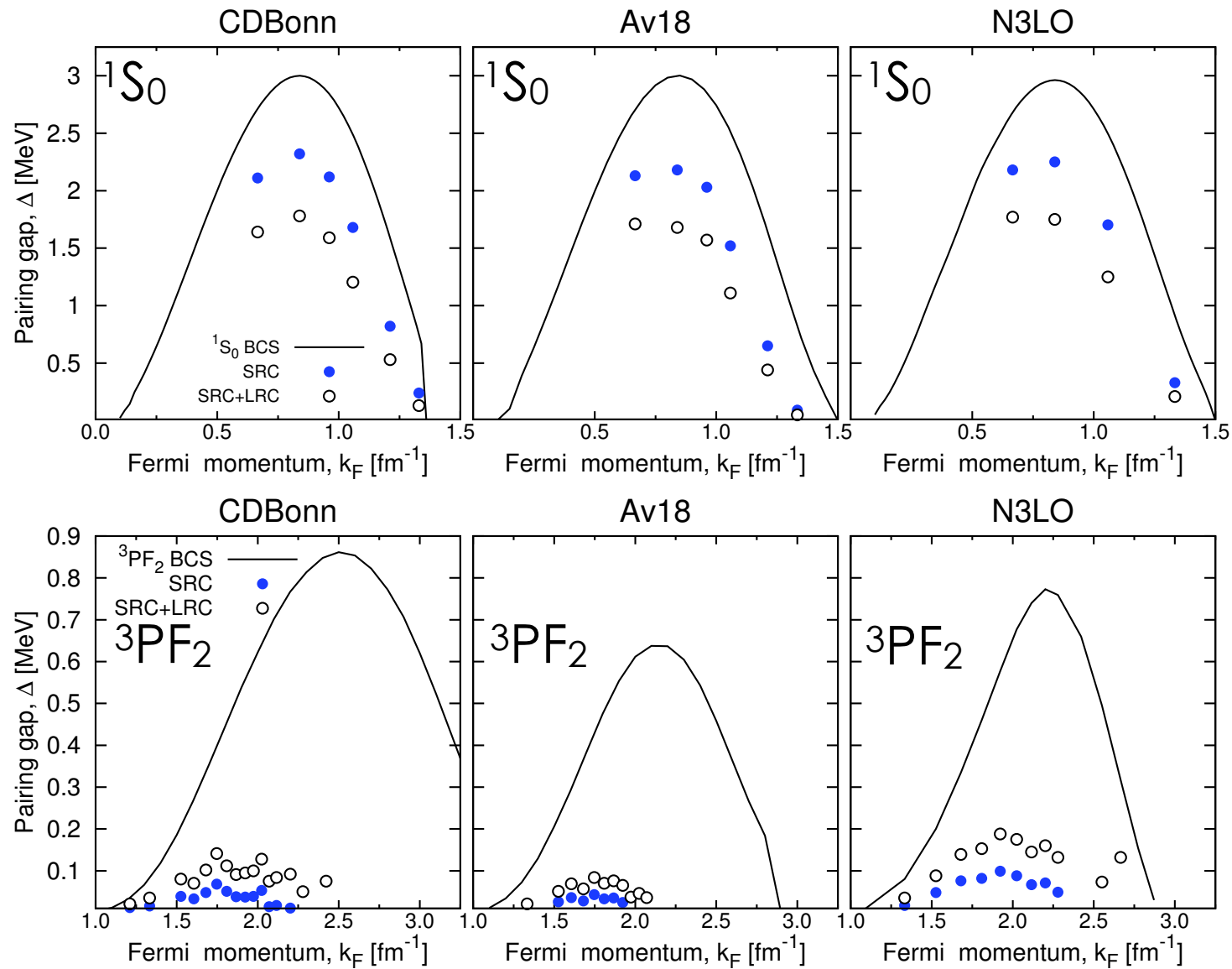
Effective Landau parameters

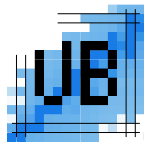


$$\langle 1\bar{1} | \mathcal{V} | 1\bar{1} \rangle = \frac{1}{4} \sum_{2,2'} \sum_{S,T} (-)^S (2S+1) \langle 12 | G_{ST}^{\text{ph}} | 1'2' \rangle_A \langle 2'\bar{1} | G_{ST}^{\text{ph}} | 2\bar{1}' \rangle_A \Lambda(22')$$

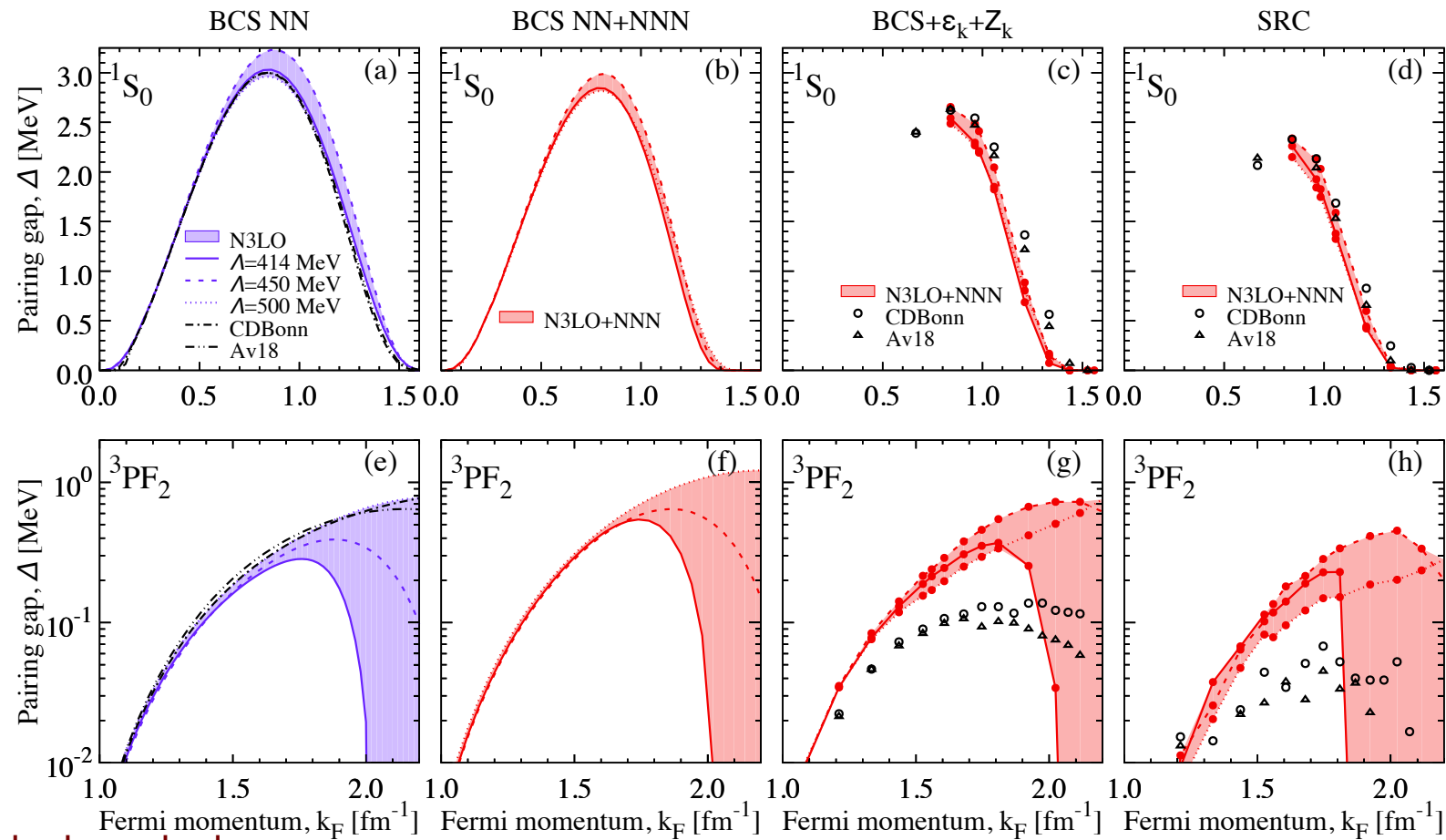
$$\Lambda_{ST}(q) = \frac{\Lambda_{ST}^0(q)}{1 - \Lambda_{ST}^0(q) \times F_{ST}}$$

- Bare NN potential only is **not** the only possible interaction
- Diagram (a): nuclear interaction
- Diagram (b): in-medium interaction, density and spin fluctuations
- Diagram (c): included by Landau parameters

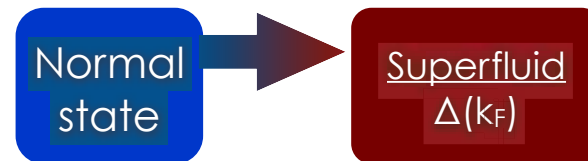
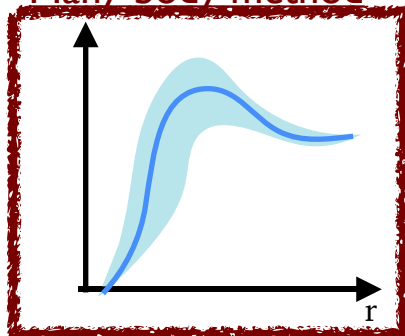


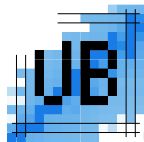


Beyond-BCS in neutron matter: SRC

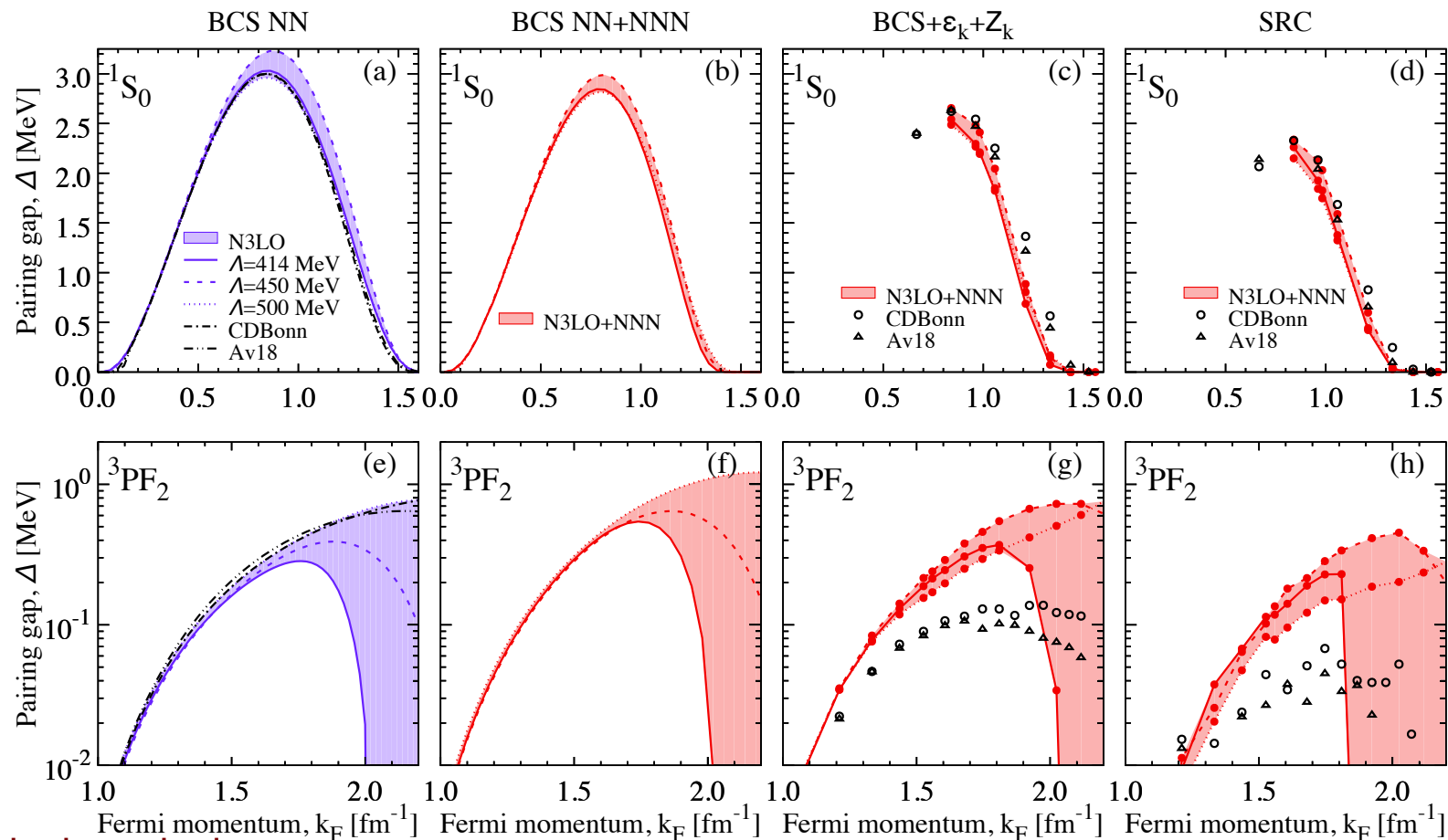


Many-body method

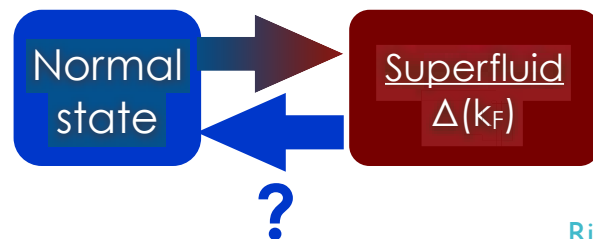
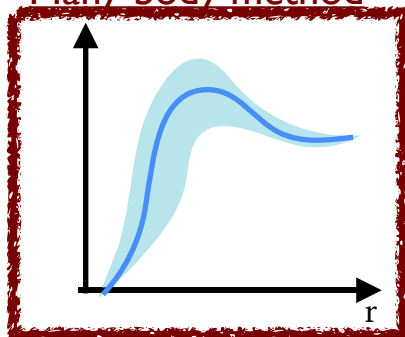




Beyond-BCS in neutron matter: SRC



Many-body method



Rios, Dickhoff, Polls, JLT **189**, 234 (2017) [arXiv:1707.04140]
Rios, Dickhoff, Polls, et al PRC **94**, 025802 (2016) [arxiv:1601.01600]



$$\bullet \text{---} = \cancel{\bullet \text{---}} = \bullet \text{---} \circlearrowright + \frac{1}{2} \bullet \text{---} \circlearrowright \text{---} \circlearrowright$$

Effective two-body force \Rightarrow **NN forces**

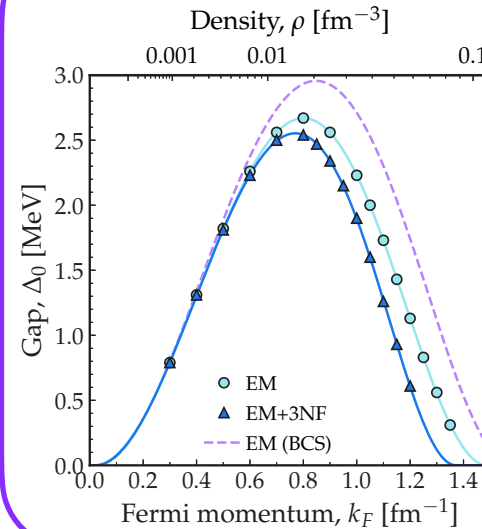
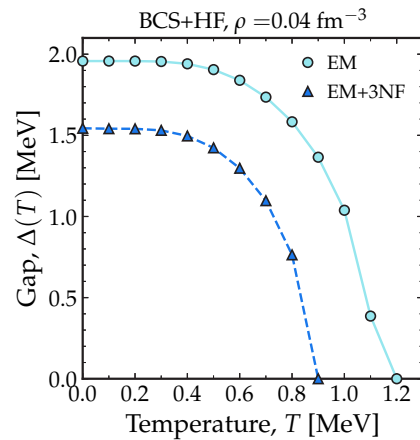
Diagrammatic equation for the propagator:

$$\text{Blue wavy line} = \text{Blue dashed line} + \text{Blue dotted line} \times \text{Black loop}$$

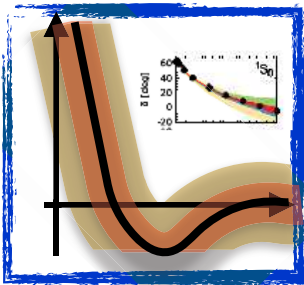
- Singlet gap: 3NF **reduce** closure
- Triplet gap: 3NF **increase** gap
- **Model dependence** to be explored
- Use **SRG** for systematics (Michael!)

Finite Temperature BCS

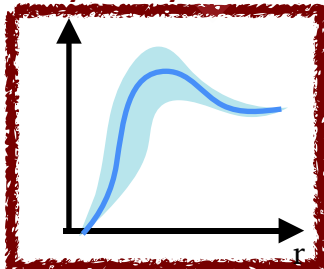
BCS+HF



Hamiltonian

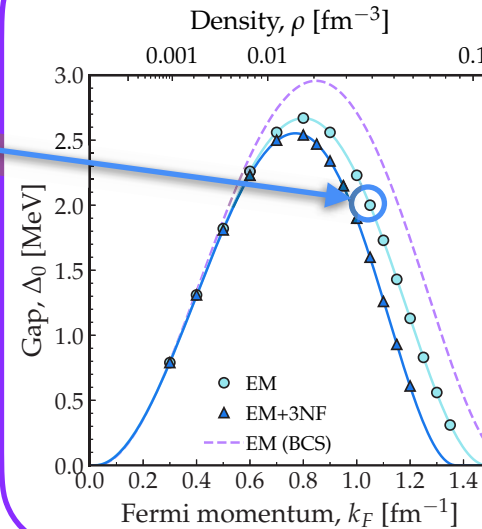
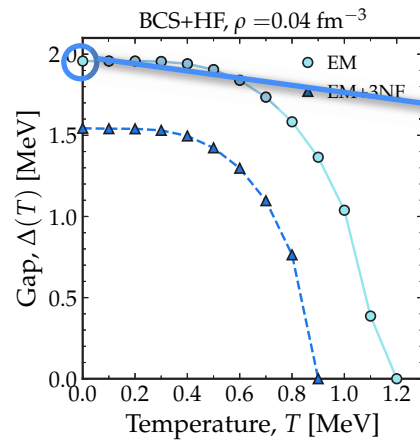


Many-body method

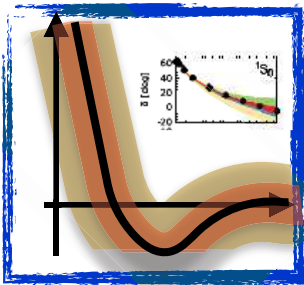


Finite Temperature BCS

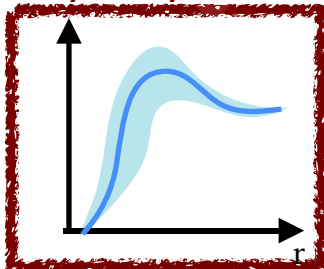
BCS+HF



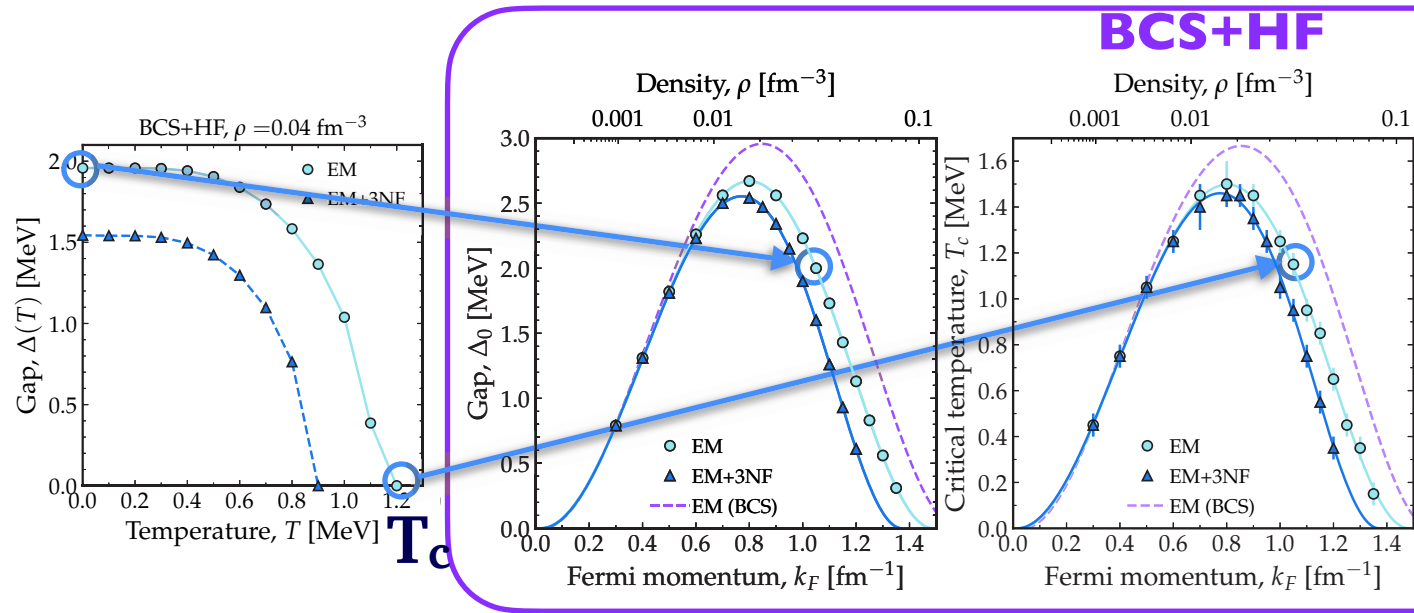
Hamiltonian



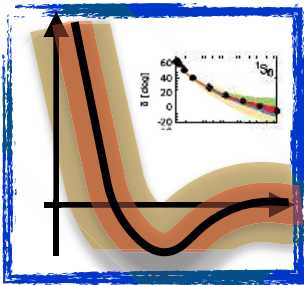
Many-body method



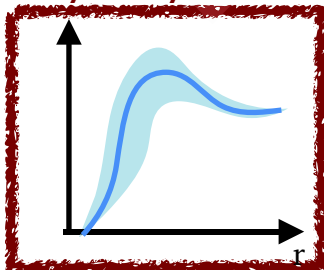
Finite Temperature BCS



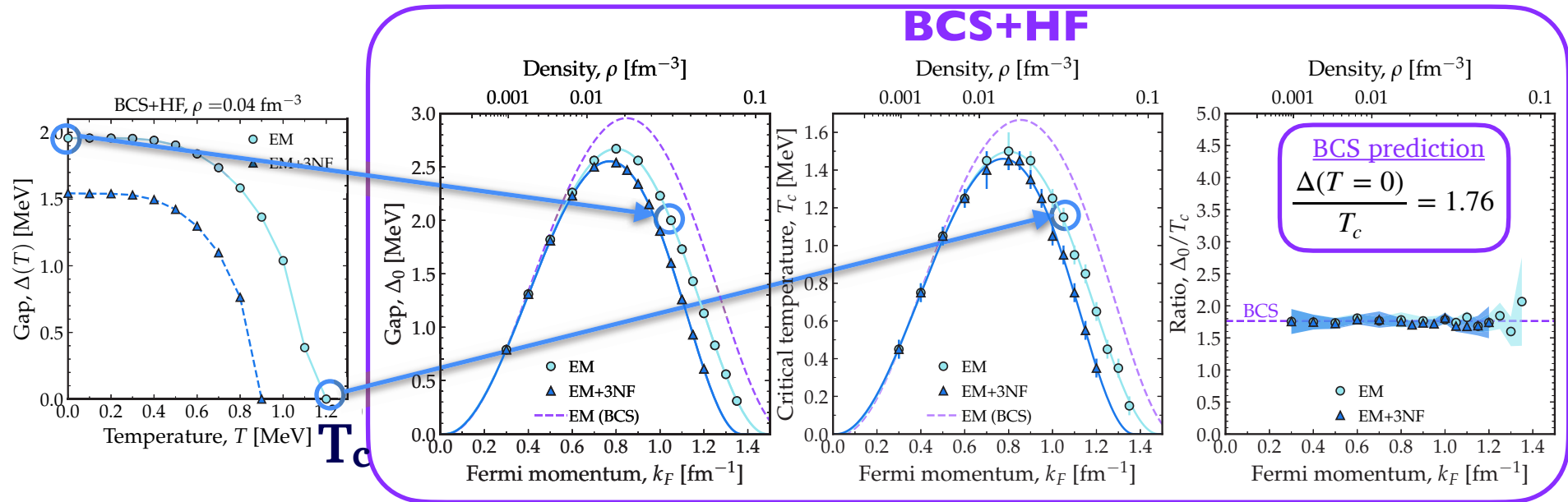
Hamiltonian



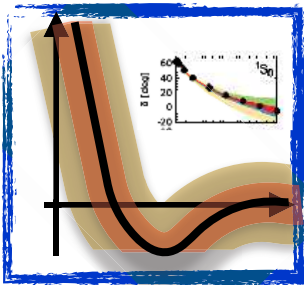
Many-body method



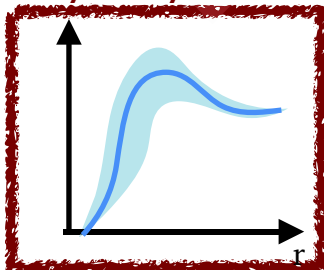
Finite Temperature BCS



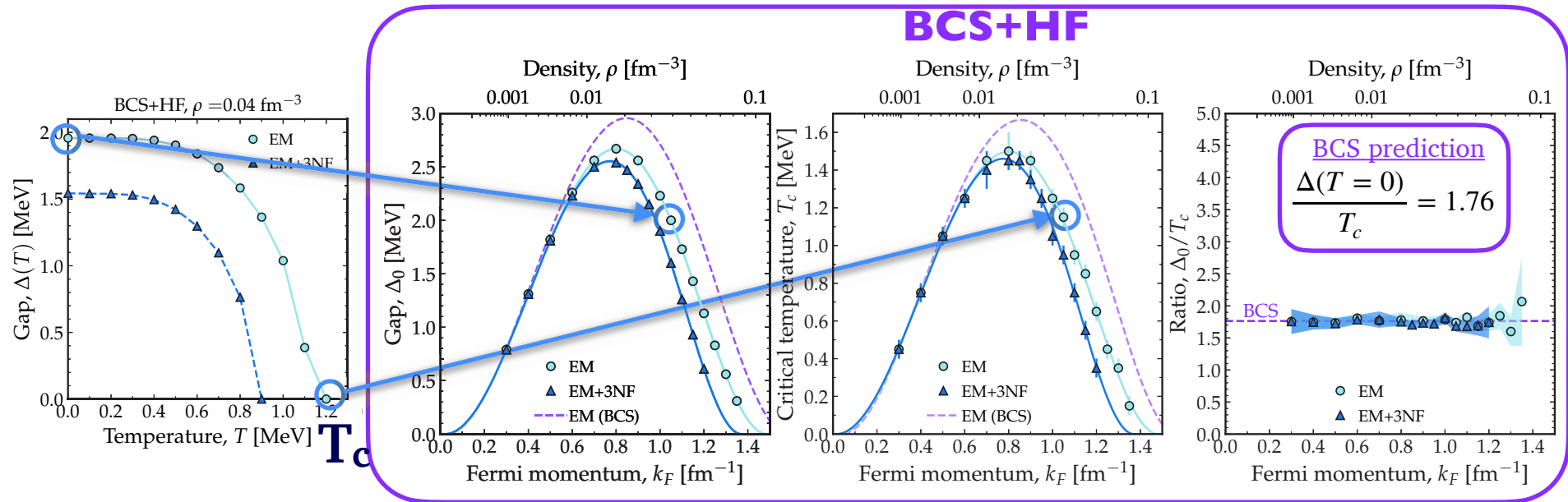
Hamiltonian



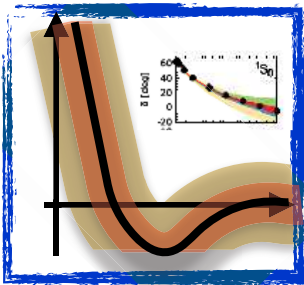
Many-body method



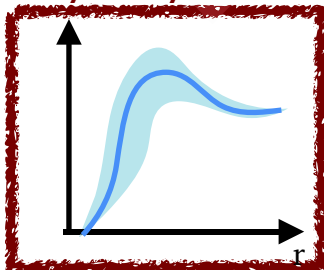
Finite Temperature BCS



Hamiltonian

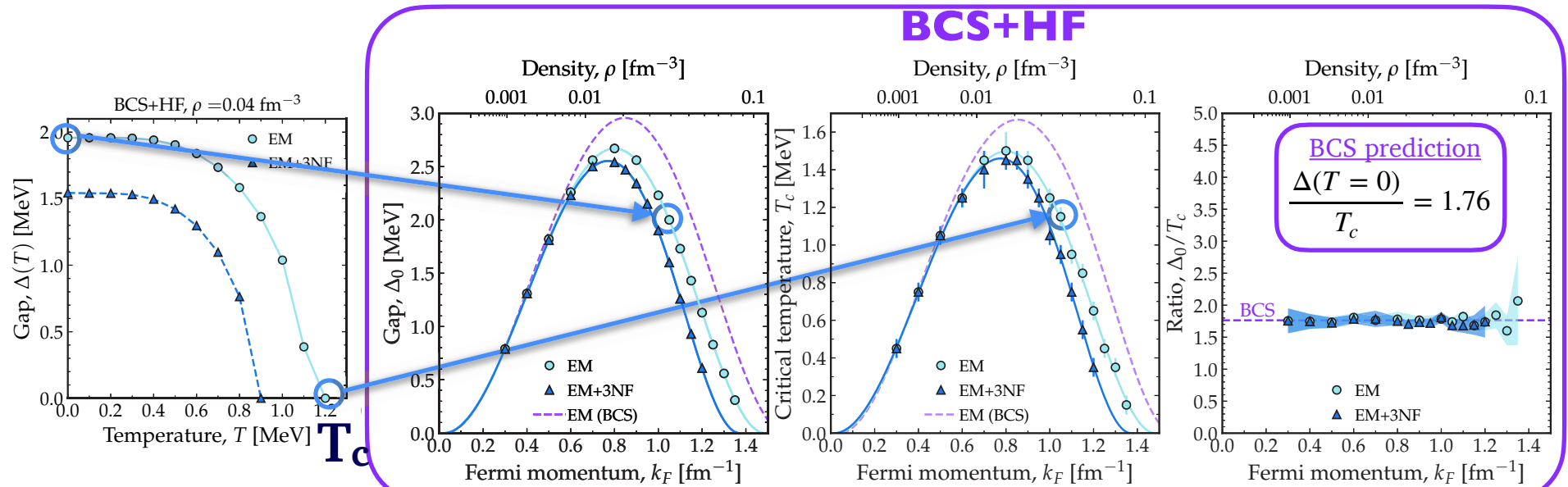


Many-body method

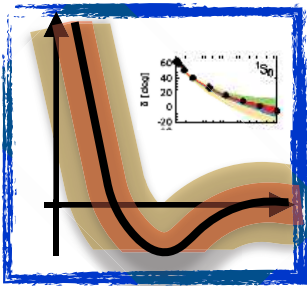


- **BCS** prediction based on **constant V**
- Full chiral **NN** interactions
- Full **HF spectrum**
- 3NFs do not change the ratio!

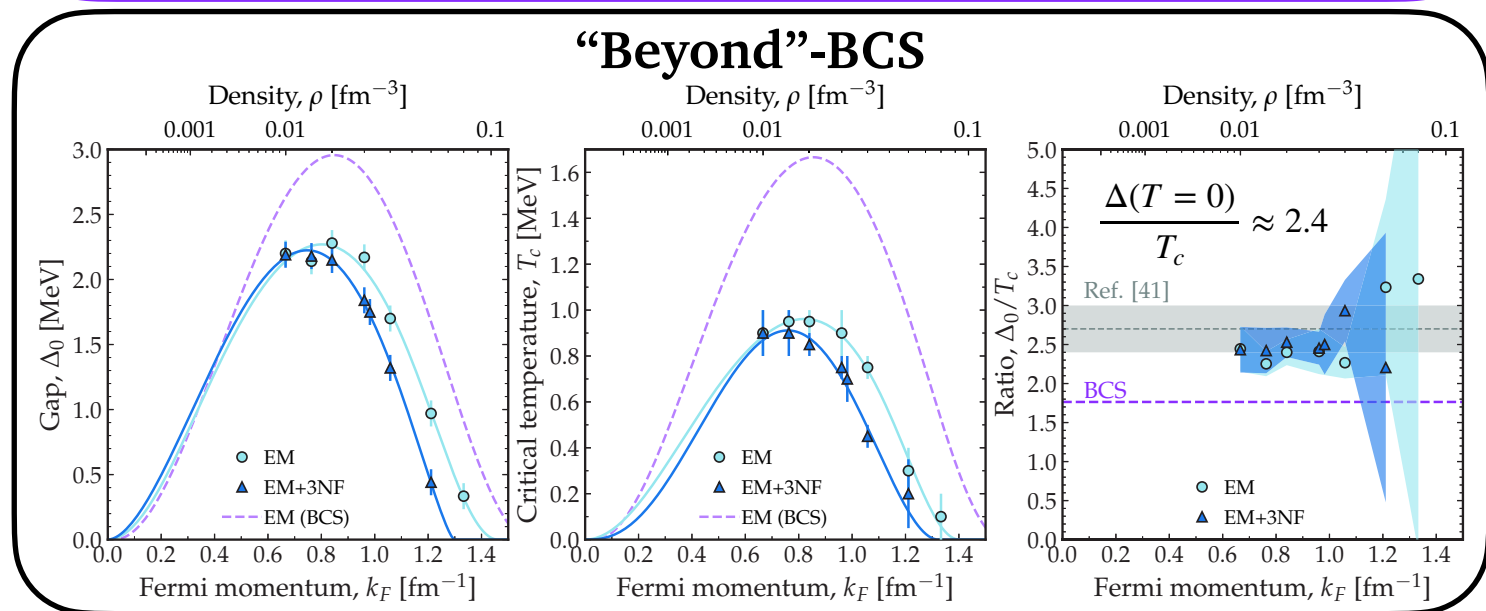
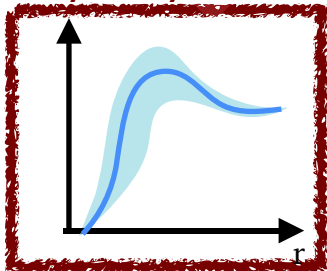
Finite Temperature beyond-BCS



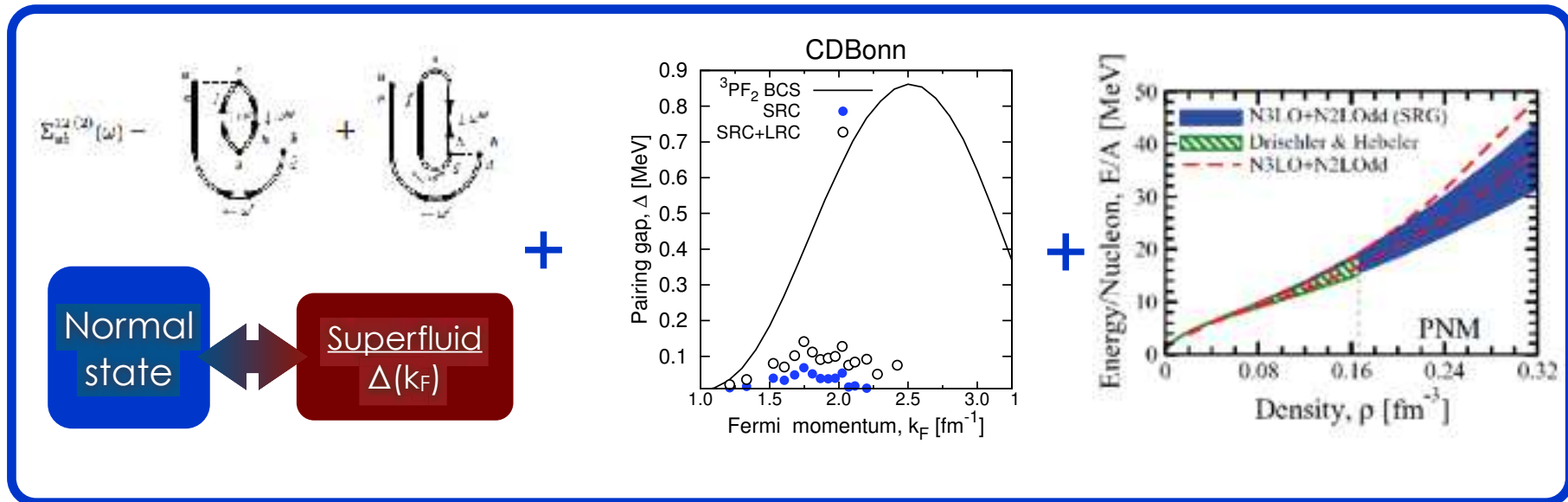
Hamiltonian



Many-body method

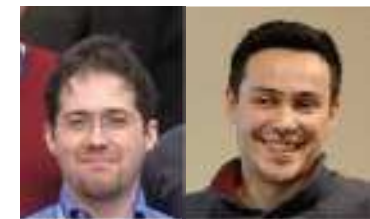


UB Ho do we go consistently beyond BCS?



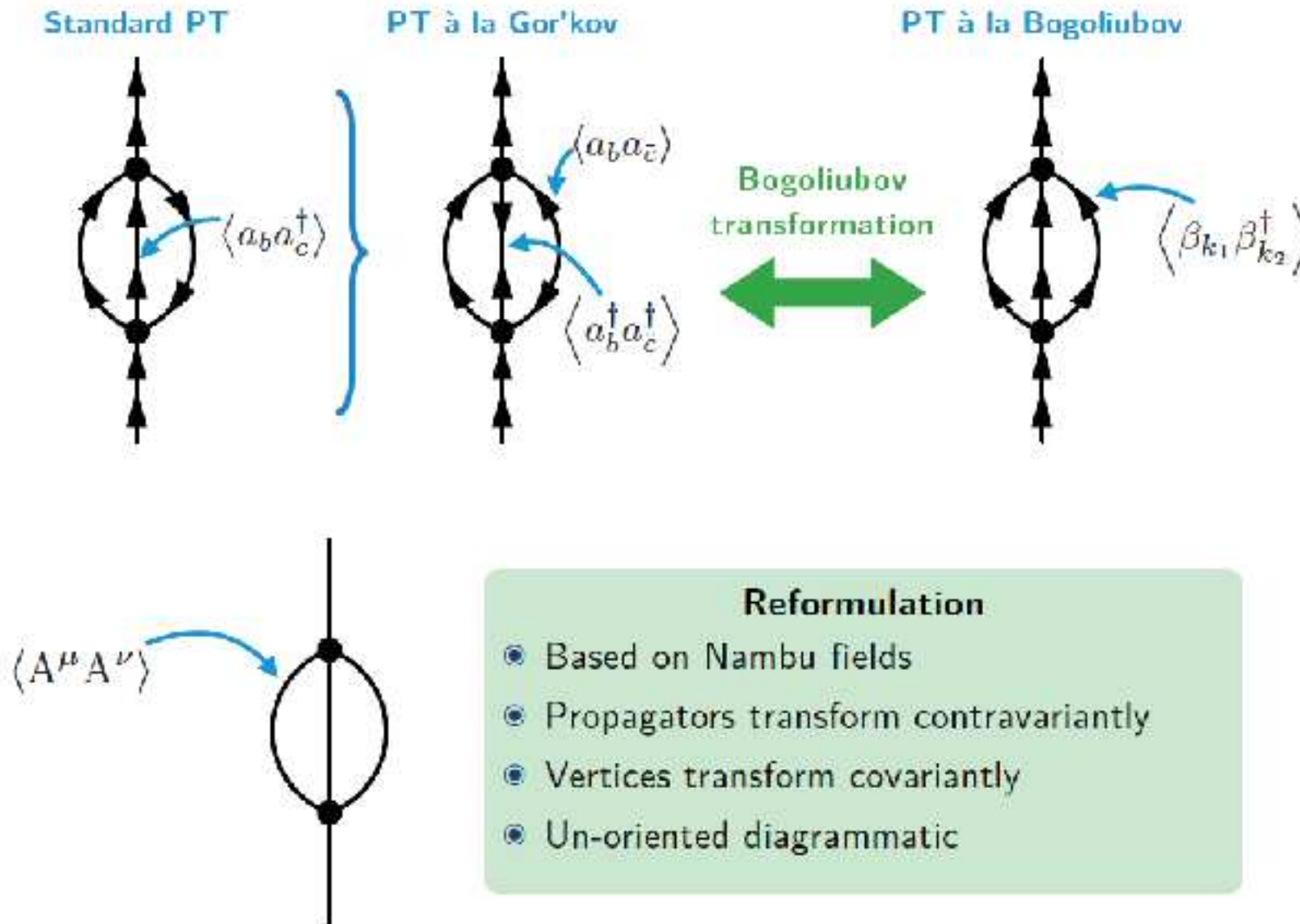
- Existing frameworks difficult to generalise
- Nambu-covariant **SCGF** technique

- Symmetry breaking ✓
- Finite temperature ✓
- Systematic expansion w diagrams ✓
- 3 nucleon forces ✓



Barbieri, Drissi

What was the issue before?



UB Nambu-Covariant Perturbation Theory

Double dimension H space

$$\mathcal{H}_1^e \equiv \mathcal{H}_1 \times \mathcal{H}_1^\dagger$$

Basis

$$\mathcal{B}^e \equiv \mathcal{B} \cup \bar{\mathcal{B}}$$

$$|b\rangle$$

$$\langle \bar{b}|$$

Elements

$$\begin{pmatrix} |\Psi_1\rangle \\ \langle \Psi'_1| \end{pmatrix}$$

$$\mu = (b, l)$$

$$\bar{\cdot} : 1 \mapsto \bar{1} = 2$$

$$\bar{\mu} \equiv (b, \bar{l})$$

$$2 \mapsto \bar{2} = 1$$

$$|b, 1\rangle \equiv \begin{pmatrix} |b\rangle \\ 0 \end{pmatrix} \quad |b, 2\rangle \equiv \begin{pmatrix} 0 \\ \langle \bar{b}| \end{pmatrix}$$

Product & metric tensor

$$g \left(\begin{pmatrix} |\Psi_1\rangle \\ \langle \Psi'_1| \end{pmatrix}, \begin{pmatrix} |\Psi_2\rangle \\ \langle \Psi'_2| \end{pmatrix} \right) \equiv \langle \Psi'_2 | \Psi_1 \rangle + \langle \Psi'_1 | \Psi_2 \rangle$$

$$g_{\mu\nu} \equiv g(|\mu\rangle, |\nu\rangle) = \delta_{\mu\bar{\nu}}$$

Nambu fields

$$\begin{aligned}
 A^{(b,1)} &\equiv a_b , & A^\mu &\equiv A^{(b,g)} = \begin{pmatrix} a_b \\ a_b^\dagger \end{pmatrix} \\
 A^{(b,2)} &\equiv \bar{a}_b , & & \\
 \bar{A}_{(b,1)} &\equiv \bar{a}_b , & \bar{A}_\mu &\equiv \bar{A}_{(b,g)} = \begin{pmatrix} a_b^\dagger & a_b \end{pmatrix} \\
 \bar{A}_{(b,2)} &\equiv a_b . & &
 \end{aligned}$$

$\bar{a}_b = a_b^\dagger \neq a_b^\dagger$

Commutator relations

(On extended indices!)

$$\begin{aligned}
 \{ A^\mu, A^\nu \} &= g^{\mu\nu} , \\
 \{ A^\mu, \bar{A}_\nu \} &= g^\mu{}_\nu , \\
 \{ \bar{A}_\mu, A^\nu \} &= g_\mu{}^\nu , \\
 \{ \bar{A}_\mu, \bar{A}_\nu \} &= g_{\mu\nu}
 \end{aligned}$$

Co- or contravariant

$$\begin{aligned}
 \bar{A}_\mu &= \sum_\nu g_{\mu\nu} A^\nu , \\
 \bar{A}_\mu &= \sum_\nu g_\mu{}^\nu \bar{A}_\nu , \\
 A^\mu &= \sum_\nu g^{\mu\nu} \bar{A}_\nu , \\
 A^\mu &= \sum_\nu g^\mu{}_\nu A^\nu .
 \end{aligned}$$

Perturbative construction

Hamiltonian partitioning

$$\Omega = \Omega_0 + \Omega_1$$

$$\Omega_0 = \frac{1}{2} \sum_{\mu\nu} U_{\mu\nu} A^\mu A^\nu$$

$$\Omega_1 = \sum_{k=1}^n \frac{1}{(2k)!} \sum_{\mu_1 \dots \mu_{2k}} v_{\mu_1 \dots \mu_{2k}}^{(k)} A^{\mu_1} \dots A^{\mu_{2k}}$$

Covariant **k-body** vertices

Fully antisymmetric vertex

Definition

$$v_{[\mu_1 \mu_2 \dots \mu_{2k-1} \mu_{2k}]}^{(k)} \equiv \frac{1}{(2k)!} \sum_{\sigma \in S_{2k}} \epsilon(\sigma) v_{\mu_{\sigma(1)} \mu_{\sigma(2)} \dots \mu_{\sigma(2k-1)} \mu_{\sigma(2k)}}^{(k)}$$

Antisymmetrisation defines a new $(0,2k)$ -tensor

Not the case in a *mixed* representation

Green's functions

Contravariant **k-body** Green's function

$$(-1)^k \mathcal{G}^{\mu_1 \dots \mu_{2k}}(\tau_1, \dots, \tau_{2k}) \equiv \left\langle T [A^{\mu_1}(\tau_1) \dots A^{\mu_{2k}}(\tau_{2k})] \right\rangle$$

$$\text{with } \langle . \rangle = \text{Tr} (. \rho) \text{ and } \rho \equiv \frac{e^{-\beta \Omega}}{\text{Tr} (e^{-\beta \Omega})}$$

Unperturbed case: $\Omega \longleftrightarrow \Omega_0$

Propagators

$$-\mathcal{G}^{\mu\nu}(\omega_p) = \begin{array}{c} \mu \\ \parallel \\ \nu \end{array} \uparrow \omega_p$$

$$-(\mathcal{G}^{(0)})^{\mu\nu}(\omega_p) = \begin{array}{c} \mu \\ | \\ \nu \end{array} \uparrow \omega_p$$

Perturbative construction

Hamiltonian partitioning

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Covariant **k-body** vertices

• **Antisymmetrisation** defines a new $(0, 2k)$ -tensor

• **Not** the case in a *mixed* representation

Green's functions

• Contravariant **k-body** Green's function

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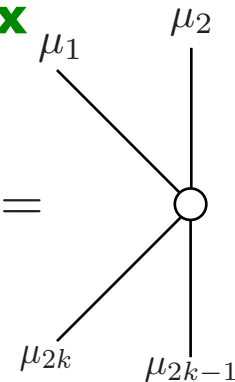
Propagators

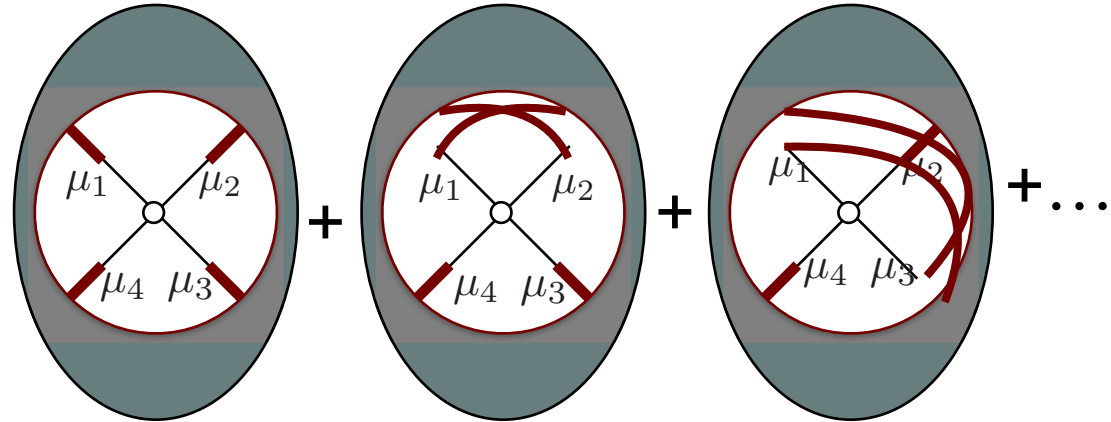
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$$-(\mathcal{G}^{(0)})^{\mu\nu}(\omega_p) = \begin{array}{c} \mu \\ | \\ \nu \end{array} \uparrow \omega_p$$

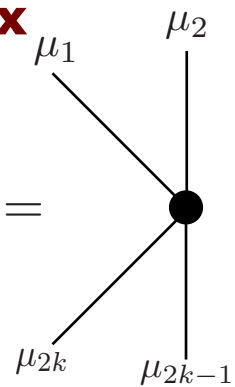
Why antisymmetric vertices?

Un-symmetrised vertex

$$v_{\mu_1 \mu_2 \dots \mu_{2k-1} \mu_{2k}}^{(k)} =$$




Antisymmetrized vertex

$$v_{[\mu_1 \mu_2 \dots \mu_{2k-1} \mu_{2k}]}^{(k)} =$$


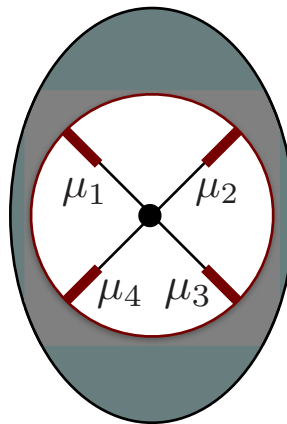


Diagram factorisation

- Derivations rely on
 - Wick theorem \Rightarrow sum over pairing
 - Sum over single-particle and Nambu indices
- Extends Hugenholtz antisymmetrisation**
- Antisym is a **one-off pre-computing** cost

Perturbative expansion

Order n graphical rules

- Draw all topologically distinct **connected unlabelled** diagrams
 - with $2k$ external legs
 - with n vertices (for order n contributions)

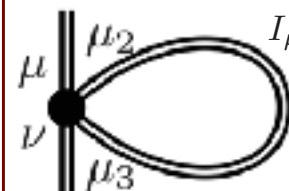
Feynman rules

1. Label vertices from 1 to n
 - S is the number of vertex labels permutations leaving the diagram invariant
2. For each line multiply by $-(\mathcal{G}^{(0)})^{\mu\nu}(\omega_e)$
3. For each k -body vertex multiply by $v_{[\mu_1 \mu_2 \dots \mu_{2k-1} \mu_{2k}]}^{(k)}$
4. Sum over each internal μ index and each independent ω_e frequency
5. Multiply by $\frac{(-1)^{n+L}}{S \times 2^T \prod_{l=2}^{l_{\max}} (l!)^m}$

Gaudin rules

- These simplify Matsubara sums
- Require **spanning trees**

Tadpoles are exceptional



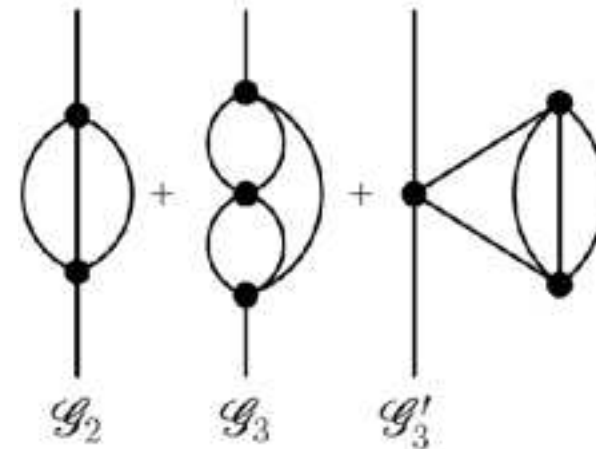
$$I_{\mu\nu} = \sum_{\mu_2 \dots \mu_{2k-1}} \frac{(-1)^k}{2^{k-1}(k-1)!} v_{[\mu \dot{\mu}_2 \dot{\mu}_3 \nu]}^{(k=2)} \times \frac{1}{\beta} \sum_{\omega_e} -\mathcal{G}^{\mu_2 \mu_3}(\omega_e) e^{-i\omega_e \eta_p}$$

- **Partially antisymmetrized** vertices needed:

$$v_{[\mu_1 \dots \dot{\mu}_x \dots \dot{\mu}_y \dots \mu_{2k}]}^{(k)} \equiv \frac{2^p p!}{(2k)!} \sum_{\sigma \in S_{2k}/S_2^p \times S_p} \epsilon(\sigma) v_{\mu_{\sigma(1)} \dots \dot{\mu}_x \dots \dot{\mu}_y \dots \mu_{\sigma(2k)}}^{(k)}$$

- p internal lines are **fixed**
- k -body **generalisation** works

HFB partitioning 3rd order



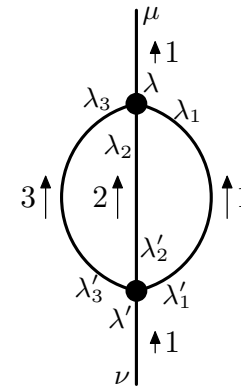
Advantages vs Gorkov

Gorkov GF

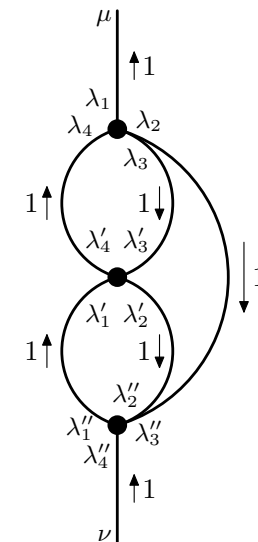
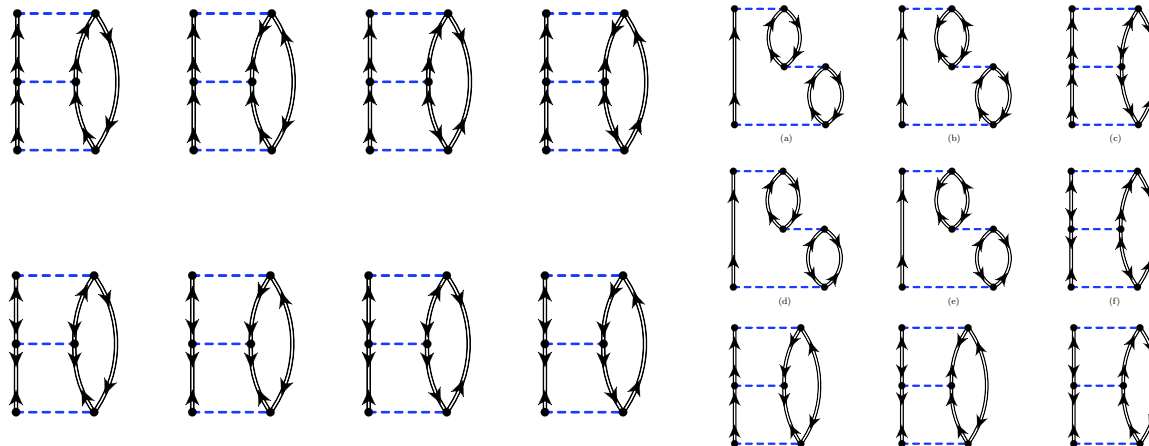
Order 2



NCGF



Order 3



UE Self-consistent Green's function resummation

Dyson equation

• Partitioning considered

$$\Omega = \underbrace{\frac{1}{2!} \sum_{\mu\nu} U_{\mu\nu} A^\mu A^\nu}_{\Omega_0} + \underbrace{\frac{1}{4!} \sum_{\alpha\beta\gamma\delta} v_{\alpha\beta\gamma\delta}^{(2)} A^\alpha A^\beta A^\gamma A^\delta}_{\Omega_1}$$

• Dyson equation

$$\mathcal{G}^{\mu\nu}(\omega_n) = \mathcal{G}^{(0)\mu\nu}(\omega_n) + \sum_{\lambda_1\lambda_2} \mathcal{G}^{(0)\mu\lambda_1}(\omega_n) \Sigma_{\lambda_1\lambda_2}(\omega_n) \mathcal{G}^{\lambda_2\nu}(\omega_n)$$

Diagrammatic expansion of

$$\Sigma_{\mu\nu}(\omega_n)$$

• with unperturbed propagators

$$\Sigma_{\mu\nu}(\omega_n) = \frac{\mathcal{J}_{\mu\nu}(\omega_n) - \mathcal{J}_{\nu\mu}(-\omega_n)}{2}$$

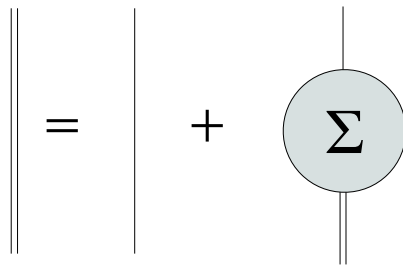
$$\mathcal{J}_{\mu\nu}(\omega_n) = \sum \text{1PI diagrams with } \mathcal{G}^{(0)}$$

• with self-consistent propagators

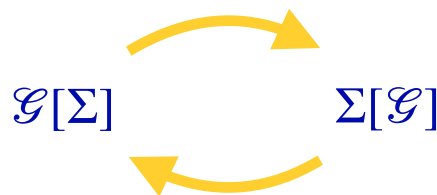
$$\Sigma_{\mu\nu}(\omega_n) = \frac{\mathcal{J}_{\mu\nu}(\omega_n) - \mathcal{J}_{\nu\mu}(-\omega_n)}{2}$$

$$\mathcal{J}_{\mu\nu}(\omega_n) = \sum \text{2PI diagrams with } \mathcal{G} \left(= \mathcal{J}_{\mu\nu}(\omega_n) \right)$$

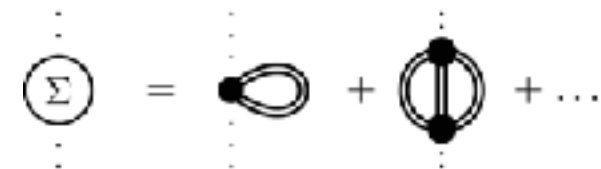
Diagrammatic representation



SCGF cycle

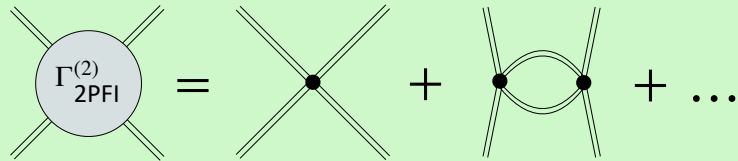


Self-energy expression



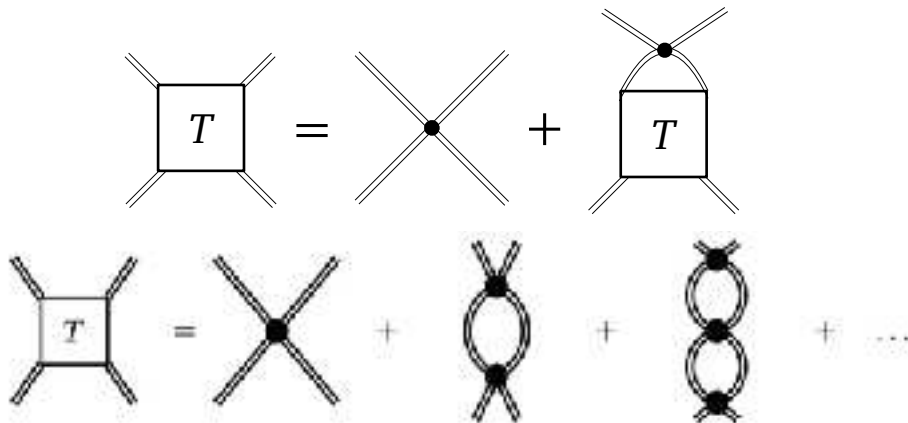
Approximations on $\Gamma_{2\text{PFI}}^{(2)}$

- Sum of all possible rungs



$$\Gamma_{2\text{PFI}}^{(2)} = \text{single vertex} + \text{bubble} + \dots$$

T-matrix $\equiv \Gamma^{(2)}$ in ladder approximation



$$T = \text{single vertex} + T \text{ with bubble}$$

$$T = \text{single vertex} + \text{bubble} + \text{chain of bubbles} + \dots$$

Ladder approximation

- Analytic/Retarded/Advanced/Sp function \Rightarrow as usual

- T-matrix equation

$$T_{MN}(Z) = V_{MN}^{(2)} + \frac{1}{2} \sum_{LL'} V_{ML}^{(2)} \Pi^{LL'}(Z) T_{L'N}(Z)$$

where

$$V_{MN}^{(2)} \equiv v_{[\mu_1 \mu_2 \nu_1 \nu_2]}^{(2)}, M \equiv (\mu_1, \mu_2) \text{ \& } N \equiv (\nu_1, \nu_2)$$

Solving the ladder

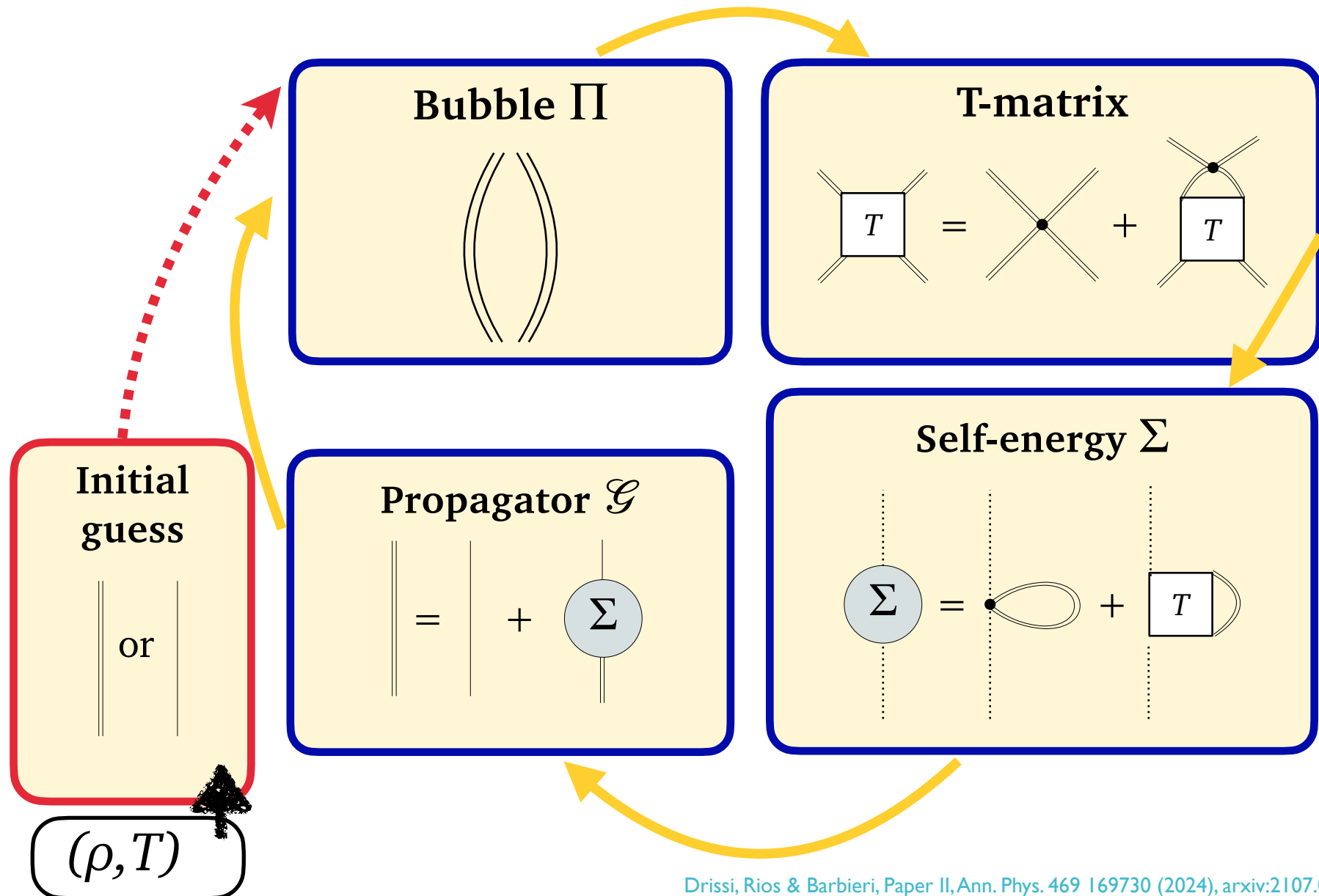
- Spectral representation

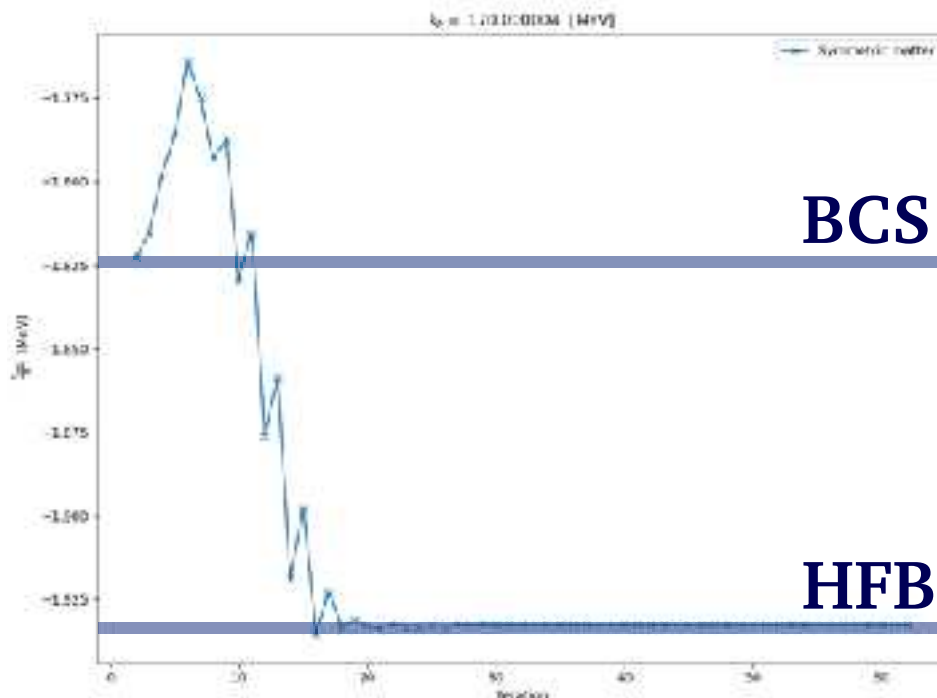
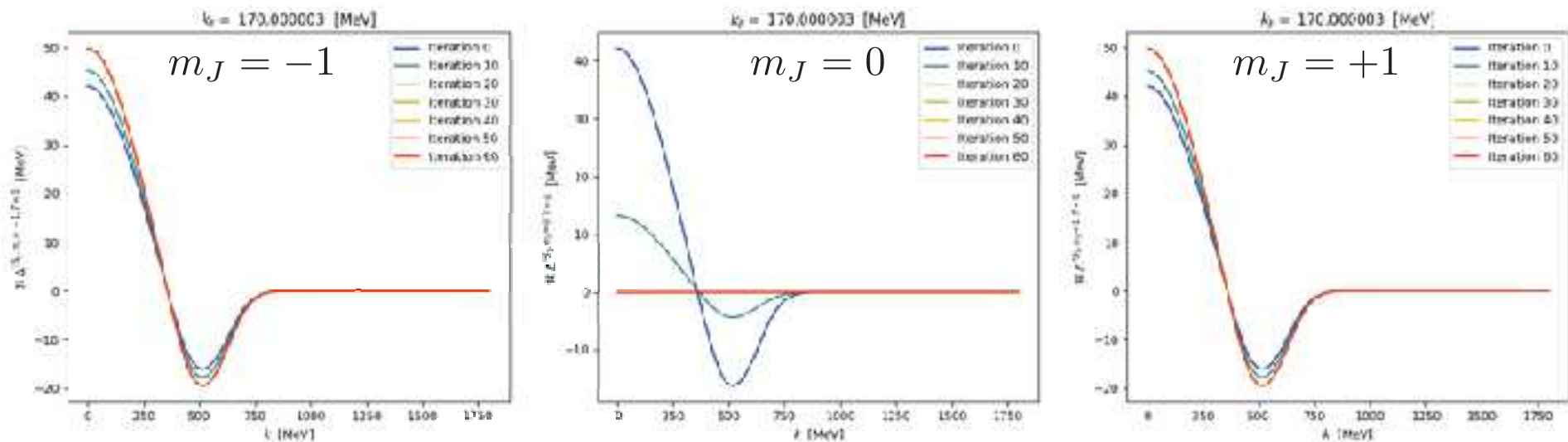
$$T_{MN}(Z) \equiv V_{MN}^{(2)} + \int_{-\infty}^{+\infty} \frac{d\Omega}{2\pi} \frac{\mathcal{T}_{MN}(\Omega)}{Z - \Omega}$$

- Solution

$$\mathcal{T}(\Omega) = iV^{(2)} \left\{ \left(gg - \frac{1}{2} \Pi^R(\Omega) V^{(2)} \right)^{-1} - \left(gg - \frac{1}{2} \Pi^A(\Omega) V^{(2)} \right)^{-1} \right\}$$

Nambu-Covariant Ladders





Preliminary!

- From BCS to HFB
- $^3\text{SD}_1$ channel
- N3LO EM
- $T=0.2$ MeV

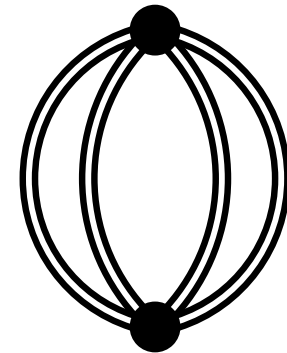
M Drissi



1) Thermal & microscopic properties

2) Nuclear uncertainty quantification

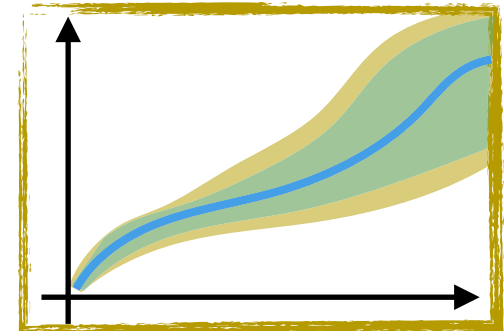
3) New superfluid extensions



Next:

Numerical implementation

Uncertainties in predictions?



Thank you!

Drissi, Rios & Barbieri, Ann. Phys. **469** 169729 (2024)
Drissi, Rios & Barbieri, Ann. Phys. **469** 169730 (2024)
Drissi & Rios, EPJA **58** 90 (2022)

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M Drissi



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