



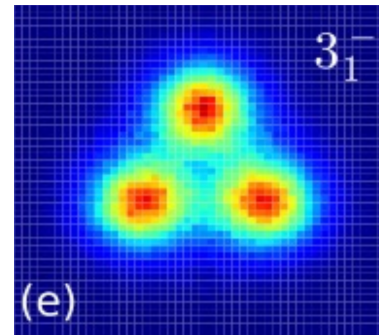
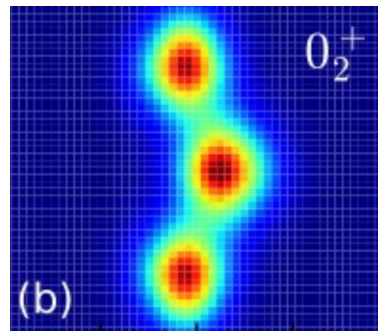
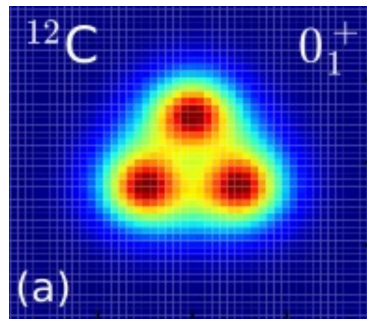
# EMERGENT GEOMETRY AND DUALITY IN THE CARBON NUCLEUS

ESNT Workshop “Espace de Structure et de Réactions Nucléaires Théorique”, CEA Paris-Saclay

04/12 2024 | TIMO A. LÄHDE (FORSCHUNGSZENTRUM JÜLICH)

# Emergent geometry and duality in the carbon nucleus

**Collaborators:** Shihang Shen, Dean Lee, Bing-Nan Lu, Ulf-G. Meißner



*Light*  
**Nuclei**  
Workshop  
CEA Paris-Saclay  
**2024**

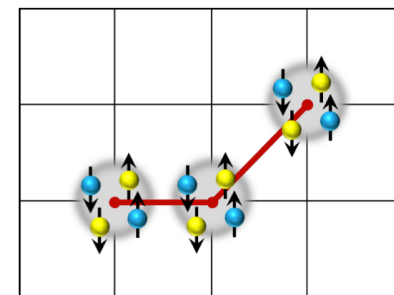


# CONTENTS

- Methods of Nuclear Lattice Effective Field Theory (NLEFT)  
— see also talk (Friday) by Lukas Bovermann
- Calculating the spectrum of  $^{12}\text{C}$  — Wigner SU(4) symmetry
- Accessing the cluster structure — Pinhole Algorithm
- Emergence of alpha clustering in  $^{12}\text{C}$  — Tomography
- Summary & Outlook

# AB INITIO STUDIES (A SAMPLE) OF $^{12}\text{C}$

- Green's function Monte Carlo → see Gandolfi, Lonardonì, Lovato, Piarulli, *Front. Phys.* **8** (2020) 117
- Monte Carlo Shell Model → see Otsuka, Abe, Yoshida, Tsunoda et al., *Nature Commun.* **13** (2022) 2234
- No-core shell model → see Choudhary, Srivastava, Gennari, Navrátil, *PRC* **107** (2023) 014309
  
- Our method — Auxiliary Field Quantum Monte Carlo (AFQMC)
  - First NLEFT work on ground and Hoyle states of  $^{12}\text{C}$ 
    - Epelbaum, Krebs, Lee, Meißner, *PRL* **106** (2011) 192501
    - Epelbaum, Krebs, Lähde, Lee, Meißner, *PRL* **109** (2012) 252501
  - Structure of the Hoyle state?
    - Freer, Fynbo, *Prog. Part. Nucl. Phys.* **78** (2014) 1
  
- Methods have progressed since 2012 — full AFQMC calculation of the  $^{12}\text{C}$  spectrum now possible



24 rotational orientations



# NUCLEAR LATTICE EFT

Frank, Brockmann (1992); Koonin, Müller, Seki, van Kolck (2000); Lee, Schäfer (2004) ...

Borasoy, Krebs, Lee, Meißner, NPA **768** (2006) 179; Borasoy, Epelbaum, Krebs, Lee, Meißner, EPJA **31** (2007) 105

- A new method for the nuclear many-body problem
- Discrete space and Euclidean time, theory of fermions (nucleons) with spin and isospin

$$V = L_s^3 \times L_t$$

- Chiral EFT potential with (smeared) contact interactions, pion exchanges + Coulomb

→ see Epelbaum, Hammer, Meißner, Rev. Mod. Phys. **81** (2009) 1773

- Physics independent of lattice spacing

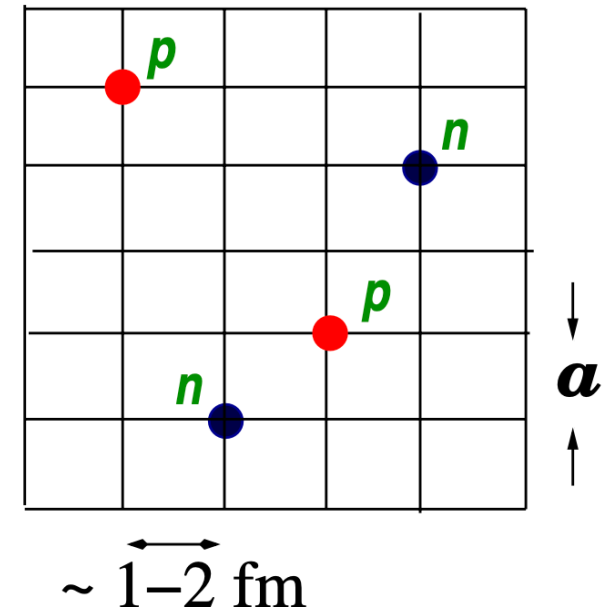
$$a = 1 \dots 2 \text{ fm} \longrightarrow p_{\text{max}} = \pi/a = 315 \dots 630 \text{ MeV}$$

Alarcón et al., EPJA **53** (2017) 83; Klein et al., EPJA **54** (2018) 121

- Strong suppression of sign oscillations due to (approximate) Wigner SU(4) spin-isospin symmetry — very important for Monte Carlo methods

Wigner, Phys. Rev. **51** (1937) 106; Mehen et al., PRL **83** (1999) 931

Chen, Lee, Schäfer, PRL **93** (2004) 242302



# EUCLIDEAN TIME PROJECTION

Lee, Prog. Part. Nucl. Phys. **63** (2009) 117

Lähde, Meißner, Springer Lecture Notes in Physics **957** (2019) 1-396

- Euclidean time projection amplitude (for many-body Hamiltonian  $H$ )

$$Z(\tau) = \langle \Psi | \exp(-H\tau) | \Psi \rangle$$

Slater determinant for  $A$  nucleons  
(more sophisticated trial states also possible)

- Transient energy

$$E(\tau) = -\frac{d}{d\tau} \ln Z(\tau), \quad E_0 = \lim_{\tau \rightarrow \infty} E(\tau)$$

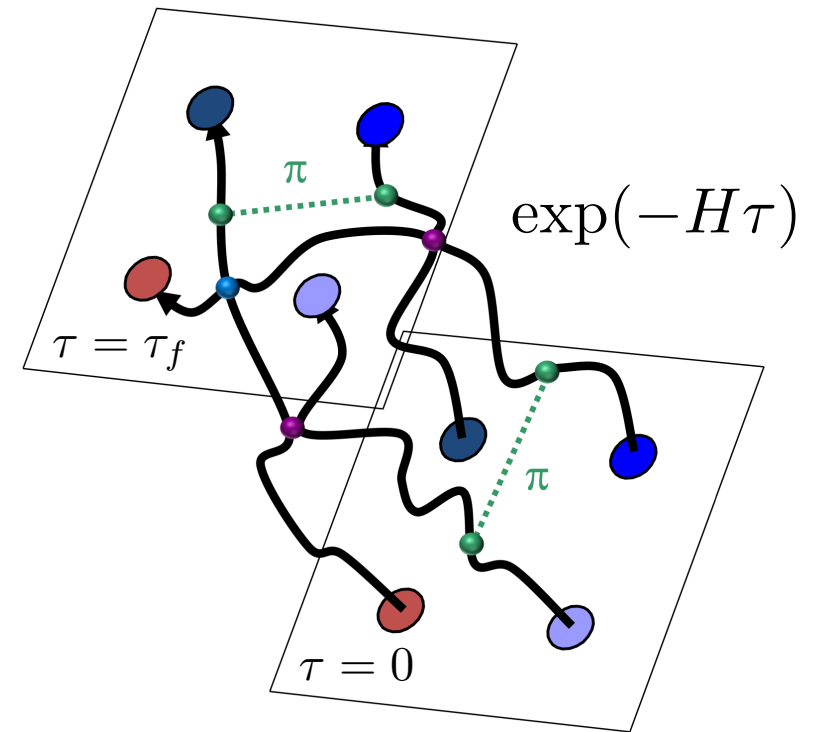
Ground state obtained for large times  
(usually reached by extrapolation)

- Expectation values of operators

$$Z^{\mathcal{O}}(\tau) = \langle \Psi | \exp(-H\tau/2) \mathcal{O} \exp(-H\tau/2) | \Psi \rangle$$

$$\langle \Psi_0 | \mathcal{O} | \Psi_0 \rangle = \lim_{\tau \rightarrow \infty} \frac{Z^{\mathcal{O}}(\tau)}{Z(\tau)}$$

Normal-ordered operators



# TRANSFER MATRIX METHOD

Lee, Prog. Part. Nucl. Phys. **63** (2009) 117

Lähde, Meißner, Springer Lecture Notes in Physics **957** (2019) 1-396

- Hamiltonian evolution divided into (small) transfer matrix steps

$$Z(L_t) = \langle \Psi | M^{L_t} | \Psi \rangle$$

$$M = : \exp(-\alpha_t H) :, \quad H = H_{\text{kin}} + V$$

$$\alpha_t = a_t/a, \quad \tau = a_t L_t$$

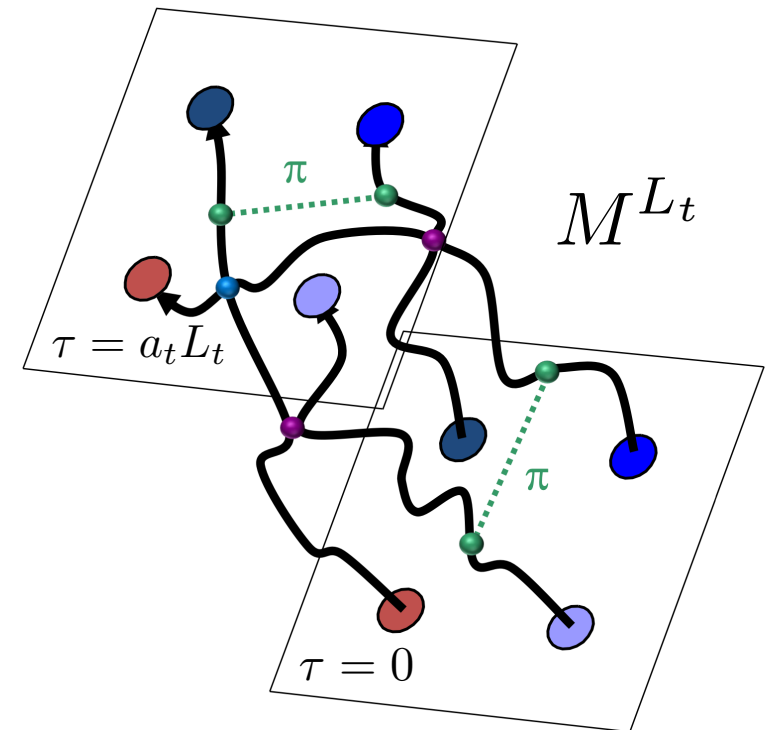
Suzuki-Trotter discretization of the Euclidean time evolution

- Transient energy and expectation values

$$E_i(L_t) = -\frac{1}{\alpha_t} \ln \frac{Z(L_t + 1)}{Z(L_t)}$$

$$Z^{\mathcal{O}}(L_t) = \langle \Psi | M^{L_t/2} \mathcal{O} M^{L_t/2} | \Psi \rangle$$

Practical simulations close to Hamiltonian limit, requires special MC algorithm → see next pages



# AUXILIARY FIELD METHOD

Lee, Prog. Part. Nucl. Phys. **63** (2009) 117

Lähde, Meißner, Springer Lecture Notes in Physics **957** (2019) 1-396

- Interactions between nucleons are represented by a path integral over auxiliary and pion fields

$$\exp\left(-\frac{C}{2}\rho(\mathbf{n})^2\right) = \frac{1}{\sqrt{2\pi}} \int ds \exp\left(-\frac{s^2}{2} + \sqrt{C}\rho(\mathbf{n})s\right)$$

Applied on every lattice site (many types of Hubbard-Stratonovich transformations possible)

- Auxiliary Field  
Quantum Monte Carlo (AFQMC)

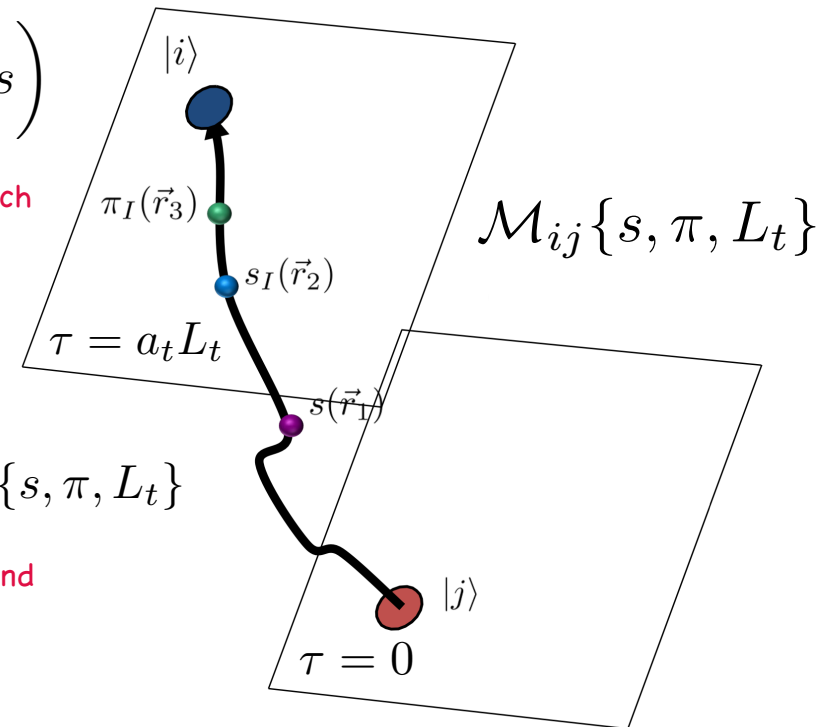
$$Z(L_t) = \int \mathcal{D}s \mathcal{D}\pi \langle \Psi | M^{L_t} \{s, \pi\} | \Psi \rangle = \int \mathcal{D}s \mathcal{D}\pi \det \mathcal{M} \{s, \pi, L_t\}$$

Each nucleon is evolved independently in a fluctuating background

- Fermion correlation matrix

$$\mathcal{M}_{kl} \{s, \pi, L_t\} = \langle \phi_k | M^{L_t} \{s, \pi\} | \phi_l \rangle$$

The trial state is defined in terms of the single-nucleon components of the Slater determinant  
→ Great simplification (at the price of introducing a path integral)



# MONTE CARLO ALGORITHMS

Lu, Li, Elhatisari, Lee, Epelbaum, Meißner, PLB **797** (2019) 134863

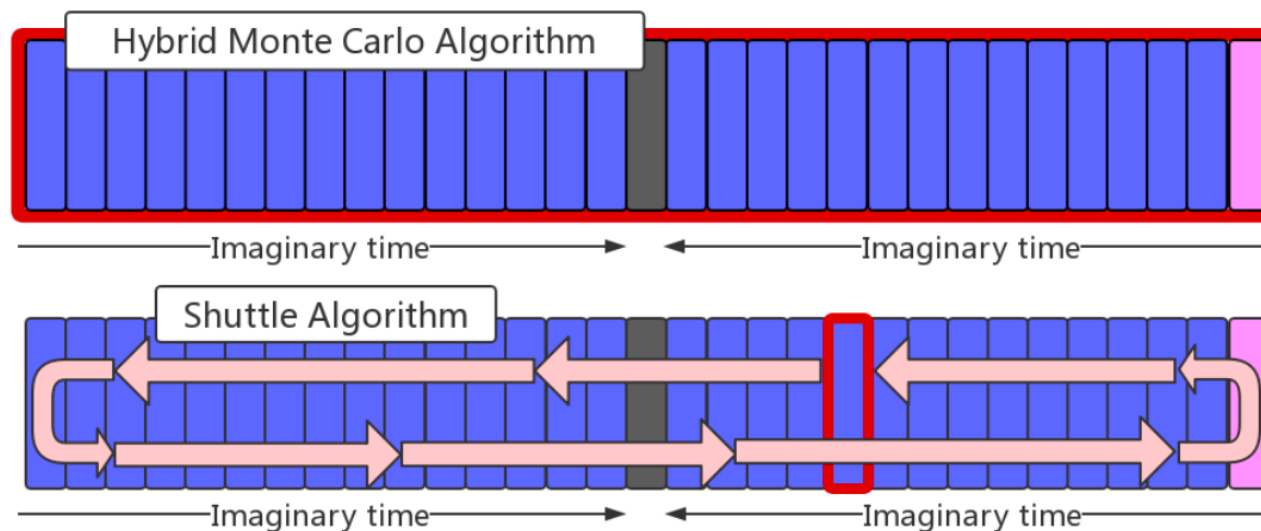
Lähde, Meißner, Springer Lecture Notes in Physics **957** (2019) 1-396

- Hybrid Monte Carlo (HMC) algorithm — *à la* Lattice QCD

Global Molecular Dynamics updates, slow for large number of time slices

- “Shuttle algorithm” — designed specifically for NLEFT (also further acceleration with GPU computing)

Heat bath updates one time slice at a time, efficient for larger number of time slices



First time step: projection onto cubic irreps ( $A_{1+}$ ,  $E_+$ ,  $T_{2+}$ , etc.)

- Very favorable computational scaling  $\propto A^{2\dots 3}$



# SPECTRUM OF $^{12}\text{C}$

Shen, Lähde, Lee, Meißner, EPJA **57** (2021) 276 [arXiv:2106.04834[nucl-th]]

- To compute excited states and transition rates, use Euclidean time projection with coupled channels

Rokash, Pine, Elhatisari, Epelbaum et al., PRC **92** (2015), 054612

Elhatisari, Lee, Rupak, Epelbaum et al., Nature **528** (2015) 111

$$Z_{a,b}(L_t) = \langle \Psi_a | M^{L_t} | \Psi_b \rangle$$

- “Adiabatic transfer matrix” — introduced in the cluster Hamiltonian treatment of alpha-alpha scattering)

$$M_{qq'}^{(a)}(L_t) = \sum_{q''} Z_{qq''}^{-1}(L_t) Z_{q''q'}(L_t + 1)$$

$$\lambda_i(L_t) = \exp(-\alpha_t E_i(L_t))$$

$$E_i(L_t) = -\ln(\lambda_i(L_t))/\alpha_t$$

Typically 2-4 channels per calculation,  
with different choices of trial states → next page

Eigenvalues of adiabatic transfer matrix  
give transient energies

- Angular momentum/parity projection (cubic irreps)
- Extrapolation to infinite Euclidean time (multi-exponential fit)

Lähde, Epelbaum, Krebs, Lee, Meißner, J. Phys. G **42** (2015) 034012

# SPECTRUM OF $^{12}\text{C}$

Shen, Lähde, Lee, Lu, Meißner, Nature Commun. **14** (2023) 2777 [arXiv:2202.13596[nucl-th]]

- Wigner SU(4) symmetric interaction, with three-nucleon force (3NF)

Lu, Li, Elhatisari, Lee et al., PLB **797** (2019) 134863

$$H = H_{\text{kin}} + \frac{C_2}{2!} \sum_{\mathbf{n}} : \tilde{\rho}(\mathbf{n})^2 : + \frac{C_3}{3!} \sum_{\mathbf{n}} : \tilde{\rho}(\mathbf{n})^3 :$$

$$\tilde{\rho}(\mathbf{n}) = \sum_i \tilde{a}_i^\dagger(\mathbf{n}) \tilde{a}_i(\mathbf{n}) + s_L \sum_i \sum_{|\mathbf{n}'-\mathbf{n}|=1} \tilde{a}_i^\dagger(\mathbf{n}') \tilde{a}_i(\mathbf{n}')$$

$$\tilde{a}_i^{(\dagger)}(\mathbf{n}) = a_i^{(\dagger)}(\mathbf{n}) + s_{NL} \sum_{|\mathbf{n}'-\mathbf{n}|=1} a_i^{(\dagger)}(\mathbf{n}')$$

Sign problem largely suppressed  
Neutron matter well described  
Ground state properties of light and medium-mass nuclei well described

- 4 parameters to determine:

- $C_2, C_3$  — ground states energies of  $^4\text{He}$  and  $^{12}\text{C}$
- $s_L$  — radius of  $^{12}\text{C}$  around 2.4 fm
- $s_{NL}$  — optimal E0 and E2 transition rates

In the sense of Chiral EFT  
→ smearings are regulators of EFT

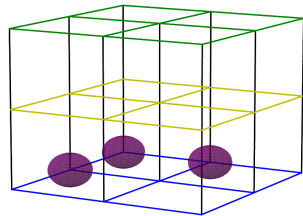
- Basis for high-fidelity NLEFT calculations (currently up to N3LO in the chiral EFT expansion) Elhatisari, Bovermann, Ma, Epelbaum et al., Nature **630** (2024) 59

# SPECTRUM OF $^{12}\text{C}$

Shen, Lähde, Lee, Meißner, EPJA **57** (2021) 276 [arXiv:2106.04834[nucl-th]]

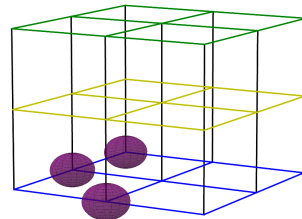
- Alpha cluster configurations of Gaussian wave packets

S1



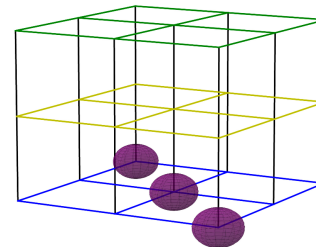
isosceles right triangle

S2



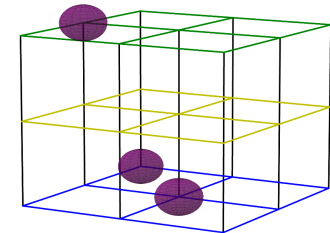
"bent-arm" shape

S3



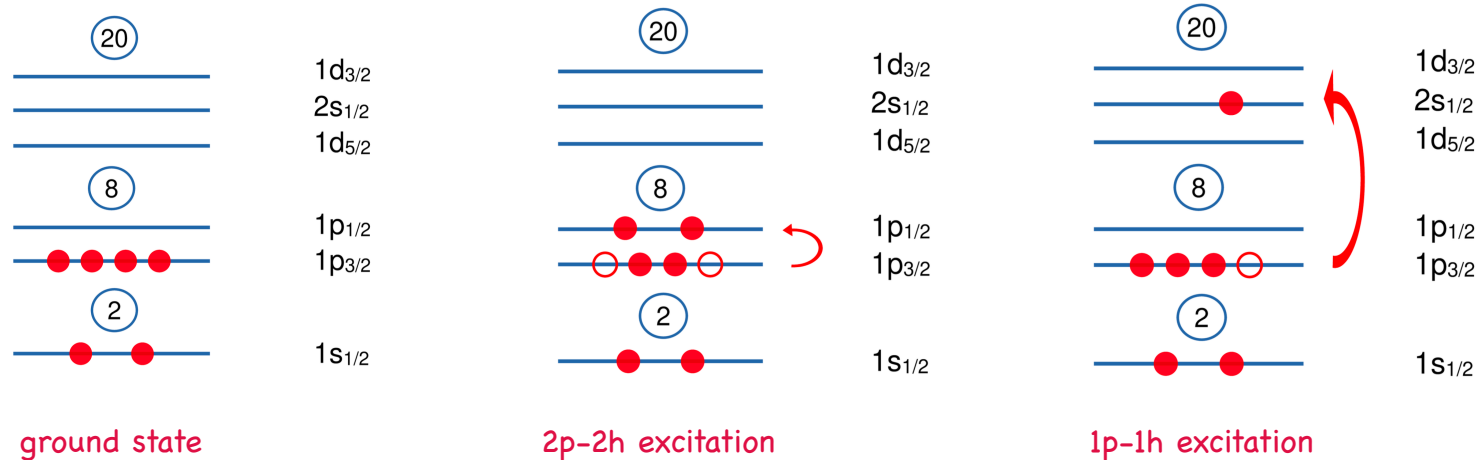
linear chain (diagonal)

S4



isosceles acute triangle

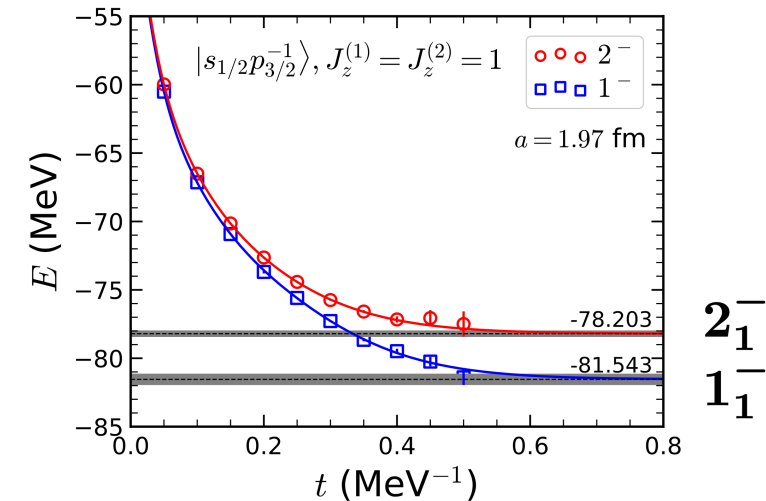
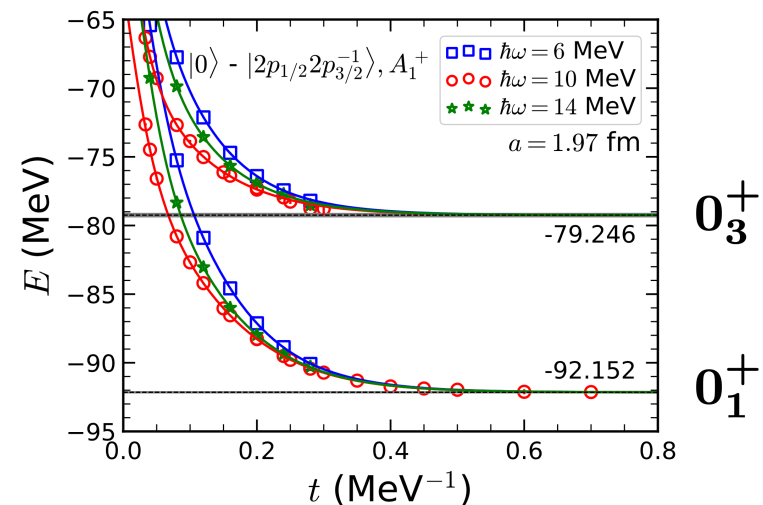
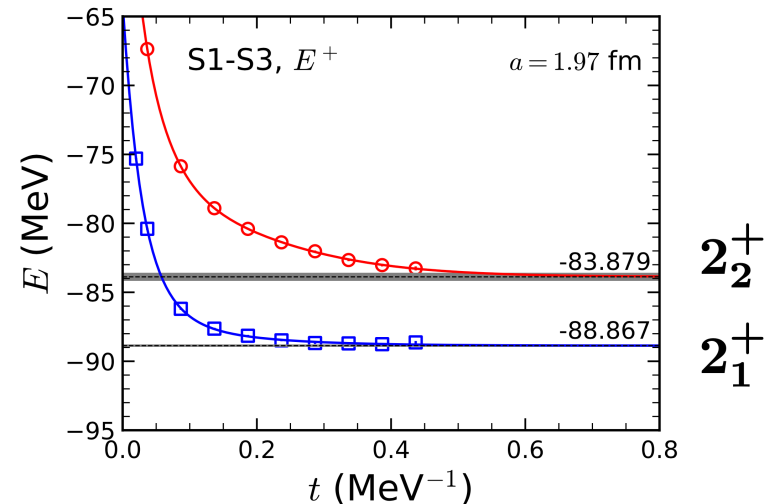
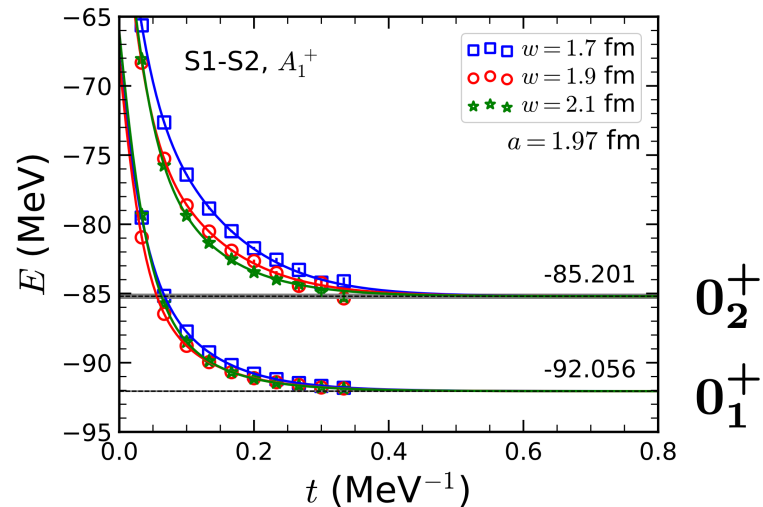
- Shell model (proton) particle-hole harmonic oscillator wave functions



# SPECTRUM OF $^{12}\text{C}$

Shen, Lähde, Lee, Meißner, EPJA **57** (2021) 276 [arXiv:2106.04834[nucl-th]]

- Transient energies (examples) with extrapolation to large projection time

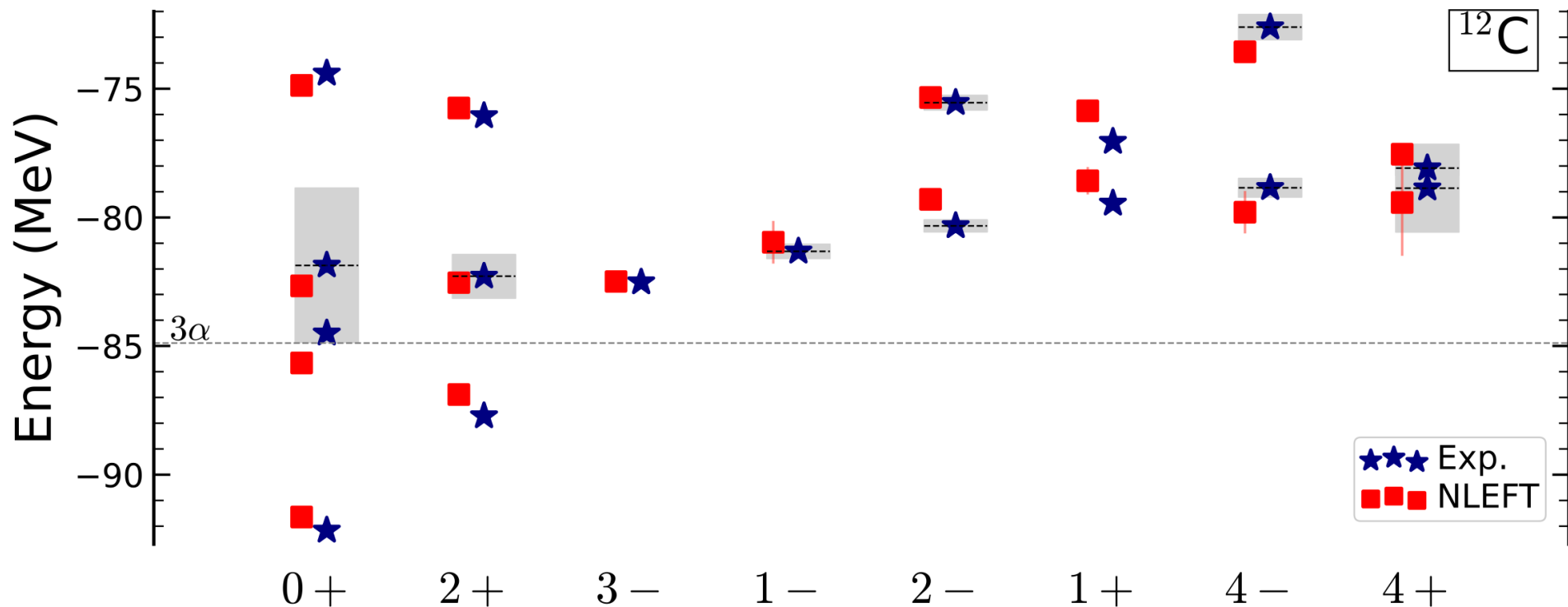


# SPECTRUM OF $^{12}\text{C}$

Shen, Lähde, Lee, Lu, Meißner, Nature Commun. 14 (2023) 2777 [arXiv:2202.13596[nucl-th]]

- All low-lying states amazingly good — resonances above the  $3\alpha$  threshold: need analysis of finite-volume spectrum

Bernard, Lage, Meißner, Rusetsky, JHEP 08 (2008) 024



- Next question — what is the structure of each state?



# TOMOGRAPHY OF $^{12}\text{C}$ STATES

Shen, Lähde, Lee, Lu, Meißner, Nature Commun. 14 (2023) 2777 [arXiv:2202.13596[nucl-th]]

- Insert a screen of “pinholes” with spin/isospin, that allow nucleons with corresponding spin/isospin to pass

Elhatisari, Epelbaum, Krebs, Lähde et al., PRL 119, 222505 (2017)

Lu, Li, Elhatisari, Lee et al., Phys. Lett. B 797, 134863 (2019)

$$Z_{a,b}(\mathbf{i}, \mathbf{j}, \mathbf{n}; L_t) = \langle \Psi_a | M^{L_t/2} \rho_A(\mathbf{i}, \mathbf{j}, \mathbf{n}) M^{L_t/2} | \Psi_b \rangle$$

$$\rho_A(\mathbf{i}, \mathbf{j}, \mathbf{n}) = : \rho(i_1, j_1, \mathbf{n}_1) \cdots \rho(i_A, j_A, \mathbf{n}_A) :$$

$$\rho(i, j, \mathbf{n}) = a_{i,j}^\dagger(\mathbf{n}) a_{i,j}(\mathbf{n})$$

Locations and spin/isospin updated with MC  
→ see next page

- Originally introduced for proton and neutron distributions in AFQMC (relative to the center of mass of the  $A$ -nucleon system)

$$\langle \rho(r) \rangle = \frac{1}{A!} \sum_{\mathbf{n}} \langle \Psi | M^{L_t/2} \rho_A(\mathbf{n}) M^{L_t/2} | \Psi \rangle \sum_{k=1}^A \delta(r - |\mathbf{r}_k - \mathbf{R}|)$$

$$\mathbf{R} = \frac{1}{A} \sum_{k=1}^A \mathbf{r}_k$$

Radial charge density profile

- Also computationally efficient stochastic estimator

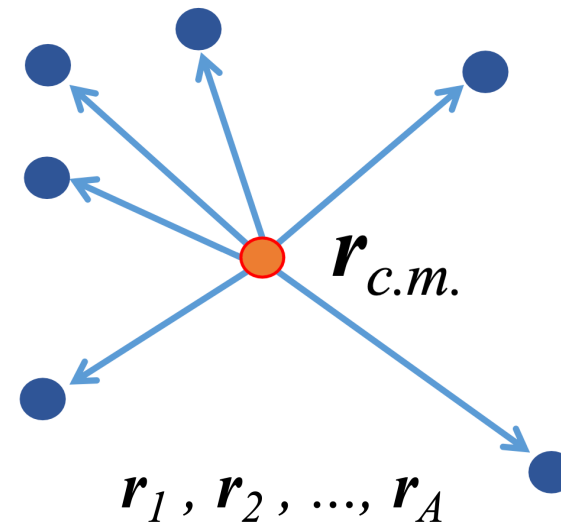
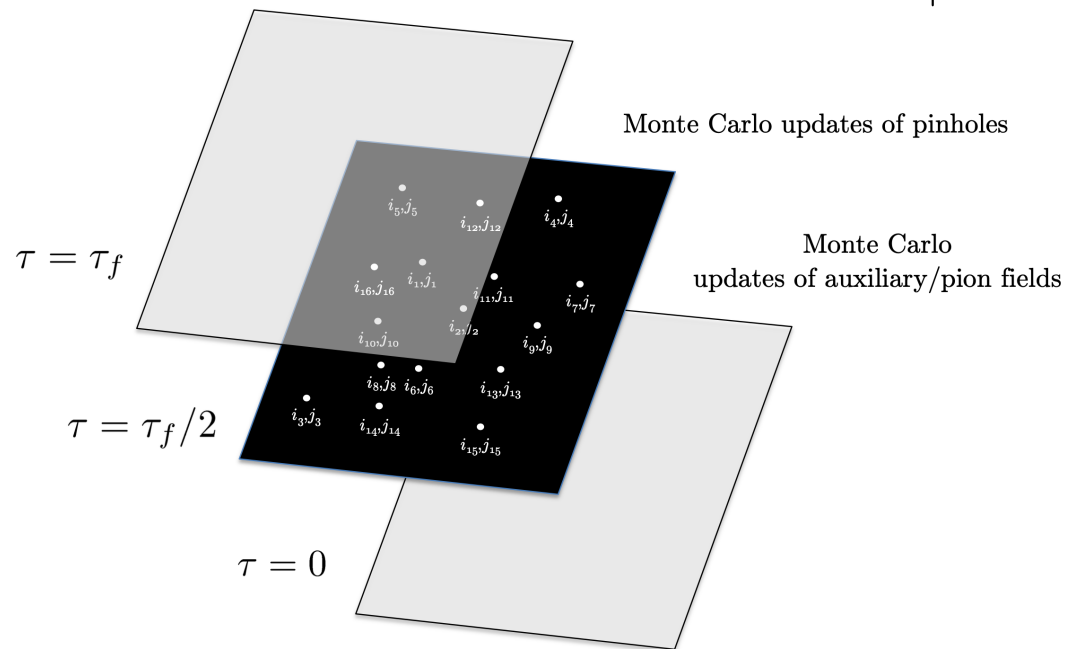
# TOMOGRAPHY OF $^{12}\text{C}$ STATES

Shen, Lähde, Lee, Lu, Meißner, Nature Commun. 14 (2023) 2777 [arXiv:2202.13596[nucl-th]]

- Metropolis updates of pinhole positions/spin/isospin

$$Z_{a,b}(\mathbf{i}, \mathbf{j}, \mathbf{n}; L_t) = \int \mathcal{D}s \mathcal{D}\pi \langle \Psi_a | M^{L_t/2} \{s, \pi\} \rho_A(\mathbf{i}, \mathbf{j}, \mathbf{n}) M^{L_t/2} \{s, \pi\} | \Psi_b \rangle$$

$$A(s, \pi; \mathbf{i}, \mathbf{j}, \mathbf{n}; L_t) = \left| \langle \Psi_a | M^{L_t/2} \{s, \pi\} \rho_A(\mathbf{i}, \mathbf{j}, \mathbf{n}) M^{L_t/2} \{s, \pi\} | \Psi_b \rangle \right|$$

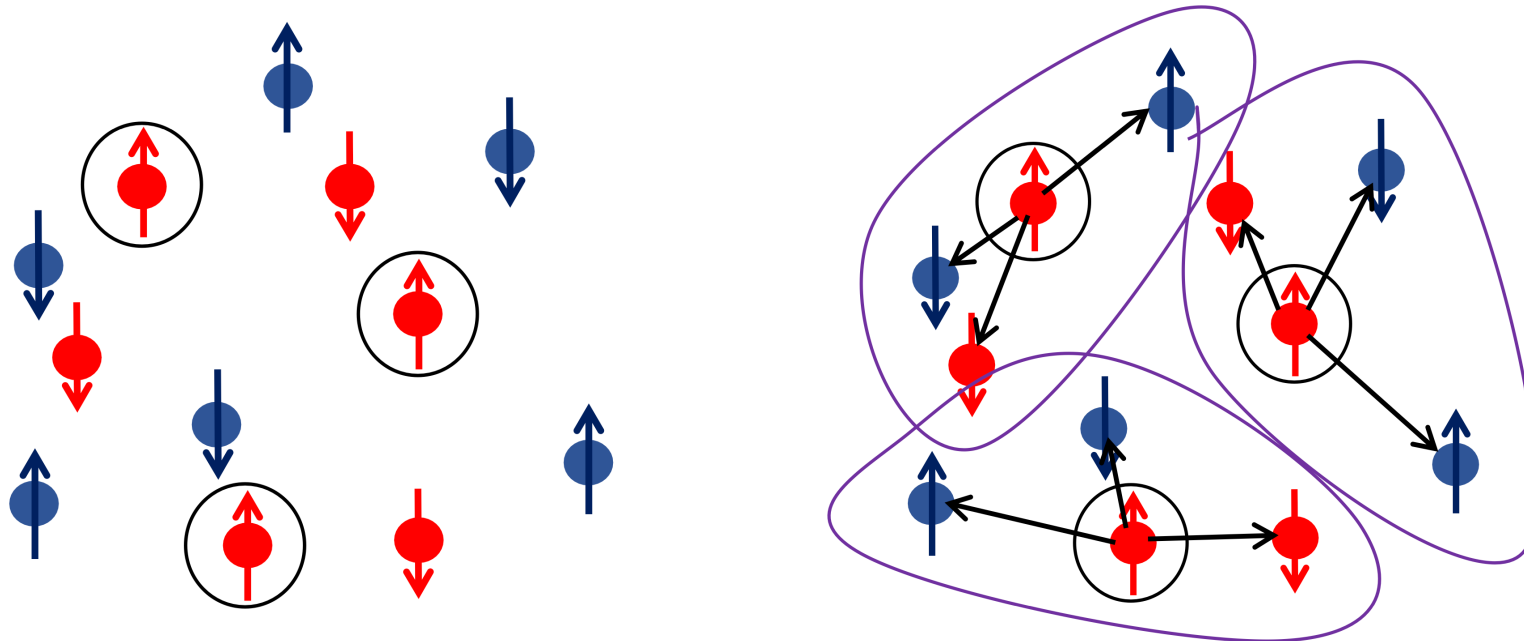


- Ensembles of configurations of  $A$  nucleons (relative to center of mass)

# TOMOGRAPHY OF $^{12}\text{C}$ STATES

Shen, Lähde, Lee, Lu, Meißner, *Nature Commun.* **14** (2023) 2777 [arXiv:2202.13596[nucl-th]]

- Search for alpha clusters using pinhole configurations:
  1. Identify 3 spin-up protons
  2. Find the closest possible candidate of the 3 other species

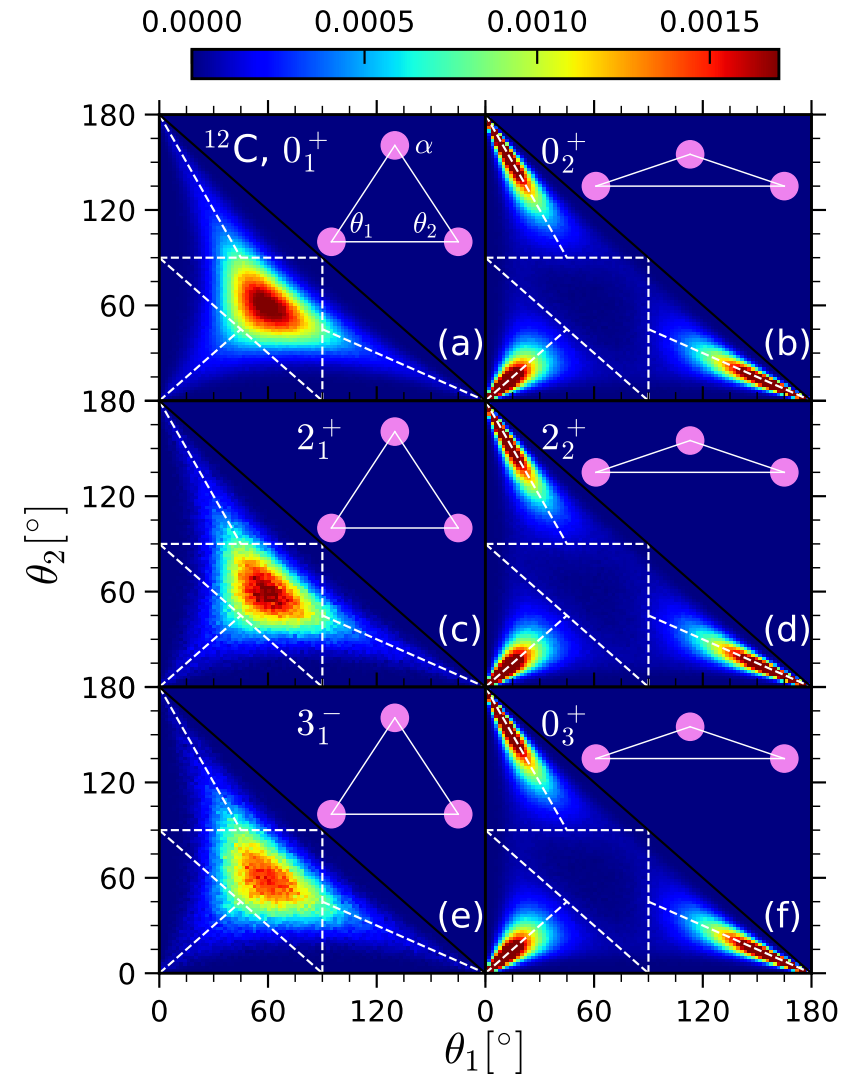
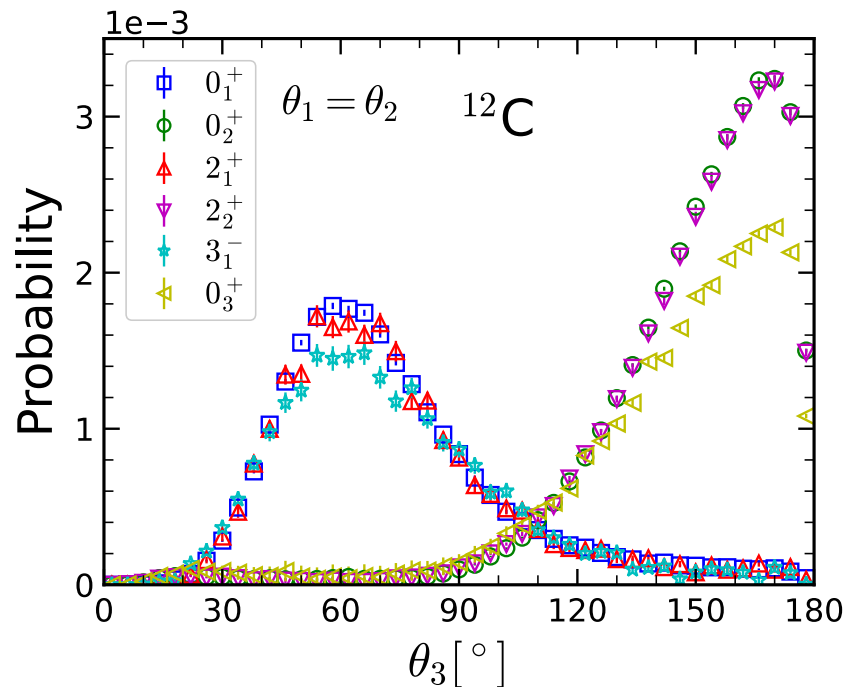


- Cluster radii (1.65 ... 1.71 fm) similar to free alpha particles (1.63 fm)

# TOMOGRAPHY OF $^{12}\text{C}$ STATES

Shen, Lähde, Lee, Lu, Meißner, Nature Commun. 14 (2023) 2777 [arXiv:2202.13596[nucl-th]]

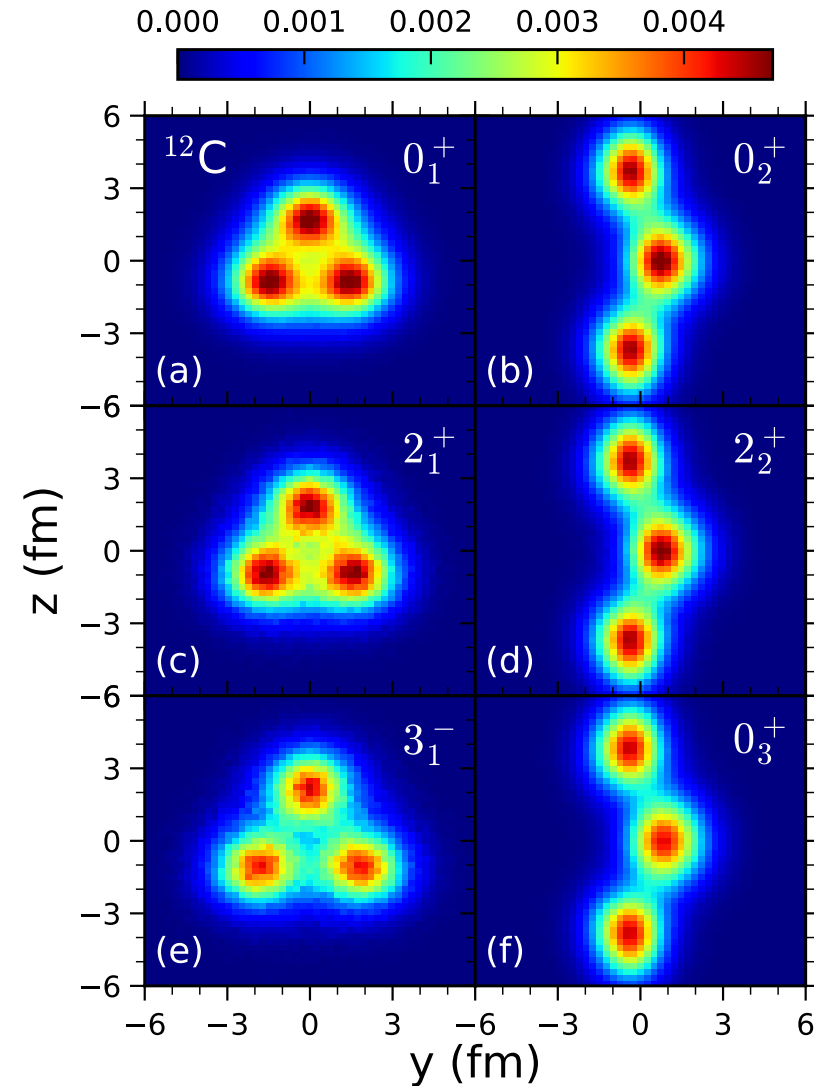
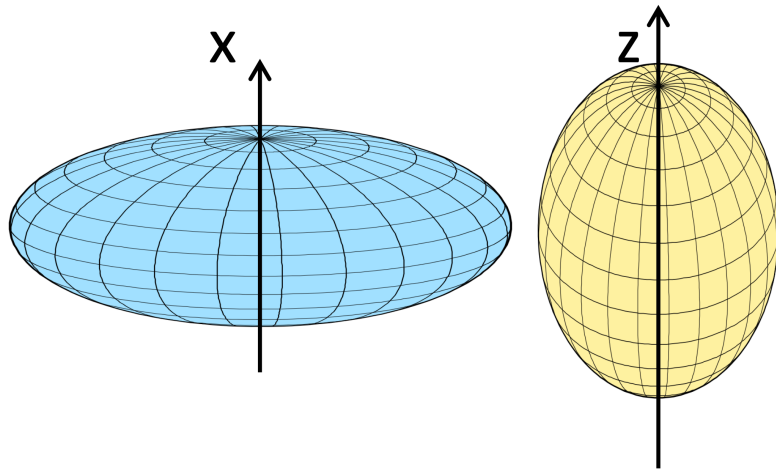
- Study angles between alpha clusters
- Two distinct structures emerge:
  1. Equilateral triangular
  2. Obtuse large-angle triangular



# TOMOGRAPHY OF $^{12}\text{C}$ STATES

Shen, Lähde, Lee, Lu, Meißner, Nature Commun. 14 (2023) 2777 [arXiv:2202.13596[nucl-th]]

- Equilateral-type (Ground) states:
  1. Align shortest axis to x
  2. Rotate one  $\alpha$  to ( $y = 0, z > 0$ )
  3. Random 120 degree rotation
- Obtuse-type (Hoyle) states:
  1. Align longest axis to z
  2. Rotate central  $\alpha$  to ( $x=0, y > 0$ )

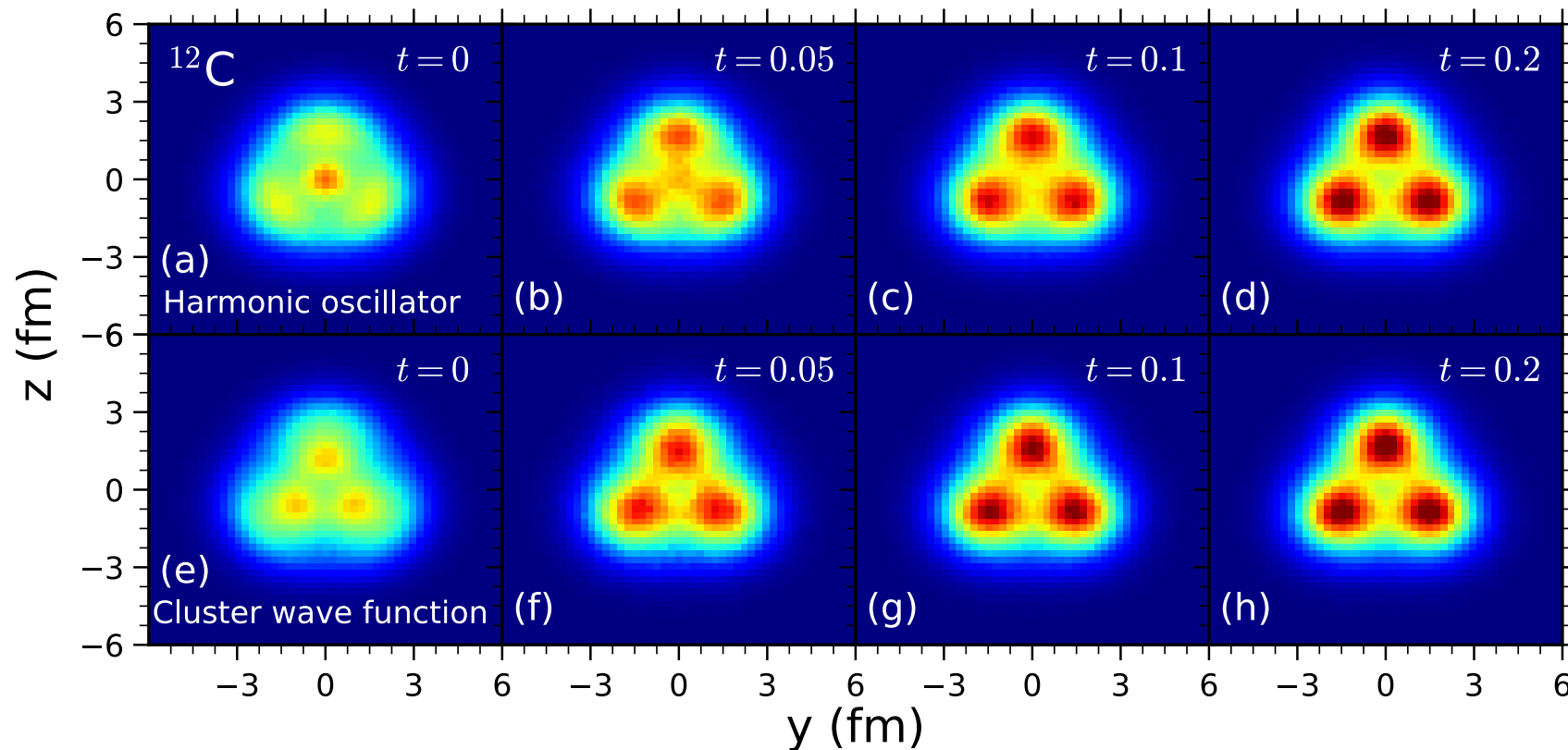




# EMERGENT ALPHA CLUSTERING

Shen, Lähde, Lee, Lu, Meißner, Nature Commun. 14 (2023) 2777 [arXiv:2202.13596[nucl-th]]

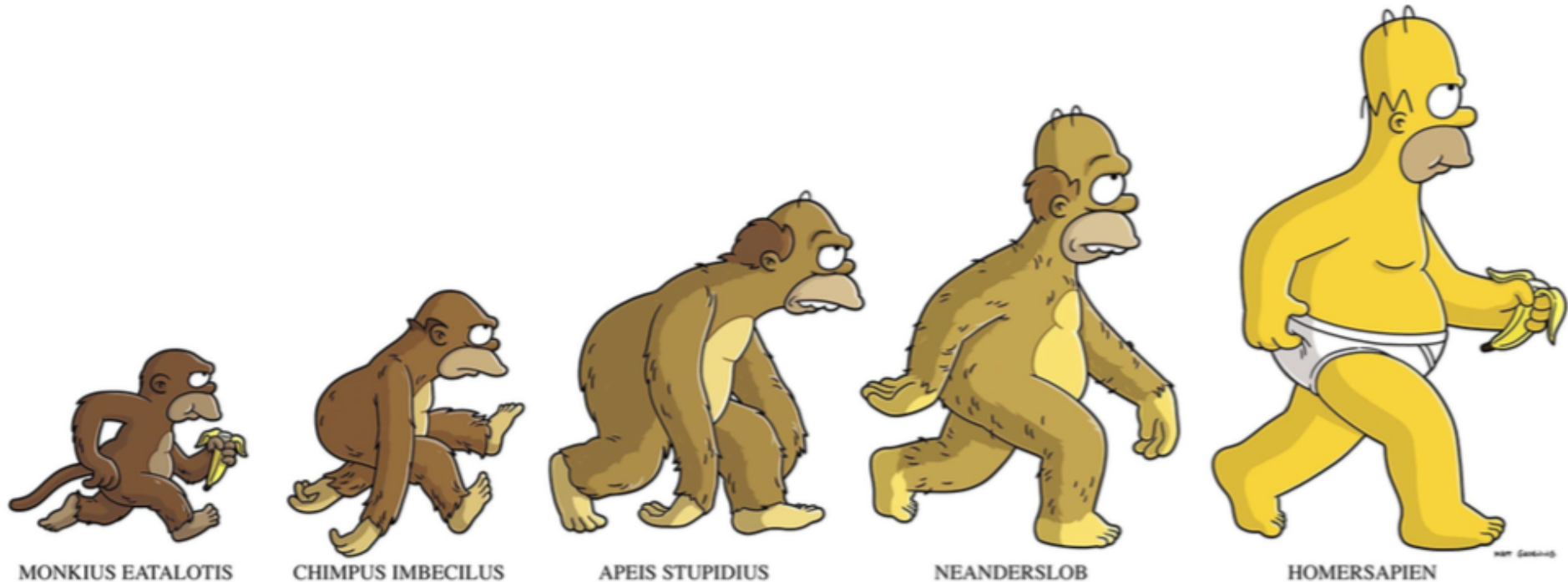
- Alpha cluster structure emerges naturally  
— not a built-in assumption



- After sufficiently large Euclidean projection time  
— memory of trial state erased

# EMERGENT ALPHA CLUSTERING

Shen, Lähde, Lee, Lu, Meißner, Nature Commun. 14 (2023) 2777 [arXiv:2202.13596[nucl-th]]



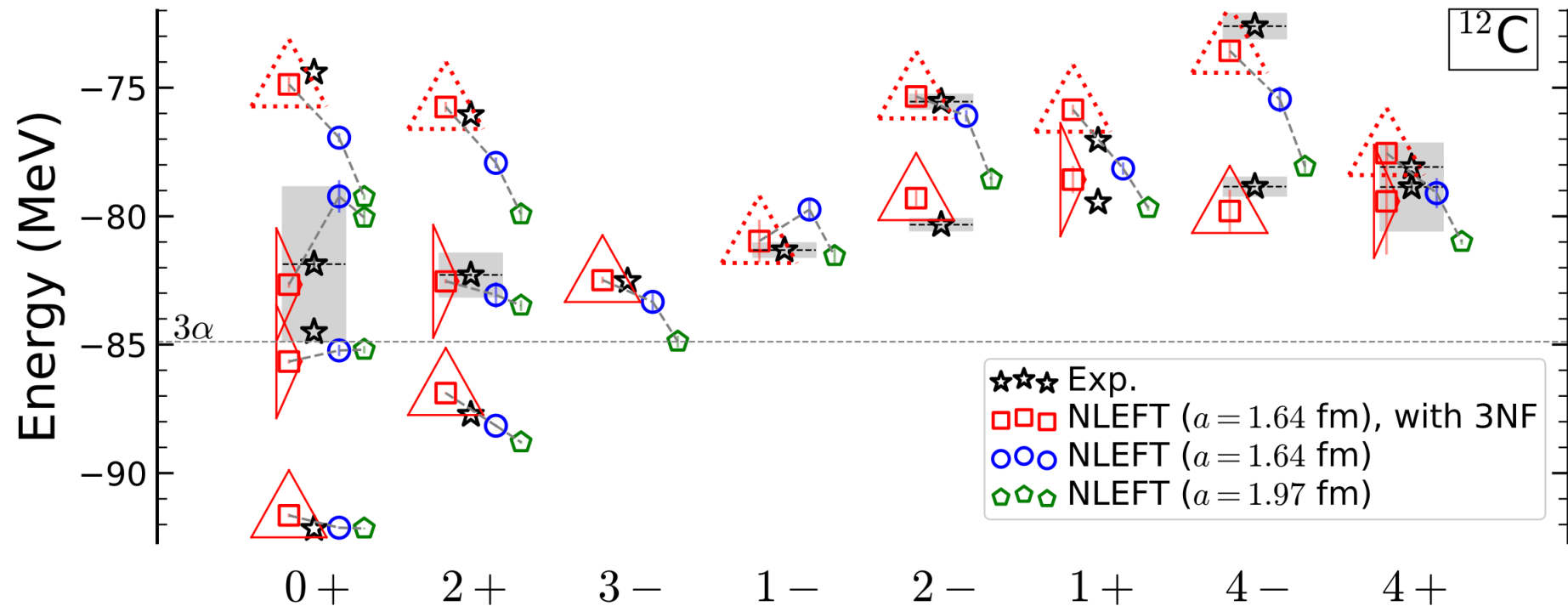
HOMERSAPIEN

Euclidean time projection

# CLUSTER/SHELL-MODEL DUALITY

Shen, Lähde, Lee, Lu, Meißner, Nature Commun. 14 (2023) 2777 [arXiv:2202.13596[nucl-th]]

- Lowest states are predominantly alpha cluster states with equilateral or obtuse (“bent-arm”) triangular arrangement



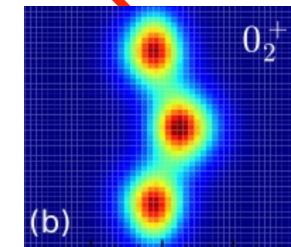
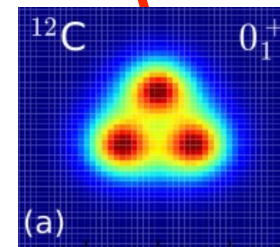
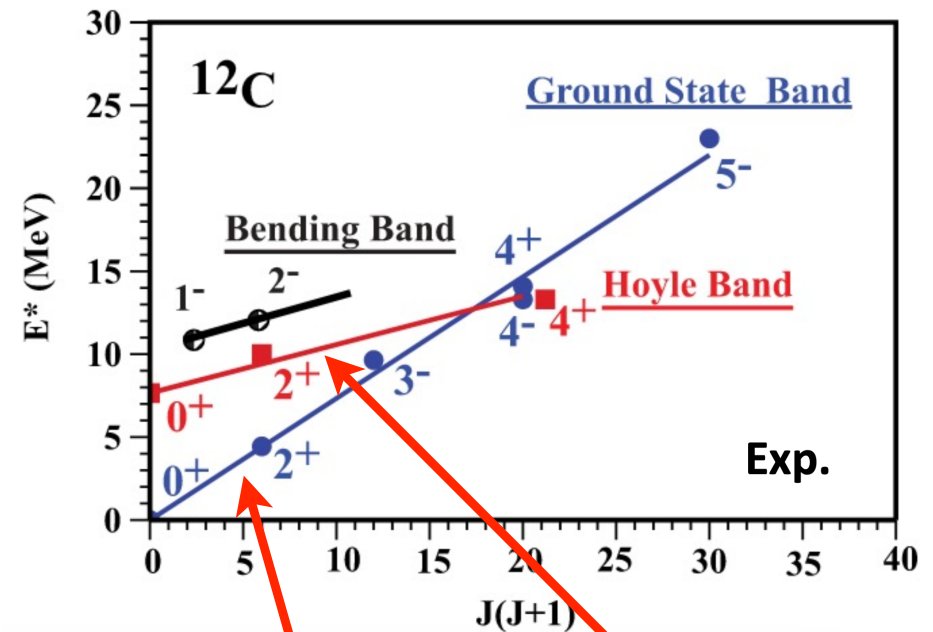
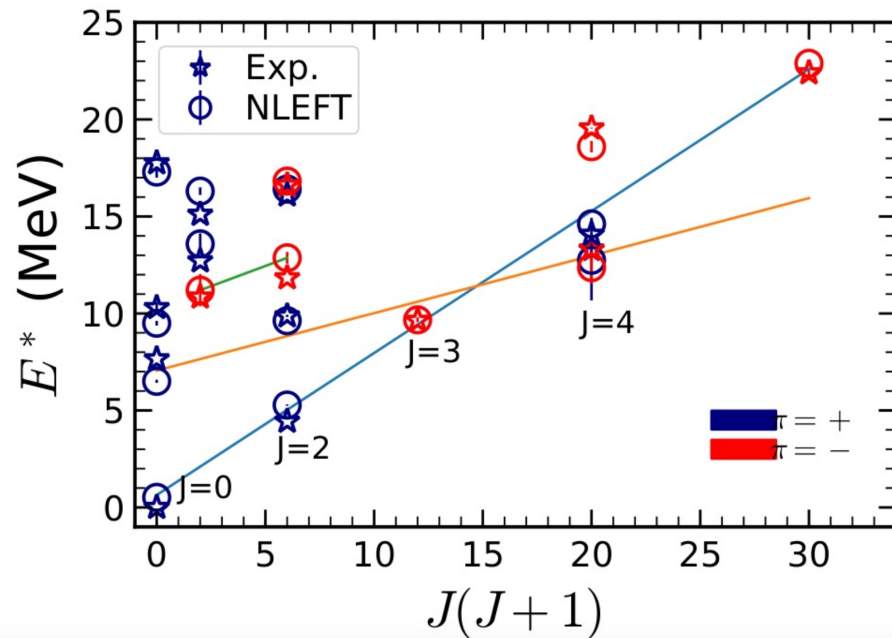
- Dashed triangles — strong 1p-1h component in the wave function, preference for shell-model trial states

# BAND STRUCTURE

Bijker, Iachello, PRC **61** (2000) 067305

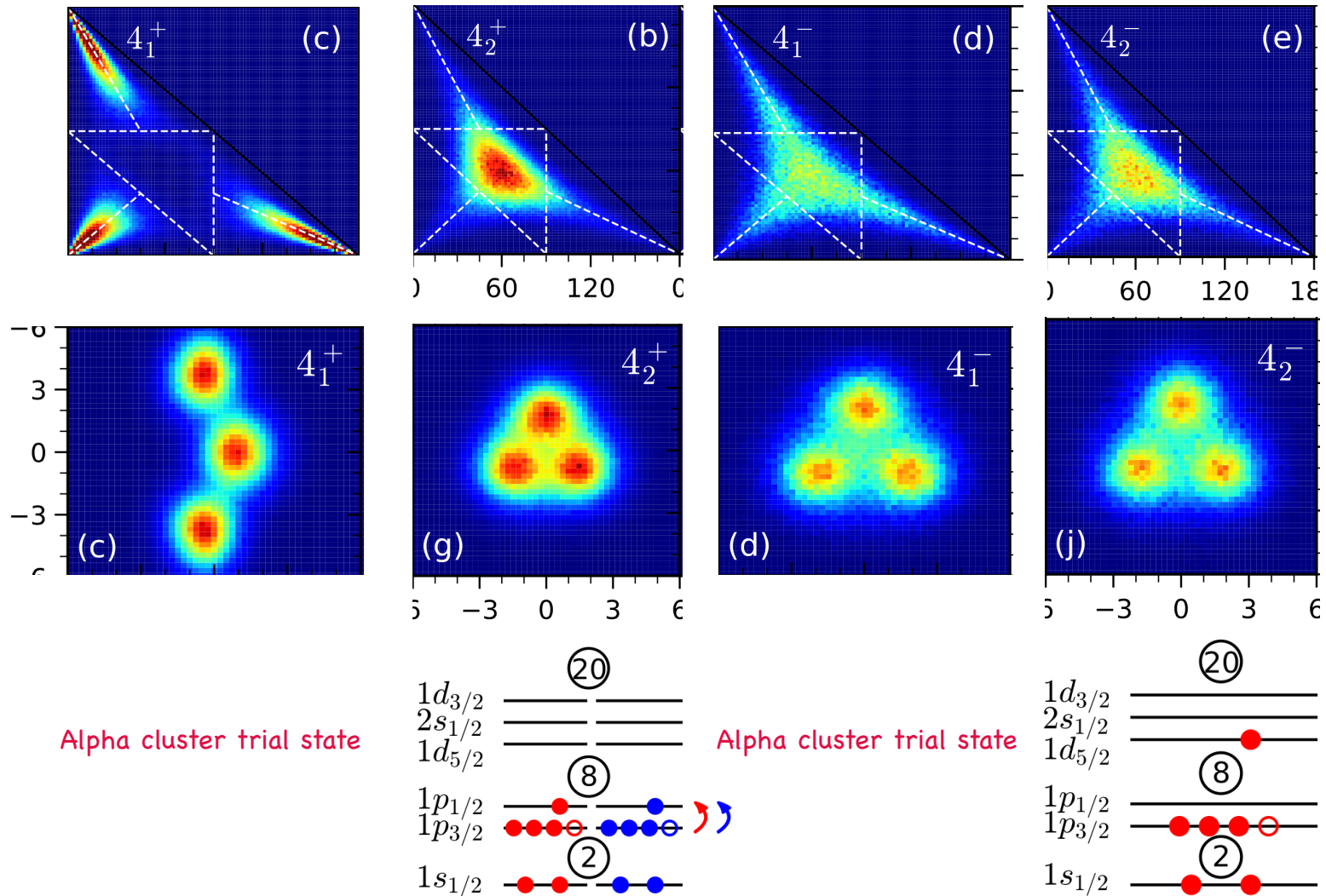
Marín-Lámbarri, Bijker, Freer, Gai et al., PRL **113** (2014) 012502

Shen, Lähde, Lee, Lu, Meißner, Nature Commun. **14** (2023) 2777 [arXiv:2202.13596[nucl-th]]



# HIGHER EXCITED STATES (4+, 4-)

Shen, Lähde, Lee, Lu, Meißner, Nature Commun. 14 (2023) 2777 [arXiv:2202.13596[nucl-th]]





# SUMMARY & OUTLOOK

- Successful description of low-lying  $^{12}\text{C}$  spectrum with Wigner SU(4) symmetric interaction (starting point for higher-order NLEFT)
- Successful characterization of the emergent alpha cluster geometry in the low-lying  $^{12}\text{C}$  states
- Now possible — NLEFT calculations to N3LO in chiral EFT (see also talk by Lukas Bovermann)  
*Elhatisari, Bovermann, Ma, Epelbaum et al., Nature 630 (2024) 59*
- Ab initio studies of Be isotopes from  $^7\text{Be}$  to  $^{12}\text{Be}$   
*Shen, Elhatisari, Lee, Meißner, Ren [arXiv:2411.14935 [nucl-th]]*
- Nucleon distribution and shell closure in  $^{22}\text{Si}$   
*Zhang, Elhatisari, Meißner, Shen [arXiv:2411.17462 [nucl-th]]*
- Studies of  $^{16}\text{O}$  and  $^{20}\text{Ne}$   
*Giacalone, Bally, Nijs, Shen et al., [arXiv:2402.05995 [nucl-th]]*  
*Giacalone, Zhao, Bally, Shen et al., [arXiv:2405.20210 [nucl-th]]*