

Clustering in nuclei at finite temperature

Elias Khan with M. Davies, J.-P. Ebran, F. Mercier, P. Stevenson, E. Yüksel,



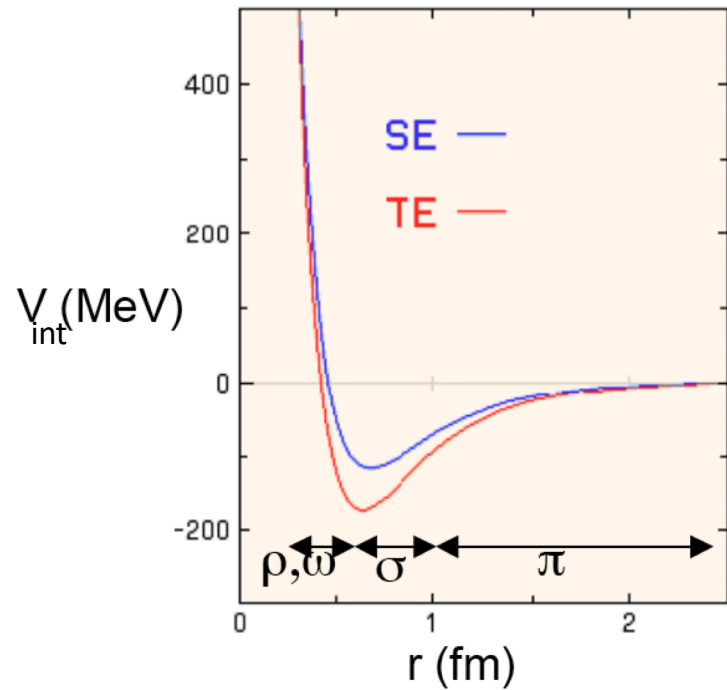
Questions addressed

What is the impact of temperature on nuclear clusterization ?

What is the (ρ, T) diagram for clusterization ?

Method: covariant EDF

V and S potentials



$$\left\{ p \frac{1}{2\tilde{M}(r)} p + W(r) + V_{ls}(r) l \cdot s \right\} \varphi_i = \epsilon_i^{NR} \varphi_i$$

$$S \approx -400 \text{ MeV}$$

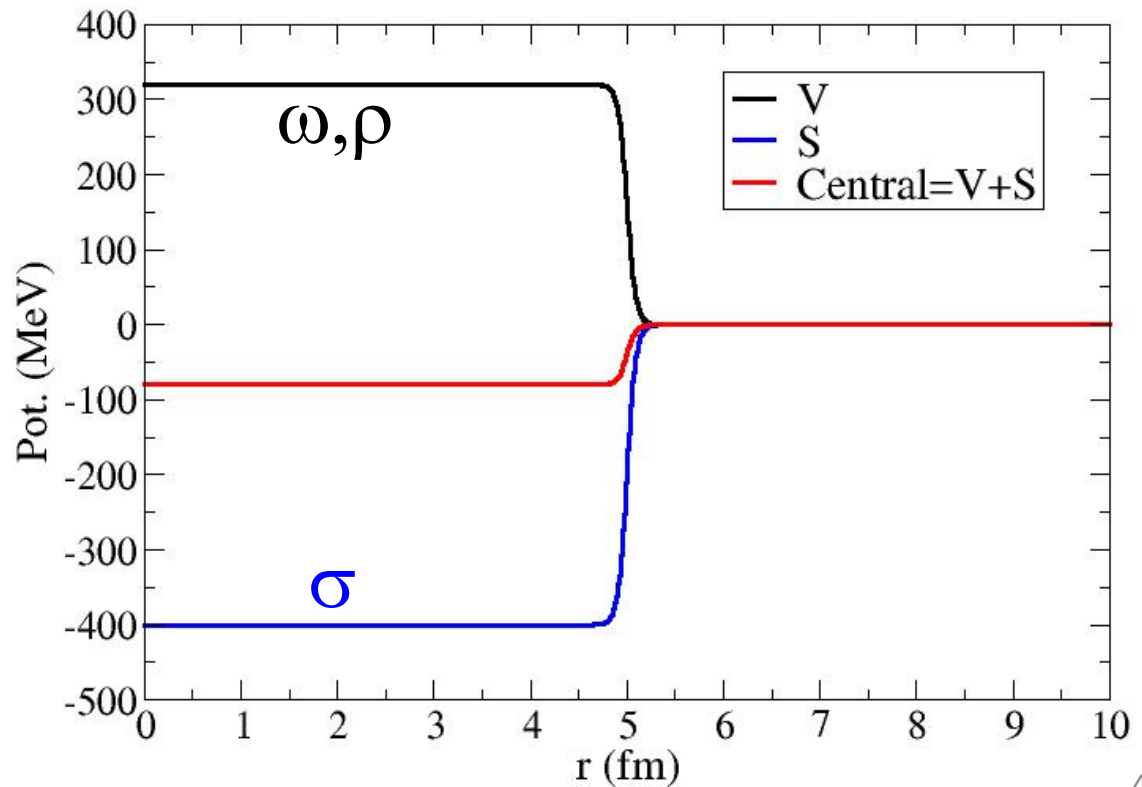
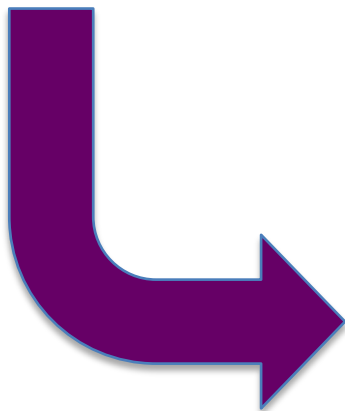
$$V \approx 320 \text{ MeV}$$



$$V_0 \approx 80 \text{ MeV}$$

$$W(r) = [V + S](r)$$

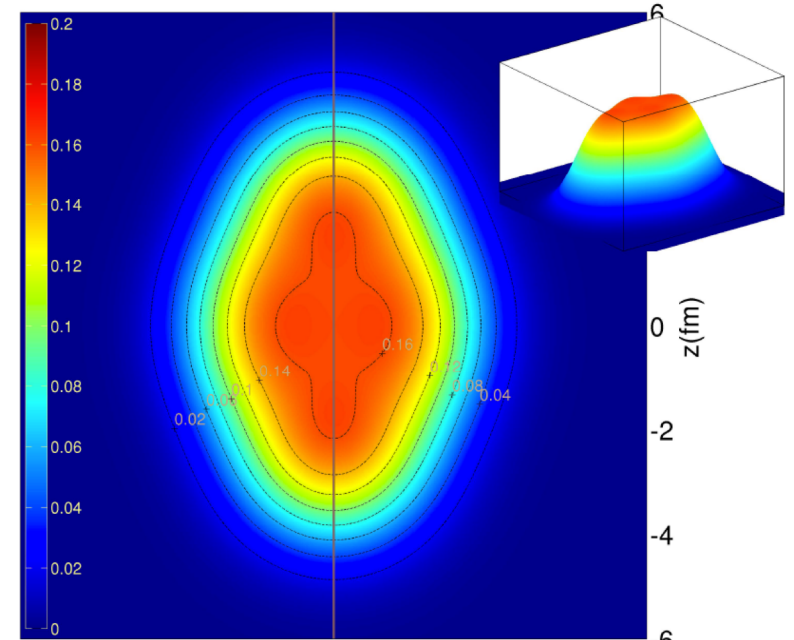
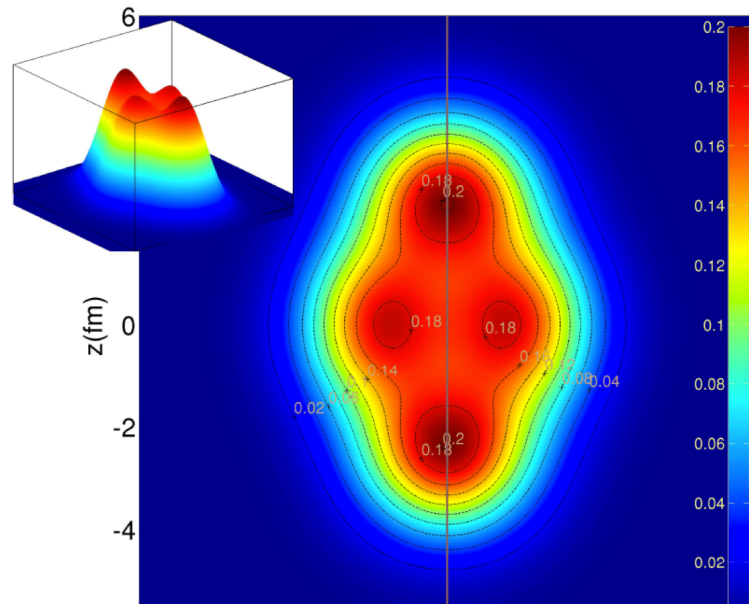
$$V_{ls}(r) = \frac{1}{2\tilde{M}^2(r)} \frac{1}{r} \frac{d}{dr} (V - S)$$



Nucleonic density: covariant EDF clusterizes

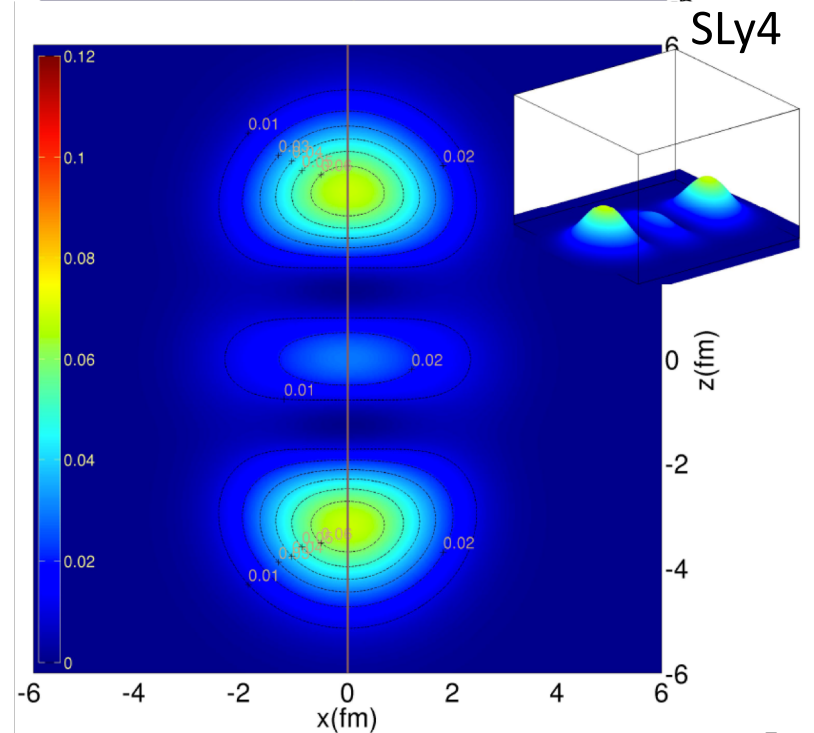
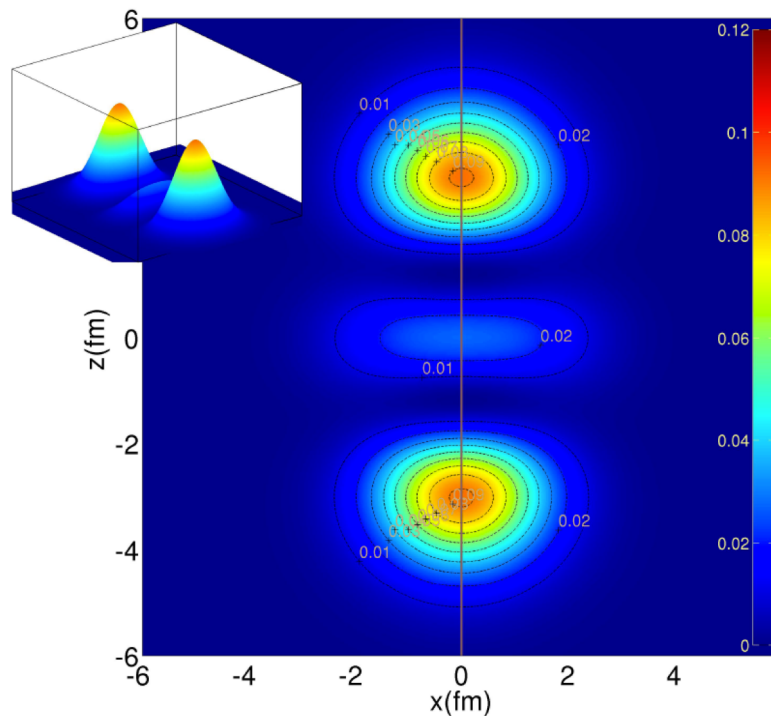
^{20}Ne

Total density

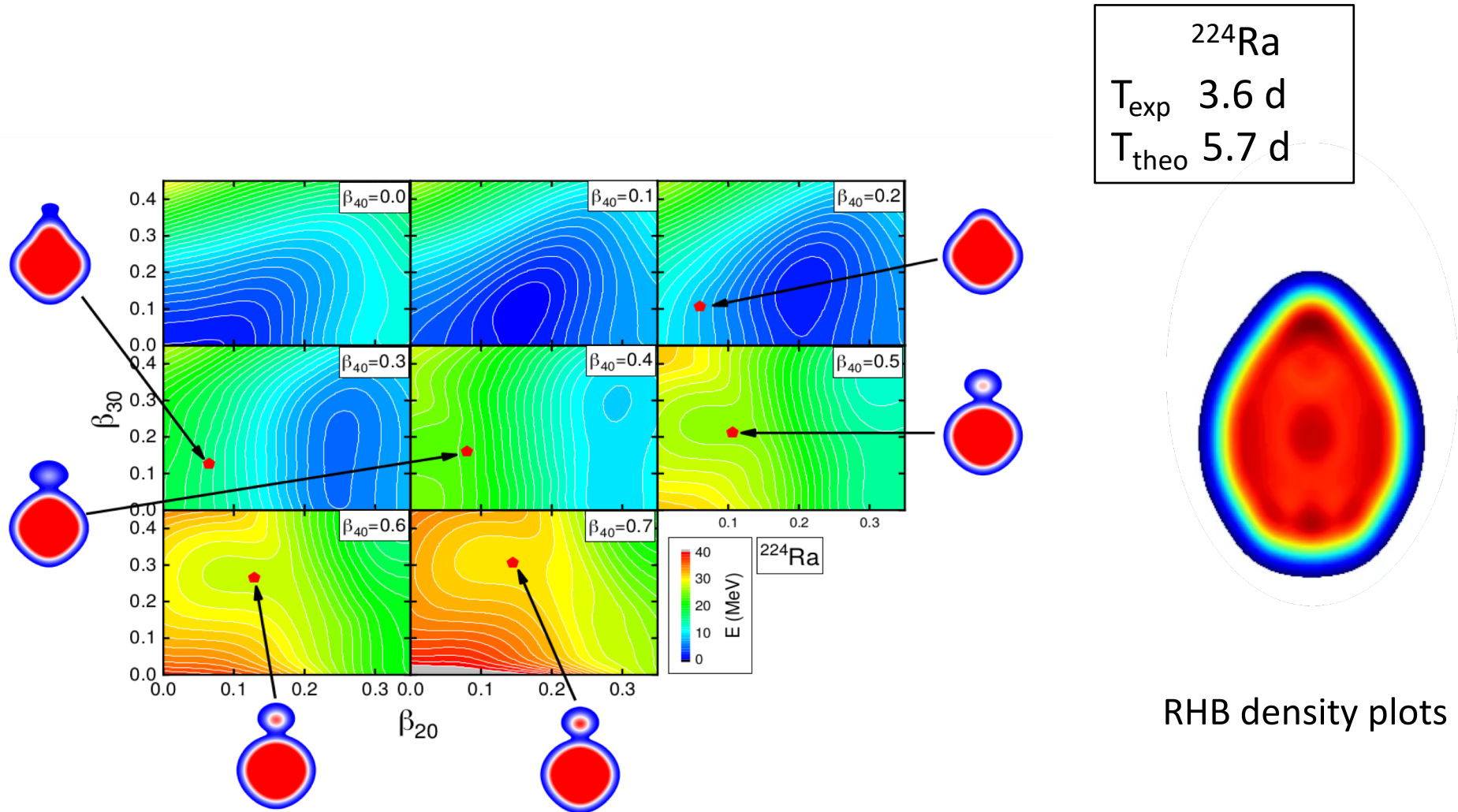


DDME2

Last occupied state

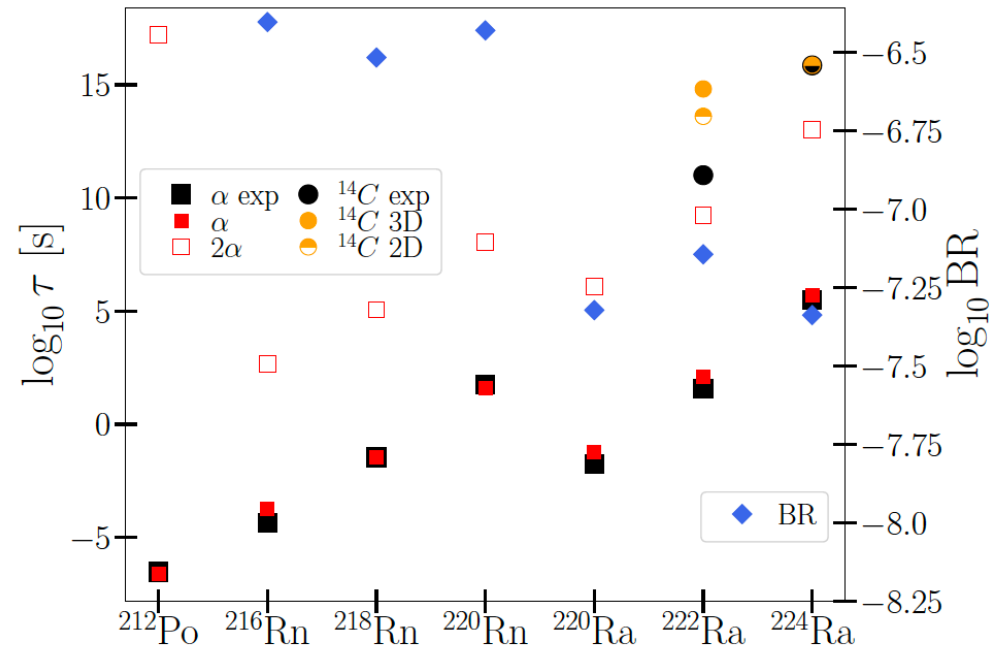


α decay in ^{224}Ra



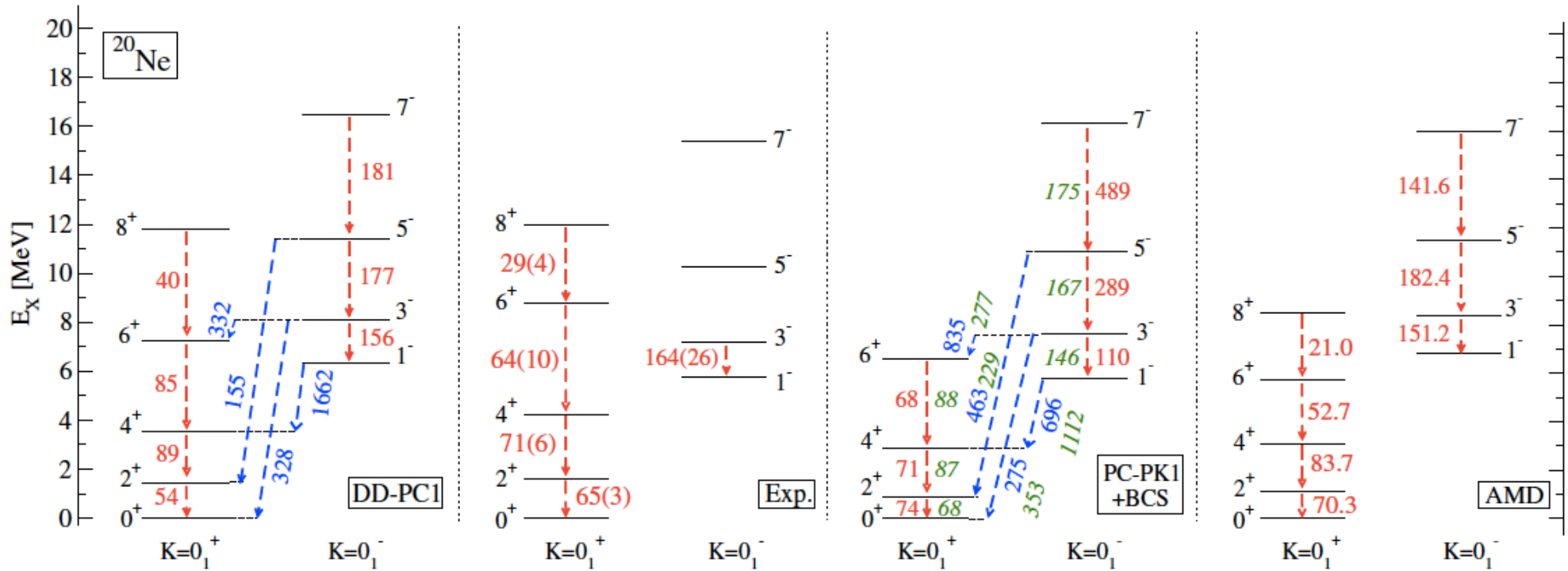
For heavy nuclei, the hexadecapole constraint β_4 is necessary to form the neck for alpha-particle emission

Relevant constraints



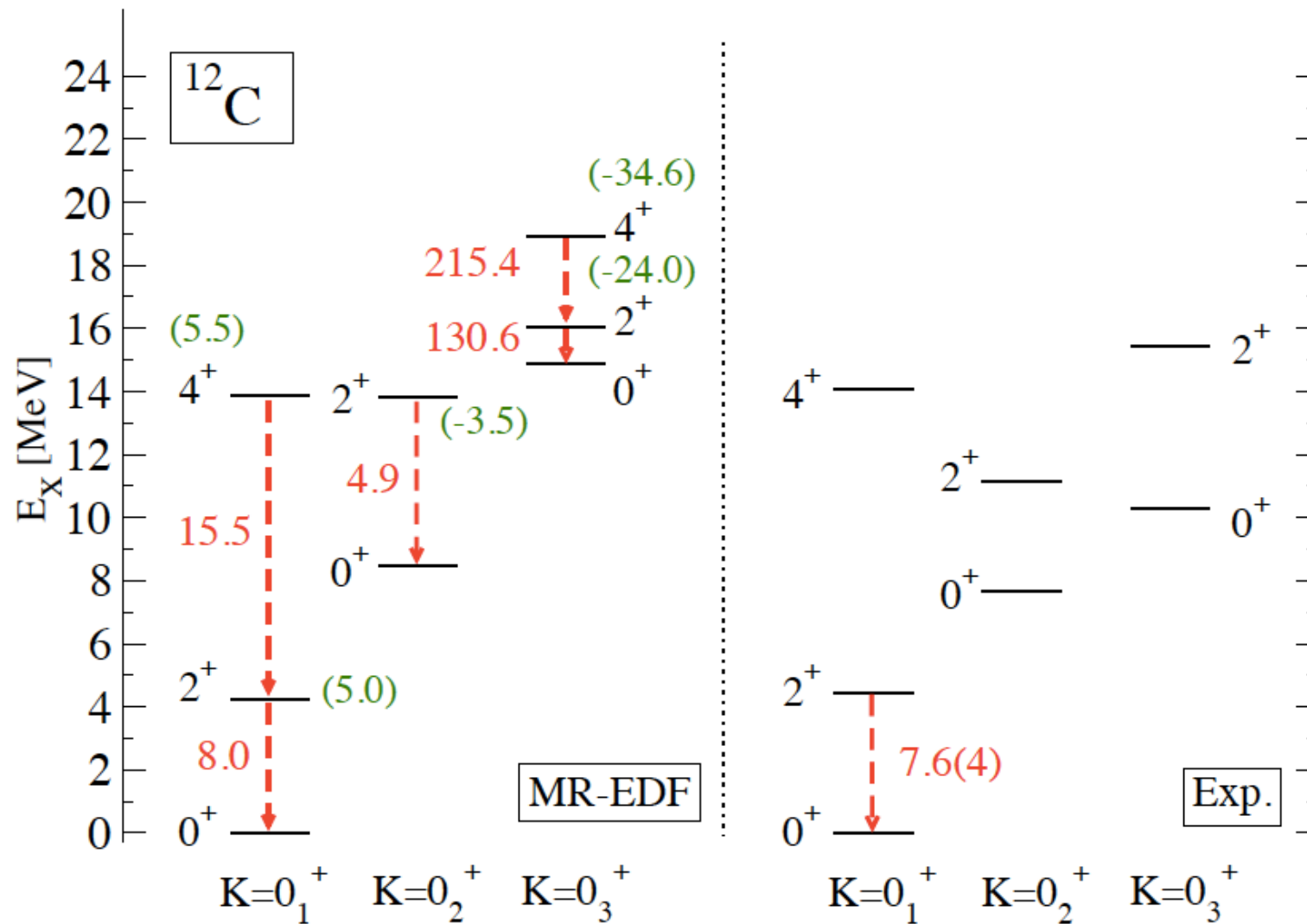
	Fission	Cluster	Alpha	Double alpha
β_{20}	×	×	×	×
β_{30}	×	×	×	
β_{40}			×	×

Comparison with the exp. spectrum on ^{20}Ne

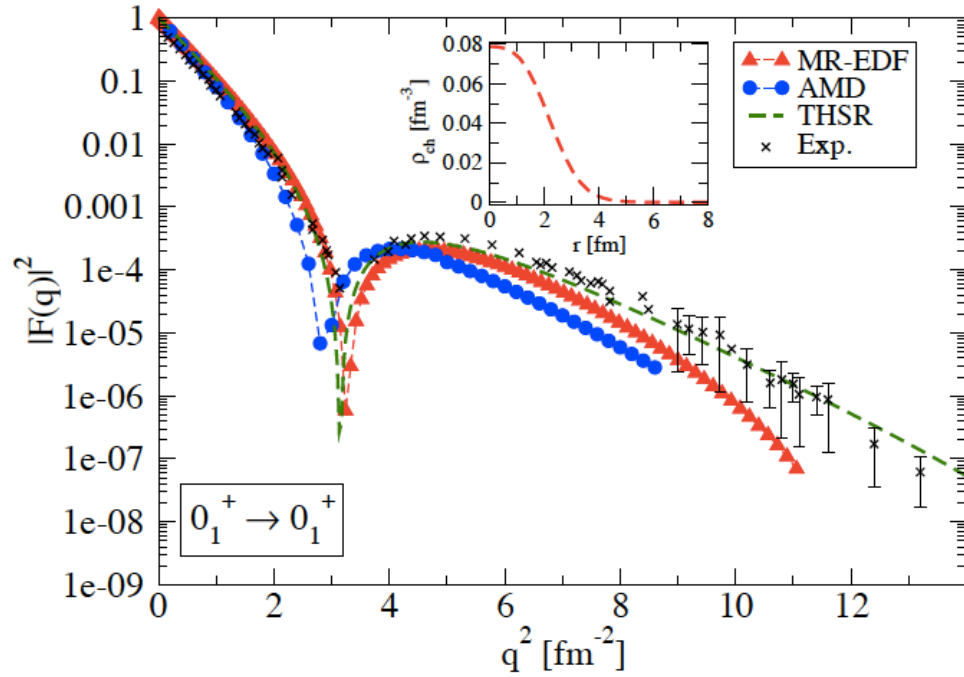


- GCM on top of axially symmetric /reflection asymmetric RHB (DD-PC1) :
- Angular momentum, parity and particle number projections

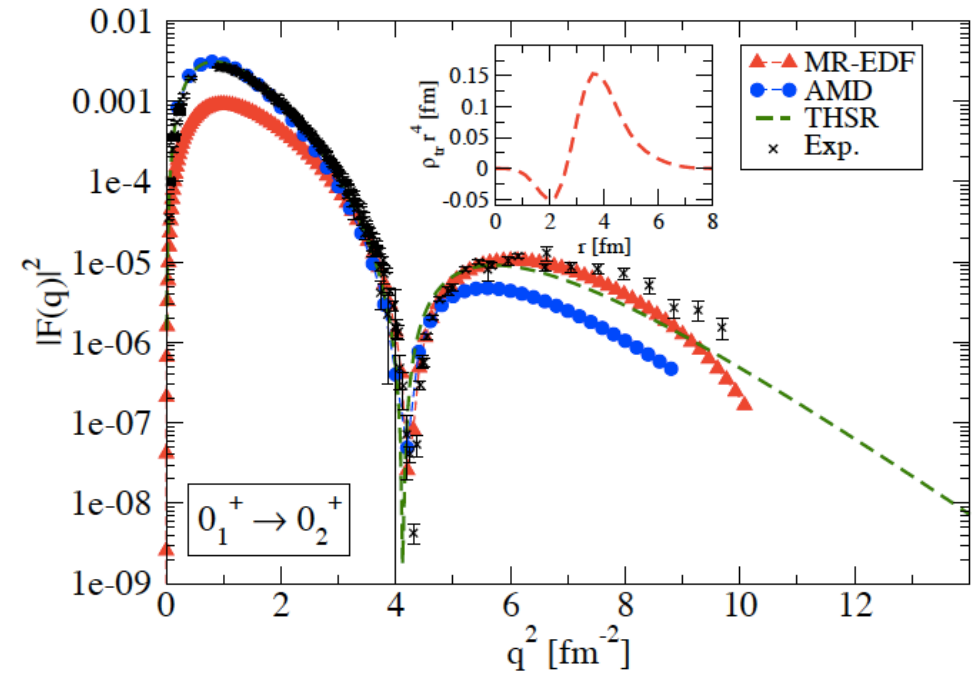
Comparison with exp. spectrum on ^{12}C



Comparison with exp. form factors on ^{12}C

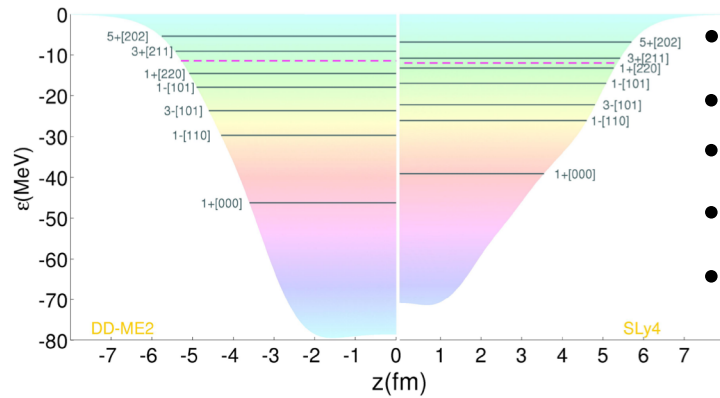


g.s.



Hoyle

Origins of nuclear clustering: control parameters of a QPT



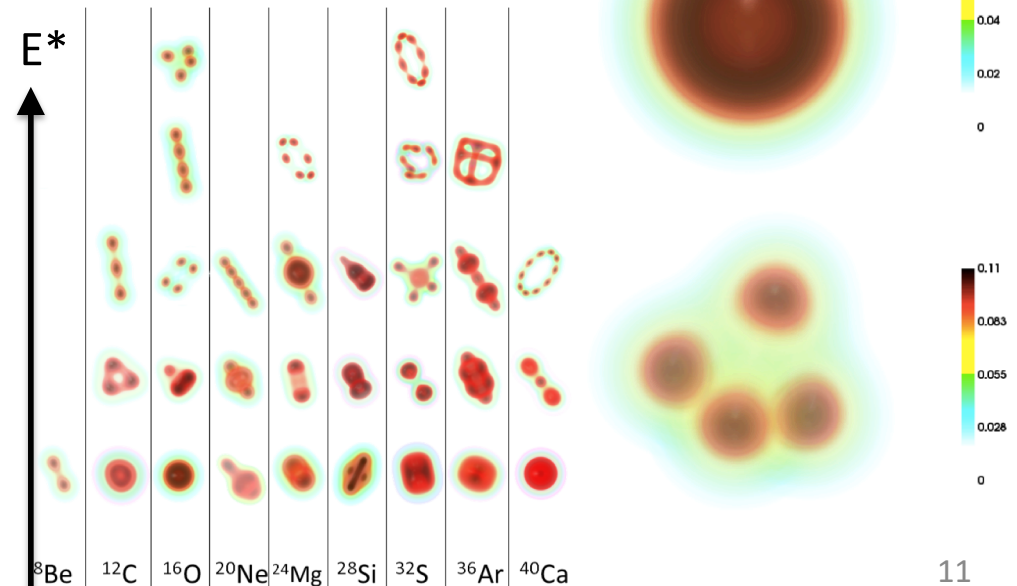
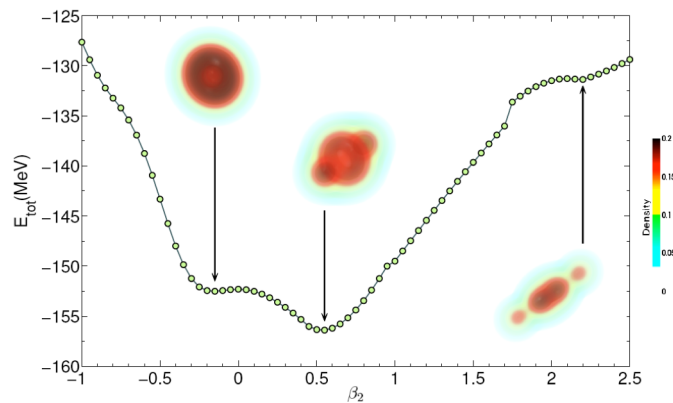
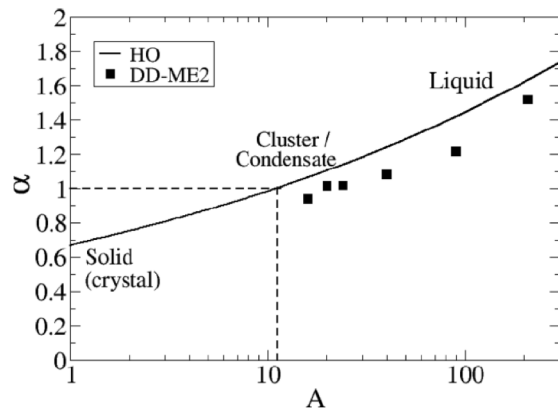
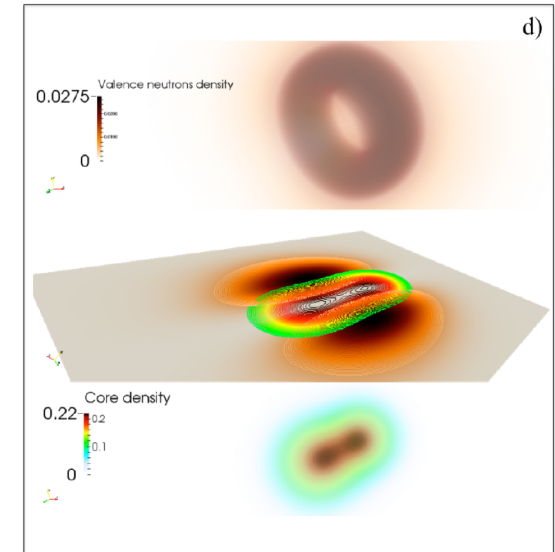
- Depth of the confining potential
- Heavy vs. Light nuclei
- Deformation / excitation energy
- Density
- Neutron excess

J.-P. Ebran, E. Khan, T. Niksic, D. Vretenar, Nature 487(2012)341

PRC87(2013)044307

PRC90(2014)054329

PRC89(2014)031303(R)

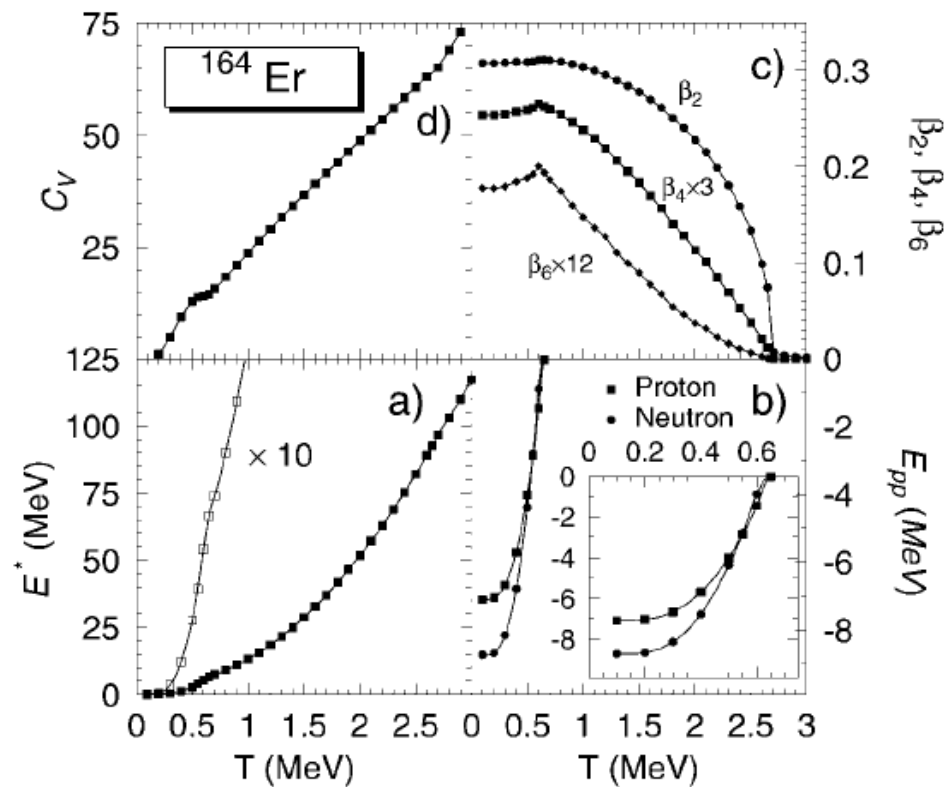


What is the impact of temperature on nuclear clustering ?

Effect of temperature on nuclei

Increasing the temperature :

- Critical temperature for pairing collapse
- Critical temperature for deformation collapse



J.L. Egidio, L.M. Robledo, V. Martin, PRL 85, 27 (2000)

What happens for temperature effect on clustering: critical temperature ?

Finite-temperature-deformed-RHB

$$\Omega = E - TS - \lambda N$$

$$\begin{pmatrix} h_D - \lambda - m & \Delta \\ -\Delta^* & -h_D^* + \lambda + m \end{pmatrix} \begin{pmatrix} U_k \\ V_k \end{pmatrix} = E_k \begin{pmatrix} U_k \\ V_k \end{pmatrix}$$

$$\rho_s = \sum_{E_k > 0} V_k^\dagger \gamma^0 (1 - f_k) V_k + U_k^T \gamma^0 f_k U_k^*$$

$$\rho_v = \sum_{E_k > 0} V_k^\dagger (1 - f_k) V_k + U_k^T f_k U_k^*,$$

$$\rho_{tv} = \sum_{E_k > 0} V_k^\dagger \tau_3 (1 - f_k) V_k + U_k^T \tau_3 f_k U_k^*,$$

$$f_k = \frac{1}{1 + e^{\beta E_k}}$$

Taking into account statistical fluctuations

- Quantum fluctuations: not considered here
- Statistical fluctuations:

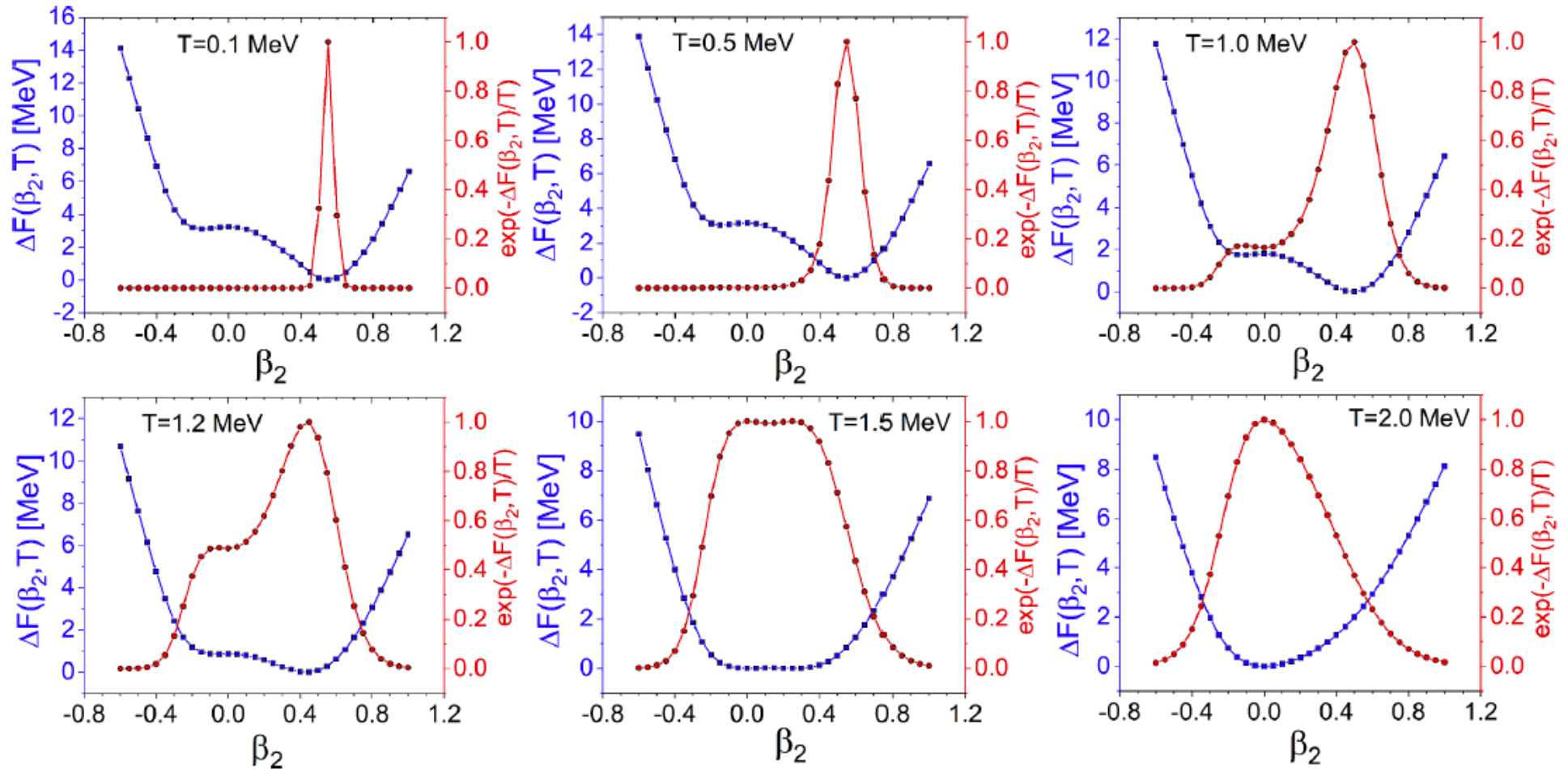
$$\bar{O} = \frac{\int d\beta_2 O(\beta_2, T) \exp(-\Delta F(\beta_2, T)/T)}{\int d\beta_2 \exp(-\Delta F(\beta_2, T)/T)}$$

with $\Delta F(\beta_2, T) = F(\beta_2, T) - F_{\min}(T)$

$F(\beta_2, T) = E(\beta_2, T) - TS(\beta_2, T)$: free energy

$S = -k_B \sum_k [f_k \ln f_k + (1 - f_k) \ln(1 - f_k)]$: entropy

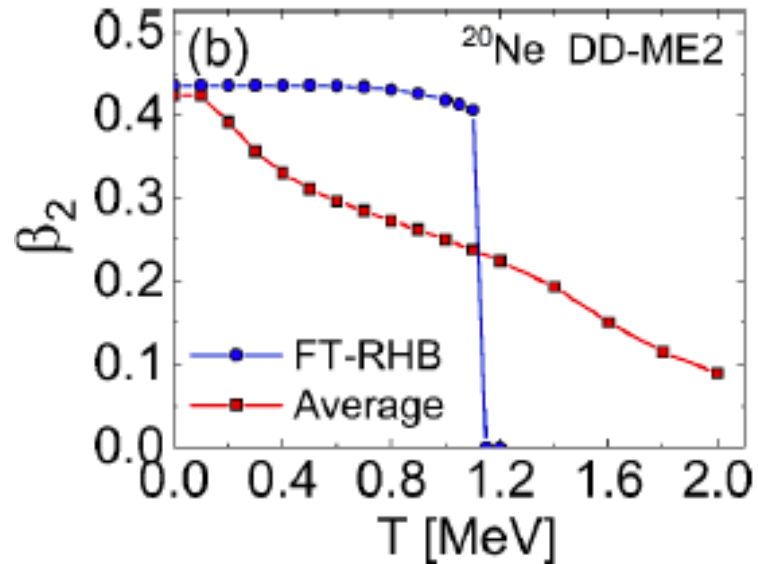
Taking into account statistical fluctuations



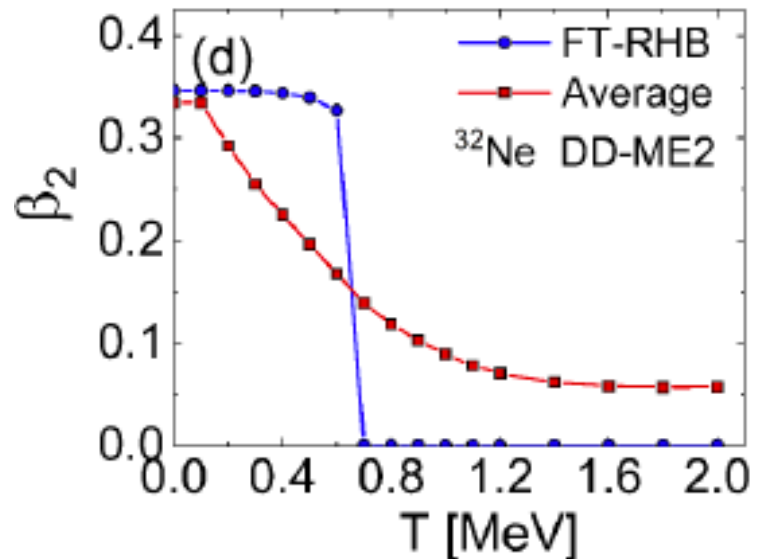
^{20}Ne

Mixing of the states increases with temperature

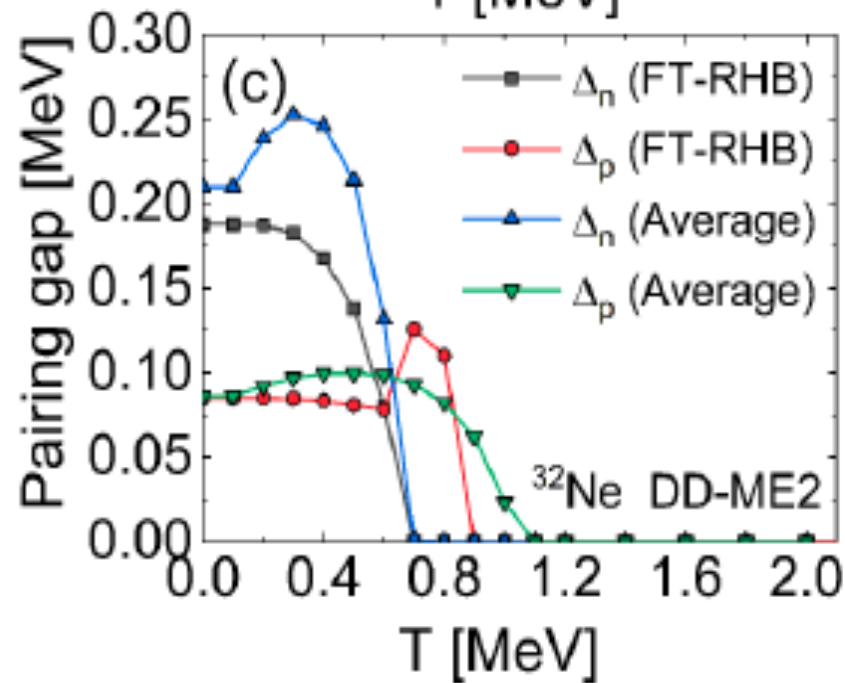
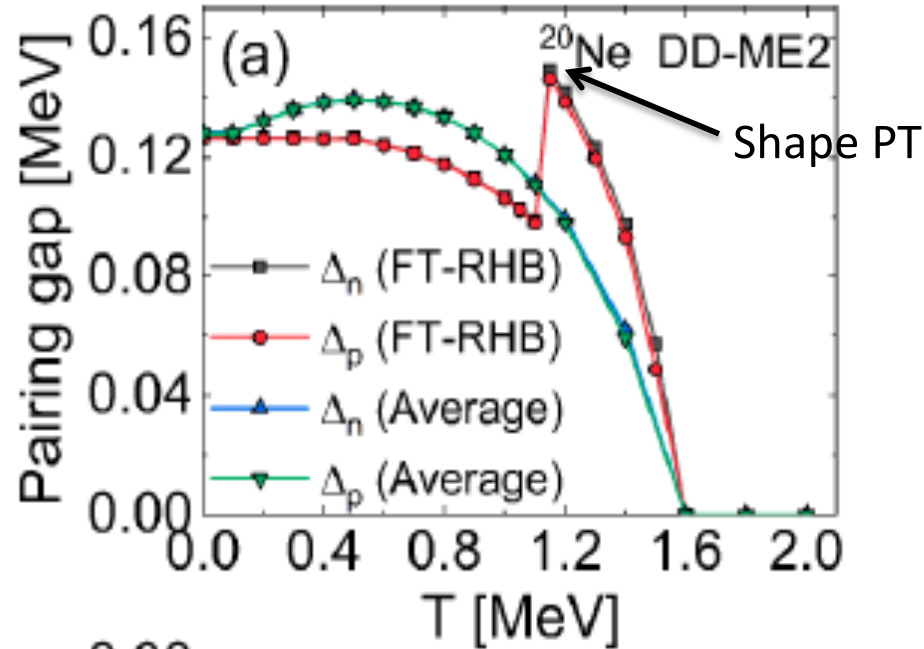
Shape phase transition



Smoothing of the transition due to the mixing

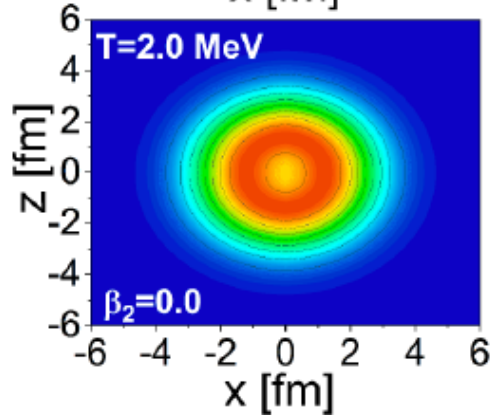
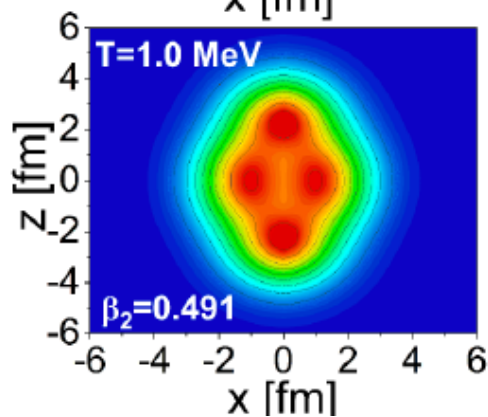
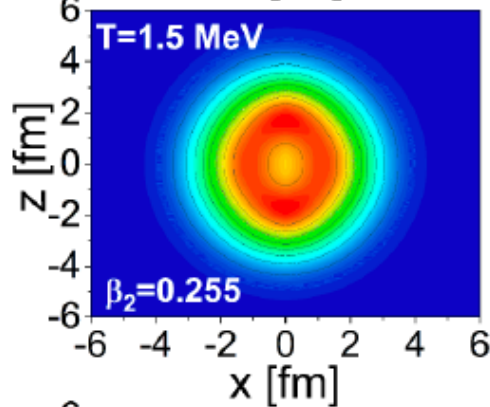
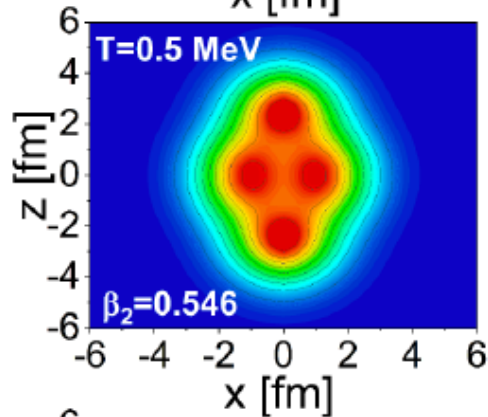
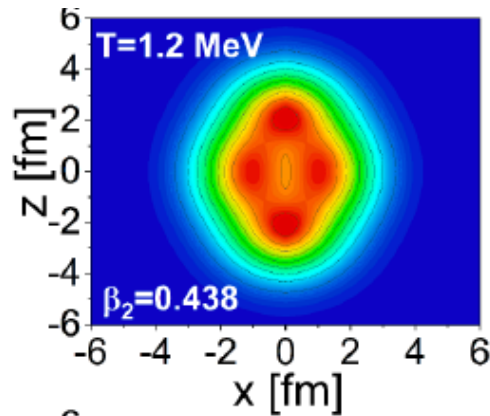
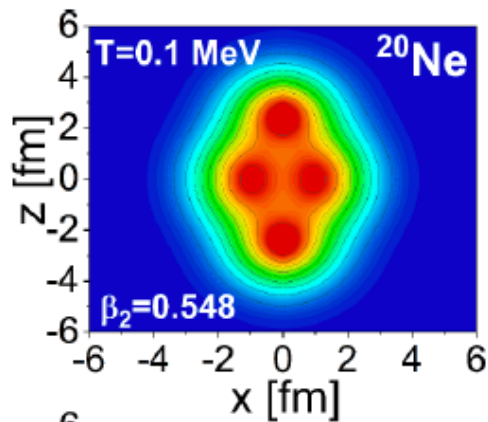


Pairing phase transition



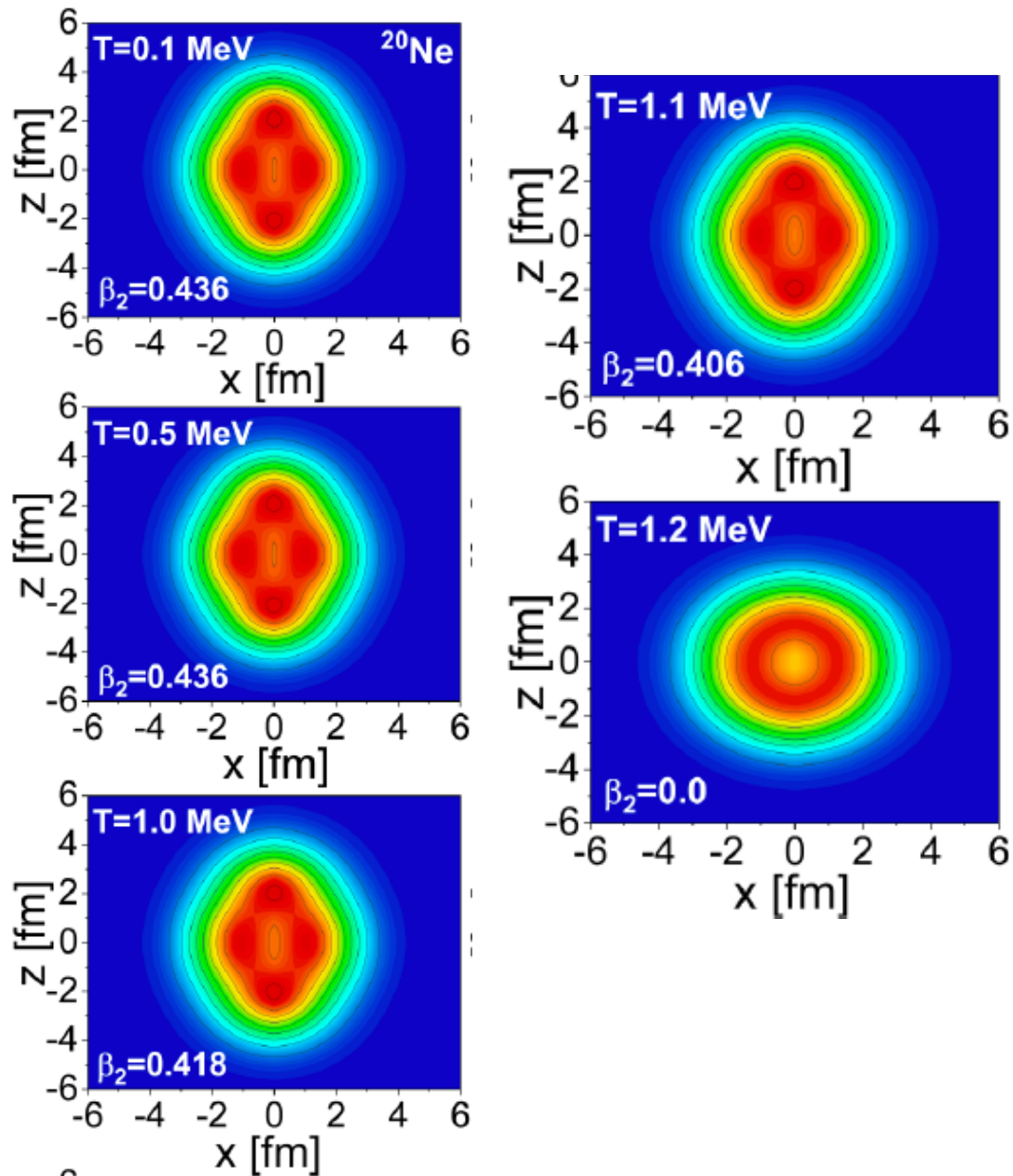
Effect of mixing

Temperature effect on clusterization in nuclei



Weak pairing: gradual vanishing of clusterization

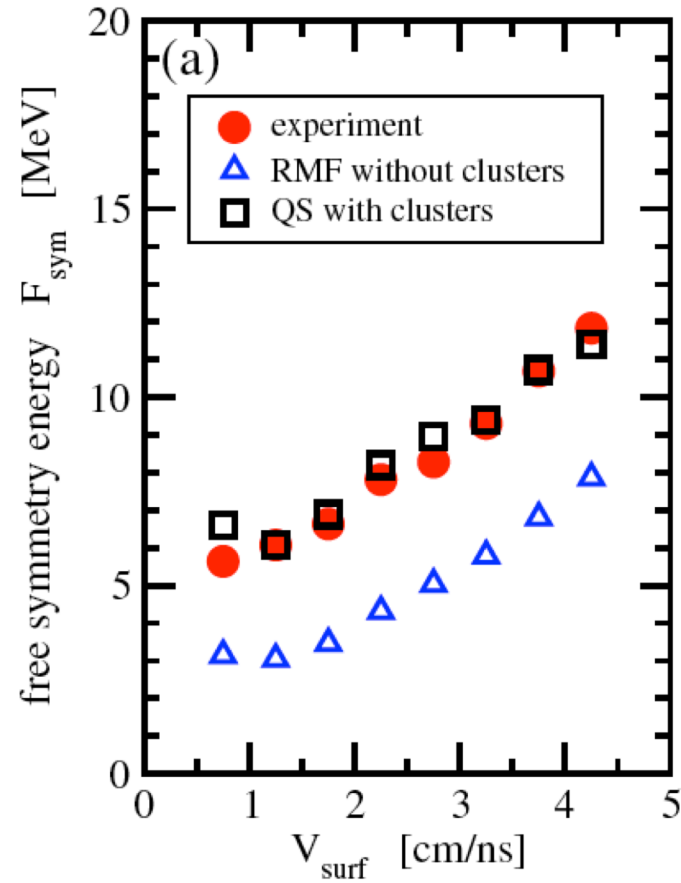
Temperature effect on clusterization in nuclei



Stronger pairing: pairing synchronizes the cluster and shape phase transitions

On the way to a (ρ, T) diagram:
what is the impact of density on nuclear clustering ?

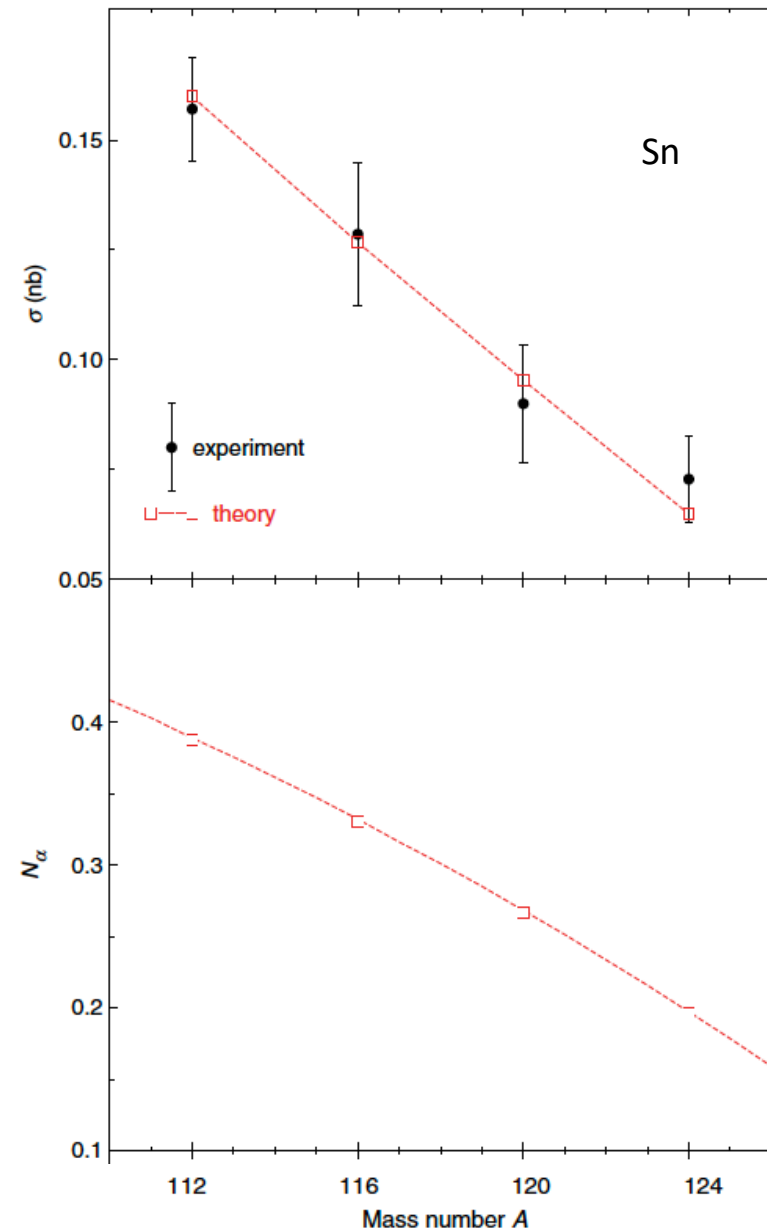
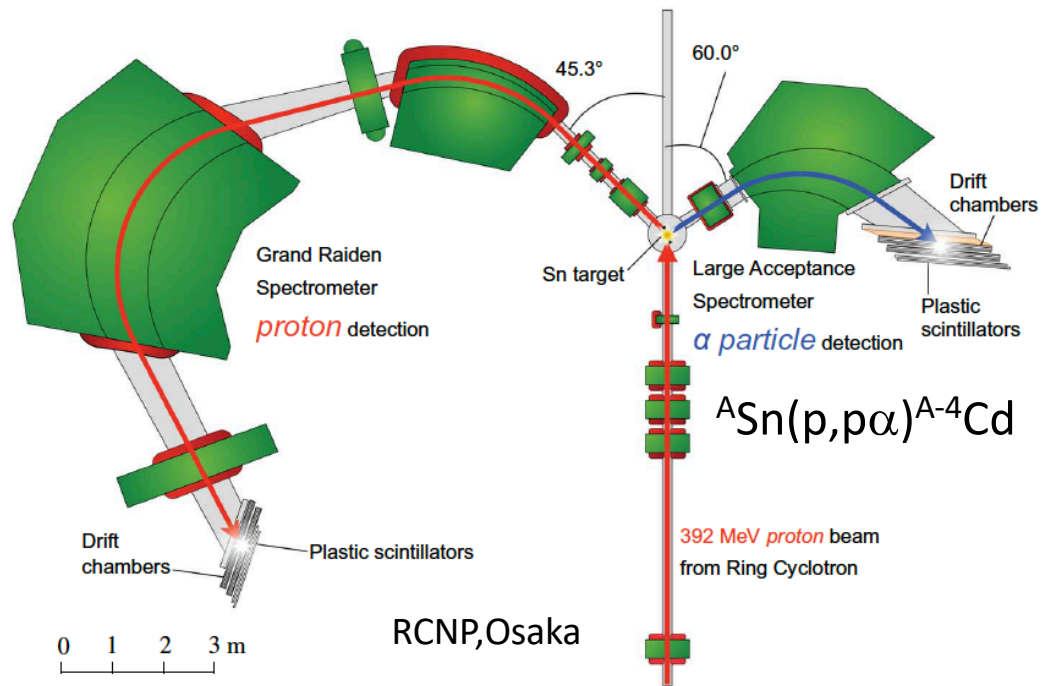
Clusters in low density nuclear matter



J.B. Natowicz et al. PRL104(2010)202501

Clusters in EoS better describe experiment
Data from heavy ion collision

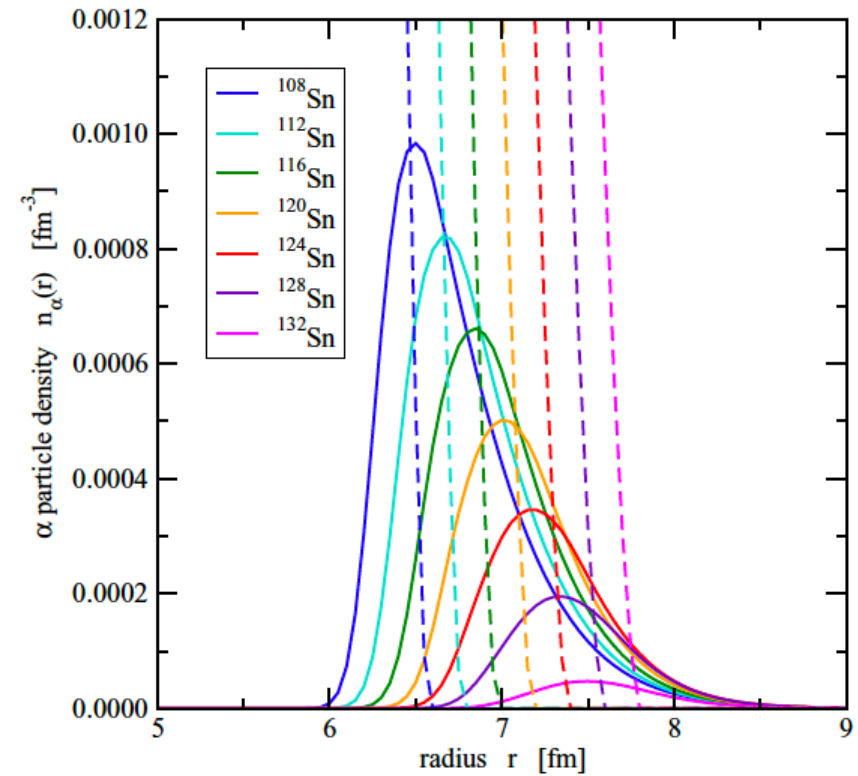
Clusters in low-density part of nuclei ?



Interpretation

$$\begin{aligned}
 \mathcal{L} = & \sum_{i=n,p,t,h} \bar{\psi}_i (\gamma_\mu i D_i^\mu - M_i) \psi_i + \frac{1}{2} (i D_\alpha^\mu \varphi_\alpha)^* \\
 & \times (i D_{\alpha\mu} \varphi_\alpha) - \frac{1}{2} \varphi_\alpha^* M_\alpha^2 \varphi_\alpha + \frac{1}{4} (i D_d^\mu \varphi_d - i D_d^\nu \varphi_d^\mu)^* \\
 & \times (i D_{d\mu} \varphi_{d\nu} - i D_{d\nu} \varphi_{d\mu}) - \frac{1}{2} \varphi_d^{\mu*} M_d^2 \varphi_{d\mu} \\
 & + \frac{1}{2} (\partial^\mu \sigma \partial_\mu \sigma - m_\sigma^2 \sigma^2 - \frac{1}{2} G^{\mu\nu} G_{\mu\nu} + m_\omega^2 \omega^\mu \omega_\mu \\
 & - \frac{1}{2} \vec{H}^{\mu\nu} \cdot \vec{H}_{\mu\nu} + m_\rho^2 \vec{\rho}^\mu \cdot \vec{\rho}_\mu), \quad (3)
 \end{aligned}$$

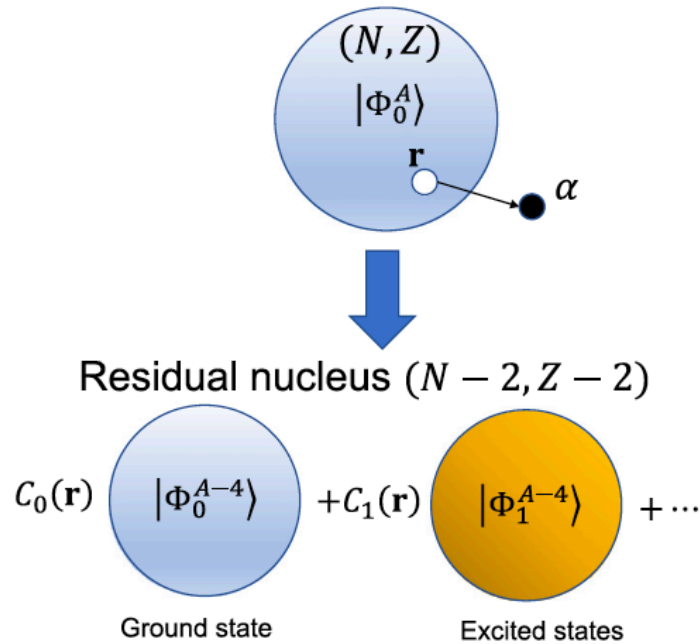
Explicit inclusion of t, ³He, α, d in the Lagrangian



α density appears in the surface of nuclei
(low density-Mott effect)

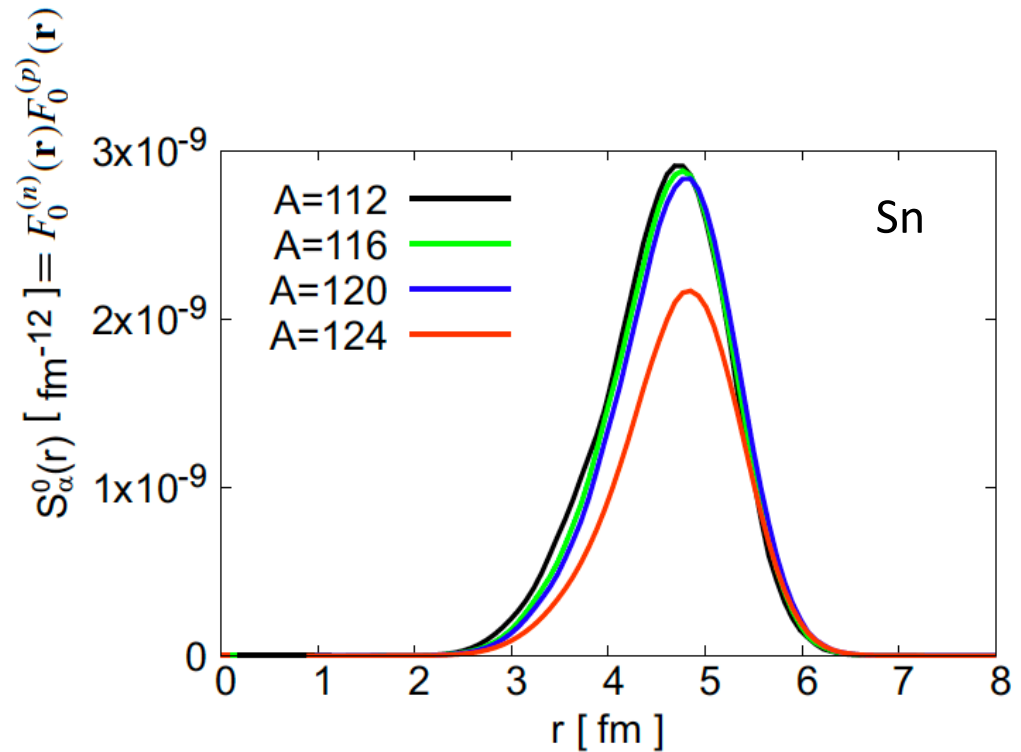
S. Typel, PRC 89(2014) 064321

The local α removal strength



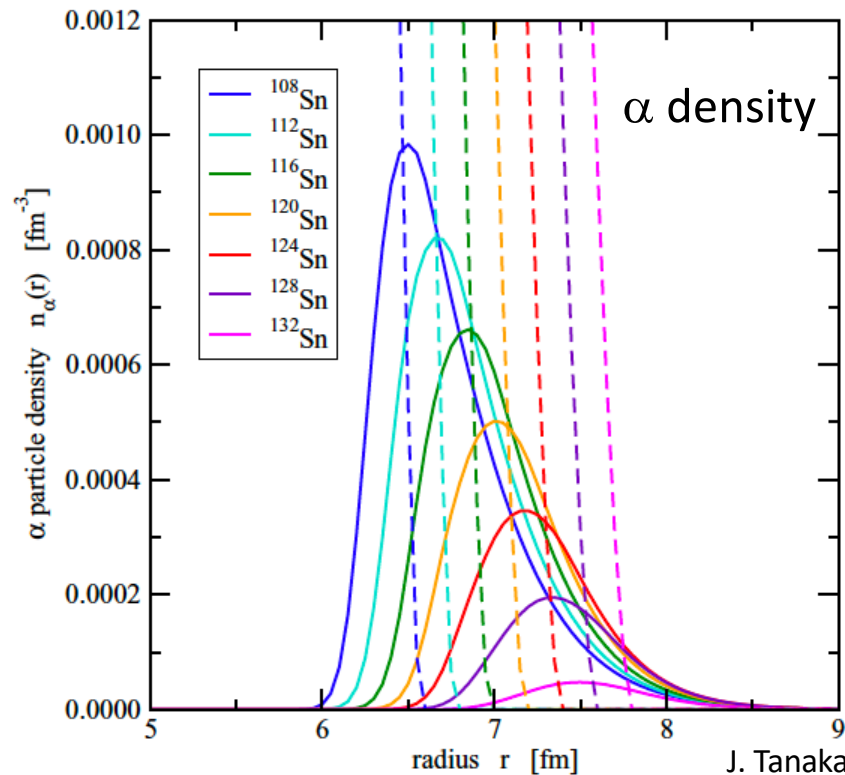
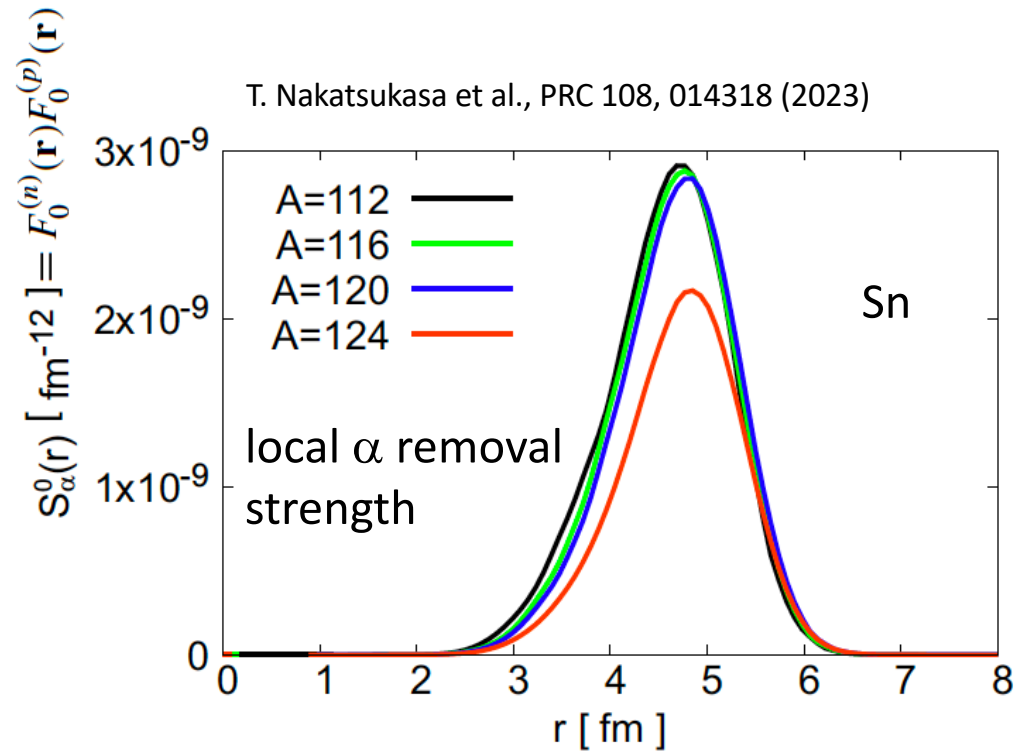
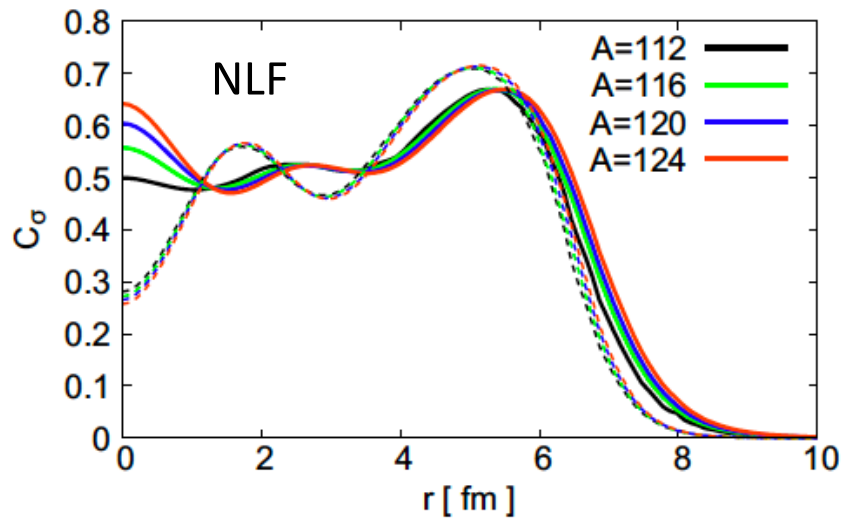
$$F_k^{(q)}(\mathbf{r}) = \left| \langle \Phi_k^{N_q-2} | \hat{\psi}_\uparrow^{(q)}(\mathbf{r}) \hat{\psi}_\downarrow^{(q)}(\mathbf{r}) | \Phi_0^{N_q} \rangle \right|^2$$

towards the g.s.: $F_0(\mathbf{r}) = |\kappa(\mathbf{r})|^2$
 If no pairing: removal of 2 particles



Peaked at the surface

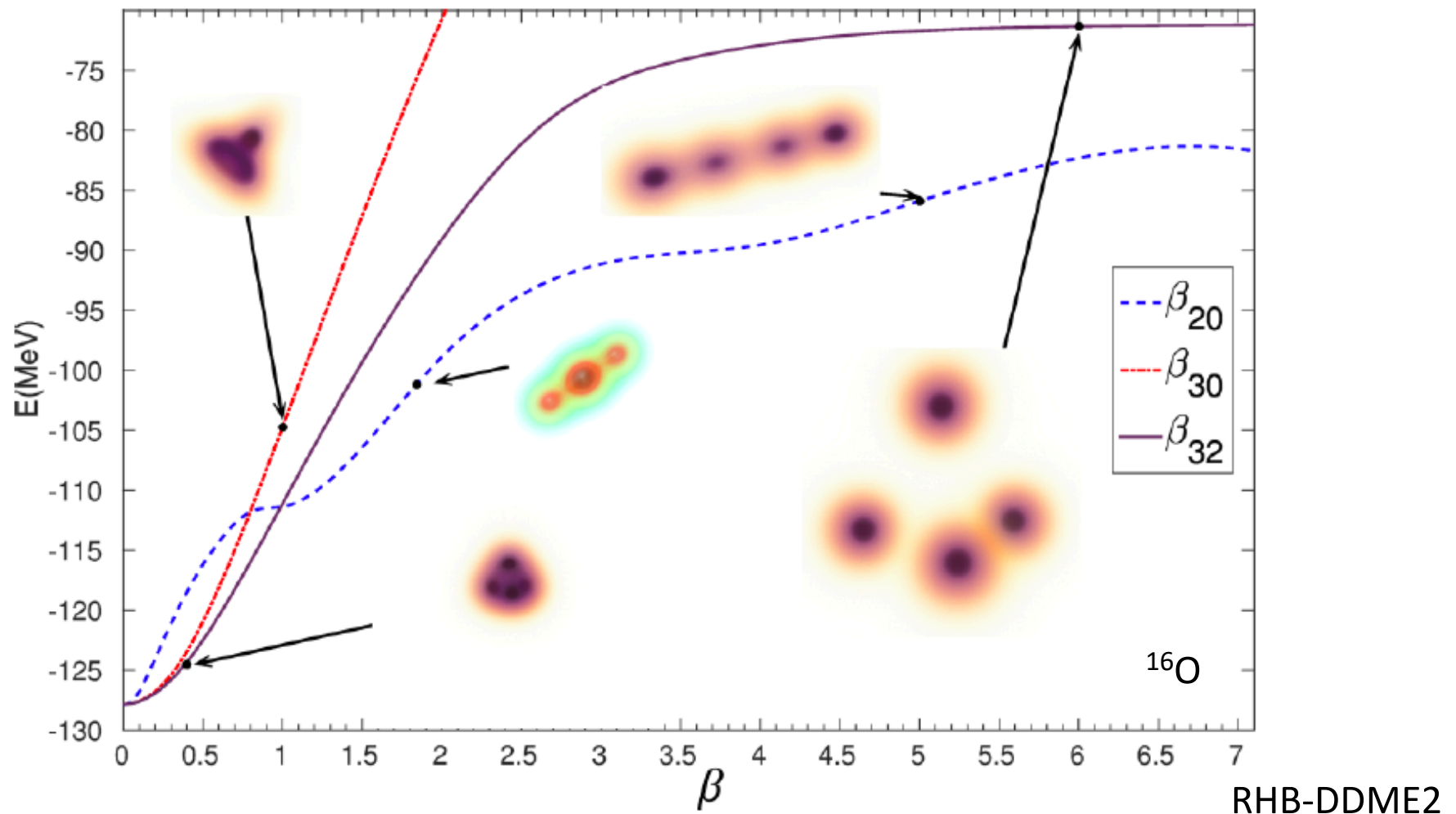
Comparison with the NLF and the α density



- NLF has a different profile inside the nuclei
- α density peaked at larger r values ($\rho_0/100$)
- See a general discussion in

E. Khan, L. Heitz, F. Mercier, J.-P. Ebran, PRC 106 (2022) 064330

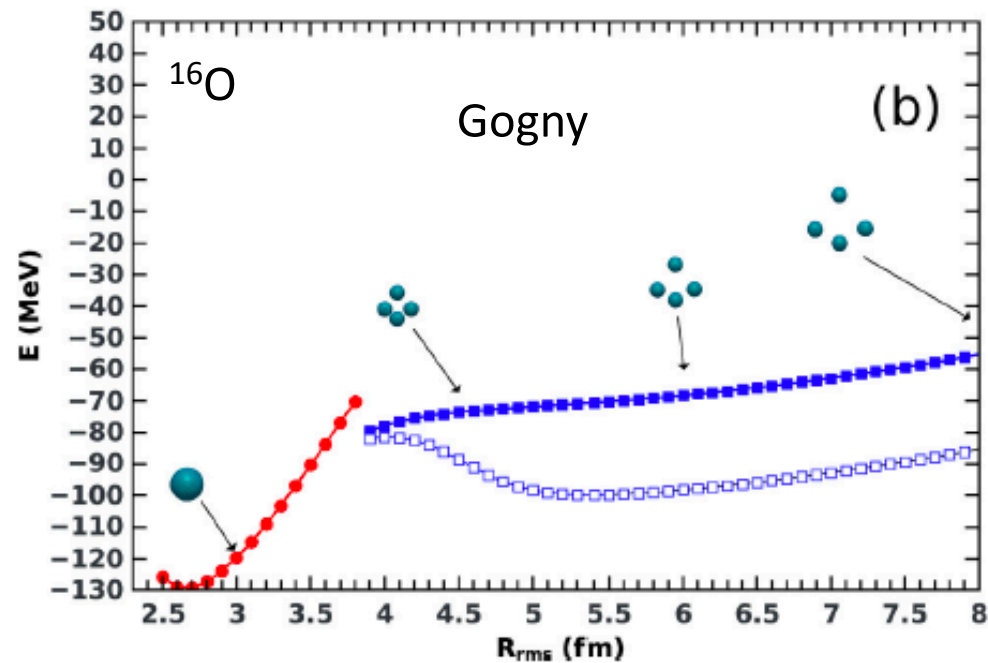
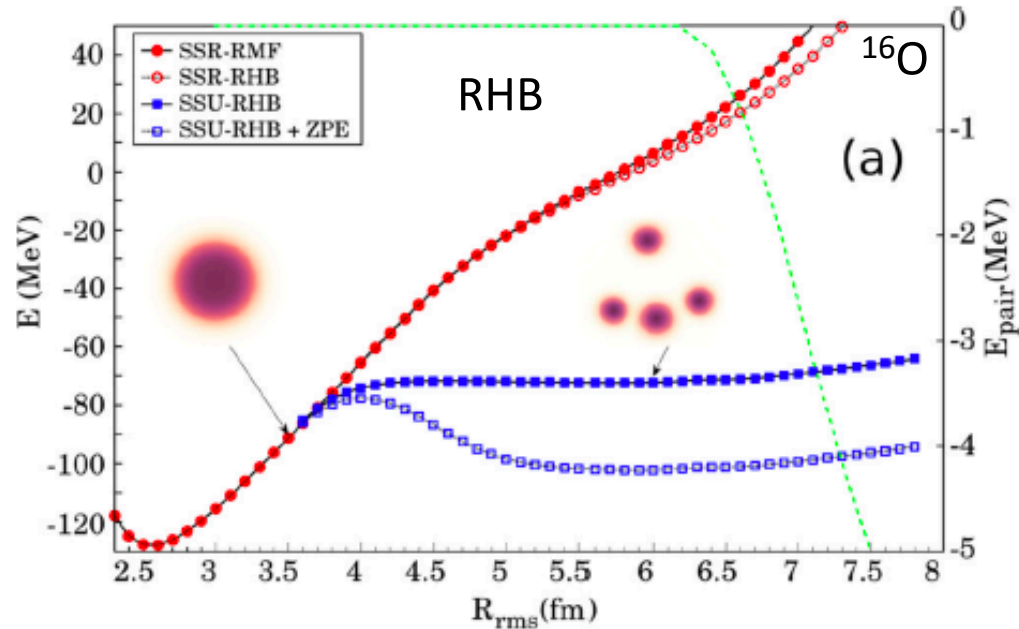
Large deformations lead to clusterization



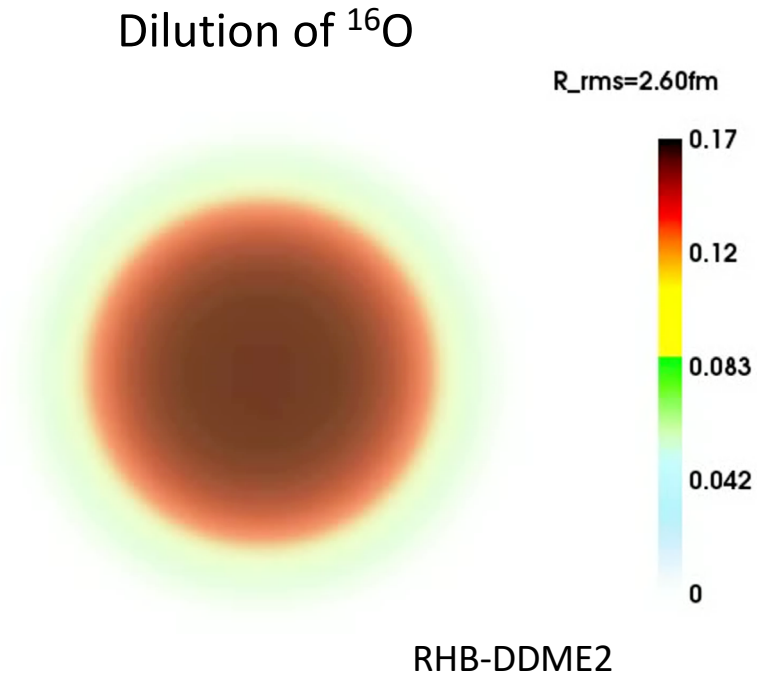
J.-P. Ebran, M. Girod, E. Khan, R. D. Lasserri, and P. Schuck, Phys. Rev. C **102**, 014305 (2020)

Large deformation \longrightarrow lowering the density ?

Clusters in low density nuclear matter



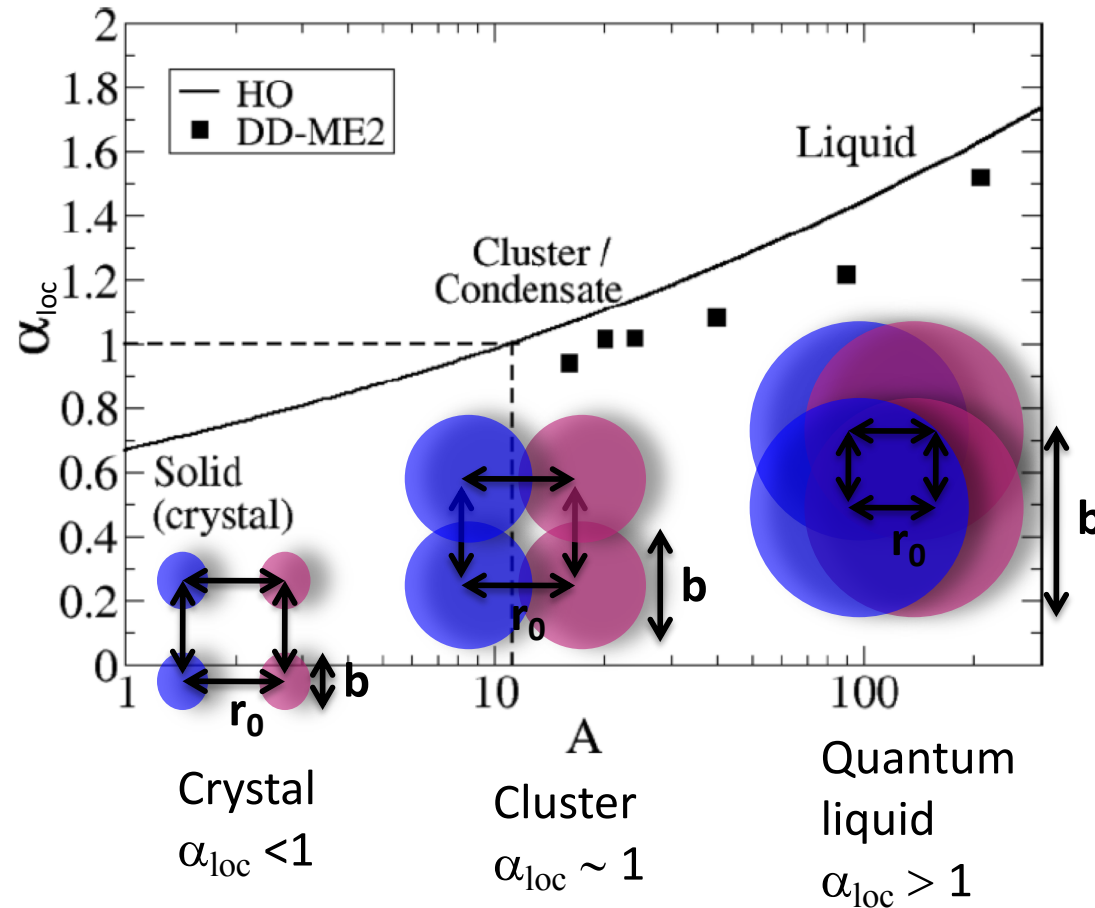
Clusters in low density nuclear matter



See also: P. Schuck and M. Girod Phys. Rev. Lett. 111, 132503 (2013)

(ρ, A) phase diagram

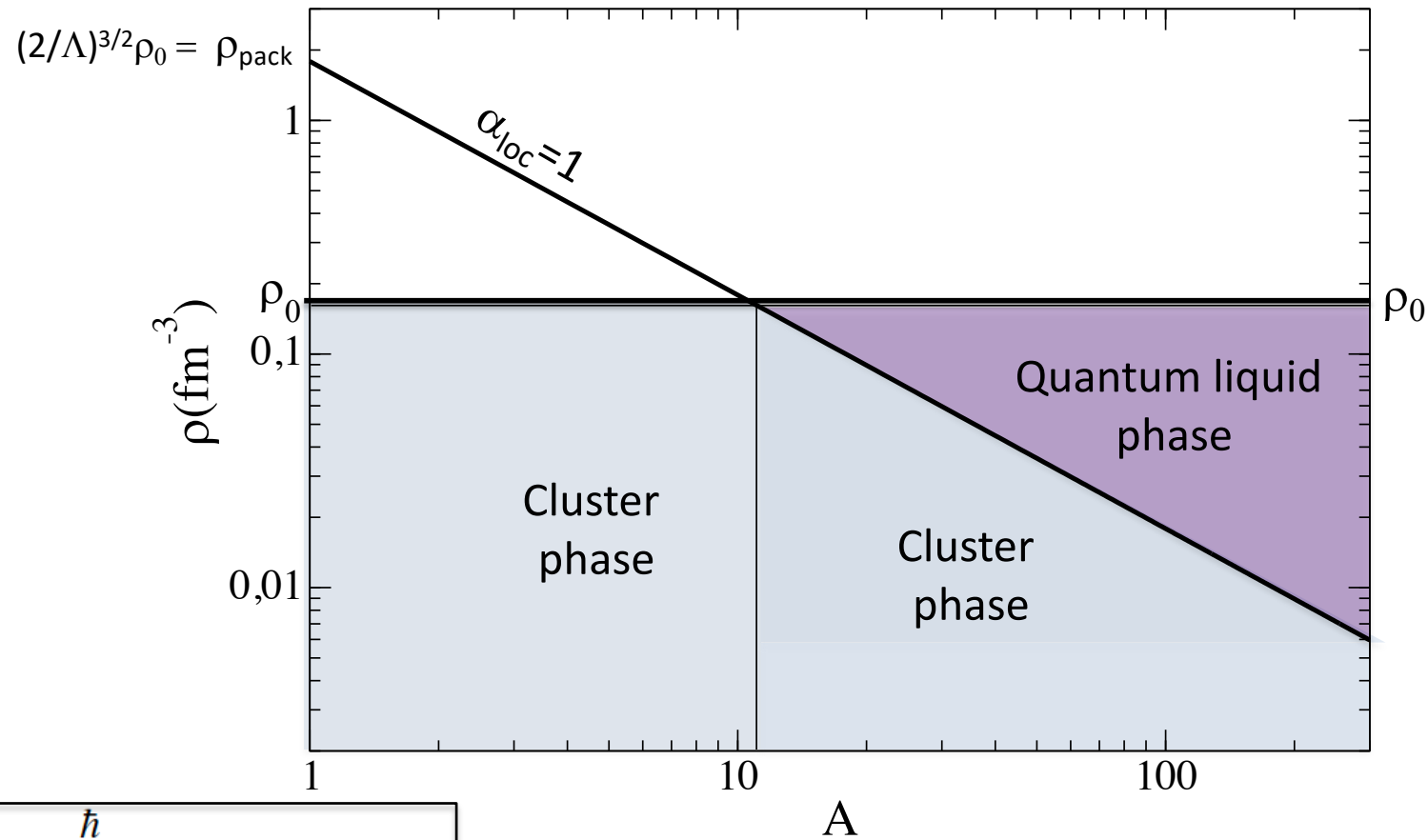
A global indicator for localization



HO
and
RHB calc.

$$\alpha_{loc} = \frac{2\Delta r}{r_0} \simeq \frac{b}{r_0} = \frac{\sqrt{\hbar} A^{1/6}}{(2mV_0 r_0^2)^{1/4}}$$

Density vs. number of nucleons as control parameters of the cluster phase in nuclei



$$r_l = \frac{\hbar}{\sqrt{2mV_0}}$$

$$\alpha_{\text{loc}} = \sqrt{\frac{r_l}{\bar{r}}} A^{1/6} = \left(\frac{\rho}{\rho_l} A \right)^{1/6}$$

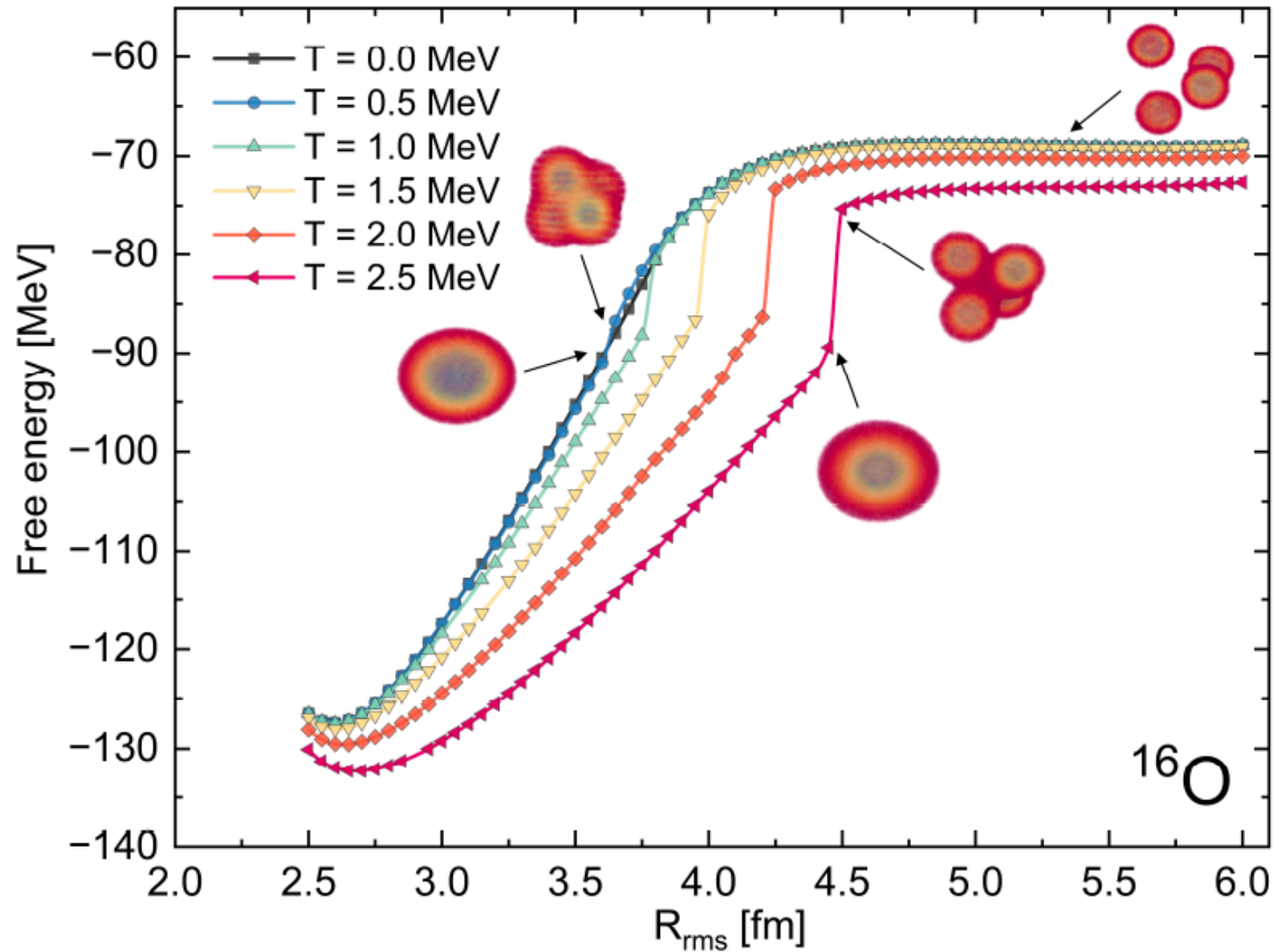
$$\rho A \lesssim \rho_l = \left(\frac{2}{\Lambda} \right)^{3/2} \rho_0 \simeq 10\rho_0$$

J.P. Ebran, L. Heitz, E. Khan, PRC 110(2024)044311

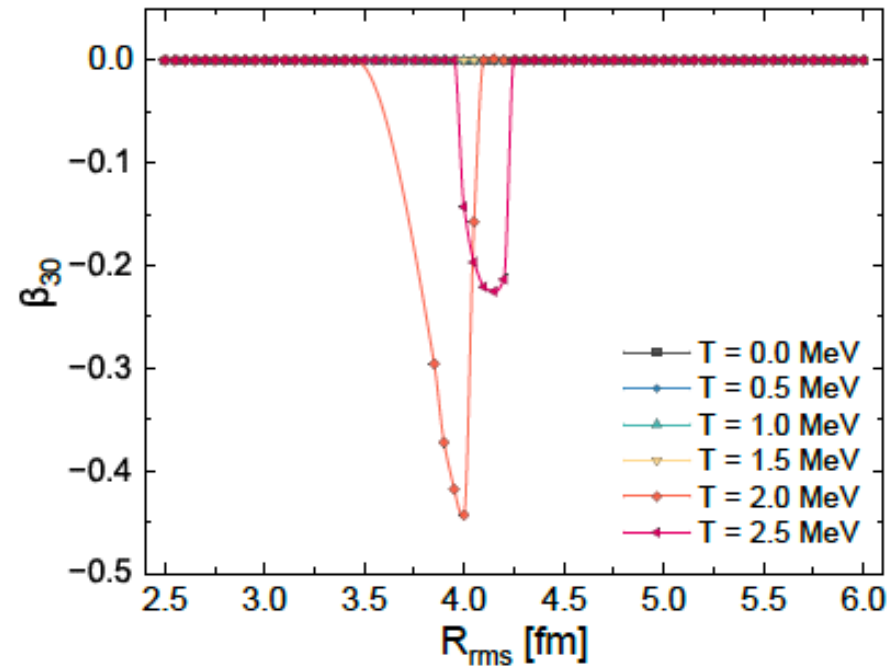
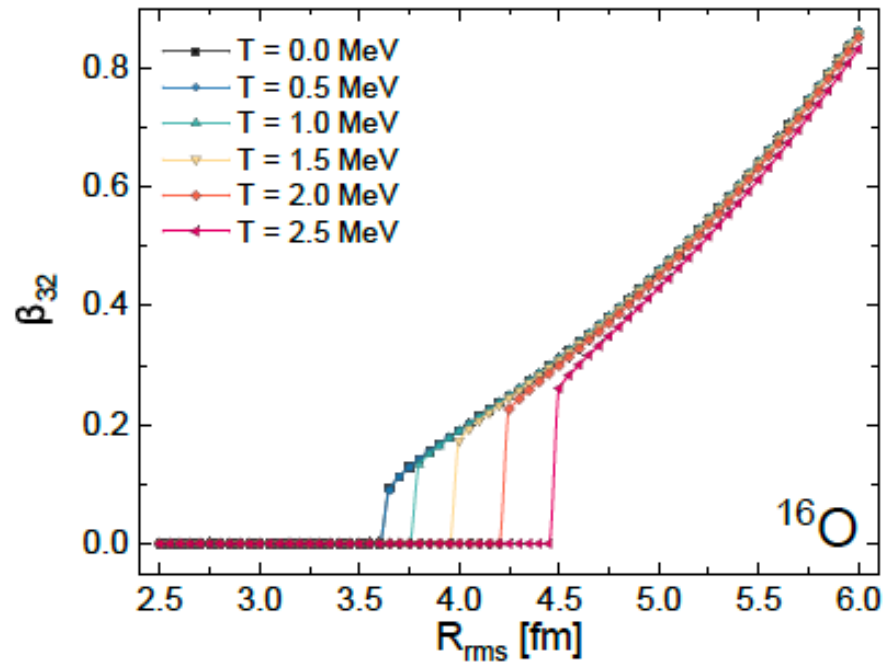
Driven by the quantality and related to the packing density

A (ρ, T) phase diagram ?

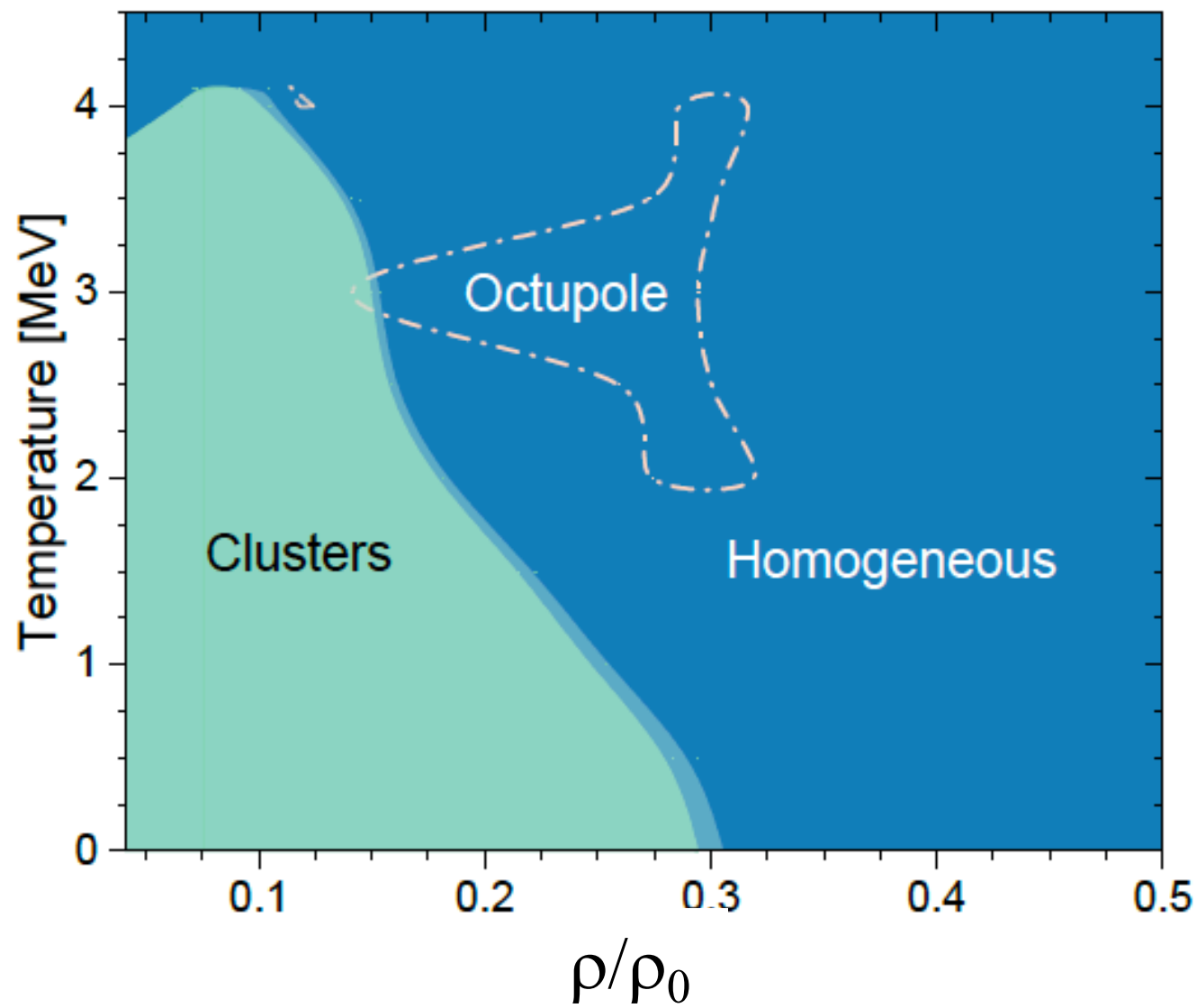
Dilution at finite temperature



Occurrences of tetrahedral and octupole deformations



(ρ, T) diagram



Summary and outlook

- Nuclear clusterization is a phase which can be controlled by
 - deformation
 - density
 - temperature
 - number of nucleons
 - excitation energy
- Useful for clusterization in Heavy Ion Collisions: density and temperature
- Quantum fluctuations in addition to statistical ones ?
- Comparison (limit) with nuclear matter calculations ?