



西南大學

Study of **Nuclear Fission Dynamics** Within **Relativistic Density Functional Theory**

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2024.12

Where is Chongqing



Peking

Shanghai

Chongqing



Beautiful modern city





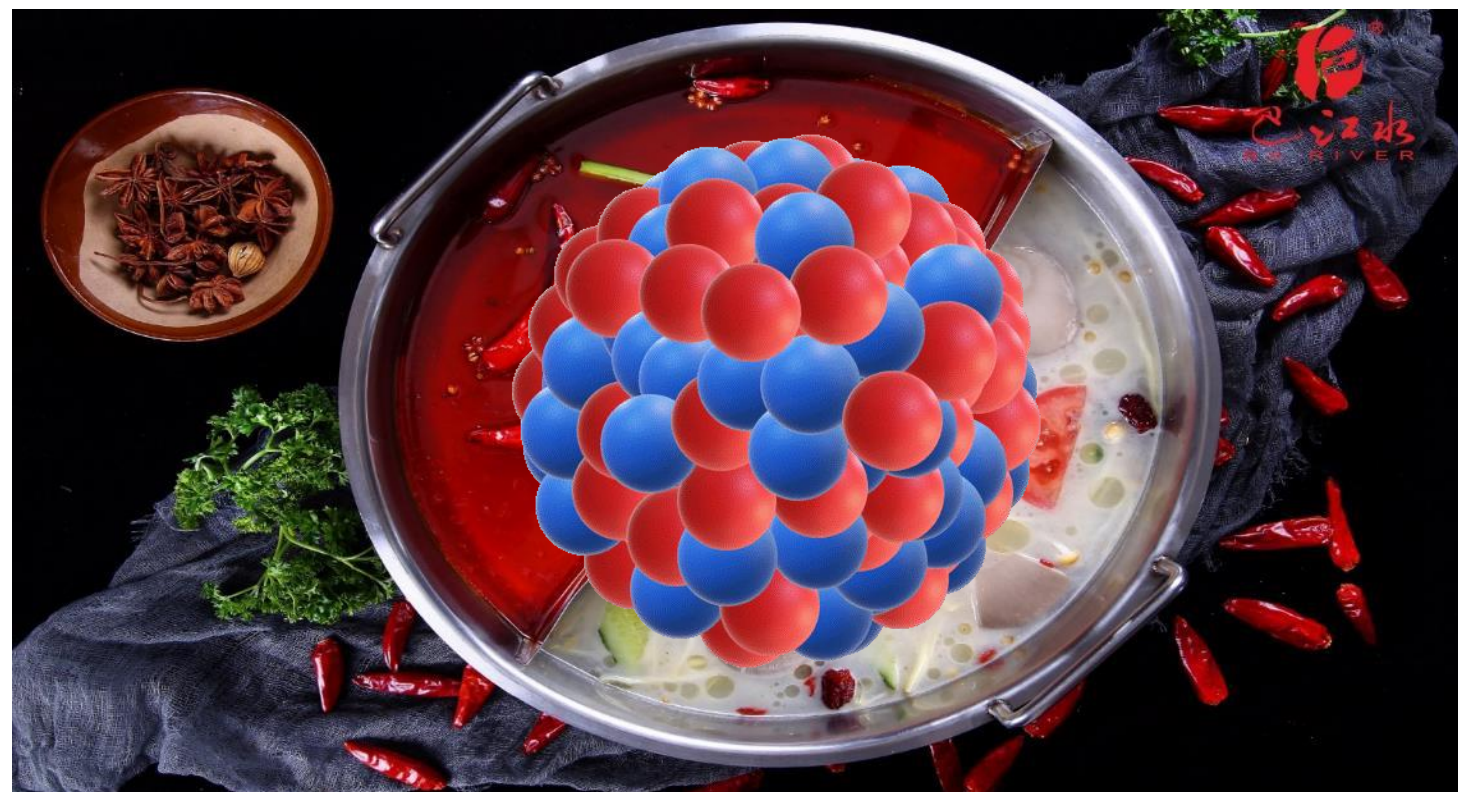
Delicious food...



Hot spring...



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2. Theoretical Framework

- RDFT based on Two Center Harmonic Oscillator (TCHO) basis
- Time Dependent Generator Coordinate Method

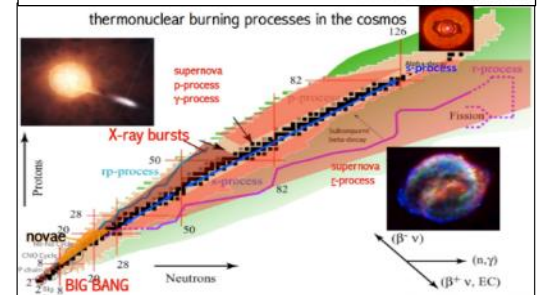
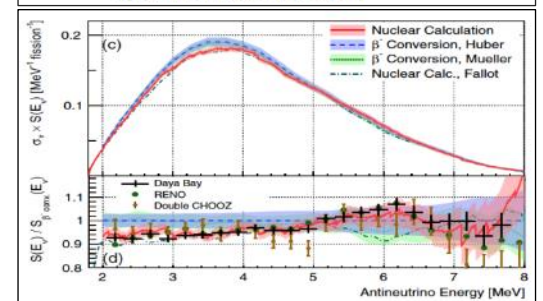
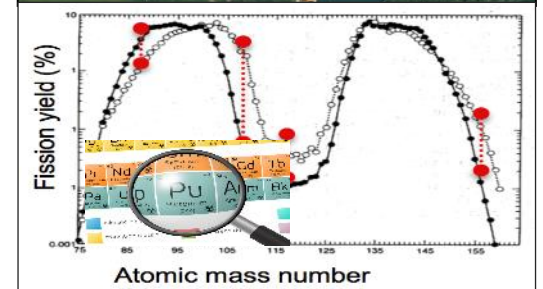
3. Results and Discussion

- Asymmetric fission of ^{180}Hg
- Fission fragment angular momenta
- Multi-dimensional collective space

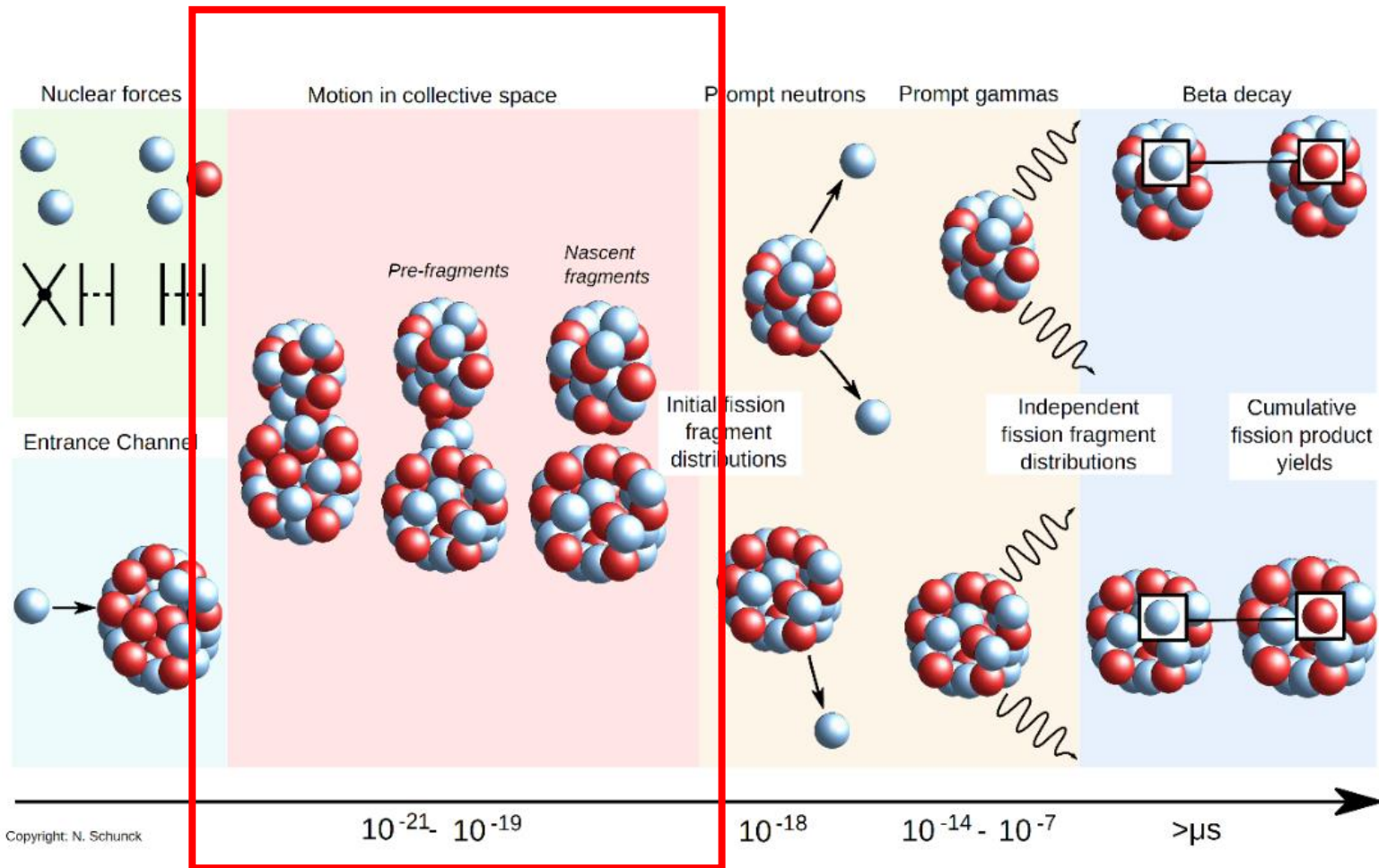
4. Summary & Outlook

Application of Nuclear Fission

- **Reactor design and operation**
 - reactivity calculations, fuel and reactor core management, reactor safety (IND and CUM)
- **Reprocessing of spent fuel and nuclear waste management**
 - Fission product inventory, decay heat calculation (IND)
- **Safeguard**
 - monitor fission products for the verification (CUM)
- **Fundamental physics**
 - understanding fission physics, r-process nucleosynthesis, antineutrino anomaly (FY), fission angular momentum ...



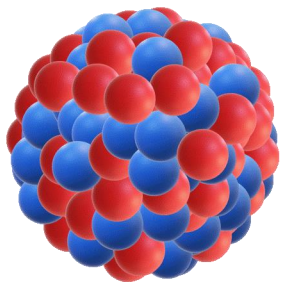
Overview of the Nuclear fission process



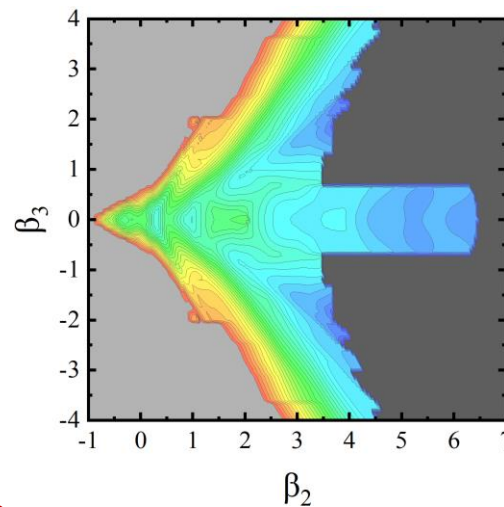
Neutron-induced fission process

- Nuclear fission: an atomic nucleus splits into two or more fragments.
- The process involves **complex correlation, fluctuation, dissipation ...**

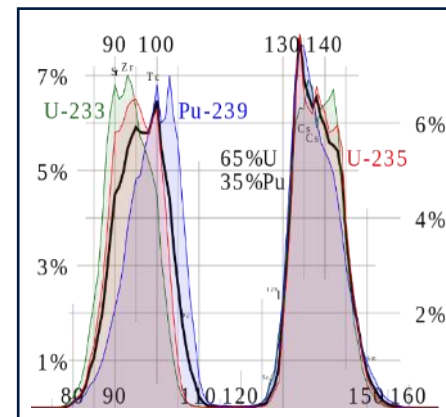
Frontiers of fission



Dynamics

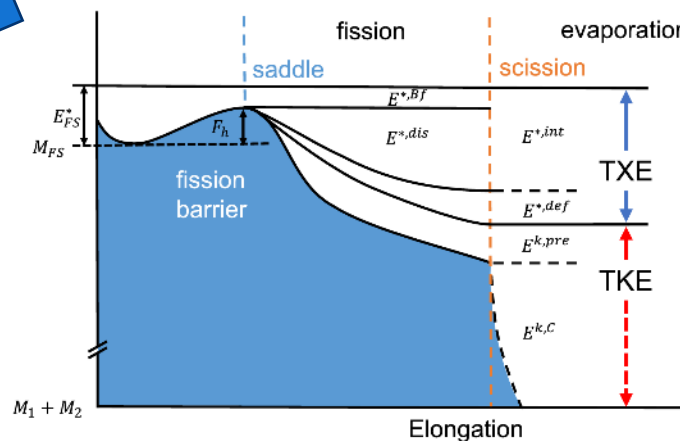
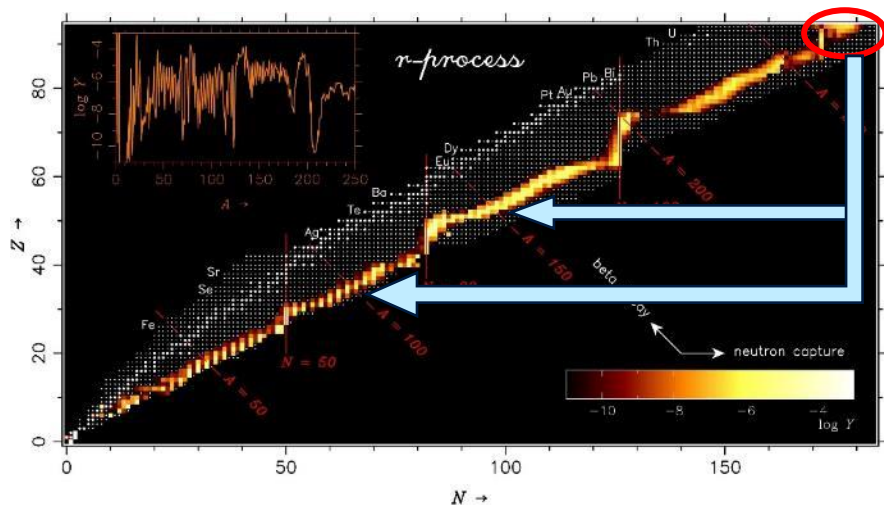
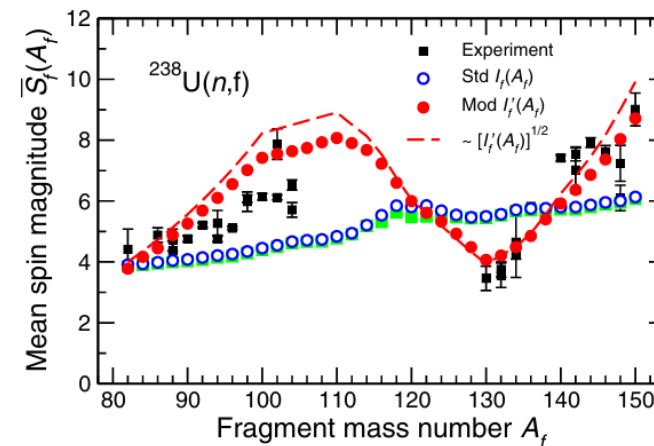


Angular momentum



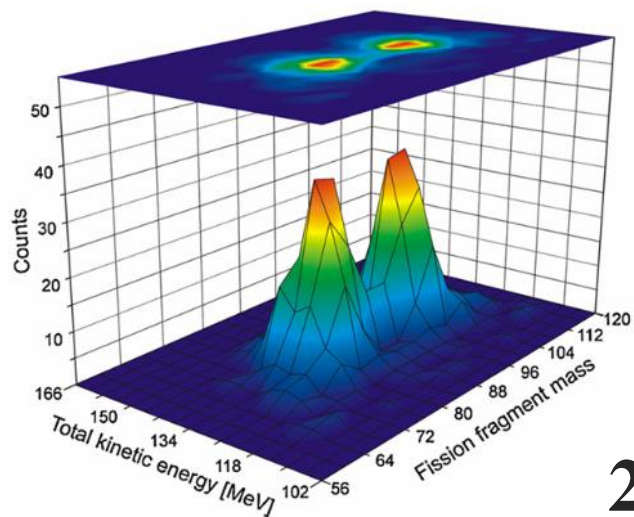
Intersection with astrophysics

Dissipation



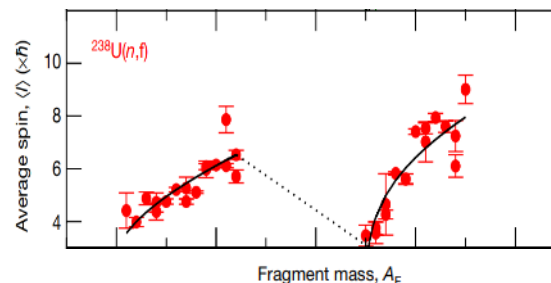
Progress in fission: Experiments & Observations

A. N. Andreyev *et al.* PRL 105, 252502(2010)



New type asymmetric fission in the sub-lead region

J. N. Wilson, *et al.*, Nature 590, 566 (2021)



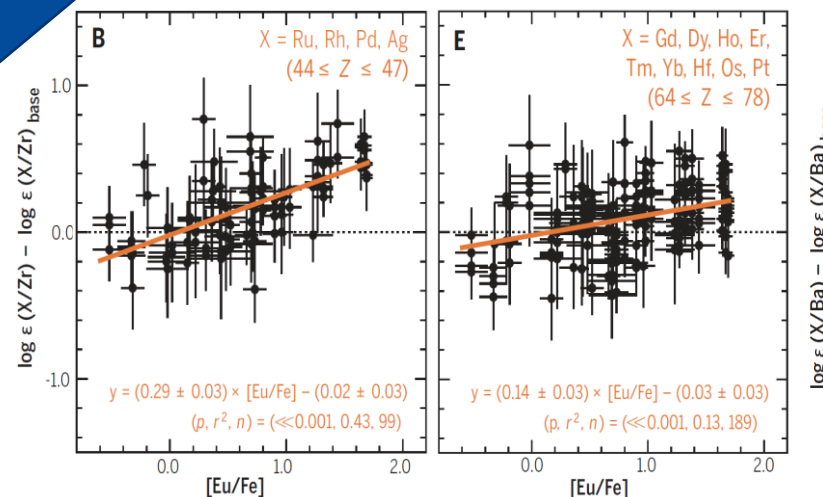
Fragments' angular momenta

2021

2010

2023

Roederer *et al.*, Science 382, 1177–1180(2023)



Fission of transuranic nuclei is involved in the r process

➤ Statistical model

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➤ Macro-micro model

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C. Ishizuka, X. Zhang, M.D. Usang, et al., PRC101, 011601(R)(2020).
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➤ Microscopic model (DFT)

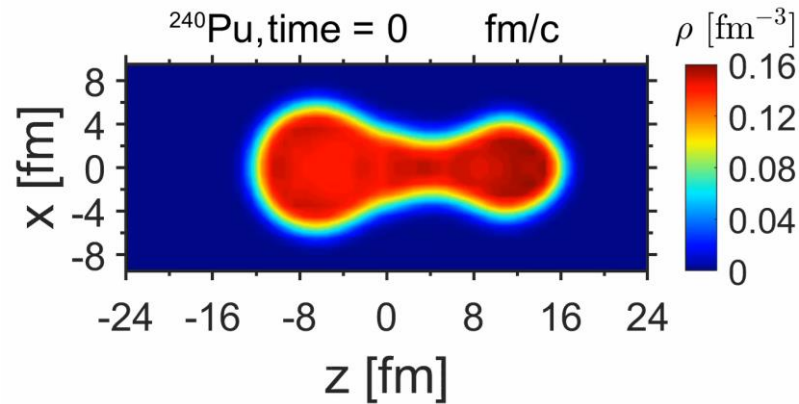
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➤ Machine learning

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➤ Time-Dependent Density Functional Theory (TDDFT)

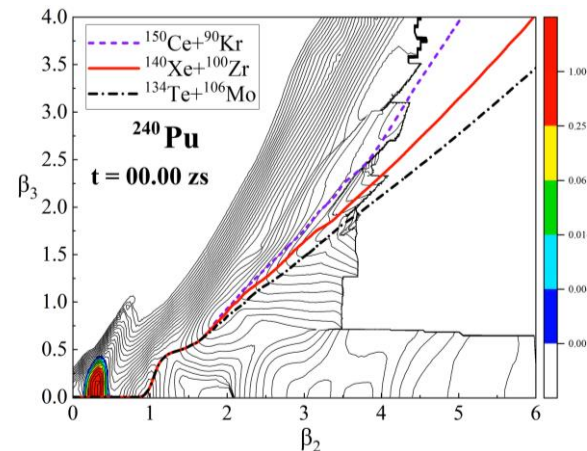
Time evolution of intrinsic spatial wave functions (densities)



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➤ Time-Dependent Generator Coordinate Method (TDGCM)

Time evolution of the probability wave packet in collective space.



J.F. Berger M. Girod, and D. Gogny, CPC63, 365 (1991)
 H. Goutte, J. F. Berger, P. Casoli, *et al.*, PRC71,024316 (2005)
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4. Summary & Outlook



➤ Relativistic energy density functional

$$\begin{aligned}
 E_{\text{CDF}} &= \int d\mathbf{r} \varepsilon_{\text{CDF}}(\mathbf{r}) \\
 &= \sum_k \int d\mathbf{r} v_k^2 \bar{\psi}_k(\mathbf{r}) (-i\gamma\nabla + m) \psi_k(\mathbf{r}) \\
 &+ \int d\mathbf{r} \left(\frac{\alpha_S}{2} \rho_S^2 + \frac{\beta_S}{3} \rho_S^3 + \frac{\gamma_S}{4} \rho_S^4 + \frac{\delta_S}{2} \rho_S \Delta \rho_S \right. \\
 &+ \frac{\alpha_V}{2} j_\mu j^\mu + \frac{\gamma_V}{4} (j_\mu j^\mu)^2 + \frac{\delta_V}{2} j_\mu \Delta j^\mu + \frac{e}{2} \rho_p A^0 \\
 &\left. + \frac{\alpha_{TV}}{2} \vec{j}_{TV}^\mu \cdot (\vec{j}_{TV})_\mu + \frac{\delta_{TV}}{2} \vec{j}_{TV}^\mu \cdot \Delta (\vec{j}_{TV})_\mu \right)
 \end{aligned}$$

$\rho_S, j^\mu, \vec{j}_{TV}^\mu$: Density and currents of nucleons

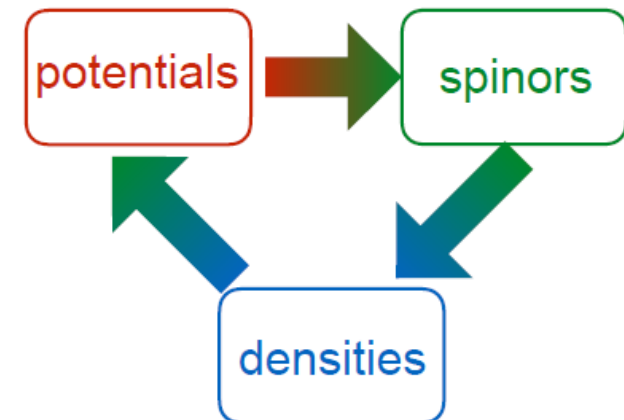
$\alpha, \beta, \gamma, \delta$: Coupling constants (PC-PK1, DD-PC1, etc.)

➤ Single particle Dirac equation:

$$\begin{pmatrix} m+V+S & \sigma \cdot p \\ \sigma \cdot p & -(m-V-S) \end{pmatrix} \begin{pmatrix} f \\ g \end{pmatrix} = \varepsilon \begin{pmatrix} f \\ g \end{pmatrix}$$

Scalar potential and vector potential

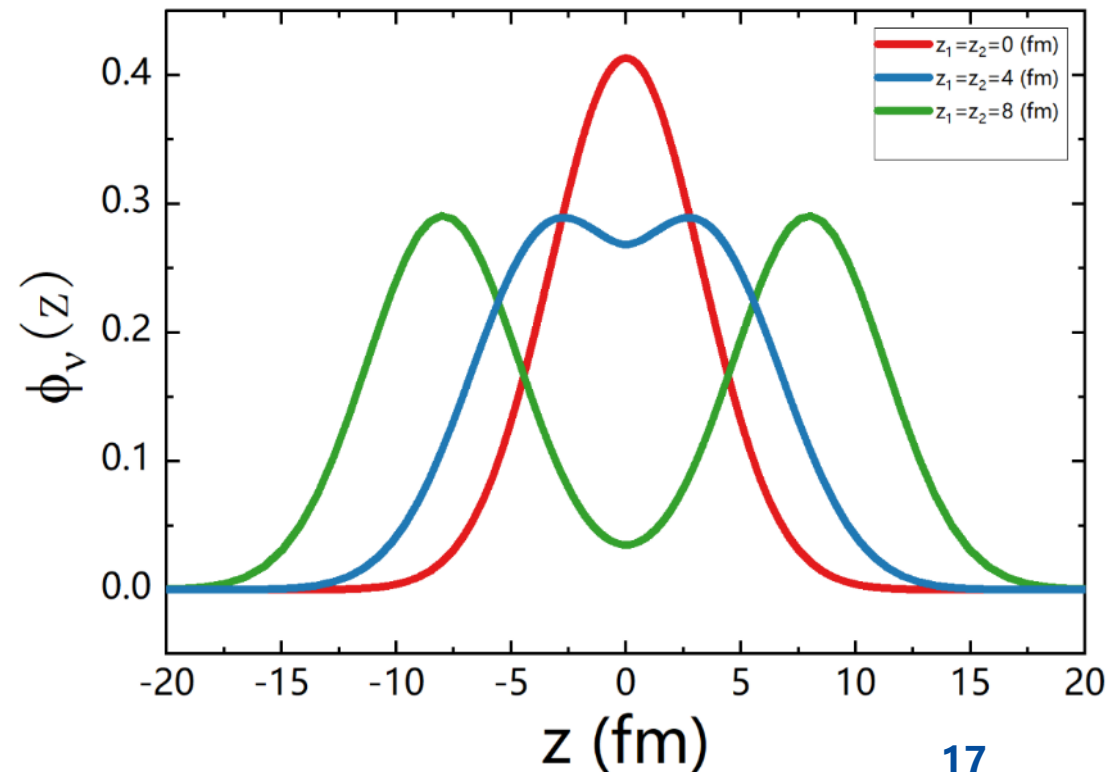
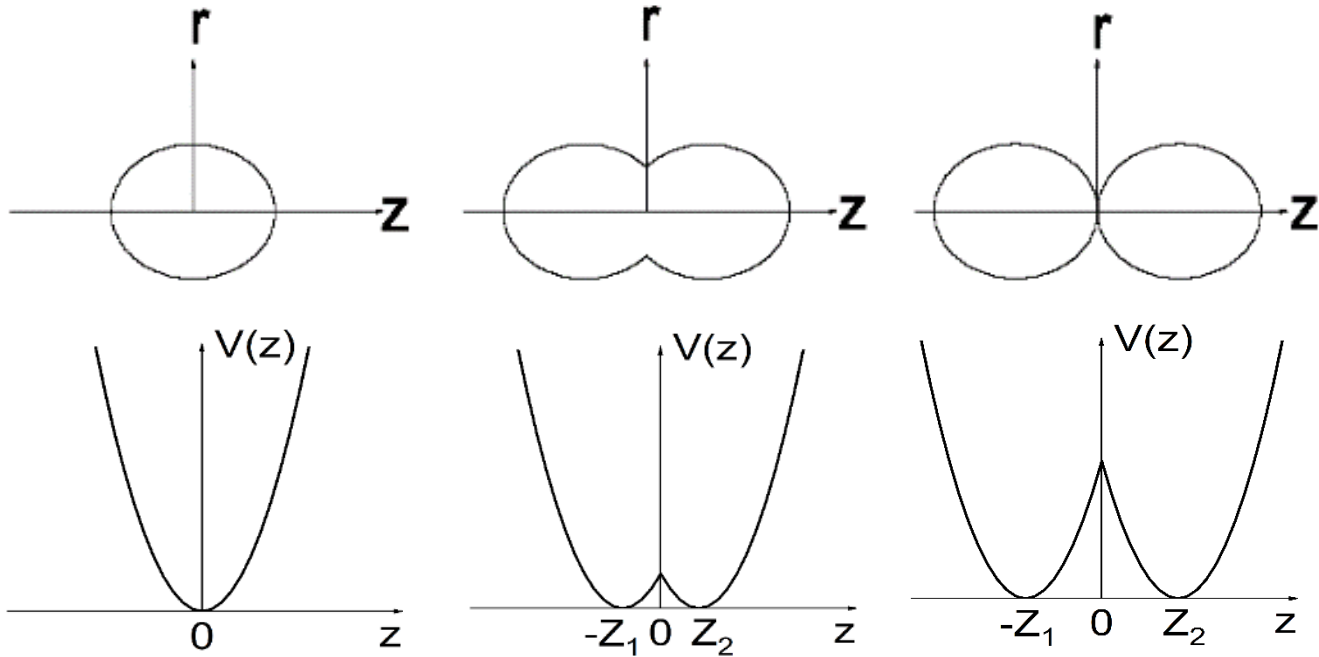
$$\begin{aligned}
 S(\mathbf{r}) &= \alpha_S \rho_S + \beta_S \rho_S^2 + \gamma_S \rho_S^3 + \delta_S \Delta \rho_S, \\
 V^\mu(\mathbf{r}) &= \alpha_V j^\mu + \gamma_V (j_\nu j^\nu) j^\mu + \delta_V \Delta j^\mu - e A^\mu \frac{1-\tau_3}{2} \\
 &+ \tau_3 (\alpha_{TV} j_{TV}^\mu + \delta_{TV} \Delta j_{TV}^\mu),
 \end{aligned}$$



Two Center Harmonic Oscillator basis

➤ Potential of Two Center Harmonic Oscillator (TCHO) basis

$$V(r_{\perp}, z) = V(r_{\perp}) + V(z) = \frac{1}{2}M\omega_{\perp}^2 r_{\perp}^2 + \begin{cases} \frac{1}{2}M\omega_1^2(z + z_1)^2, & z < 0, \\ \frac{1}{2}M\omega_2^2(z - z_2)^2, & z \geq 0, \end{cases} \xrightarrow[\text{WF}]{\text{Eigen}} \Phi(z, r_{\perp}, \varphi) = \phi_v(z) \phi_{n_r}^{m_l}(r_{\perp}) \Theta(\varphi) \chi_{m_s}$$

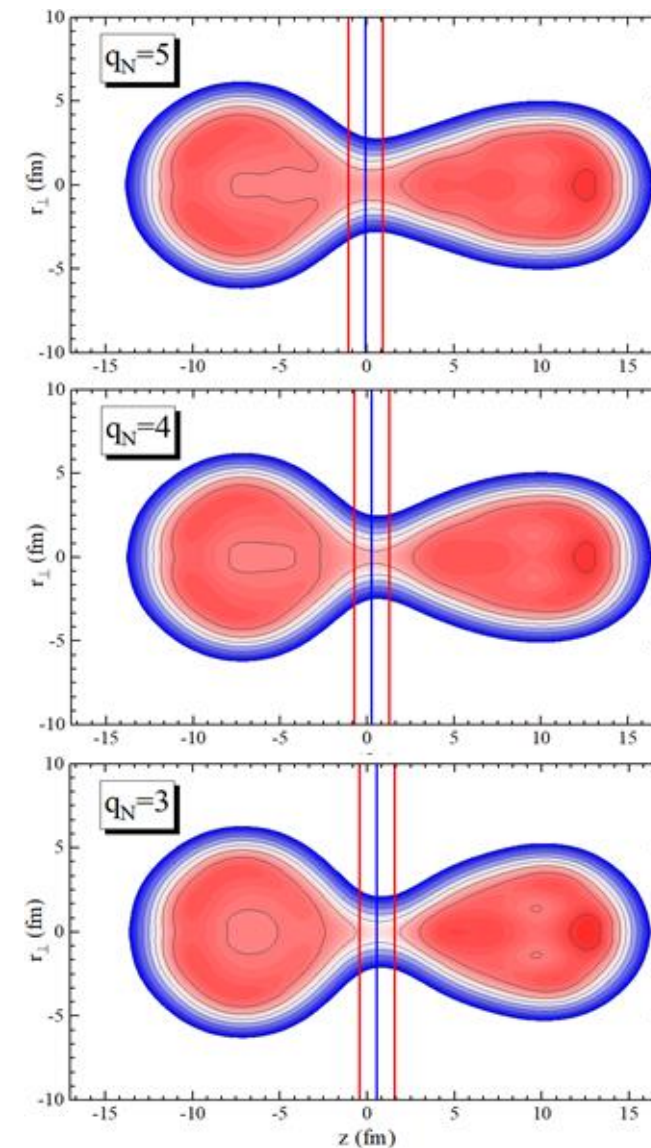
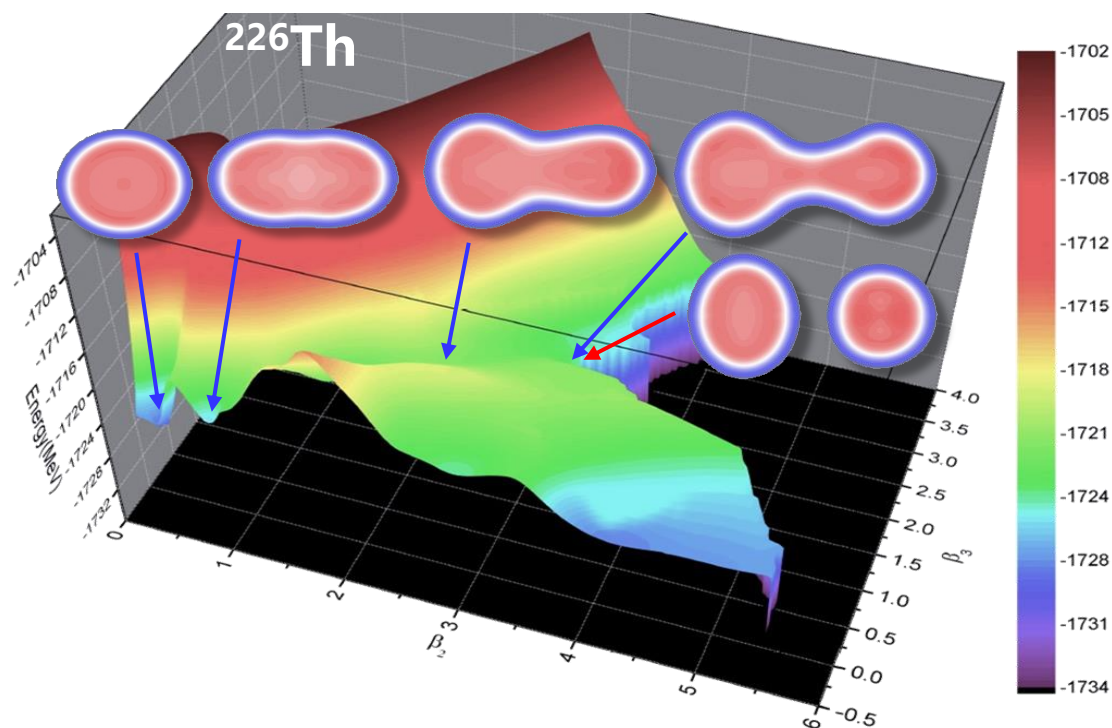


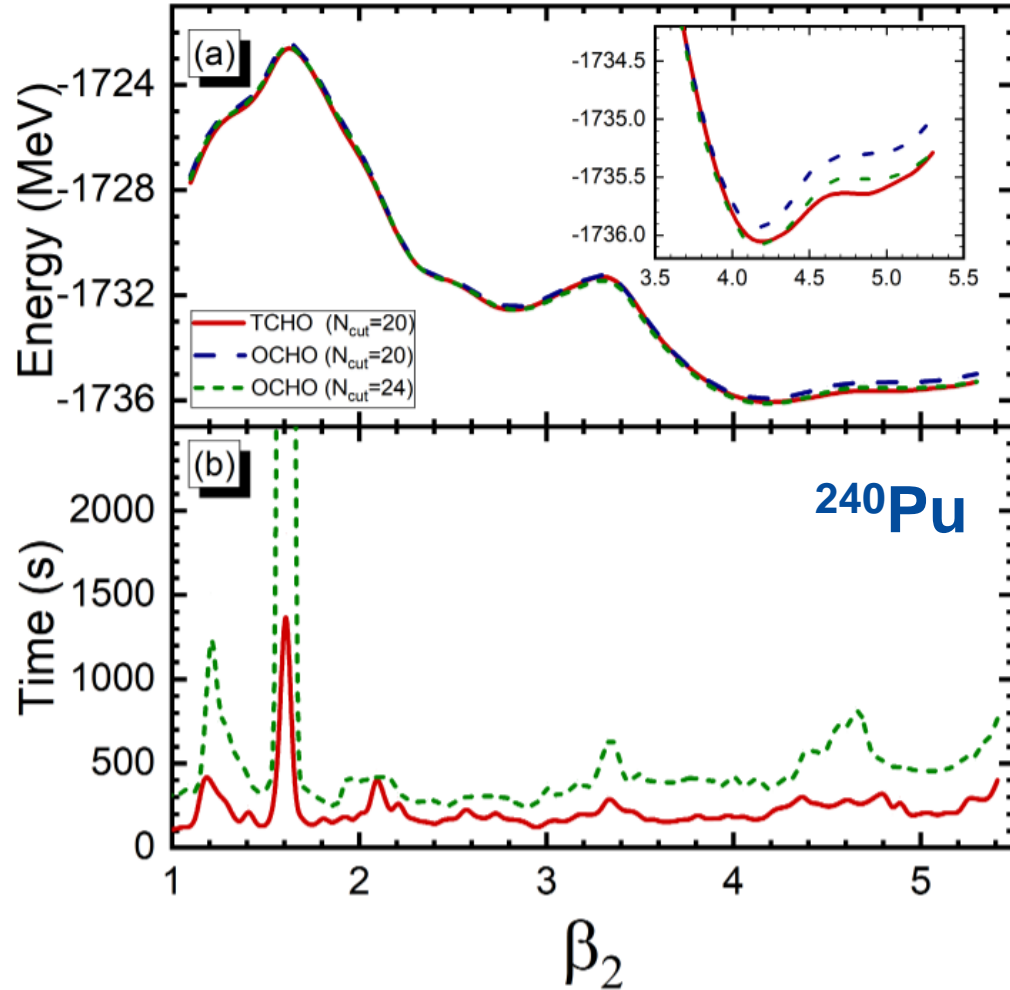
Multi-dimensional constraints for PES

$$\langle E_{RDFT} \rangle + \sum_{k=2,3} C_k (\langle \hat{Q}_k \rangle - q_k)^2 + C_N (\langle \hat{Q}_N \rangle - q_N)^2$$

$\hat{Q}_{2,3}$ — Quadrupole, octupole moment operator

$\hat{Q}_N = \exp[-(z - z_N)^2 / a_N^2]$ z_N is the neck, $a_N = 1 \text{ fm}$

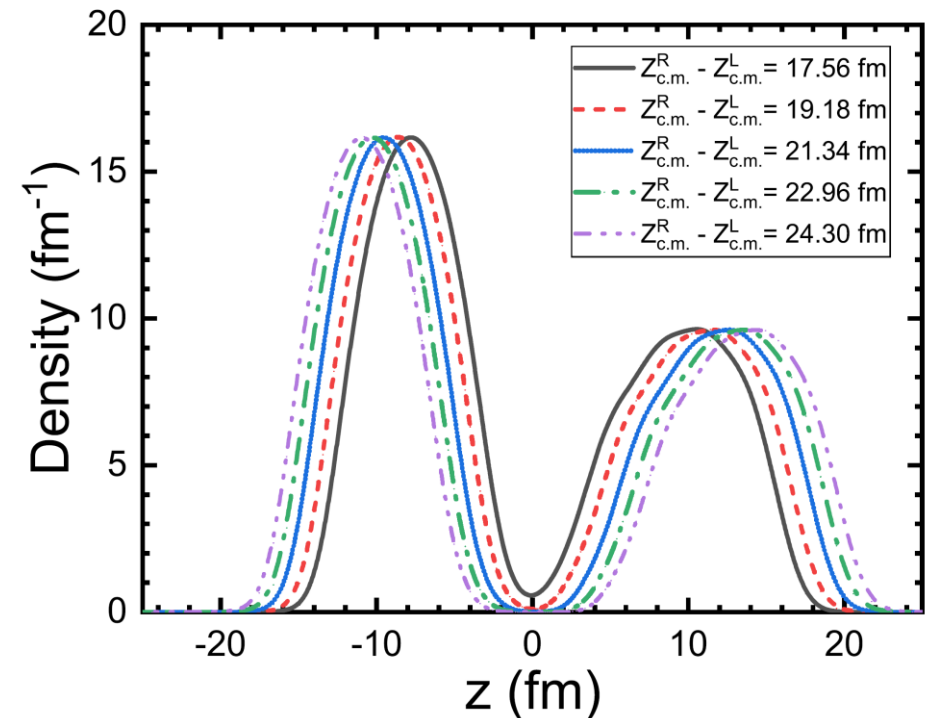




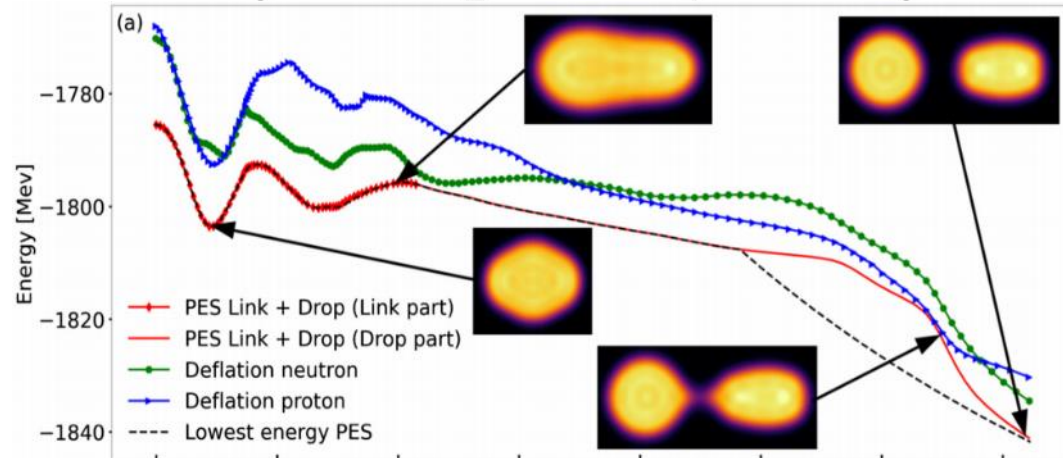
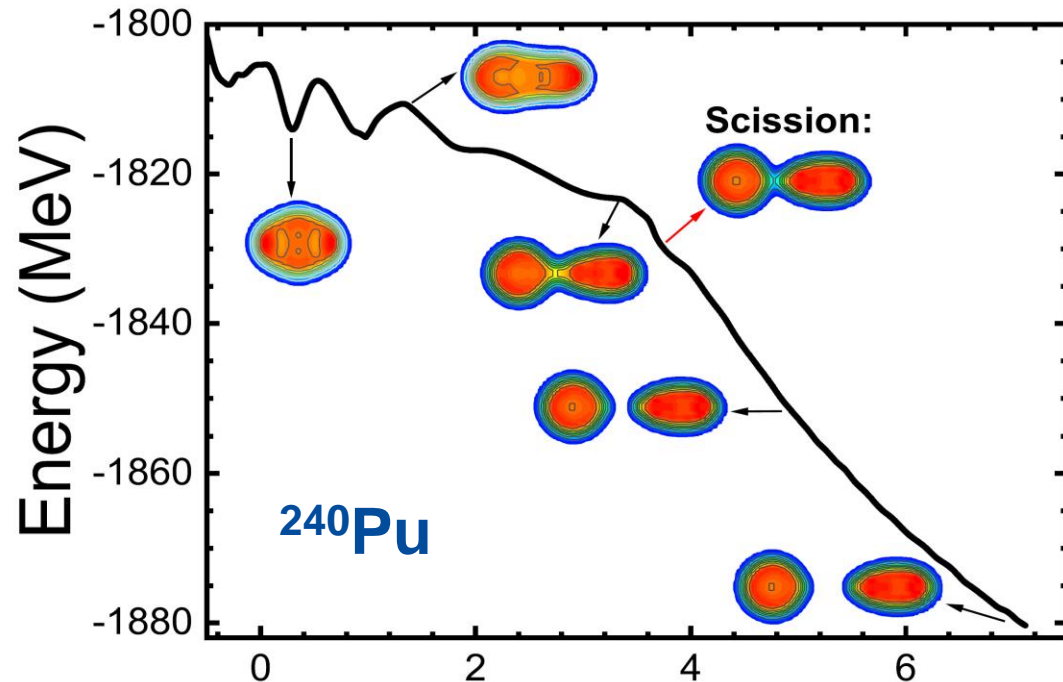
➤ **Better accuracy and efficiency**

➤ **ρ constraint for post-scission**

$$\langle E_{\text{tot}} \rangle + C_{\rho} \left(\int \rho(r_{\perp}, z) - \rho_0(r_{\perp}, z) dr \right)^2$$



Advantages of RDFT-TCHO



➤ Mass tensor calculated in cranking approximation:

$$\mathcal{B} = \mathcal{M}_{(1)}^{-1} \mathcal{M}_{(2)} \mathcal{M}_{(1)}^{-1}$$

$$\mathcal{M}_{(n),kl} = \sum_{i,j} \frac{\langle i | \hat{Q}_k | j \rangle \langle j | \hat{Q}_l | i \rangle}{(E_i + E_j)^n} (u_i v_j + v_i u_j)^2$$

Girod & Grammaticos, NPA330, 40 (1979)

Provides the microscopic inputs for the dynamical study involving pre-post-scission configurations.

➤ Time Dependent Generator Coordinate Method (TDGCM+GOA)

$$i\hbar \frac{\partial}{\partial t} g(\beta_2, \beta_3, t) = \left[-\frac{\hbar^2}{2} \sum_{kl} \frac{\partial}{\partial \beta_k} B_{kl}(\beta_2, \beta_3) \frac{\partial}{\partial \beta_l} + V(\beta_2, \beta_3) \right] g(\beta_2, \beta_3, t)$$

➤ Probability current

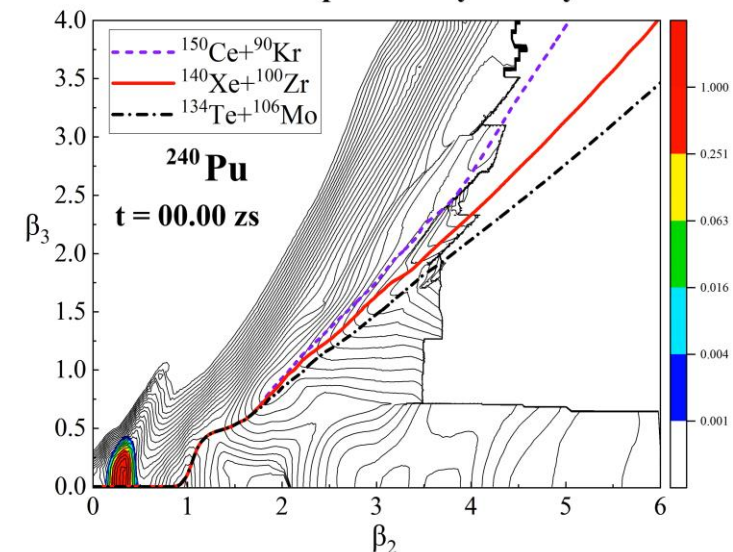
$$J_k(\beta_2, \beta_3, t) = \frac{\hbar}{2i} \sum_{l=2}^3 B_{kl}(\beta_2, \beta_3) \left[g^*(\beta_2, \beta_3, t) \frac{\partial g(\beta_2, \beta_3, t)}{\partial \beta_l} - g(\beta_2, \beta_3, t) \frac{\partial g^*(\beta_2, \beta_3, t)}{\partial \beta_l} \right]$$

➤ Fission yield distribution

$$Y(A) = \int_0^T dt \vec{J}(\beta_2, \beta_3, t) \cdot \vec{n} ds$$

Tao, Zhao, ZPLi, Niksic, Vretenar, PRC96, 024319 (2017)

Evolution of collective probability density distribution





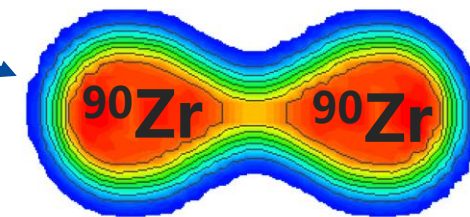
Result 1

Asymmetric fission of ^{180}Hg



New type asymmetric fission of ^{180}Hg

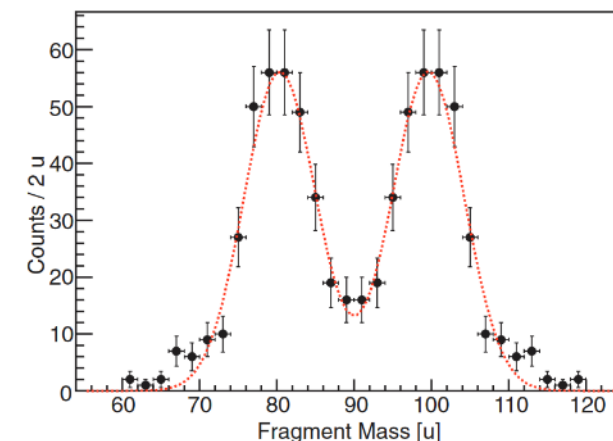
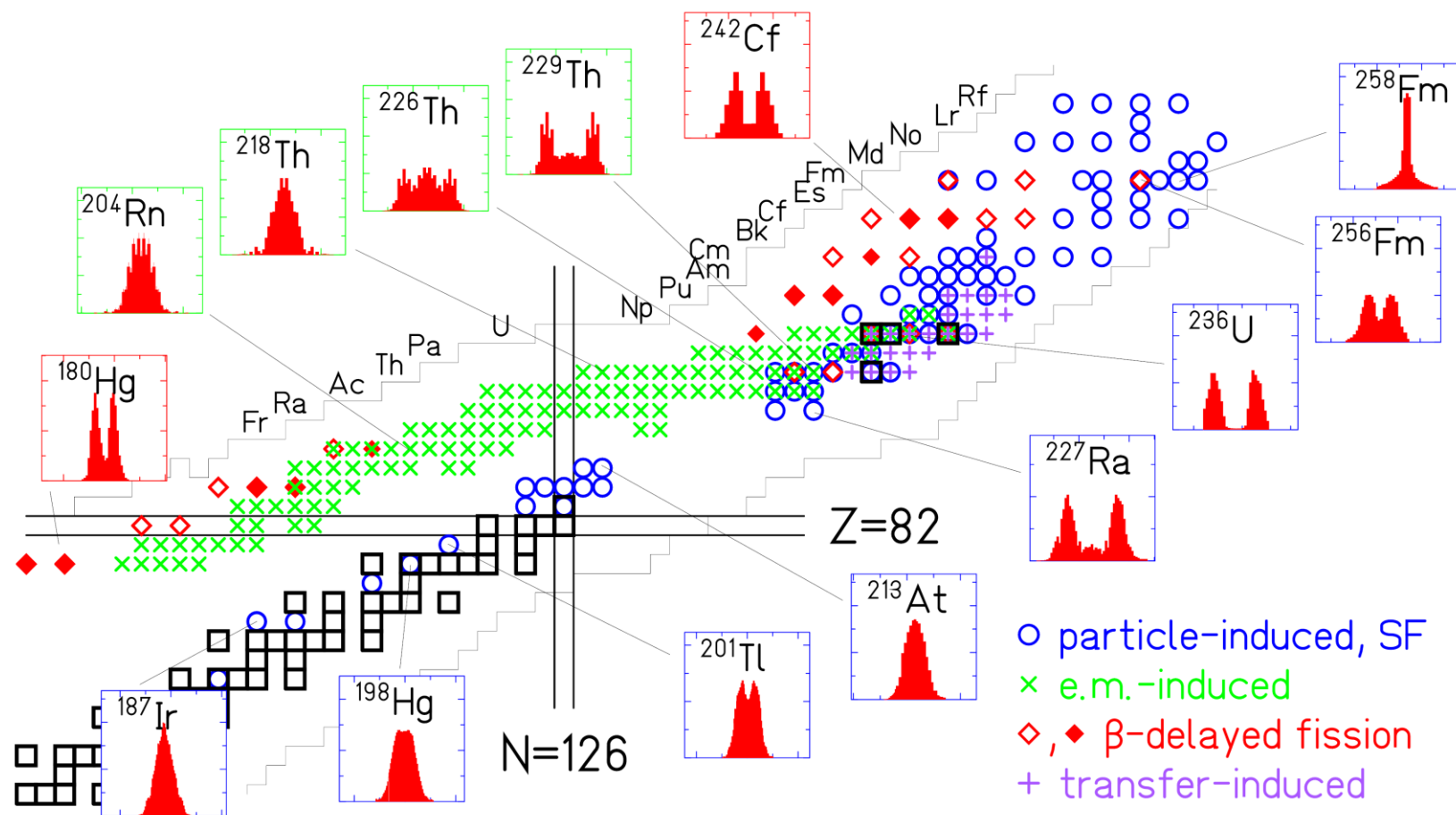
“double magic”



↑ predicted

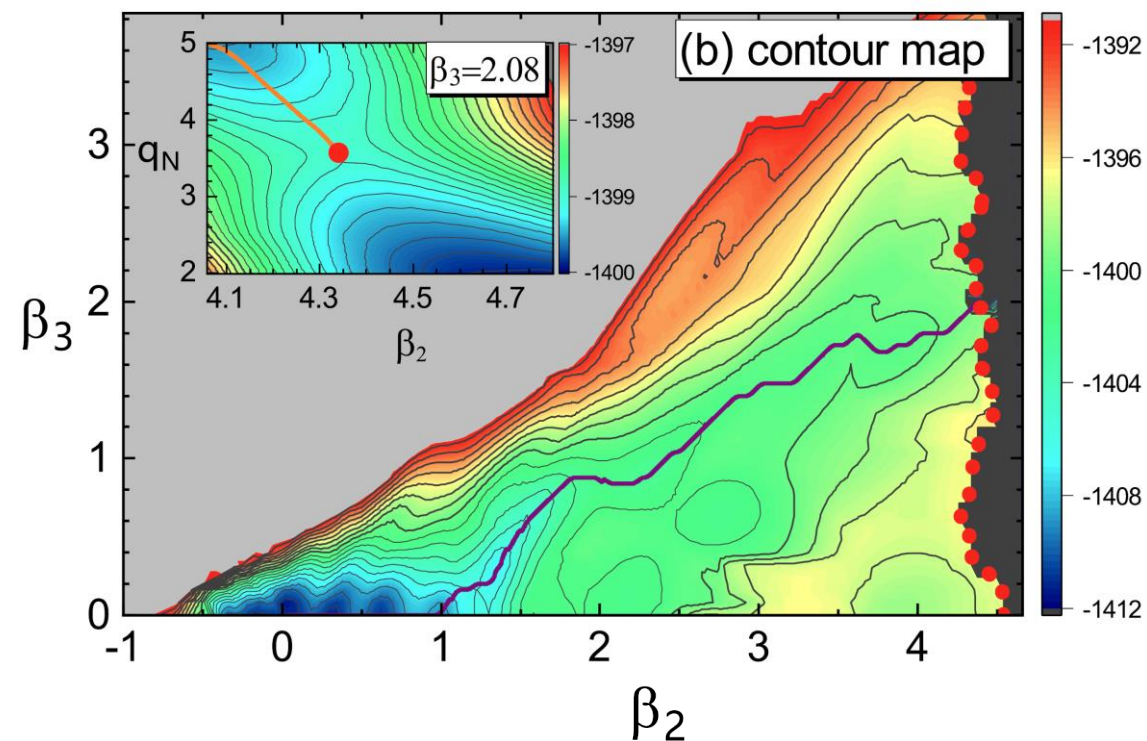
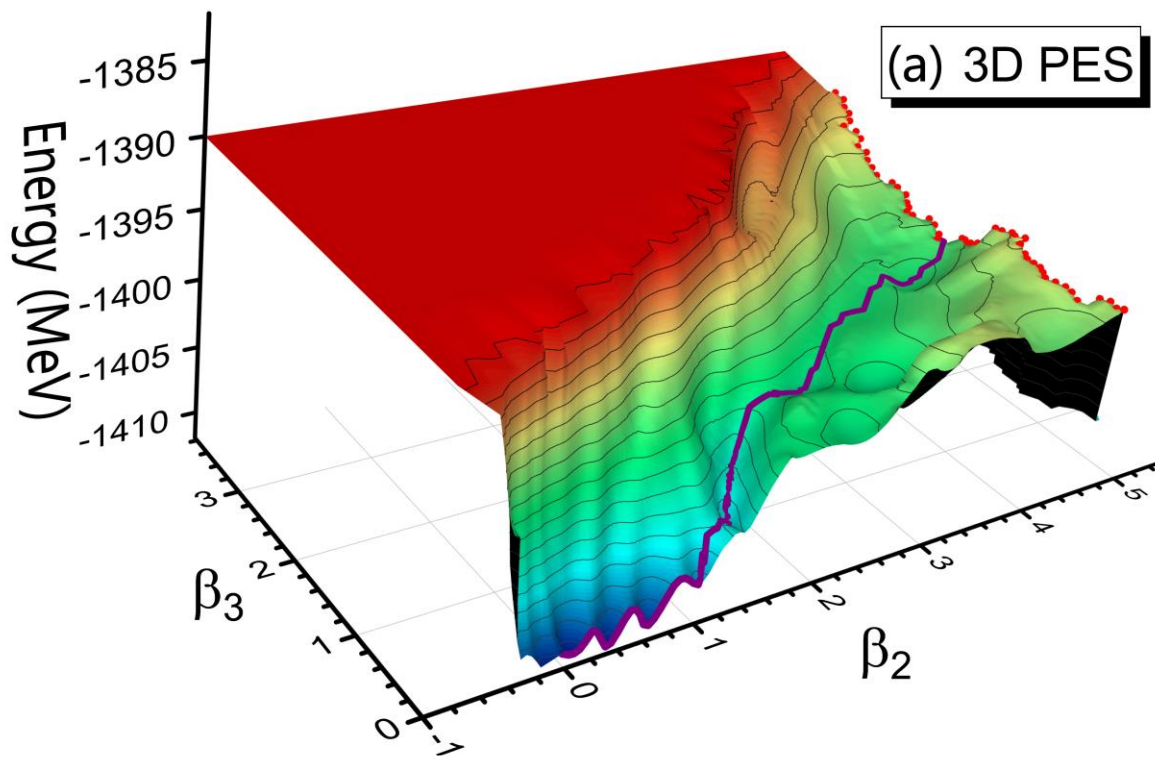
^{180}Hg

↓ observed



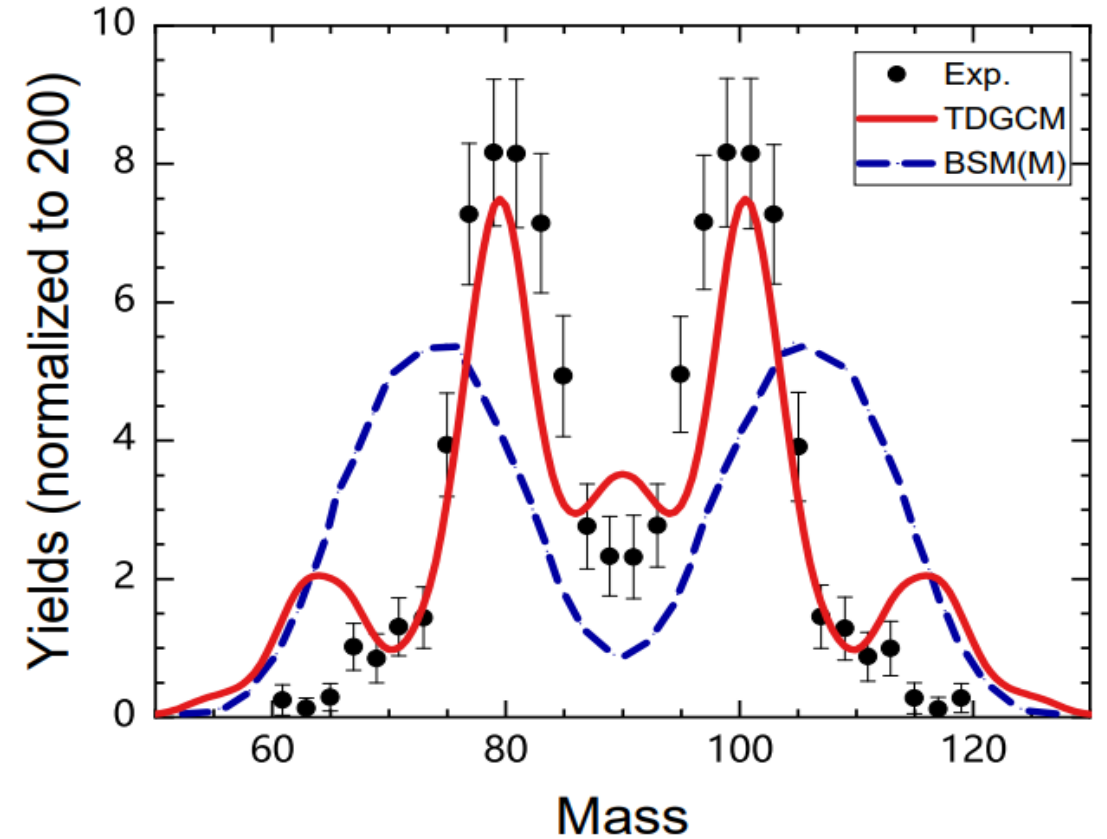
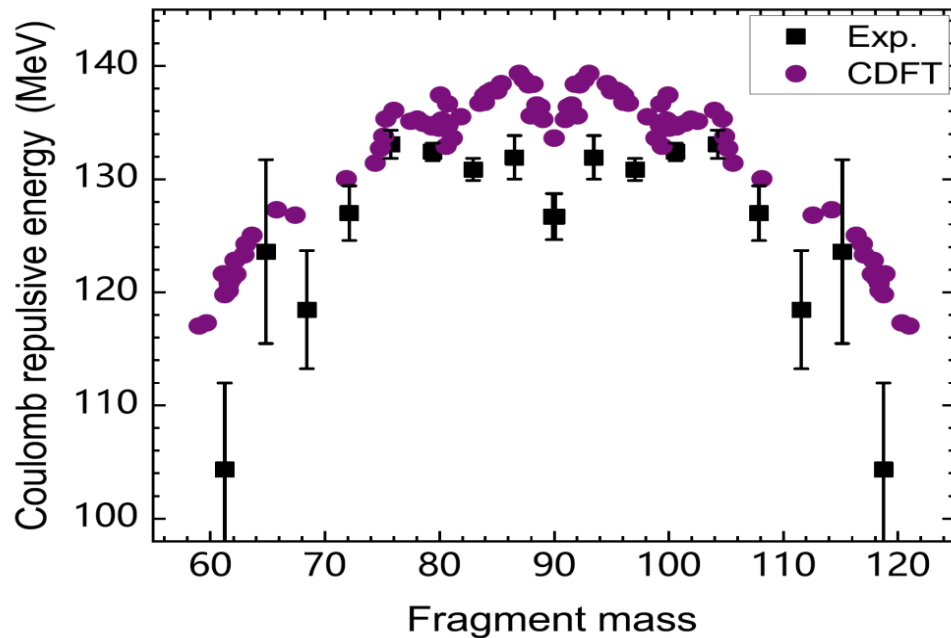
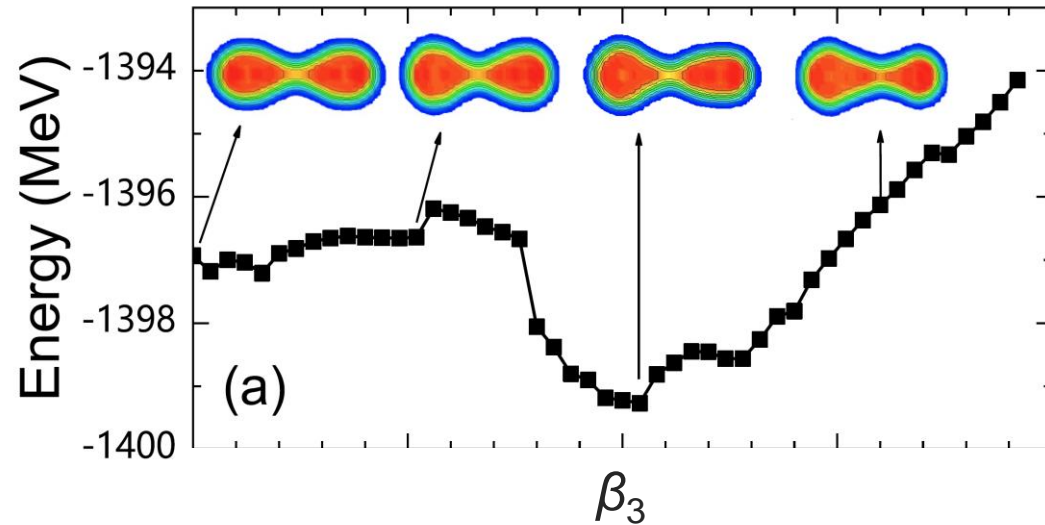
Andreyev, Nishio & Schmidt, Rep. Prog. Phys. (2018)

Andreyev et al., PRL (2010)



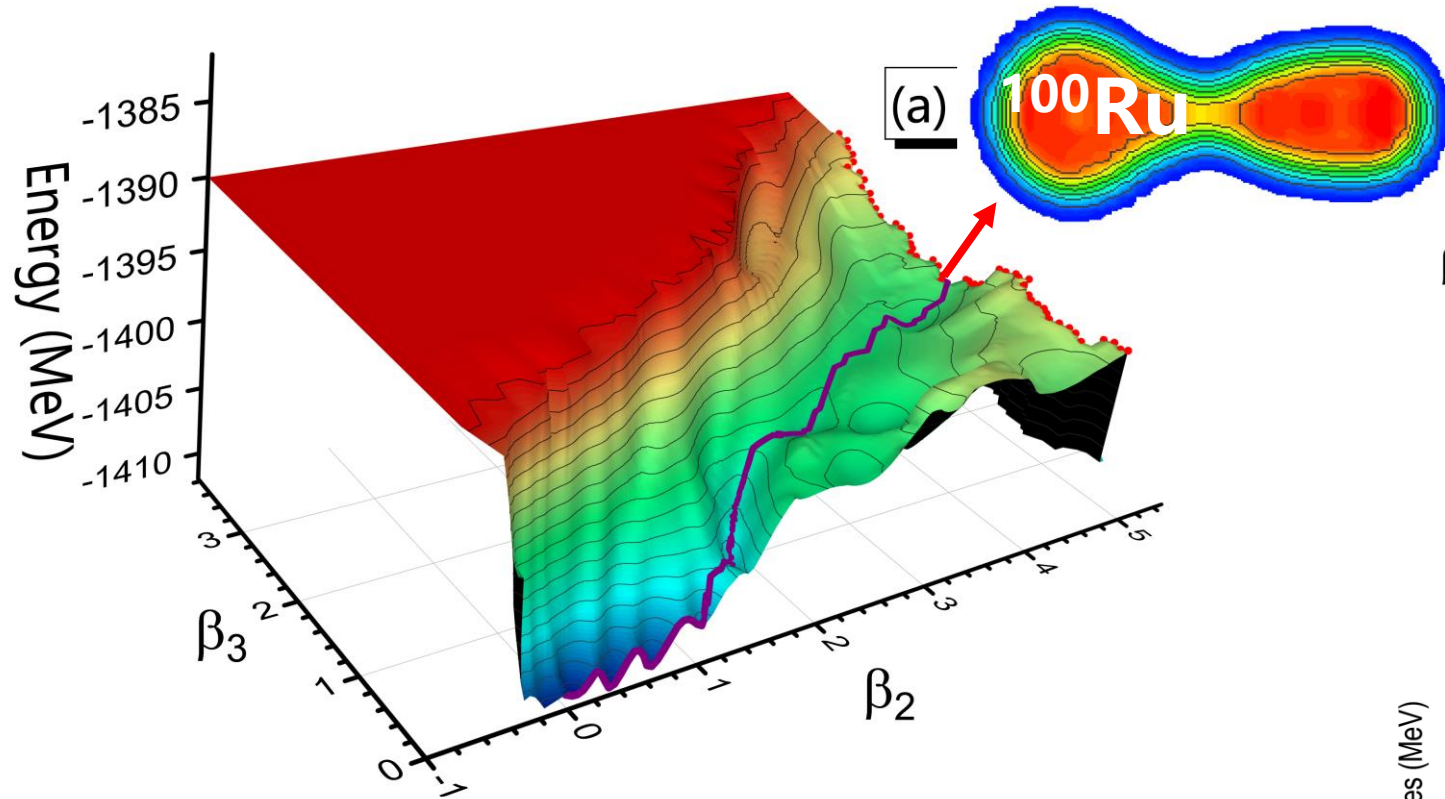
- Lower asymmetric fission valley is predicted
- Scission is defined as the saddle in the β_2 - q_N plane

Scission, TKE, and Yield distributions

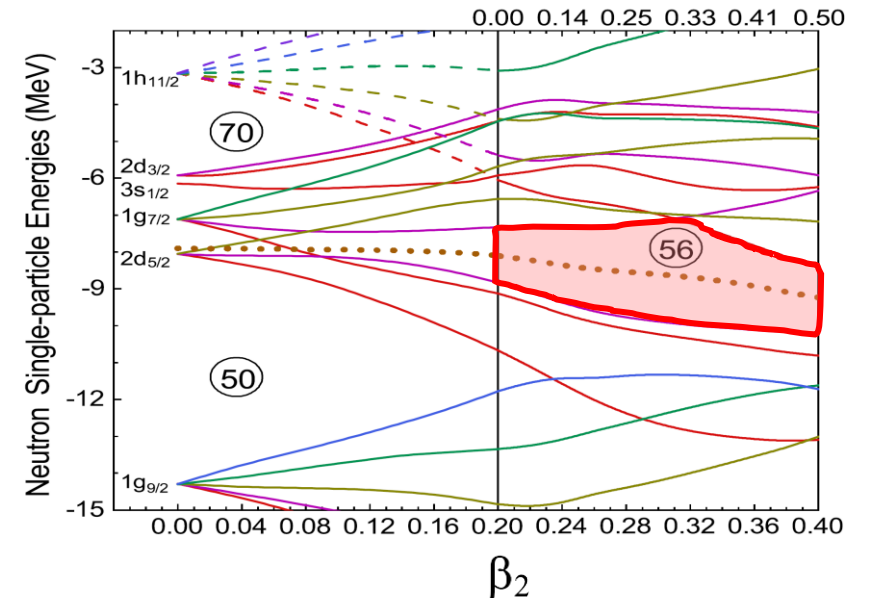
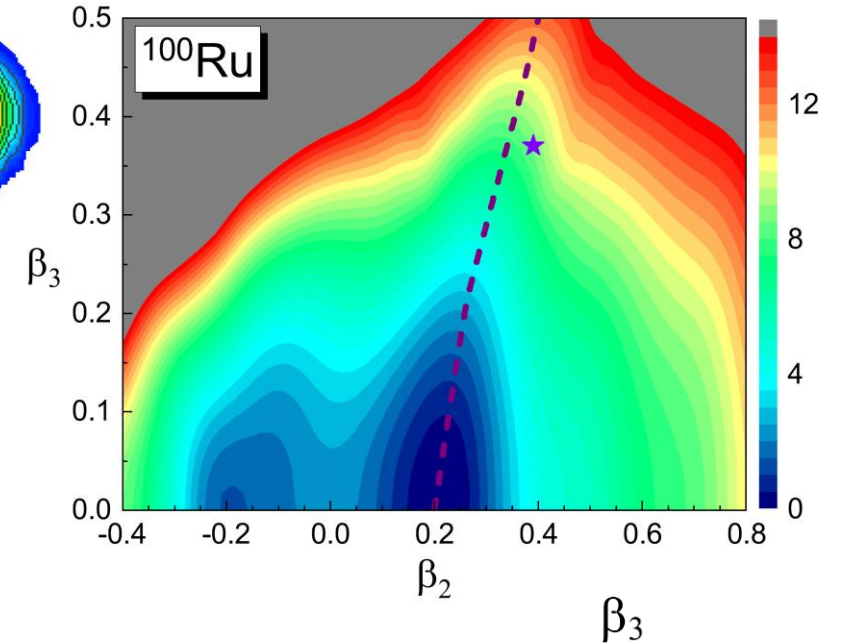


- Detailed structure of TKE
- Reproduce asymmetric fission

Mechanics for asymmetric fission



- Heavy fragment $\sim ^{100}\text{Ru}$: octupole deformed
- N=56 octupole shell of heavy fragment play crucial role



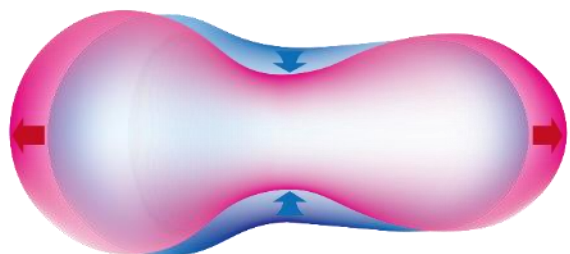


Result 2

Fission Fragment Angular Momenta

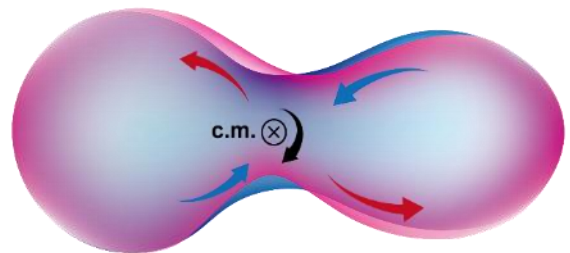


➤ Collective Hamiltonian:



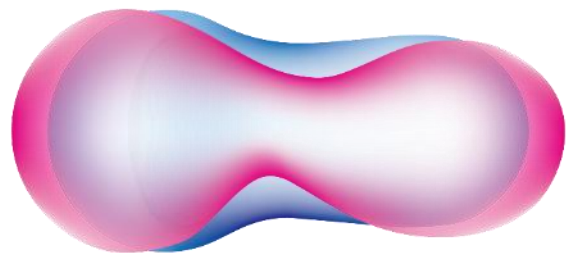
Vibration

$$\hat{H} = -\frac{\hbar^2}{2} \sum_{p,q=2,3} \frac{\partial}{\partial \beta_p} B_{pq}^{-1} \frac{\partial}{\partial \beta_q} + V(\beta_2, \beta_3)$$



Rotation

$$+ \sum_i \frac{\hat{S}_i^2}{2I_i} + \sum_{i < j} \chi_2 \hat{Q}_2^i \cdot \hat{Q}_2^j \quad (i, j = L, H, \Lambda)$$



Vibration
+Rotation

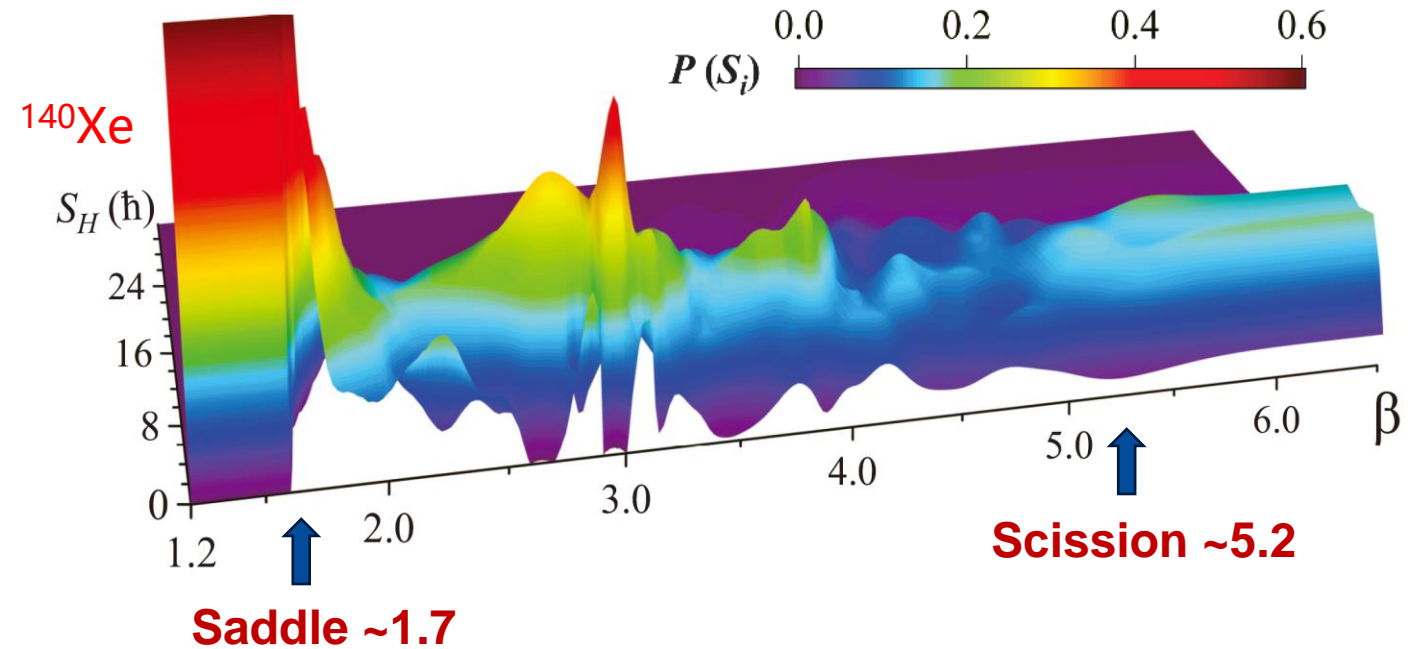
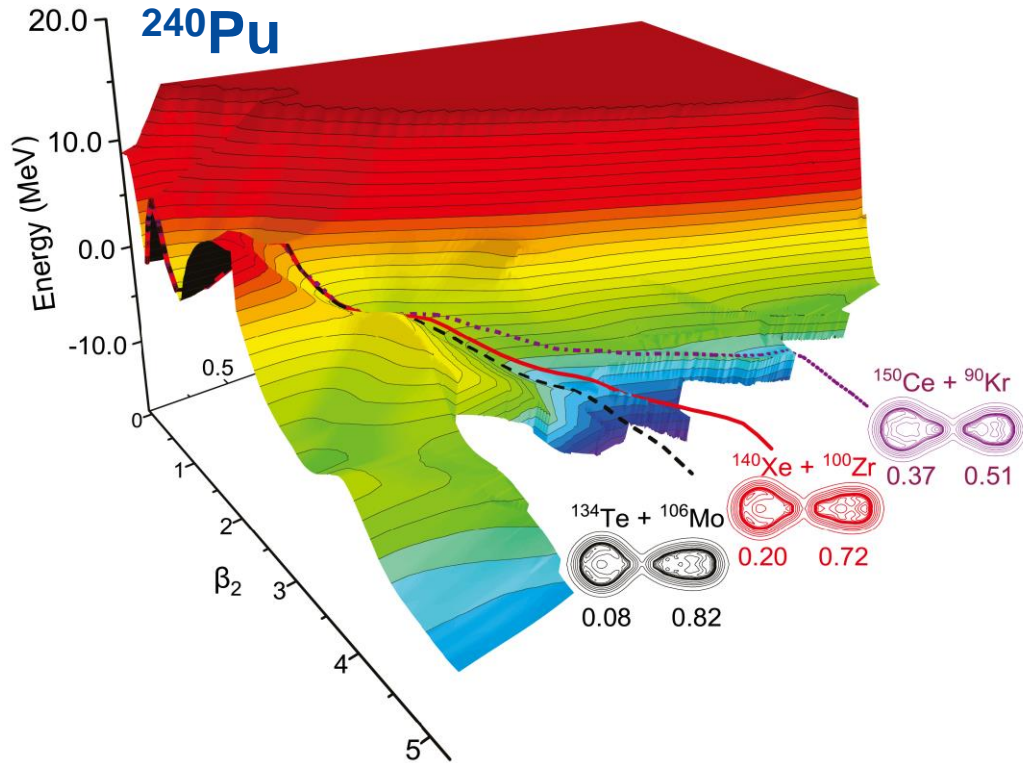
$$\hat{Q}_{2\mu}^i = D_{\mu 0}^2 q_2^i = D_{\mu 0}^2 \cdot \frac{3}{4\pi} A_i R_0^2 \beta_2^i$$

➤ Wave function of the system:

$$\sum_n \phi_n(\beta_2, \beta_3, t) |S_L S_H (S_{LH}) S_\Lambda; JM\rangle \quad n \equiv (S_L, S_H, S_\Lambda)$$

Scheme of **TD-MRV** model

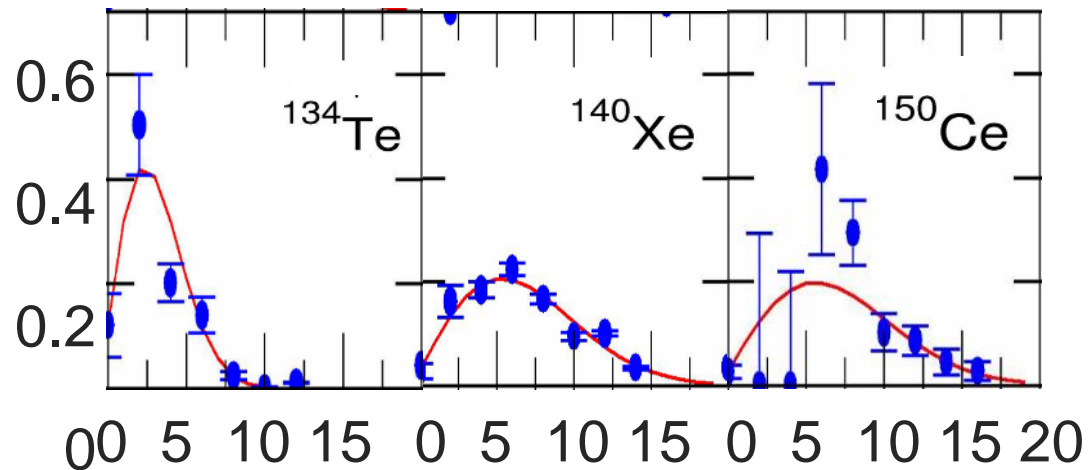
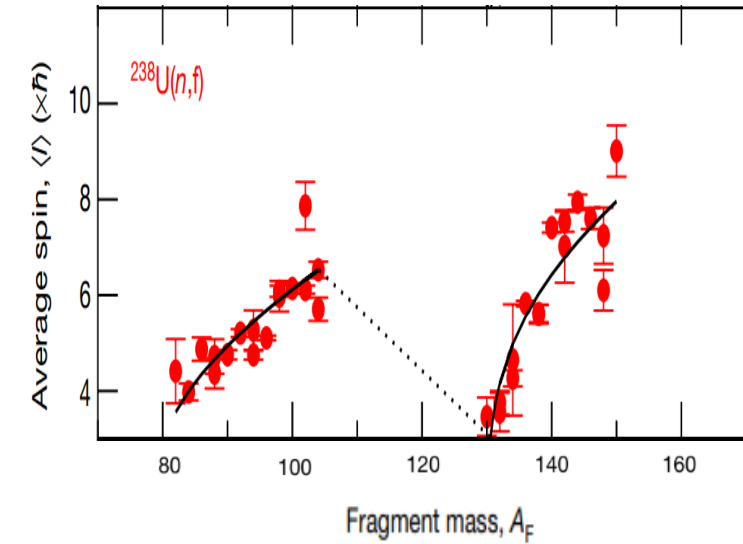
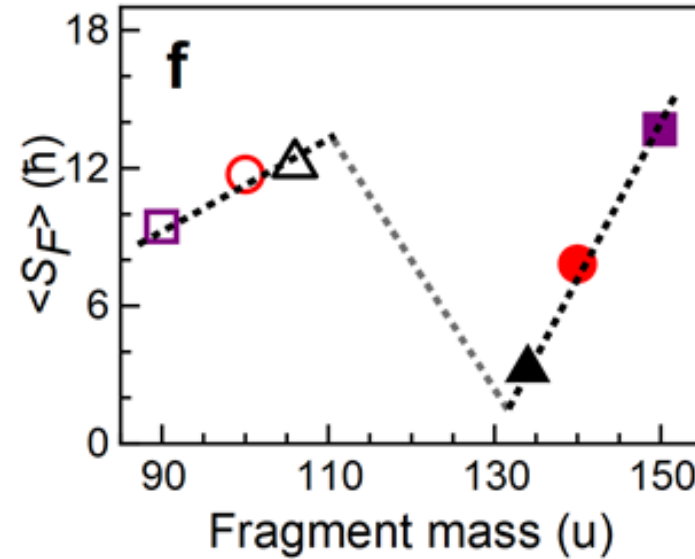
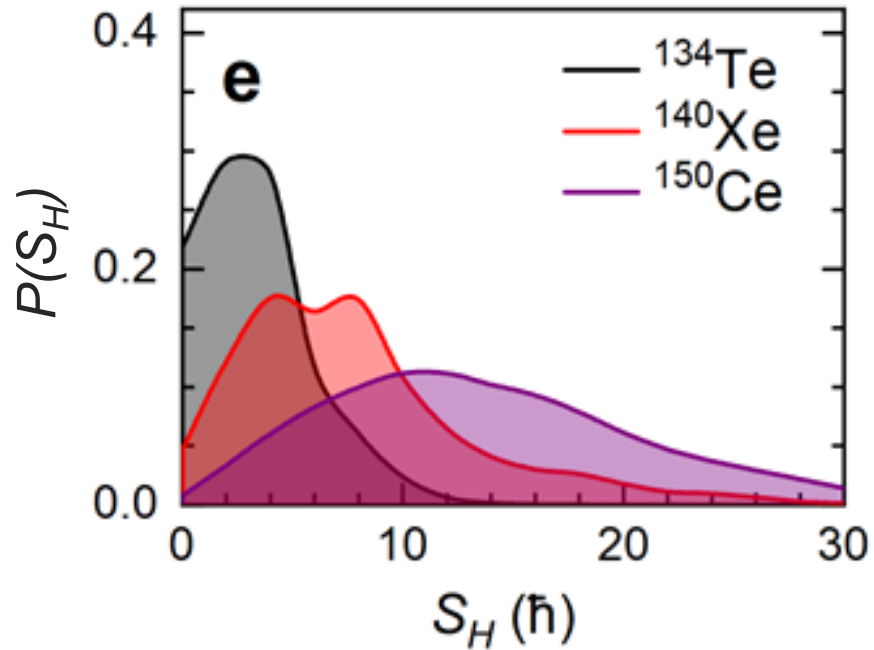
Evolution of Fragment Angular Momenta (FAM)



- From saddle to scission: **Rapid generation** and **chaotic evolution** of FAM
- After scission: **Vary slightly till to a stable pattern** when the fragments are well separated.

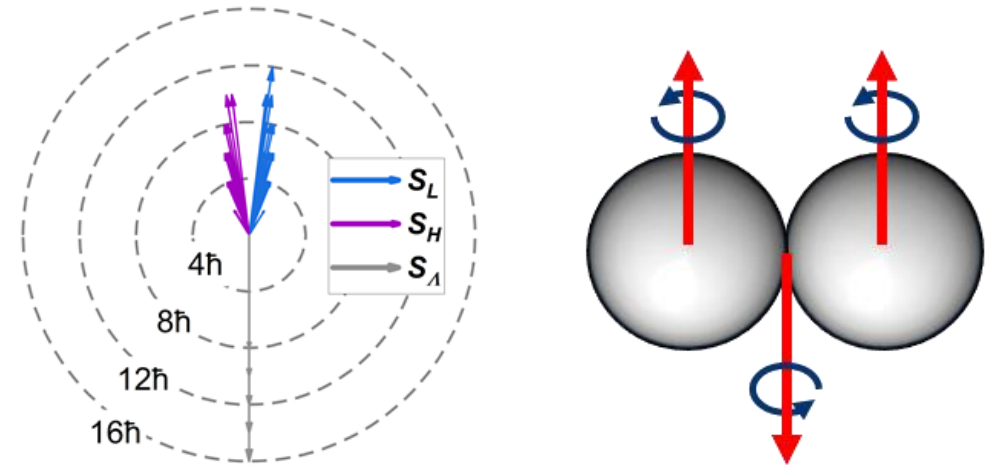
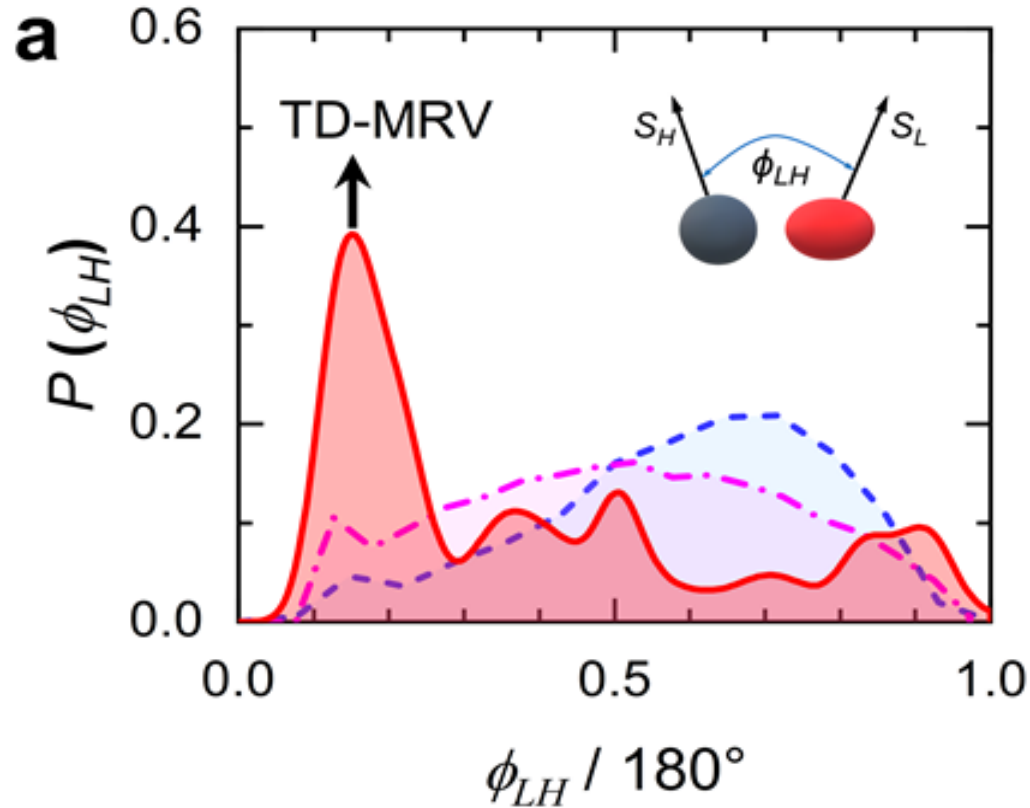
Distributions of FAM

J. N. Wilson, et al., Nature 590, 566 (2021)



- With the increase of A_F , the $P(S_H)$ increases rapidly: **2 \hbar ~10 \hbar** .
- The mass dependence of average FAM presents a clear **sawtooth-like pattern**.

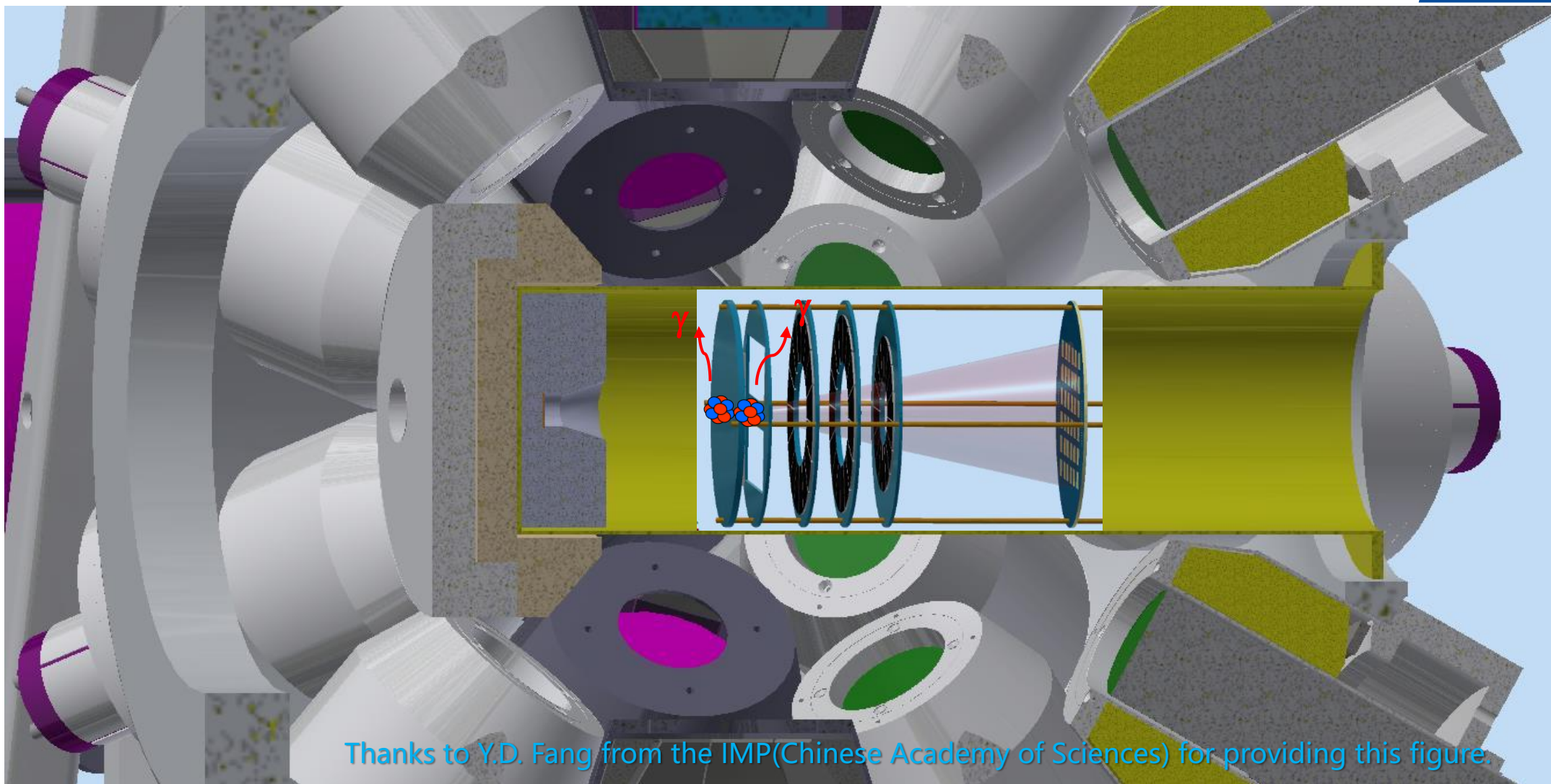
Zhou, Chen, Li, Smith & ZPLi, arXiv:2311.06177



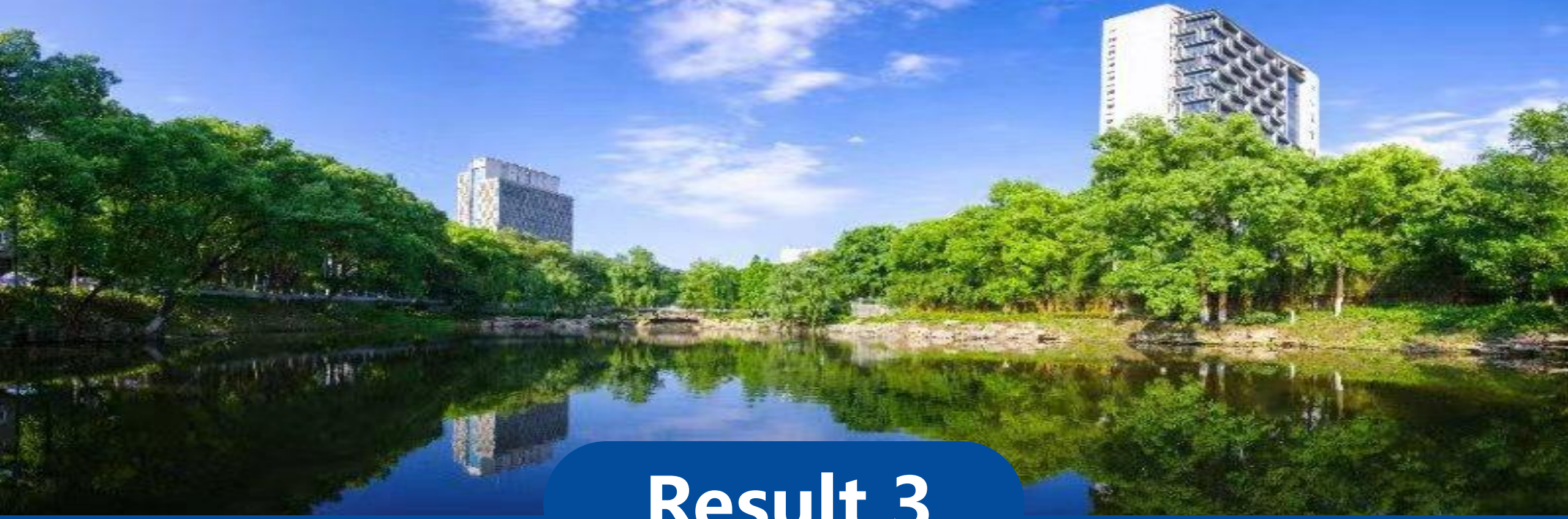
Dominated by **Wriggling mode**

- Opening angle displays strong correlations at **$\sim 30^\circ$, 90° , and 160°** .
- Differ from those of AMP (blue) and statistical model (magenta), which **did not consider quantum shape fluctuations**.

Measurement of correlations of FAM



We have carried out the experimental measurements for correlations of γ related to FAM from ^{252}Cf source in IMP, Lanzhou, China.



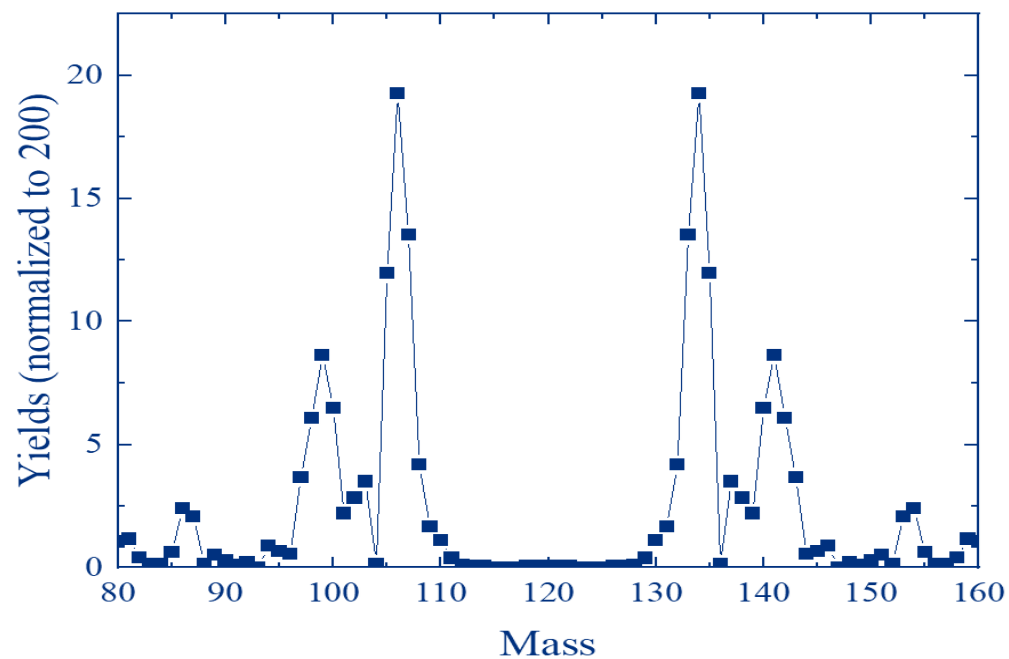
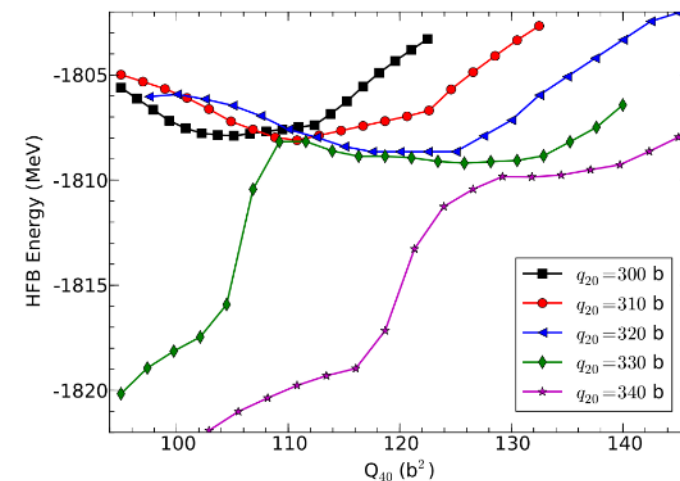
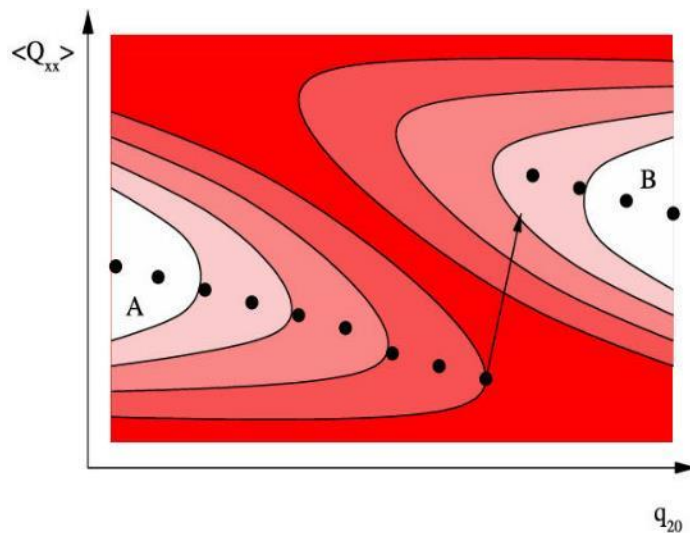
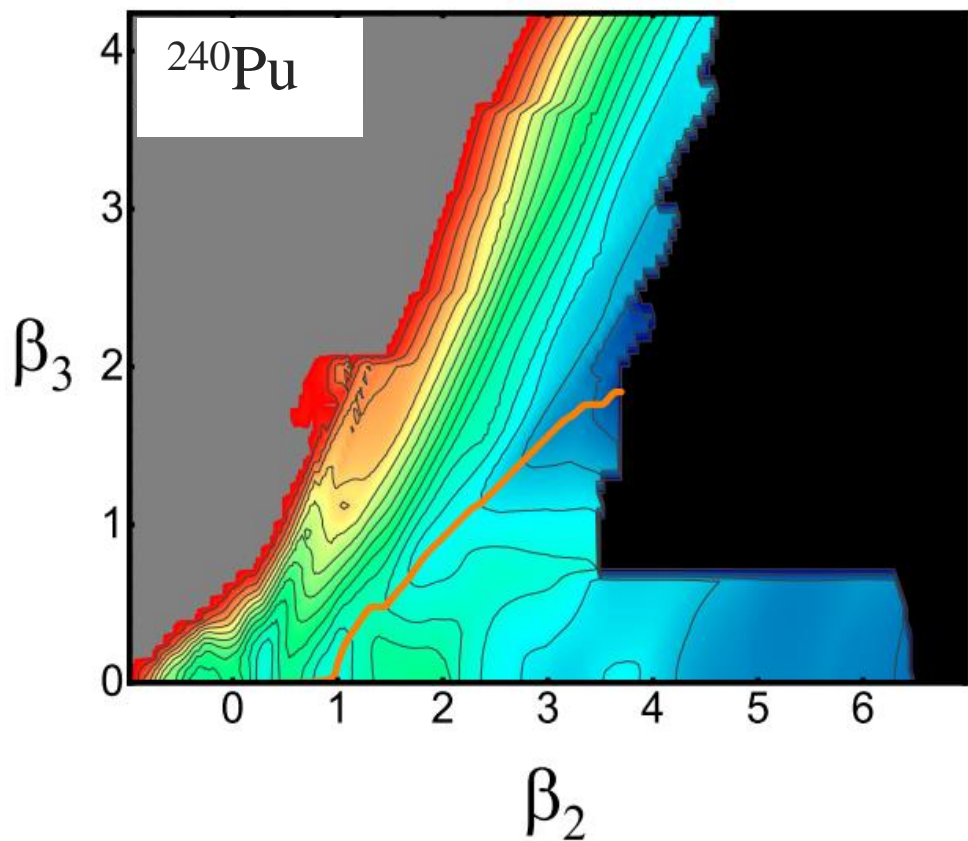
Result 3

Multi-dimensional collective space



Limitation of 2D space

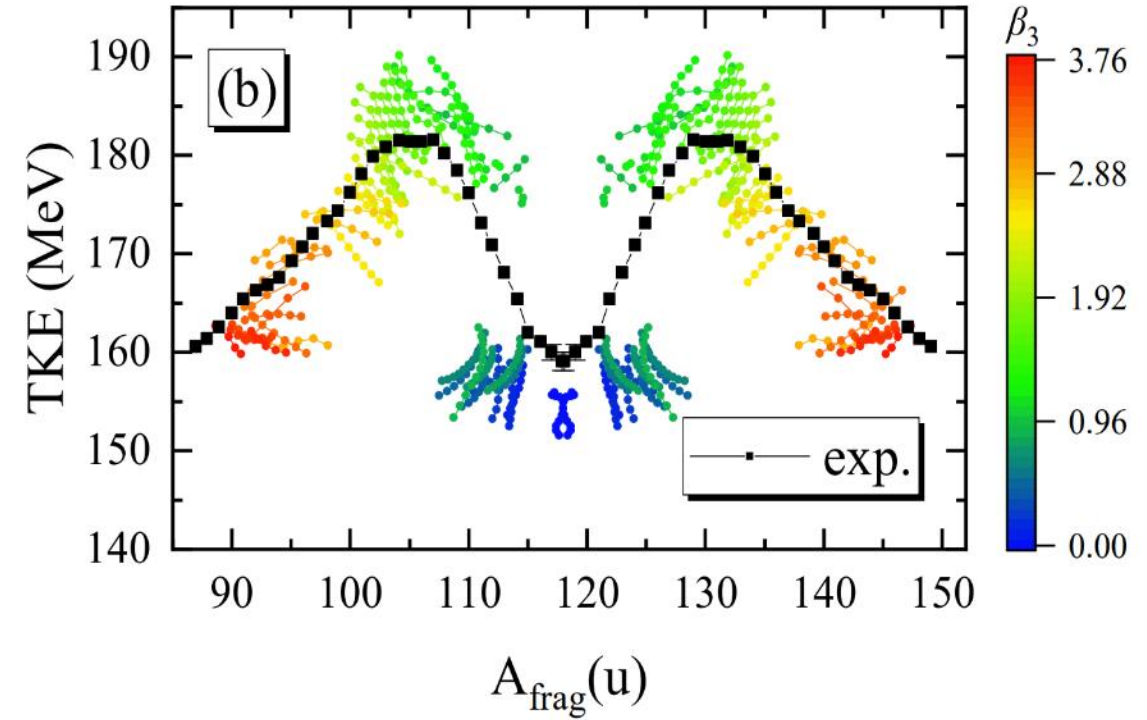
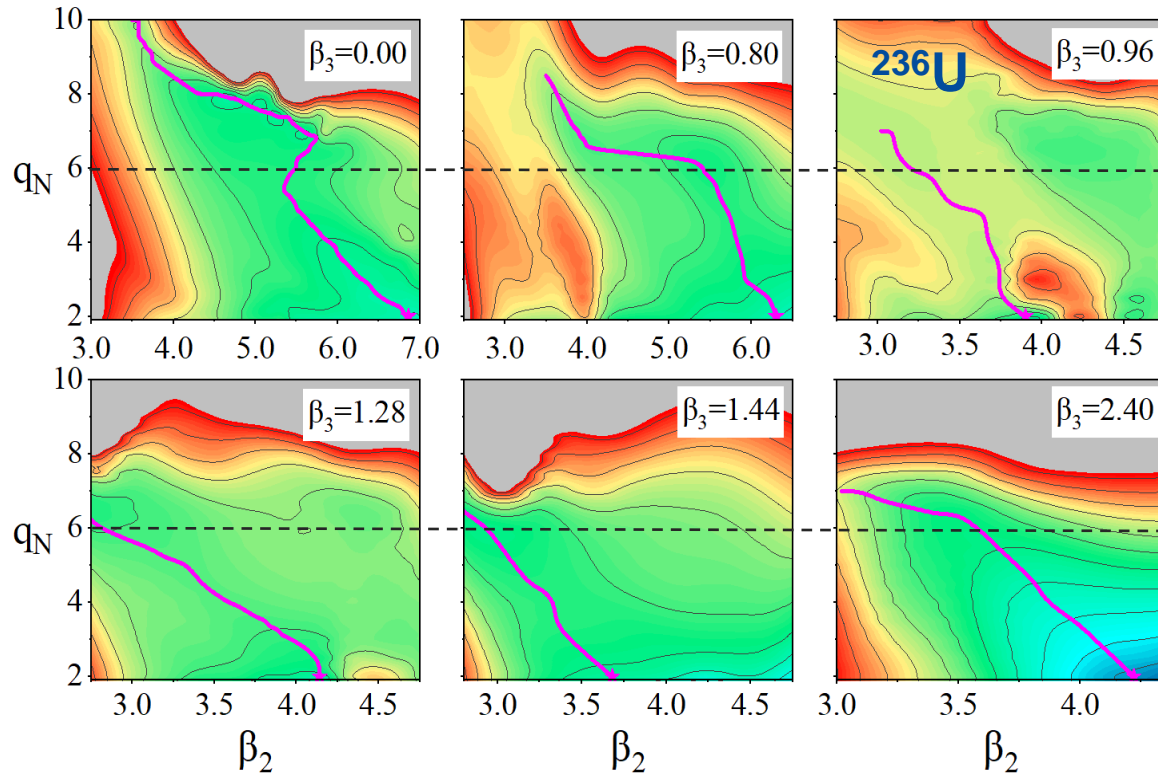
Discontinuity!



Dubray, Noël & Regnier, CPC2012;
Zdeb, Warda & Robledo, PRC2021;

.....

3D PES with nucleon number at neck (q_N)



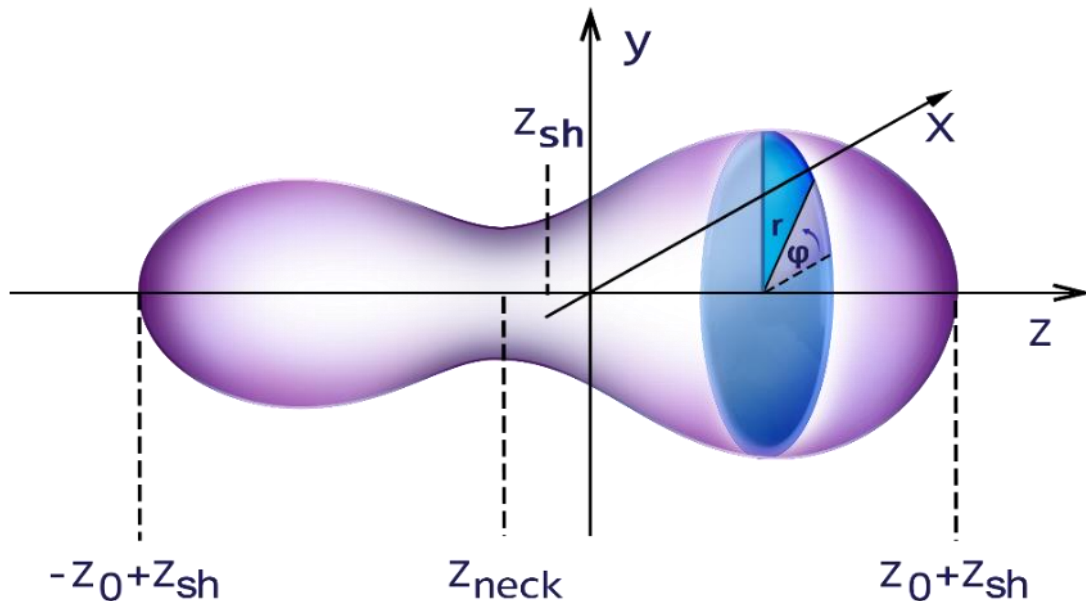
- $q_N < 6$, PES is very flat: large fluctuations
- Coexistence of fission modes

- Few to ~10 MeV fluctuation in TKE
- Few to ~10 nucleons fluctuation in A_{frag}

q_N cannot be defined globally!!

Zhou, Li, Chen, Chen & ZPLi, Chin. Phys. C (2023)

Nuclear Shapes in Fourier Expansion



➤ Expand $r(z, \varphi)$ instead of $R(\theta, \varphi)$

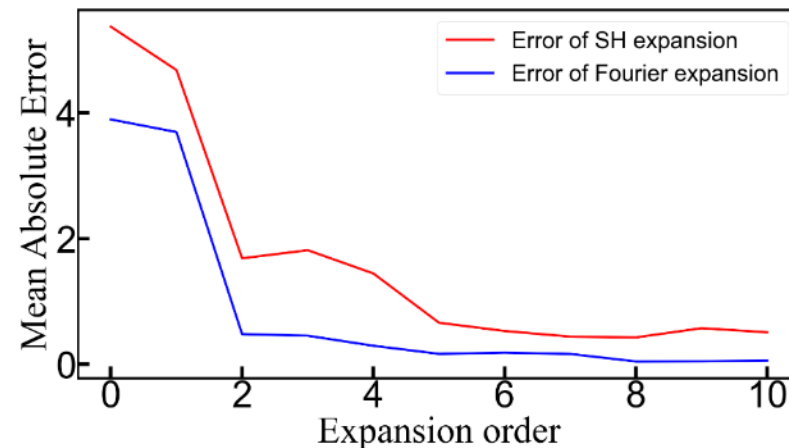
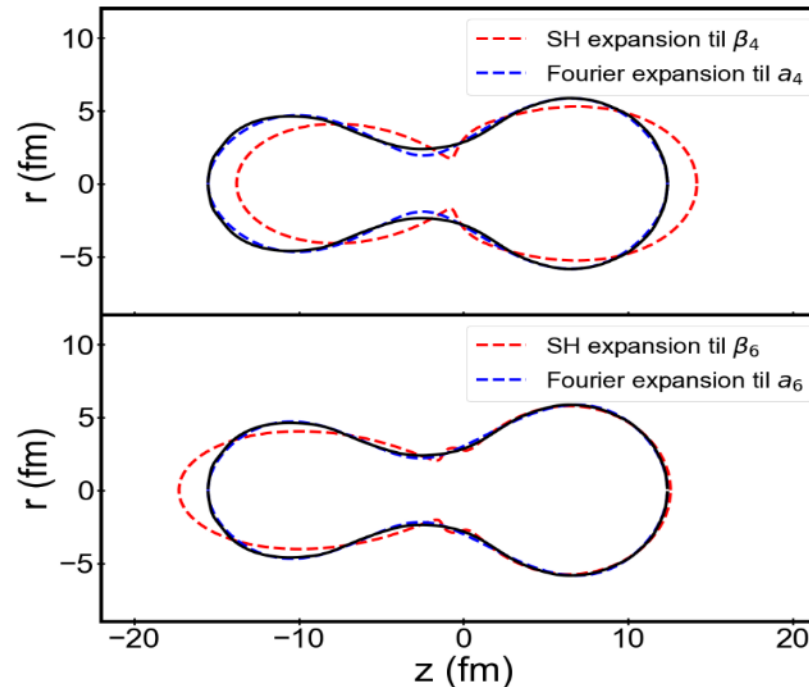
$$r^2(z) = R_0^2 \sum_{n=1}^{\infty} \left[a_{2n} \cos\left(\frac{(2n-1)\pi}{2} \frac{z - z_{sh}}{z_0}\right) + a_{2n+1} \sin\left(\frac{2n\pi}{2} \frac{z - z_{sh}}{z_0}\right) \right]$$

z_0 : Half length of atomic nucleus along z axis

z_{sh} : The geometric center along the z axis

K. Pomorski et al. Acta Phys. Pol. B Supl. 13, 361(2020)

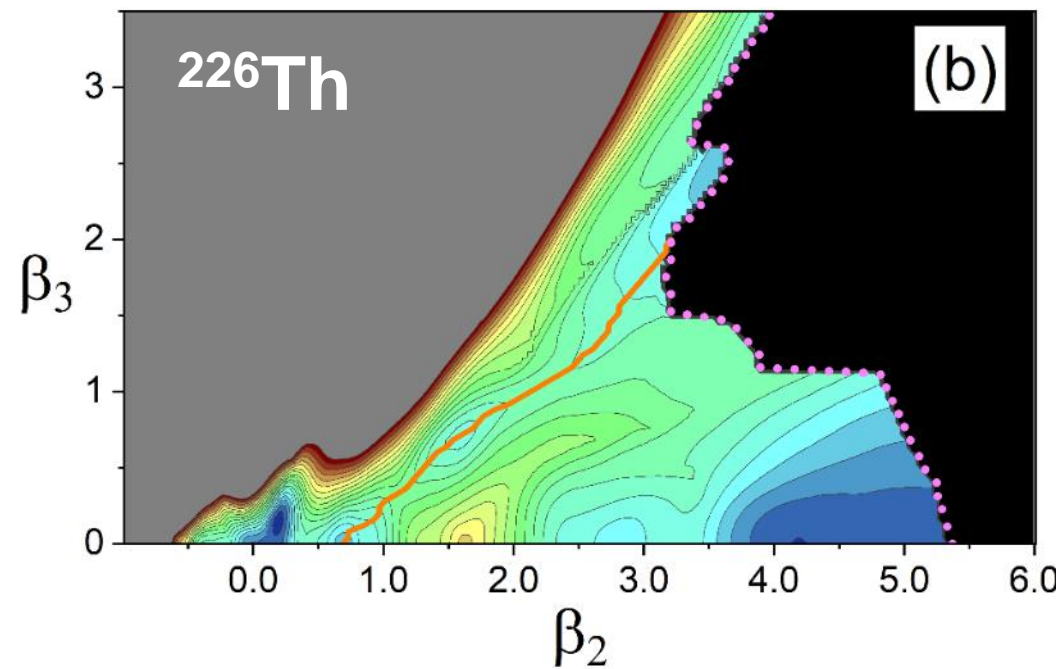
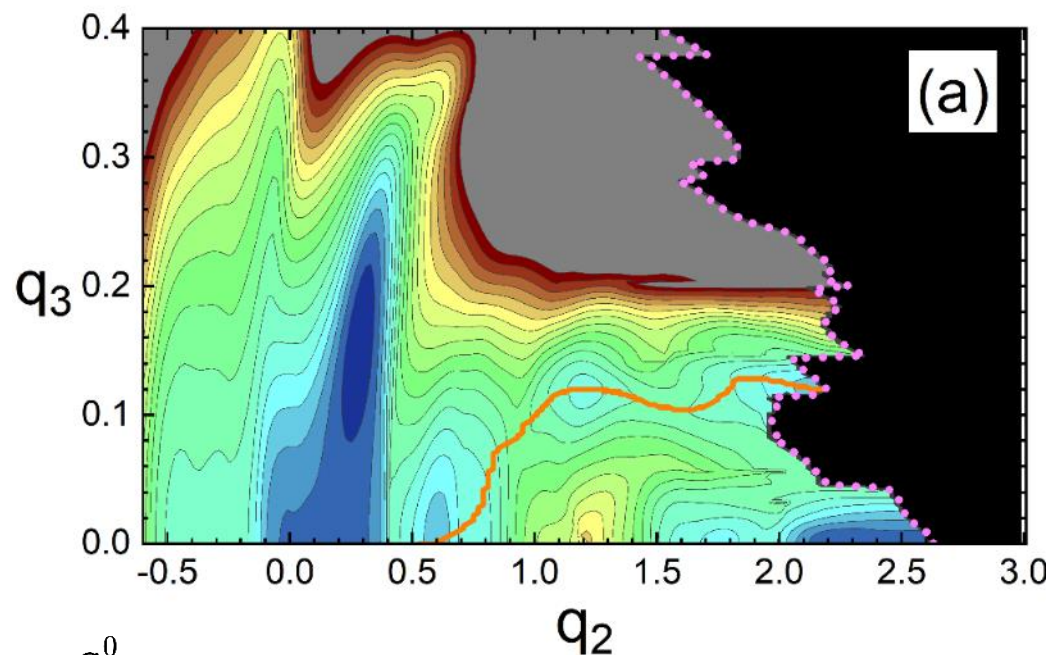
➤ Expansion convergency



$$a_{2n} = \int \frac{4R_0}{3Az_0} \cos\left(\frac{(2n-1)\pi(z-z_{sh})}{2z_0}\right) \rho_{RDFT}(r,z) r dr dz$$

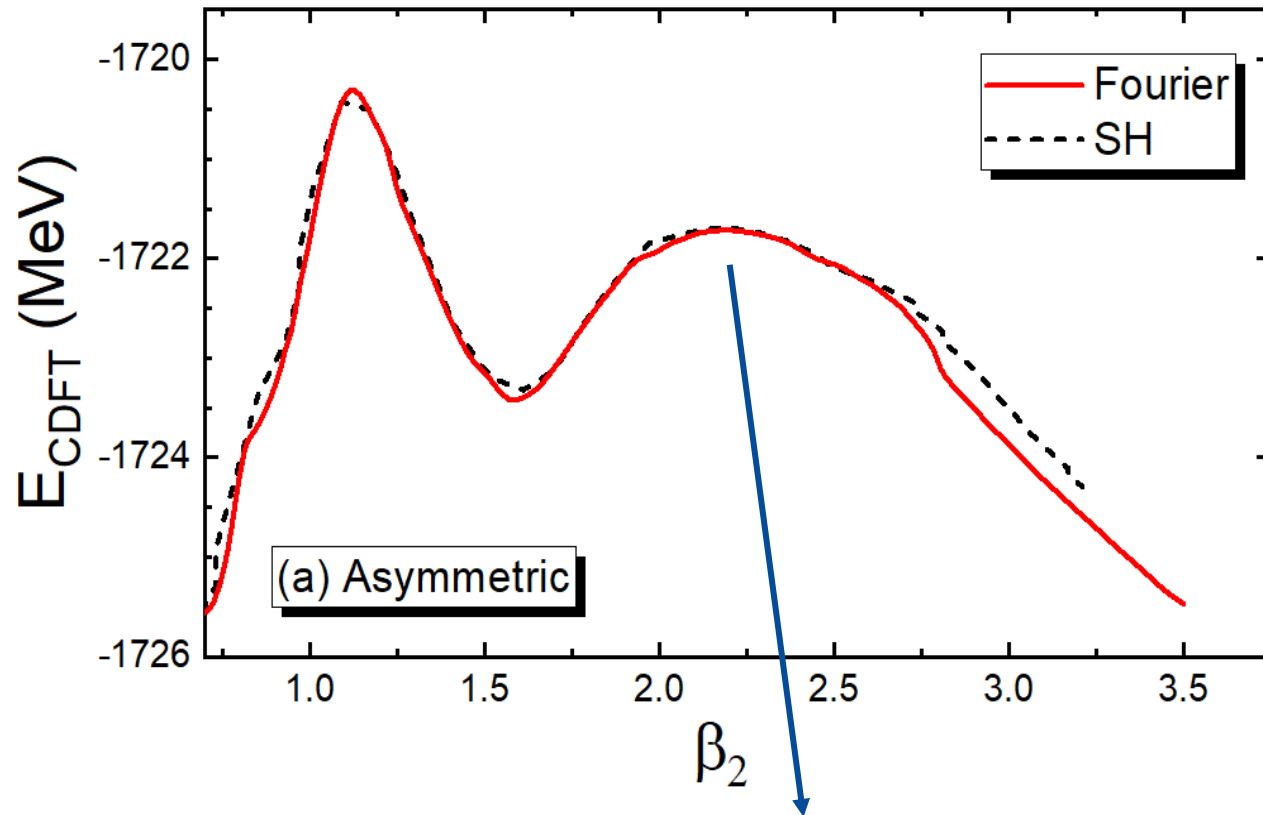
$$a_{2n+1} = \int \frac{4R_0}{3Az_0} \sin\left(\frac{n\pi(z-z_{sh})}{z_0}\right) \rho_{RDFT}(r,z) r dr dz$$

➤ Deformation operator



$$q_2 = \frac{a_2^0}{a_2} - \frac{a_2}{a_2^0} \quad q_3 = a_3$$

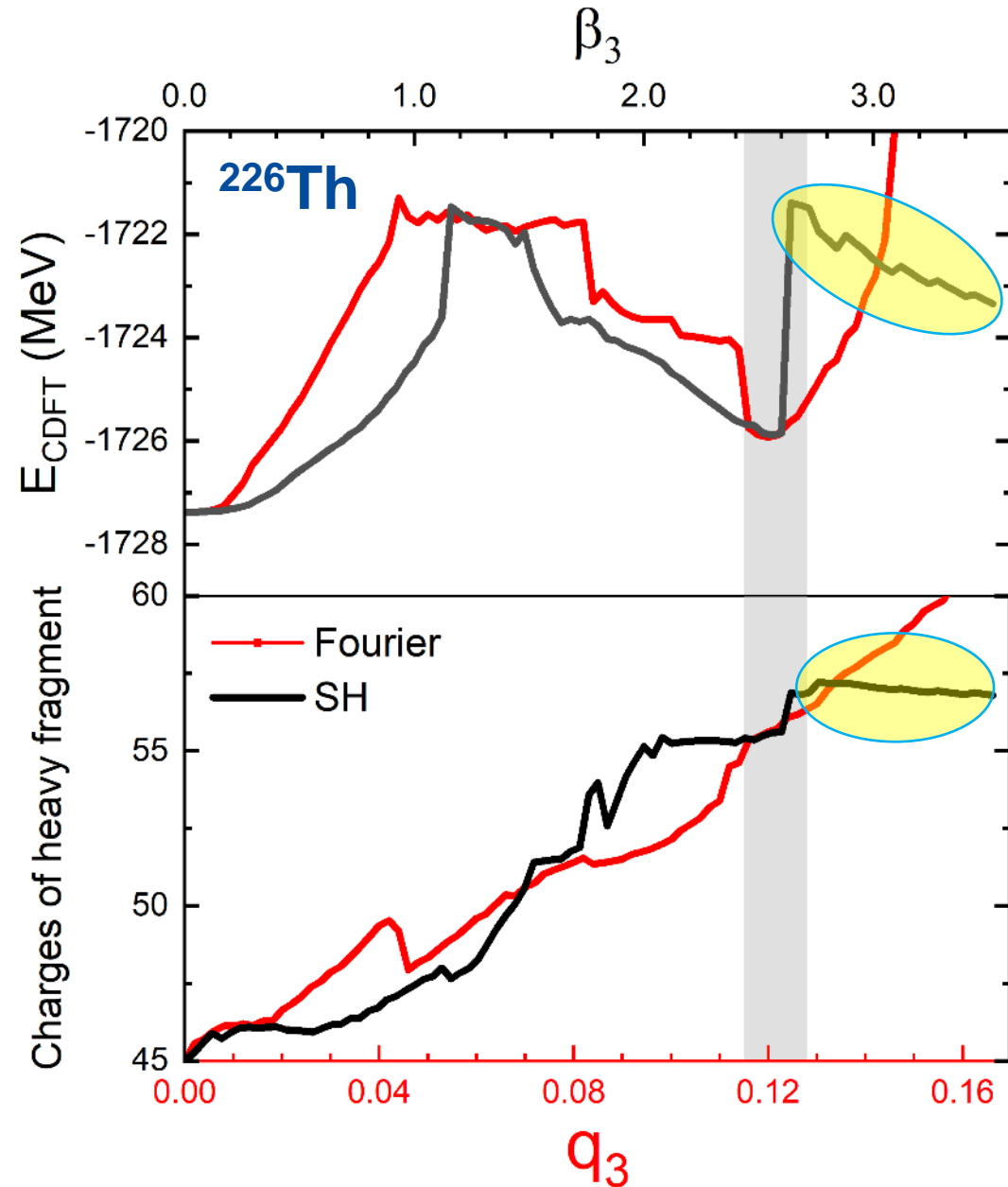
Convergency for expansion orders



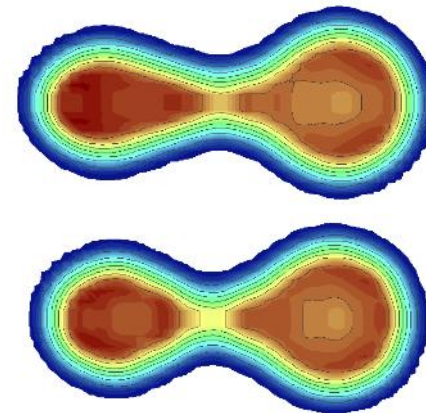
- Similar PECs before $\beta_2 = 2.5$
- More reasonable configurations for larger deformations

n	2	3	4	5	6	7
a_n	0.4865	0.1050	-0.2607	-0.0165	-0.0096	-0.0131
β_n	2.2329	-1.0275	3.5623	-3.1828	6.2295	-7.2577

Comparison of scission configurations



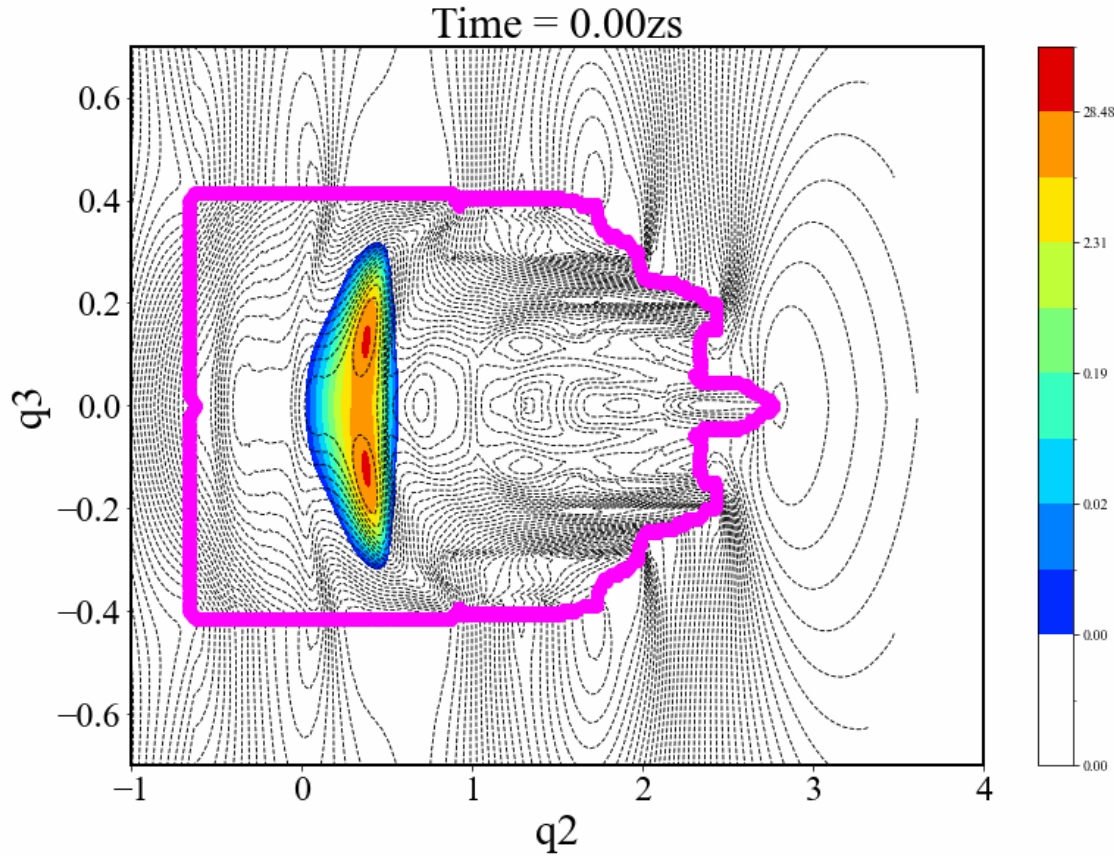
- Symmetric and asymmetric fission valleys are consistent
- **Nearly linear relation between Z_F (A_F) and q_3** , but not β_3 : Jumps and plateaus
- Not so reasonable in large (β_2, β_3).



Spherical Harmonic

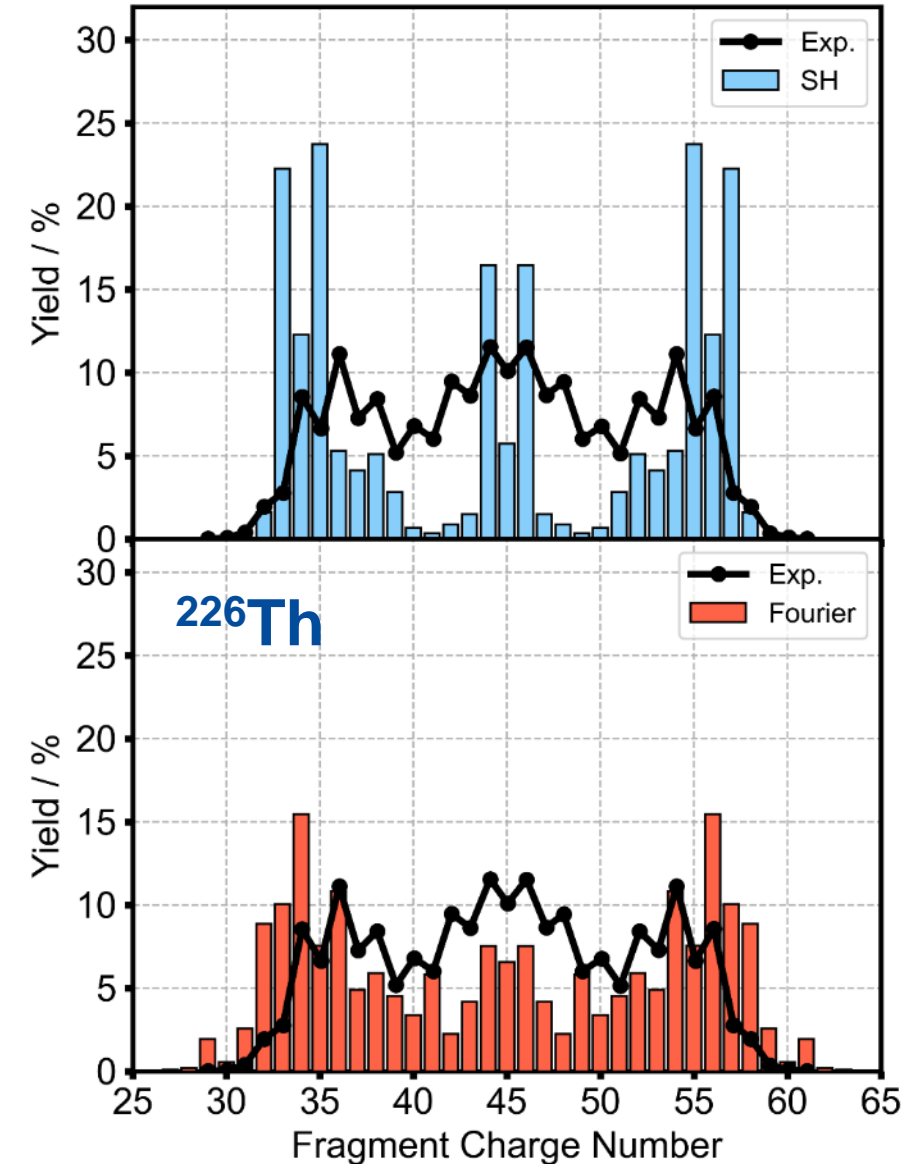
Fourier

Fission fragment charge distribution

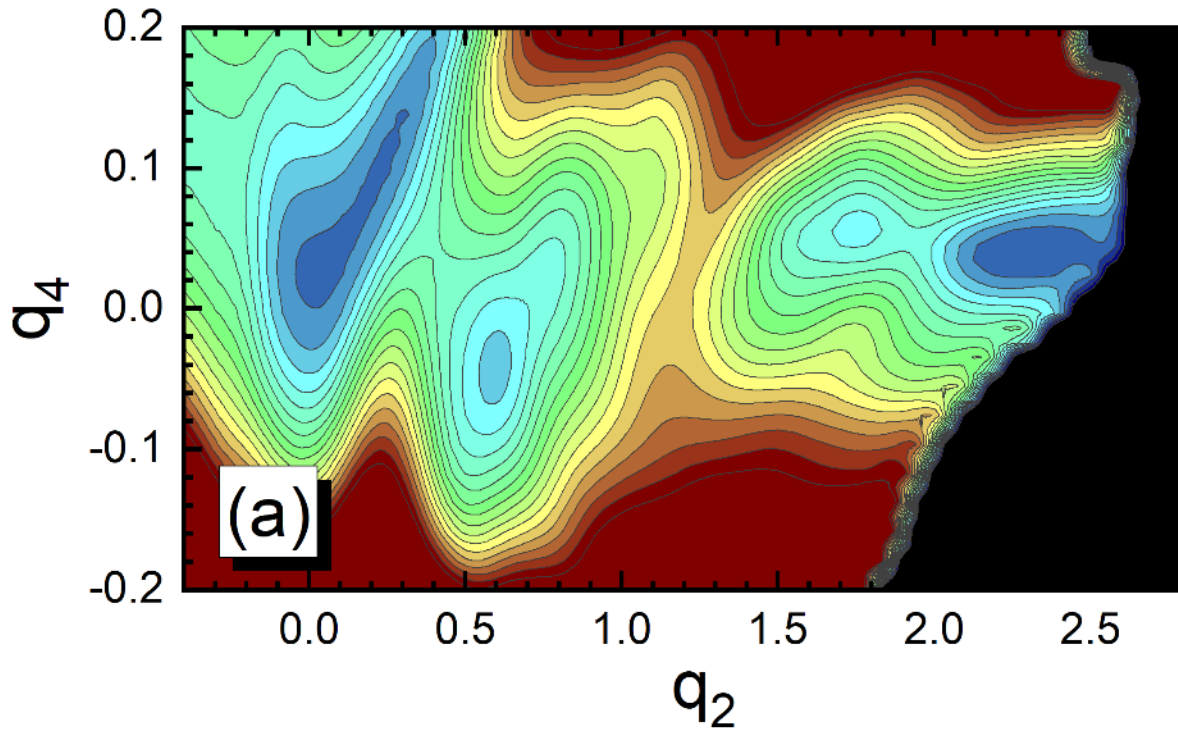


- **More fluctuations** in fragment charge distribution.
- The peak height and position that still need to be improved in further 3D calculation.

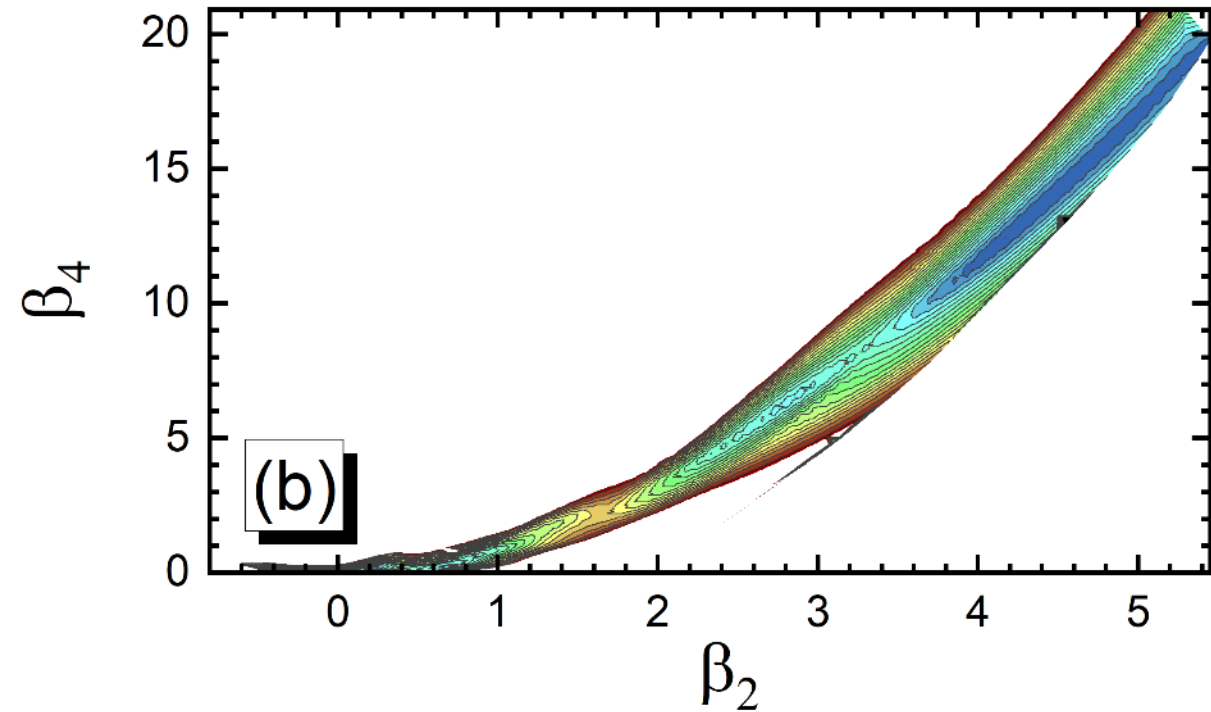
K.-H. Schmidt, et al. Nucl. Phys. A665, 221 (2001)



PES in (q_2, q_4) plane with $q_3 = 0$



PES in (β_2, β_4) plane with $\beta_3 = 0$



- PES in (q_2, q_4) plane is smooth; Scission configurations are easy to obtain
- While for (β_2, β_4) plane, β_4 has strong correlation with β_2 .

◆ Relativistic DFT based on TCHO basis

- Better accuracy and efficiency; possibility for post-scission

◆ Based on RDFT-TCHO & TDGCM+GOA

- Asymmetric fission of ^{180}Hg : $N=56$ octupole shell gap
- Sawtooth-like distribution of FAM; Strong correlation at $\varphi_{\text{LH}} \sim 30^\circ$
- Fourier shapes: fast convergence; better scission configurations ...

- **Outlook:** TDGCM based on 3D Fourier shapes; $Y(A, Z)$ & $P(S_L, S_H)$
Finite temperature; dissipation ...

Curriculum Vitae

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Thanks for your attention!

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Pengcheng Lab.

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