

Microscopic description of collective inertia and fission path for spontaneous fission

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Spontaneous fission—subbarrier fission

Collective inertia

Density functional theory for fission

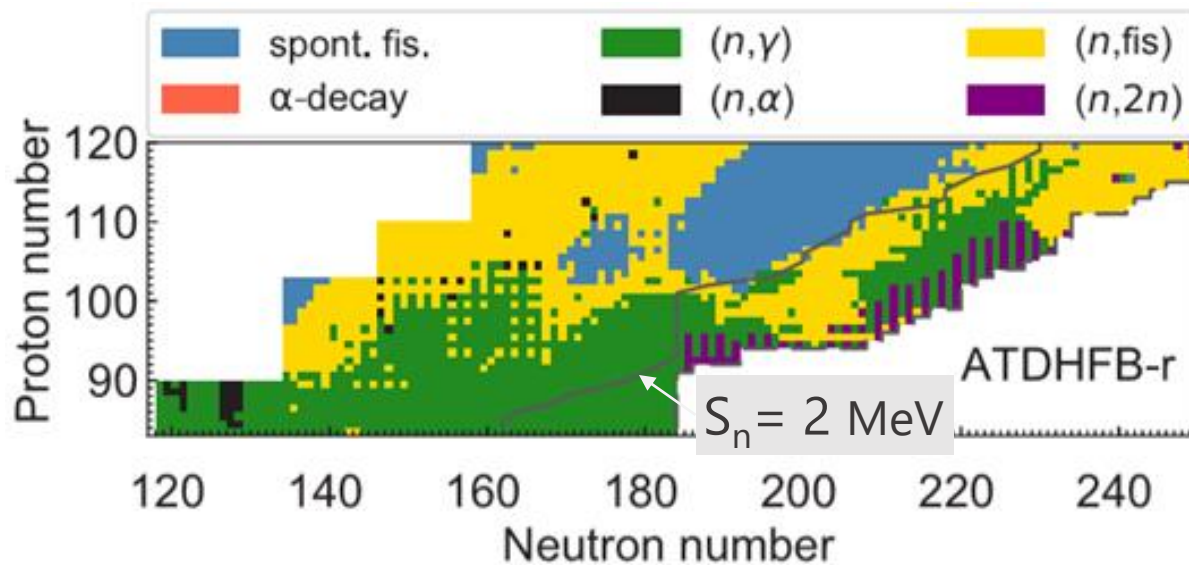
Result: Fission for one dimension

On-going: Fission for two dimensions

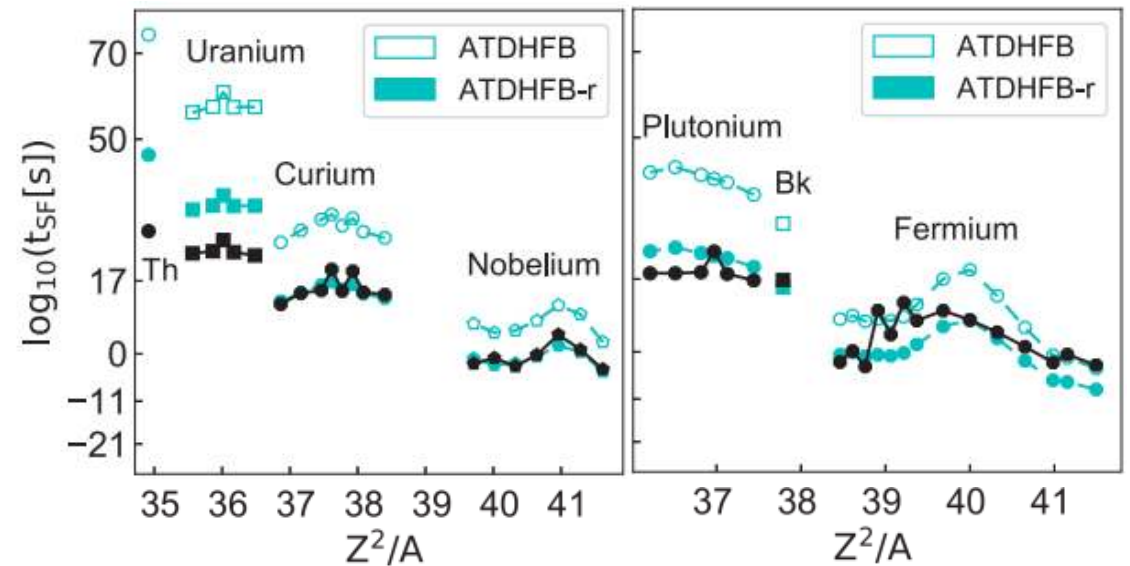
Nucleosynthesis network calculation under the r-process environment

($T = 0.9$ GK, $n_n = 1.0 \times 10^{28} \text{ cm}^{-3}$)

S.A. Giuliani, G. Martinez-Pinedo, L.M. Robledo, PRC 97, 034323 (2018)



Fission half-life



BCPM DFT

Spontaneous fission half-life

Neutron-induced reactions

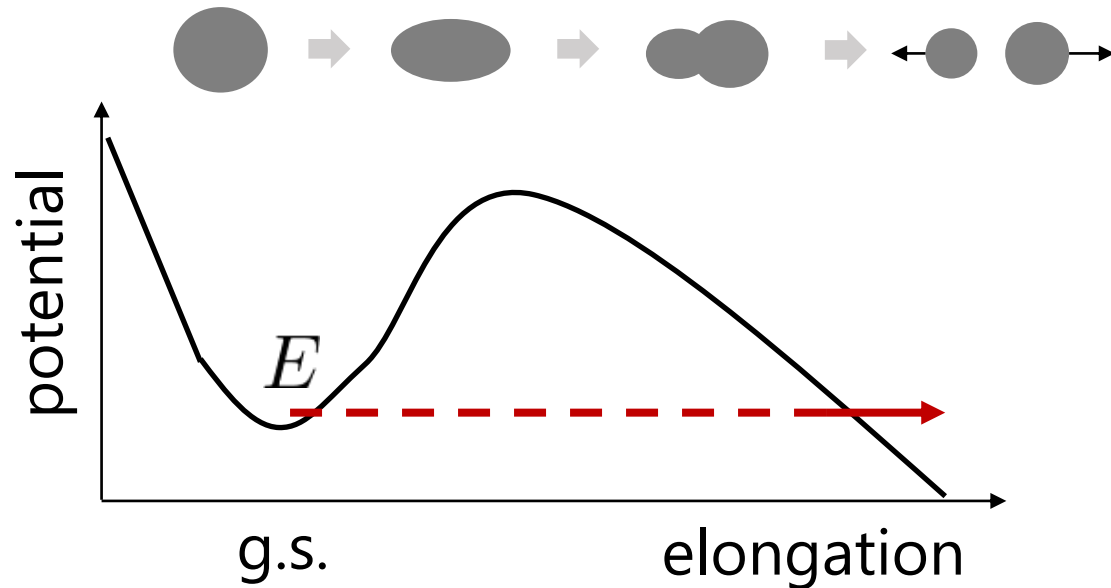
α decay rate

Correction to inertia

$$\mathcal{M}_r = \alpha \mathcal{M}$$

Sadhukhan et al, PRC88, 064314 (2013), etc.

Fission barrier



Spontaneous fission

➔ **Tunneling in quantum many-body system**

WKB approximation

Fission half-life

$$T_{1/2} = \ln 2 / (nP) \quad P = \frac{1}{1 + \exp(2S)}$$

Action S

$$S = \int_{s_1}^{s_2} ds \sqrt{2M(s)(V(s) - E)} / \hbar$$

s Coordinate (Q_{20} etc.)

V Potential

M **Collective inertia**

Fission observable

- Fission half-life
- Mass distribution



Collective inertia

characterizes the kinetic energy

Goal: correctly evaluate collective inertia M for fission half-life

$$S = \int_{s_1}^{s_2} ds \sqrt{2M(s)(V(s) - E)}/\hbar$$

- To develop the method for not only collective potential (fission barrier) but also **collective inertia** along a fission path
- To include **dynamical effects** of fission on collective inertia

Quantum many-body system

→ non-interacting reference system in one-body potential

DFT is an exact theory

Its functional form is unknown

Gogny, Skyrme, Covariant, etc.

$$\mathcal{E}(\rho, \kappa, \kappa^*)$$

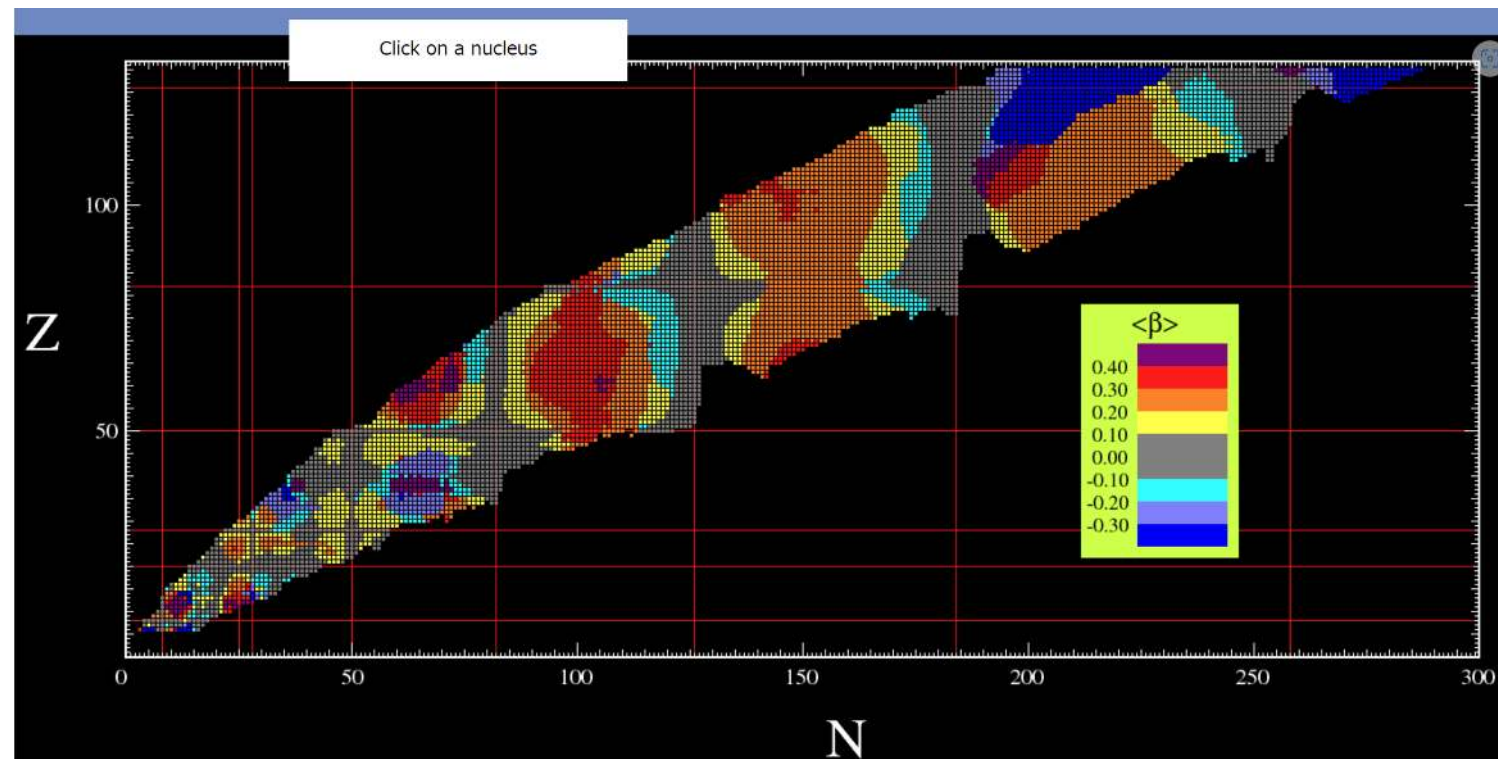
Good description for

- ✓ Ground state property
- ✓ Shell effect
- ✓ Deformation

Bender, Heenen, Reinhard, Rev. Mod. Phys. 75 (2003) 121

Nakatsukasa, Matsuyanagi, Matsuo, Yabana, RMP88(2016) 045004

https://www-phynu.cea.fr/science_en_ligne/carte_potentiels_microscopiques/carte_potentiel_nucleaire_eng.htm



DFT + Cranking approximation

Skyrme, Gogny, Relativistic EDFs

Low computation cost

Problem: Neglect dynamical effects (time-odd terms)

Prochniak et al., NPA730 (2004) 59
Delaroche et al., PRC81 (2010) 014303
Baran et al., PRC84, 054321 (2011)
Sadhukhan et al, PRC88, 064314 (2013)
Giuliani and Robledo, PLB 787, 134 (2018)

$$\mathcal{M}^{\text{PC}} = \frac{1}{2} [M^{(1)}]^{-1} M^{(3)} [M^{(1)}]^{-1}$$
$$M_{ij}^{(n)} = \sum_{\mu < \nu} \frac{\langle \phi(s) | \hat{s}_i | \mu\nu \rangle \langle \mu\nu | \hat{s}_j^\dagger | \phi(s) \rangle}{(E_\mu + E_\nu)^n}$$

Note: Dynamical effects reproduce the collective inertia for translational motion

Perturbative cranking approximation

$$\mathcal{M}^{\text{PC}} = \frac{1}{2} [M^{(1)}]^{-1} M^{(3)} [M^{(1)}]^{-1}$$

$$M_{ij}^{(n)} = \sum_{\mu < \nu} \frac{\langle \phi(s) | \hat{s}_i | \mu\nu \rangle \langle \mu\nu | \hat{s}_j^\dagger | \phi(s) \rangle}{(E_\mu + E_\nu)^n}$$

s : collective variables

ϕ : constrained HFB states

$$|\mu\nu\rangle = a_\mu^\dagger a_\nu^\dagger |\phi(s)\rangle$$

DFT + Cranking approximation

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Our method: Local QRPA

Hinohara et al., PRC84 (2011) 061302; 85 (2012) 024323
Sato, Hinohara, NPA849 (2011) 53
Yoshida, Hinohara, PRC83 (2011) 061302

Include dynamical effects by QRPA

High computation cost

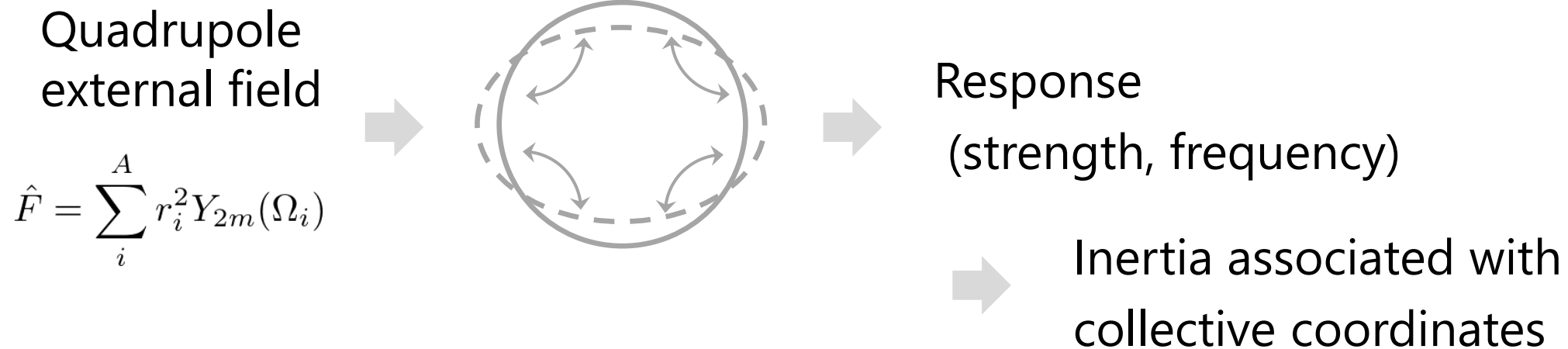


Finite amplitude method

Nakatsukasa et al., PRC76, 024318(2007)
Avogadro & Nakatsukasa, PRC84, 014314(2011)

$$T = \frac{1}{2} M \left(\frac{dQ_{20}}{dt} \right)^2 \quad Q_{\lambda\mu} = r^\lambda Y_{\lambda\mu} \quad \text{Collective surface vibrations}$$

QRPA: linear response to an external field



Local QRPA for vibrational mass at each CHFB state

Hinohara et al., PRC82 (2010) 064313

$$\delta\langle\phi(s)|[\hat{H}_M(s), \hat{Q}^i(s)] - \frac{1}{i}\hat{P}^i(s)|\phi(s)\rangle = 0$$

$$\delta\langle\phi(s)|[\hat{H}_M(s), \frac{1}{i}\hat{P}^i(s)] - C_i(s)\hat{Q}^i(s)|\phi(s)\rangle = 0$$

s : deformation parameters

$$\begin{pmatrix} A & B \\ B^* & A^* \end{pmatrix} \begin{pmatrix} Q^i \\ -Q^{i*} \end{pmatrix} = \frac{1}{i} \begin{pmatrix} P^i \\ -P^{i*} \end{pmatrix}$$

$$\begin{pmatrix} A & B \\ B^* & A^* \end{pmatrix} \begin{pmatrix} P^i \\ -P^{i*} \end{pmatrix} = iC_i \begin{pmatrix} Q^i \\ -Q^{i*} \end{pmatrix}$$

at deformation s

- ➔ Low-lying collective modes
- Eigen-frequency $\hat{Q}^i, \hat{P}^i, C_i = \Omega_i^2$
- ➔ $M(s)$ **Collective inertia for quadrupole vibration**

$$\frac{\partial s_m}{\partial q^i} = \langle\phi(s)|[\hat{s}_m, \frac{1}{i}\hat{P}_i]|\phi(s)\rangle$$

$$M(s) = \frac{\partial q^1}{\partial s_1} \frac{\partial q^1}{\partial s_1} \quad s_1 = r^2 Y_{20}$$

Finite amplitude method (FAM)

Nakatsukasa et al., PRC76 (2007) 024318
 Avogadro & Nakatsukasa, PRC84, 014314

- ✓ Small computational cost
- ✓ equivalent to QRPA response

Constrained HFB along the fission path

Constrained on Q_{20}

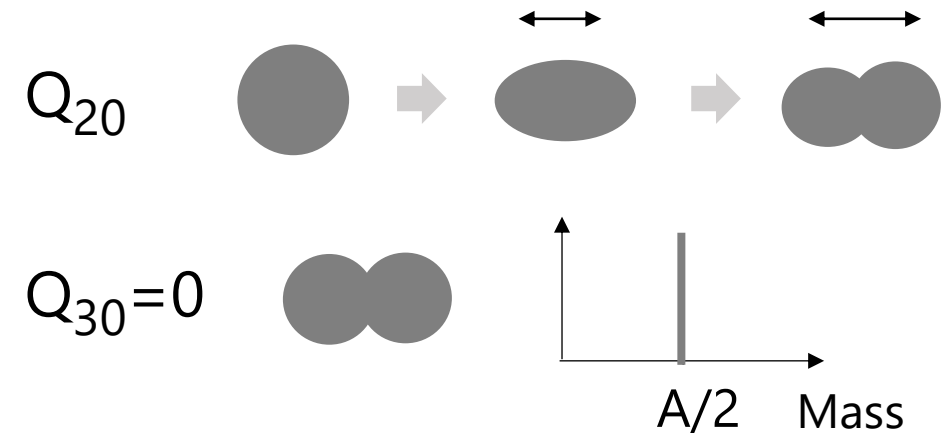
Fix $Q_{30}=0$

→ symmetric fission path only

Box: 26 fm x 26 fm x 39 fm

Volume pairing, SkM* EDF

$E_{QP} \approx 60$ MeV



Local QRPA on constrained HFB states

Select the most collective mode among QRPA solutions

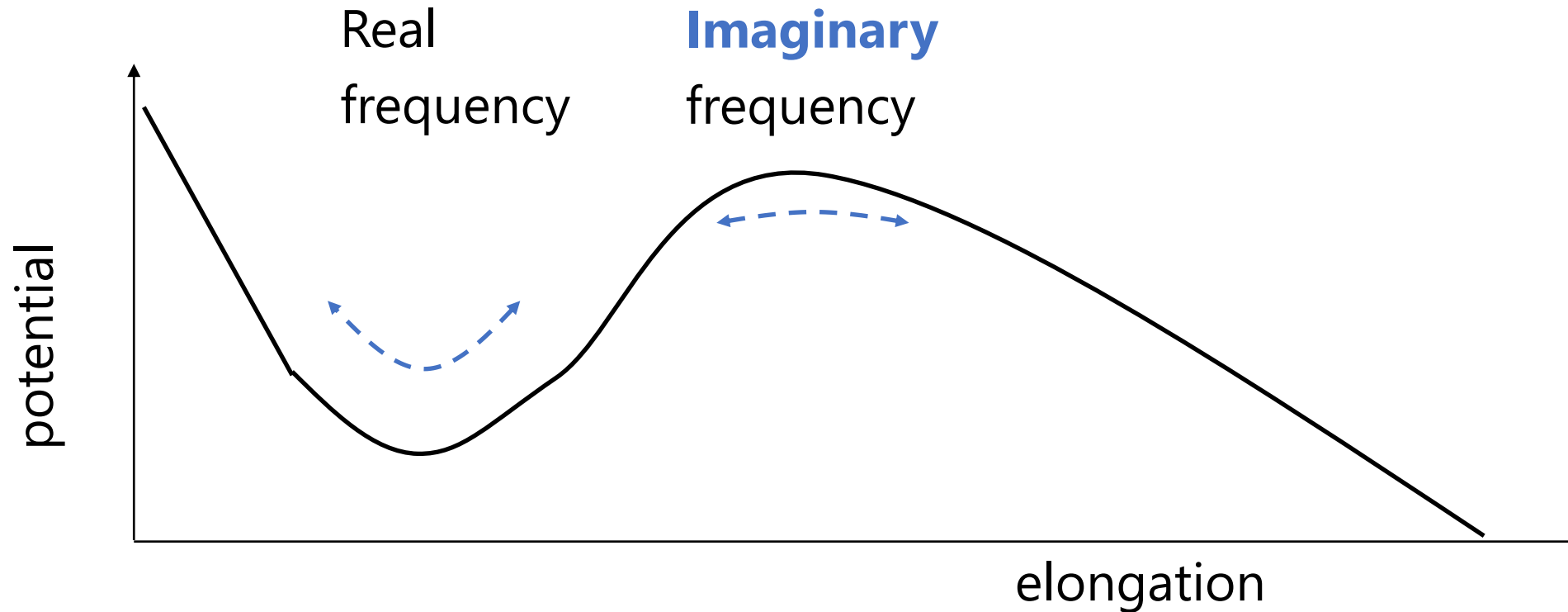
Washiyama, Hinohara, Nakatsukasa, PRC103, 014306 (2021)

Computation time: (1000 hours+500 hours) x 50 states

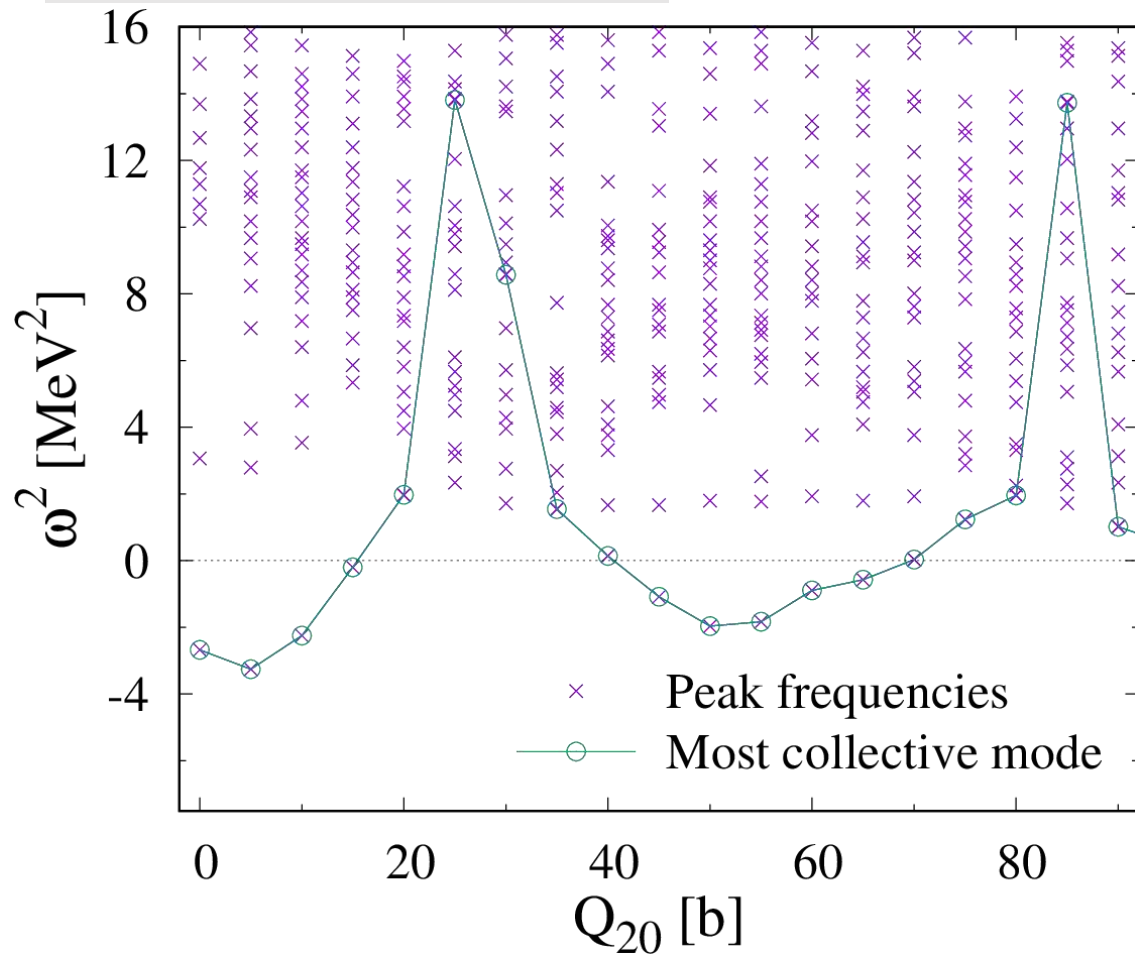
This research used computational resources of OFP & Wisteria/BDEC-01 Odyssey (Univ. Tokyo), provided by the Multidisciplinary Cooperative Research Program in CCS, Univ. Tsukuba.

Local QRPA = QRPA on top of constrained DFT state

Eigen-frequency \longleftrightarrow Curvature of the potential



^{240}Pu , axial states

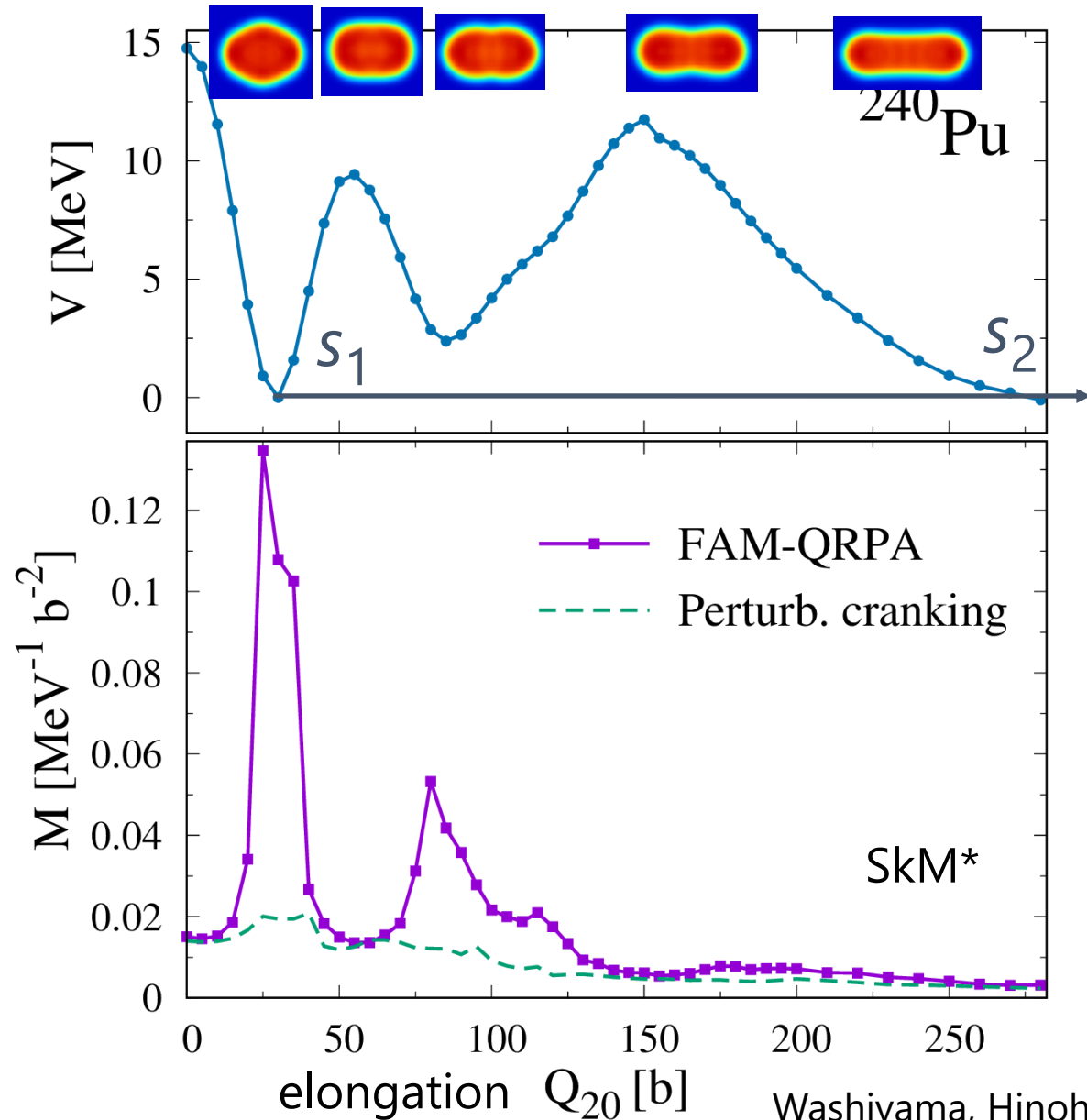


Select the most collective mode from many QRPA solutions in $\omega^2 < 16 \text{ MeV}^2$



Select the smallest collective inertia M = the largest strength

Result: Collective inertia along fission path



- ✓ Mass symmetric fission
 - ➔ Too high second barrier
- ✓ Large change in M by QRPA
- ✓ Smooth in cranking M

Action integral S

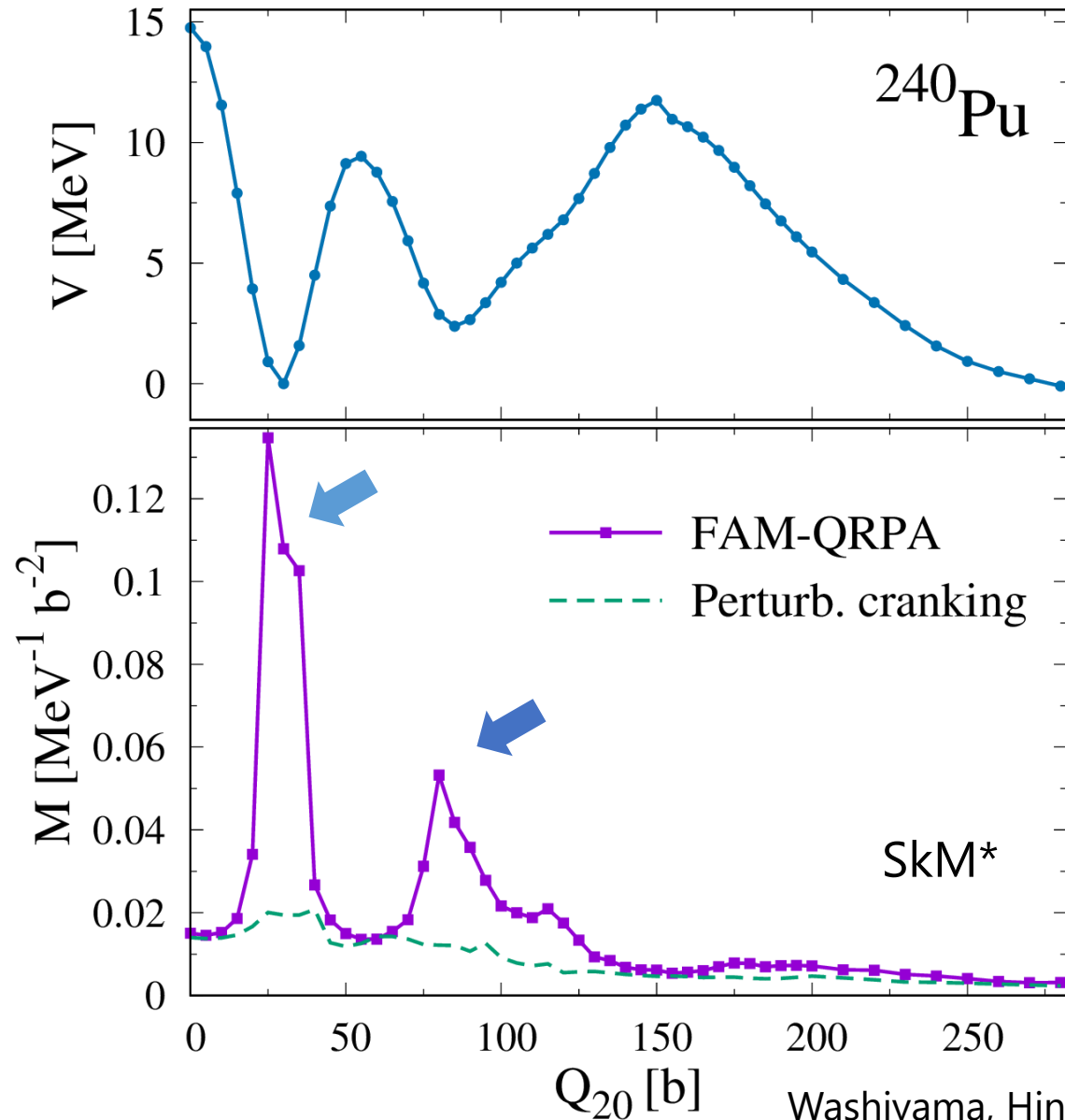
$$S = \int_{s_1}^{s_2} ds \sqrt{2M(s)(V(s) - E)}/\hbar$$

QRPA $S = 82.0$

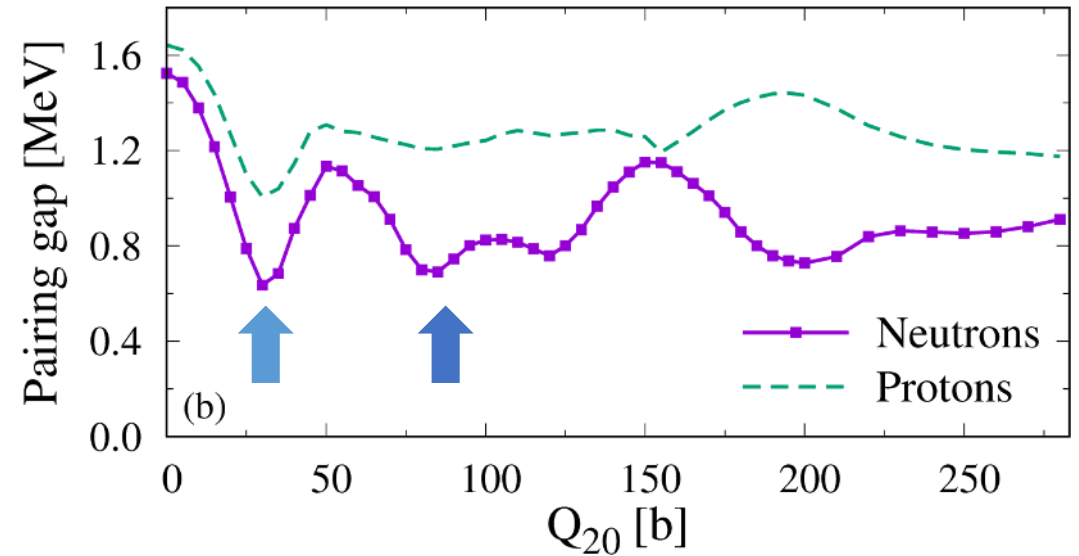
Cranking $S = 62.0$

↷ M

60000 CPU hours
OpenMP + MPI



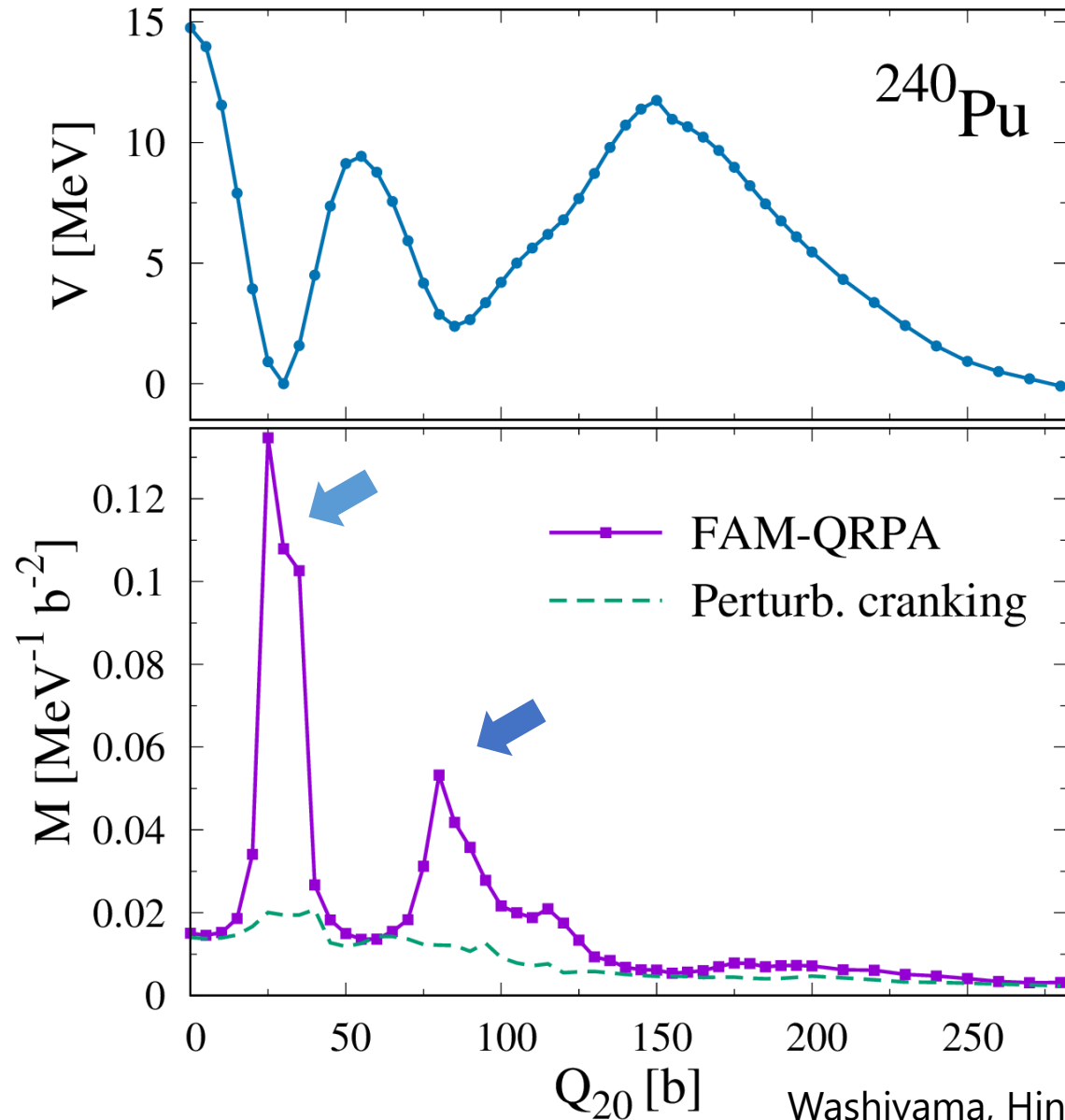
Pairing gap



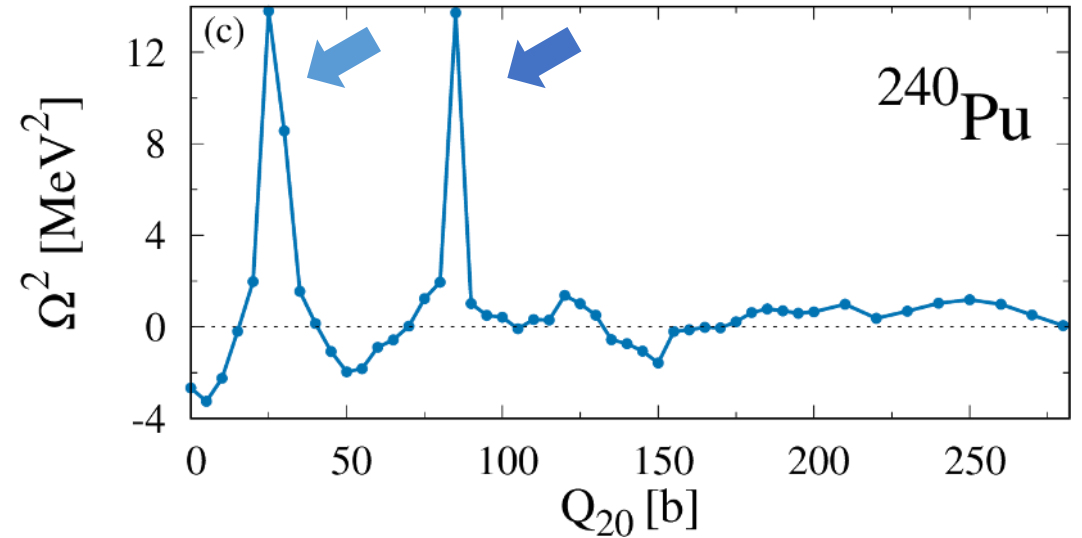
Pair gap ↘ ↔ **Inertia** ↗

$$\mathcal{M} \propto \Delta^{-2}$$

Inertia pairing gap

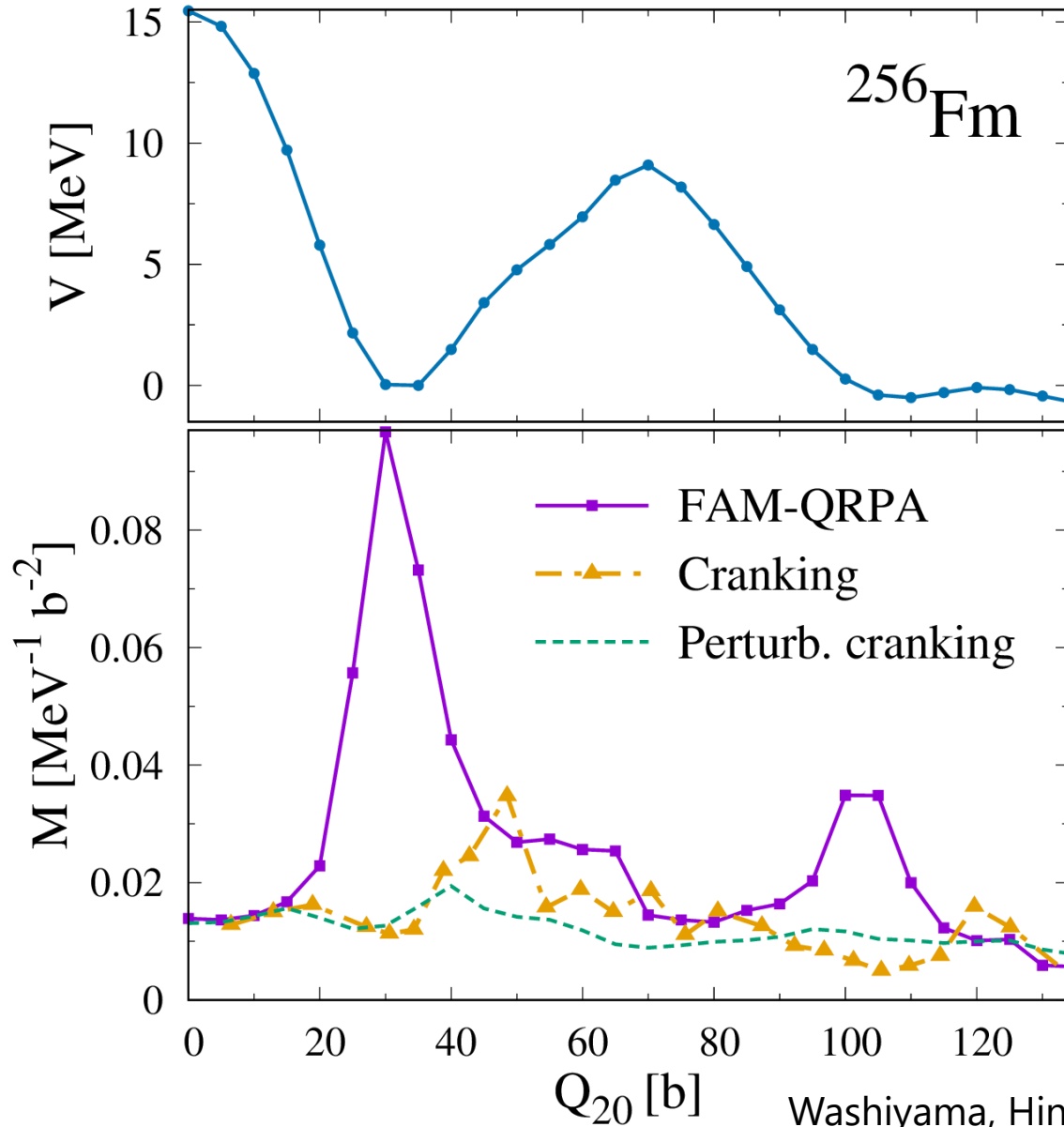


Eigen-frequency of LQRPA



$\Omega^2 \nearrow \longleftrightarrow$ **Inertia** \nearrow

Ω can be imaginary near the fission barrier



Non-perturbative cranking approximation

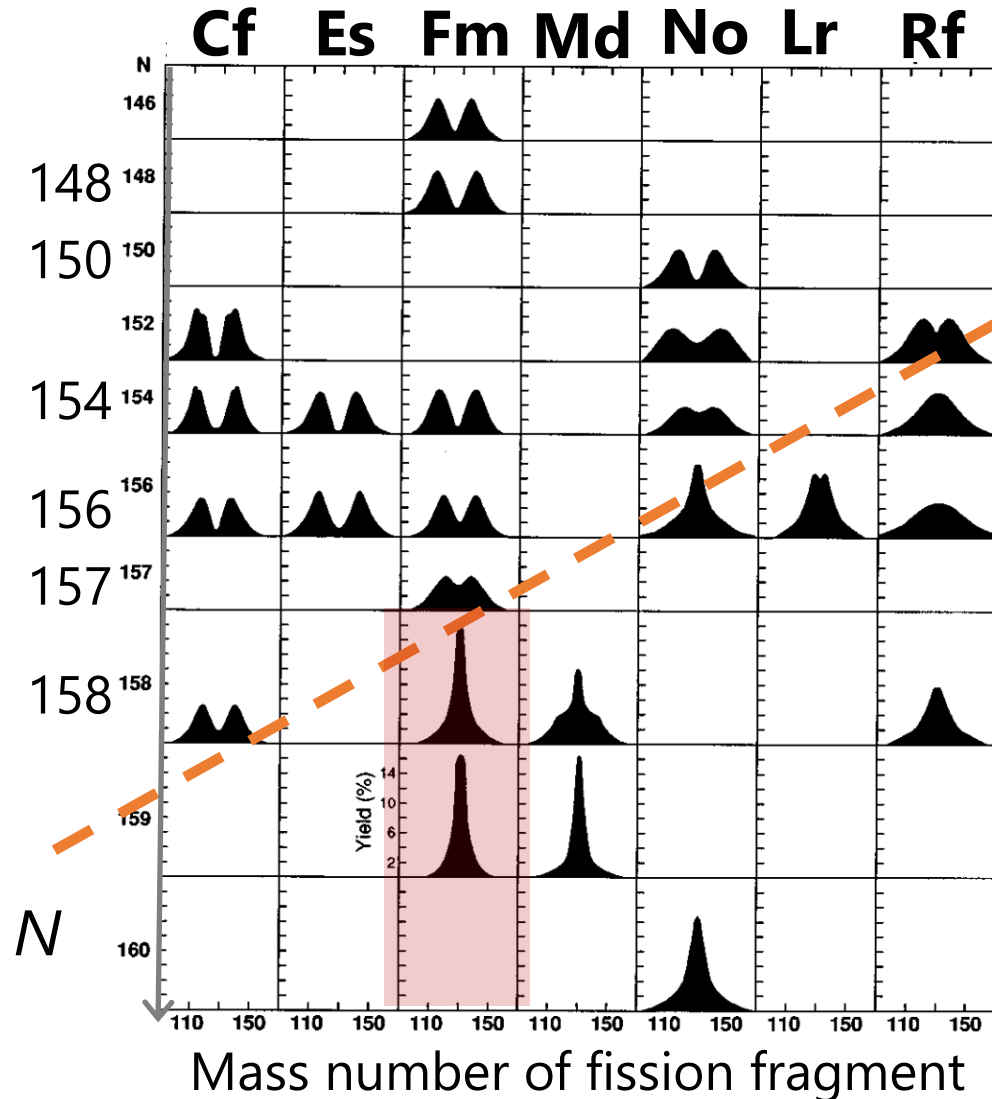
Baran et al., PRC84, 054321 (2011)

- Peak structure
- Larger inertia in QRPA

non-perturbative cranking

$$\frac{F^i}{\dot{s}_i} = U^\dagger \frac{\partial \rho}{\partial s_i} V^* + U^\dagger \frac{\partial \kappa}{\partial s_i} U^* - V^\dagger \frac{\partial \rho^*}{\partial s_i} U^* - V^\dagger \frac{\partial \kappa^*}{\partial s_i} V^*,$$

Mass distribution of fission fragments



Asymmetric fission

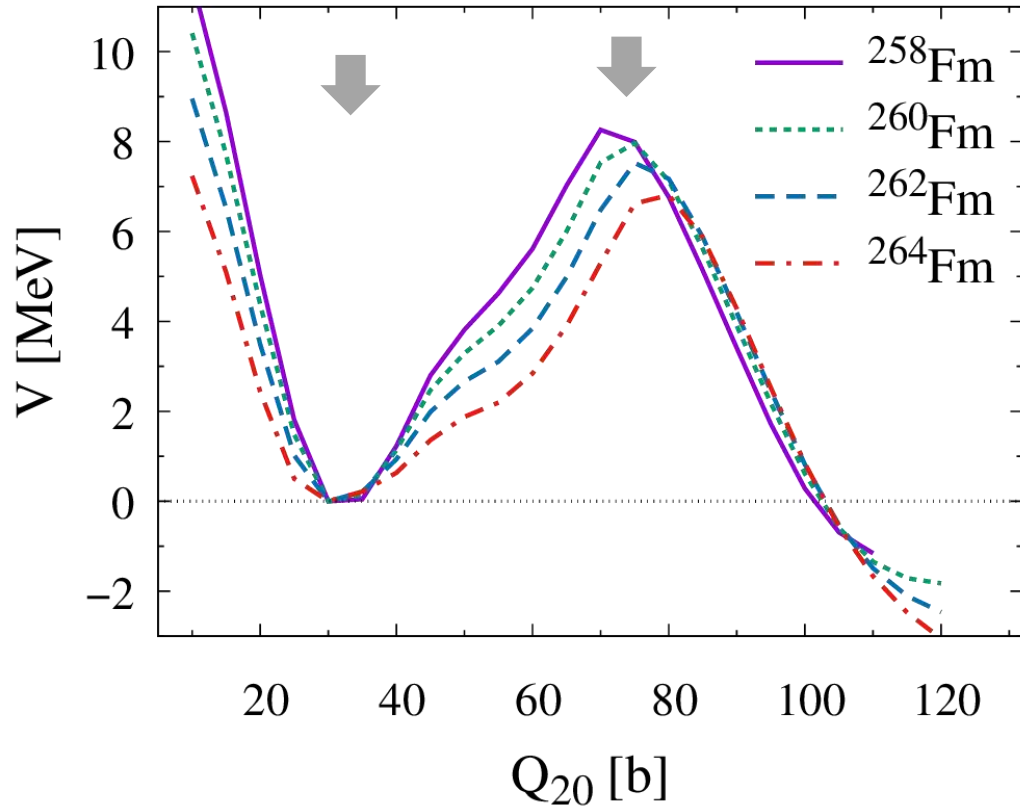
Symmetric fission

Mass symmetric fission in
Fm isotopes with $N \geq 158$

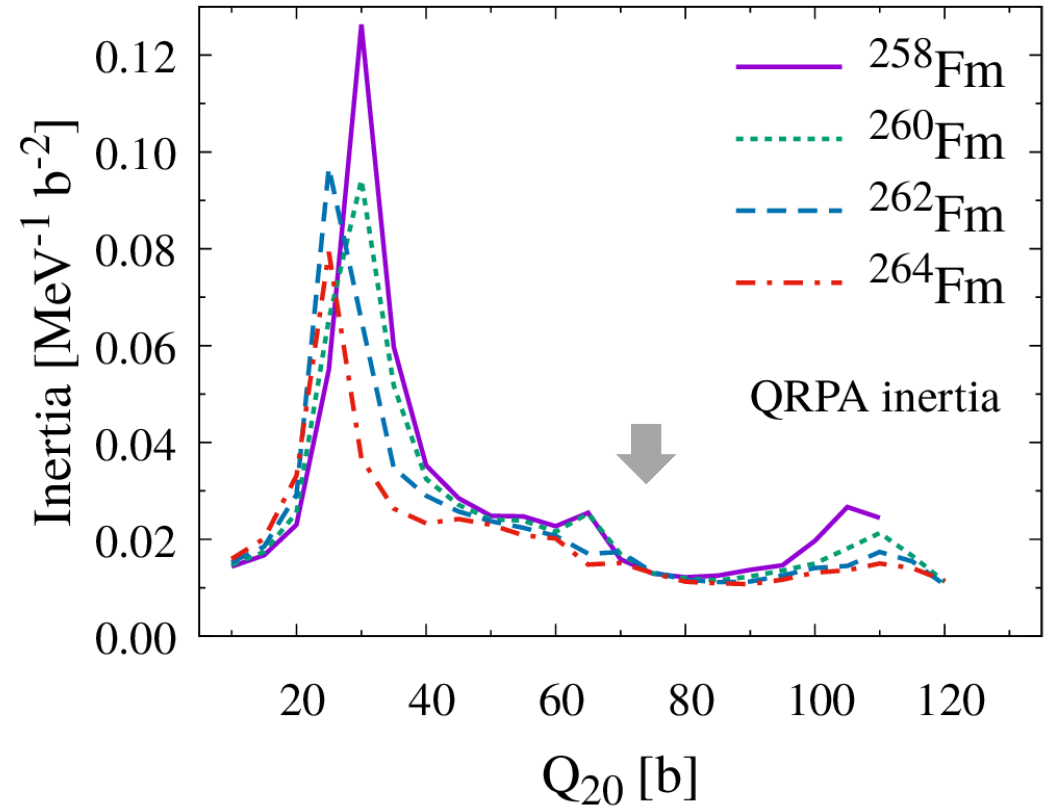
Calculation without Q_{30} may be sufficient
for fission half-life

M.R. Lane et al., PRC53, 2893 (1996)

Fission barrier

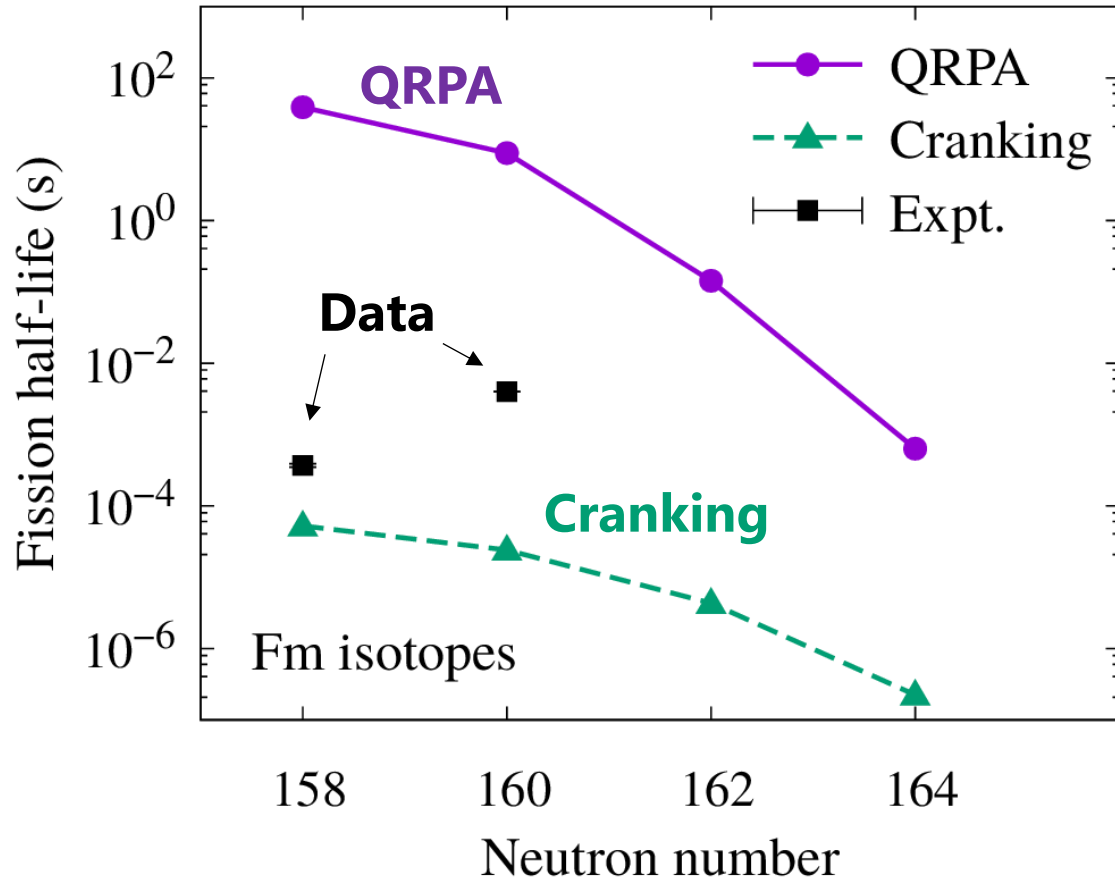


Inertia



- Ground state $Q \sim 30 \text{ b}$ ($\beta \sim 0.26$)
- $^{258}\text{Fm} > ^{260}\text{Fm} > ^{262}\text{Fm} > ^{264}\text{Fm}$

Isotope dependence of half-life



$$T_{1/2} = \ln 2 / (nP)$$

$$n = 10^{20.38} \text{ s}^{-1} \quad P = \frac{1}{1 + \exp(2S)}$$

A. Baran, Phys. Lett. B 76, 8 (1978)

Sadhukhan et al, PRC88, 064314 (2013), etc.

Larger V & $M \rightarrow$ longer half-life

Difference in inertia \rightarrow Half-life

Overestimation

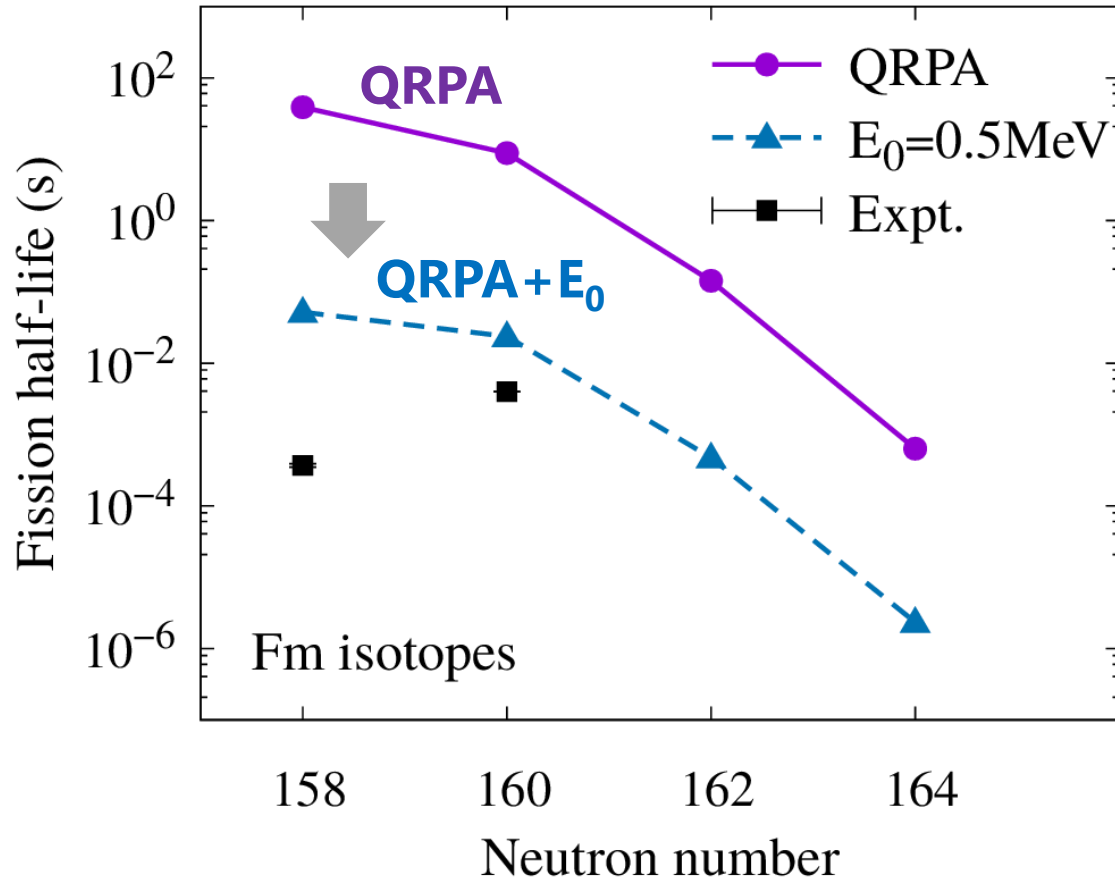
Triaxial shape

Zero-point correction

Data from N. E. Holden and D. C. Hoffman,
Pure Appl. Chem. 72, 1525 (2000)

Washiyama, EPJ Web of Conf. 306, 01026 (2024)

Zero-point correction, E_0



$$S = \int_{s_1}^{s_2} ds \sqrt{2M(s)(V(s) - E)/\hbar}$$

$$E \rightarrow E + E_0$$

Zero-point correction $E_0 = 0.5 \text{ MeV}$

→ Decrease half-life



Fission half-life is sensitive to fission barrier, collective inertia, and zero-point correction energy

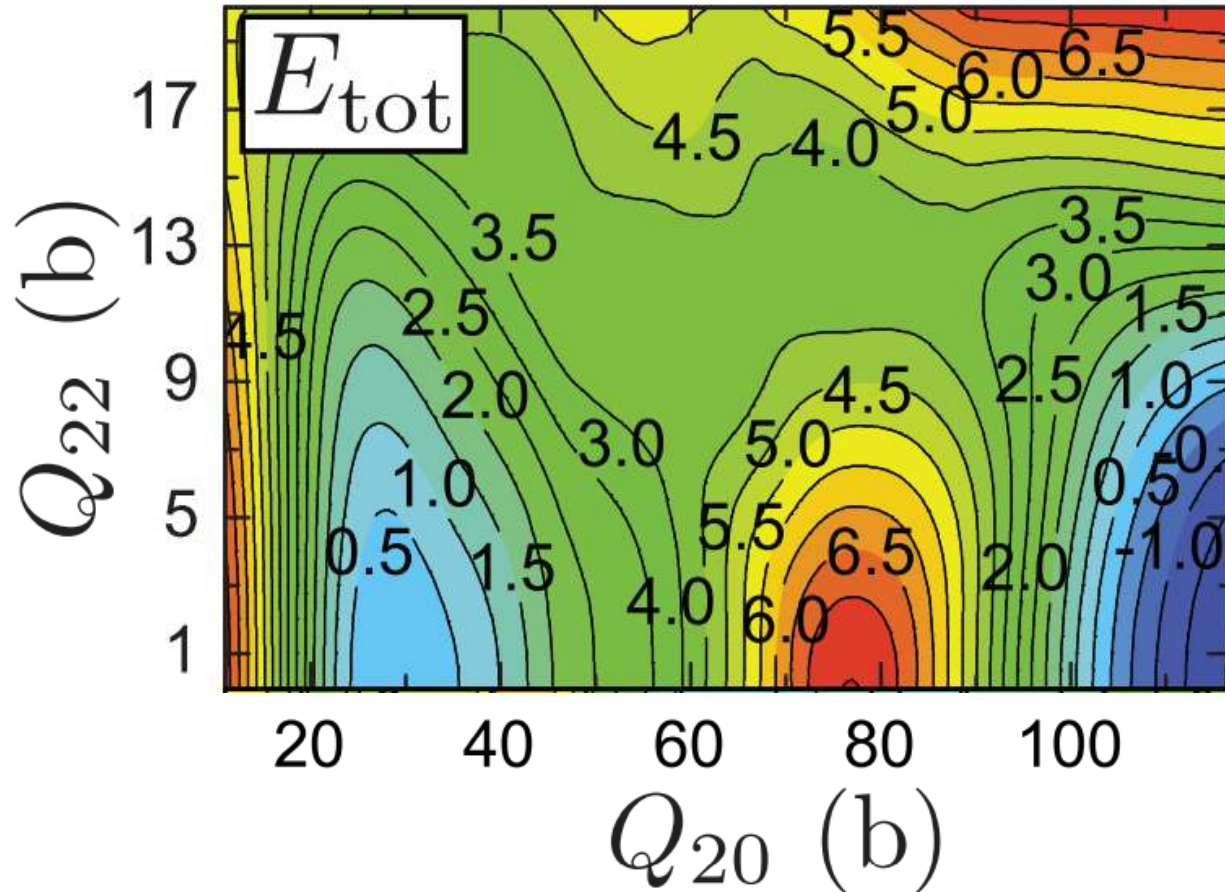
Data from N. E. Holden and D. C. Hoffman,
Pure Appl. Chem. 72, 1525 (2000)

On going work,

Spontaneous fission in two-dimensional collective space
—Determination of the fission path and fission half-life

Description of SF in ^{264}Fm

Sadhukhan et al, PRC88, 064314 (2013)



Fission path = Minimizing the action

$$S = \int_{s_1}^{s_2} ds \sqrt{2M(s)(V(s) - E)}/\hbar$$

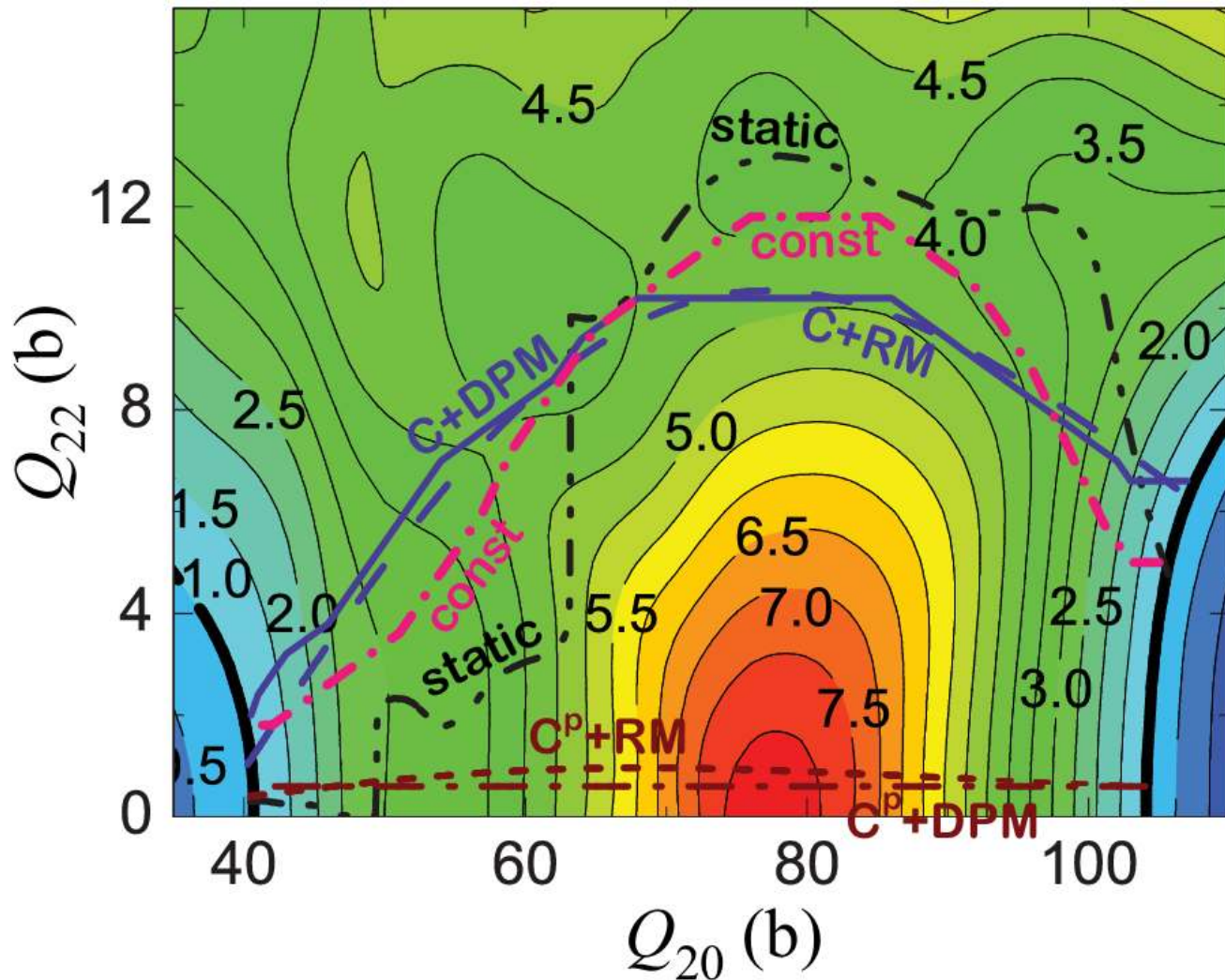
Multi-dimensional case

$$\mathcal{M}_{\text{eff}}(s) = \sum_{ij} \mathcal{M}_{ij} \frac{dq_i}{ds} \frac{dq_j}{ds}$$

Inertia tensor

Description of SF in ^{264}Fm

Sadhukhan et al, PRC88, 064314 (2013)



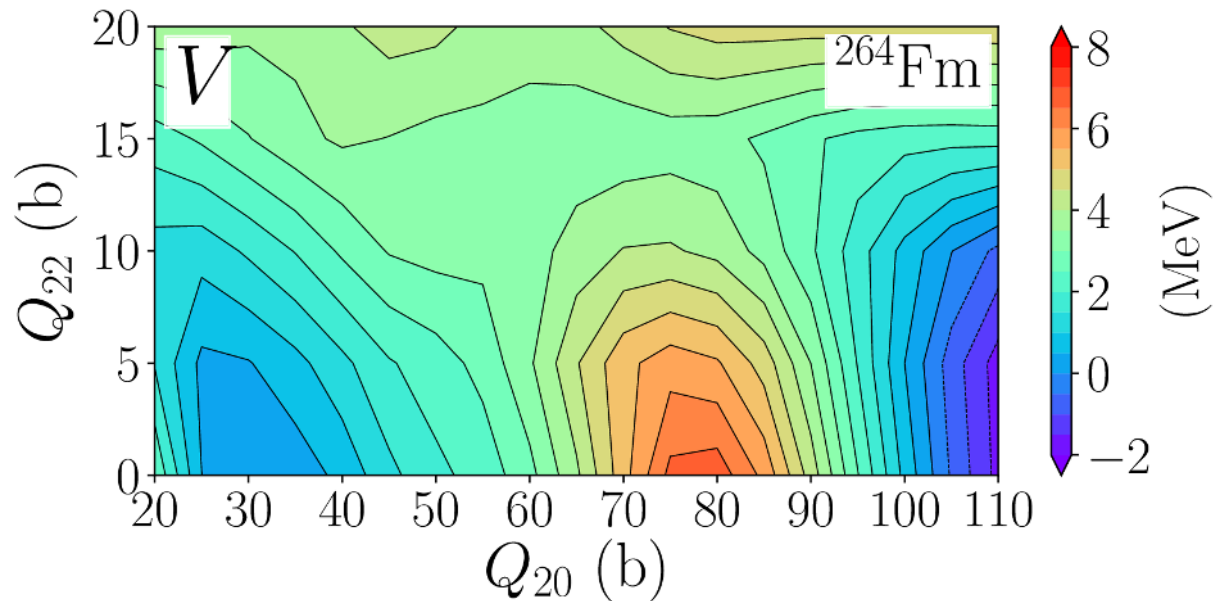
Fission path = Minimizing the action

$$S = \int_{s_1}^{s_2} ds \sqrt{2M(s)(V(s) - E)/\hbar}$$

Spontaneous fission path strongly depends on the choice of collective inertia

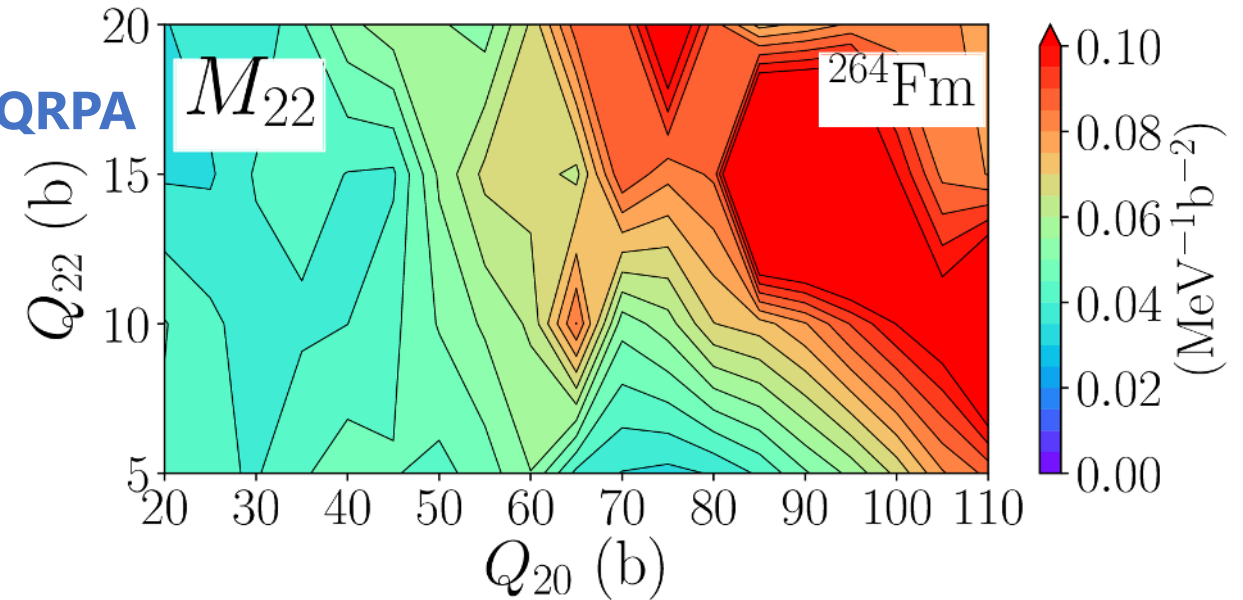
Q_{22} : Triaxial shape

$$Q_{22} \propto r^2 (Y_{22} + Y_{2-2})$$

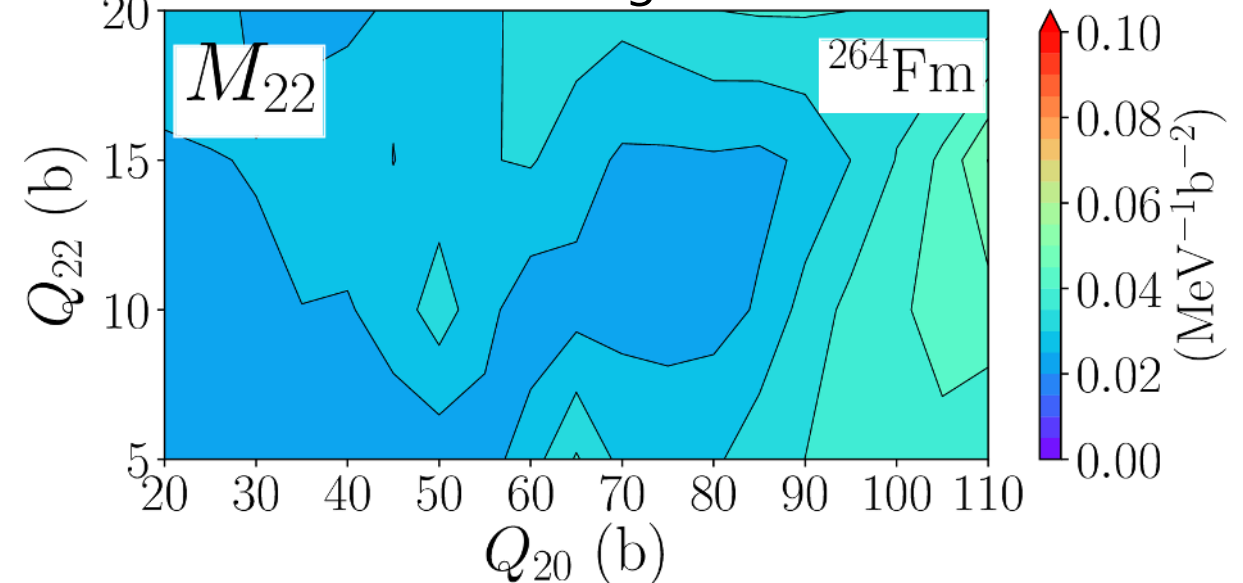


$$M_{22} = M_{Q_{22}} Q_{22}$$

Local QRPA

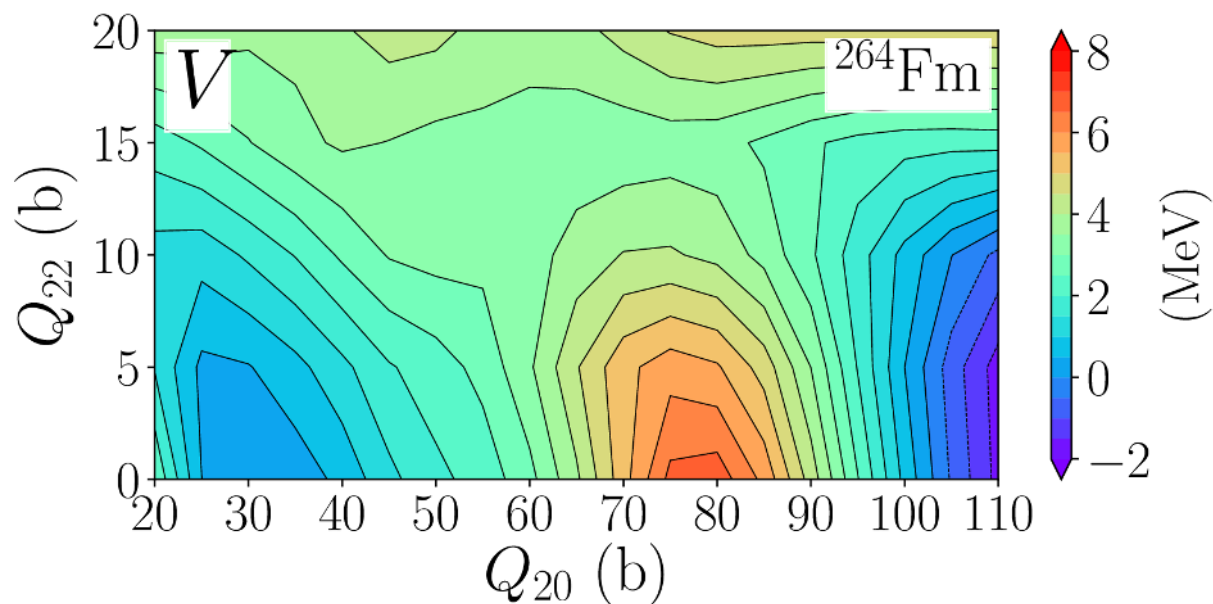


Perturbative cranking



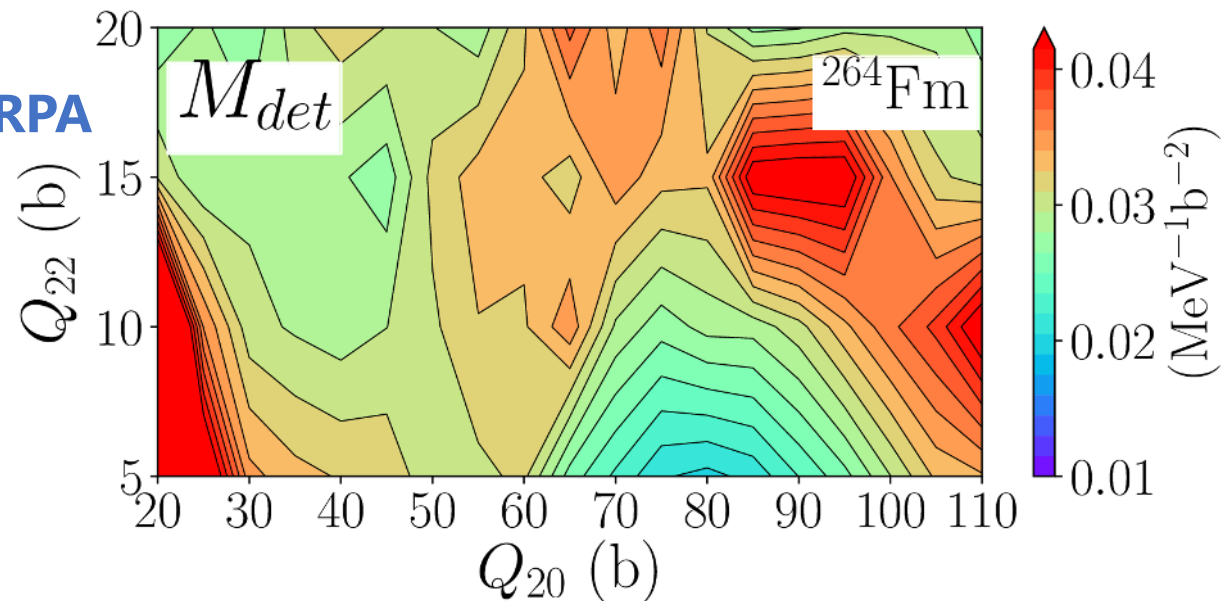
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$$Q_{22} \propto r^2 (Y_{22} + Y_{2-2})$$

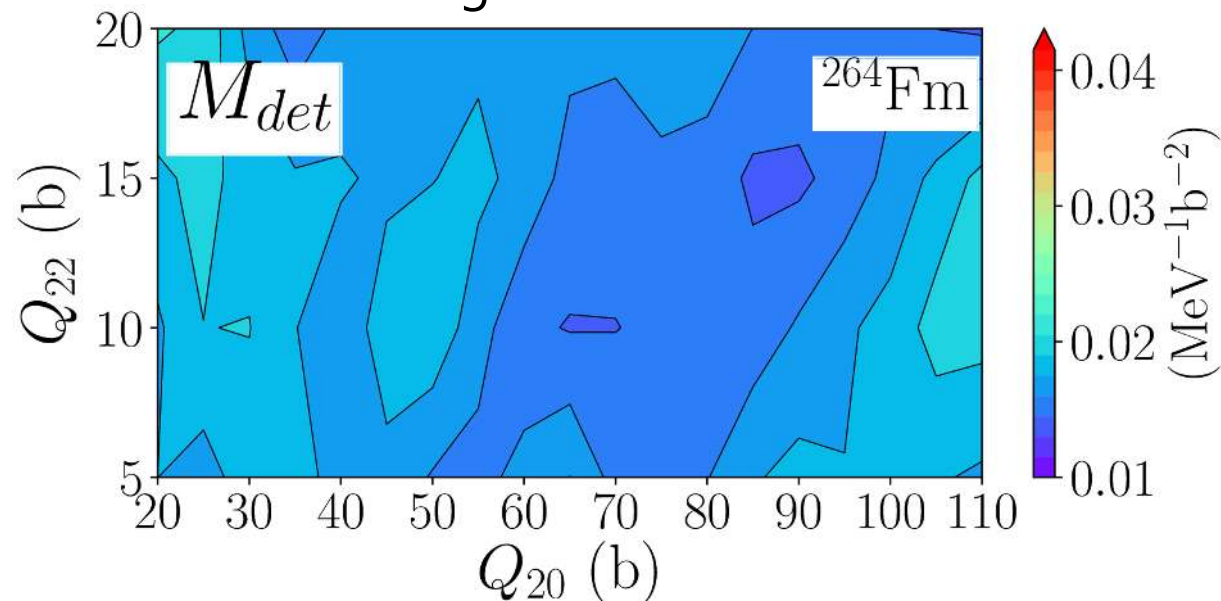


$$M_{det} := \sqrt{M_{00}M_{22} - M_{02}^2}$$

Local QRPA



Perturbative cranking



Spontaneous fission by DFT

Description of collective inertia by DFT + Local QRPA

Correct description on dynamical residual effects

Pronounced peaks at g.s. and fission isomer

$$M_{\text{QRPA}} > M_{\text{cranking}}$$

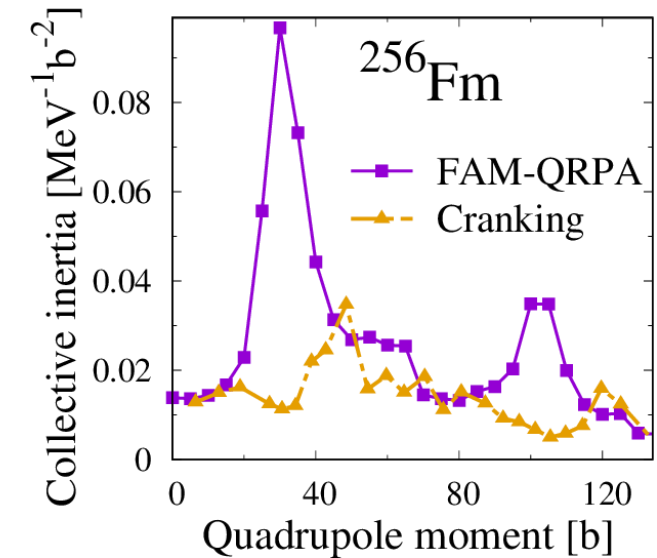
Fission half-life in Fm isotopes

Strong dependence of choice of inertia

Zero-point correction

On-going works

Fission paths in multidimensional space



Half-life

