

Mid-Mass Low-Lying

Spectroscopy via the In-Medium No-Core Shell Model

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Ab Initio Nuclear Structure Toolbox

Many-Body Solution

Diagonalization & Decoupling
NCSM, IM-SRG, IM-NCSM,...

Pre-Conditioning

Similarity RG Transform
Basis Optimization

Post-Processing

Model-Space Extrapolation
Uncertainty Quantification

Chiral EFT Inputs

Interactions & Currents
NN, 3N, YN, YNN,...

$$H |\Psi_n\rangle = E_n |\Psi_n\rangle$$

Uncertainty Quantification

all theory calculations are affected by systematic uncertainties (some also by statistical uncertainties)

in ab initio methods uncertainties result in a controlled way from truncations

variation of all truncation parameters gives access to systematic UQ

■ **chiral EFT uncertainties**

- assess convergence of observables in dependence of chiral truncation order
- convergence affected by regulator scheme and scale, degrees of freedom
- additional uncertainties from fit to experimental data

■ **many-body uncertainties**

- assess convergence of observables with model-space truncation
- convergence affected by basis choice, truncation scheme, Hamiltonian

No-Core Shell Model

No-Core Shell Model

Barrett, Vary, Navrátil, Maris, Nogga, Roth,...

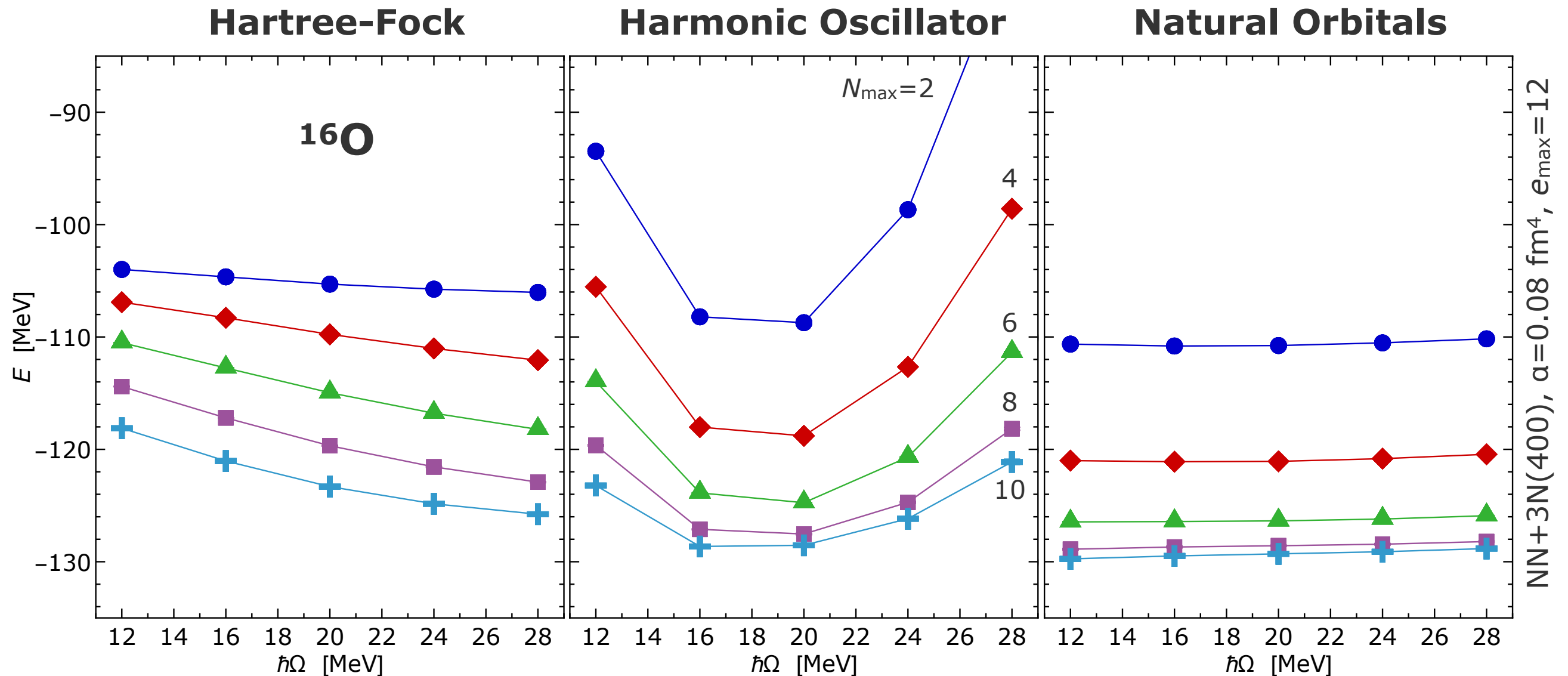
no-core shell model is
universal and powerful ab initio approach for
light nuclei (up to $A \approx 25$)

- **idea**: solve eigenvalue problem of Hamiltonian represented in model space of HO Slater determinants truncated w.r.t. HO excitation energy $N_{\max} \hbar \Omega$

$$\left(\begin{array}{c} \text{Matrix of HO Slater determinants} \end{array} \right) \begin{pmatrix} \vdots \\ C_{i'}^{(n)} \\ \vdots \end{pmatrix} = E_n \begin{pmatrix} \vdots \\ C_i^{(n)} \\ \vdots \end{pmatrix}$$

NCSM Convergence: Energies

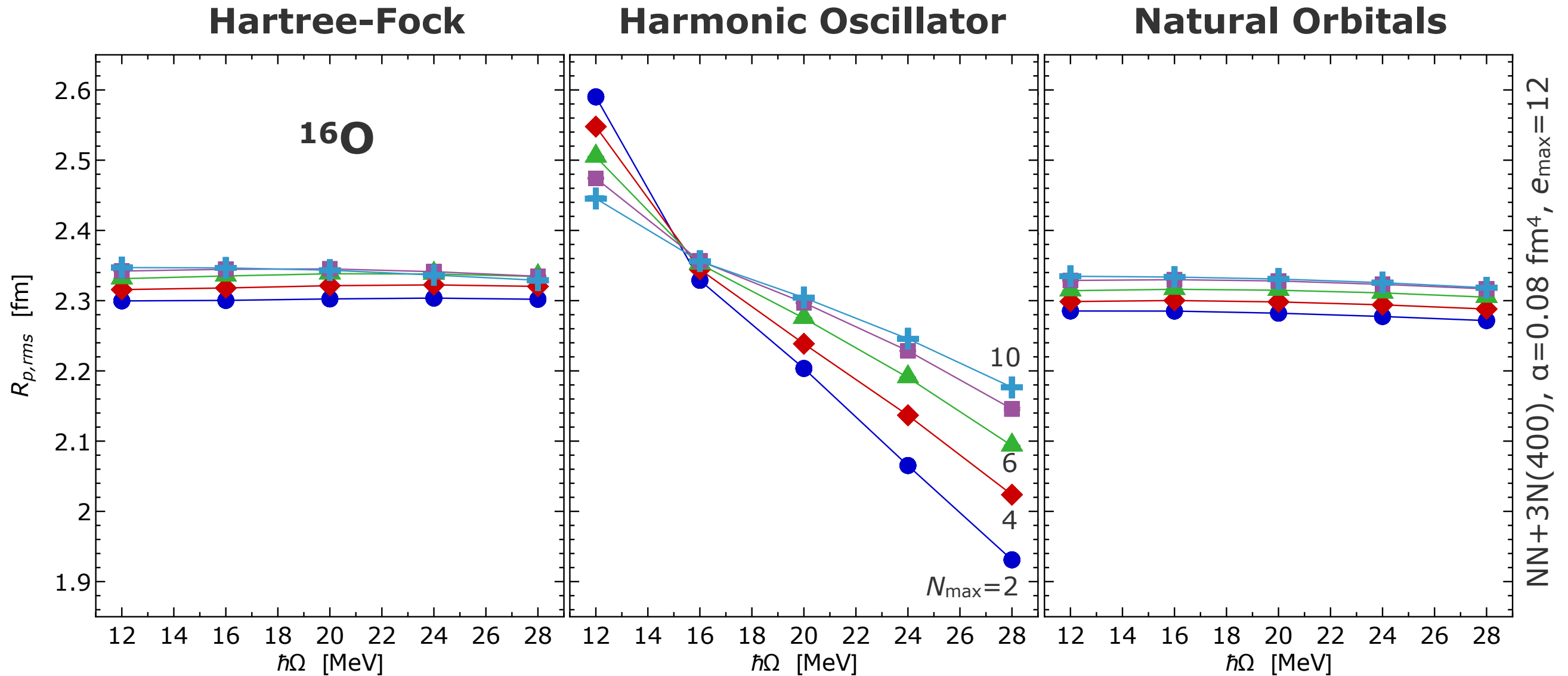
Tichai, Müller, Vobig, Roth; PRC 99, 034321 (2019)



- natural-orbital basis **eliminates frequency dependence** and **accelerates convergence** of NCSM

NCSM Convergence: Radii

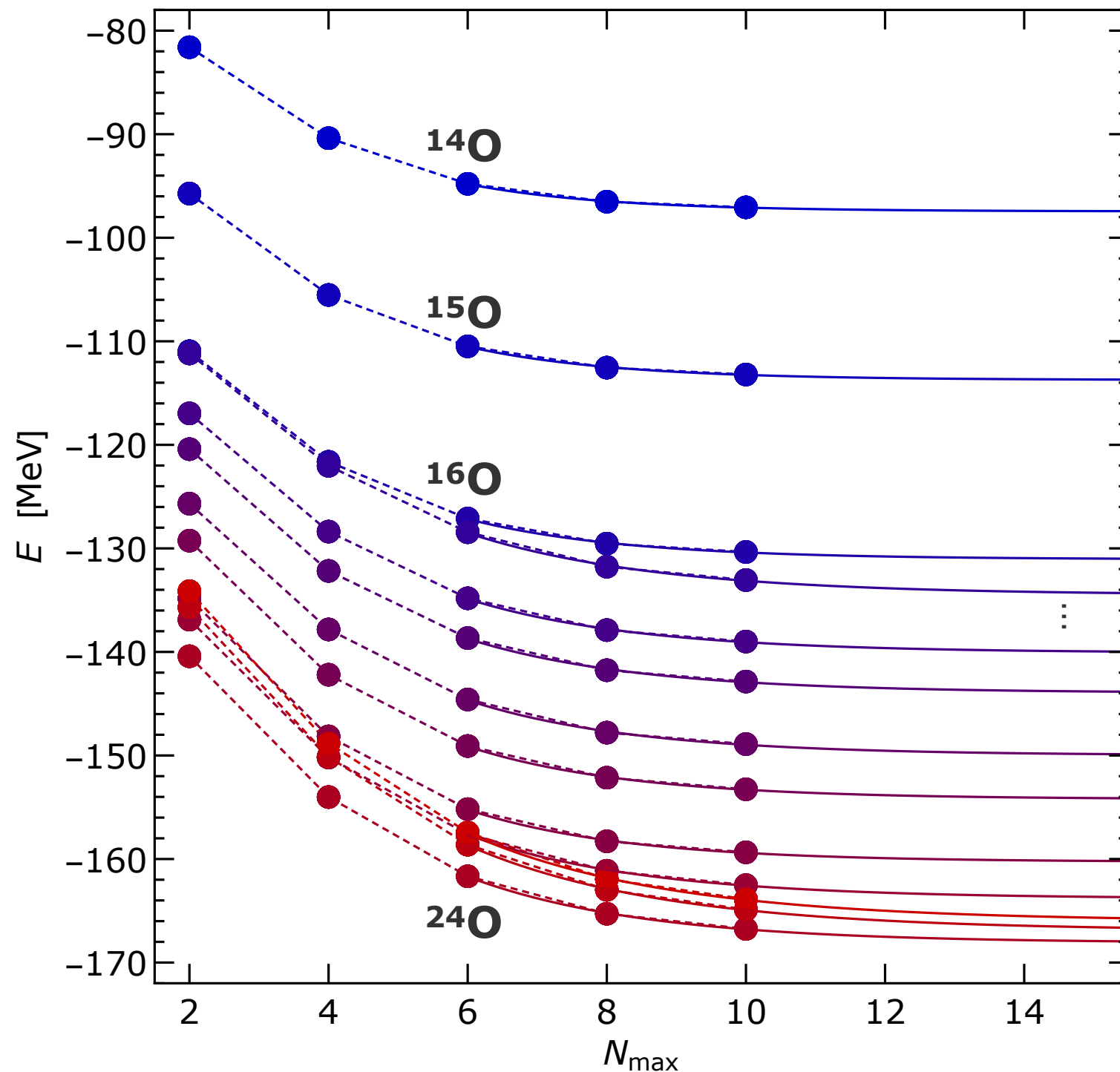
Tichai, Müller, Vobig, Roth; PRC 99, 034321 (2019)



- MBPT natural-orbital basis **eliminates frequency dependence** and **accelerates convergence** of NCSM

Oxygen Isotopes

Tichai, Müller, Vobig, Roth; PRC 99, 034321 (2019)

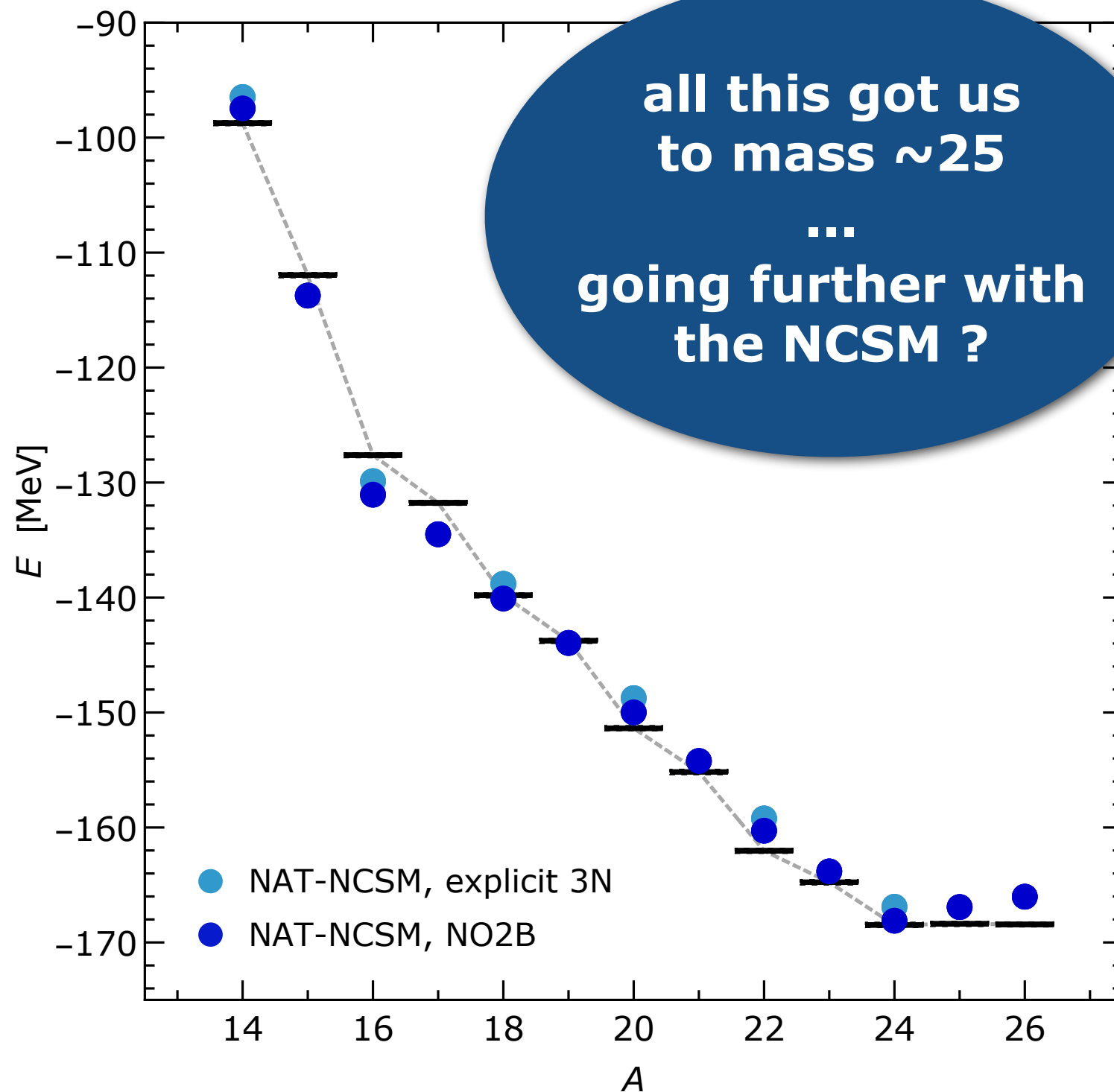


- excellent convergence with natural-orbital basis for all oxygen isotopes

chiral NN+3N
 $\Lambda_{3N}=400$ MeV
 $\alpha=0.08$ fm⁴
 $\hbar\Omega=20$ MeV
 $e_{\text{max}}=12$

Oxygen Isotopes

Tichai, Müller, Vobig, Roth; PRC 99, 034321 (2019)



- excellent convergence with natural-orbital basis for all oxygen isotopes
- very good agreement with experimental systematics and dripline
- NO2B instead of explicit 3N causes $\sim 1\%$ overbinding

In-Medium NCSM

No-Core Shell Model

Barrett, Vary, Navrátil, Maris, Nogga, Roth,...

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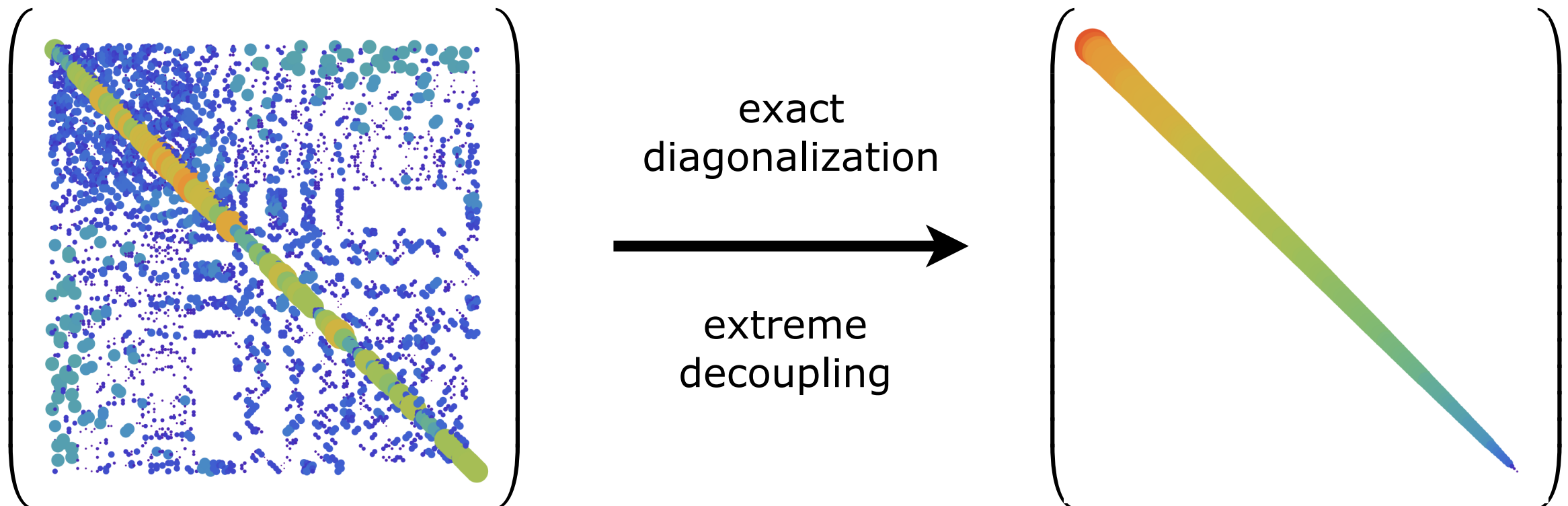
$$\left(\begin{array}{c} \text{Matrix of dots with a diagonal band of orange and green dots} \end{array} \right) \begin{pmatrix} \vdots \\ C_{i'}^{(n)} \\ \vdots \end{pmatrix} = E_n \begin{pmatrix} \vdots \\ C_i^{(n)} \\ \vdots \end{pmatrix}$$

No-Core Shell Model

Barrett, Vary, Navrátil, Maris, Nogga, Roth,...

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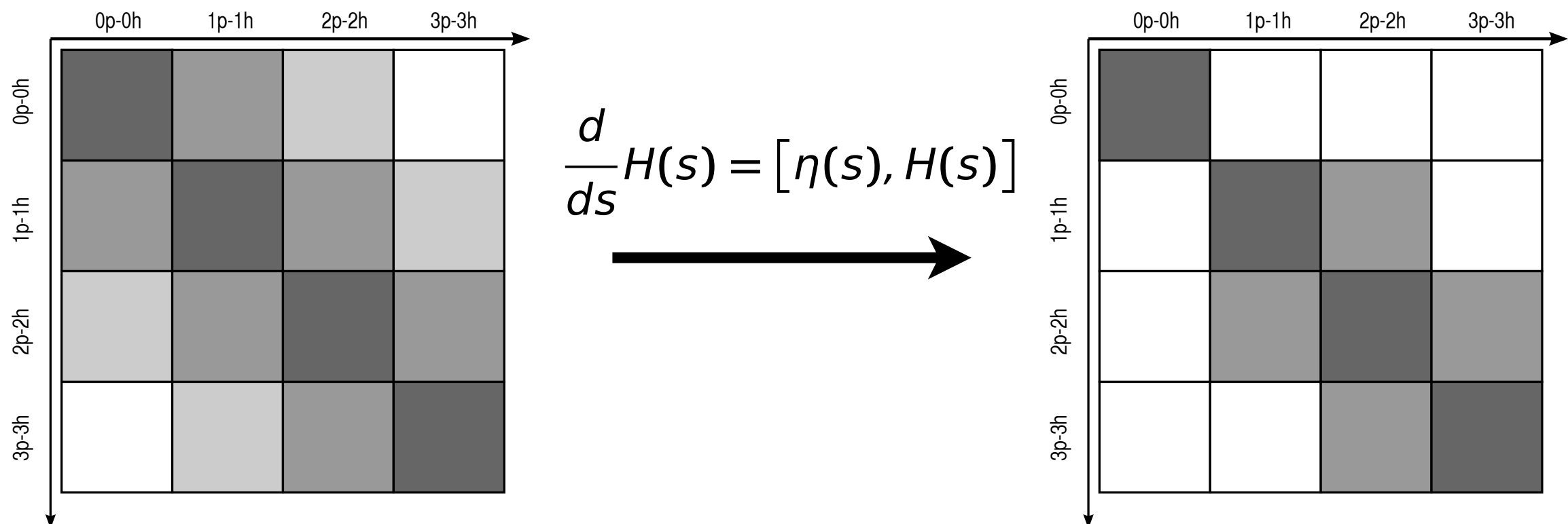


In-Medium SRG

Tsukiyama, Bogner, Schwenk, Hergert,...

decouple reference state from
excitations by a unitary transformation of
Hamiltonian and other operators

- use IM-SRG to decouple single-determinant reference state for particle-hole excitations, 0p0h matrix-element gives ground-state energy

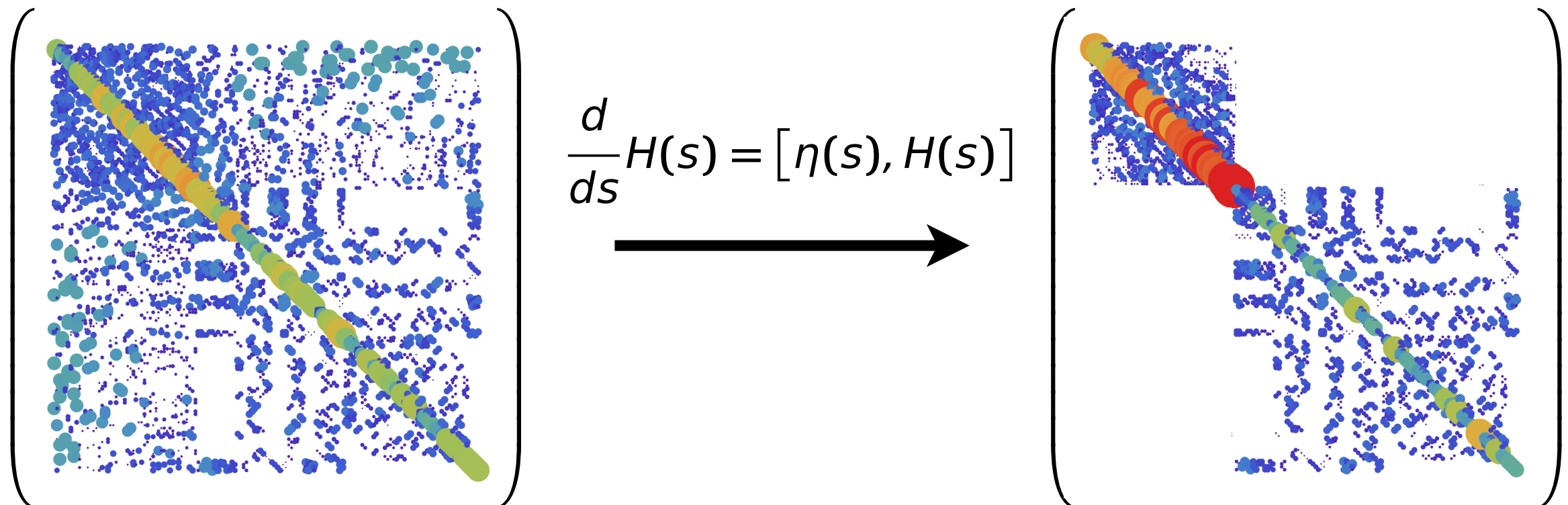


Multi-Reference In-Medium SRG

Hergert, Gebrerufael, Vobig, Mongelli, Roth,...

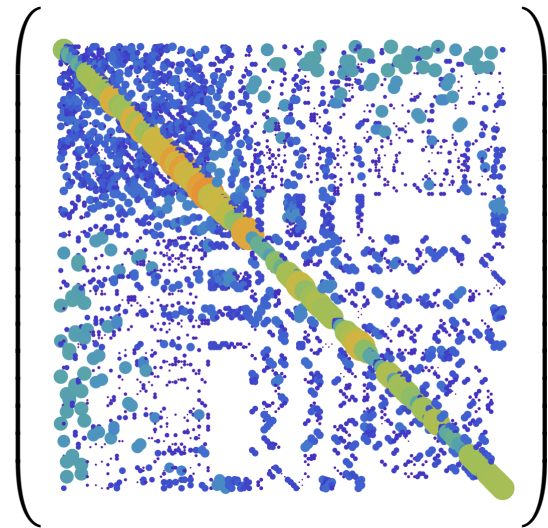
decouple reference state from
excitations by a unitary transformation of
Hamiltonian and other operators

- **idea**: use multi-reference formulation of IM-SRG to decouple reference space for rest of model space, i.e., block diagonalize A -body Hamiltonian

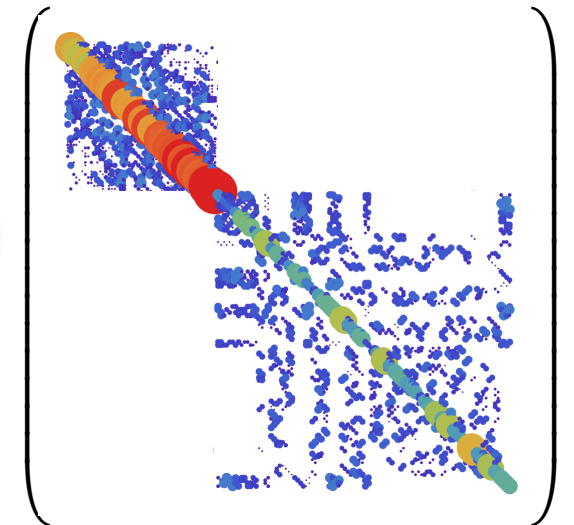


Multi-Reference In-Medium SRG

Hergert, Gebrerufael, Vobig, Mongelli, Roth,...



use SRG flow equations for multi-reference normal-ordered Hamiltonian to decouple reference space



$$\frac{d}{ds}H(s) = [\eta(s), H(s)]$$

[Kutzelnigg & Mukherjee, 1997]

- Hamiltonian and generator in normal order with respect to multi-determinant reference state, omit residual three-body piece

$$H(s) = E(s) + \sum_{ij} f_j^i(s) \tilde{A}_j^i + \frac{1}{4} \sum_{ijkl} \Gamma_{kl}^{ij}(s) \tilde{A}_{kl}^{ij} + \cancel{\frac{1}{36} \sum_{ijklmn} W_{lmn}^{ijk}(s) \tilde{A}_{lmn}^{ijk}}$$

- define generator to suppress off-diagonal contributions that couple reference state to ph excitations

$$\eta(s) = [H(s), H^d(s)] = [H^{od}(s), H^d(s)]$$

In-Medium NCSM

NCSM
reference state

- ground-state from NCSM at small N_{\max} as reference state for multi-reference IM-SRG
- access to all open-shell nuclei and systematically improvable

MR-IM-SRG
decoupling

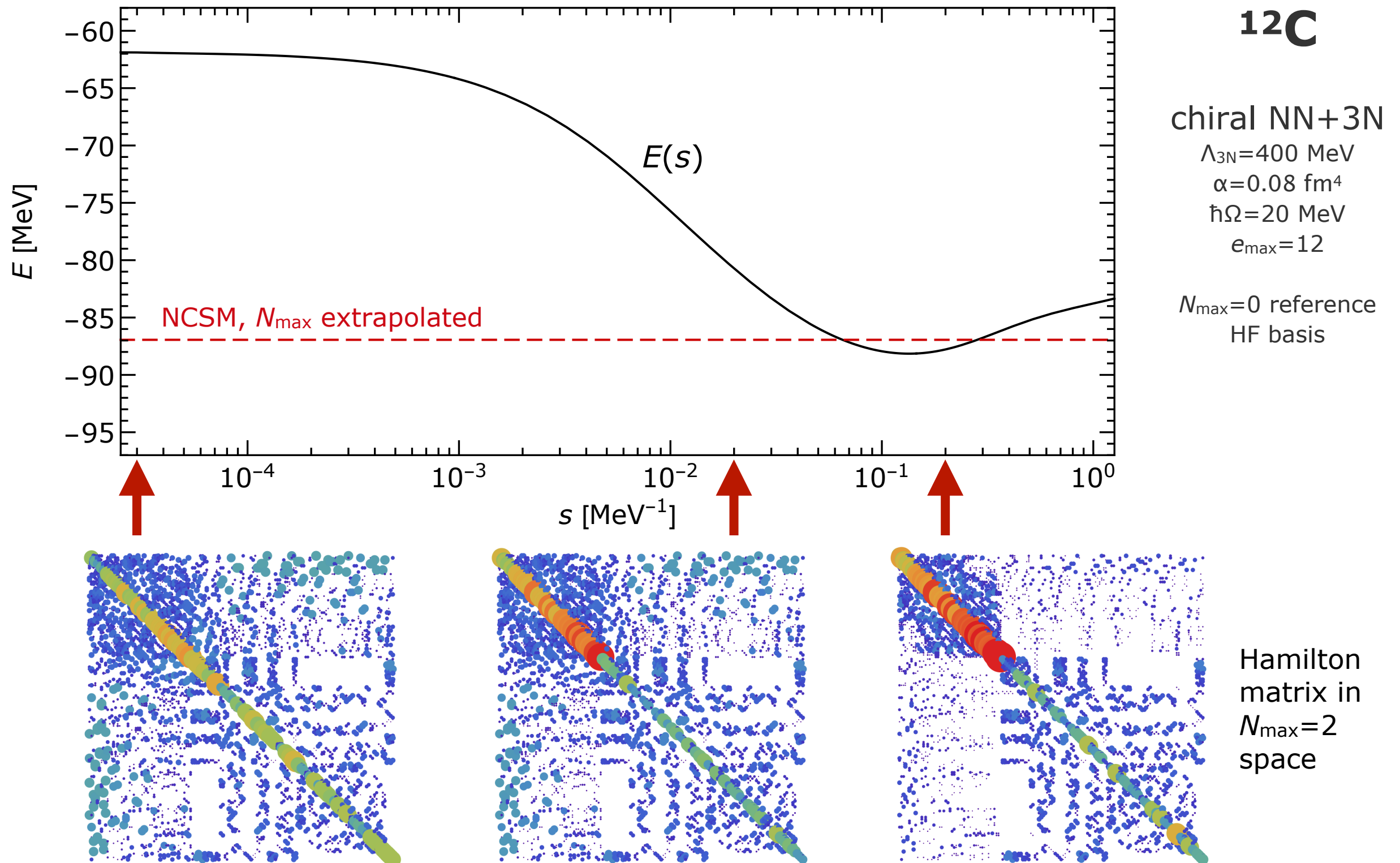
- IM-SRG evolution of multi-reference normal-ordered Hamiltonian and other operators
- decoupling of particle-hole excitations, i.e., pre-diagonalization in many-body space

NCSM
many-body solution

- use in-medium evolved Hamiltonian for a subsequent NCSM calculation
- access to ground and excited states and full suite of observables

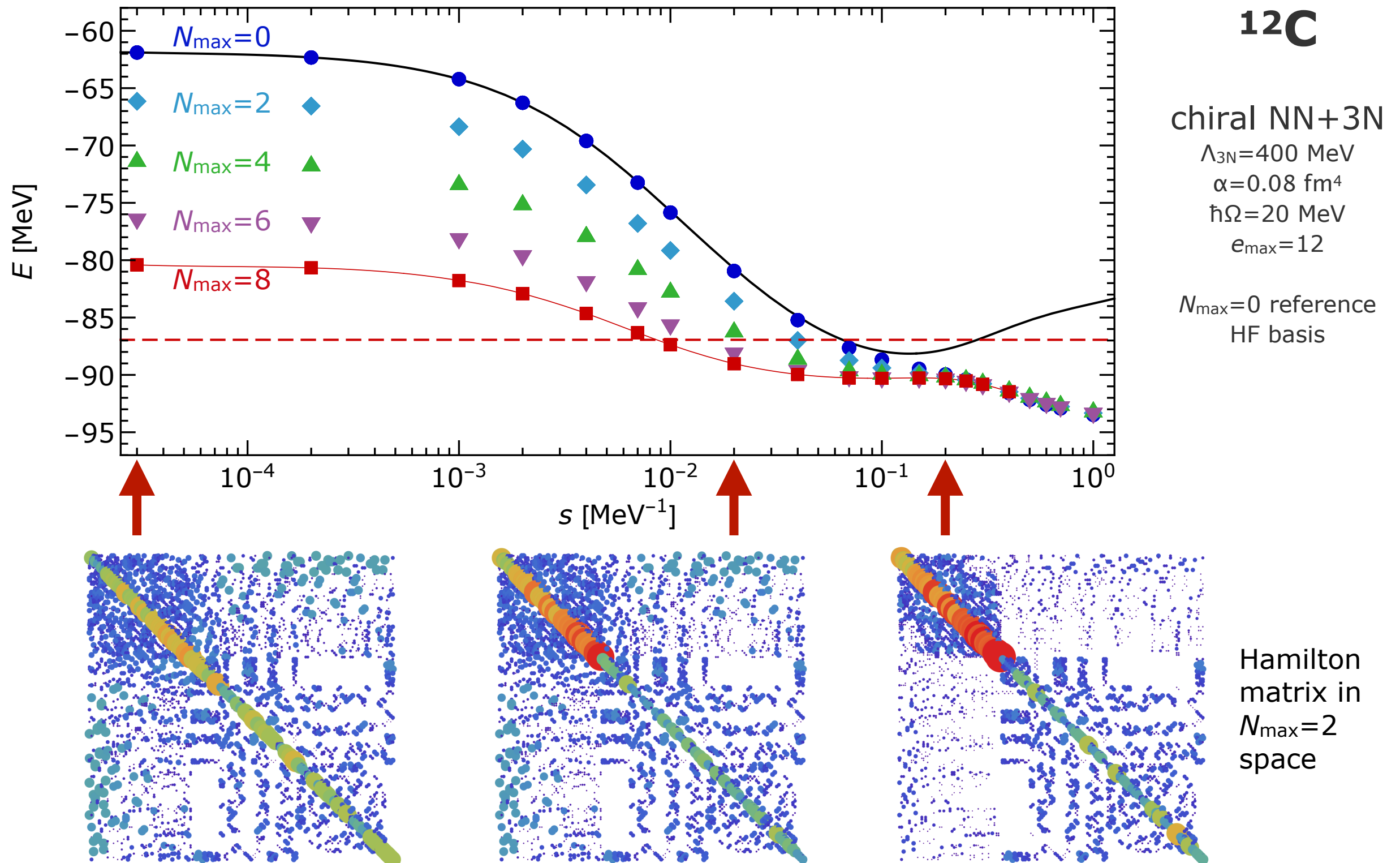
In-Medium NCSM: Flow Evolution

Gebrerufael, Vobig, Hergert, Roth; PRL 118, 152503 (2017)



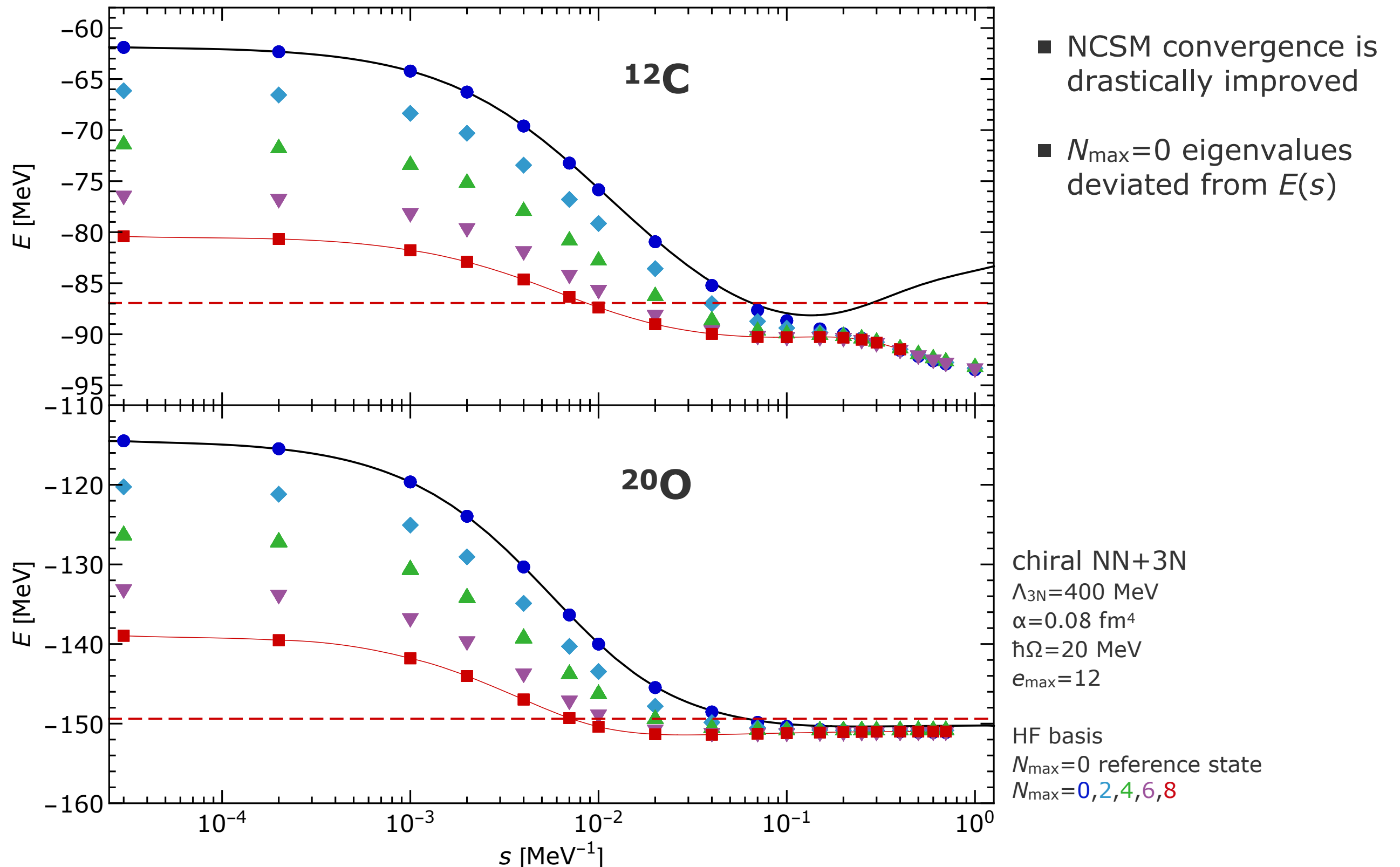
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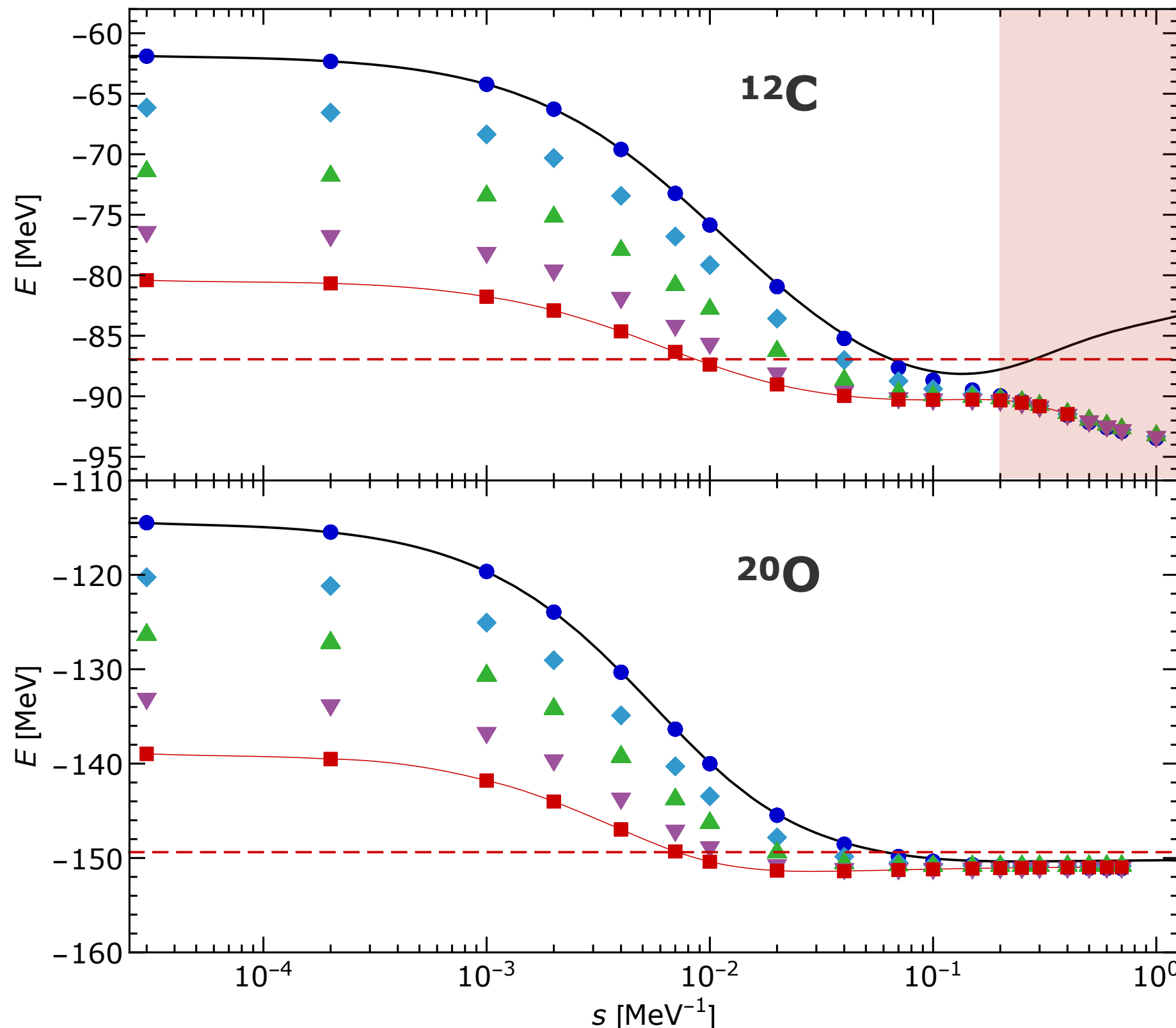
In-Medium NCSM: Flow Evolution

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In-Medium NCSM: Flow Evolution

Gebrerufael, Vobig, Hergert, Roth; PRL 118, 152503 (2017)



- NCSM convergence is drastically improved
- $N_{\text{max}}=0$ eigenvalues deviated from $E(s)$
- large- s behaviour can be improved by slight modification of generator

chiral NN+3N

$\Lambda_{3N}=400$ MeV

$\alpha=0.08$ fm⁴

$\hbar\Omega=20$ MeV

$e_{\text{max}}=12$

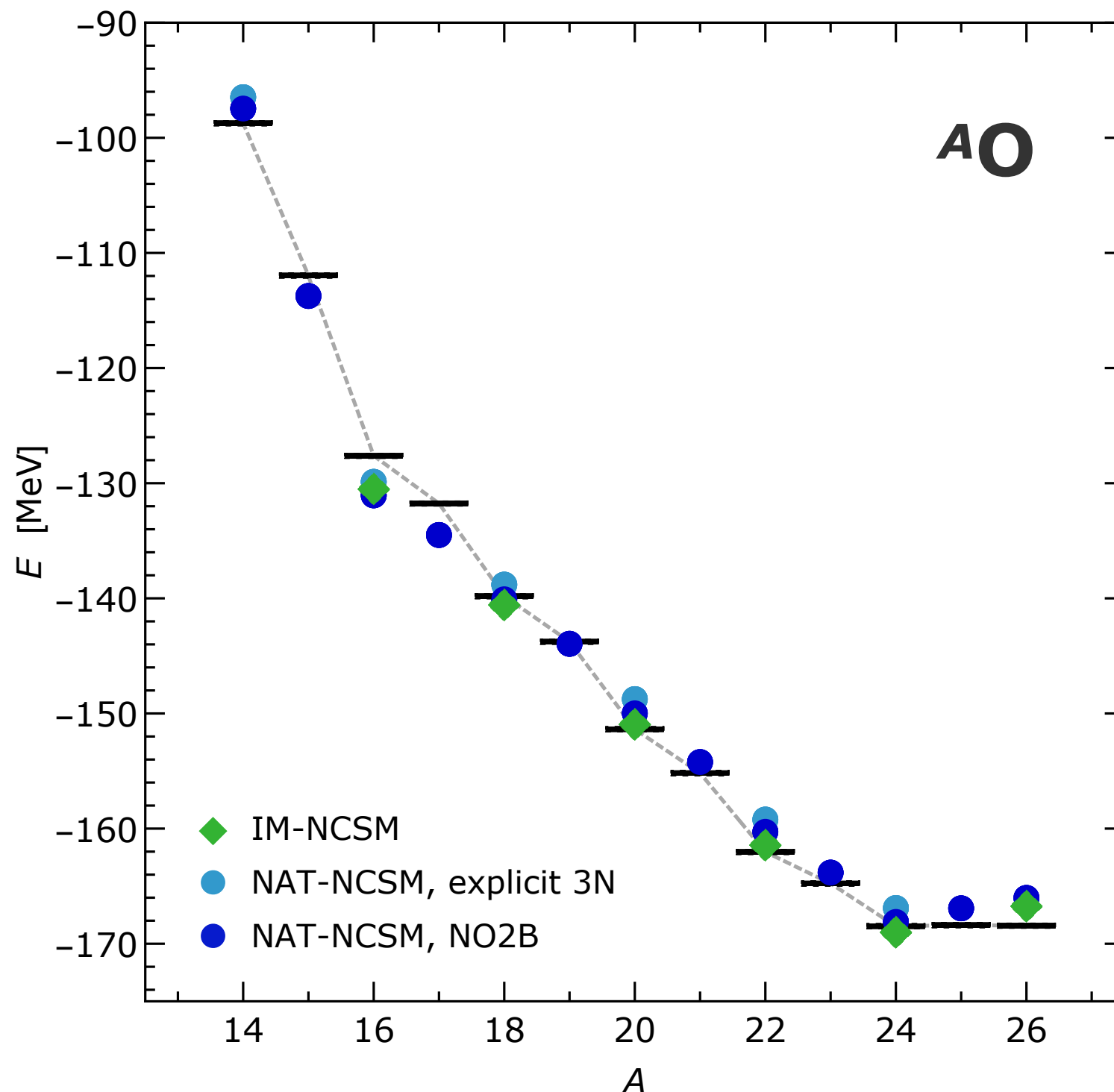
HF basis

$N_{\text{max}}=0$ reference state

$N_{\text{max}}=0, 2, 4, 6, 8$

IM-NCSM: Oxygen Isotopes

Gebrerufael, Vobig, Hergert, Roth; PRL 118, 152503 (2017)



- excellent agreement with direct NCSM
- IM-SRG evolution limited to $J=0$ reference states and thus even-mass isotopes
- odd-mass nuclei via simple particle attachment or removal in final NCSM run

chiral NN+3N

$\Lambda_{3N}=400$ MeV

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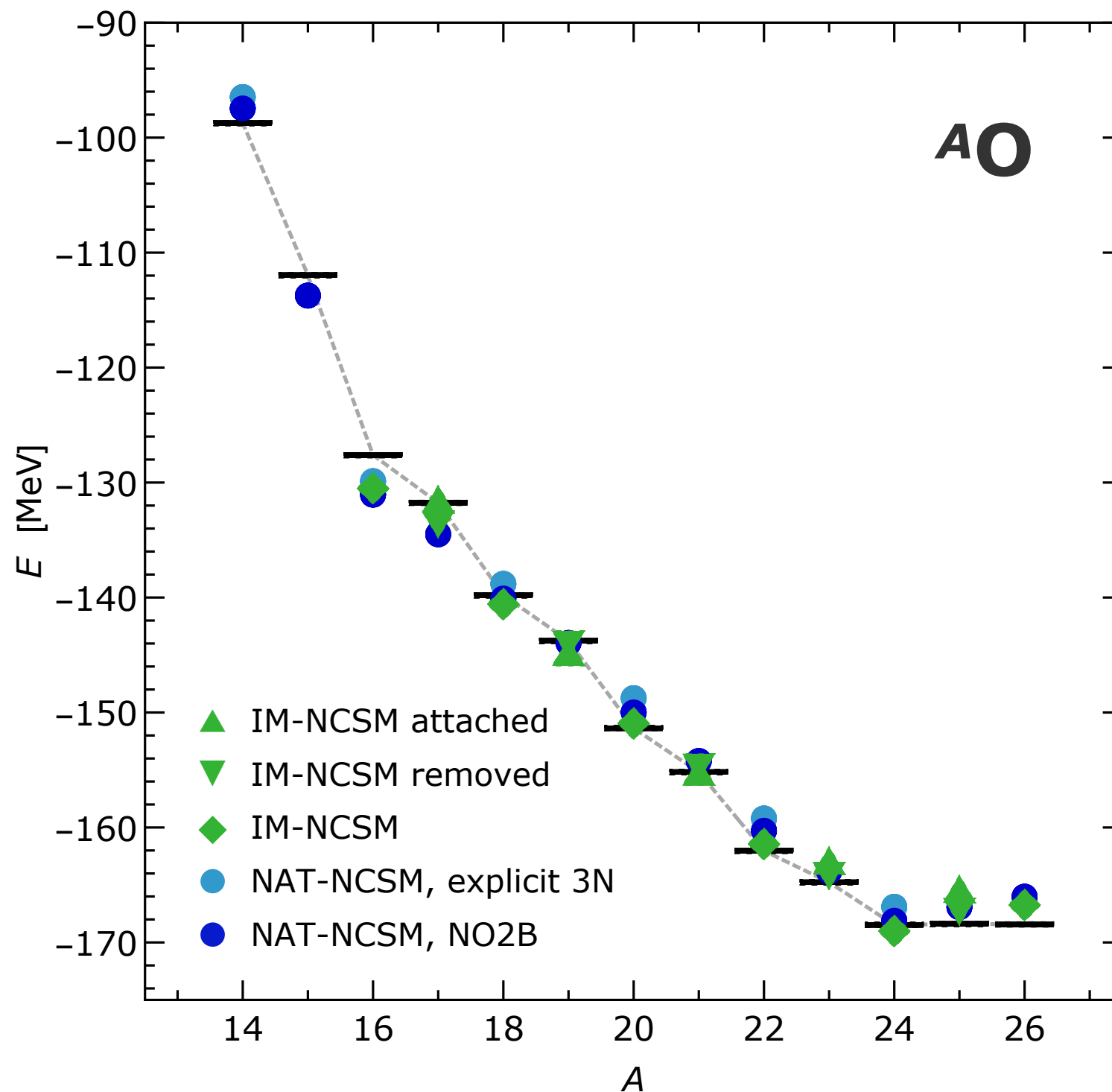
$\hbar\Omega=20$ MeV

$e_{\max}=12$

HF basis

$N_{\max}=0$ reference

IM-NCSM: Oxygen Isotopes



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$e_{\max}=12$

HF basis

$N_{\max}=0$ reference

In-Medium NCSM: Uncertainties

- IM-SRG evolution induces additional uncertainties due to the **truncation** of all normal-ordered operators **at the two-body level**...
...NO2B, IM-SRG(2), IM-SRG(M2)
- explicit inclusion of normal-ordered three-body terms is prohibitive in large model spaces (for now)
- probe accuracy of NO2B approximation through variation of flow parameter & reference space truncation
- **uncertainty quantification protocol:** perform IM-NCSM calculation for...
 - different reference space truncations: $N_{\max}^{\text{ref}} = 0, 2, 4$
 - different flow parameters: $s_{\text{sat}}, s_{\text{sat}}/2$
 - different model-space truncations: $N_{\max} = 0, 2, 4, 6, \dots$...maximum difference to next-smaller control parameters gives estimate for many-body uncertainty

Applications with Nonlocal $NN+3N$ Interactions

Family of Non-Local NN+3N Interactions

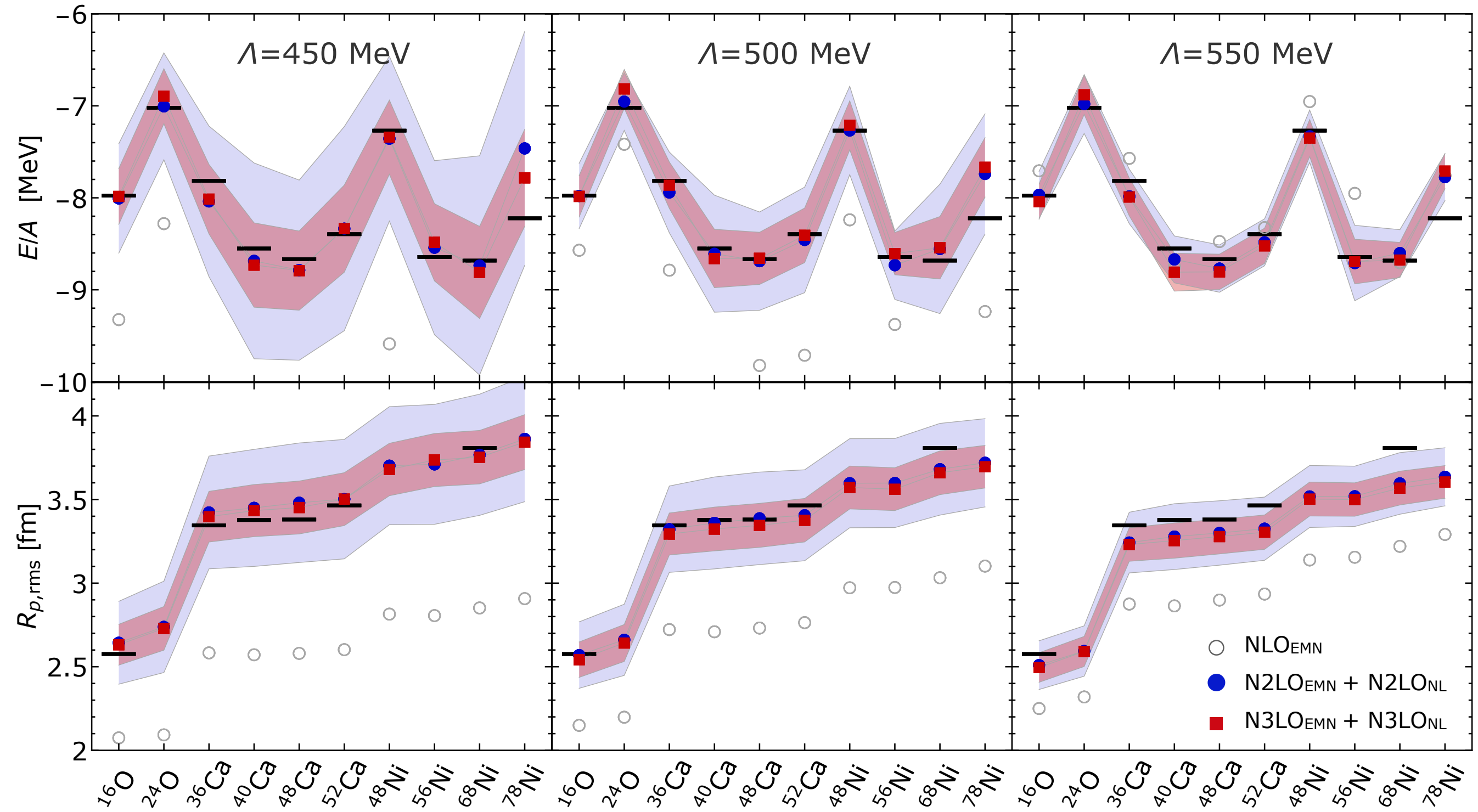
	NN	3N	4N
LO			
NLO			
N2LO			
N3LO			
N4LO			

- NN+3N interactions with same chiral orders and regulator scheme and scale
- up to N3LO with different cutoff values
- low-energy constants fit to $A=2,3$ data and ^{16}O ground-state energy

T. H  ther et al., PLB 808, 135651 (2020)

Medium-Mass Nuclei

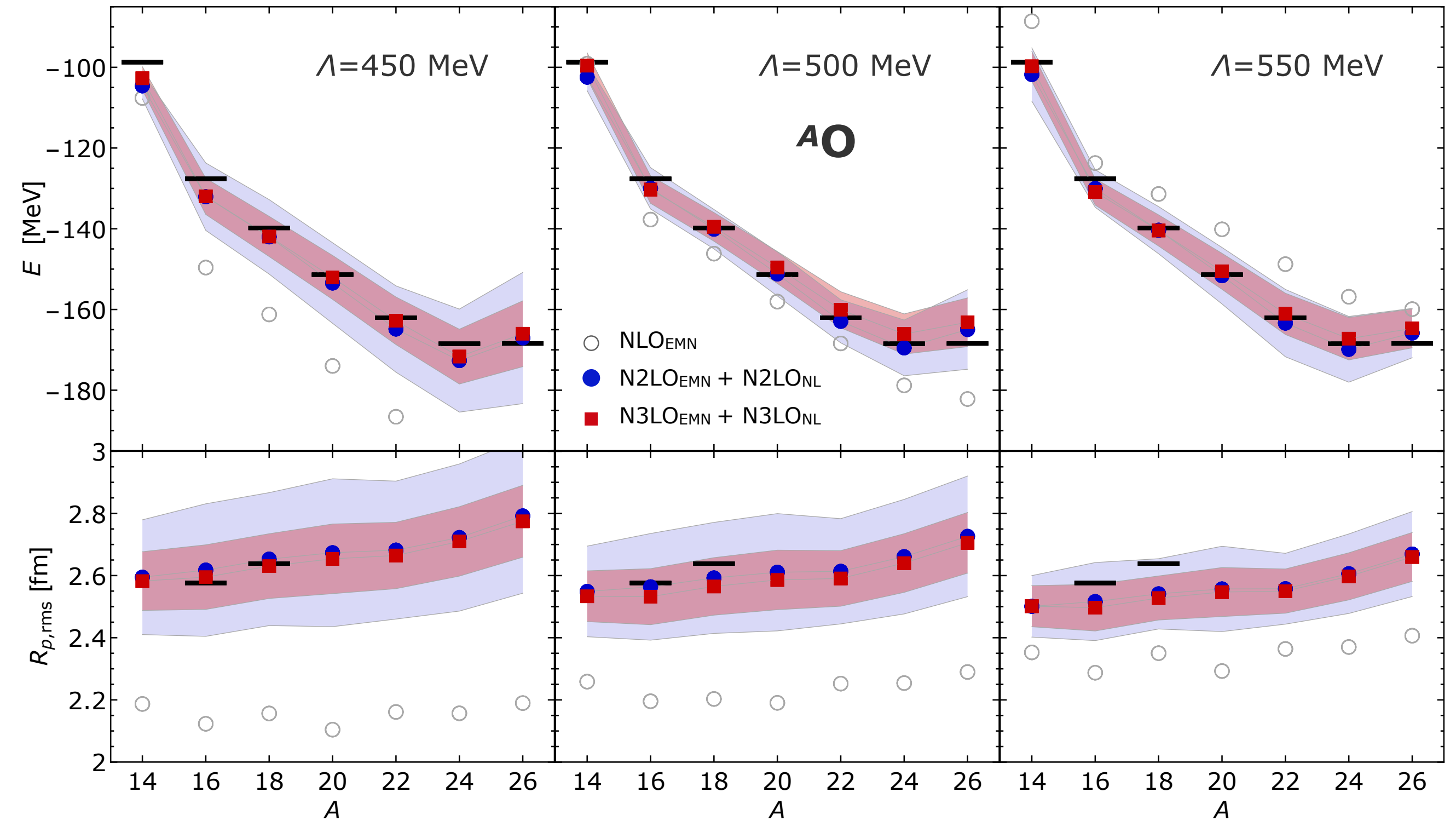
Hüther et al.; PLB 808, 135651 (2020)



IM-SRG(M2), natural orbitals, $\hbar\Omega=20$ MeV, $\alpha=0.04$ fm⁴, $e_{\text{max}}=12$, $E_{3\text{max}}=16$
 error bands show interaction + many-body uncertainties

Oxygen Isotopic Chain

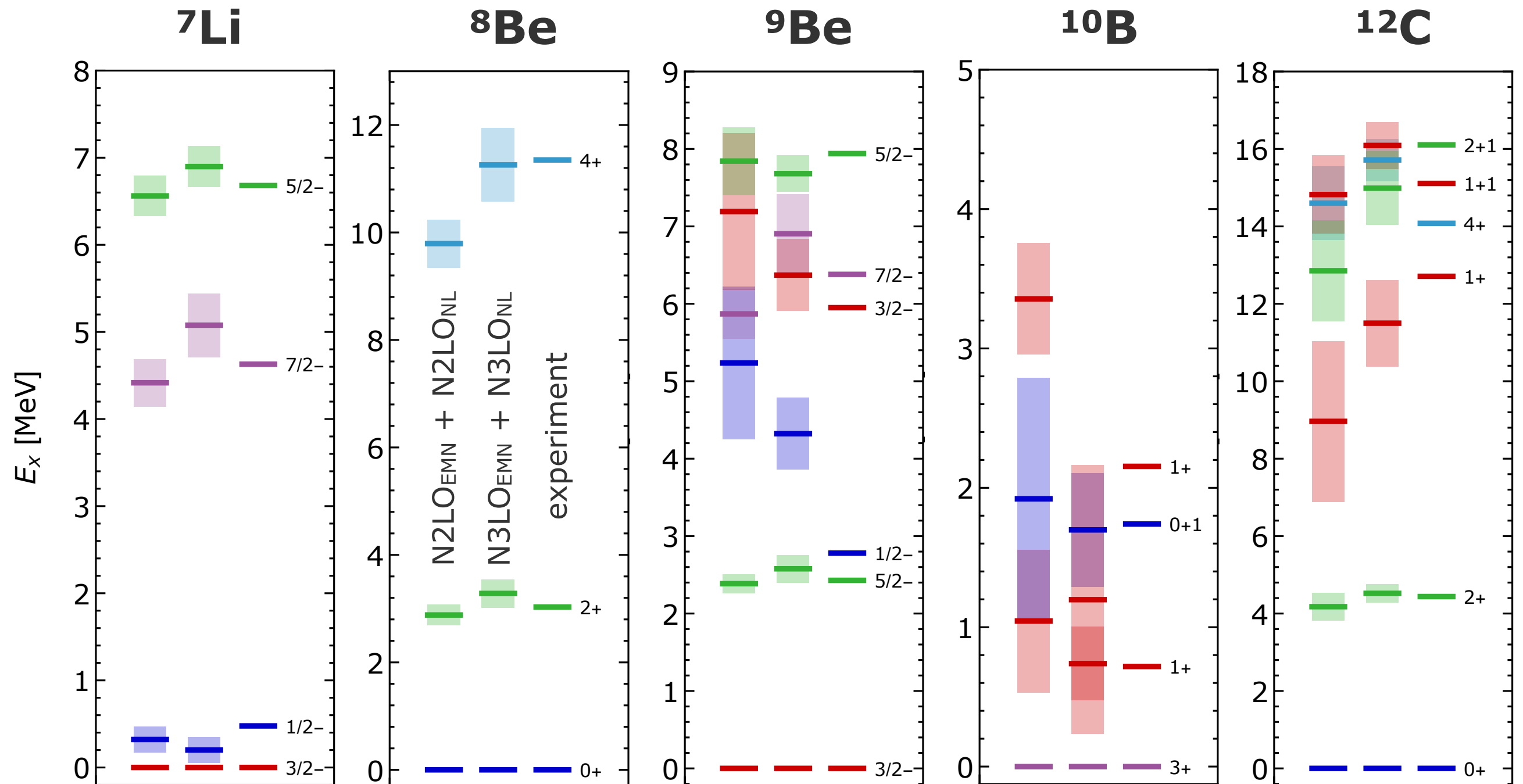
Hüther et al.; PLB 808, 135651 (2020)



IM-NCSM, natural orbitals, $\hbar\Omega=20$ MeV, $\alpha=0.04$ fm⁴, $e_{\max}=12$, $E_{3\max}=14$, $N_{\text{ref}}=2$
 error bands show interaction + many-body uncertainties

p-Shell Spectra

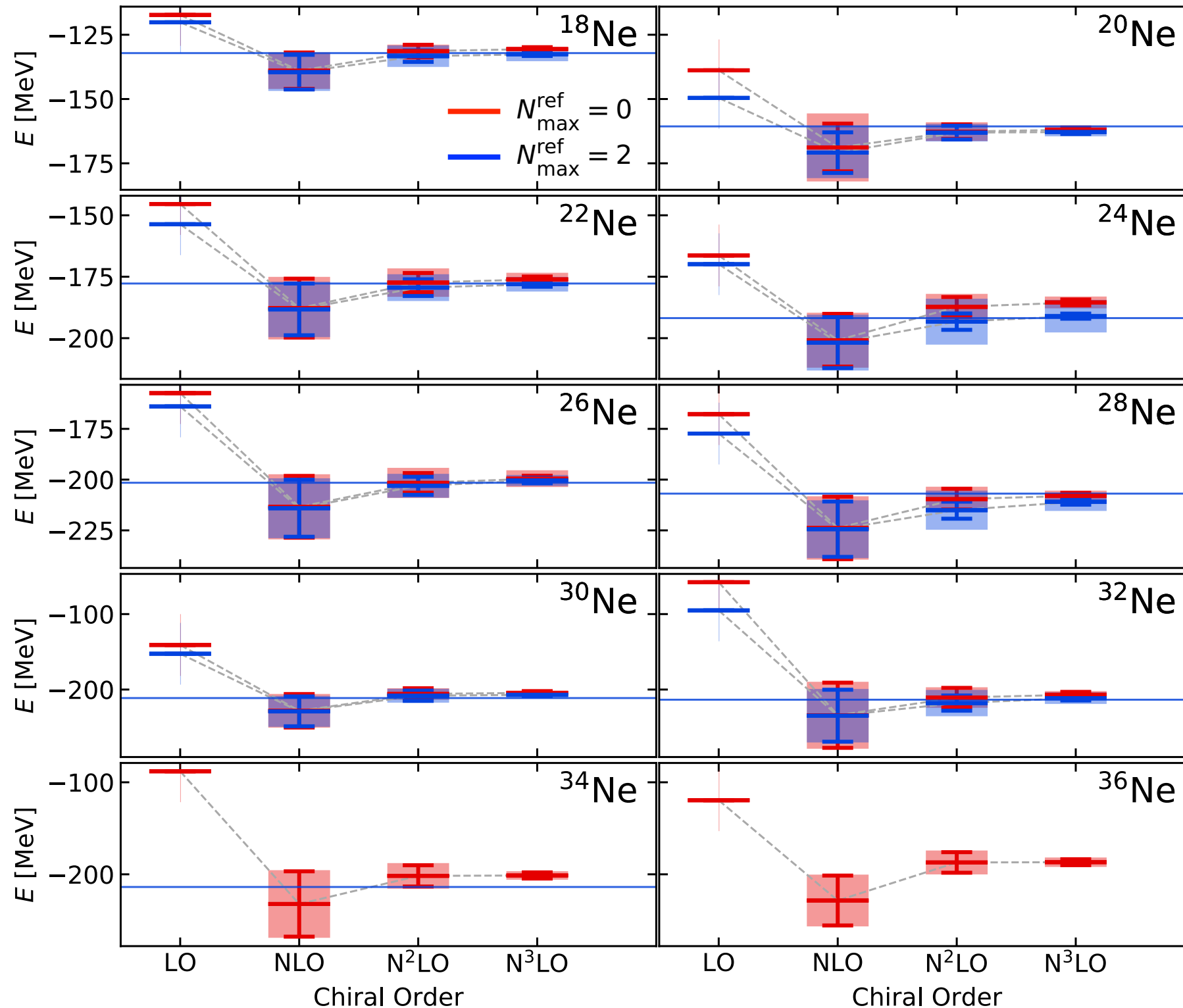
Hüther et al.; PLB 808, 135651 (2020)



NCSM/IM-NCSM, $\Lambda=500$ MeV, $\hbar\Omega=20$ MeV
error bands show interaction uncertainties

Applications: Neon Isotopes

Ground-State Energies



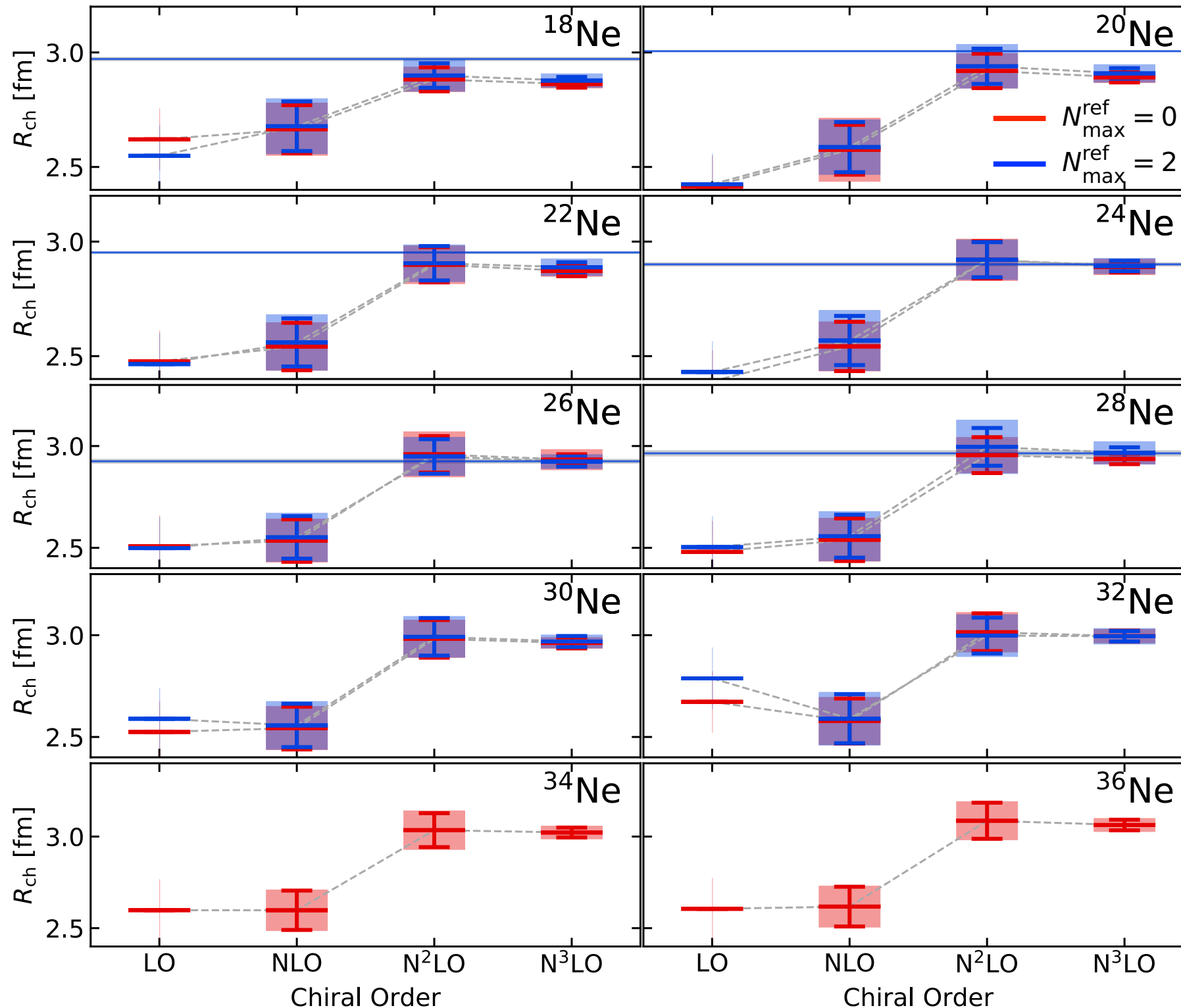
- amazing reproduction of experimental energies for all isotopes
- uncertainties under control

$\Lambda = 500 \text{ MeV}$
 $\alpha = 0.04 \text{ fm}^4$
 $\hbar\Omega = 20 \text{ MeV}$
 $e_{\text{max}} = 12$
 NAT basis
 $N_{\text{max}}^{\text{ref}} = 0, 2$
 $N_{\text{max}} = 4$

error bars:
 68% interaction
 uncertainties

error bands:
 interaction +
 many-body
 uncertainties

Charge Radii



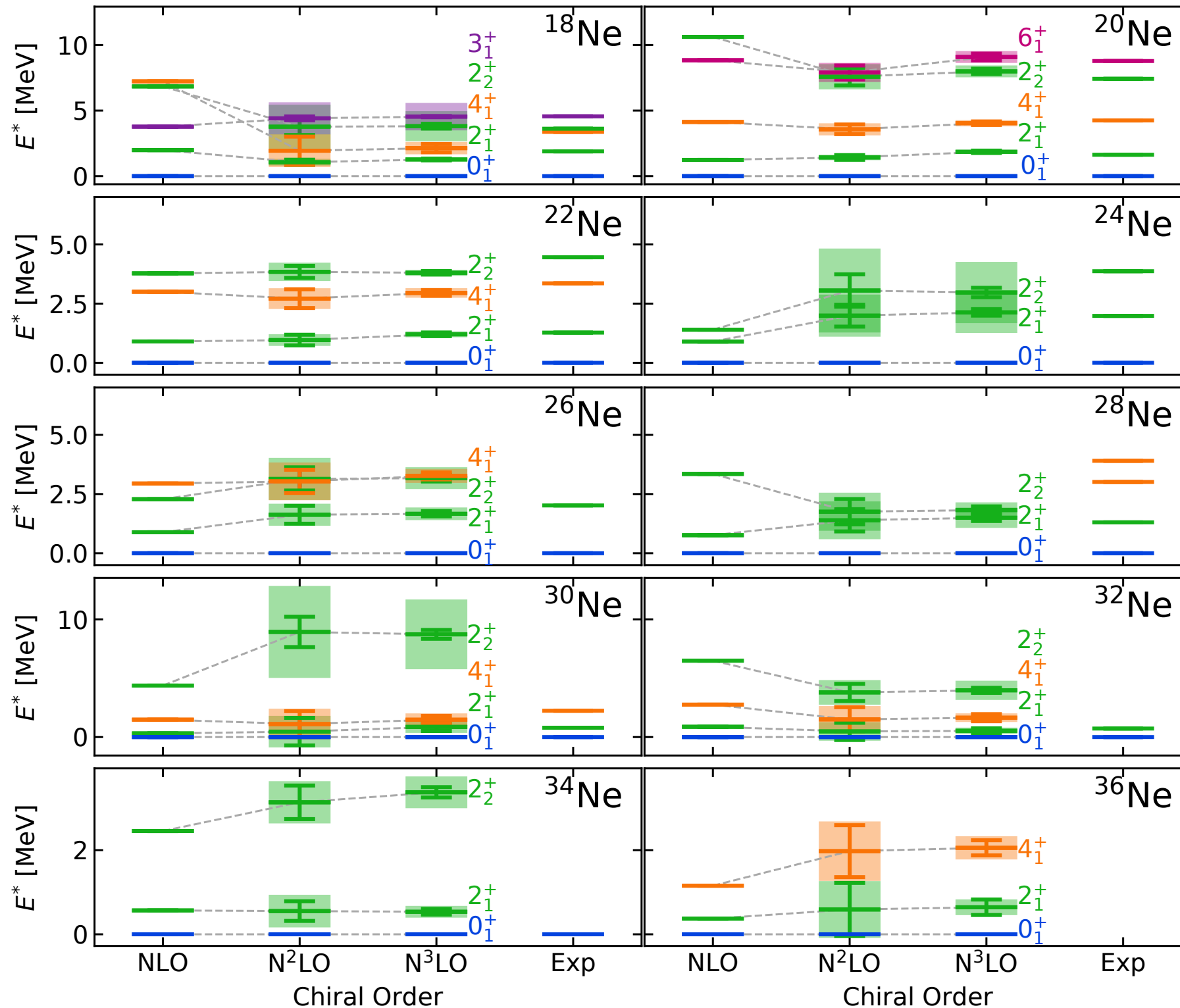
- excellent description of radii, slight underestimation for light isotopes
- stable results in N^2 LO and N^3 LO

$\Lambda = 500 \text{ MeV}$
 $\alpha = 0.04 \text{ fm}^4$
 $\hbar\Omega = 20 \text{ MeV}$
 $e_{\max} = 12$
 NAT basis
 $N_{\max}^{\text{ref}} = 0, 2$
 $N_{\max} = 4$

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Excitation Energies



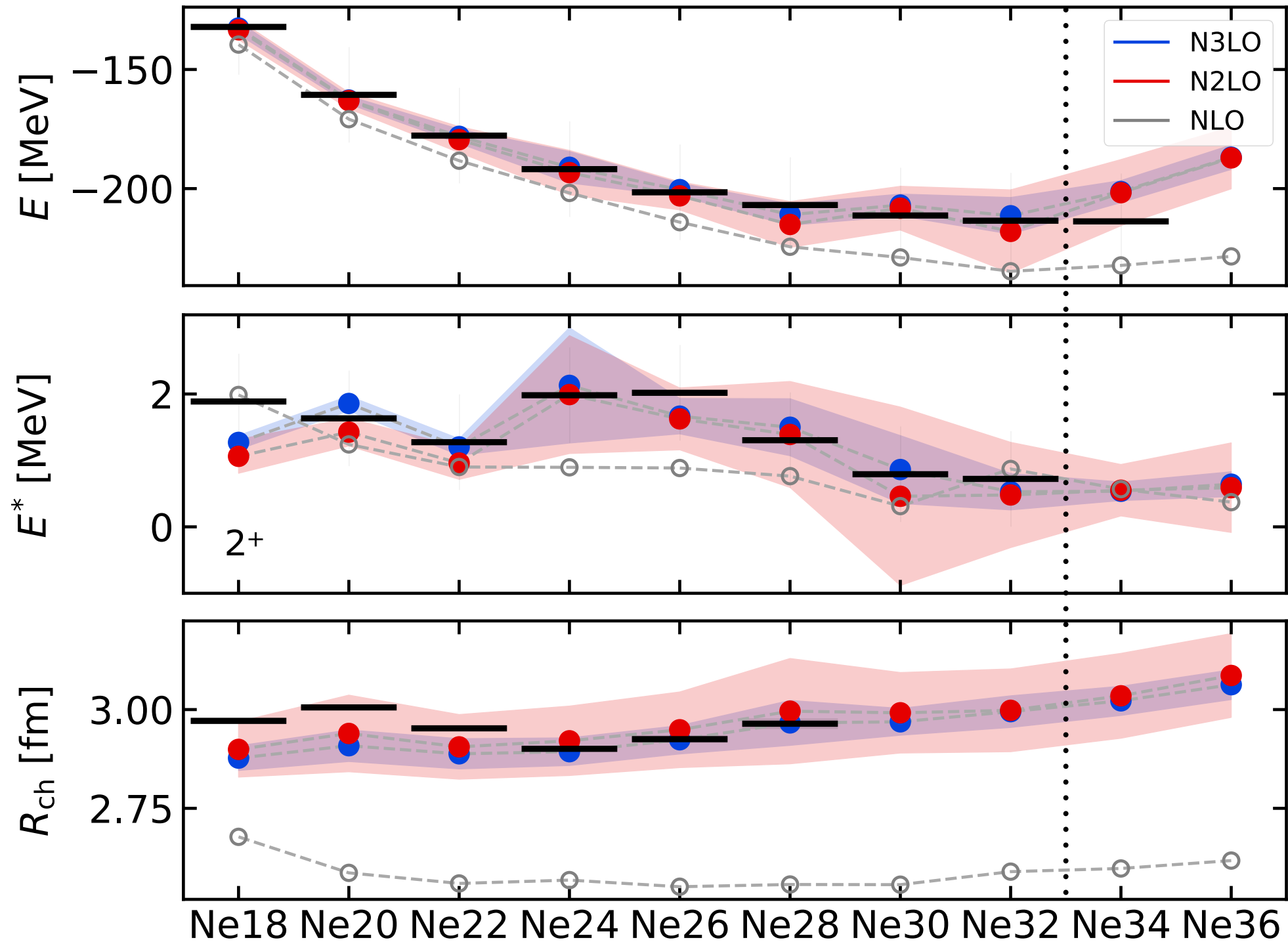
■ excellent description of excitation spectra

$\Lambda = 500 \text{ MeV}$
 $\alpha = 0.04 \text{ fm}^4$
 $\hbar\Omega = 20 \text{ MeV}$
 $e_{\text{max}} = 12$
 NAT basis
 $N_{\text{max}}^{\text{ref}} = 2$
 $N_{\text{max}} = 4$

error bars:
 68% interaction
 uncertainties

error bands:
 interaction +
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Neon Isotopes: Systematics



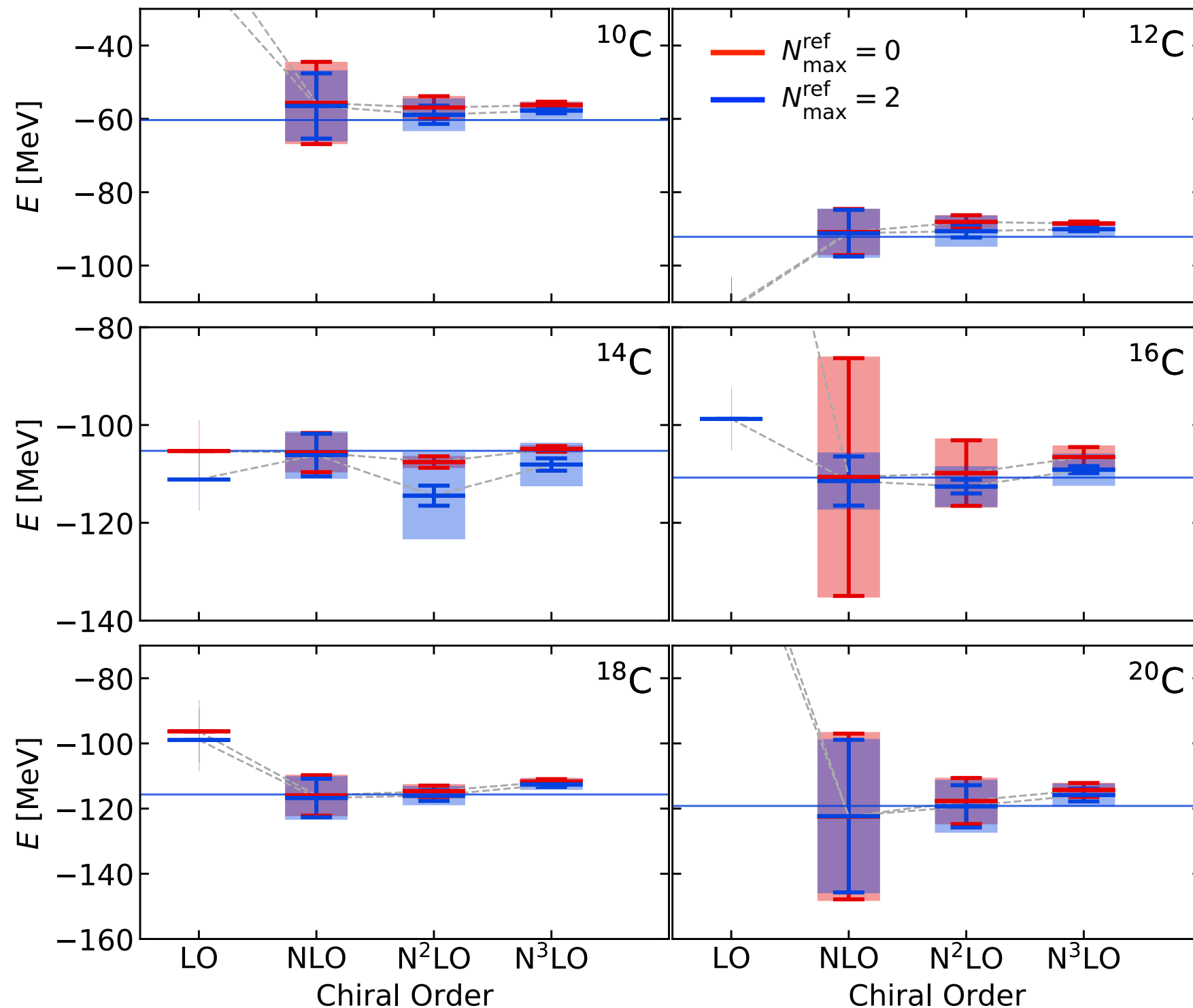
error bands:
interaction +
many-body
uncertainties

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 $\alpha = 0.04$ fm⁴
 $\hbar\Omega = 20$ MeV
 $e_{max} = 12$
NAT basis

$N_{max}^{ref} = 2, 0$
 $N_{max} = 4, 2$

Applications: Carbon Isotopes

Ground-State Energies



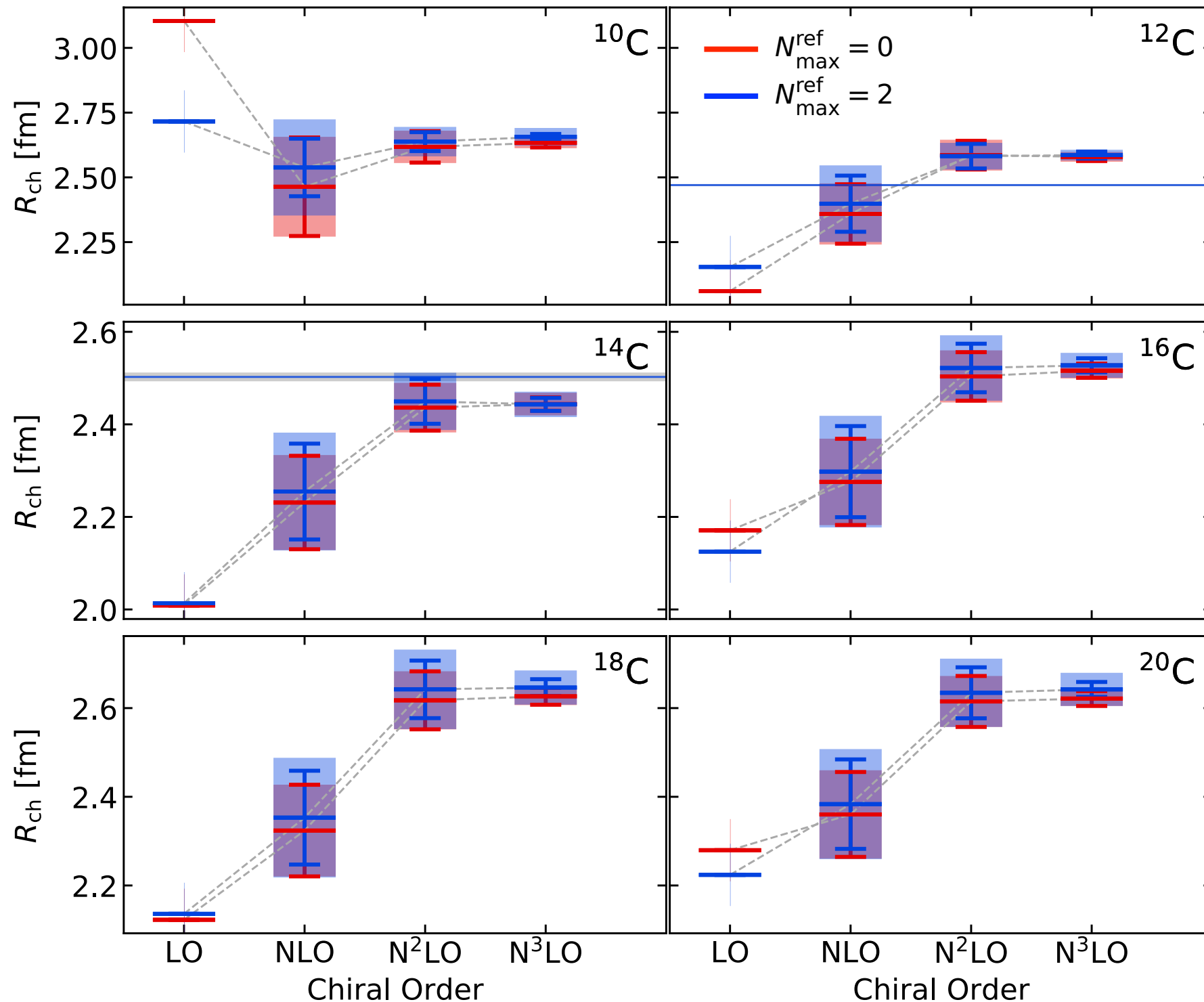
- good reproduction of experimental ground-state energies
- uncertainties under control

$\Lambda = 500$ MeV
 $\alpha = 0.04$ fm⁴
 $\hbar\Omega = 20$ MeV
 $e_{\max} = 12$
 NAT basis
 $N_{\max}^{\text{ref}} = 0, 2$
 $N_{\max} = 4$

error bars:
 68% interaction
 uncertainties

error bands:
 interaction +
 many-body
 uncertainties

Charge Radii



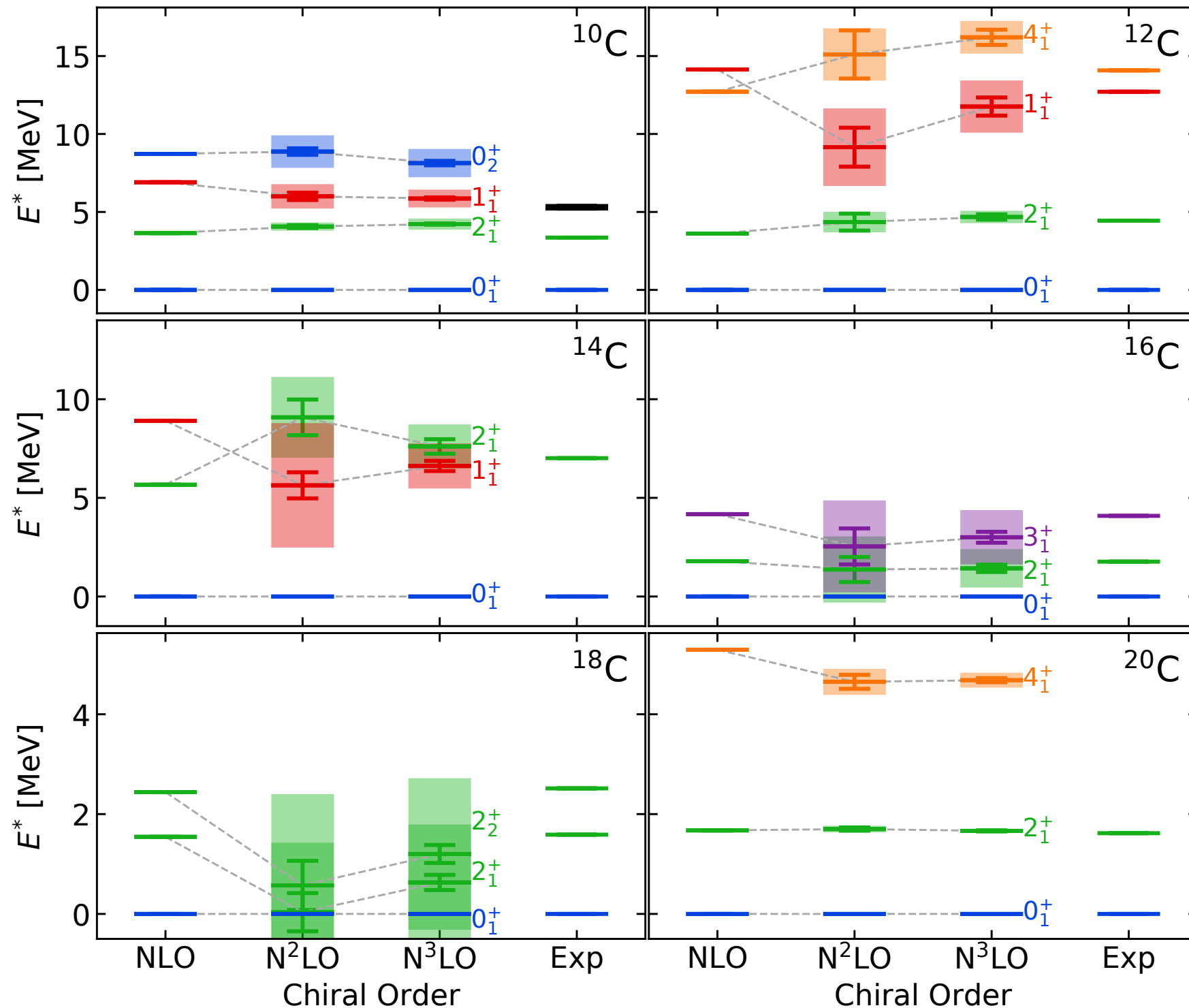
- slight overprediction of ^{12}C and slight underprediction of ^{14}C charge radius
- stable results in N^2LO and N^3LO

$\Lambda = 500 \text{ MeV}$
 $\alpha = 0.04 \text{ fm}^4$
 $\hbar\Omega = 20 \text{ MeV}$
 $e_{\text{max}} = 12$
 NAT basis
 $N_{\text{max}}^{\text{ref}} = 0, 2$
 $N_{\text{max}} = 4$

error bars:
 68% interaction
 uncertainties

error bands:
 interaction +
 many-body
 uncertainties

Excitation Spectra



- good overall description of spectra, except cluster states (not shown)

- stable results, uncertainties strongly state dependent

$\Lambda = 500 \text{ MeV}$
 $\alpha = 0.04 \text{ fm}^4$
 $\hbar\Omega = 20 \text{ MeV}$

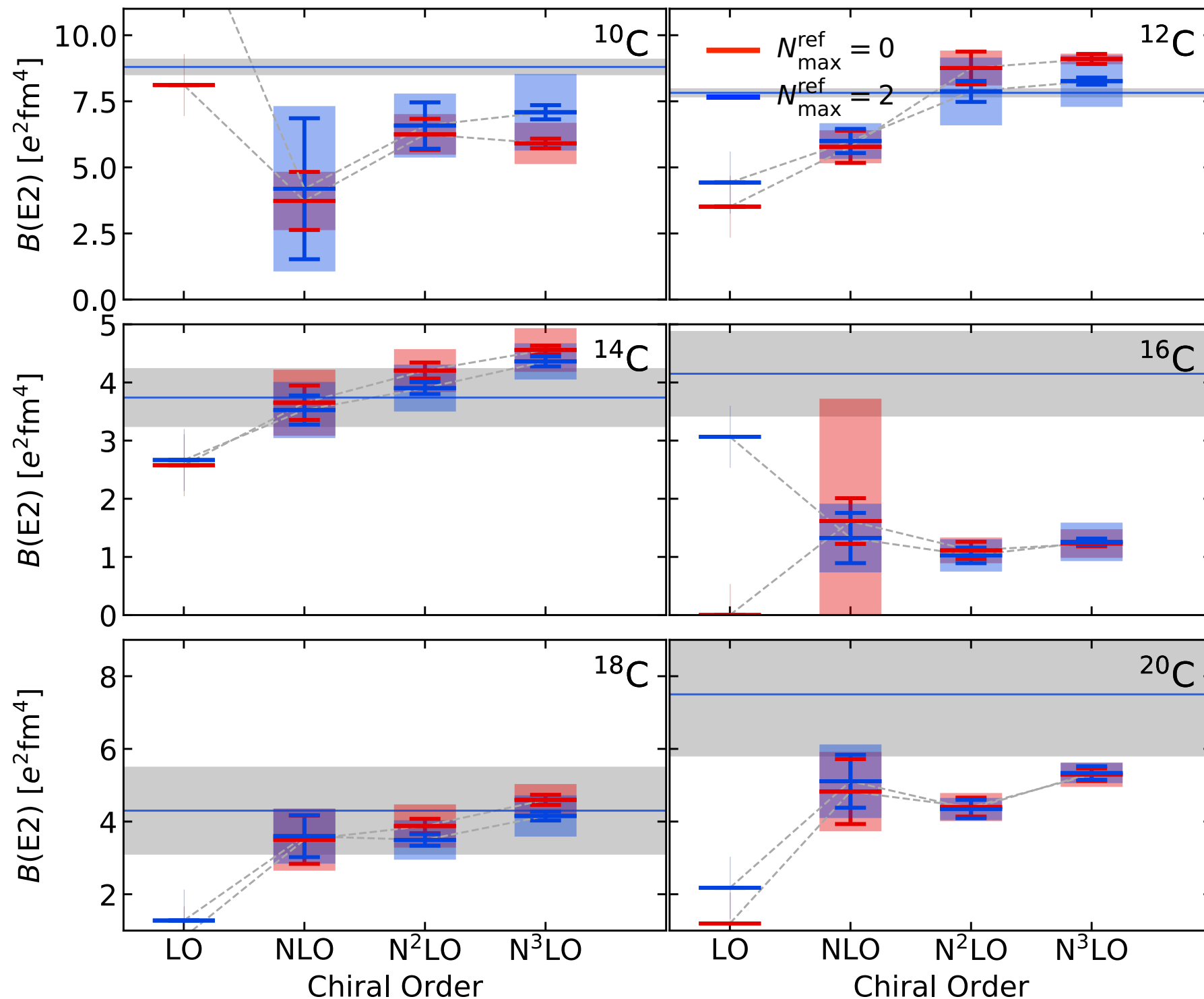
$e_{\text{max}} = 12$
 NAT basis

$N_{\text{max}}^{\text{ref}} = 2$
 $N_{\text{max}} = 4$

error bars:
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error bands:
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B(E2, $2^+ \rightarrow 0^+$) Transition Strength



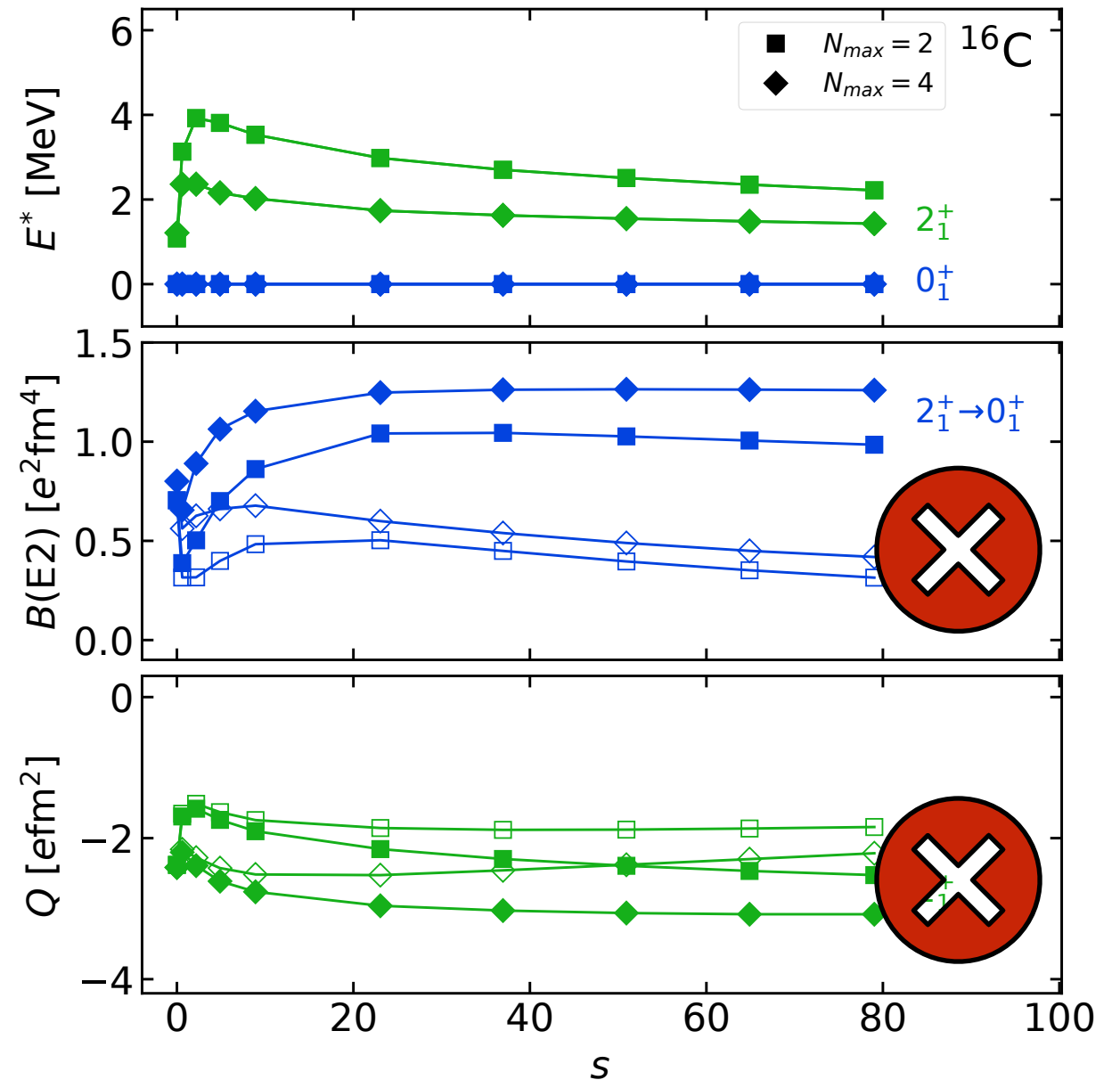
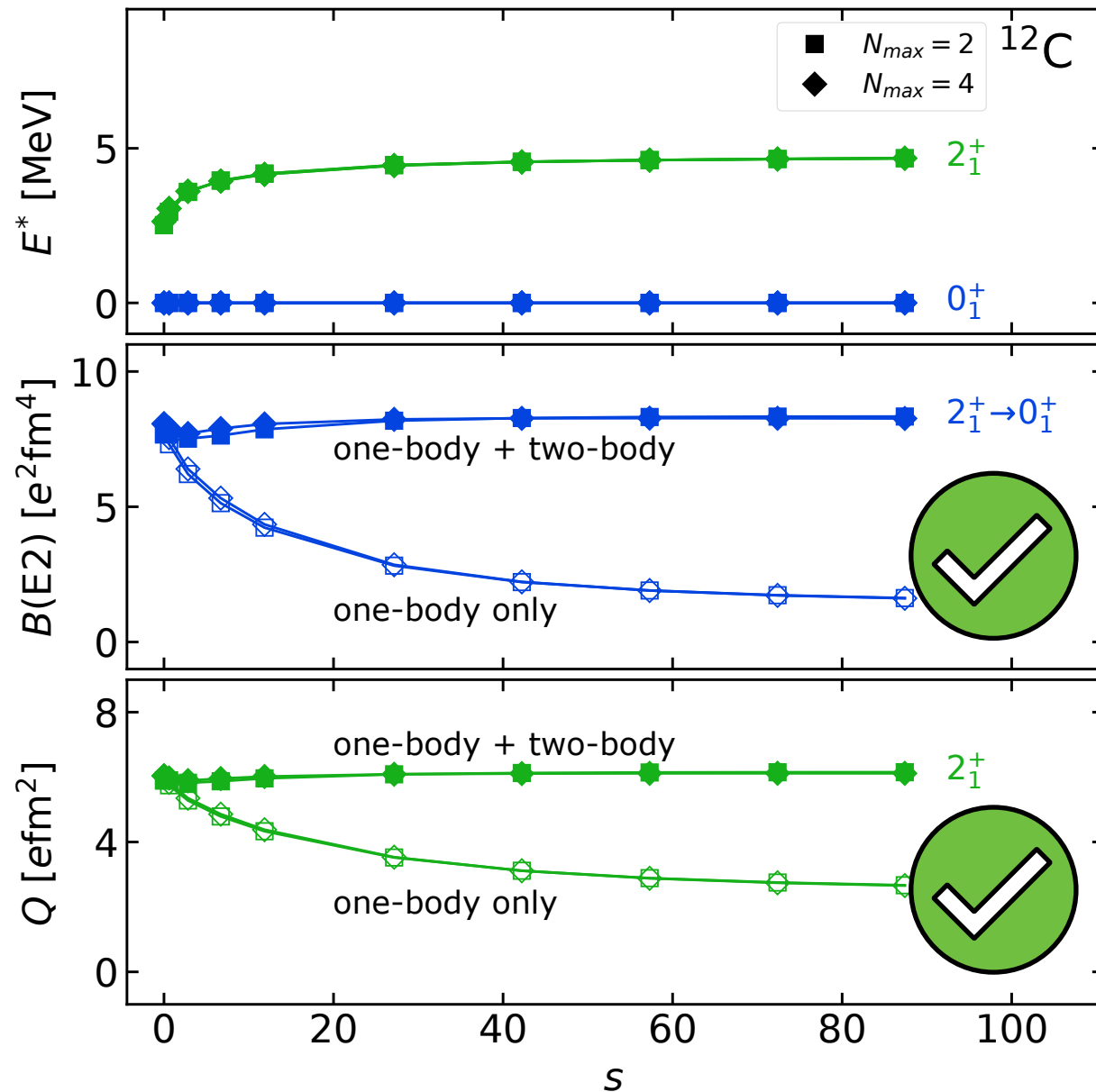
- agreement with experiment within uncertainties
- exception ^{16}C !

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 $\alpha = 0.04 \text{ fm}^4$
 $\hbar\Omega = 20 \text{ MeV}$
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 NAT basis
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Hierarchy Inversion



- IM-SRG evolution of E2 operator generates dominant (induced) two-body contribution... what about three-body and beyond?
- not a problem, if reference space contains the relevant static correlations

Next Stage: Active-Space IM-CI

■ **limitations of IM-NCSM setup**

- beyond sd-shell, the HO-based N_{\max} truncation does not make sense
- benefit from optimization of reference space to accommodate specific correlations

■ adopt a **more general CI strategy** for the definition of the reference space

- quantum chemistry: restricted active-space CI methods
- partitioning of single-particle orbits: hole - active - particle
- huge flexibility in model-space design

■ **perturbative corrections** to account for complete particle space

- use second-order MCPT with CI eigenstate as 'unperturbed' reference
- demonstrated successfully with the NCSM-PT *[Tichai et al., PLB 786, 448 (2018)]*

In-Medium CI for Collective Excitations

Collective Excitations

- traditional description via **RPA-type equations of motion methods**

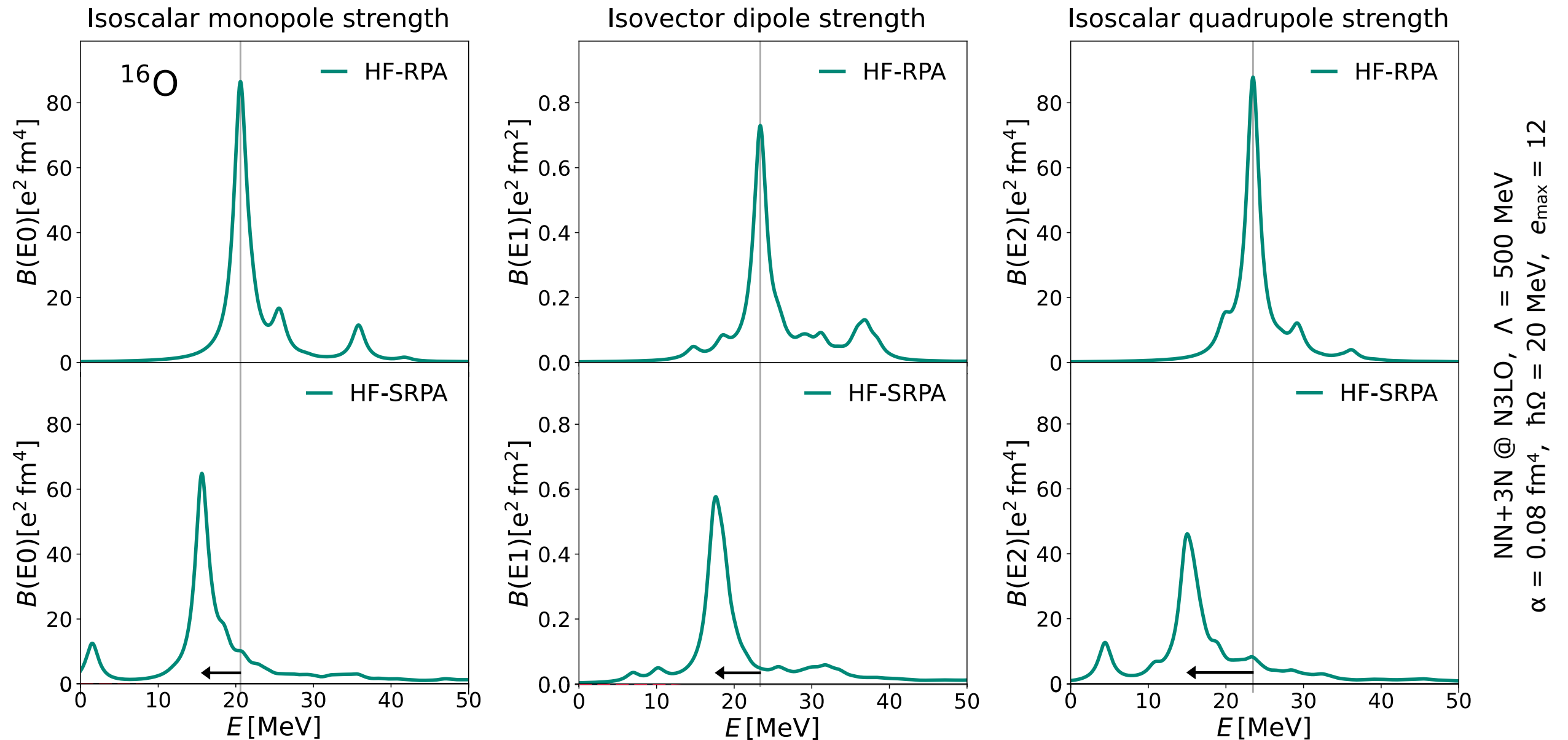
$$|E_\omega\rangle = Q_\omega^\dagger |\text{RPA}\rangle \qquad Q_\omega |\text{RPA}\rangle = 0$$

- phonon operator expanded in **particle-hole scheme** including "backwards" terms that imply **ground-state correlations**

$$\begin{aligned} Q_\omega^\dagger = & \sum_{p,h} \left(\chi_{ph}^\omega a_p^\dagger a_h - \gamma_{ph}^\omega a_h^\dagger a_p \right) && \textbf{RPA} \\ & + \sum_{pp',hh'} \left(\chi_{php'h'}^\omega a_p^\dagger a_h a_{p'}^\dagger a_{h'} - \gamma_{php'h'}^\omega a_h^\dagger a_p a_{h'}^\dagger a_{p'} \right) && \textbf{SRPA} \\ & + \dots \end{aligned}$$

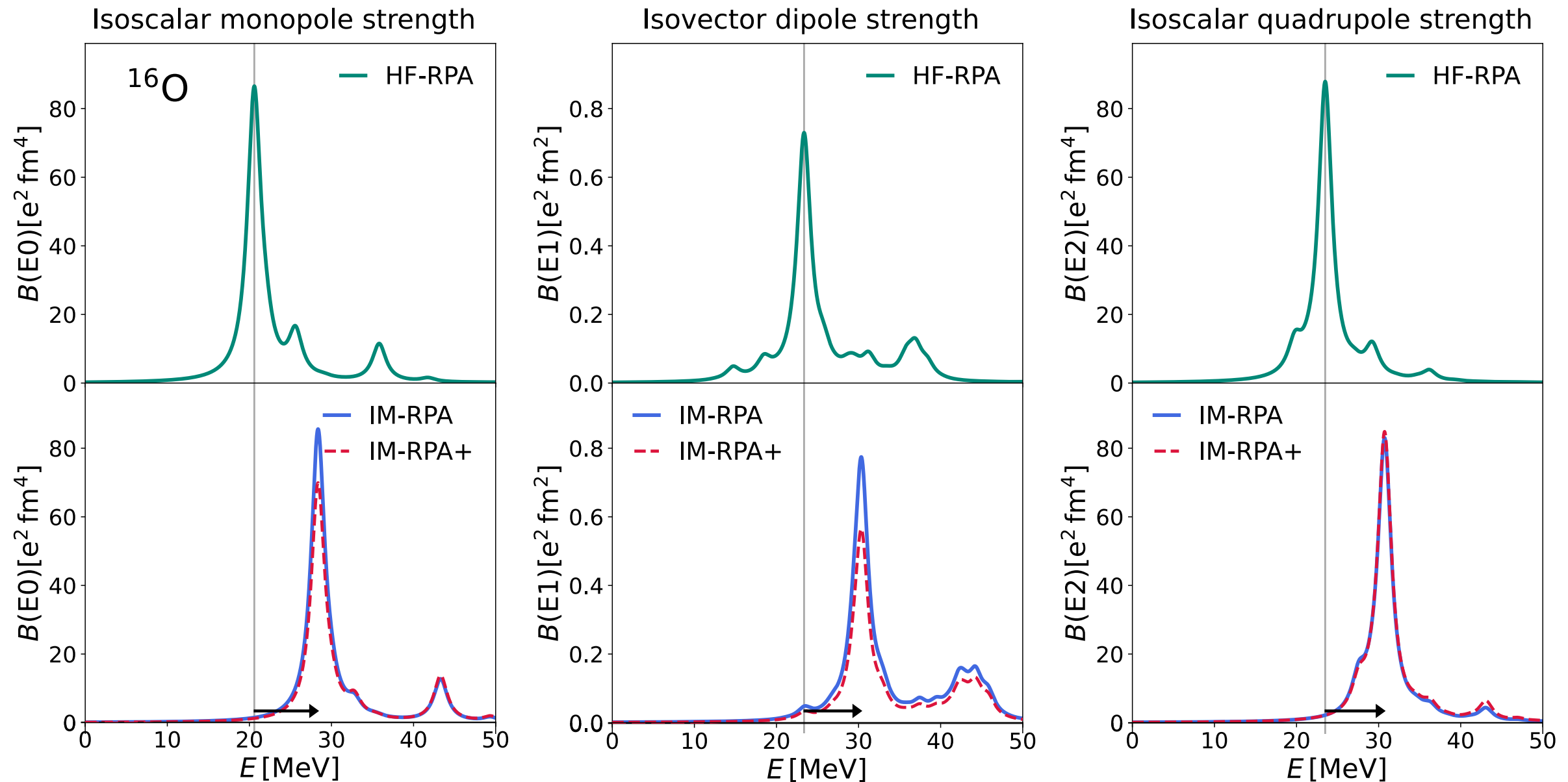
- might look old-fashioned, but is a systematic expansion scheme which can be used in an ab initio framework
- can we use RPA in connection **with IM-SRG evolved Hamiltonians** and what do we gain... decoupling targets ground state, not high-lying excitations

Standard HF-RPA and HF-SRPA



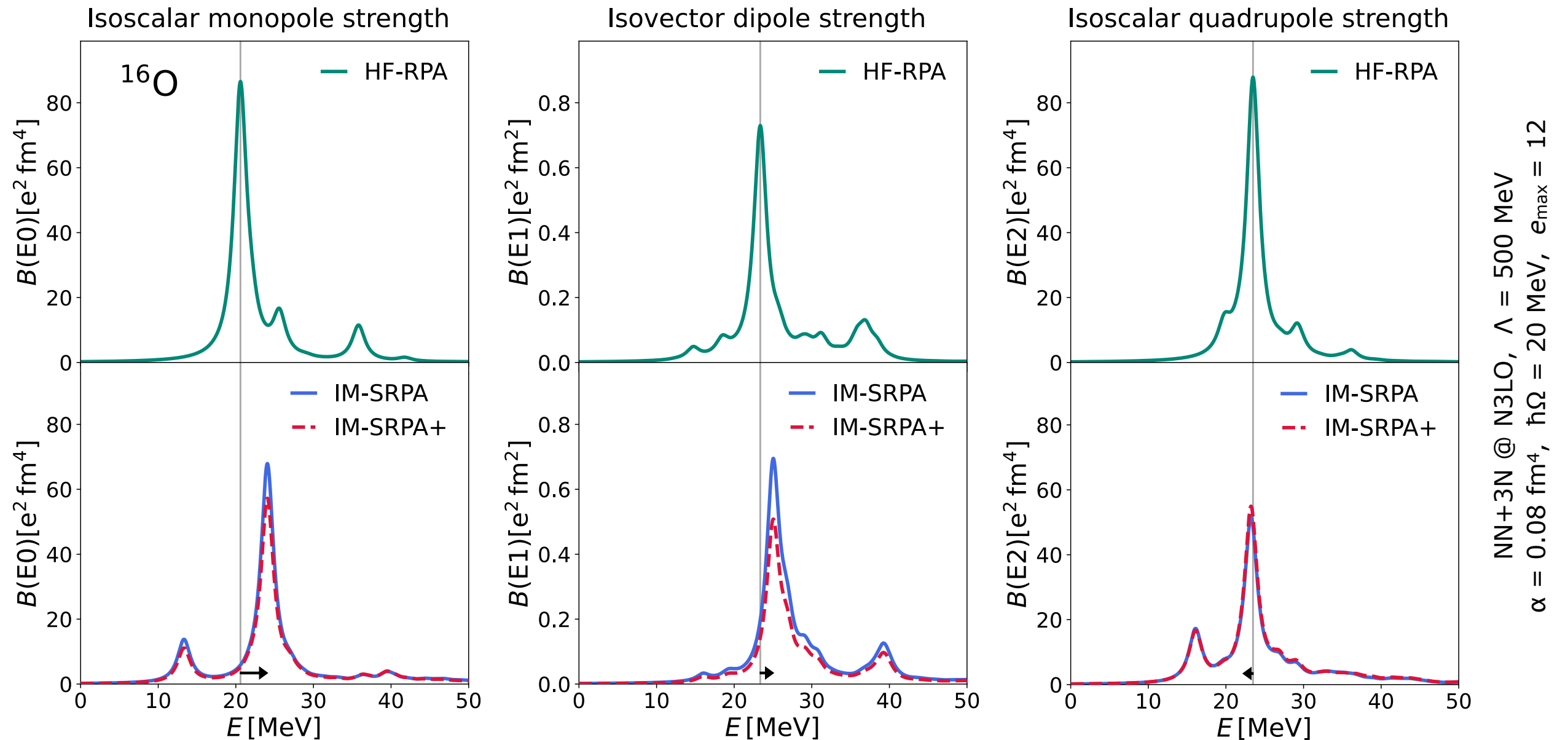
- **SRPA induces unphysical downward shift of resonances**
- 2p2h correlations are included for excited states, but ground state remains at HF level in quasi-boson approximation

SRPA Combined with IM-SRG



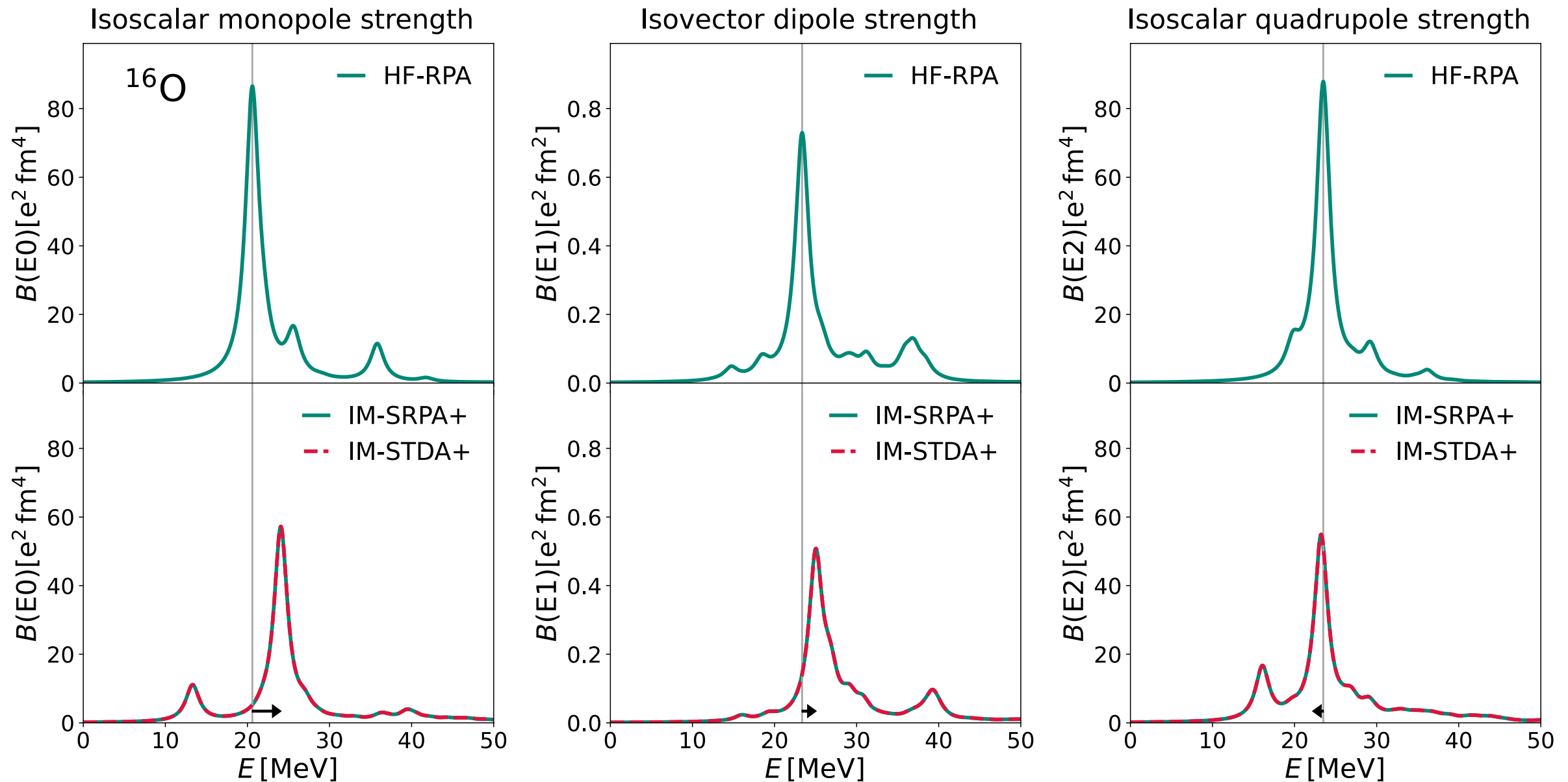
- **IM-SRG evolved Hamiltonian shifts resonances up in energy**
- decoupling designed to account for all ground-state correlations, not so for excited states

SRPA Combined with IM-SRG



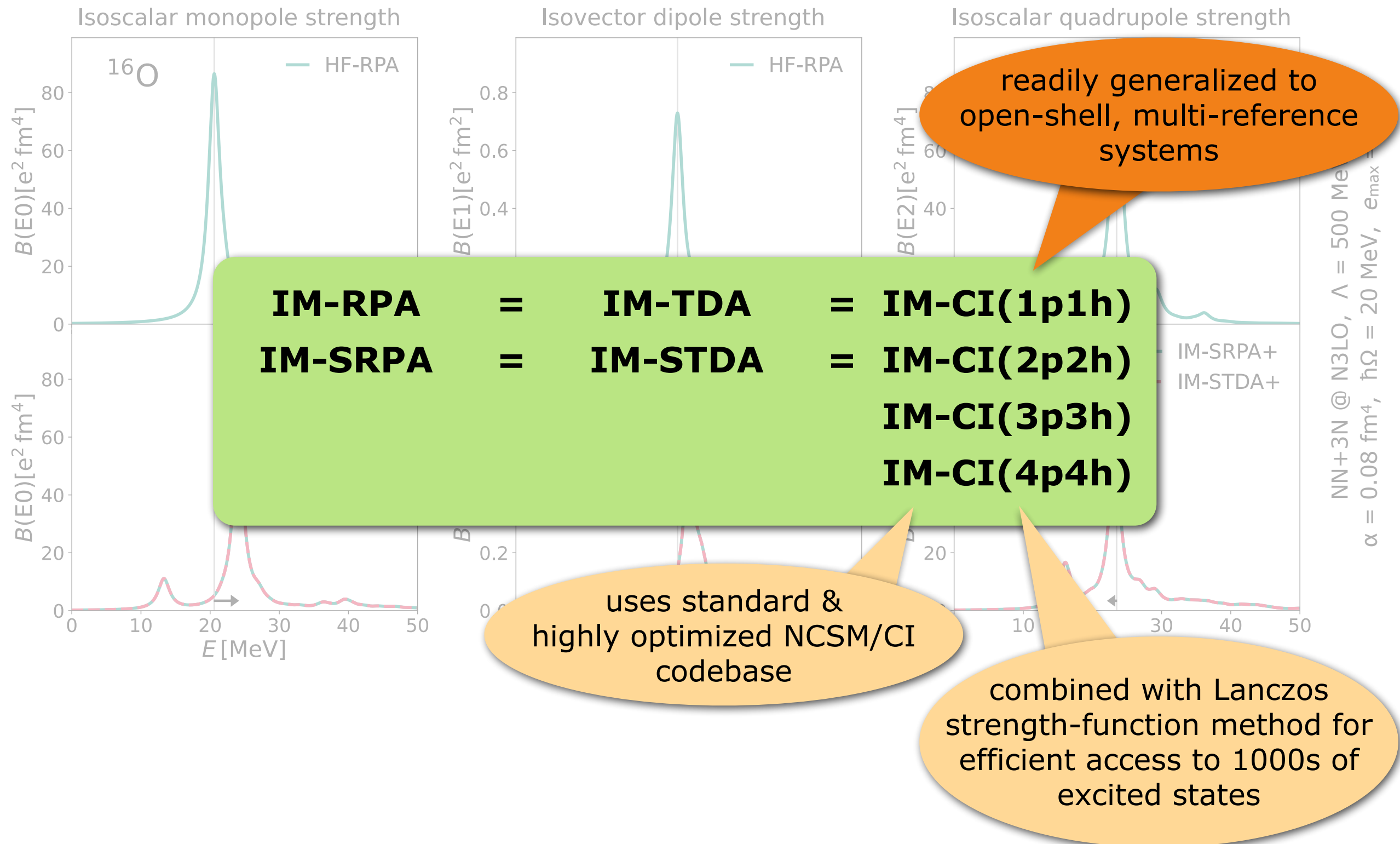
- **SRPA with IM-SRG evolved Hamiltonian moves resonances down**
- building 2p2h correlations explicitly into excited states, ground state is decoupled from the outset

IM-SRPA versus IM-STDA



- **simple (S)TDA calculation becomes equivalent to (S)RPA**
- IM-SRG decoupling suppresses $V_{\text{pp-hh}}$ matrix elements that make up the B -matrix, leading to vanishing "backward" amplitudes Y

IM for Collective Excitations



Epilogue

■ thanks to my group and my collaborators

- M.L. Agel, T. Gesser, P. Lehnung, M. Knöll, L. Mertes,
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& K. Hebeler, A. Tichai
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- P. Navrátil
TRIUMF, Vancouver
- H. Hergert
NSCL / Michigan State University
- J. Vary, P. Maris
Iowa State University
- E. Epelbaum, H. Krebs & the LENPIC Collaboration
Universität Bochum, ...



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Exzellente Forschung für
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