

REARRANGEMENT COLLISIONS AND INTER-CLUSTER POTENTIALS

IN THE ZERO-RANGE LIMIT

“Effective Field theory and Strong interaction with accurate error estimation”

12.4.2024 – CEA-ESNT

לוראנצו כונטסי	פאבון באלדאראמה
L. Contessi	M.P. Valderrama
Université Paris-Saclay, CNRS-IN2P3	Beihang University

צאוראב מונדאל	ראכשאנדה גוצואמא	ואראת ראהא
S. Mondal	R. Goswami	U. Raha
	IIT Guwahati	

⊗

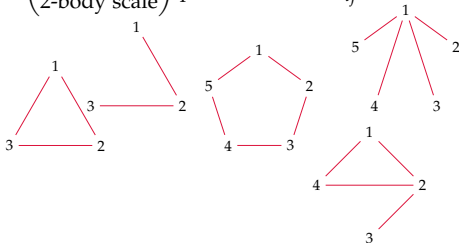
יוהנס קירשר
J. Kirscher
SRM University AP

The problem: What $\left(\begin{array}{c} \text{few-body} \\ \text{non-relativistic} \\ \text{quantum} \end{array} \right)$ -complexity can a

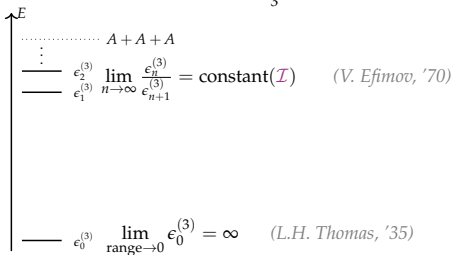
"The mind is localized and extends forever to infinity. The body is extended and remains localized." (R. Descartes)

universal/most versatile/elemental $\hat{=}$ $\left(\begin{array}{c} \text{no} \\ \text{2-body scale} \end{array} \right)$ pair interaction v_{ij} evoke?

$$H = \sum_{i=1}^A t_i + \sum_{\{ij\} \in \mathcal{I}} v_{ij}$$



$$v_{ij} = \begin{cases} c_0 \cdot \delta(|\mathbf{r}_{ij}|) \\ c_0 \cdot e^{-c_1 r_{ij}^2} \\ c_0 \cdot |\mathbf{r}_{ij}|^{-1} \cdot e^{-c_1 |\mathbf{r}_{ij}|} - c_3 \cdot |\mathbf{r}_{ij}|^{-6} \\ c_0 \cdot |\mathbf{r}_{ij}|^{-12} + c_1 \cdot |\mathbf{r}_{ij}|^{-6} \\ \vdots \end{cases} \Rightarrow$$



A NUCLEAR FEW-BODY THEORY.¹

$$\text{EFT}(\pi) = \begin{array}{cccc} \text{Degrees} & & \text{Naïve dim.} & \text{"Constructive"} \\ \text{of freedom} & \oplus & \text{analysis} & \text{scales} \\ \text{Symmetries} & & & \end{array}$$

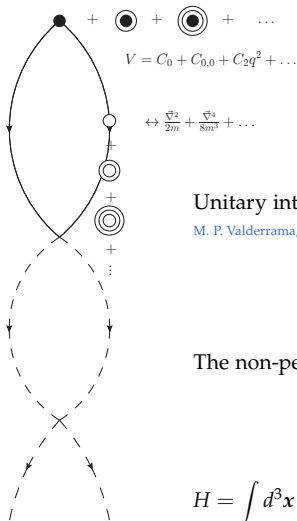
¹P. F. Bedaque, J.-W. Chen, H. W. Hammer, D. B. Kaplan, U. van Kolck, G. Rupak, M. J. Savage (199x-201y)

A UNIVERSAL FEW-BODY THEORY.¹

$$\text{EFT}(\not{\mu}) = \underbrace{\text{Degrees of freedom} \oplus \text{Symmetries} \oplus \text{Naïve dim. analysis} \oplus \text{“Constructive” scales}}_{\Rightarrow \text{systematic expansion of } \mathcal{L} \text{ and Amplitude}}$$

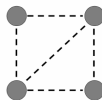
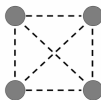
¹P. F. Bedaque, J.-W. Chen, H. W. Hammer, D. B. Kaplan, U. van Kolck, G. Rupak, M. J. Savage (199x-201y)

A NUCLEAR FEW-BODY THEORY.



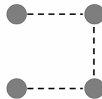
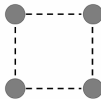
"full": $\mathcal{I}_{full} = \{(i, j) : i < j \leq 4\}$

"circle-slash": $\mathcal{I}_{cs} = \{(1, 2), (2, 3), (3, 4), (1, 4), (2, 4)\}$



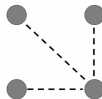
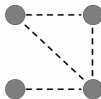
"circle": $\mathcal{I}_c = \{(i, \text{mod}(i, N) + 1) : i \leq 4\}$

"line": $\mathcal{I}_l = \{(i, i + 1) : i < 4\}$



"loner": $\mathcal{I}_4 = \{(1, 2), (2, 3), (1, 3), (3, 4)\}$

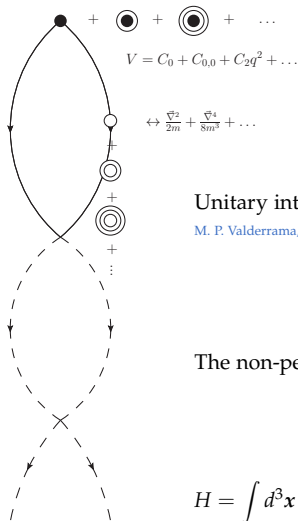
"star": $\mathcal{I}_{s^*} = \{(1, i) : i \neq 1\}$



$$H = \int d^3x \sum_i \left[\psi_i^\dagger \left(-\frac{\nabla^2}{2m} \right) \psi_i + C \cdot \sum_j \psi_i^\dagger \psi_j^\dagger \psi_i \psi_j + D \cdot \sum_k \psi_i^\dagger \psi_j^\dagger \psi_k^\dagger \psi_i \psi_j \psi_k \right]$$

$j \in \{V_{ij} : |a_{ij}| \rightarrow \infty\}$
 $j, k \in \{\geq 2 \text{ scattering lengths} \rightarrow \infty\}$

A UNIVERSAL FEW-BODY THEORY.

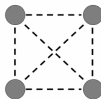


Unitary interaction graphs

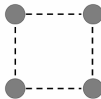
M. P. Valderrama, L. Contessi, JK (2021)

The non-perturbative leading order:

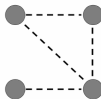
"full": $\mathcal{Z}_{full} = \{(i, j) : i < j \leq 4\}$



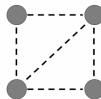
"circle": $\mathcal{Z}_{circle} = \{(i, \text{mod}(i, N) + 1) : i \leq 4\}$



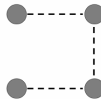
"loner": $\mathcal{Z}_4 = \{(1, 2), (2, 3), (1, 3), (3, 4)\}$



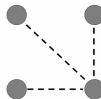
"circle-slash": $\mathcal{Z}_{cs} = \{(1, 2), (2, 3), (3, 4), (1, 4), (2, 4)\}$



"line": $\mathcal{Z}_4 = \{(i, i + 1) : i < 4\}$



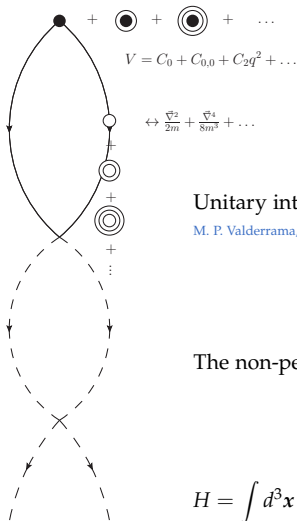
"star": $\mathcal{Z}_{star} = \{(1, i) : i \neq 1\}$



$$H = \int d^3x \sum_i \left[\psi_i^\dagger \left(-\frac{\nabla^2}{2m} \right) \psi_i + C \cdot \sum_{j \in \{V_{ij} : |a_{ij}| \rightarrow \infty\}} \psi_i^\dagger \psi_j^\dagger \psi_i \psi_j + D \cdot \sum \psi_i^\dagger \psi_j^\dagger \psi_k^\dagger \psi_i \psi_j \psi_k \right]$$

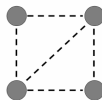
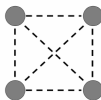
$j, k \in \{\geq 2 \text{ scattering lengths} \rightarrow \infty\}$

A NUCLEAR FEW-BODY THEORY.



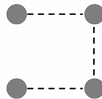
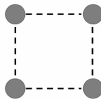
"full": $\mathcal{I}_{full} = \{(i, j) : i < j \leq 4\}$

"circle-slash": $\mathcal{I}_{cs} = \{(1, 2), (2, 3), (3, 4), (1, 4), (2, 4)\}$



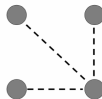
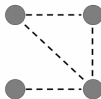
"circle": $\mathcal{I}_c = \{(i, \text{mod}(i, N) + 1) : i \leq 4\}$

"line": $\mathcal{I}_l = \{(i, i + 1) : i < 4\}$



"loner": $\mathcal{I}_4 = \{(1, 2), (2, 3), (1, 3), (3, 4)\}$

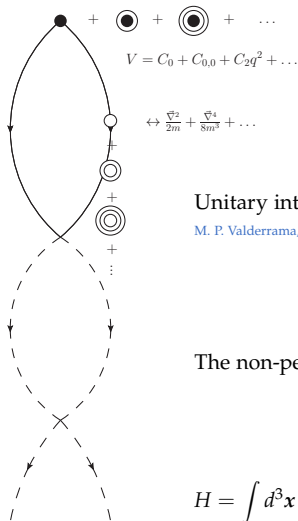
"star": $\mathcal{I}_{s^*} = \{(1, i) : i \neq 1\}$



$$H = \int d^3x \sum_i \left[\psi_i^\dagger \left(-\frac{\nabla^2}{2m} \right) \psi_i + C \cdot \sum_j \psi_i^\dagger \psi_j^\dagger \psi_i \psi_j + D \cdot \sum_k \psi_i^\dagger \psi_j^\dagger \psi_k^\dagger \psi_i \psi_j \psi_k \right]$$

$j \in \{V_{ij} : |a_{ij}| \rightarrow \infty\}$
 $j, k \in \{\geq 2 \text{ scattering lengths} \rightarrow \infty\}$

A UNIVERSAL FEW-BODY THEORY.

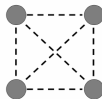


Unitary interaction graphs

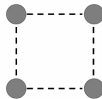
M. P. Valderrama, L. Contessi, JK (2021)

The non-perturbative leading order:

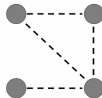
"full": $\mathcal{I}_{full} = \{(i, j) : i < j \leq 4\}$



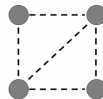
"circle": $\mathcal{I}_{circle} = \{(i, \text{mod}(i, N) + 1) : i \leq 4\}$



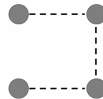
"loner": $\mathcal{I}_4 = \{(1, 2), (2, 3), (1, 3), (3, 4)\}$



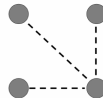
"circle-slash": $\mathcal{I}_{cs} = \{(1, 2), (2, 3), (3, 4), (1, 4), (2, 4)\}$



"line": $\mathcal{I}_4 = \{(i, i + 1) : i < 4\}$



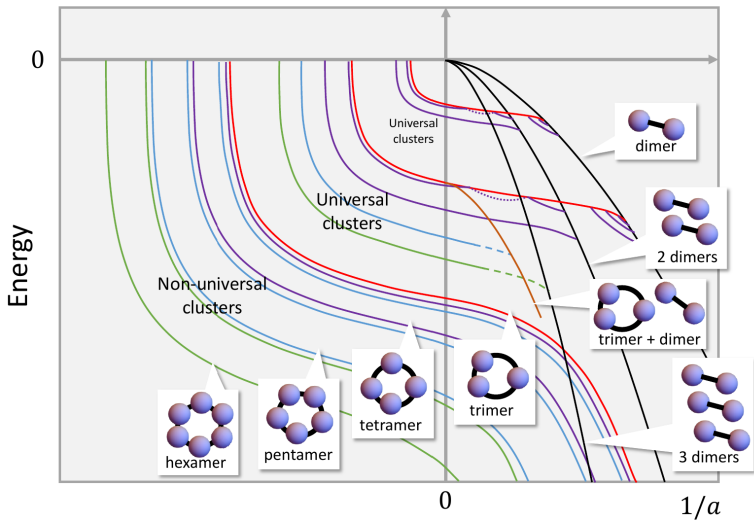
"star": $\mathcal{I}_{star} = \{(1, i) : i \neq 1\}$



$$H = \int d^3x \sum_i \left[\psi_i^\dagger \left(-\frac{\nabla^2}{2m} \right) \psi_i + C \cdot \sum_{j \in \{V_{ij} : |a_{ij}| \rightarrow \infty\}} \psi_i^\dagger \psi_j^\dagger \psi_i \psi_j + D \cdot \sum \psi_i^\dagger \psi_j^\dagger \psi_k^\dagger \psi_i \psi_j \psi_k \right]$$

$j, k \in \{\geq 2 \text{ scattering lengths} \rightarrow \infty\}$

An open problem for few-/many-body theories:
 $3+n$ spectra near/in the unitary limit.



THE REFINED¹ RESONATING GROUP METHOD²

THE ATOMIC NUCLEUS AS A SET OF MOLECULES.



$$\Psi = \hat{A} \left\{ \sum_i \phi(A_i)\phi(B_i)F_i(\mathbf{R}_i) + \sum_j \phi(A_j)\phi(B_j)\phi(C_j)F_j(\mathbf{R}_{1j}, \mathbf{R}_{2j}) \right. \\ \left. \sum_k \phi(A_k)\phi(B_k)\phi(C_k)\phi(D_k)F_k(\mathbf{R}_{1k}, \mathbf{R}_{2k}, \mathbf{R}_{3k}) \right. \\ \left. + \dots + \sum_m c_m \chi_m \right\} Z(\mathbf{R}_{c.m.})$$

$$\delta\Psi \stackrel{!}{=} \delta F$$

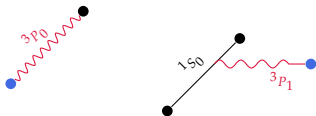
$$\Rightarrow \left(-\frac{\hbar^2}{2\mu} \nabla_{\mathbf{R}'}^2 + V_{\text{direct}}(\mathbf{R}') - E \right) F(\mathbf{R}') + \int K(\mathbf{R}', \mathbf{R}'') F(\mathbf{R}'') d\mathbf{R}'' = 0$$

¹H. M. Hofmann, K. Wildermuth, Y. C. Tang, H. H. Hackenbroich, D. R. Thompson, uvm.

²J. A. Wheeler PR 52 (1937)

THE REFINED¹ RESONATING GROUP METHOD²

THE ATOMIC NUCLEUS AS A SET OF MOLECULES.



$$\Psi = \hat{A} \left\{ \sum_i \phi(A_i)\phi(B_i)F_i(\mathbf{R}_i) + \sum_j \phi(A_j)\phi(B_j)\phi(C_j)F_j(\mathbf{R}_{1j}, \mathbf{R}_{2j}) \right. \\ \left. \sum_k \phi(A_k)\phi(B_k)\phi(C_k)\phi(D_k)F_k(\mathbf{R}_{1k}, \mathbf{R}_{2k}, \mathbf{R}_{3k}) \right. \\ \left. + \dots + \sum_m c_m \chi_m \right\} Z(\mathbf{R}_{c.m.})$$

$$\delta\Psi \stackrel{!}{=} \delta F$$

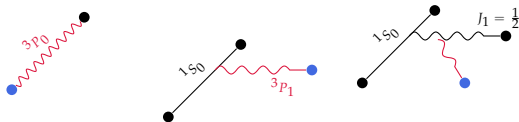
$$\Rightarrow \left(-\frac{\hbar^2}{2\mu} \nabla_{\mathbf{R}'}^2 + V_{\text{direct}}(\mathbf{R}') - E \right) F(\mathbf{R}') + \int K(\mathbf{R}', \mathbf{R}'') F(\mathbf{R}'') d\mathbf{R}'' = 0$$

¹H. M. Hofmann, K. Wildermuth, Y. C. Tang, H. H. Hackenbroich, D. R. Thompson, uvm.

²J. A. Wheeler PR 52 (1937)

THE REFINED¹ RESONATING GROUP METHOD²

THE ATOMIC NUCLEUS AS A SET OF MOLECULES.



$$\Psi = \hat{A} \left\{ \sum_i \phi(A_i)\phi(B_i)F_i(\mathbf{R}_i) + \sum_j \phi(A_j)\phi(B_j)\phi(C_j)F_j(\mathbf{R}_{1j}, \mathbf{R}_{2j}) \right. \\ \left. \sum_k \phi(A_k)\phi(B_k)\phi(C_k)\phi(D_k)F_k(\mathbf{R}_{1k}, \mathbf{R}_{2k}, \mathbf{R}_{3k}) \right. \\ \left. + \dots + \sum_m c_m \chi_m \right\} Z(\mathbf{R}_{c.m.})$$

$$\delta\Psi \stackrel{!}{=} \delta F$$

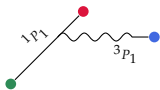
$$\Rightarrow \left(-\frac{\hbar^2}{2\mu} \nabla_{\mathbf{R}'}^2 + V_{\text{direct}}(\mathbf{R}') - E \right) F(\mathbf{R}') + \int K(\mathbf{R}', \mathbf{R}'') F(\mathbf{R}'') d\mathbf{R}'' = 0$$

¹H. M. Hofmann, K. Wildermuth, Y. C. Tang, H. H. Hackenbroich, D. R. Thompson, uvm.

²J. A. Wheeler PR 52 (1937)

THE REFINED¹ RESONATING GROUP METHOD²

THE ATOMIC NUCLEUS AS A SET OF MOLECULES.



$$\Psi = \hat{A} \left\{ \begin{aligned} & \sum_i \phi(A_i)\phi(B_i)F_i(\mathbf{R}_i) + \sum_j \phi(A_j)\phi(B_j)\phi(C_j)F_j(\mathbf{R}_{1j}, \mathbf{R}_{2j}) \\ & \sum_k \phi(A_k)\phi(B_k)\phi(C_k)\phi(D_k)F_k(\mathbf{R}_{1k}, \mathbf{R}_{2k}, \mathbf{R}_{3k}) \\ & + \dots + \sum_m c_m \chi_m \end{aligned} \right\} Z(\mathbf{R}_{c.m.})$$

$$\delta\Psi \stackrel{!}{=} \delta F$$

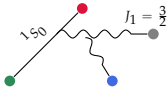
$$\Rightarrow \left(-\frac{\hbar^2}{2\mu} \nabla_{\mathbf{R}'}^2 + V_{\text{direct}}(\mathbf{R}') - E \right) F(\mathbf{R}') + \int K(\mathbf{R}', \mathbf{R}'') F(\mathbf{R}'') d\mathbf{R}'' = 0$$

¹H. M. Hofmann, K. Wildermuth, Y. C. Tang, H. H. Hackenbroich, D. R. Thompson, uvm.

²J. A. Wheeler PR 52 (1937)

THE REFINED¹ RESONATING GROUP METHOD²

THE ATOMIC NUCLEUS AS A SET OF MOLECULES.



$$\Psi = \hat{A} \left\{ \begin{aligned} & \sum_i \phi(A_i)\phi(B_i)F_i(\mathbf{R}_i) + \sum_j \phi(A_j)\phi(B_j)\phi(C_j)F_j(\mathbf{R}_{1j}, \mathbf{R}_{2j}) \\ & \sum_k \phi(A_k)\phi(B_k)\phi(C_k)\phi(D_k)F_k(\mathbf{R}_{1k}, \mathbf{R}_{2k}, \mathbf{R}_{3k}) \\ & + \dots + \sum_m c_m \chi_m \end{aligned} \right\} Z(\mathbf{R}_{c.m.})$$

$$\delta\Psi \stackrel{!}{=} \delta F$$

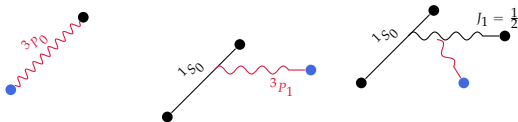
$$\Rightarrow \left(-\frac{\hbar^2}{2\mu} \nabla_{\mathbf{R}'}^2 + V_{\text{direct}}(\mathbf{R}') - E \right) F(\mathbf{R}') + \int K(\mathbf{R}', \mathbf{R}'') F(\mathbf{R}'') d\mathbf{R}'' = 0$$

¹H. M. Hofmann, K. Wildermuth, Y. C. Tang, H. H. Hackenbroich, D. R. Thompson, uvm.

²J. A. Wheeler PR 52 (1937)

THE REFINED¹ RESONATING GROUP METHOD²

THE ATOMIC NUCLEUS AS A SET OF MOLECULES.



$$\Psi = \hat{A} \left\{ \sum_i \phi(A_i)\phi(B_i)F_i(\mathbf{R}_i) + \sum_j \phi(A_j)\phi(B_j)\phi(C_j)F_j(\mathbf{R}_{1j}, \mathbf{R}_{2j}) \right. \\ \left. \sum_k \phi(A_k)\phi(B_k)\phi(C_k)\phi(D_k)F_k(\mathbf{R}_{1k}, \mathbf{R}_{2k}, \mathbf{R}_{3k}) \right. \\ \left. + \dots + \sum_m c_m \chi_m \right\} Z(\mathbf{R}_{c.m.})$$

$$\delta\Psi \stackrel{!}{=} \delta F$$

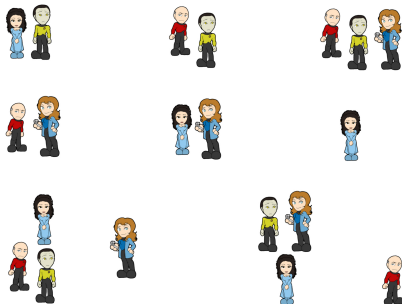
$$\Rightarrow \left(-\frac{\hbar^2}{2\mu} \nabla_{\mathbf{R}'}^2 + V_{\text{direct}}(\mathbf{R}') - E \right) F(\mathbf{R}') + \int K(\mathbf{R}', \mathbf{R}'') F(\mathbf{R}'') d\mathbf{R}'' = 0$$

¹H. M. Hofmann, K. Wildermuth, Y. C. Tang, H. H. Hackenbroich, D. R. Thompson, uvm.

²J. A. Wheeler PR 52 (1937)

THE REFINED RESONATING GROUP METHOD

THE ATOMIC NUCLEUS AS A SET OF MOLECULES.

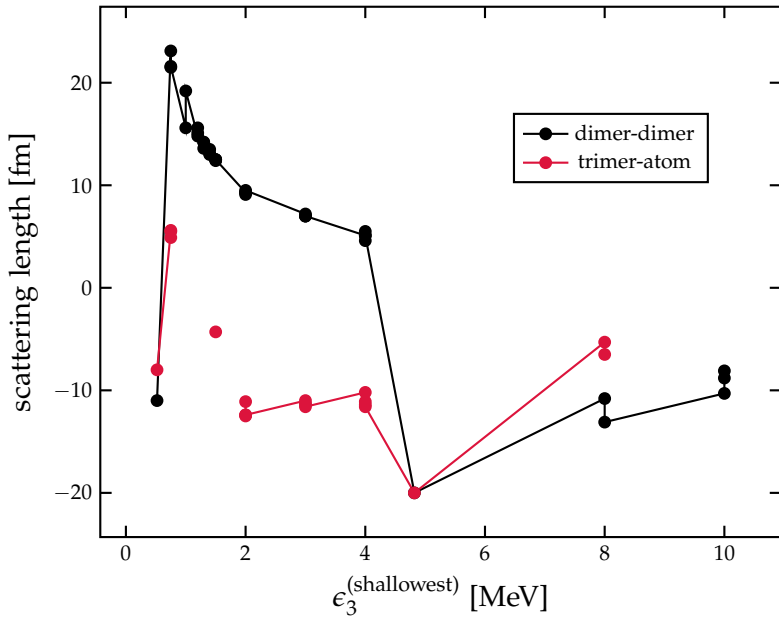


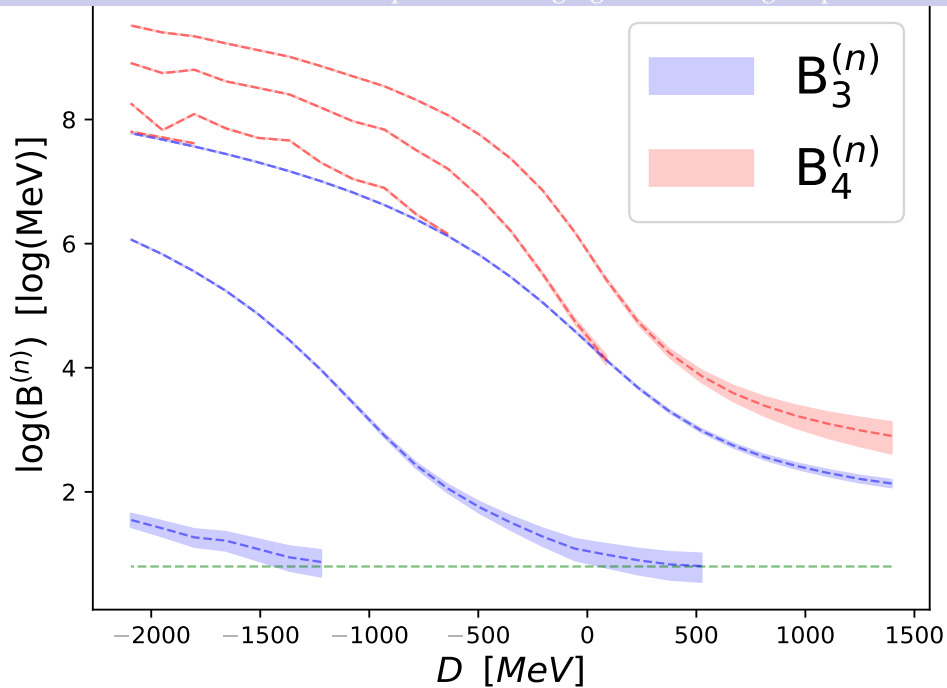
John Wheeler's idea (1937):

[...] It was as if, at a party, all the tall people clustered together at one moment, with all the short people in another cluster; then at the next moment [...] four groups formed, consisting of guests from the north, east, west, and south parts of the city; and so on, [...]

4-DISTINGUISHABLE-PARTICLE SCATTERING ($J^\pi = 0^+$)

FOR (ISO)SPIN INDEPENDENT V_{ij} AND $B(2) = 0.5 \text{ MeV}$ AND $m \approx 1 \text{ GeV}$ AND $4 \text{ fm}^{-2} \lesssim \lambda \lesssim 10 \text{ fm}^{-2}$

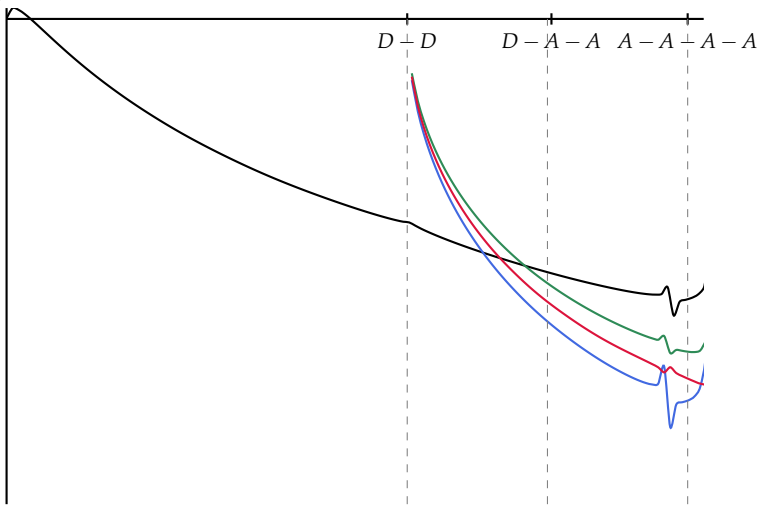




THE UNITARY 4-BODY SCATTERING PROBLEM

(2-FRAGMENT APPROXIMATION)

$E_{\text{c.m.}}$ [MeV]



δ [Deg]

INTER-CLUSTER POTENTIALS FROM CONSTITUENTS.

? MATCHING EFT'S WITH DIFFERENT DoF's

cluster degree(s) of freedom

≡

element of the spectrum of the EFT at the desired order:

$$\hat{H}_{\text{EFT}}\phi_A^{(n)} = e_A^{(n)}\phi_A^{(n)}$$

e.g. contact leading order:

$$\hat{V} = C_0(\Lambda) \sum_{i<j}^X \delta_\Lambda^{(3)}(\mathbf{r}_i - \mathbf{r}_j) + D_1(\Lambda) \sum_{\substack{i<j<k \\ \text{cyclic}}}^X \delta_\Lambda^{(3)}(\mathbf{r}_i - \mathbf{r}_j) \delta_\Lambda^{(3)}(\mathbf{r}_j - \mathbf{r}_k)$$

retain in(ter) cluster indistinguishability:

$$\hat{\mathcal{A}} = \hat{\mathcal{A}}_{AB}\hat{\mathcal{A}}_A\hat{\mathcal{A}}_B = \mathbb{1} + \sum_{P \in S_{A+B}} \text{sign}(P)\hat{P}$$

variation in the relative motion **between** the composites:

$$\begin{aligned} & \left(\hat{T}_{\mathbf{R}} - E_{\text{rel}} + \mathbb{N}^{-1} \langle \phi_A \phi_B | \hat{V} | \phi_A \phi_B \rangle \right) \chi(\mathbf{R}) \\ -\mathbb{N}^{-1} \int d\mathbf{R}' & \left[\langle \phi_A \phi_B | (\hat{T}_{\mathbf{R}} - E_{\text{rel}} + \hat{V}) \hat{P} \{ | \phi_A \phi_B \rangle \delta_\Lambda^{(3)}(\mathbf{R} - \mathbf{R}') \} \right] \chi(\mathbf{R}') = 0 \end{aligned}$$

INTER-CLUSTER POTENTIALS FROM CONSTITUENTS.

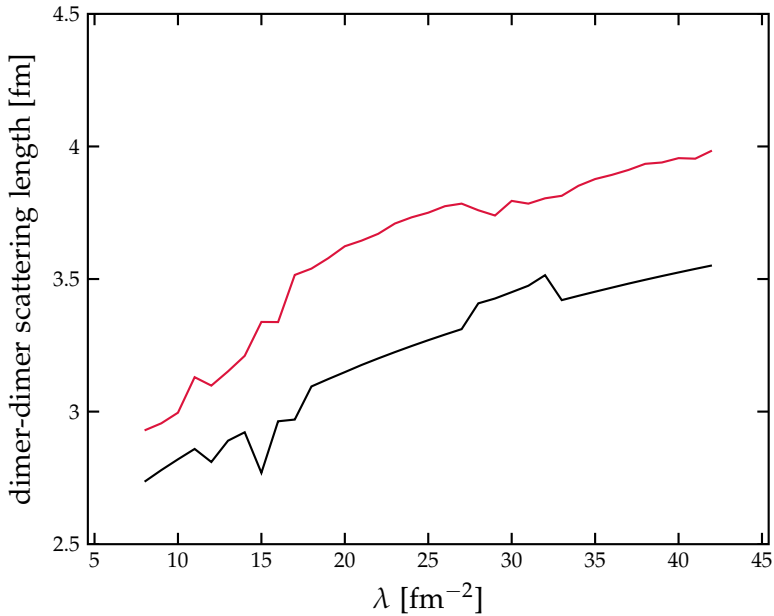
[?] MATCHING EFT'S WITH DIFFERENT DoF'S

$$\sum_{n=1}^{N_{\text{loc}}} \hat{\eta}_n e^{-w_n \mathbf{R}^2} \chi(\mathbf{R}) - \sum_{n=1}^{N_{\text{n-loc}}} \int \left\{ \hat{\zeta}_n e^{-a_n \mathbf{R}^2 - b_n \mathbf{R} \cdot \mathbf{R}' - c_n \mathbf{R}'^2} \right\} \chi(\mathbf{R}') d\mathbf{R}' = 0$$

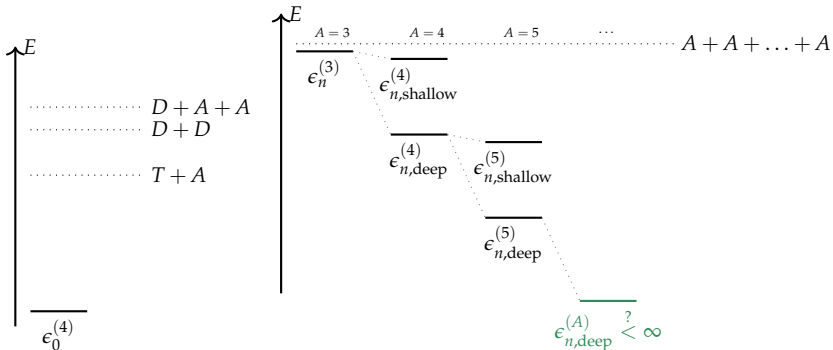
with $\hat{\eta}_n, \hat{\zeta}_n, w_n, a_n, b_n, c_n$ dependent upon $C(\lambda), D(\lambda), \alpha, \lambda, E, A, B$.

DIMER-DIMER SCATTERING – DoF: DIMER

UNIVERSAL/REGULATOR-INDEPENDENT WAVE FUNCTION $B(2) \approx 2.2$ MeV



UNDERSTANDING CORRELATIONS BETWEEN few- AND (few+n)-BODY SYSTEMS



*Some replied that happiness was a complex artifact and that man's aim lay not in happiness but in the zealous fulfillment of **historical laws**. And others said that happiness was a matter of **out-and-out** struggle, which would last eternally.*

(Andrei Platonov, "Chevengur")