

# Cross-section calculation using QMC and the short-time approximation

4th International Workshop on Quantitative Challenges in Short-Range Correlations and the EMC Effect Research

CEA Paris-Saclay

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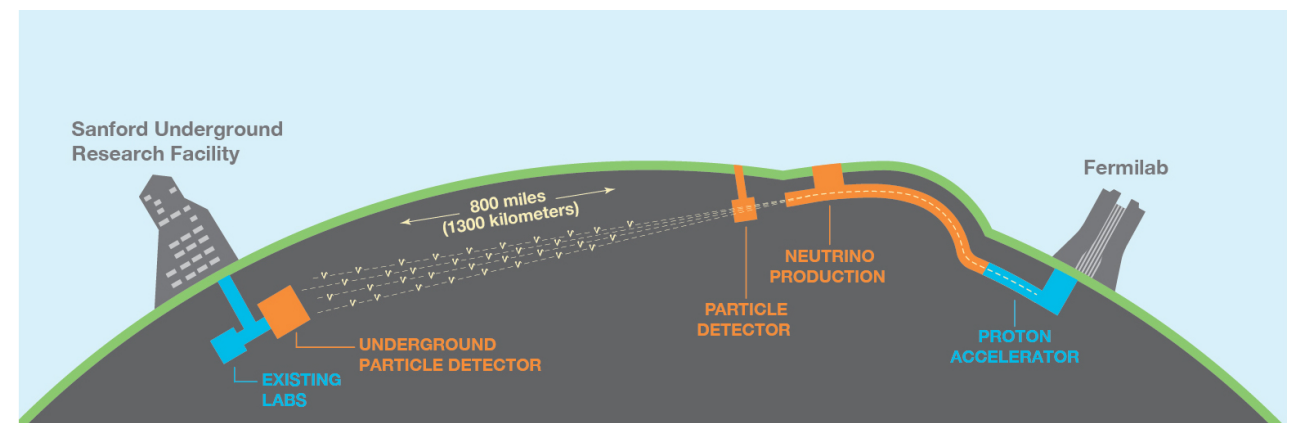
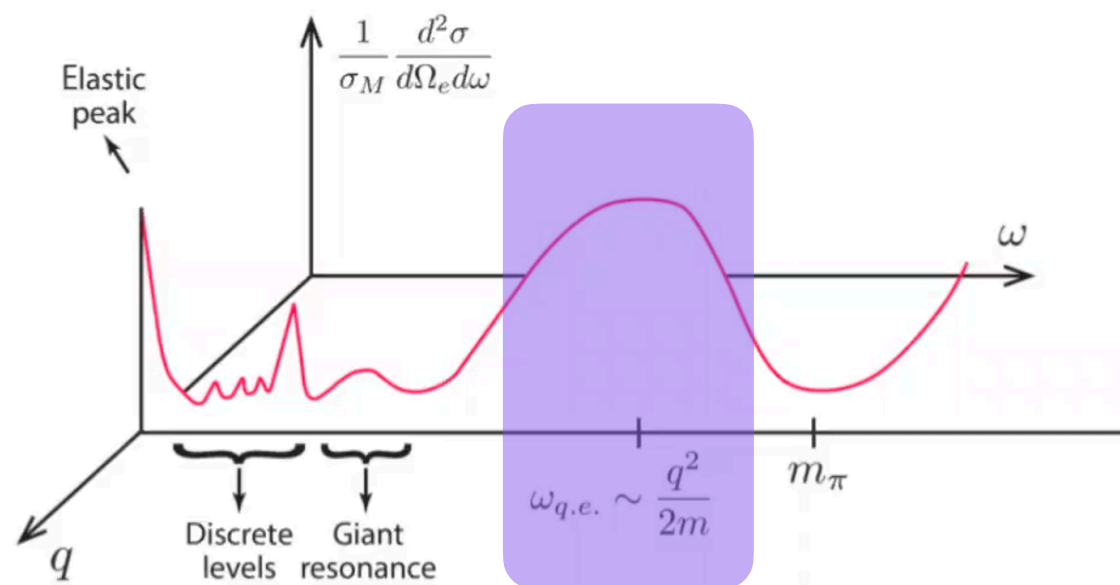
Maria Piarulli and Saori Pastore

 Washington University in St. Louis



# Electron-nucleus scattering

Theoretical understanding of nuclear effects is extremely important for neutrino experimental programs



Lepton-nucleus cross sections  $\omega \sim 10^2$  MeV

# Ab-initio description of nuclei



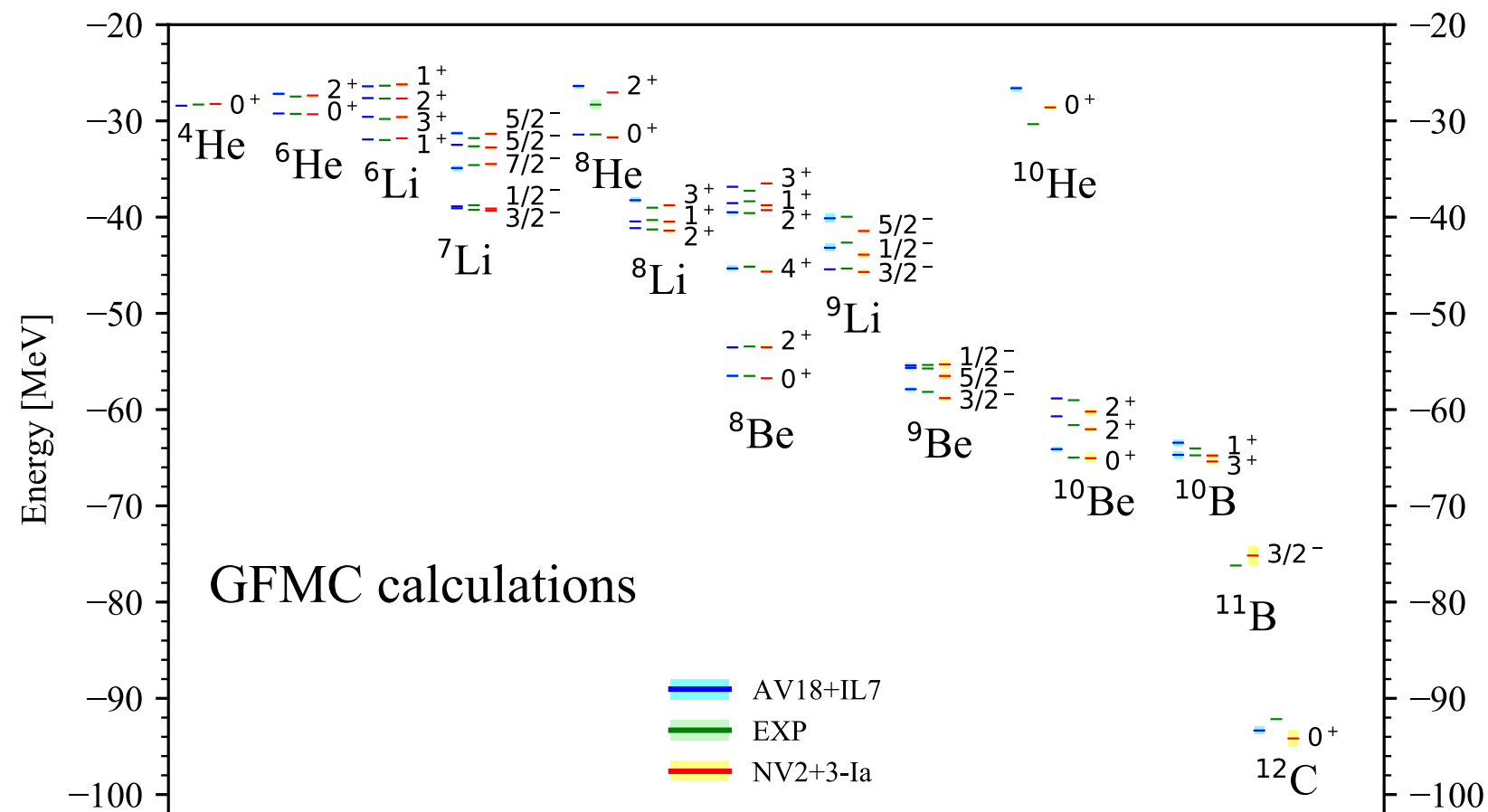
- Nuclear interaction
- Electroweak interaction of leptons with nucleons and clusters of correlated nucleons
- Computational method

# Many-body nuclear interaction



Many-body Nuclear Hamiltonian: Argonne  $v_{18}$  + Urbana IX

$$H = \sum_i T_i + \sum_{i < j} v_{ij} + \sum_{i < j < k} v_{ijk}$$





# Many-body nuclear interaction



Many-body Nuclear Hamiltonian: Argonne  $v_{18}$  + Urbana IX

$$H = \sum_i T_i + \sum_{i < j} v_{ij} + \sum_{i < j < k} v_{ijk}$$

Quantum Monte Carlo methods:

Use nuclear wave functions that minimize the expectation value of  $E$

$$E_V = \frac{\langle \psi | H | \psi \rangle}{\langle \psi | \psi \rangle} \geq E_0$$

The evaluation is performed using Metropolis sampling

# Nuclear Wave Functions



Variational wave function for nucleus in J state

$$|\psi\rangle = \mathcal{S} \prod_{i<j}^A \left[ 1 + U_{ij} + \sum_{k \neq i,j}^A U_{ijk} \right] \left[ \prod_{i<j} f_c(r_{ij}) \right] |\Phi(JMTT_3)\rangle$$

Two-body spin- and isospin-dependent correlations

$$U_{ij} = \sum_p f^p(r_{ij}) O_{ij}^p$$

$$O_{ij}^p = [1, \boldsymbol{\sigma}_i \cdot \boldsymbol{\sigma}_j, S_{ij}] \otimes [1, \boldsymbol{\tau}_i \cdot \boldsymbol{\tau}_j]$$

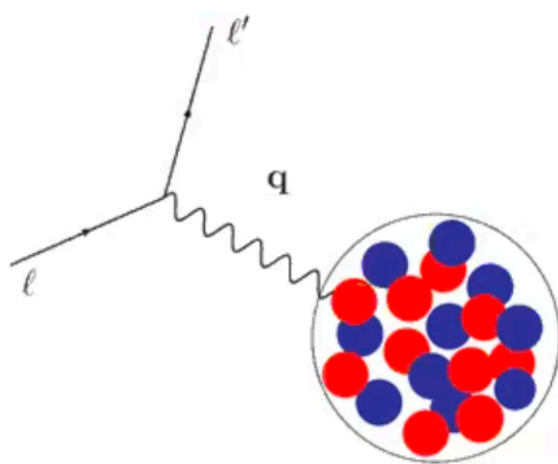
$$U_{ijk} = \epsilon v_{ijk}(\bar{r}_{ij}, \bar{r}_{jk}, \bar{r}_{ki})$$



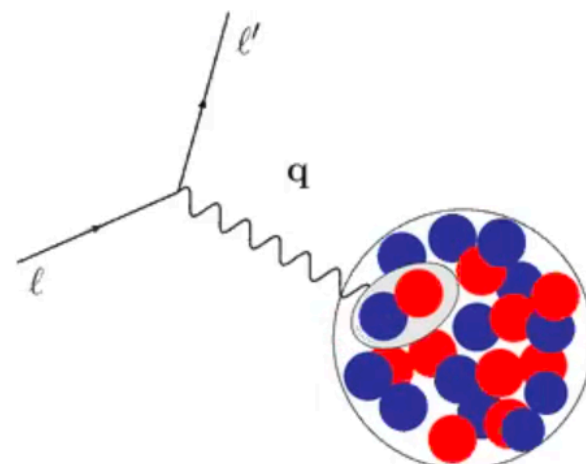
# Electromagnetic interactions

Phenomenological Hamiltonian for NN and NNN

The interaction with external probes is described in terms on one- and two-body charge and current operators



one-body



two-body

Charge operators

$$\rho = \sum_{i=1}^A \rho_i + \sum_{i<j} \rho_{ij} + \dots$$

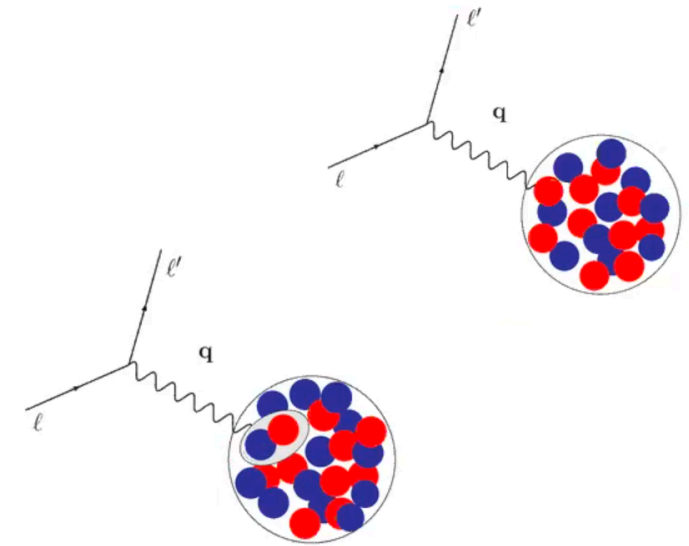
Current operators

$$\mathbf{j} = \sum_{i=1}^A \mathbf{j}_i + \sum_{i<j} \mathbf{j}_{ij} + \dots$$

# Electromagnetic interactions



- One body-currents: non-relativistic reduction of covariant nucleons' isoscalar and isovector currents
- Two-body currents: modeled on MEC currents constrained by commutation relation with the nuclear Hamiltonian (Pastore et al. PRC84(2011)024001, PRC87(2013)014006)
- Argonne v18 two-nucleon and Urbana IX potentials, together with these currents, provide a quantitatively successful description of many nuclear electroweak observables, including charge radii, electromagnetic moments and transition rates, charge and magnetic form factors of nuclei with up to  $A = 12$  nucleons

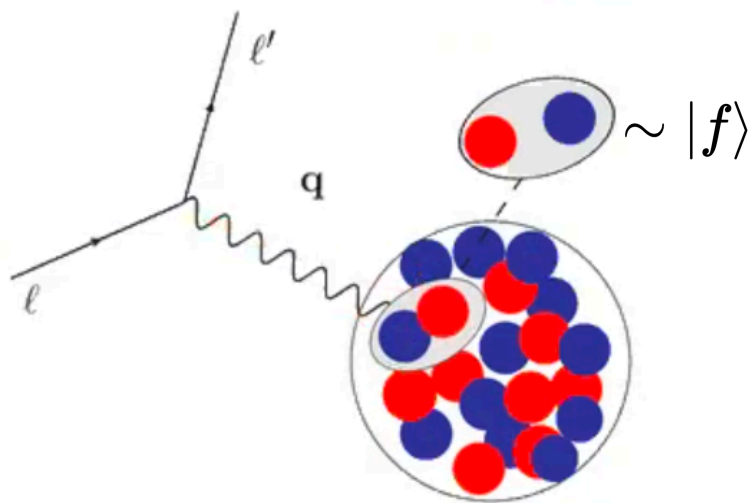




# Short-time approximation

Pastore et al. PRC101(2020)044612

Quasielastic scattering cross sections are expressed in terms of response function



Response functions

$$R_{\alpha}(q, \omega) = \sum_f \delta(\omega + E_0 - E_f) |\langle f | O_{\alpha}(\mathbf{q}) | 0 \rangle|^2$$

$$R_{\alpha}(q, \omega) = \int_{-\infty}^{\infty} \frac{dt}{2\pi} e^{i(\omega + E_i)t} \langle \Psi_i | O_{\alpha}^{\dagger}(\mathbf{q}) e^{-iHt} O_{\alpha}(\mathbf{q}) | \Psi_i \rangle$$

The sum over all final states is replaced by a two nucleon propagator

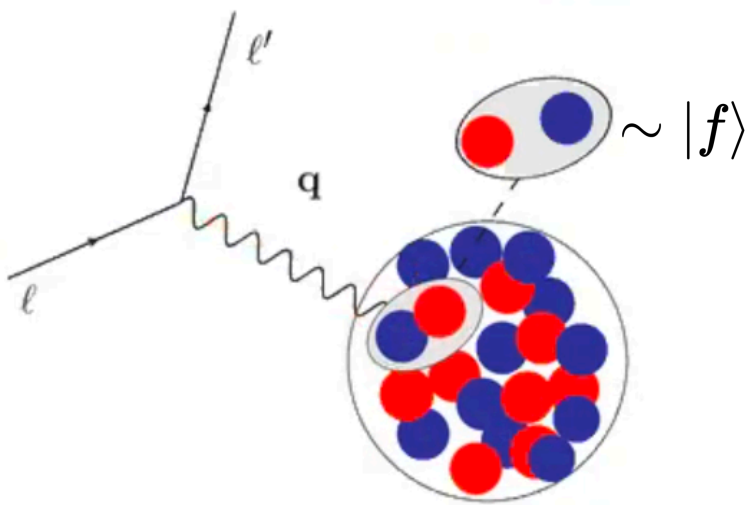
$$\begin{aligned} O^{\dagger} e^{-iHt} O &= \left( \sum_i O_i^{\dagger} + \sum_{i < j} O_{ij}^{\dagger} \right) e^{-iHt} \left( \sum_{i'} O_{i'} + \sum_{i' < j'} O_{i'j'} \right) \\ &= \sum_i O_i^{\dagger} e^{-iHt} O_i + \sum_{i \neq j} O_i^{\dagger} e^{-iHt} O_j \\ &\quad + \sum_{i \neq j} \left( O_i^{\dagger} e^{-iHt} O_{ij} + O_{ij}^{\dagger} e^{-iHt} O_i \right. \\ &\quad \left. + O_{ij}^{\dagger} e^{-iHt} O_{ij} \right) + \dots \end{aligned}$$



# Short-time approximation

Pastore et al. PRC101(2020)044612

Quasielastic scattering cross sections are expressed in terms of response function



Response functions

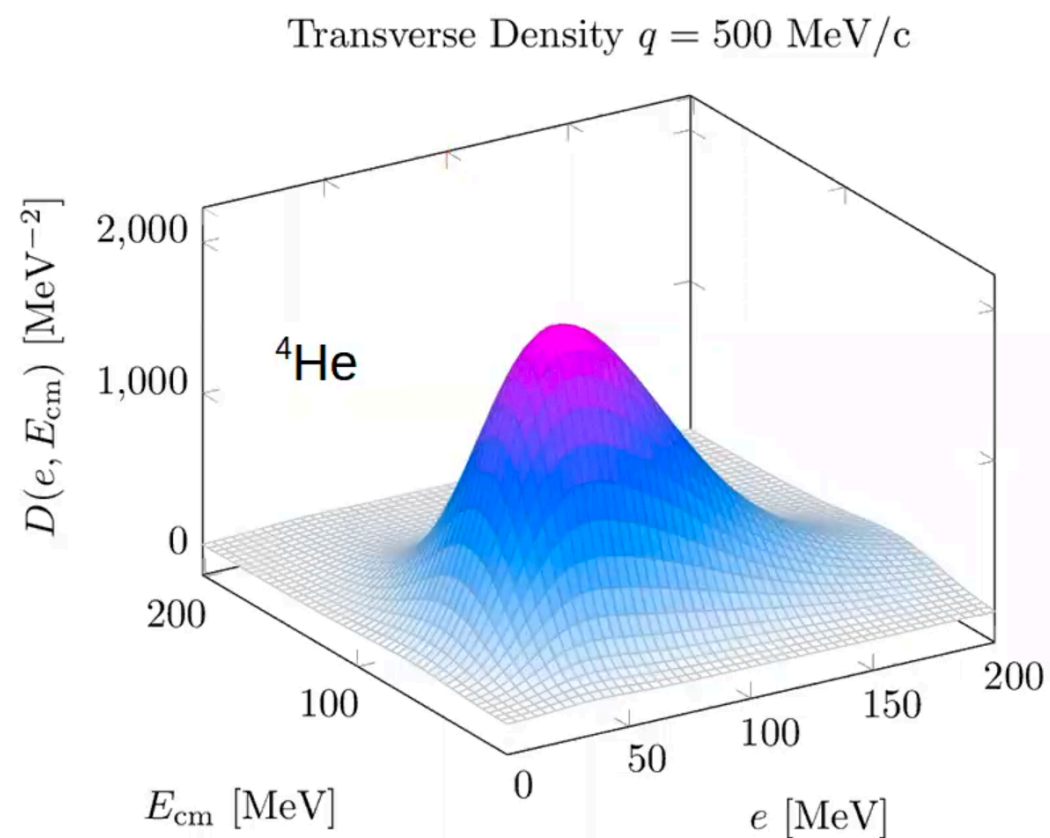
$$R_{\alpha}(q, \omega) = \sum_f \delta(\omega + E_0 - E_f) |\langle f | O_{\alpha}(\mathbf{q}) | 0 \rangle|^2$$

Response densities

$$R^{\text{STA}}(q, \omega) \sim \int \delta(\omega + E_0 - E_f) de dE_{cm} \mathcal{D}(e, E_{cm}; q)$$

STA: scattering of external probes from pairs of correlated nucleons

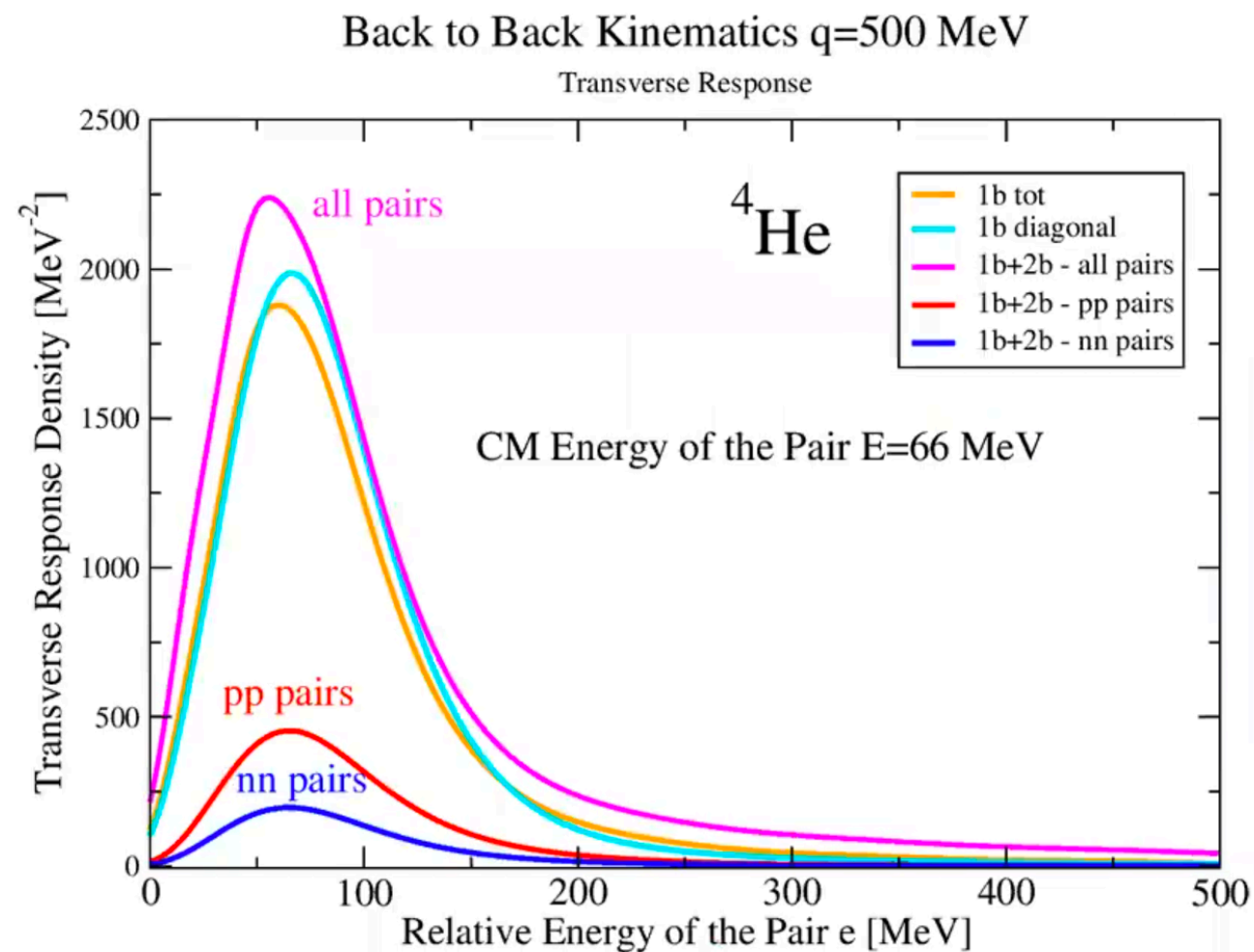
# Transverse response density



Electron scattering from  $^4\text{He}$ :

- Response density as a function of  $(E, e)$
- Give access to particular kinematics for the struck nucleon pair

# Back-to-back kinematic



We can select a particular kinematic, and assess the contributions from different particle identities



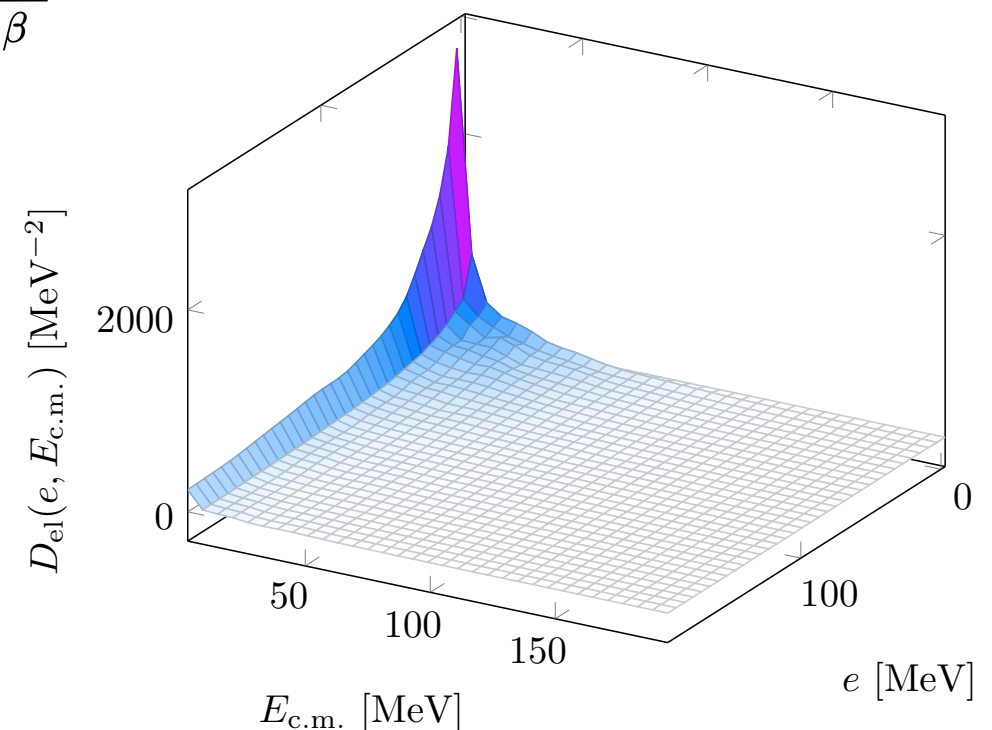
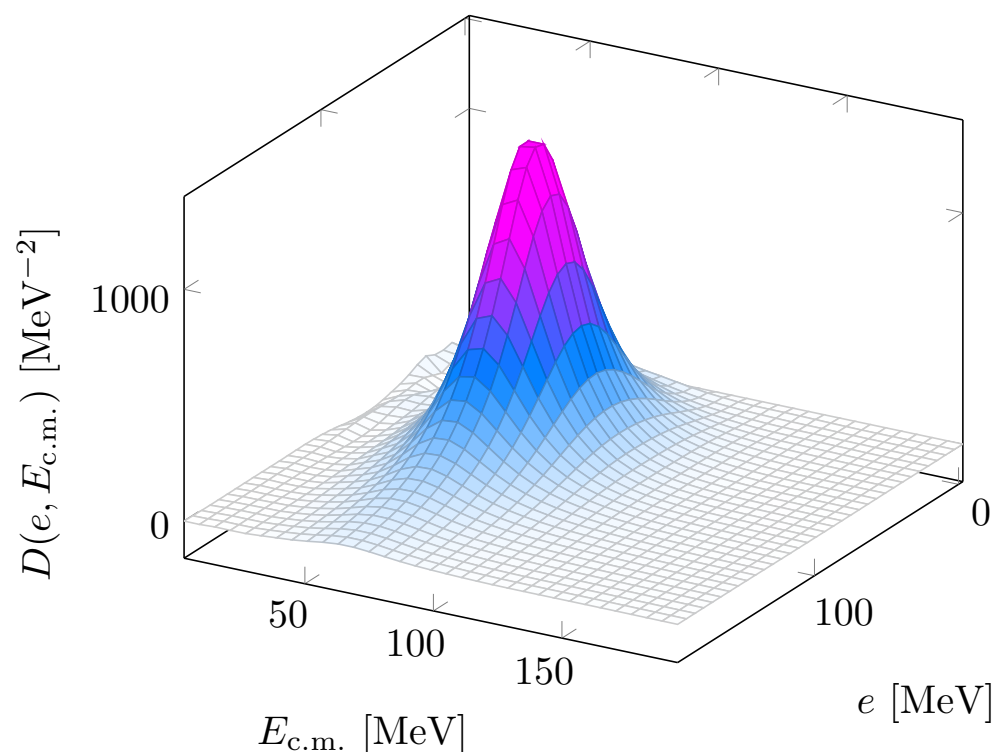


# Longitudinal response density: elastic peak removal

$$R^{\text{STA}}(q, \omega) \sim \int \delta(\omega + E_0 - E_f) de dE_{cm} \mathcal{D}(e, E_{cm}; q)$$

$$\mathcal{D}(e, E_{cm}) - \mathcal{D}_{el}(e, E_{cm})$$

$$\mathcal{D}_{el}(\mathbf{q}, \mathbf{p}', \mathbf{P}') = |\langle \Psi_0 | J(\mathbf{q}) | \Psi_0 \rangle|^2 \\ \times \sum_{\beta} \langle \Psi_0 | \Psi_2(\mathbf{p}', \mathbf{P}', \beta) \rangle \langle \Psi_2(\mathbf{p}', \mathbf{P}', \beta) | \Psi_0 \rangle$$



$^3\text{H}$  Longitudinal response at 300 MeV

# Benchmark



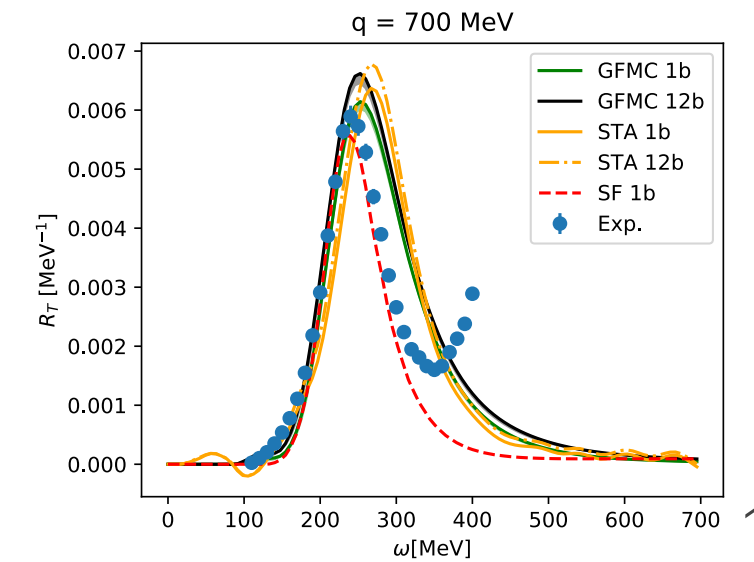
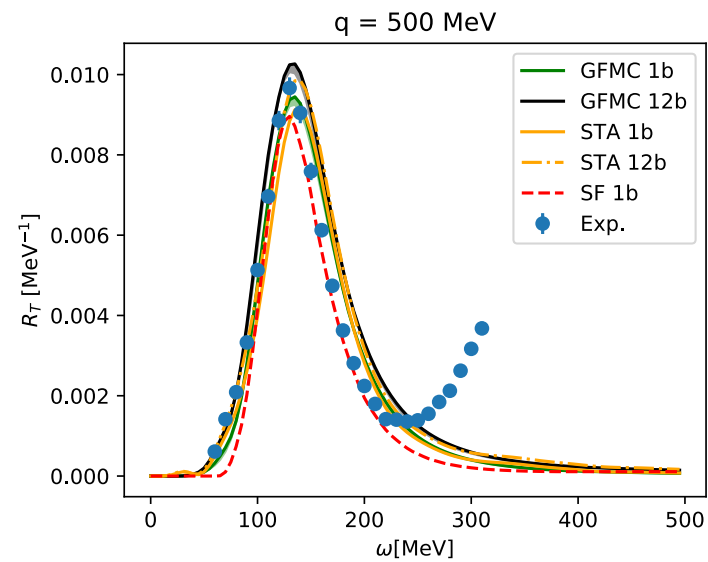
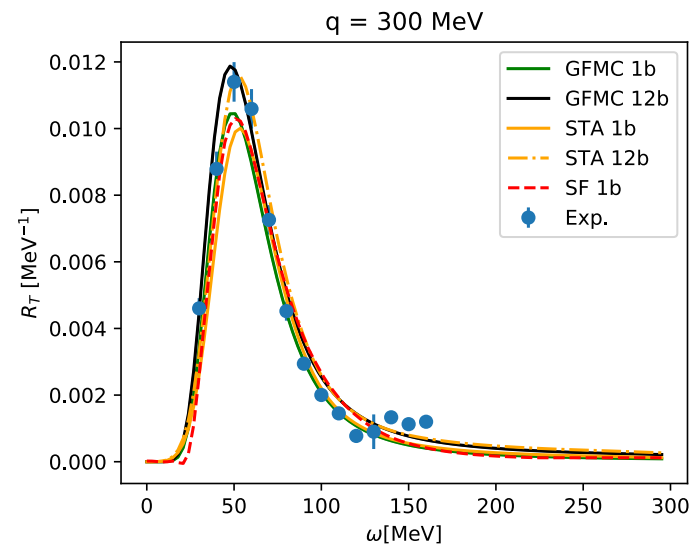
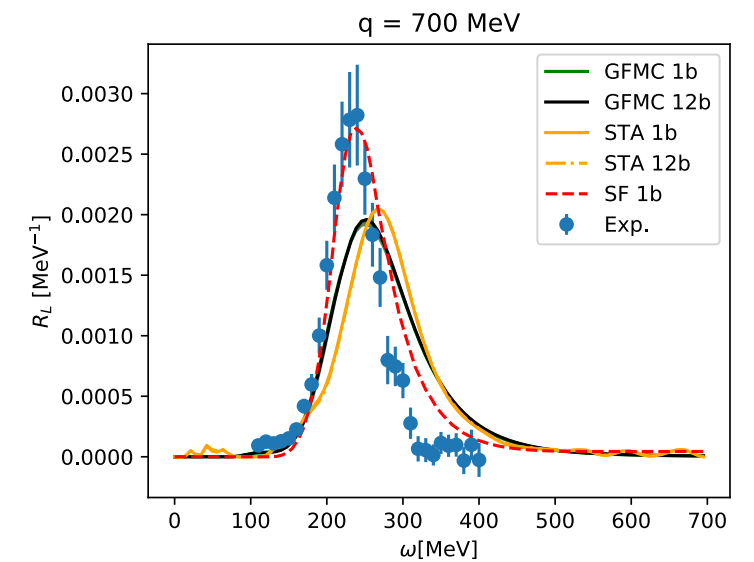
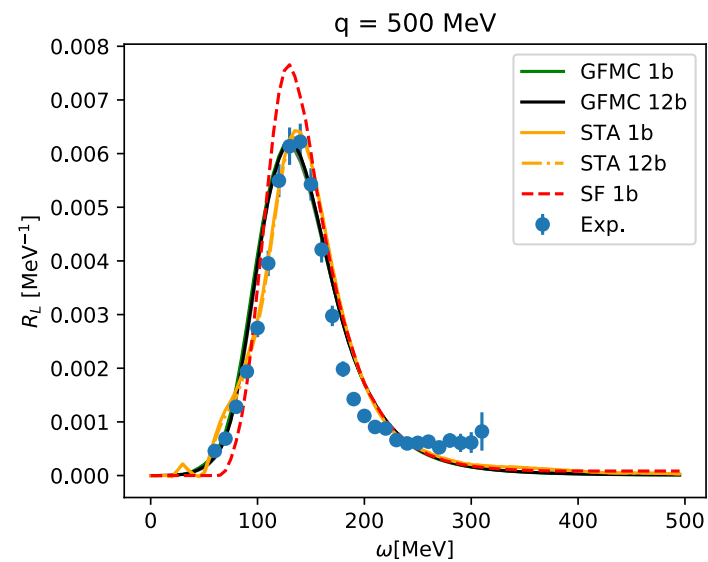
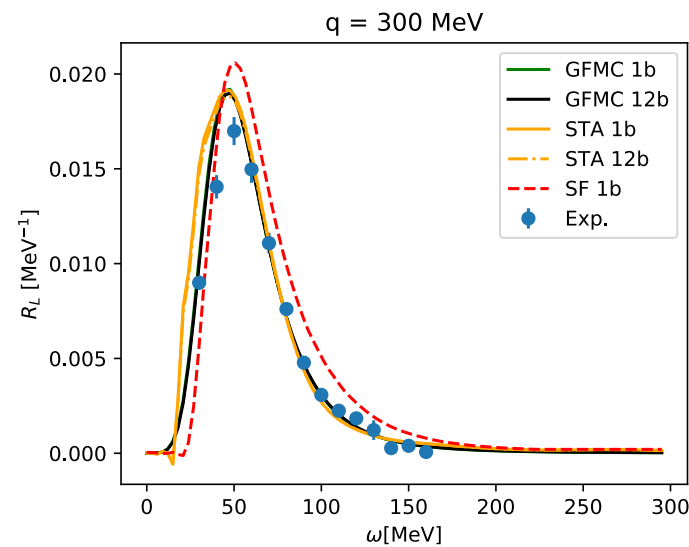
L.A., S. Pastore, N. Rocco, et al. PRC105(2022)014002

- We benchmarked three different methods based on the same description of nuclear dynamics of the initial target state
- Compared to the experimental data for the longitudinal and transverse electromagnetic response functions of  $^3\text{He}$ , and the inclusive cross sections of both  $^3\text{He}$  and  $^3\text{H}$
- Comparing the results allows for a precise quantification of the uncertainties inherent to factorization schemes

# Benchmark



Longitudinal and transverse response function in  $^3\text{He}$

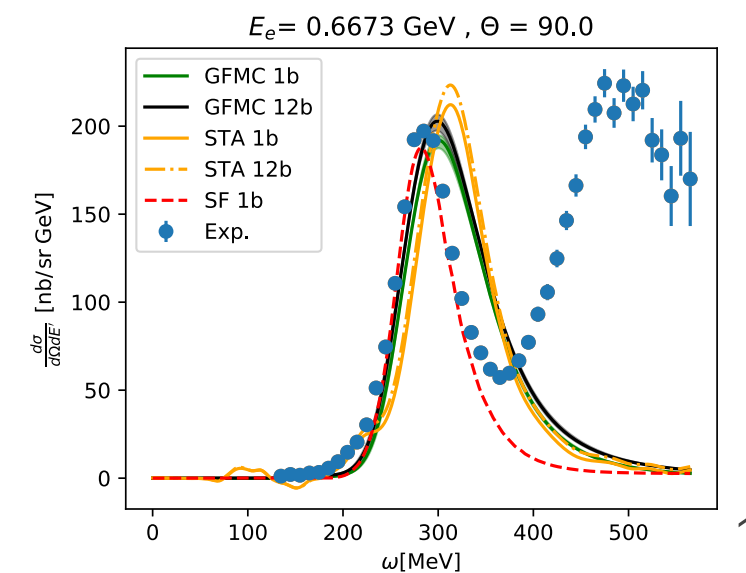
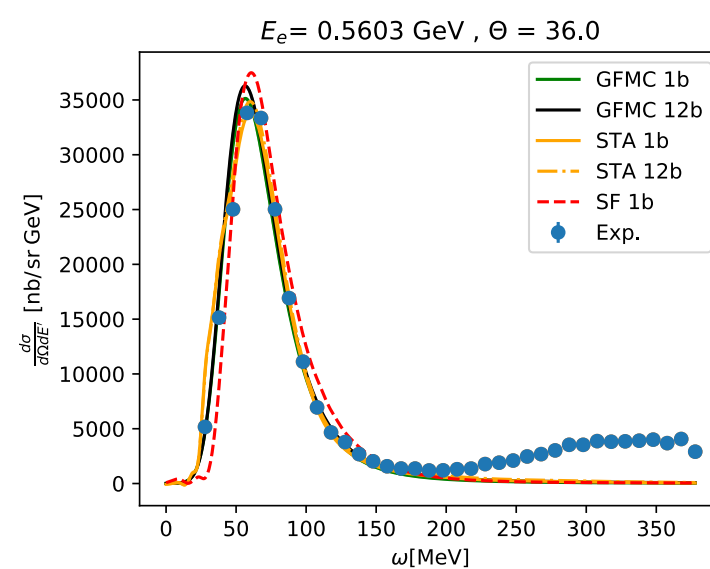
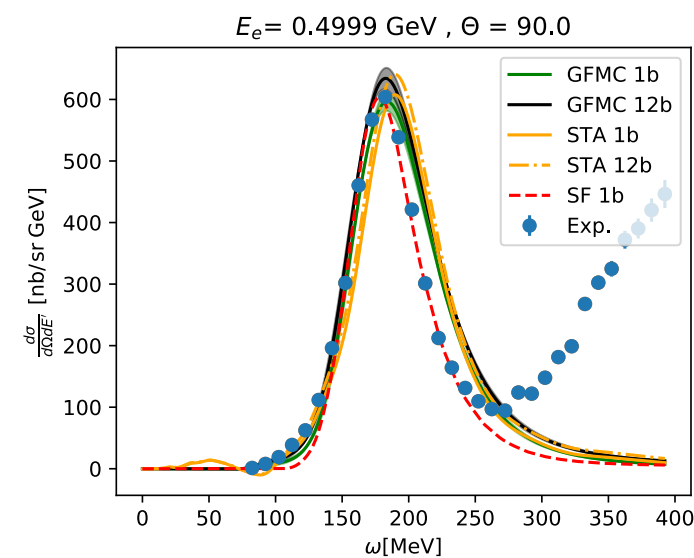
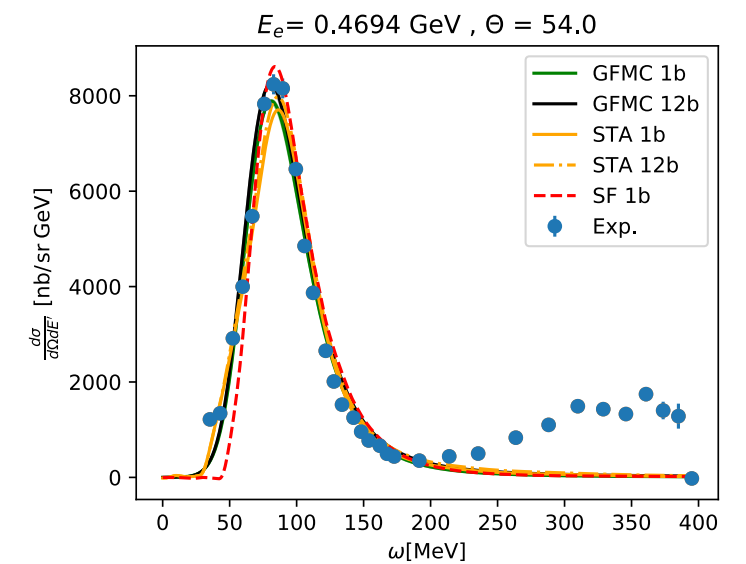
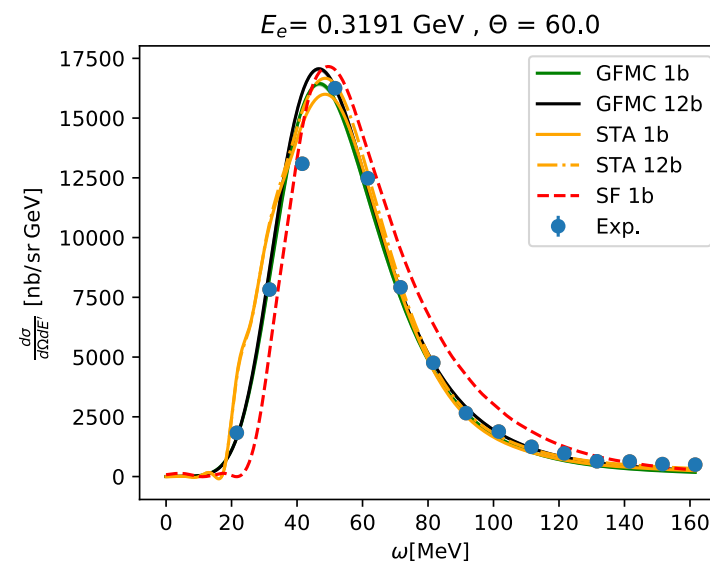
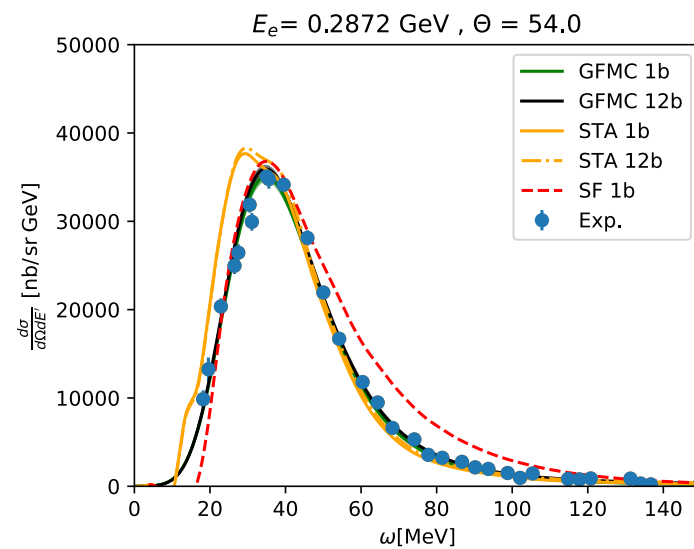


L.A., S. Pastore, N. Rocco, et al. PRC105(2022)014002

# Cross sections



$^3\text{He}$

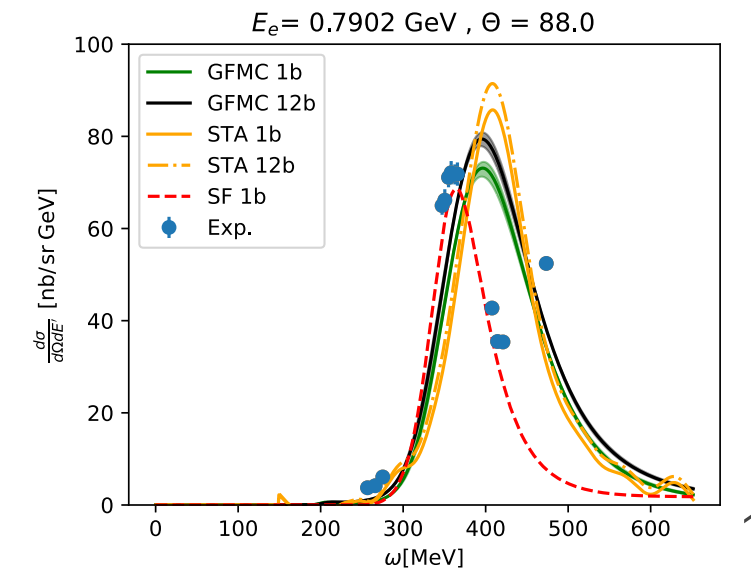
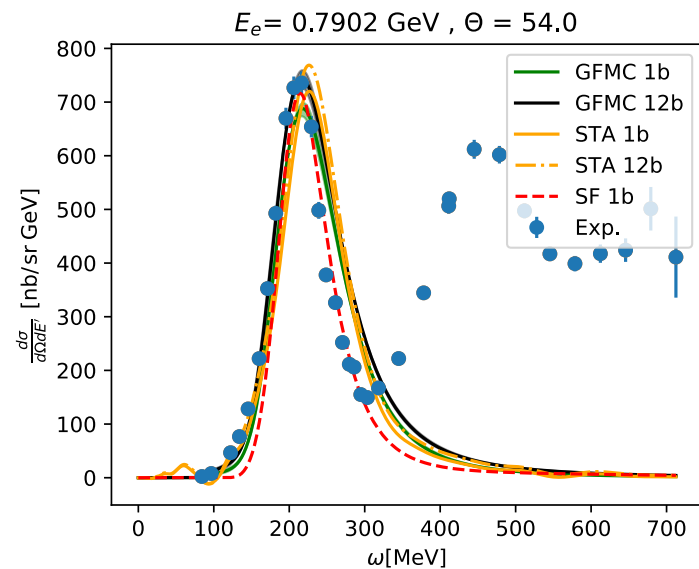
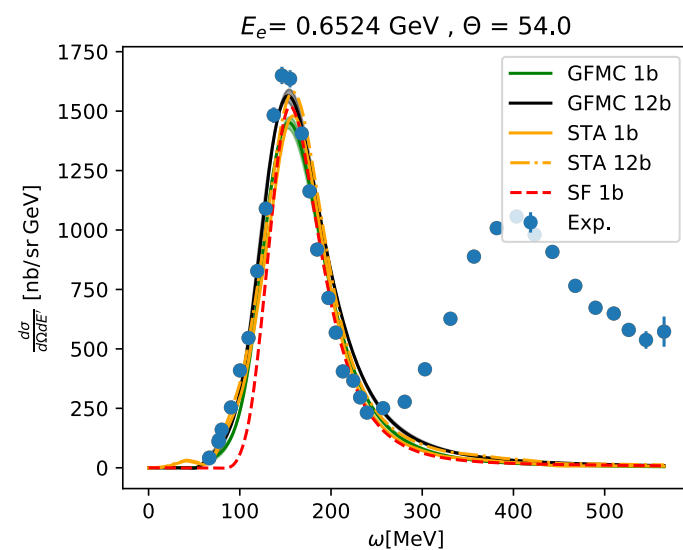
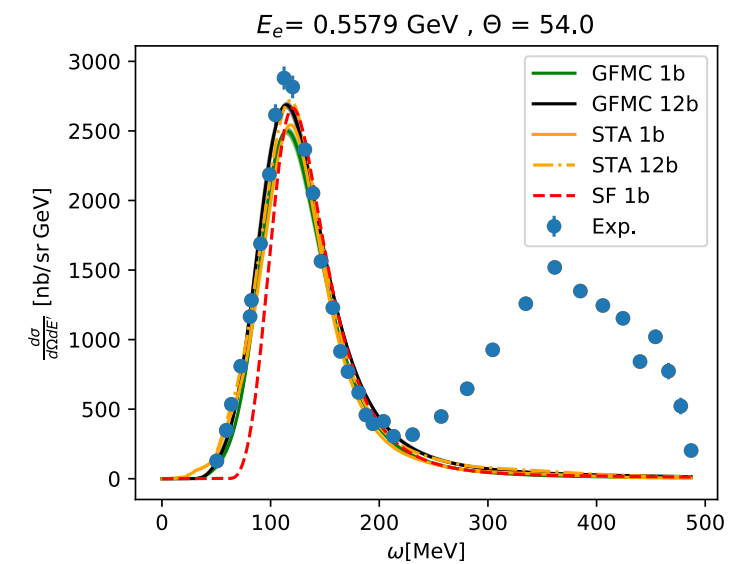
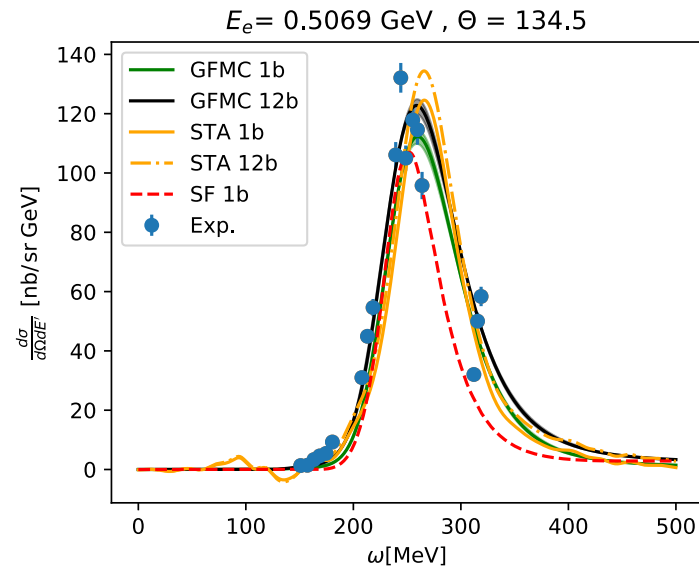
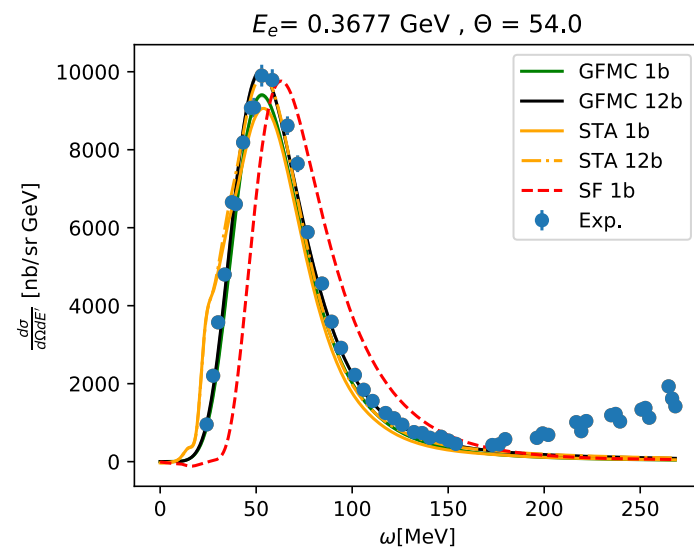


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# Cross sections



$^3\text{H}$

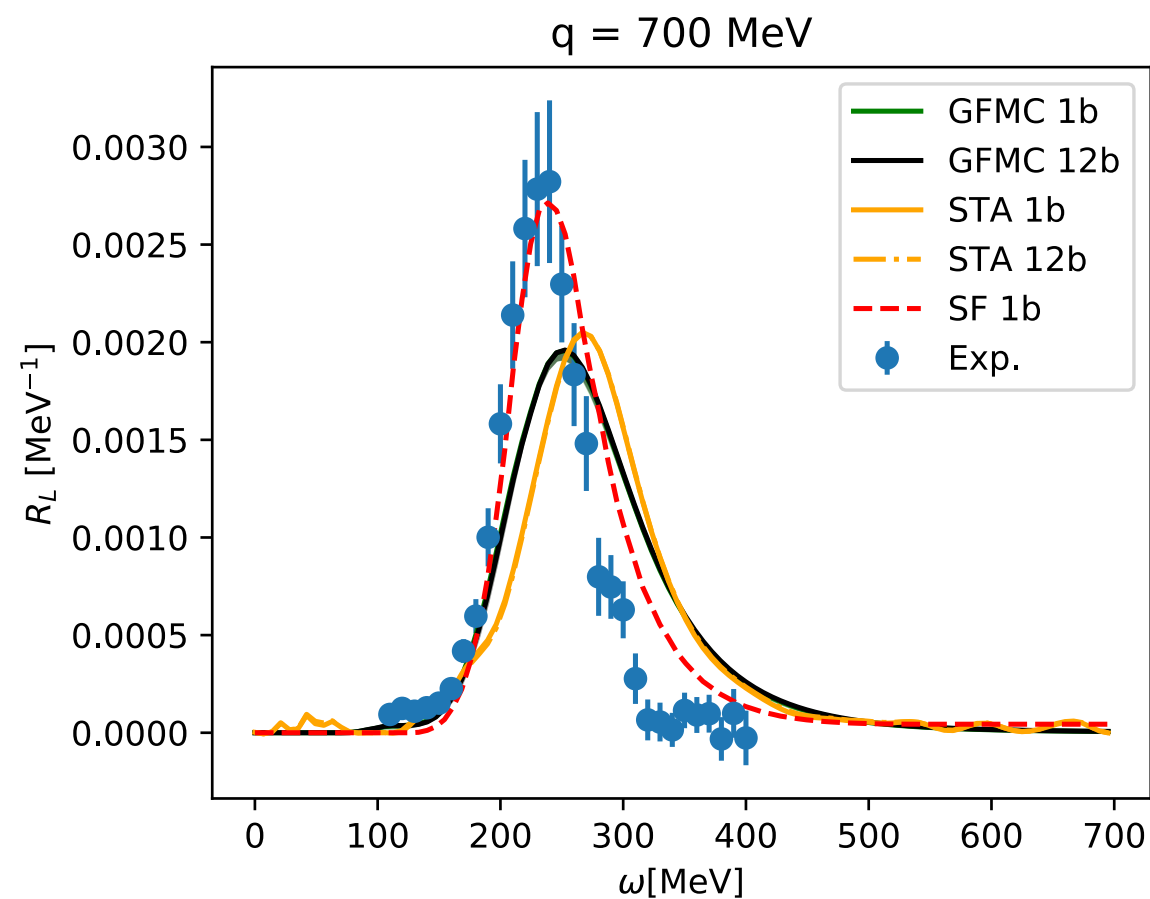


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# Relativistic corrections



Necessary to include relativistic correction at higher momentum  $q$



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# Relativistic corrections



We are currently working on including relativistic corrections within the STA formalism:

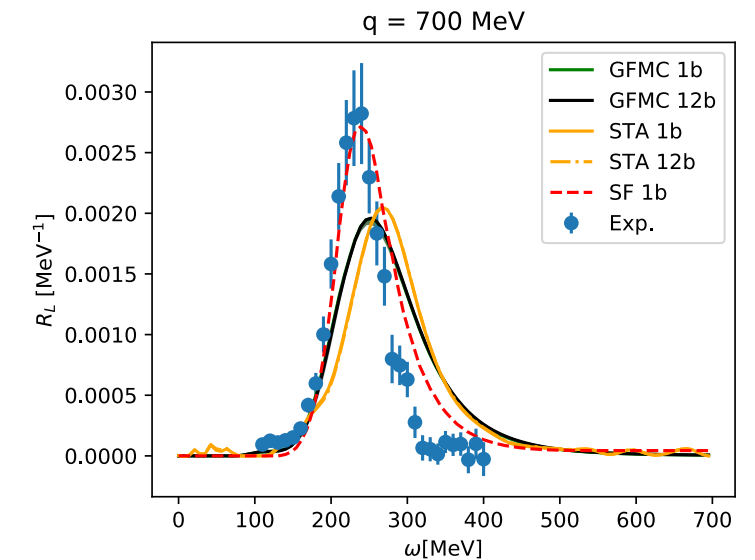
R. Weiss, S. Pastore, J. Carlson

- Relativistic kinematic
- Relativistic currents

# Relativistic corrections



- Relativistic kinematic
- Relativistic currents



$$E_{NR} \rightarrow E(\mathbf{p}, \mathbf{P}) = \sqrt{\left(\frac{\mathbf{P}}{2} - \mathbf{p}\right)^2 + m^2} + \sqrt{\left(\frac{\mathbf{P}}{2} + \mathbf{p}\right)^2 + m^2} - 2m = \sqrt{\frac{P^2}{4} - Pp \cos(\theta) + p^2 + m^2} + \sqrt{\frac{P^2}{4} + Pp \cos(\theta) + p^2 + m^2} - 2m$$

$$R_\alpha(\mathbf{q}, \omega) =$$

...

...

...

$$= \int \frac{d^3 K}{(2\pi)^3} \int \frac{d^3 k}{(2\pi)^3} \left[ \sum_{M_i} \sum_{\alpha_{A-2}, \alpha} \int d^3 R_{A-2} \langle \Psi_i | O_L^\dagger(\mathbf{q}) | \mathbf{K} \mathbf{k} \alpha, \mathbf{R}_{A-2}, \alpha_{A-2} \rangle \langle \mathbf{K} \mathbf{k} \alpha, \mathbf{R}_{A-2}, \alpha_{A-2} | O_R(\mathbf{q}) | \Psi_i \rangle \right] \delta(\omega + E_i - E(\mathbf{k}, \mathbf{K}))$$



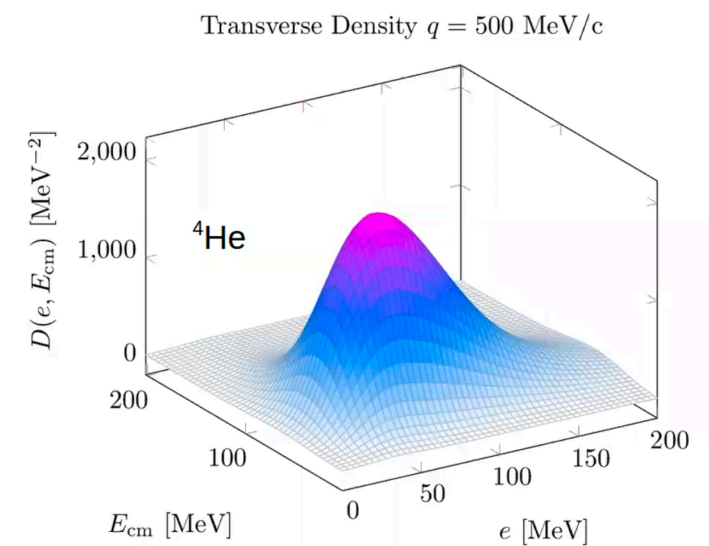
# Heavier nuclei



Computational complexity of response functions and densities:

Wave-function		${}^4\text{He}$	${}^{12}\text{C}$
Spin	$2^A$	16	4096
Isospin	$\frac{A!}{Z!(A-Z)!}$	6	924
Pairs	$A(A-1)/2$	6	66

Response densities: E, e grid



$$R_{\alpha}(q, \omega) = \int_{-\infty}^{\infty} \frac{dt}{2\pi} e^{i(\omega + E_i)t} \langle \Psi_i | O_{\alpha}^{\dagger}(\mathbf{q}) e^{-iHt} O_{\alpha}(\mathbf{q}) | \Psi_i \rangle$$

# Heavier nuclei

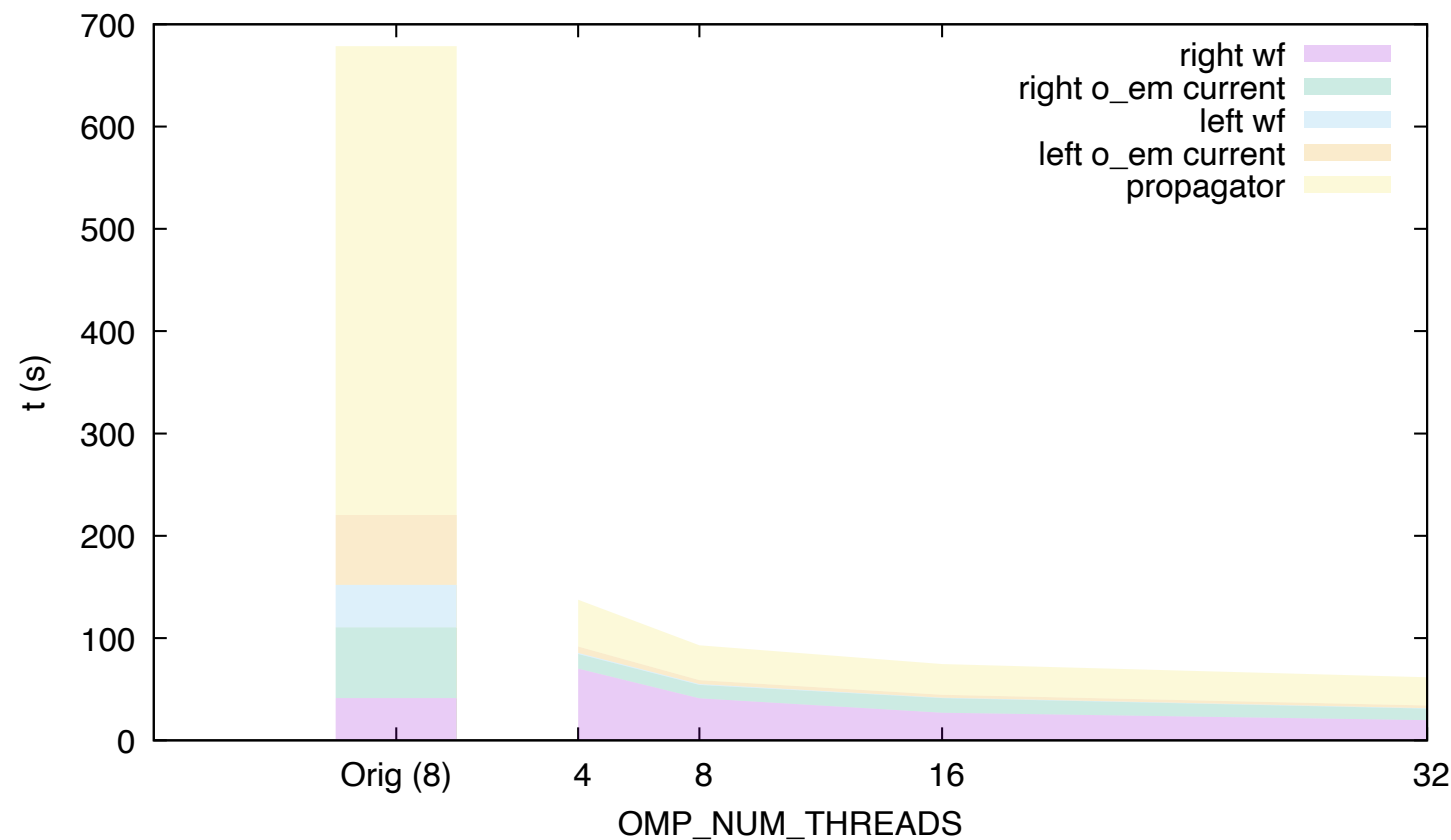


Optimization was necessary to tackle heavier nuclei

$$R_{\alpha}(q, \omega) = \int_{-\infty}^{\infty} \frac{dt}{2\pi} e^{i(\omega + E_i)t} \langle \Psi_i | O_{\alpha}^{\dagger}(\mathbf{q}) e^{-iHt} O_{\alpha}(\mathbf{q}) | \Psi_i \rangle$$

- Parallelization – MPI and OpenMP:  
Variational Monte Carlo is almost perfectly parallelizable, but with increased system size memory becomes a constrain
- Refactoring of the code
- Computational algorithms and approximations:  
variation of integration ranges ( $r$ ,  $R$ ) for struck nucleon pair

# Heavier nuclei: $^{12}\text{C}$



Optimization specific to  $^{12}\text{C}$  was needed in order to perform full response densities calculations:

- **parallelization**
- **refactoring of the code**
- **reduction of memory usage**
- computational algorithms and approximations

$$R_{\alpha}(q, \omega) = \int_{-\infty}^{\infty} \frac{dt}{2\pi} e^{i(\omega + E_i)t} \langle \Psi_i | O_{\alpha}^{\dagger}(\mathbf{q}) e^{-iHt} O_{\alpha}(\mathbf{q}) | \Psi_i \rangle$$

# Heavier nuclei



Comparison for  ${}^3H$ , 72k MC configurations, 40x40 points in  $r$ ,  $R$  integration

Original:

- ~15k core hours  
(LA et al arXiv:2108.10824)

After optimization:

- ~1.8k core hours

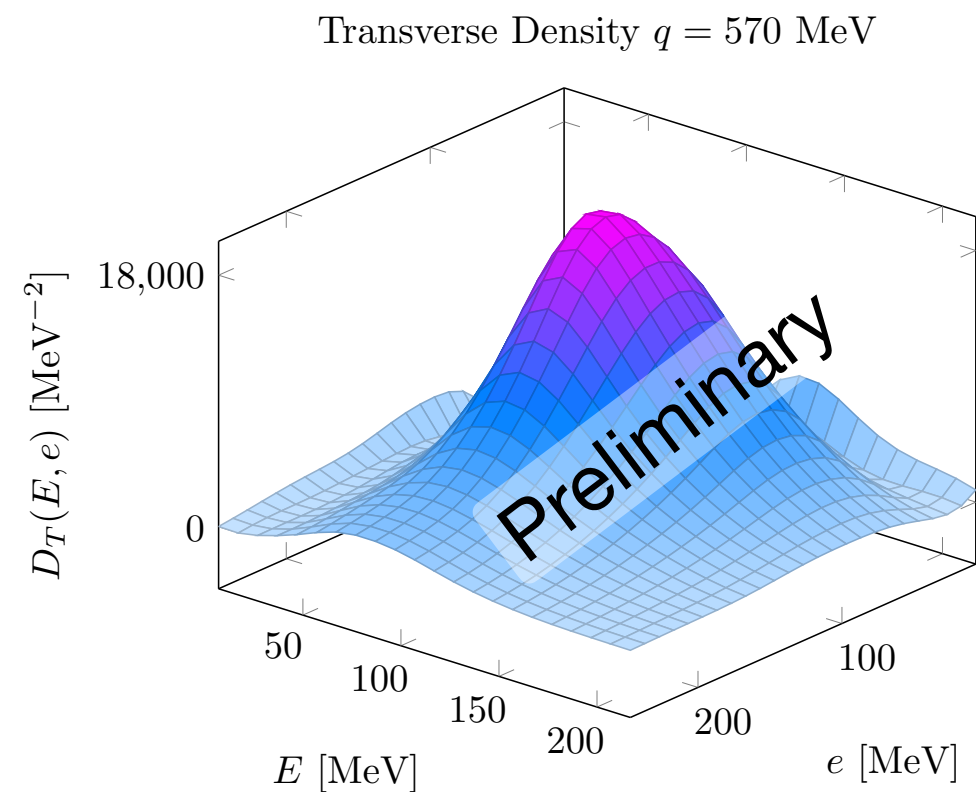
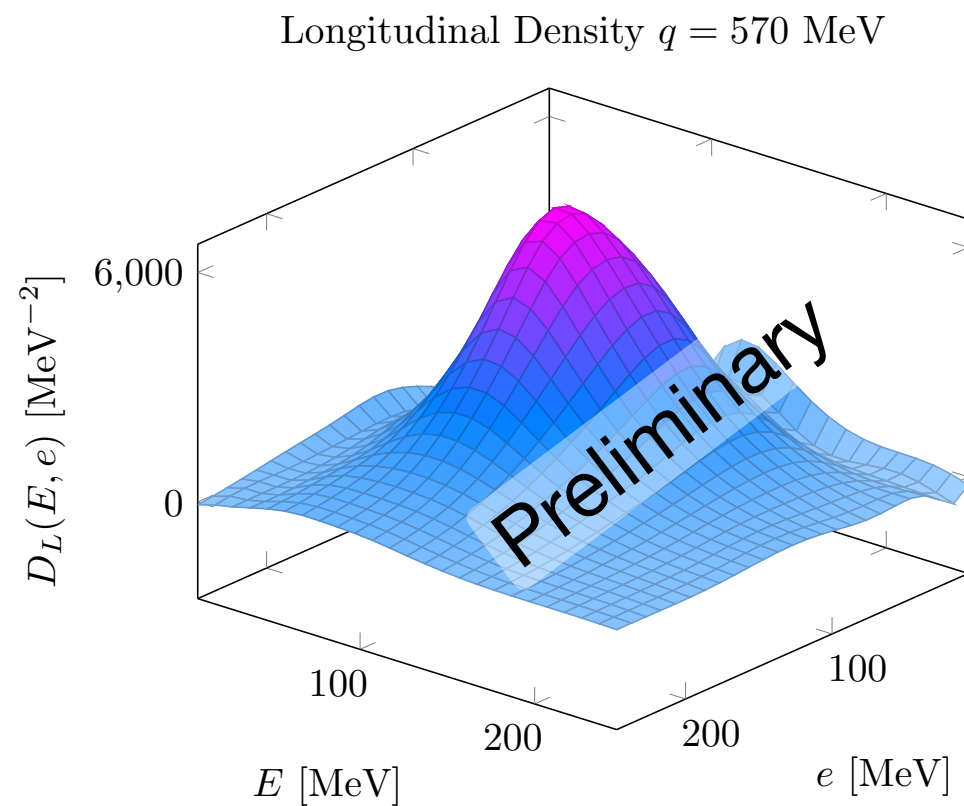
Additional approximations reduce the computation time:

$r$ ,  $R$  integration up to 8, 5 fm

reduction of grid spacing in relative energy  $e$



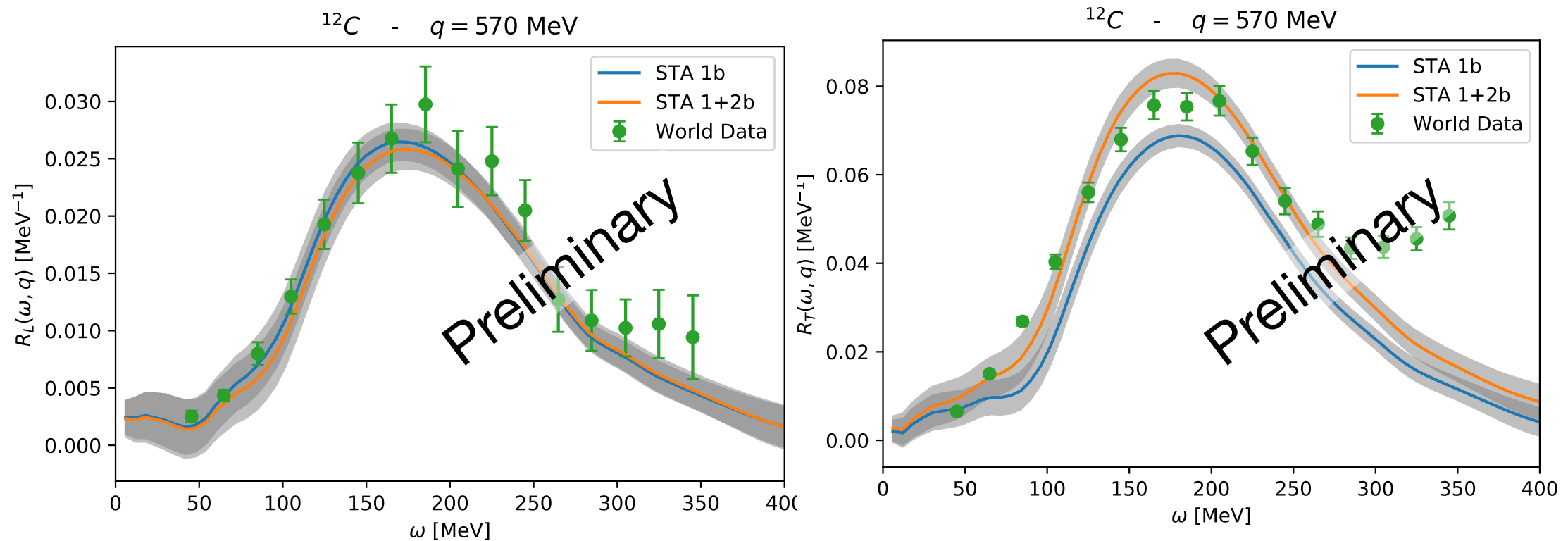
# Response densities for $^{12}\text{C}$



Longitudinal and transverse response densities in  $^{12}\text{C}$

$q = 570 \text{ MeV}$

# Response functions for $^{12}\text{C}$



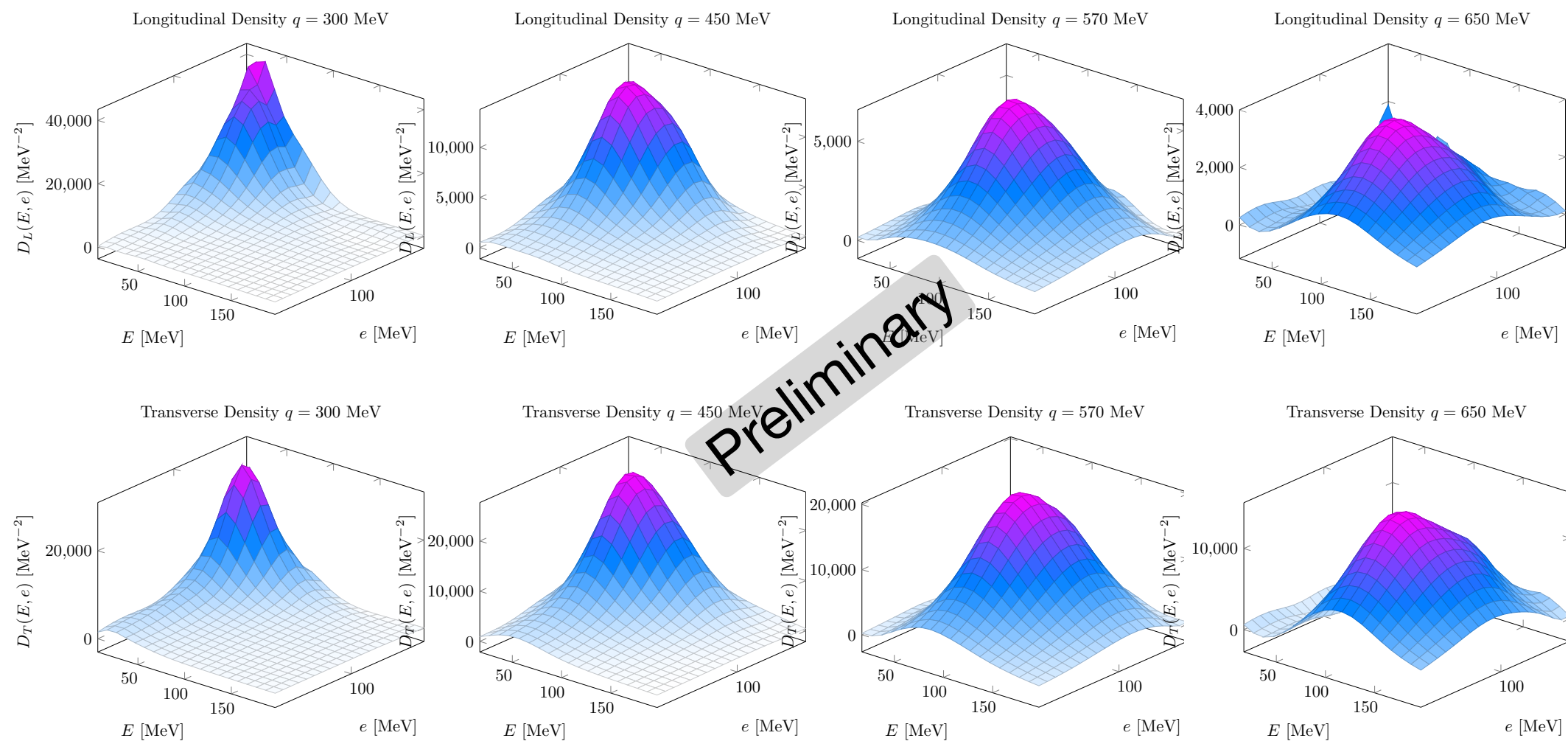
Preliminary results for longitudinal and transverse response functions in  $^{12}\text{C}$

$q = 570$  MeV

# Responses for $^{12}\text{C}$



Longitudinal and transverse response for  $300 < q < 650 \text{ MeV}$ :





# Cross sections: Interpolation schemes



- Cross sections weakly dependent on interpolation scheme in  $4\text{He}$ , but relevant in  $^{12}\text{C}$
- We tested various interpolation schemes on  $4\text{He}$ , where we can evaluate responses for an arbitrary fine grid of values of  $q$ .



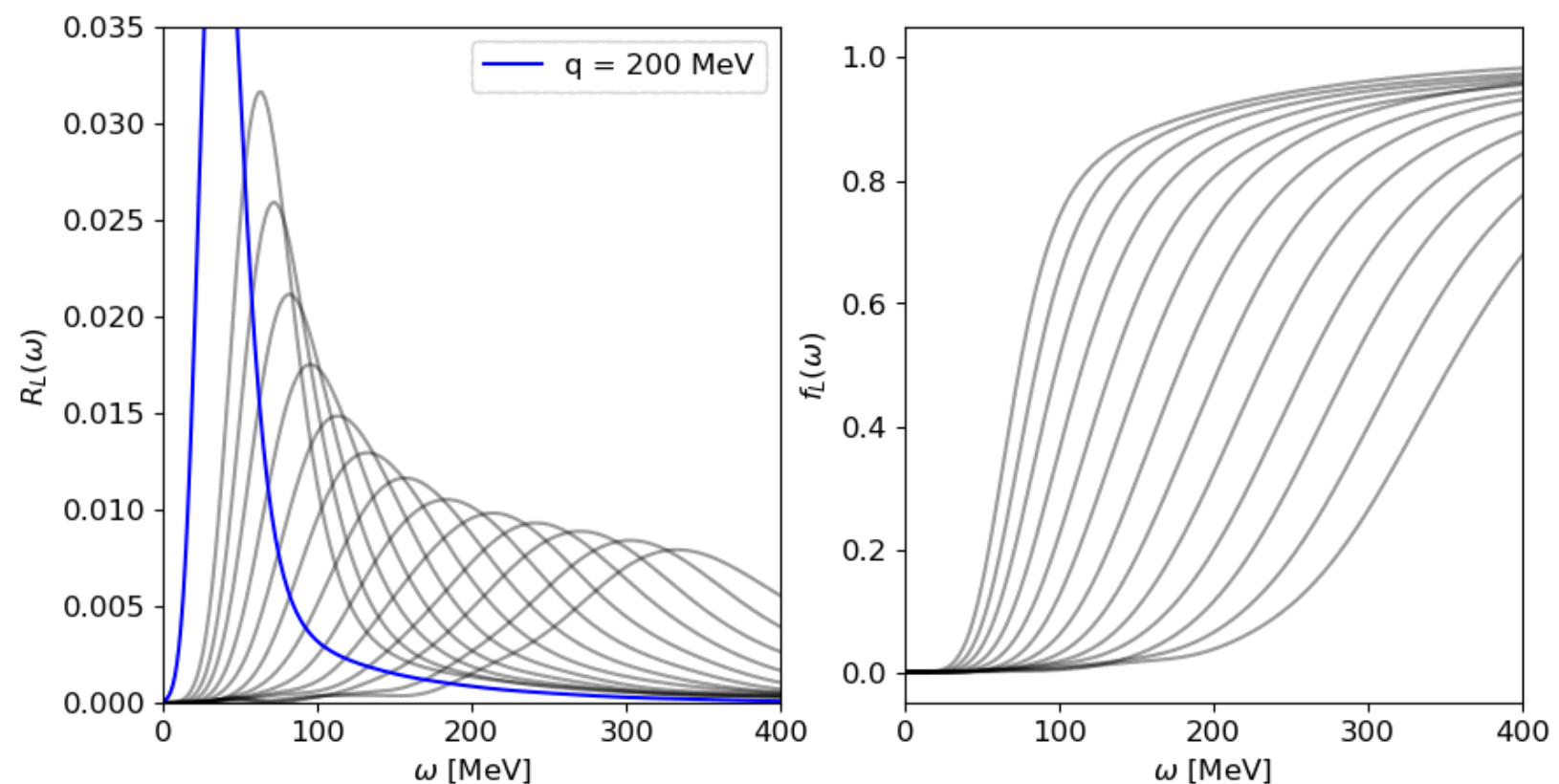
# Cross sections: Interpolation schemes



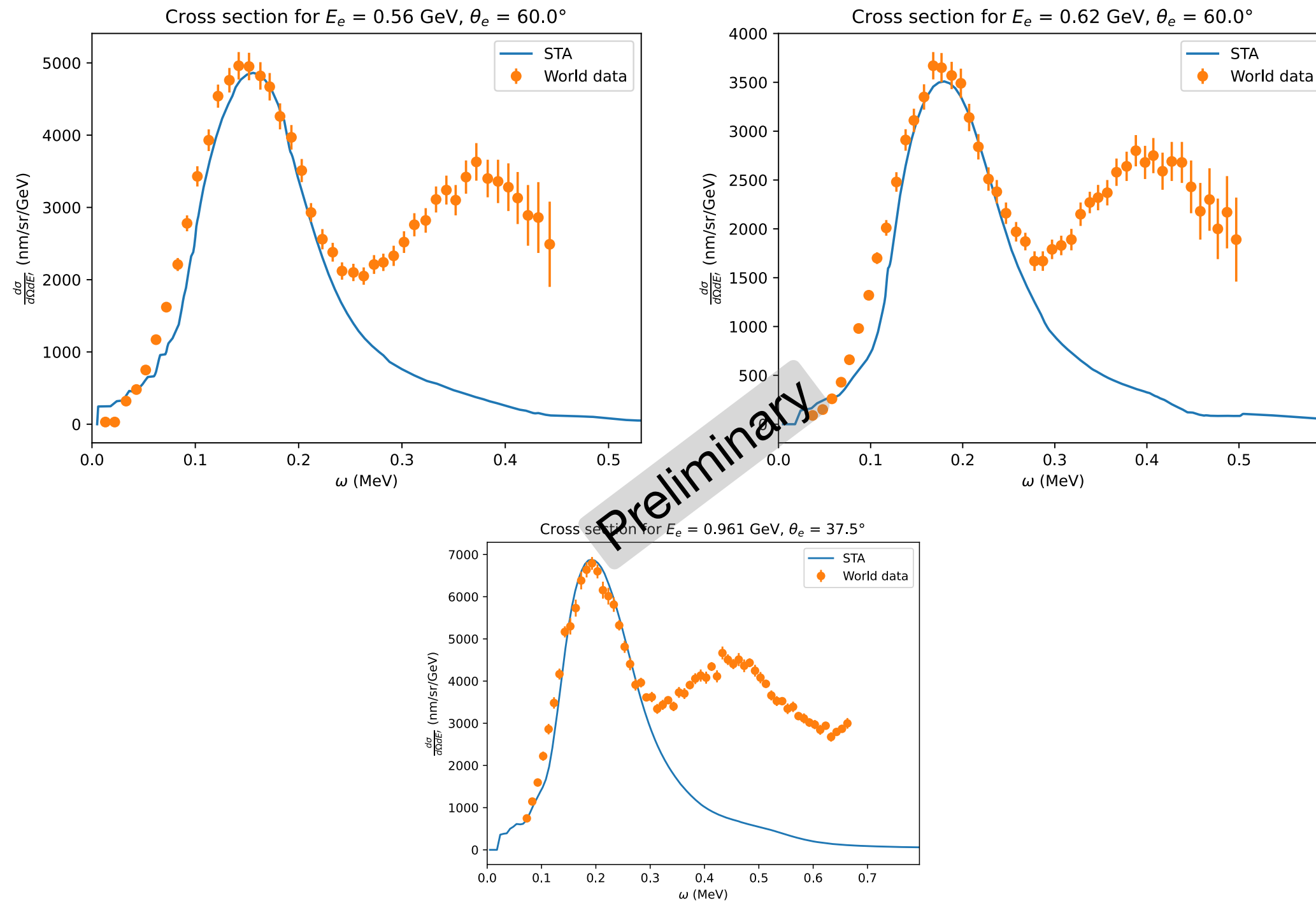
- We **interpolate** in between cumulative integrals of responses, using information from the sum rules
- Outside the range ( $q < 300$  MeV and  $q > 650$  MeV), we use **scaling functions**

$$\psi'_{\text{nr}} = \frac{m_N}{|\mathbf{q}|k_F} \left( \omega - \frac{|\mathbf{q}|^2}{2m_N} - \varepsilon \right)$$

4He, longitudinal response



# Cross sections results



# Conclusion:



- The STA responses for  $^{12}\text{C}$  are in good agreement with the data
- Given the computational complexity of evaluating cross sections, a novel interpolation scheme was adopted for the calculation of cross sections

## **EW interactions:**

- The current work on EM interactions allows for a thorough evaluation of the method, and a comparison with the abundant experimental data for electron-nucleus scattering
- Use of information from response densities in event generators.

# Collaborators:

S. Pastore, M. Piarulli

# Thank you!

Quantum Monte Carlo Group @ WashU

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Lorenzo Andreoli (PD)

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