

# Diagrammatic resummations for the in-medium similarity renormalization group

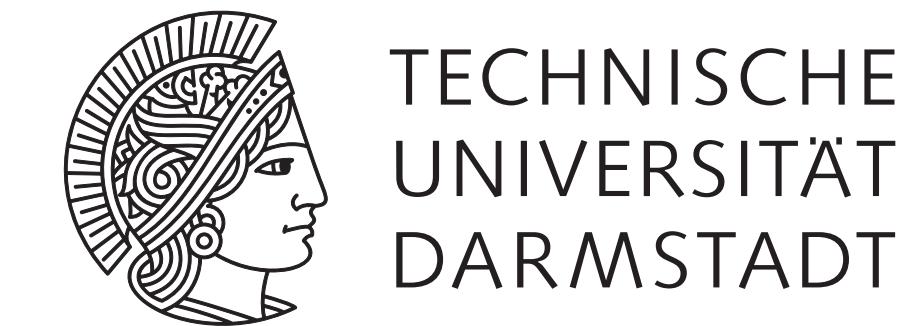


**Matthias Heinz**

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ESNT Workshop "Automated tools for many-body theory"  
June 8, 2023



# Outline

- Introduction to the IMSRG

Tsukiyama et al., PRL **106** (2011)

Hergert et al., Phys. Rep. **621** (2016)

**MH et al., PRC **103** (2021)**

- Connecting IMSRG and MBPT

Hjorth-Jensen et al., Phys. Rep. **261** (1995)

Bartlett, Shavitt, *Many-Body Methods in Chemistry and Physics* (2009)

Hergert et al., Phys. Rep. **621** (2016)

- Approaches to improving IMSRG truncations

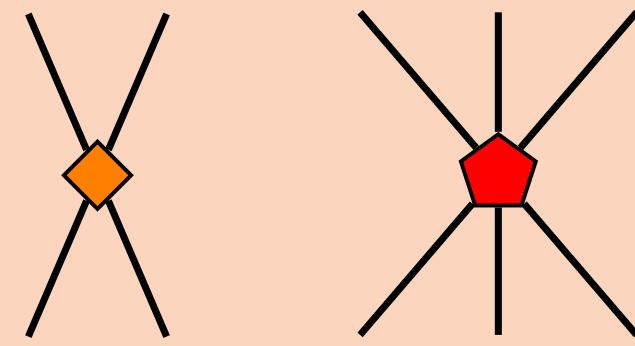
Morris, PhD Thesis, MSU (2016)

Arthuis et al., CPC **240** (2019)

# Introduction to the IMSRG

Tsukiyama et al., PRL **106** (2011)  
Hergert et al., Phys. Rep. **621** (2016)  
**MH et al.**, PRC **103** (2021)

# The basics



$$H = T_{\text{int}} + V_{\text{NN}} + V_{\text{3N}}$$

$$H|\Psi\rangle = E|\Psi\rangle$$

**Many-body solver**

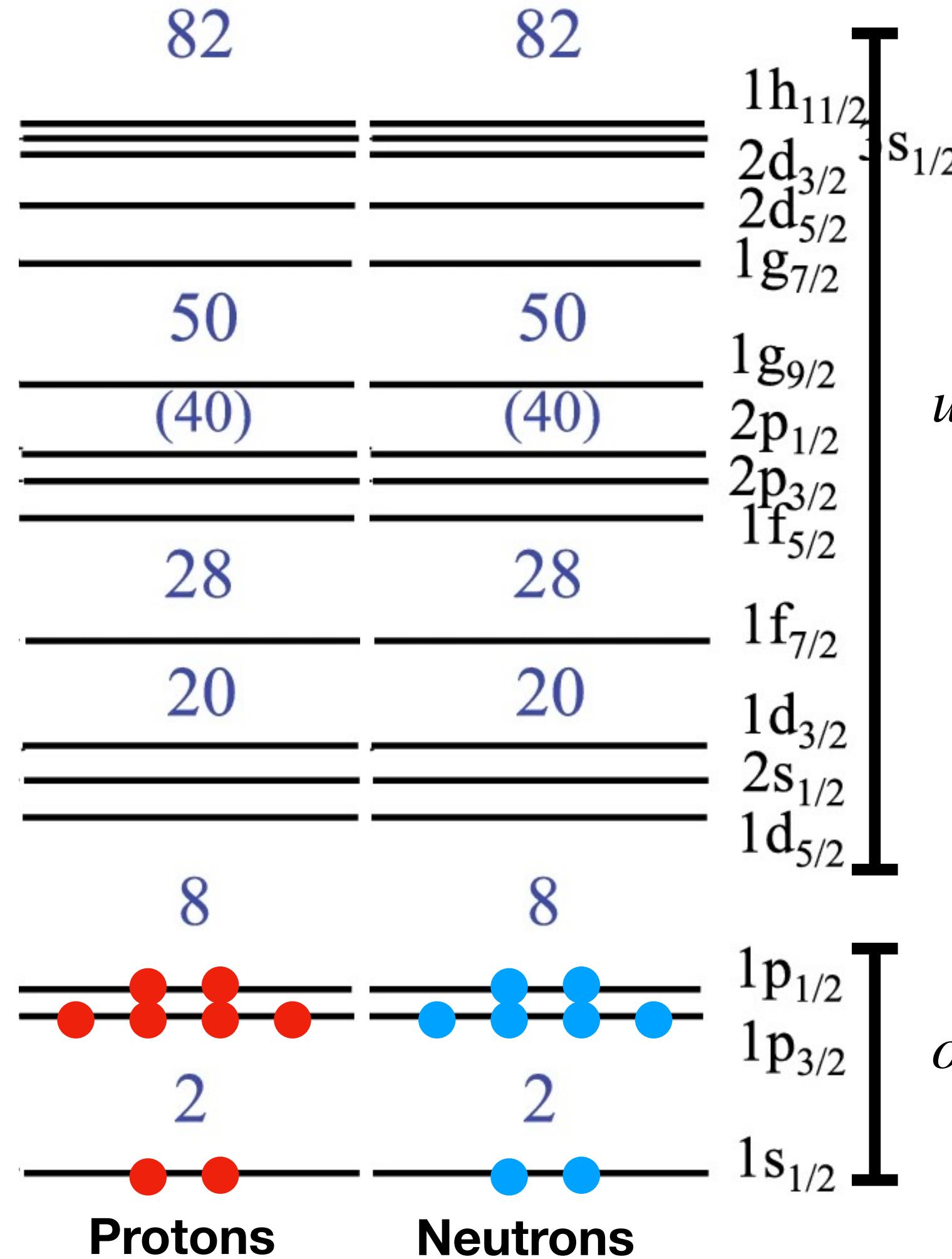
**IMSRG**

$E$

$$|\Psi\rangle = U|\Phi\rangle = e^{\Omega}|\Phi\rangle$$

$$\langle R_p^2 \rangle, \dots \quad \overline{H} \rightarrow \text{spectra}$$

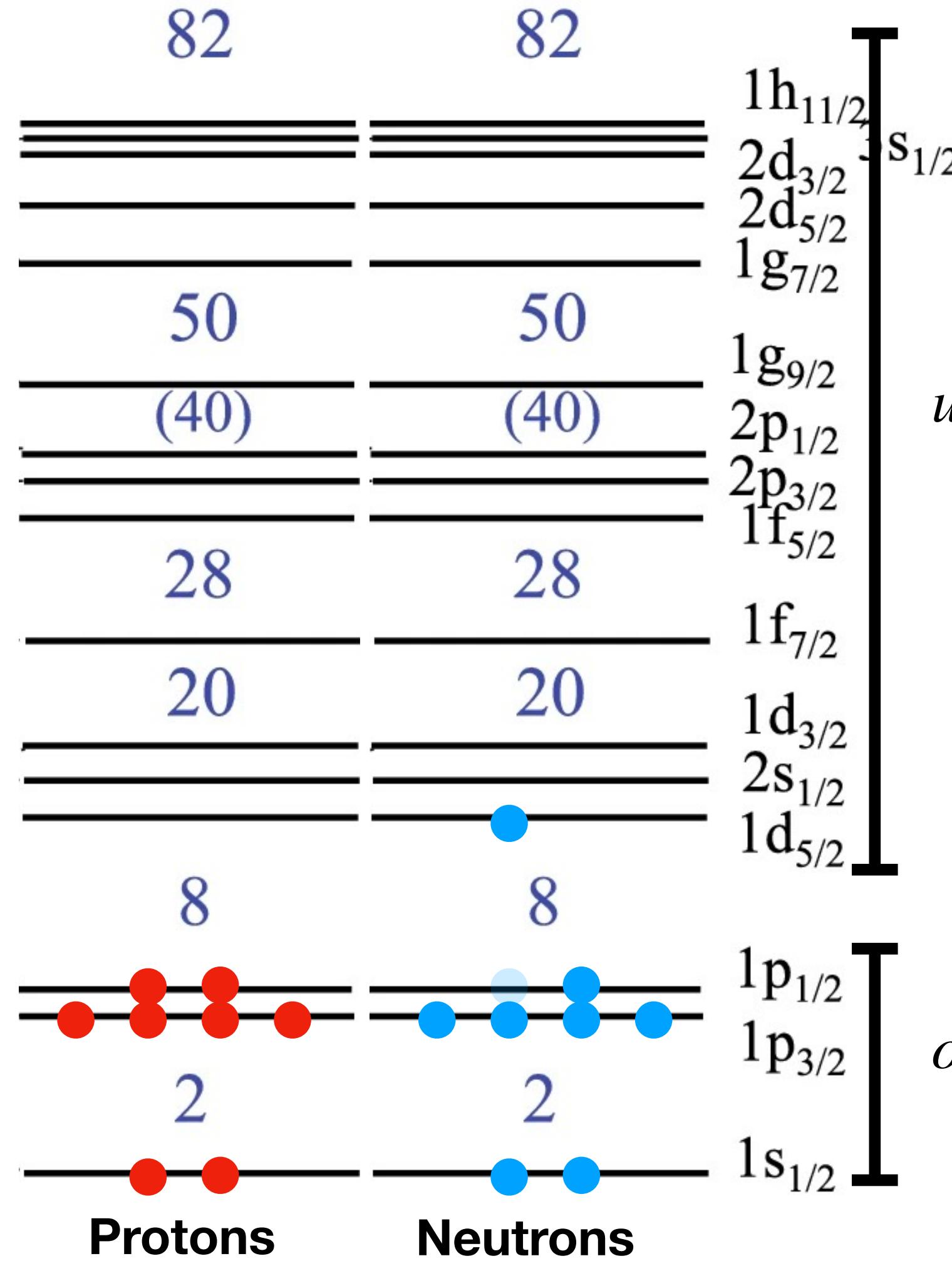
# Solving the many-body problem



$$|\Phi\rangle \sim \prod_{i=1}^o |\phi_i\rangle \quad \text{mean-field reference state}$$

Hagino et al., Found. Chem. 22 (2020)

# Solving the many-body problem

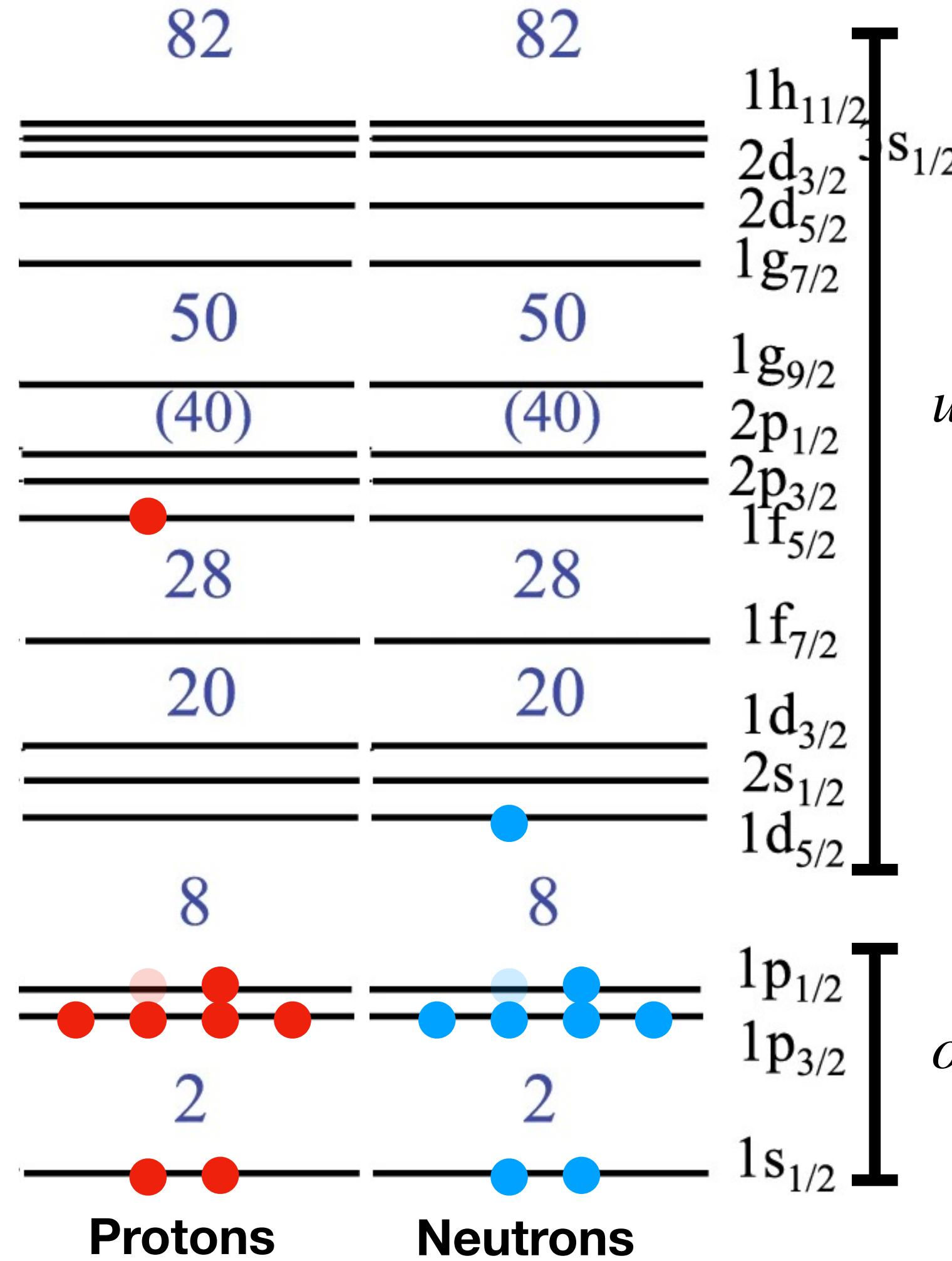


$|\Phi\rangle \sim \prod_{i=1}^o |\phi_i\rangle$  mean-field reference state

$|\Phi_i^a\rangle$  1p1h excitation  $(\binom{o}{1} \binom{u}{1})$  states)

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# Solving the many-body problem



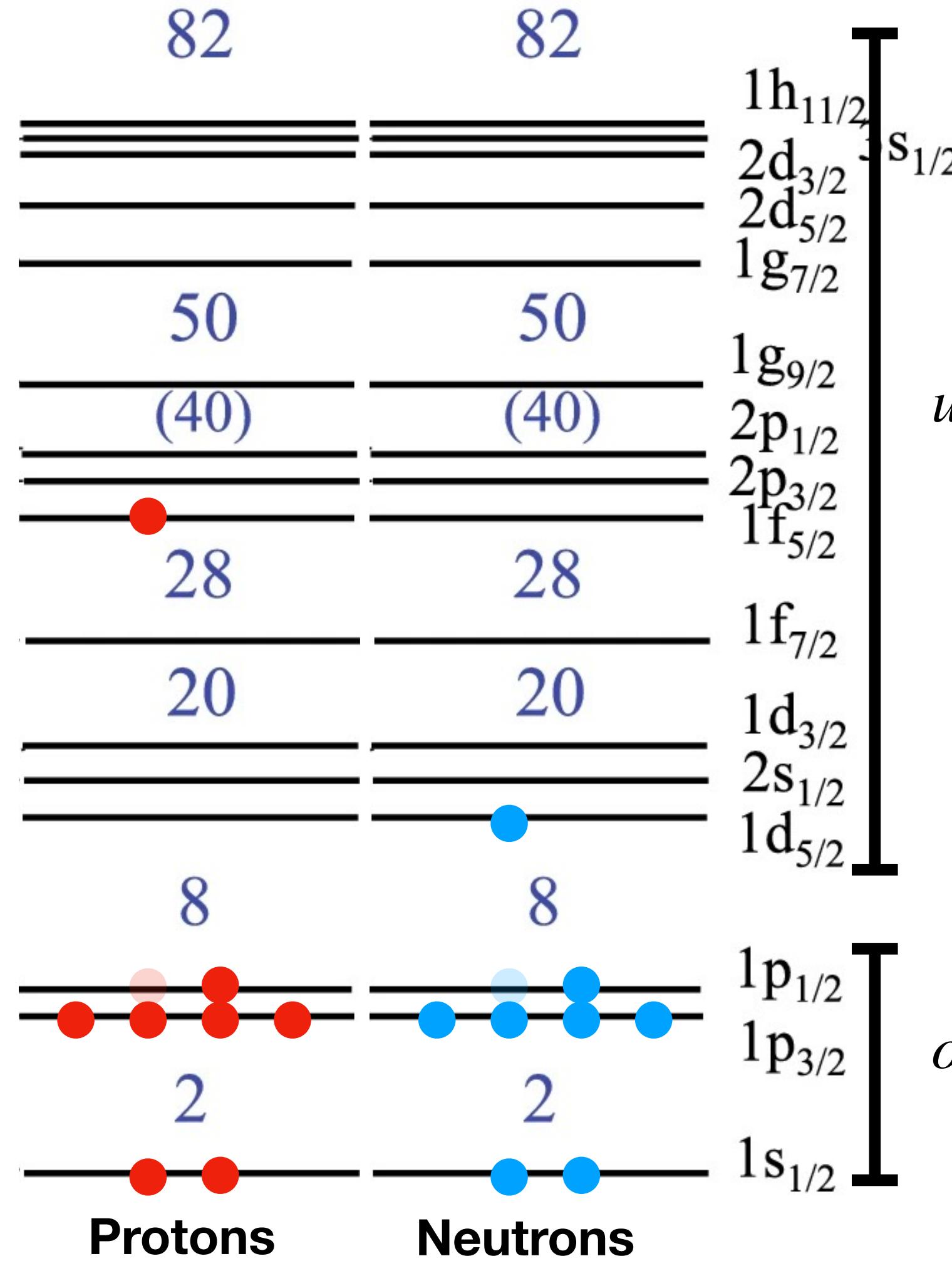
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# Solving the many-body problem



$|\Phi\rangle \sim \prod_{i=1}^o |\phi_i\rangle$  mean-field reference state

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$|\Phi_{ij}^{ab}\rangle$  2p2h excitation ( $\binom{o}{2} \binom{u}{2}$  states)

Scales factorially in  $o$  and  $u$ .

**Obvious scalability problem.**

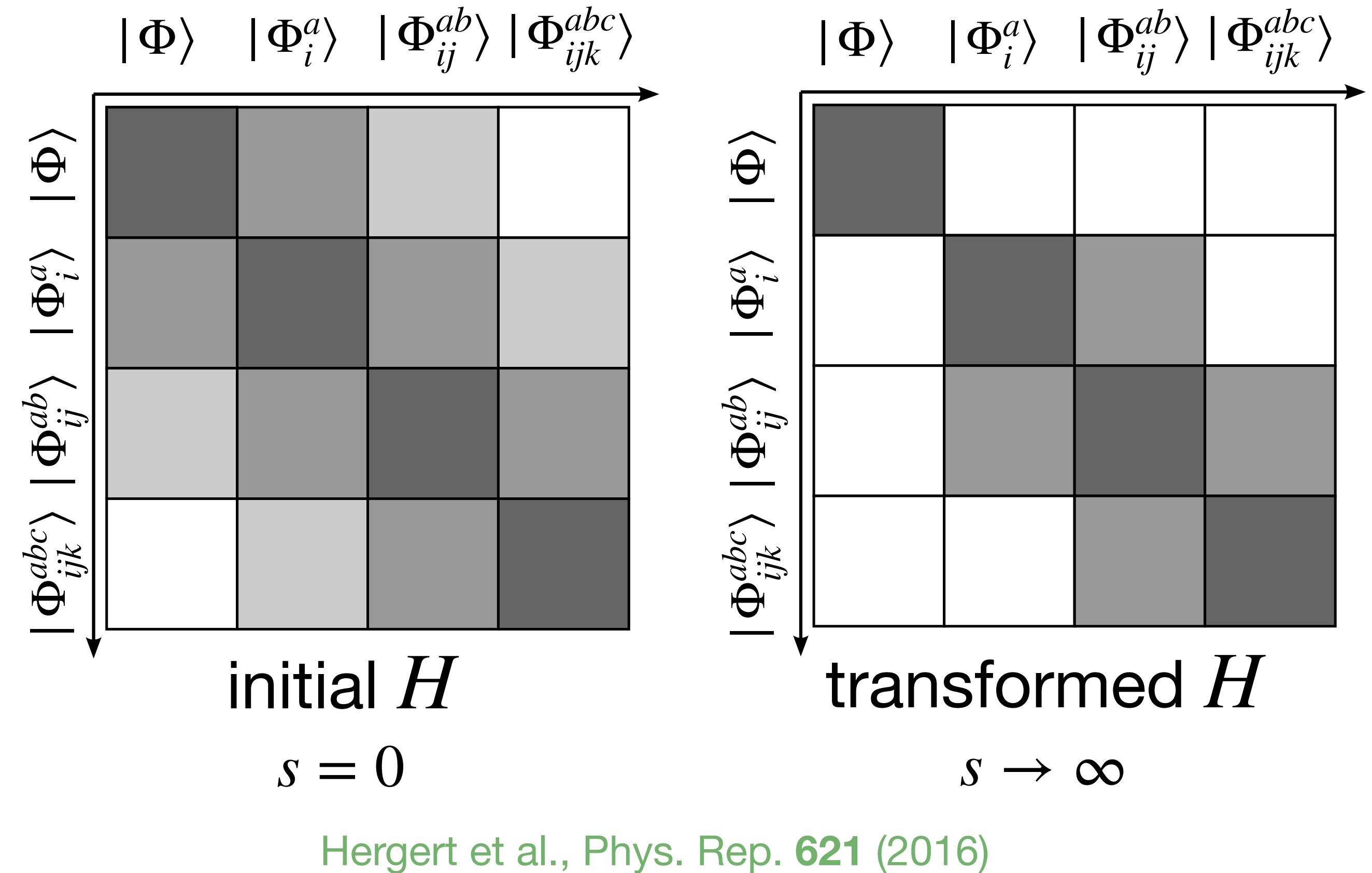
# The IMSRG

in-medium similarity renormalization group

- IMSRG generates unitary transformation of Hamiltonian

$$\frac{dH}{ds} = [\eta, H]$$

- Normal order with respect to  $|\Phi\rangle$  approximately handles **3N forces** and **induced many-body forces**



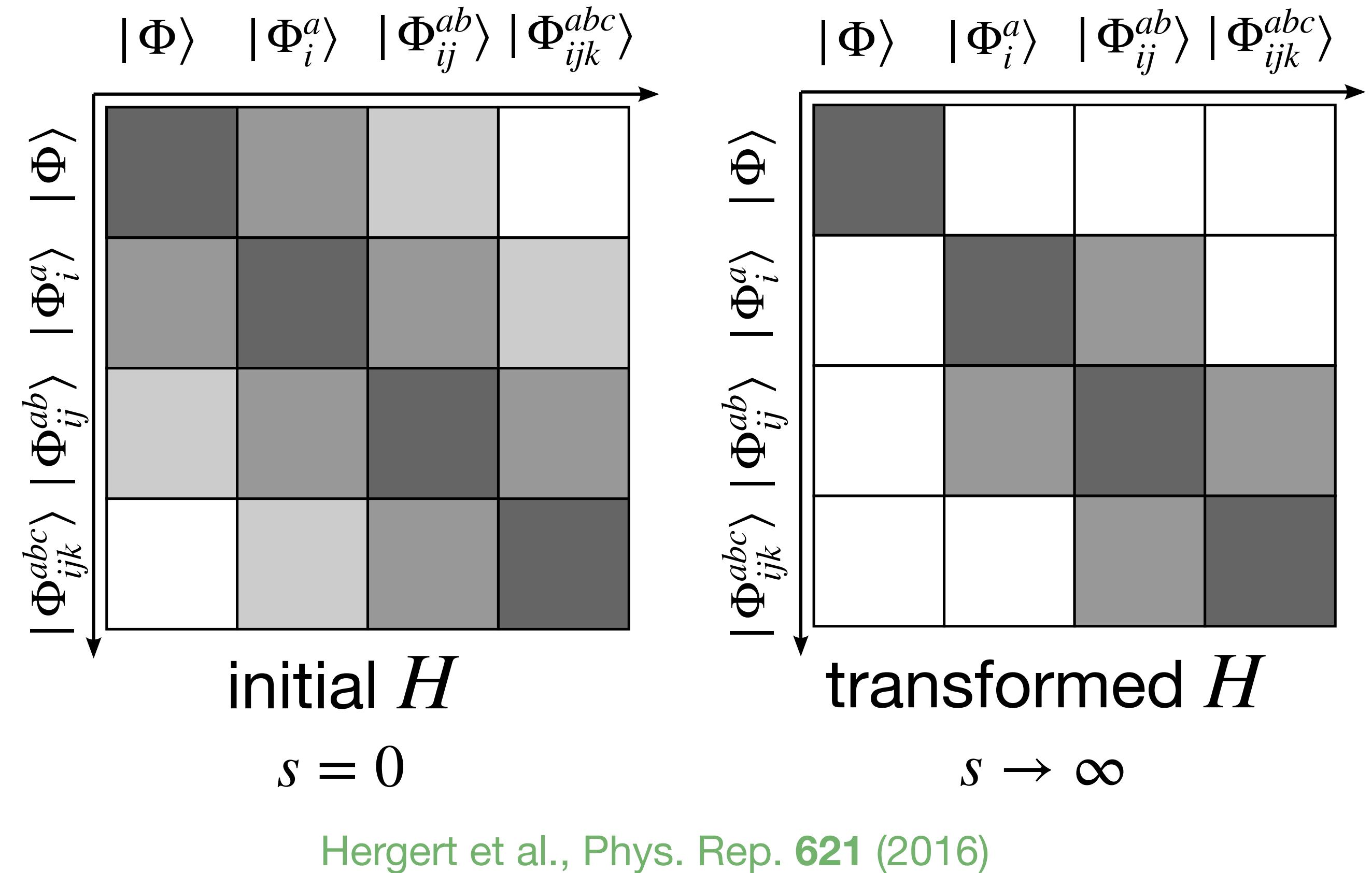
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## Truncation necessary!

- Standard = IMSRG(2)
- More refined = **IMSRG(3)** MH et al., PRC 103 (2021)

# Normal ordering

Hagen et al., PRC 76 (2007)  
Roth et al., PRL 109 (2012)  
Miyagi et al., PRC 105 (2022)  
Hebeler, ..., MH et al., PRC 107 (2023)

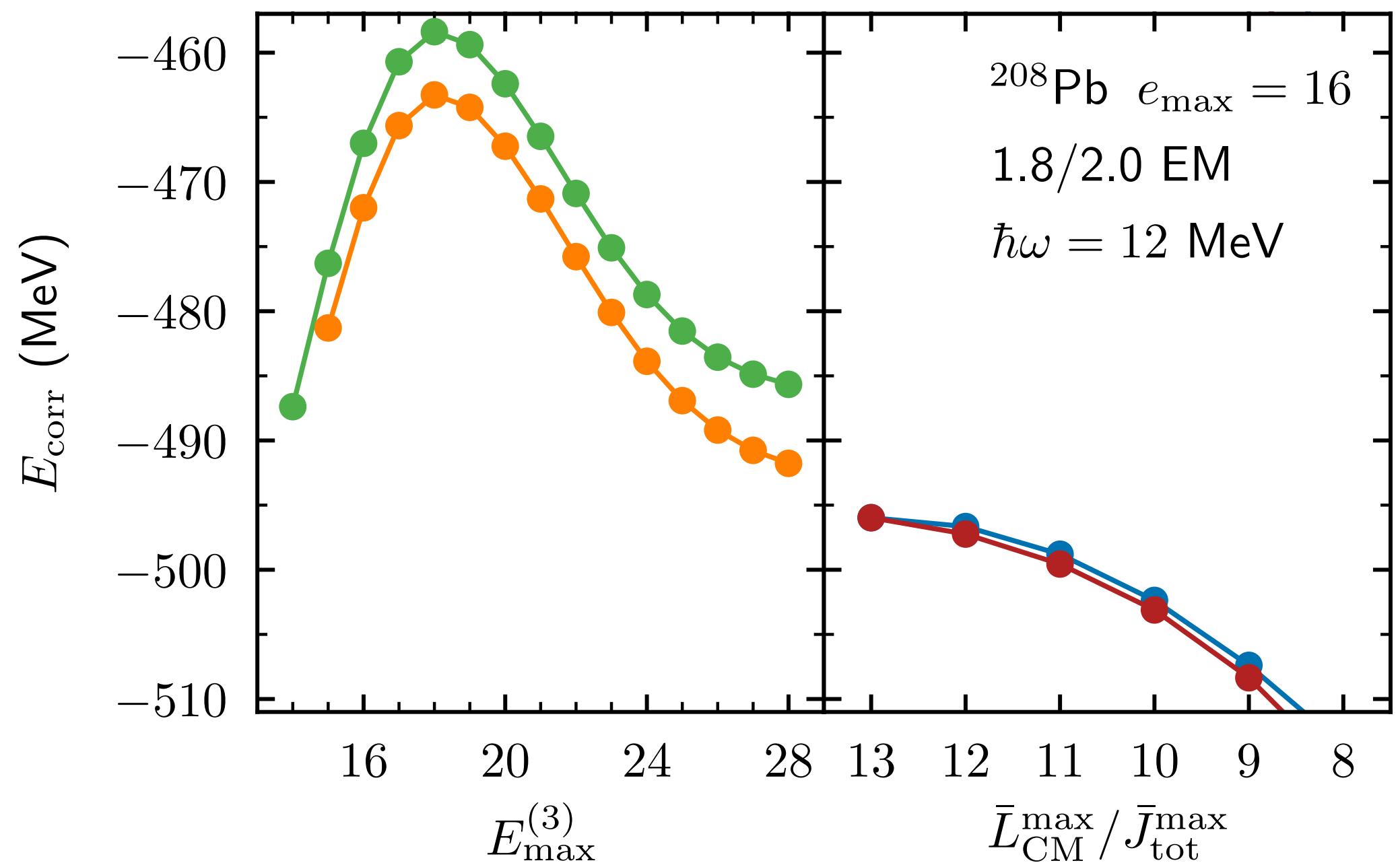
- Normal ordering w.r.t.  $|\Phi\rangle$  gives us effective interactions

$$\Gamma_{pqrs} = V_{\text{NN},pqrs} + \sum_i n_i V_{3\text{N},pqirsi}$$

Discard

$$W_{pqrstu} = V_{3\text{N},pqrstu}$$

- Single-particle representation of  $V_{3\text{N}}$  is expensive (100s of GB or TB)



Hebeler, ..., MH et al., PRC 107 (2023)

# IMSRG truncation

Generator  $\eta$  gives decoupling

$$\frac{dH(s)}{ds} = [\eta(s), H(s)]$$

$$\langle \Phi_i^a | H(\infty) | \Phi \rangle \stackrel{!}{=} 0$$

$$\langle \Phi_{ij}^{ab} | H(\infty) | \Phi \rangle \stackrel{!}{=} 0$$

...

**Truncate operators in many-body rank**

$$H(0) = E_0 + f + \Gamma$$

$$\eta(0) = \eta^{(1\text{B})} + \eta^{(2\text{B})}$$

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**Truncate operators in many-body rank**

$$H(s) = E(s) + f(s) + \Gamma(s) + W(s) + \dots$$

$$\eta(s) = \eta^{(1B)}(s) + \eta^{(2B)}(s) + \eta^{(3B)}(s) + \dots$$

$$H(0) = E_0 + f + \Gamma$$

$$\eta(0) = \eta^{(1B)} + \eta^{(2B)}$$

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**IMSRG(2)**

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MH et al., PRC 103 (2021)

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## IMSRG(3)

MH et al., PRC **103** (2021)

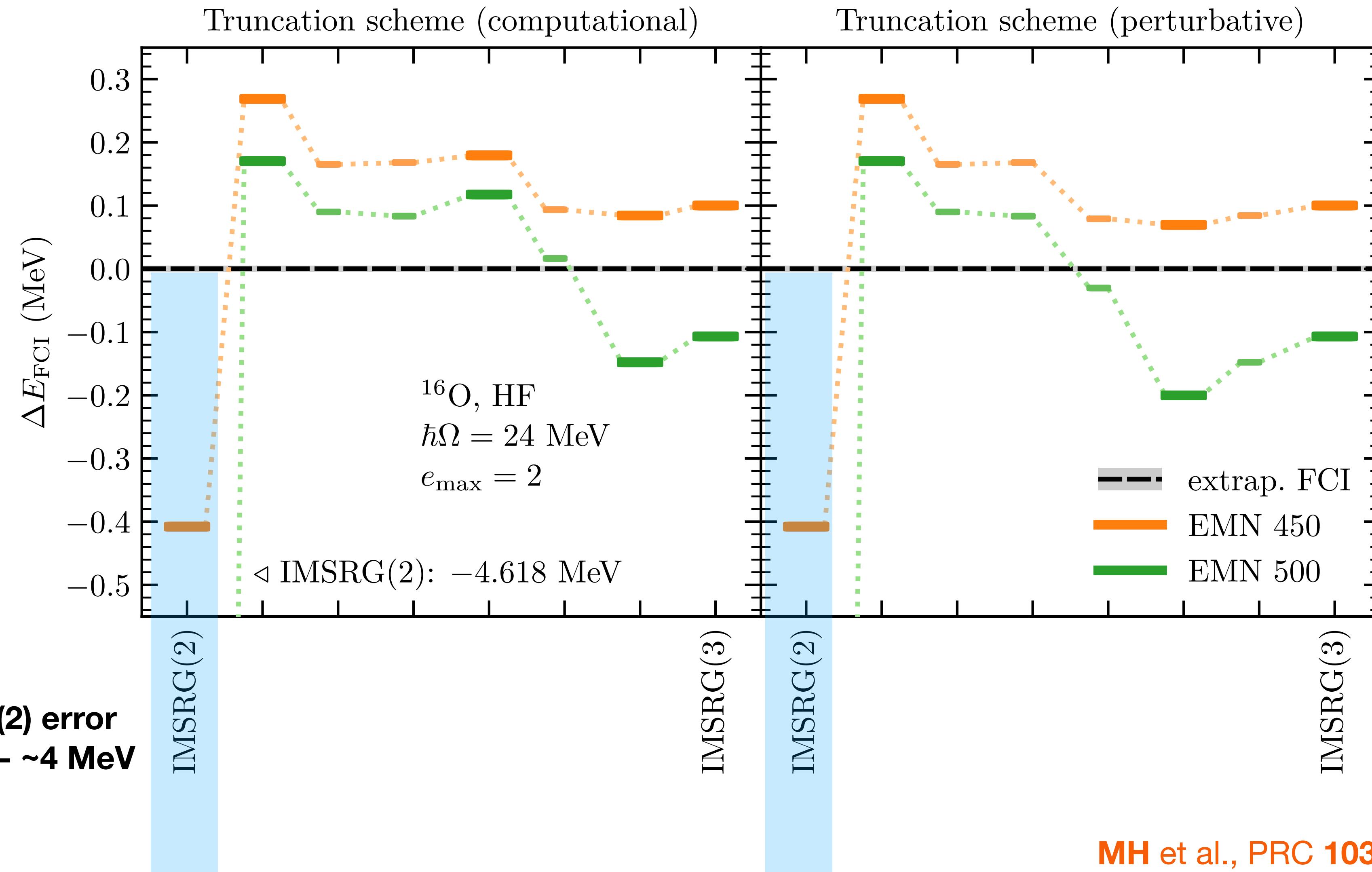
**IMSRG(3)** = more precision, but higher cost [ $\mathcal{O}(N^9)$ ,  $\mathcal{O}(N^7)$  with approx]:

- Automated derivation with Drudge
- Angular momentum coupling with AMC
- **No automatic code generation yet**

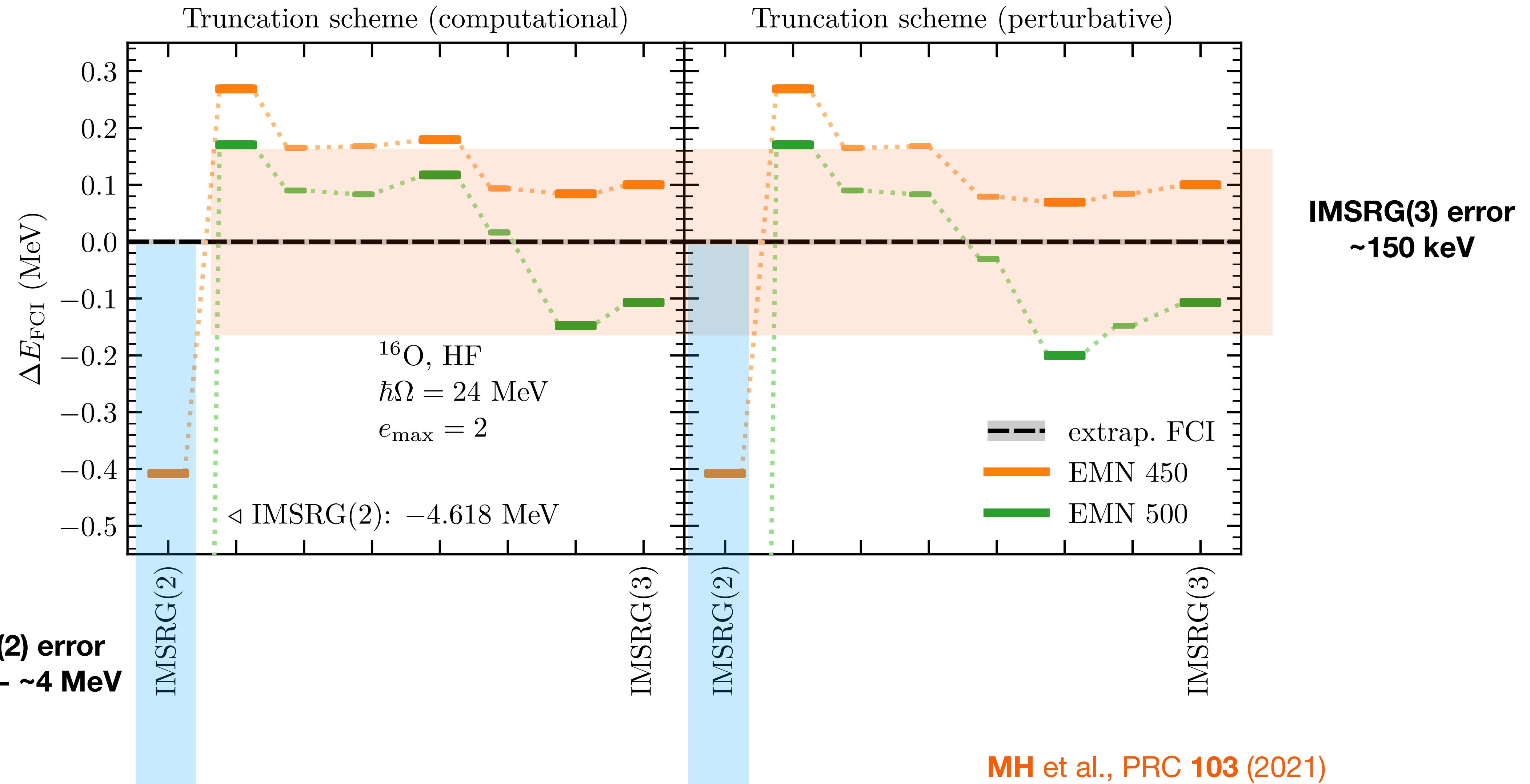
Zhao, Scuseria, <https://github.com/tschijnmo/drudge> (2021)

Tichai et al., EPJA **56** (2020)

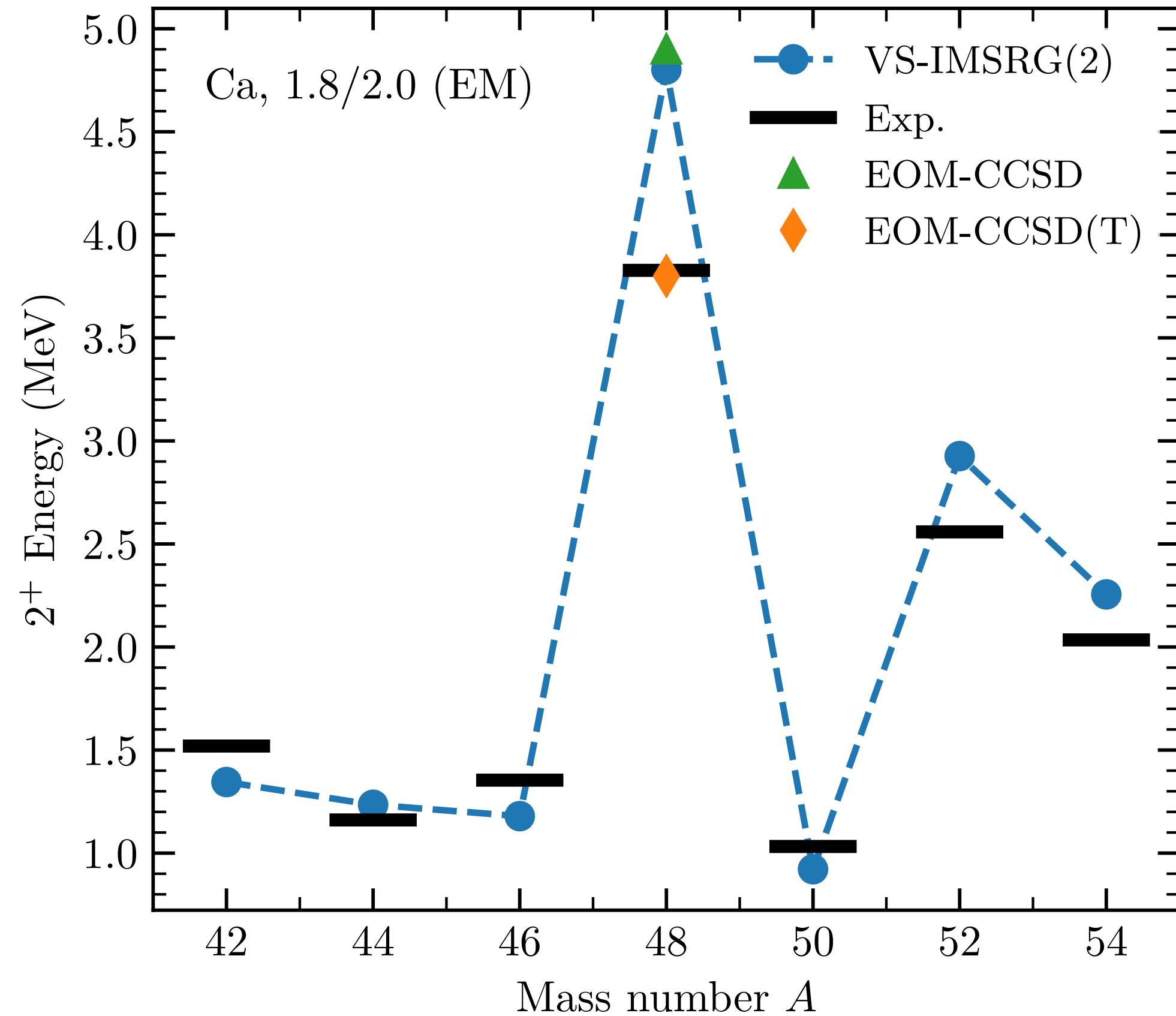
# Where does IMSRG(3) matter?



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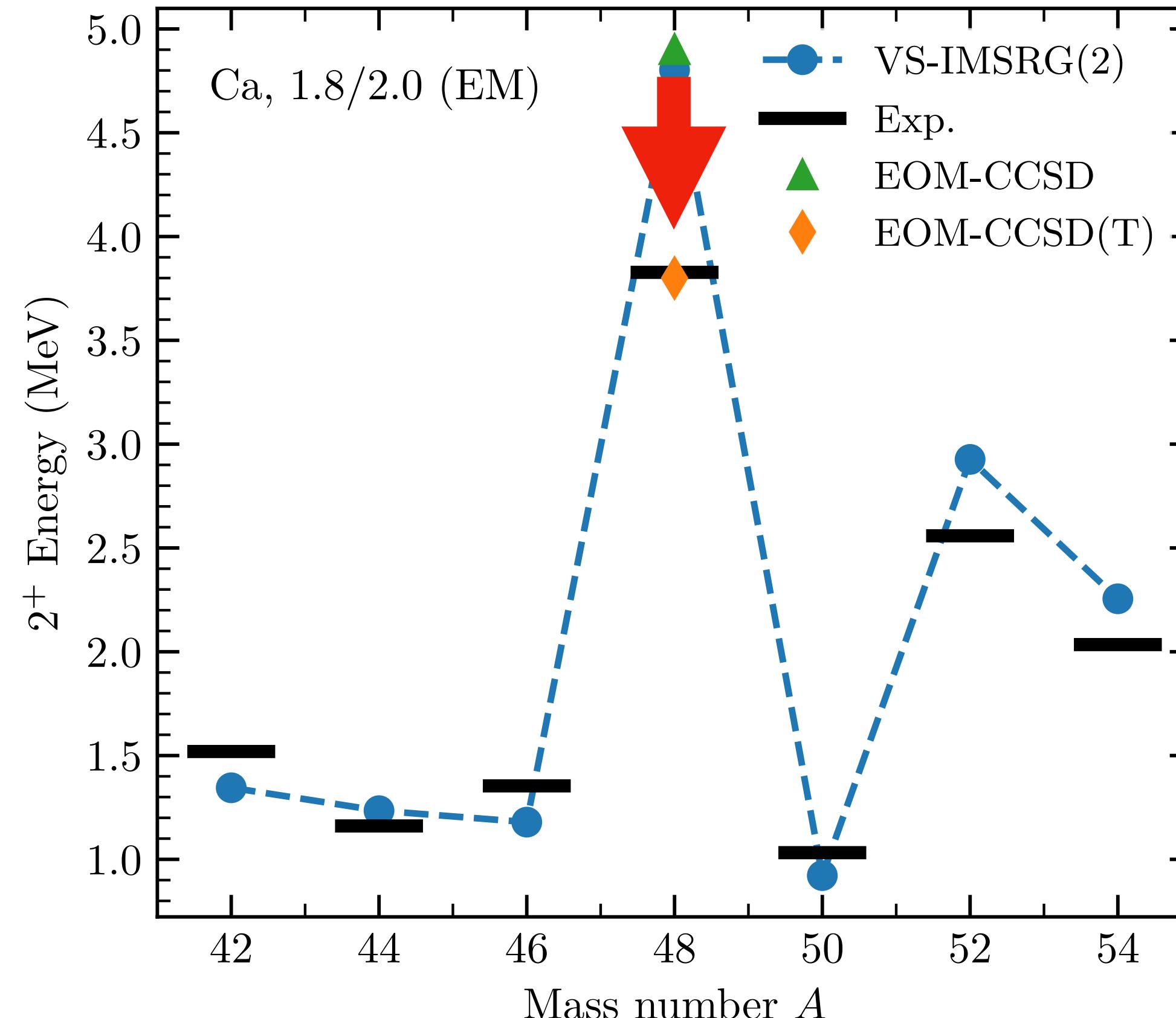
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Hagen et al., PRL 117 (2016)

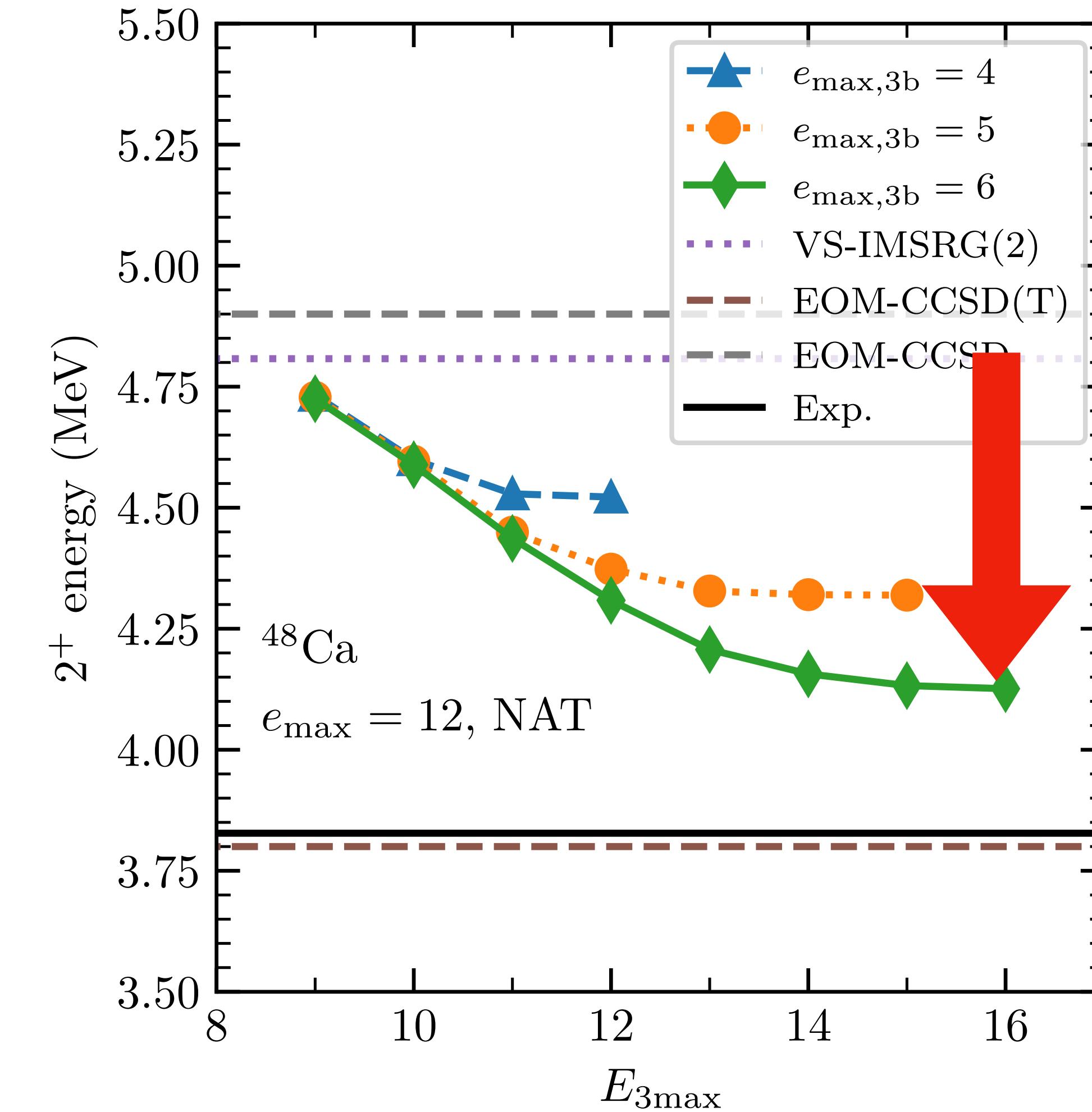
Simonis et al., PRC 96 (2017)

# Where does IMSRG(3) matter?



Hagen et al., PRL 117 (2016)

Simonis et al., PRC 96 (2017)



# Comparing IMSRG to CC

## IMSRG

- ODE solution
- Similarity transformation
- No factorization (binary contractions)
- $|\Psi\rangle = \exp(\Omega) |\Phi\rangle$ , with (for example)

$$\Omega^{(2)} = \frac{1}{4} \sum_{pqrs} \Omega_{pqrs}^{(2)} a_p^\dagger a_q^\dagger a_s a_r$$

- **Hermitian**, unitary, BCH does not truncate, but (typically) converges
- IMSRG(3) can include residual 3B interactions

## CC

- Iterative solution
  - Similarity transformation
  - Factorization important
  - $|\Psi\rangle = \exp(T) |\Phi\rangle$ , with (for example)
- $$T^{(2)} = \frac{1}{4} \sum_{abij} T_{abij}^{(2)} a_a^\dagger a_b^\dagger a_j a_i$$
- Non-Hermitian, unitary, BCH truncates after finite commutators
  - CCSDT (NO2B)  $\neq$  CC with 3B interactions

**\*\*90% correct**

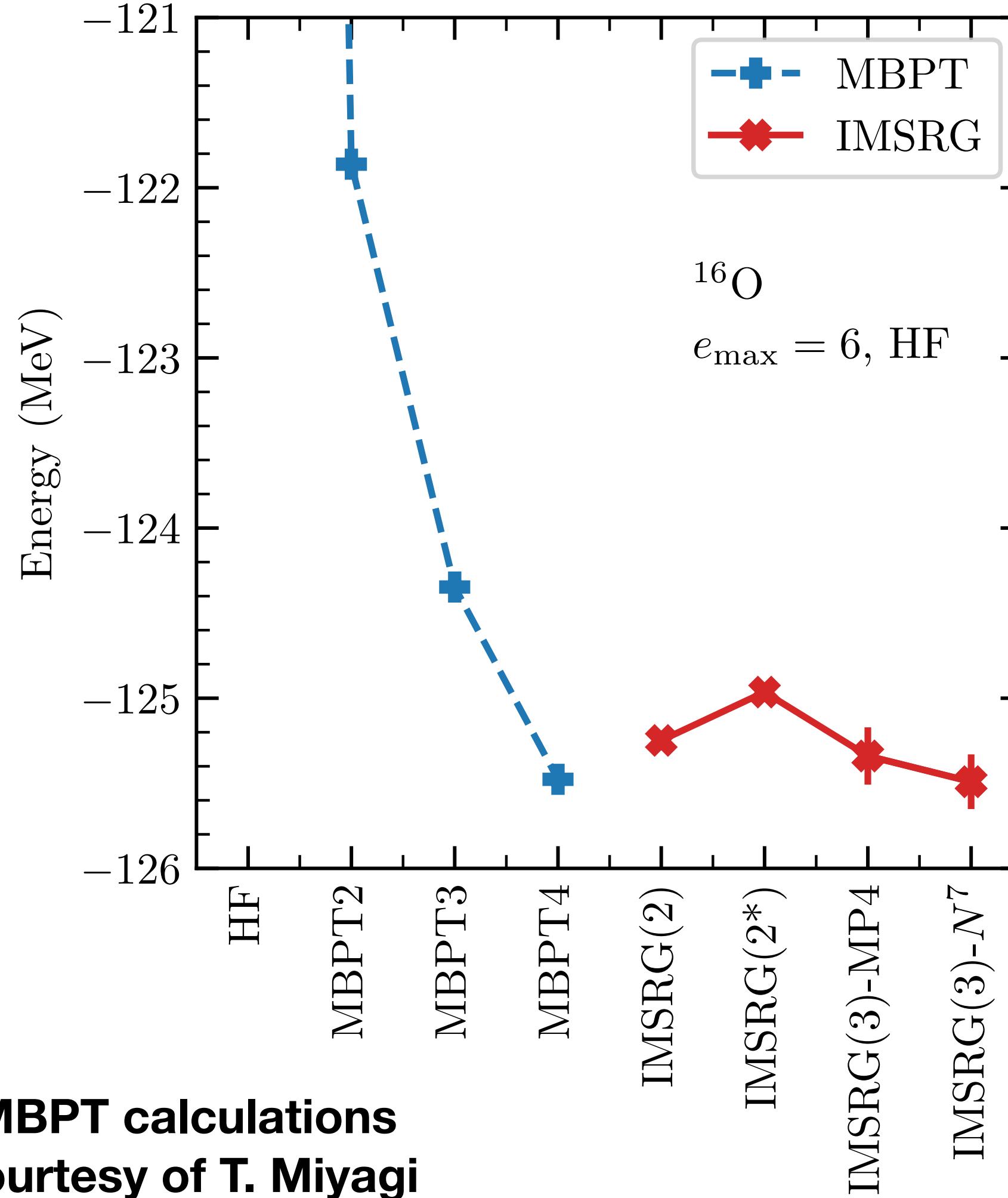
# Connecting IMSRG and MBPT

Hjorth-Jensen et al., Phys. Rep. **261** (1995)

Bartlett, Shavitt, *Many-Body Methods in Chemistry and Physics* (2009)

Hergert et al., Phys. Rep. **621** (2016)

# IMSRG truncations in detail



## IMSRG(2)

- Complete up to **MPBT3**
- Some higher-order effects + **nonperturbative ladders and rings**

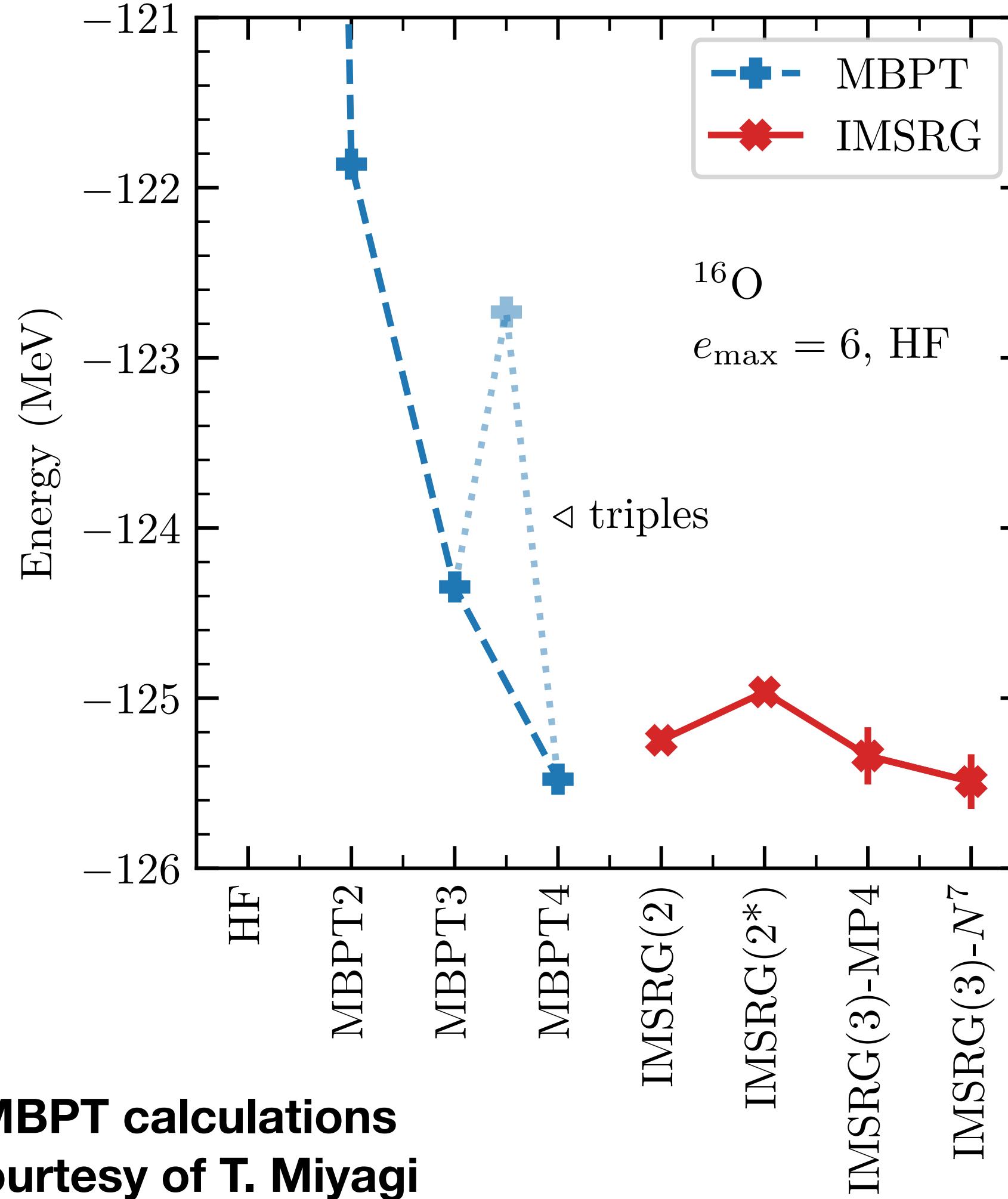
## IMSRG(3)-MP4

- Complete up **MBPT4**

## IMSRG(3)- $N^7$

- Perturbatively less important terms also included

# IMSRG truncations in detail



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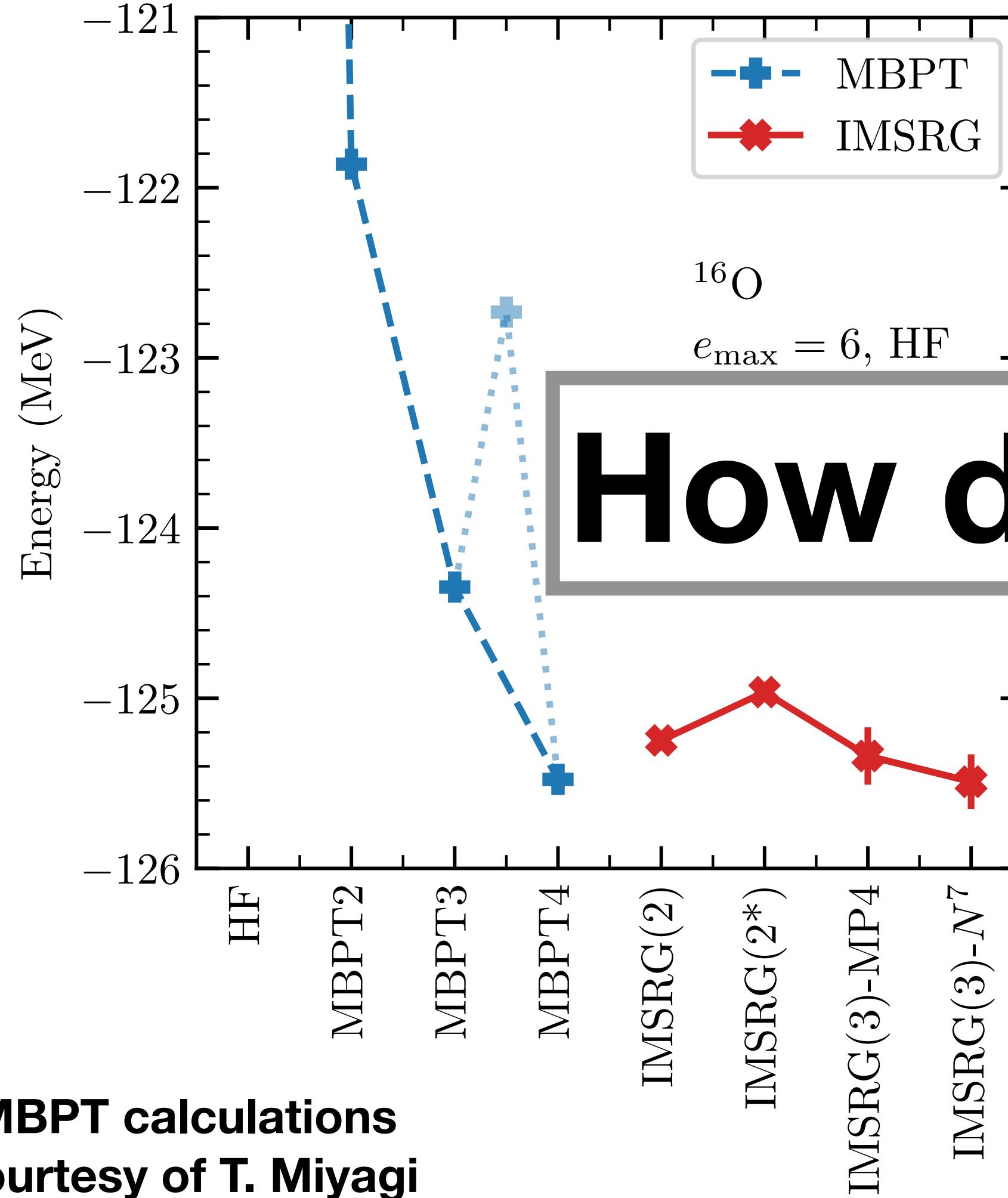
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# IMSRG truncations in detail



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- Complete up to **MPBT3**
- Some higher-order effects +

How do we know this?

## IMSRG(3)-MP4

- Complete up **MBPT4**

## IMSRG(3)- $N^7$

- Perturbatively less important terms also included

# IMSRG ingredients (to start)

$$H(s = 0) = \begin{array}{c} \text{Diagram: } \bullet \text{---} \times \\ \text{Diagram: } \bullet \text{---} \bullet \\ \text{Diagram: } \bullet \text{---} \bullet \text{---} \bullet \end{array} + \begin{array}{c} \text{Diagram: } \vee \\ \text{Diagram: } \vee \text{---} \vee \\ \text{Diagram: } \vee \text{---} \vee \text{---} \vee \end{array} + \begin{array}{c} f_{ia}/\varepsilon_i^a \\ \Gamma_{ijab}/\varepsilon_{ij}^{ab} \\ W_{ijkabc}/\varepsilon_{ijk}^{abc} \end{array}$$

Solving  $\frac{dH}{ds} = [\eta(s), H(s)]$  generates new diagrams out of  $\eta$  and  $H$  vertices

# IMSRG ingredients (to start)

$$H(s=0) = \boxed{\begin{array}{c} \text{+} \\ \text{+} \\ \text{+} \end{array}} + \begin{array}{c} \text{+} \\ \text{+} \\ \text{+} \end{array}$$

$H(s=0) =$

$\eta(s=0) =$

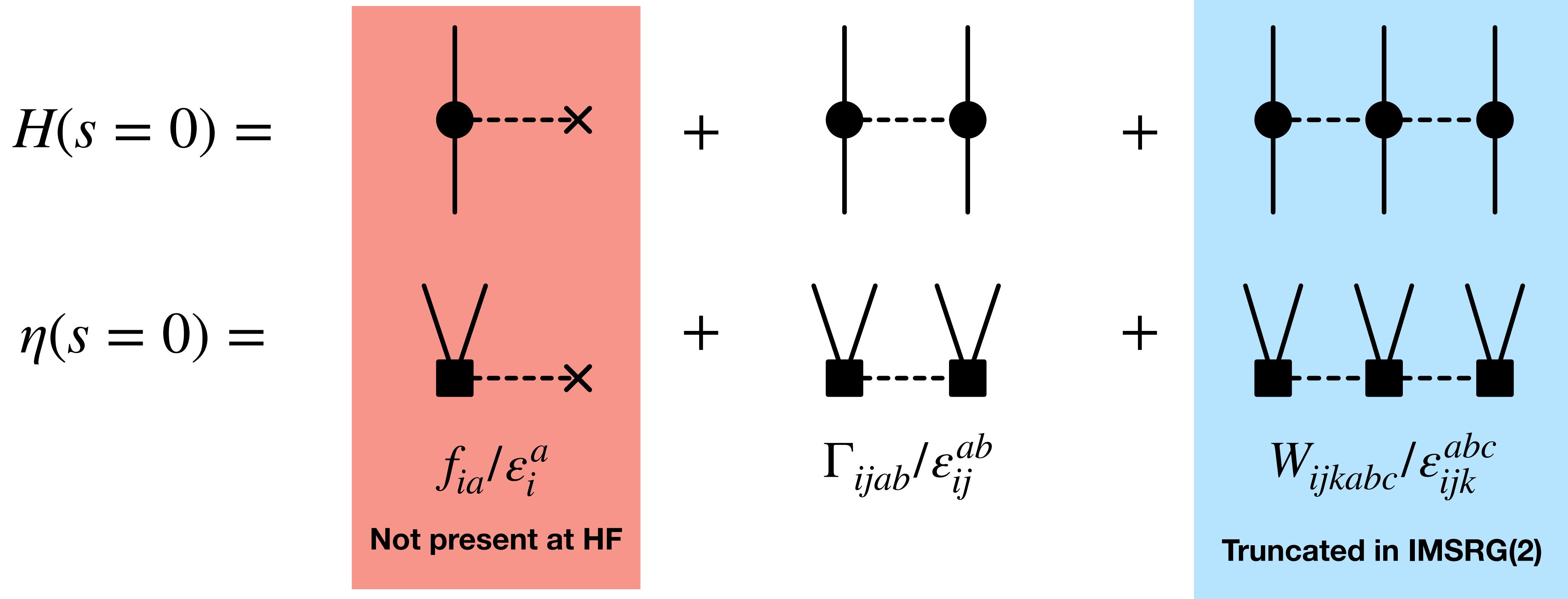
$f_{ia}/\epsilon_i^a$

$W_{ijkabc}/\epsilon_{ijk}^{abc}$

Not present at HF

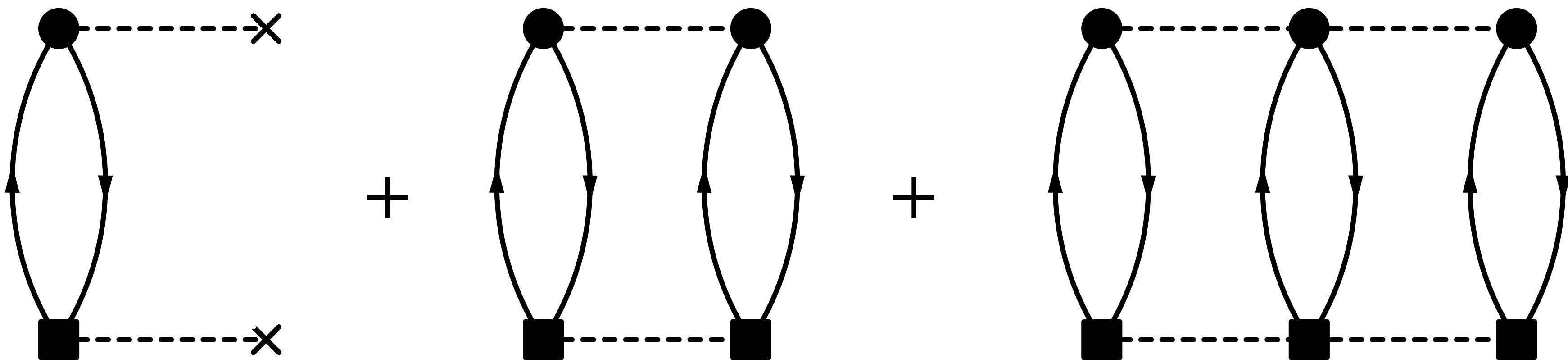
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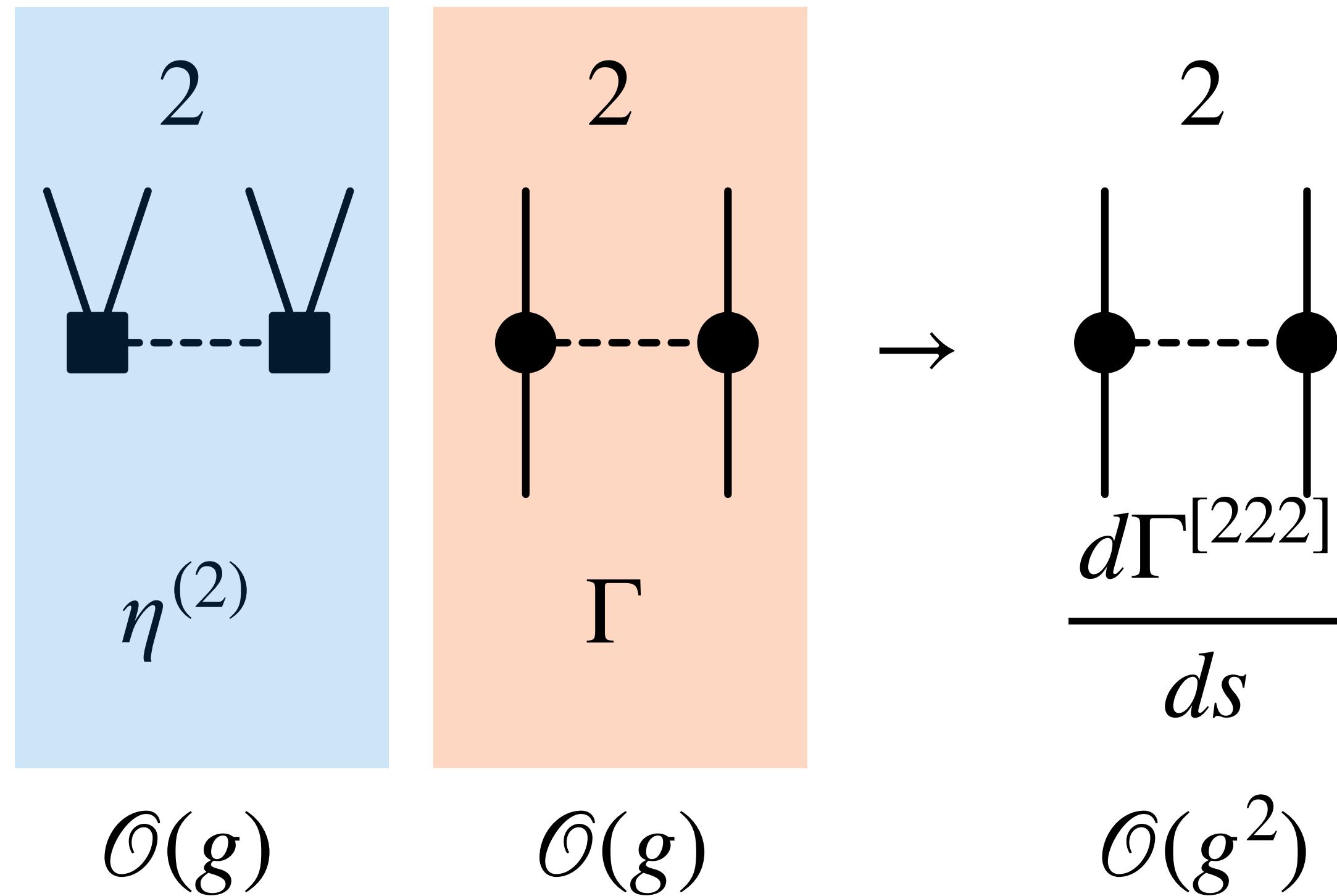
Solving  $\frac{dH}{ds} = [\eta(s), H(s)]$  generates new diagrams out of  $\eta$  and  $H$  vertices

# Energy contributions

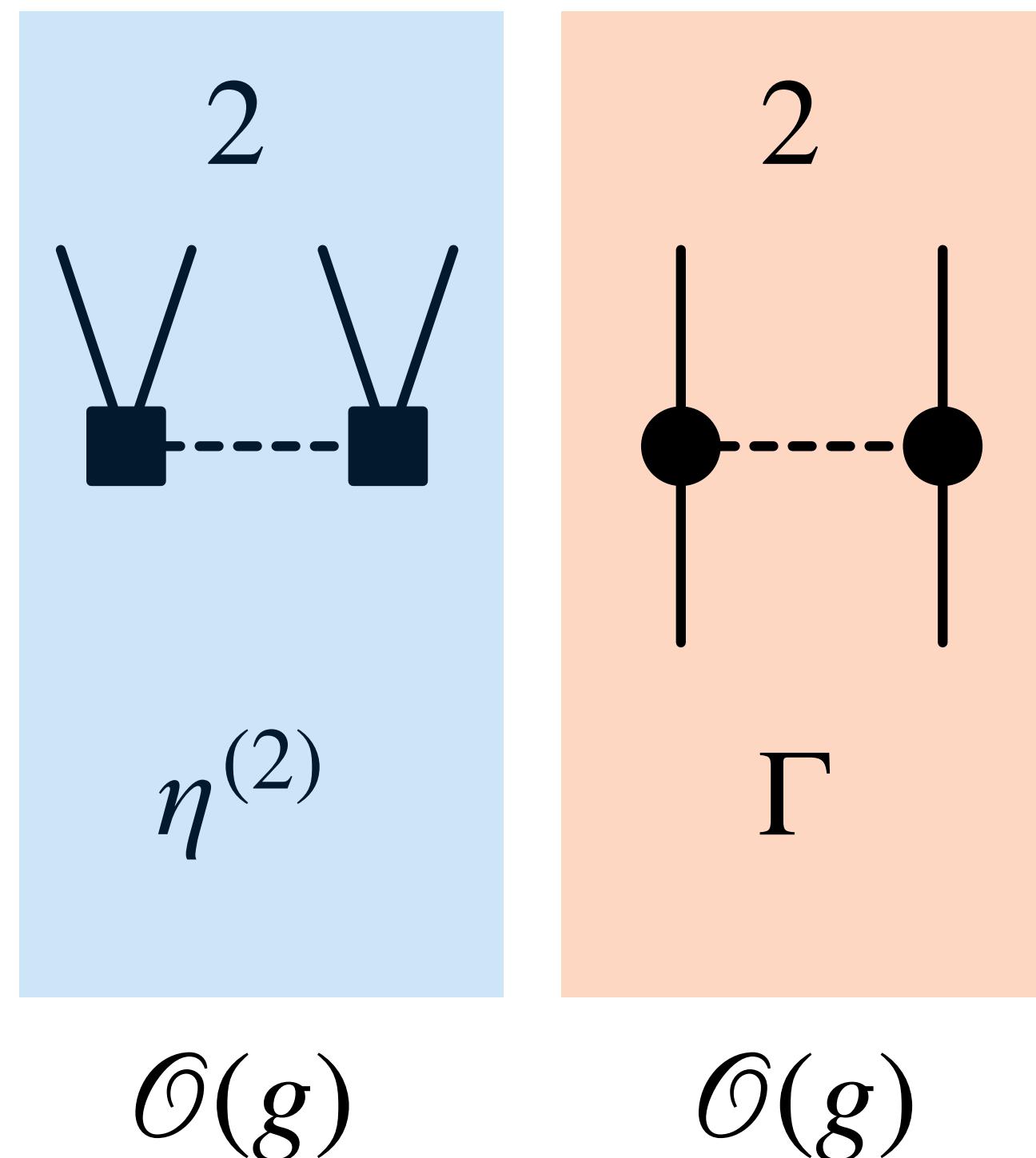
$$\frac{dH^{(0B)}}{ds} = \frac{dE}{ds} =$$


- Zero-body part produces MBPT2-like energy contributions
- Insertion of  $H(s = 0)$  and  $\eta(s = 0)$  gives MBPT2

# Effective 2B interactions



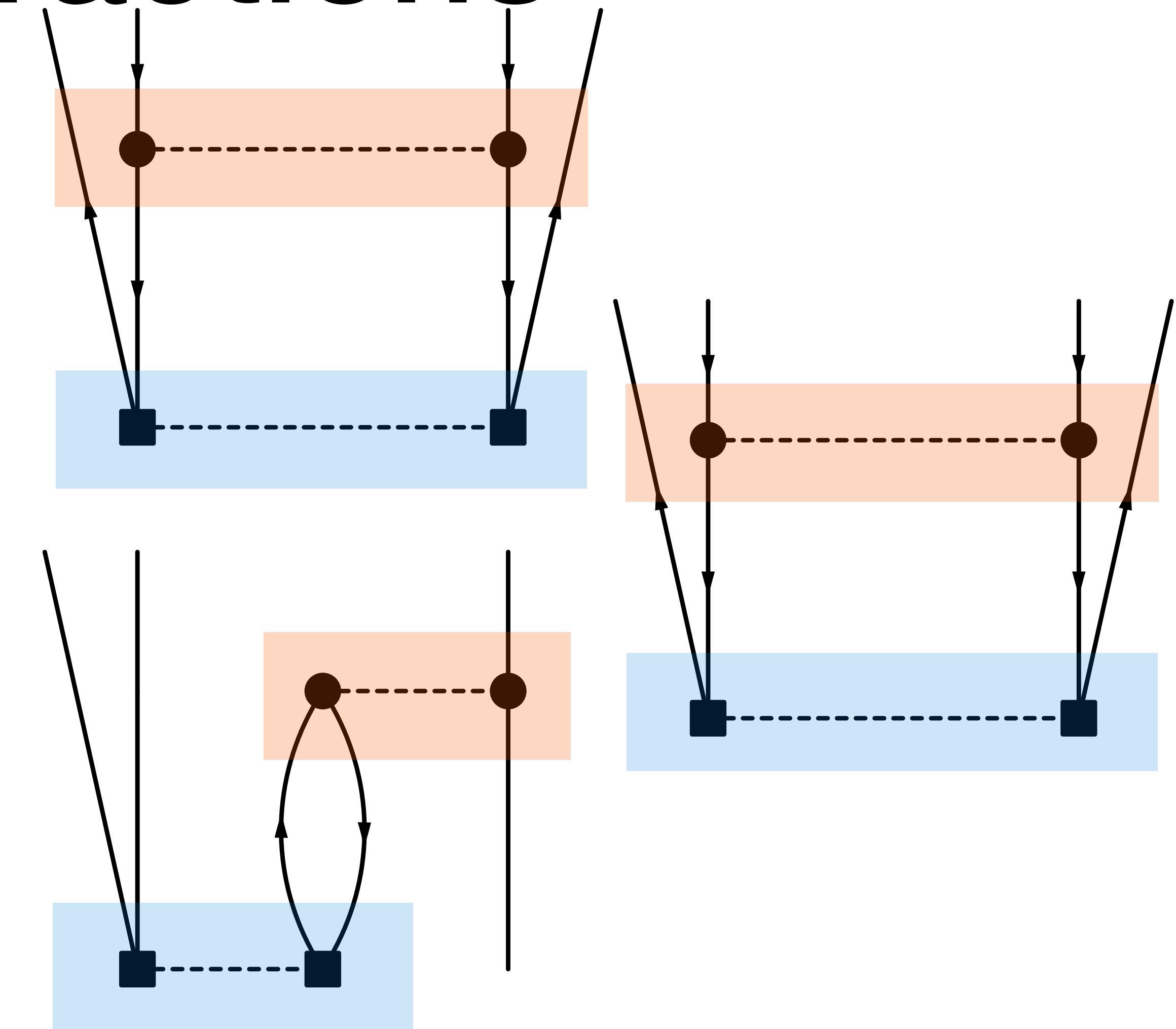
# Effective 2B interactions



→

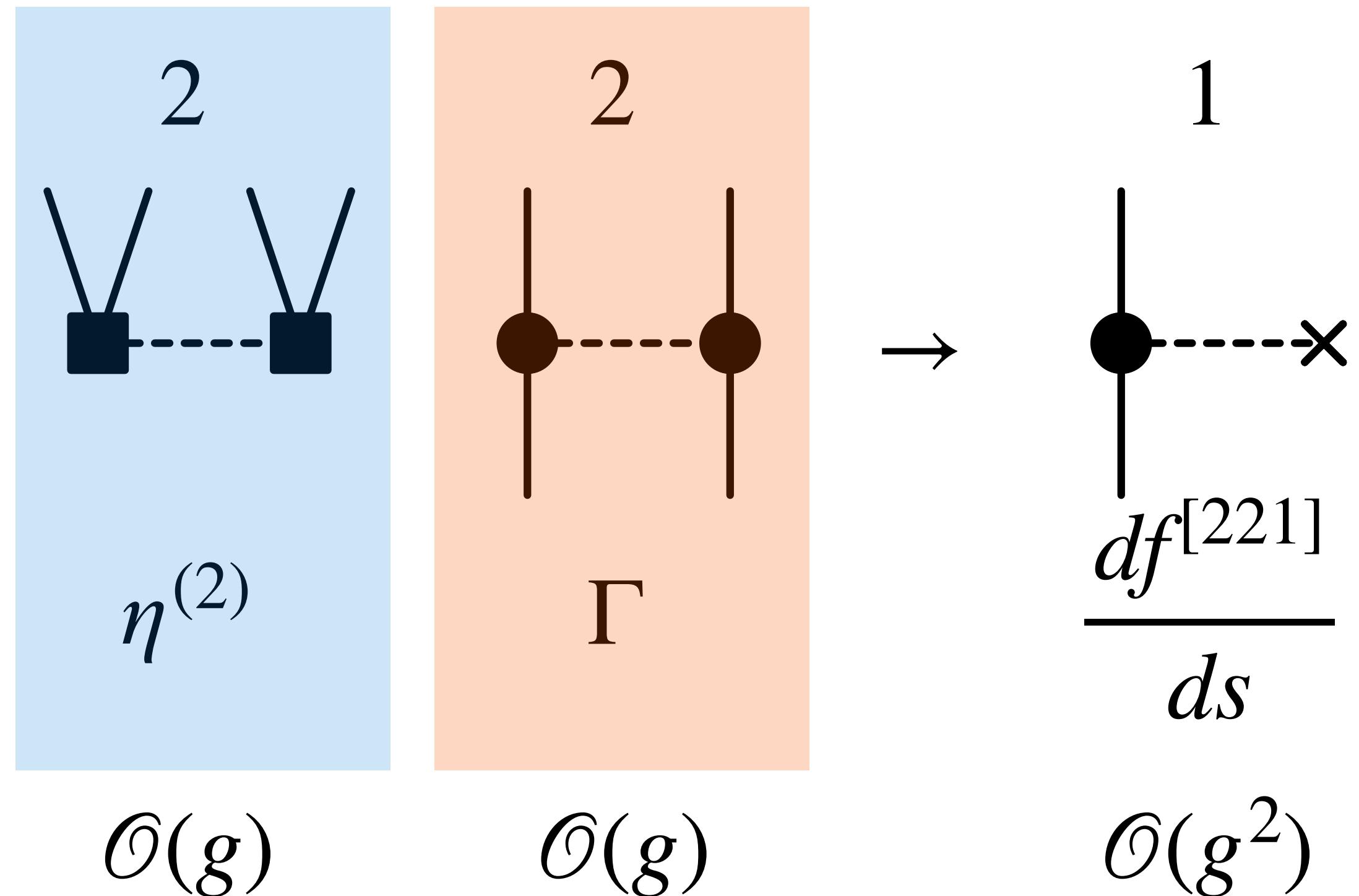
$$\frac{d\Gamma^{[222]}}{ds}$$

Diagram illustrating the effective two-body interaction generated at order  $\mathcal{O}(g^2)$ . It shows two particles, each represented by a black circle with two internal lines, labeled with index 2 above them. They are positioned on a light orange background, which is labeled  $\Gamma$ . Below the circles, the interaction order is given as  $\mathcal{O}(g^2)$ .



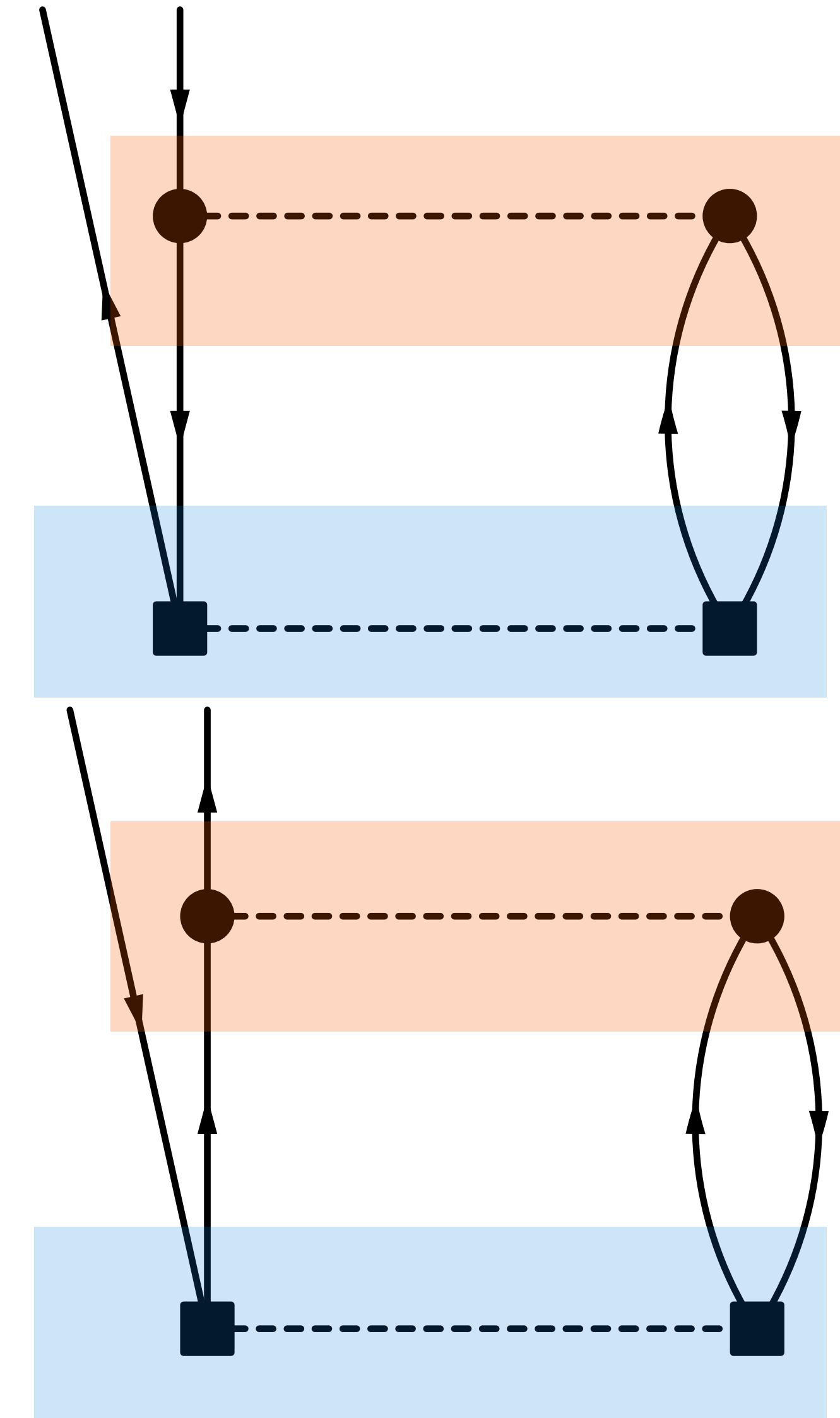
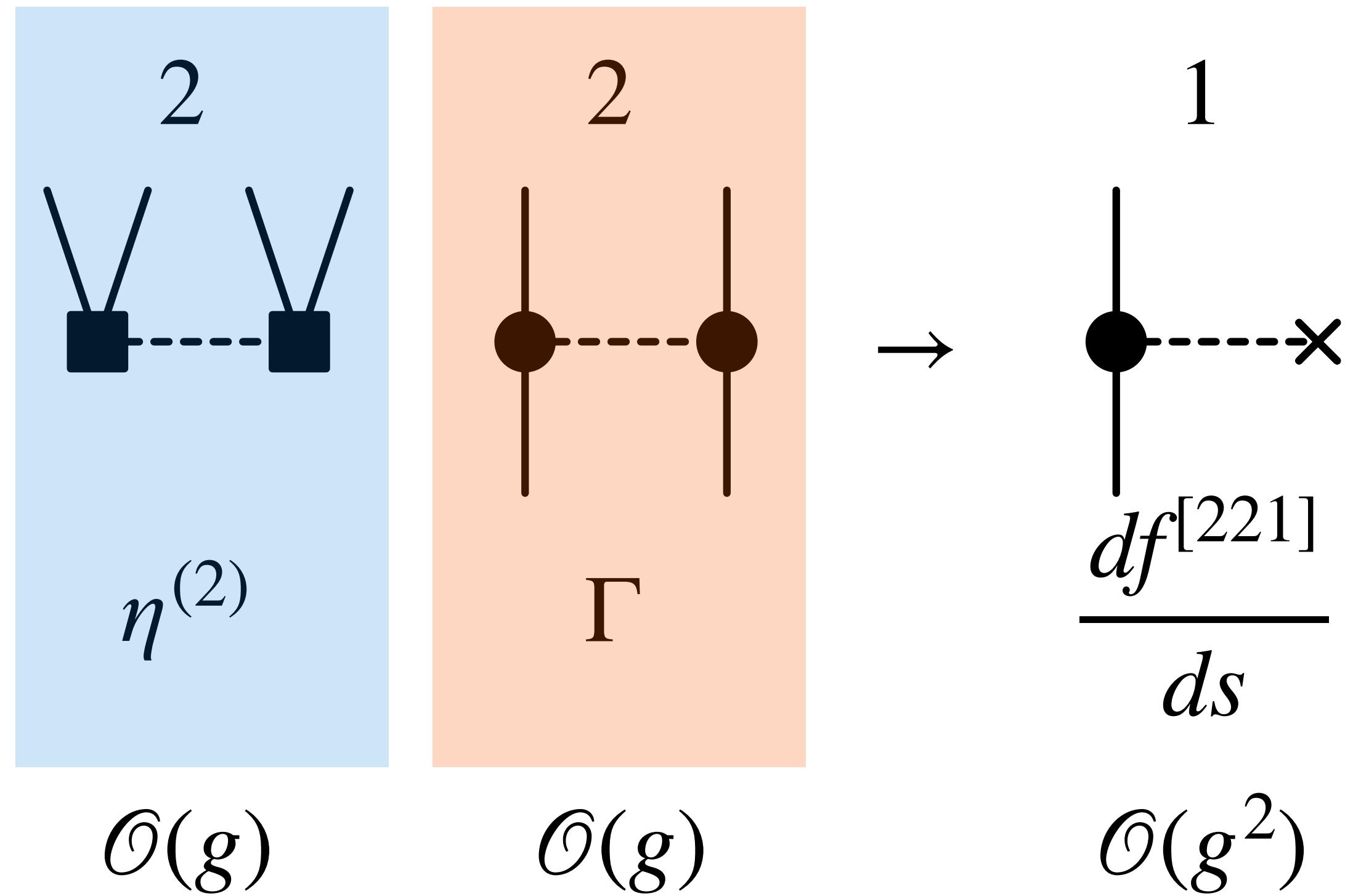
Effective 2B interactions generated to get MBPT3 right.

# Effective 1B interactions



IMSRG "induces" effective 1B interactions

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IMSRG "induces" effective 1B interactions

# Technical details

1. Consider diagrams generated by  $[\eta(s = 0), H(s = 0)]$  within truncation
  - IMSRG(2): **No induced effective 3B interactions**
2. Repeat up to desired perturbative order ( $\mathcal{O}(g^3)$ ,  $\mathcal{O}(g^4)$ )
3. Look at what diagrams are generated and compare with MBPT

Technical complications (\*\*missing 10%):

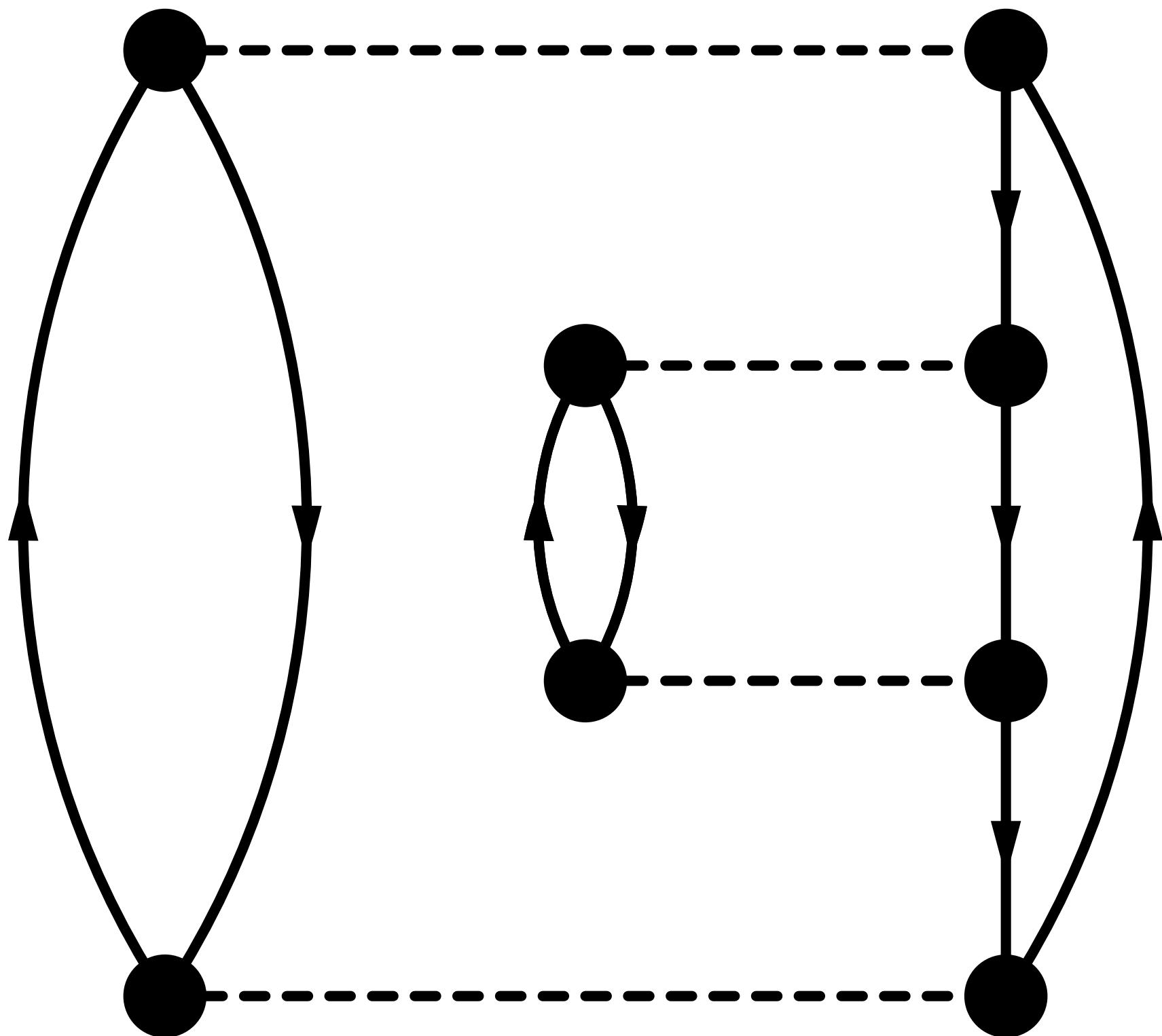
- Integrating in  $s$
- Dealing with complementary diagrams (at and beyond MBPT4)

# Improving the IMSRG(2)

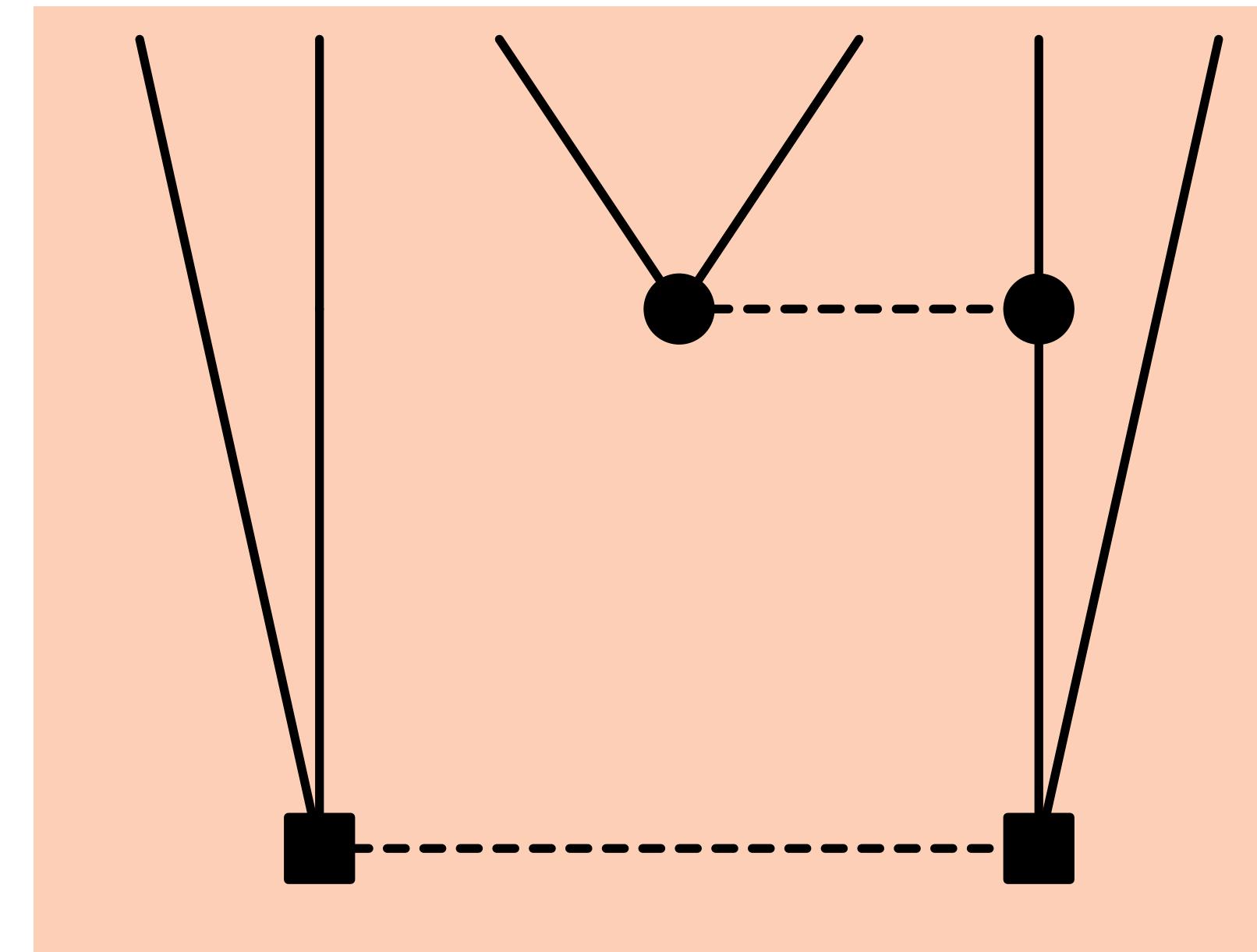
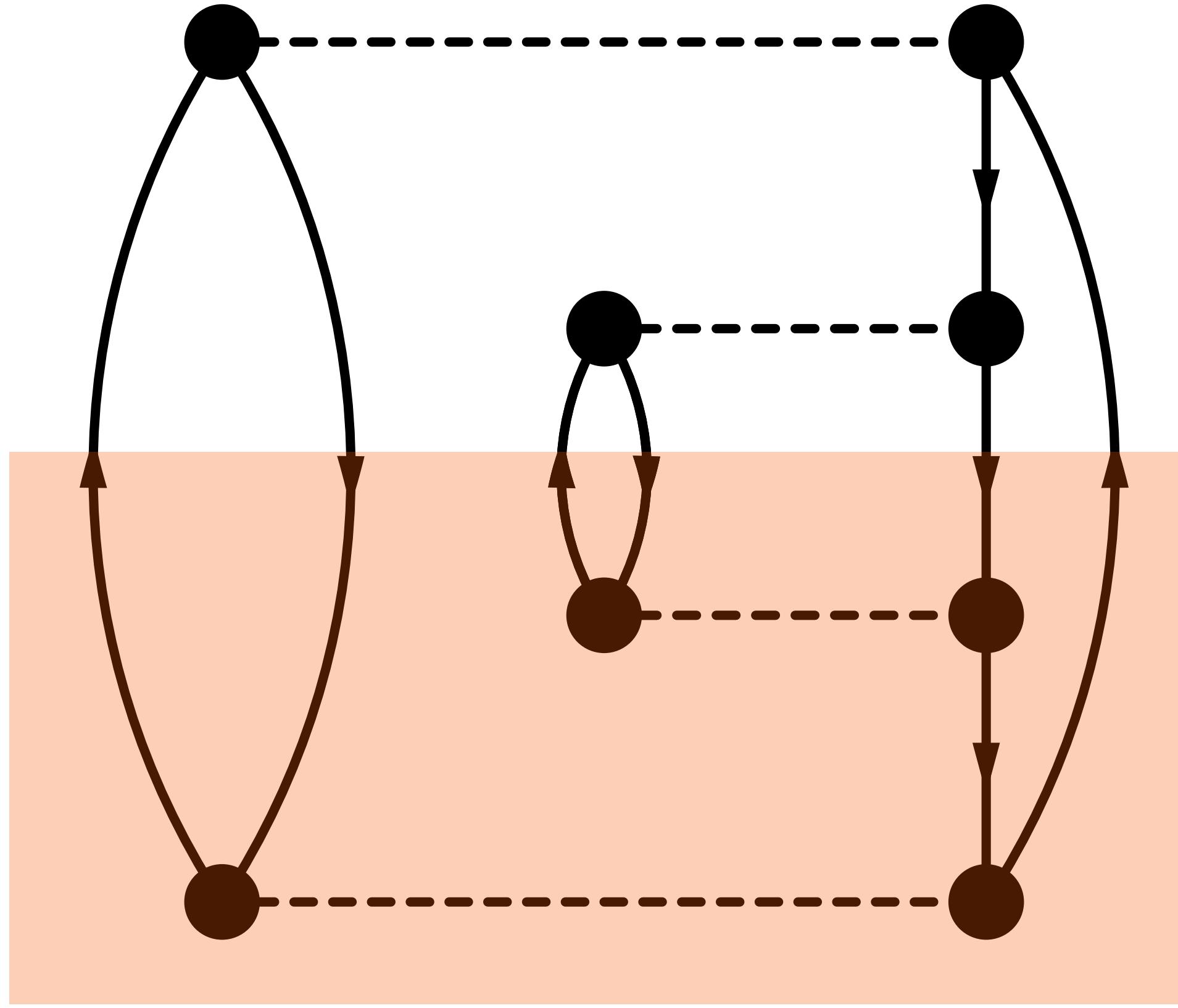
Morris, PhD Thesis, MSU (2016)

Arthuis et al., CPC **240** (2019)

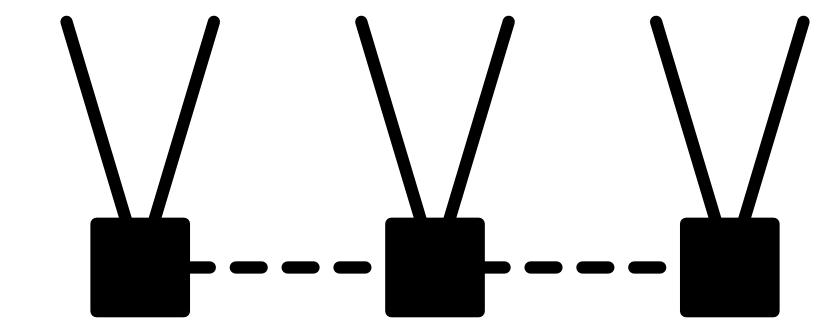
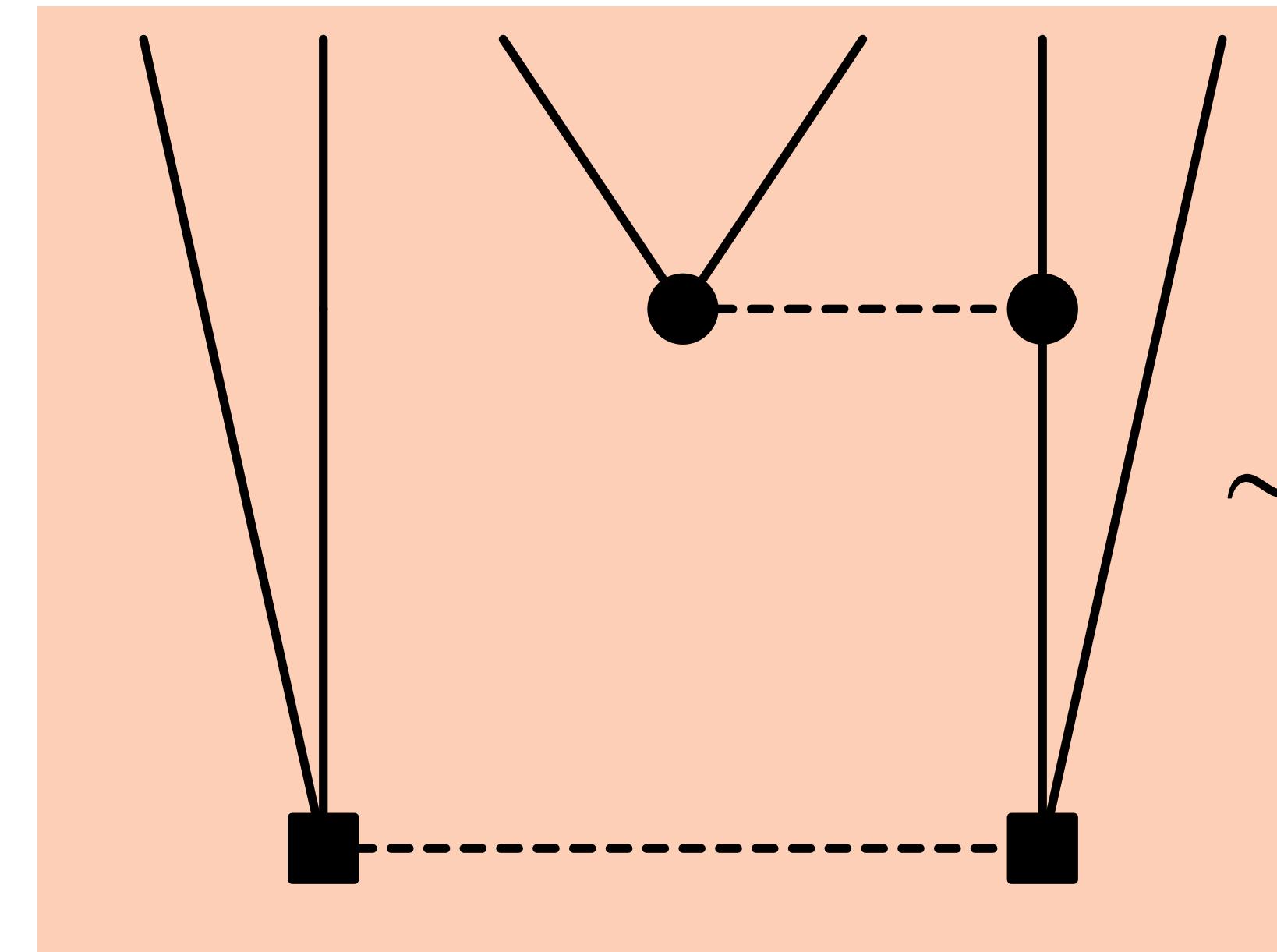
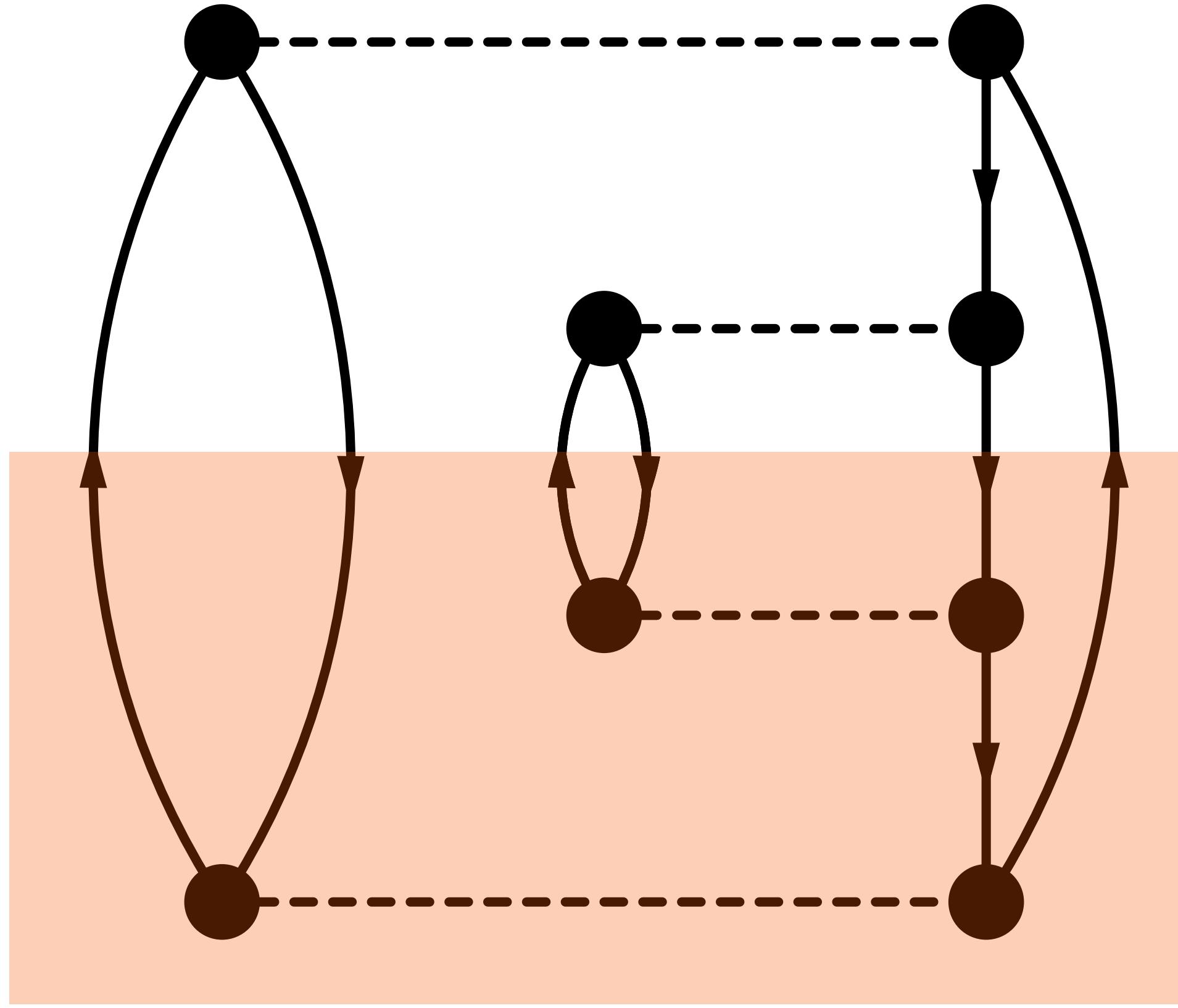
# Triples (MBPT4) missing in IMSRG(2)



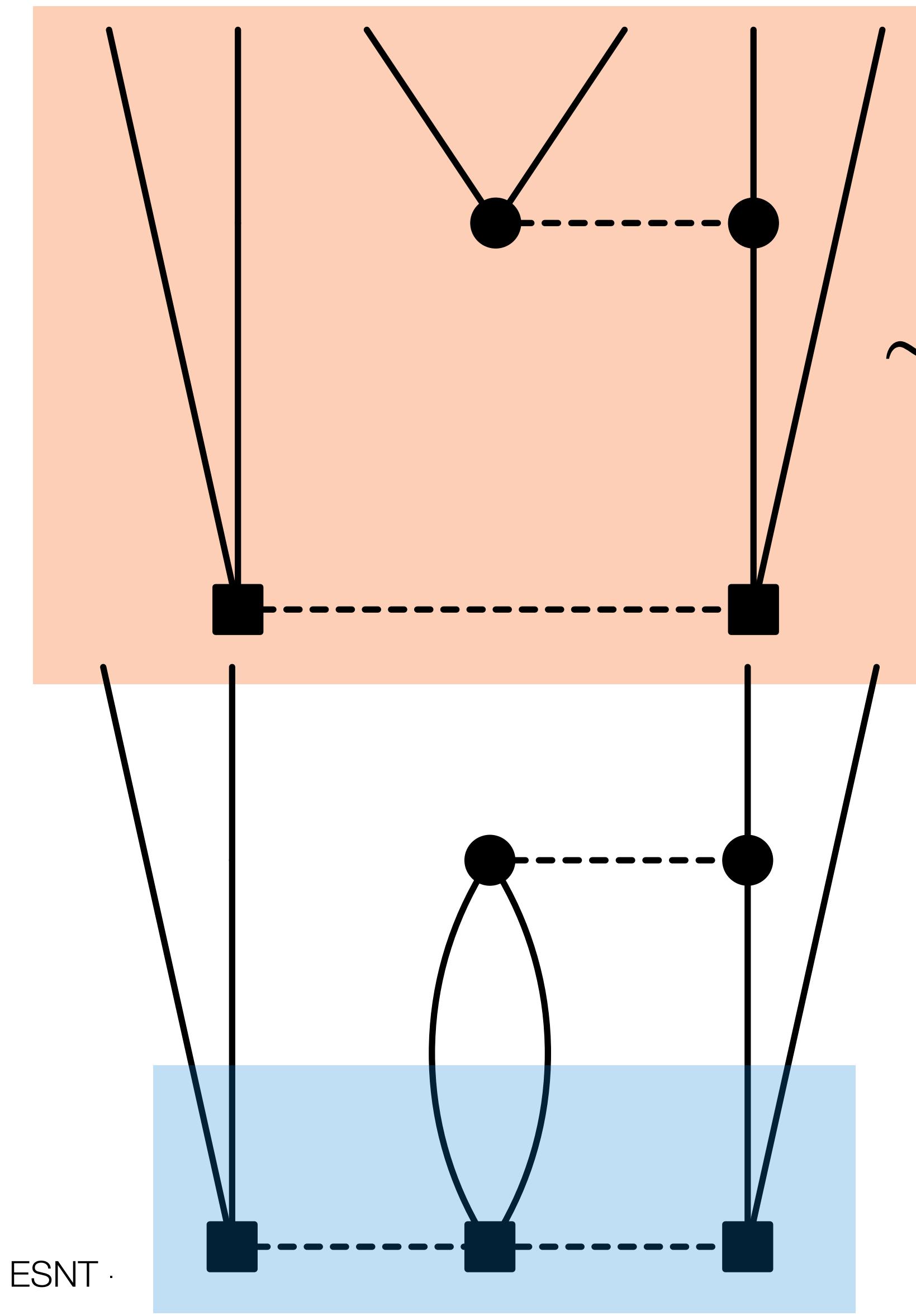
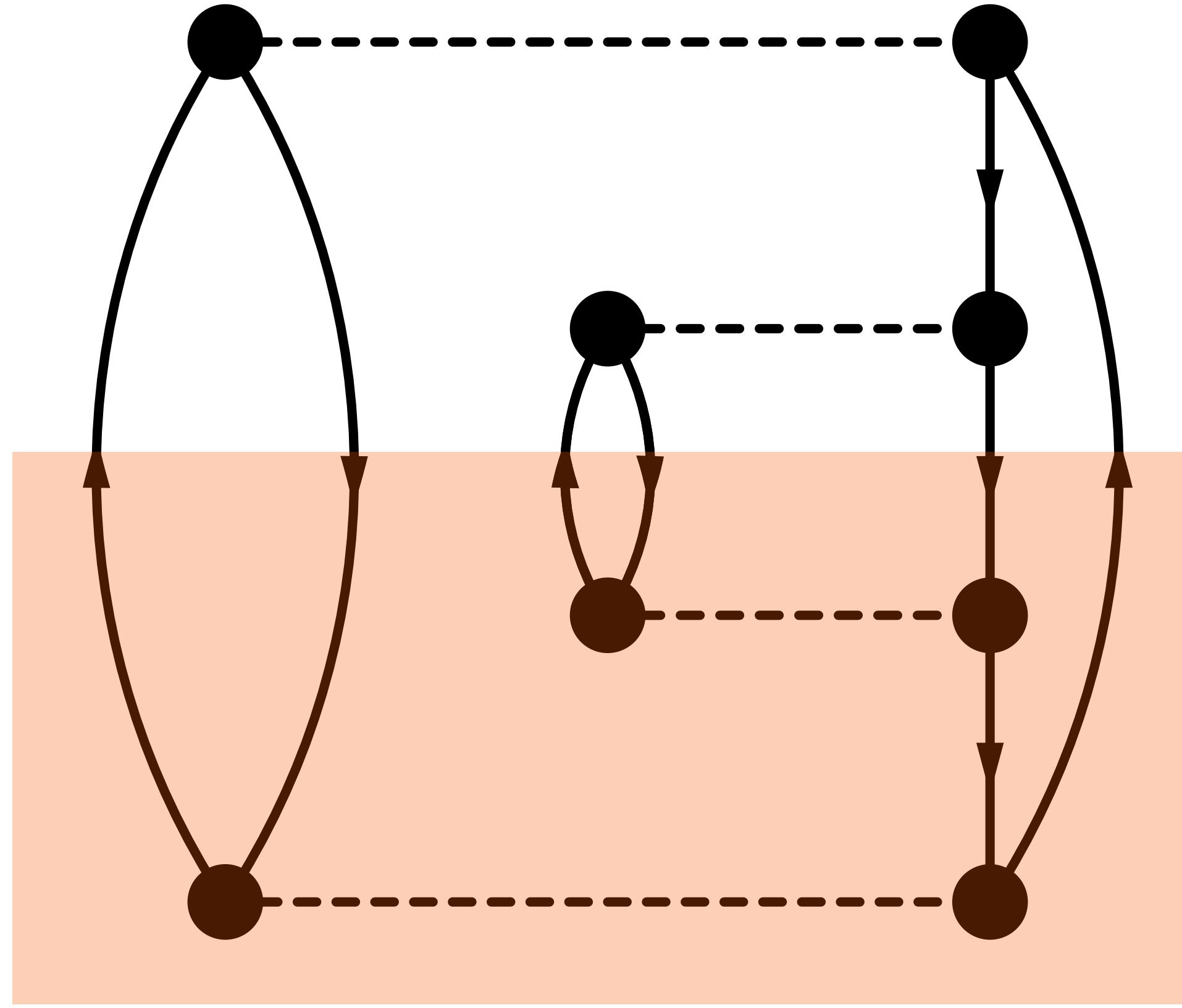
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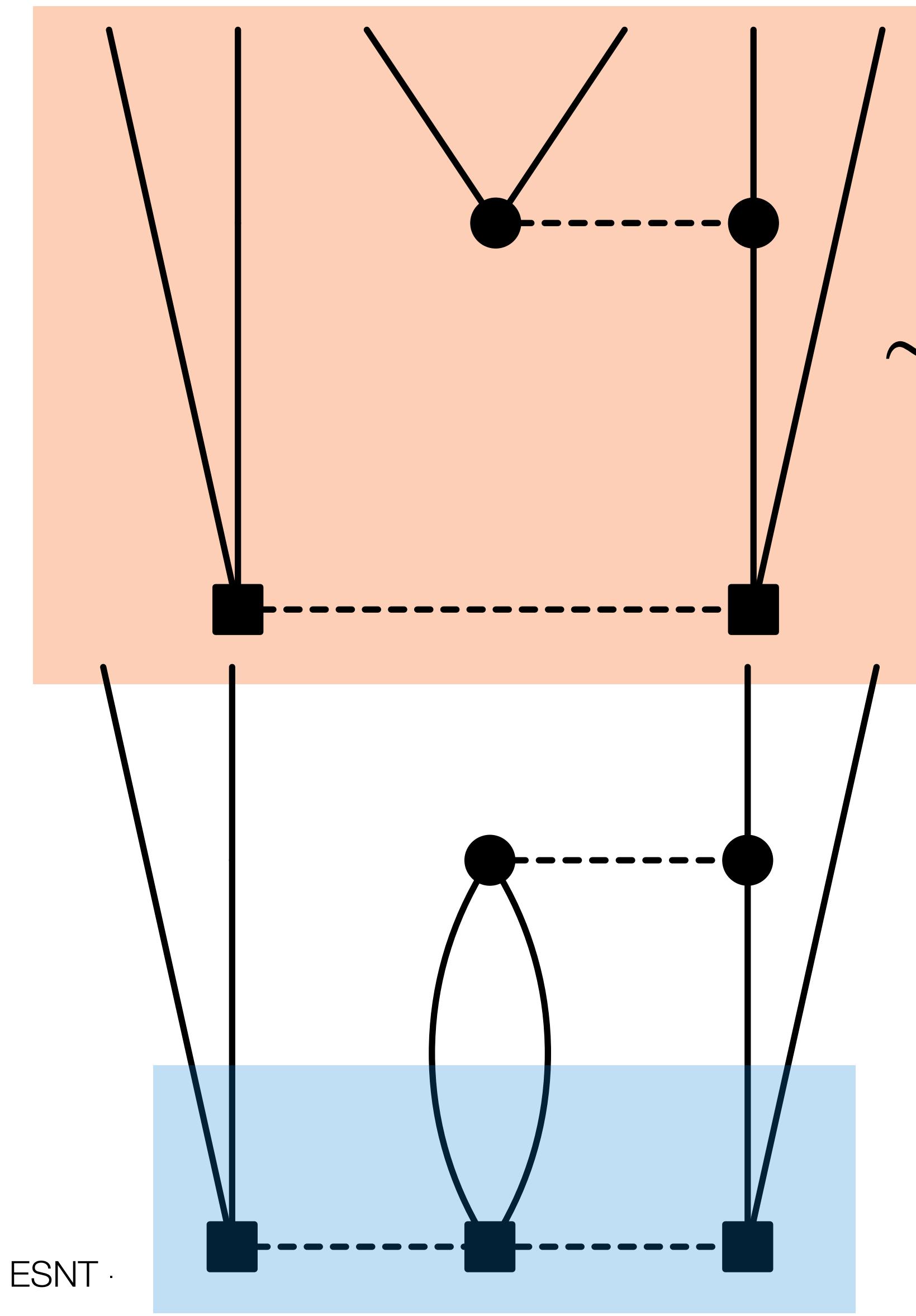
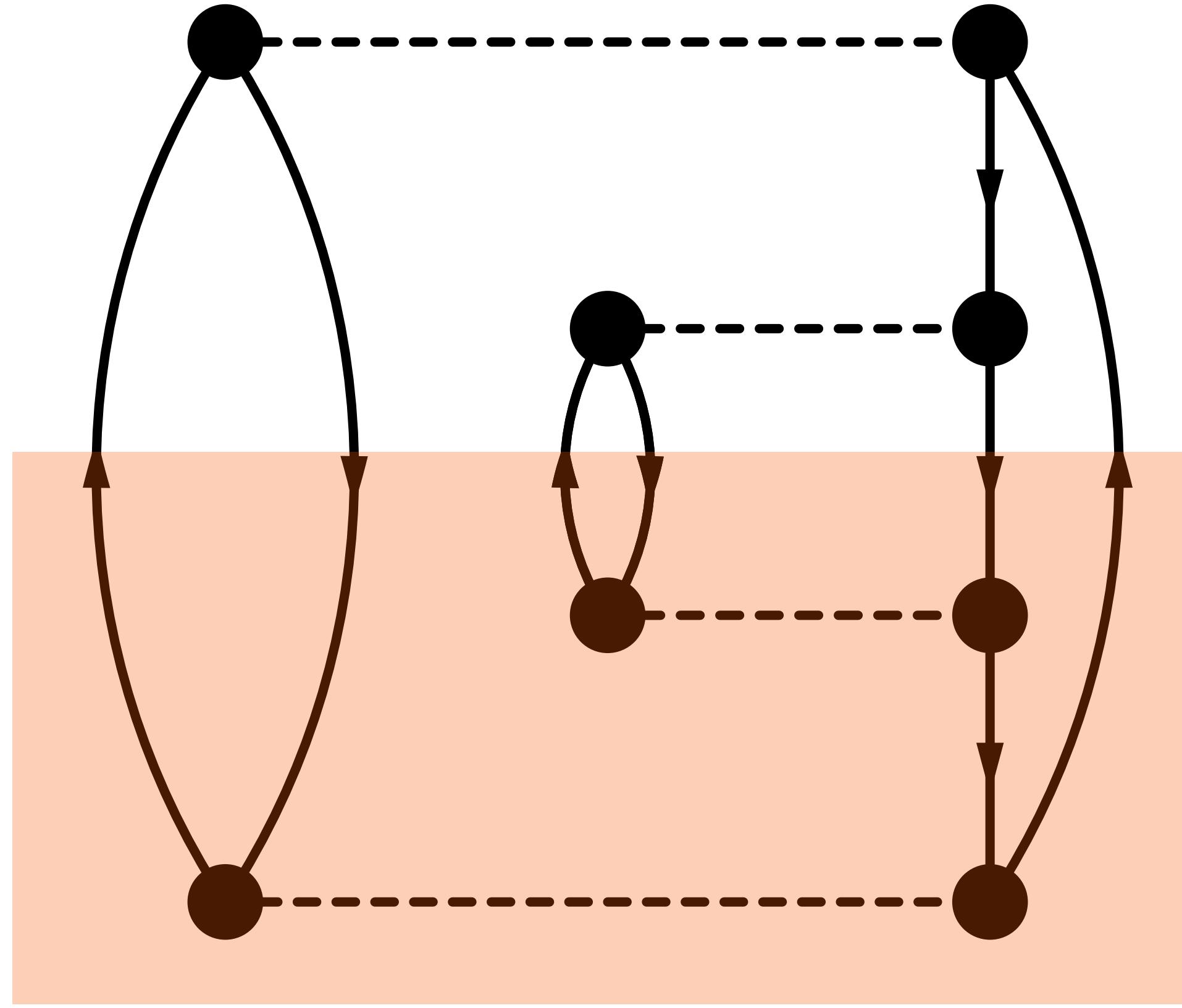
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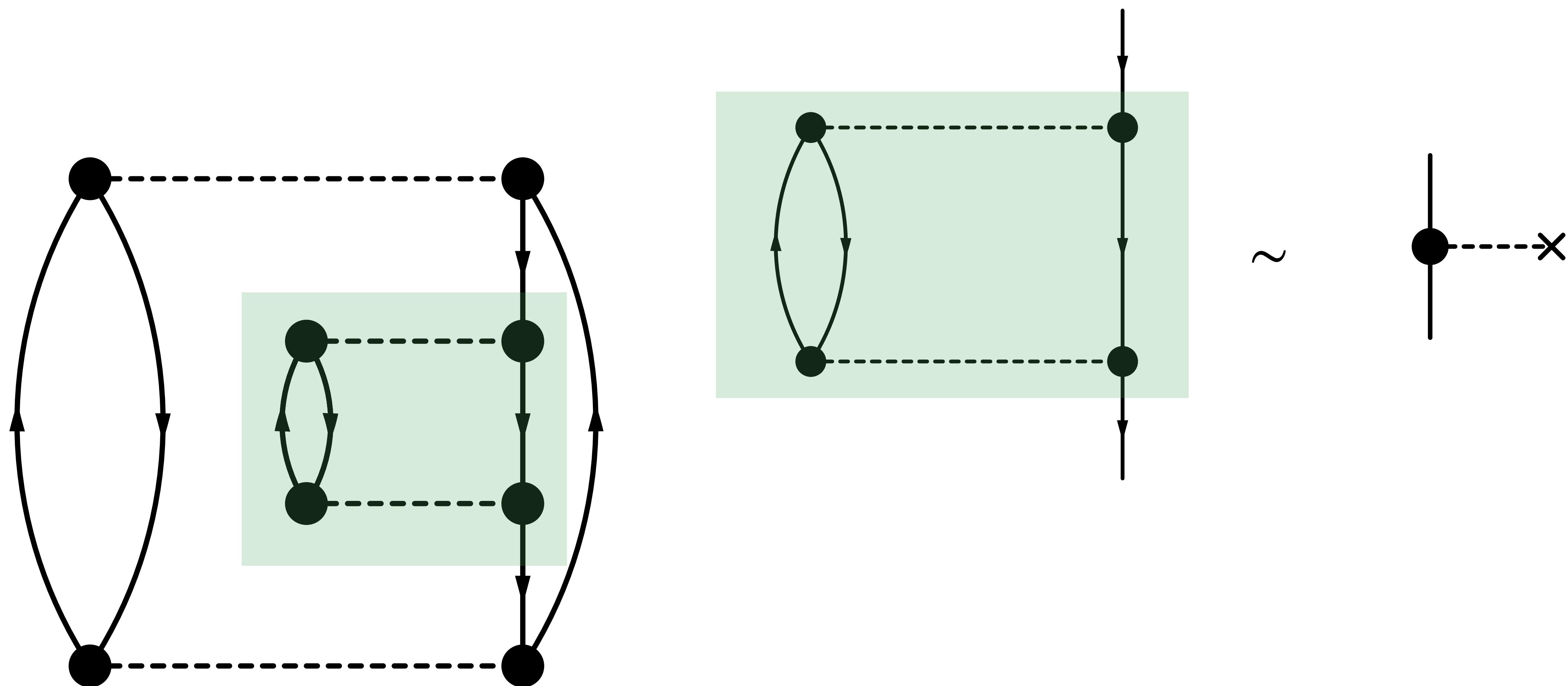
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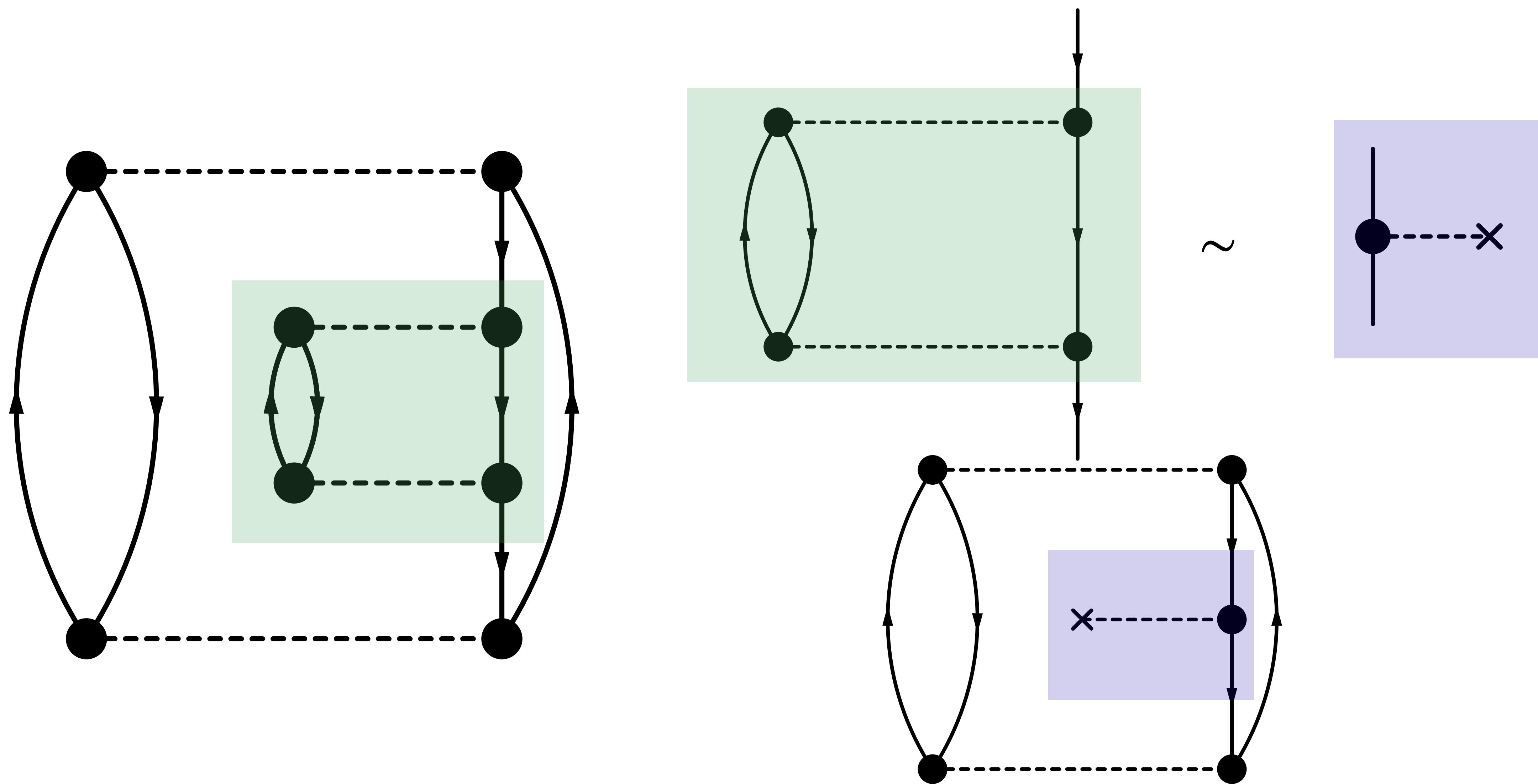
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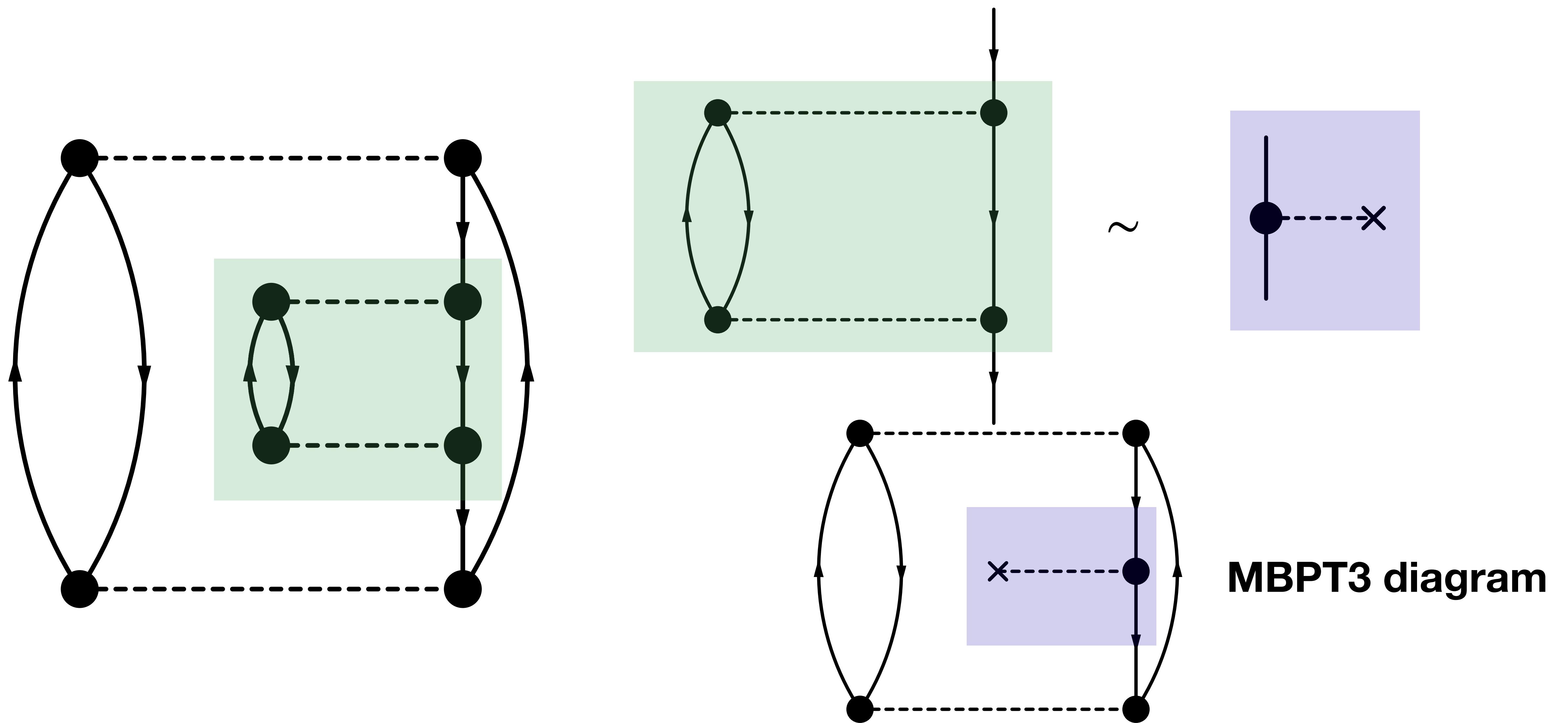
# Alternative approach: effective interactions



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# Alternative approach: effective interactions

Hjorth-Jensen et al., Phys. Rep. 261 (1995)

The diagram illustrates the equivalence of different Feynman diagrams. It consists of two parts separated by an equals sign (=). The left part shows a vertex with a solid line and a dashed line ending in an 'X'. The right part is divided into two sections by another equals sign. The first section shows a loop diagram with a dashed horizontal line connecting two vertices, each with a solid vertical line. The second section shows a vertical chain of four vertices connected by solid vertical lines, with arrows indicating flow from bottom to top. This entire right-hand side is equated to a mathematical expression involving summation and fractions.

$$= \frac{1}{2} \sum_{ija} \frac{\Gamma_{jlai} \Gamma_{ajk}}{\varepsilon_k + \varepsilon_{ij}^a}$$

# What could possibly go wrong?

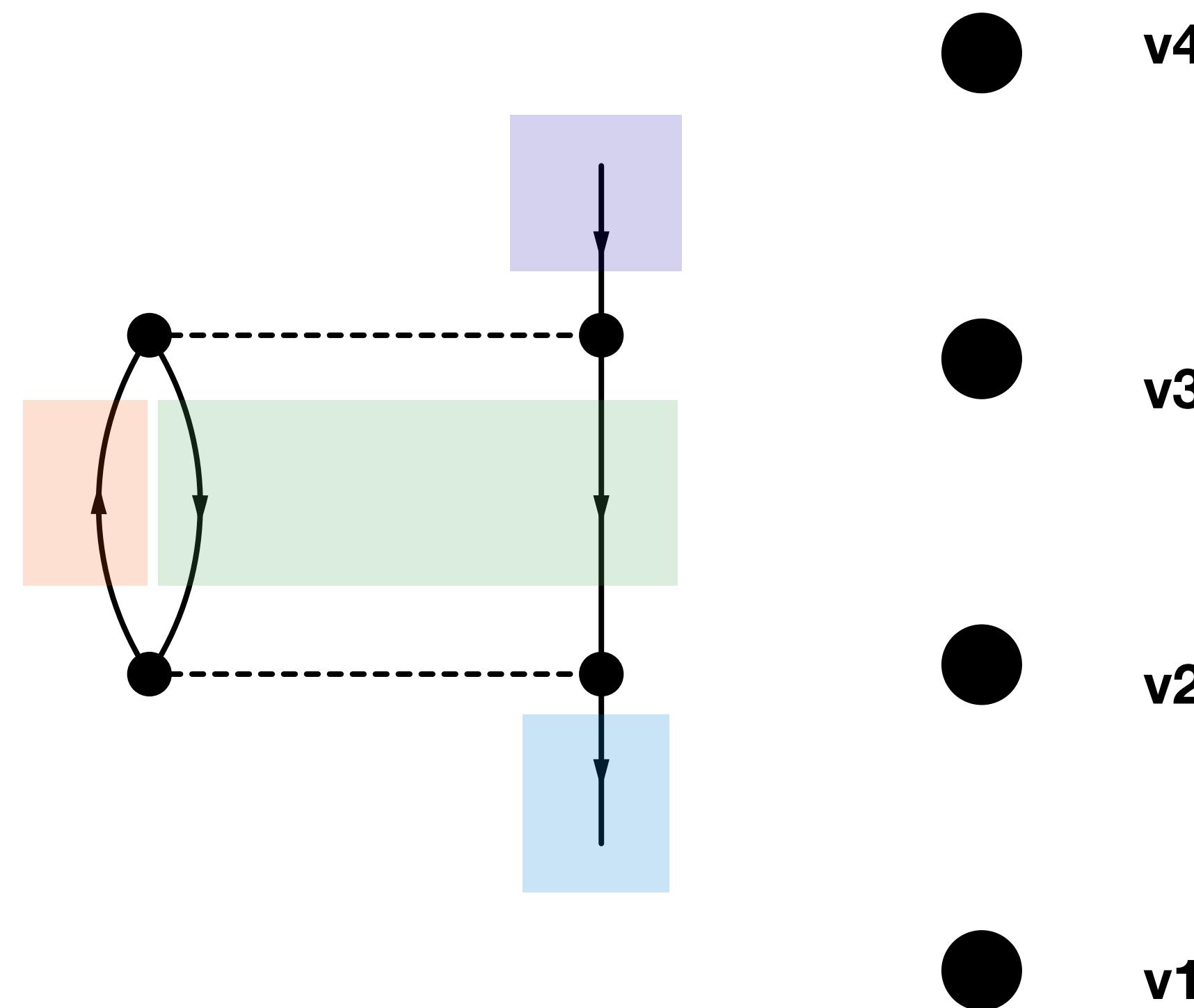
- Wrong energy denominators (unavoidable, but maybe good enough)
- **Wrong prefactor (overcounting)?**
- Fine-tuned solution, not extensible to higher IMSRG truncations

# High-order diagrams with ADG

Arthuis et al., CPC 240 (2019)

- Diagrams via adjacency matrices
- Search for submatrix corresponding to effective 1B interaction

$$\begin{pmatrix} & v_1 & v_2 & v_3 & v_4 \\ v_1 & 0 & 0 & 0 & X \\ v_2 & 1 & 0 & 1 & 0 \\ v_3 & 0 & 2 & 0 & 0 \\ v_4 & X & 0 & 1 & 0 \end{pmatrix}$$



# High-order diagrams with ADG

Arthuis et al., CPC 240 (2019)

Order	4	5	6
-------	---	---	---

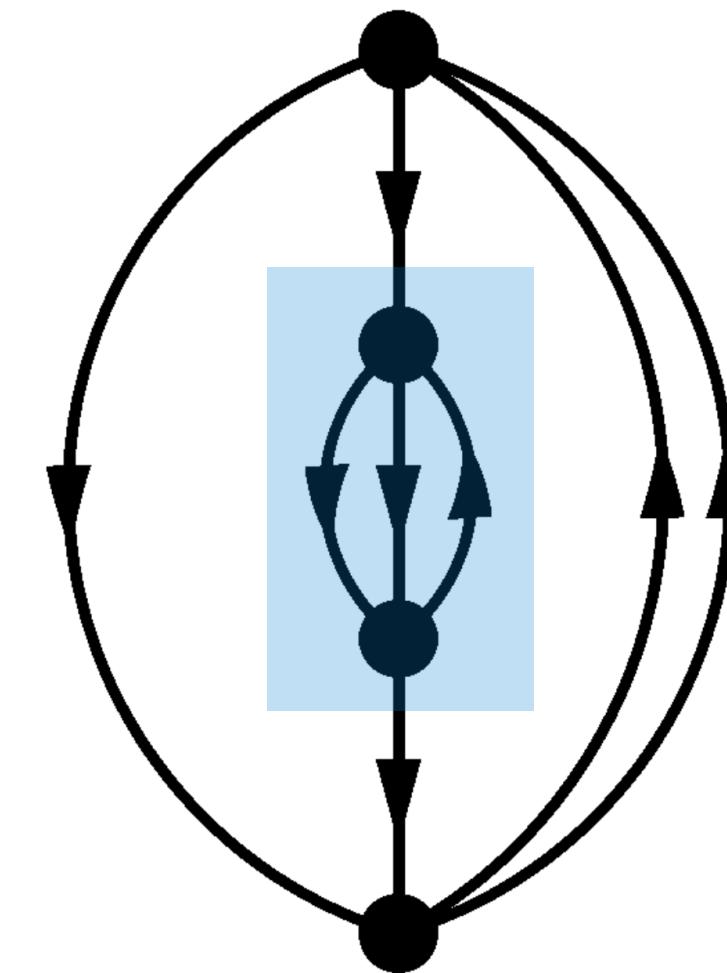
Num diags (one hit)	1	8	206
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Num diags (two hits)	2
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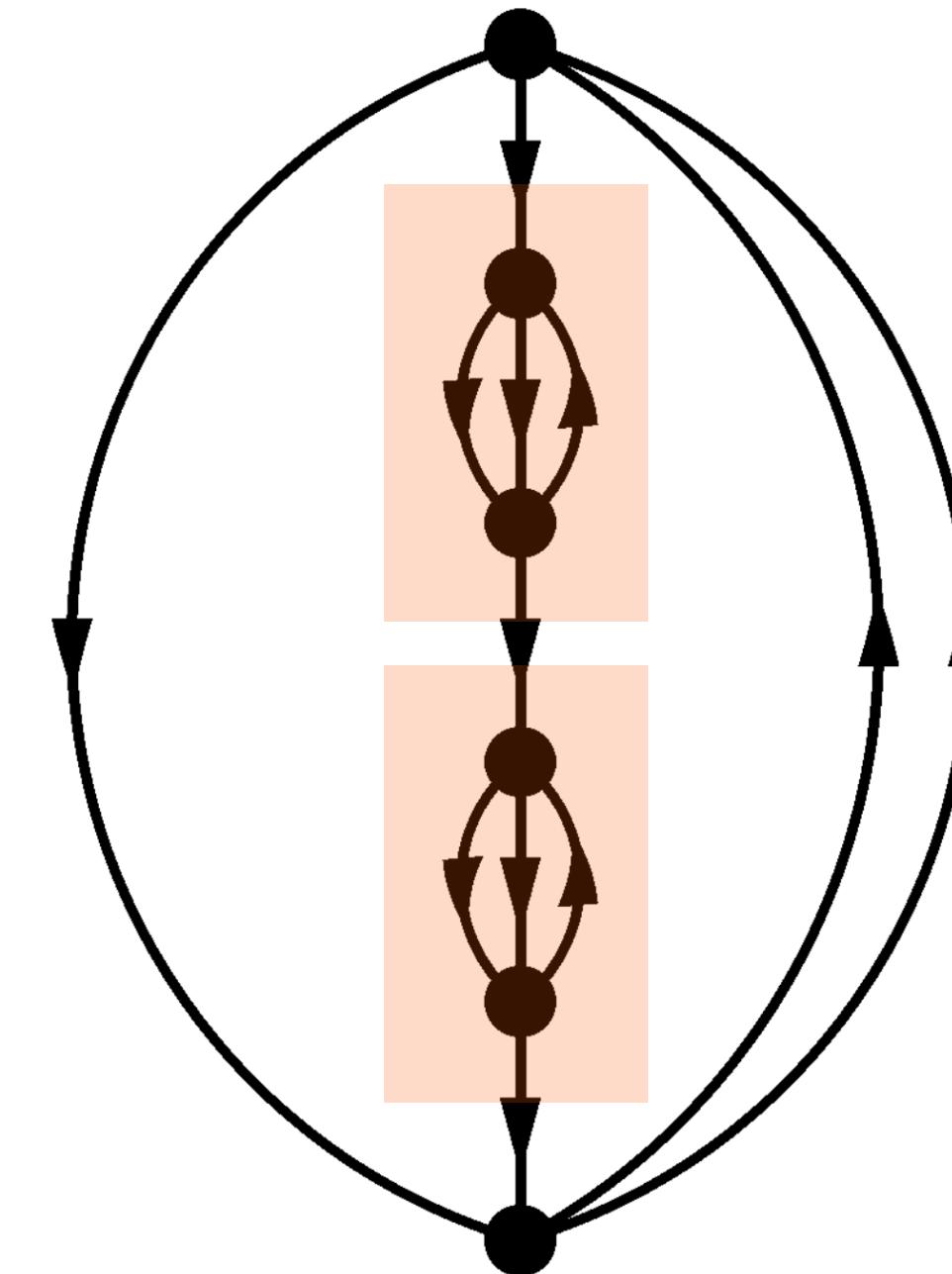


$$\frac{1}{(2!)^2} \sum \frac{v_{abnq} v_{ncop} v_{opcr} v_{qrab}}{\epsilon_{ab}^{nq} \epsilon_{cab}^{opq} \epsilon_{ab}^{qr}}$$

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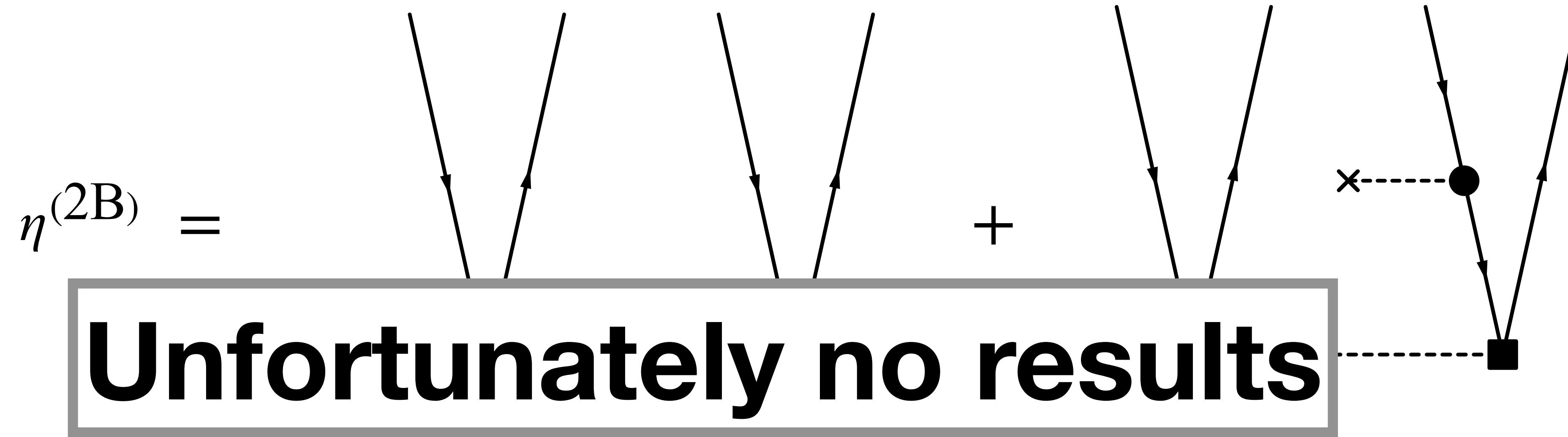
$$-\frac{1}{(2!)^3} \sum \frac{v_{abnt} v_{ncop} v_{opcq} v_{qdrs} v_{rsdu} v_{tuab}}{\epsilon_{ab}^{nt} \epsilon_{cab}^{opt} \epsilon_{ab}^{qt} \epsilon_{dab}^{rst} \epsilon_{ab}^{tu}}$$

# "Correcting" the IMSRG(2)

$$\eta^{(2B)} = \begin{array}{c} \text{Diagram 1: Two vertices connected by a dashed horizontal line, each with two outgoing lines forming a V-shape.} \\ + \\ \text{Diagram 2: Similar to Diagram 1, but the central dashed line has a cross symbol at the vertex where the two V-shapes meet.} \end{array}$$

- Improved generator has same decoupling condition  $\langle \Phi_{ij}^{ab} | H(\infty) | \Phi \rangle \stackrel{!}{=} 0$
- Gets important triples at MBPT4 approximately right
- But also higher order contributions due to nonperturbative power in IMSRG

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# Main takeaways

- IMSRG is a powerful many-body method for nuclear physics
- Perturbative improvements in IMSRG are complicated due to **ODE form**
- But they can give **insight into many-body error**
- Analysis connecting IMSRG and MBPT is complex
  - **Possibly automate perturbative analysis to get effective interactions?**
- Automated tools (ADG) can **support efforts** to validate improvements
- **TODO: Automated code generation using generic contraction engine**

# Acknowledgments



**Studienstiftung**  
des deutschen Volkes

Thanks to:

- **Jan Hoppe, Pierre Arthuis, Takayuki Miyagi, Alex Tichai**, Ragnar Stroberg, Kai Hebeler, Achim Schwenk
- TU Darmstadt "STRONGINT" group
- ORNL Nuclear Theory and **Titus Morris**
- **... and all of you for your attention**



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# Backup

# Details for matrix elements

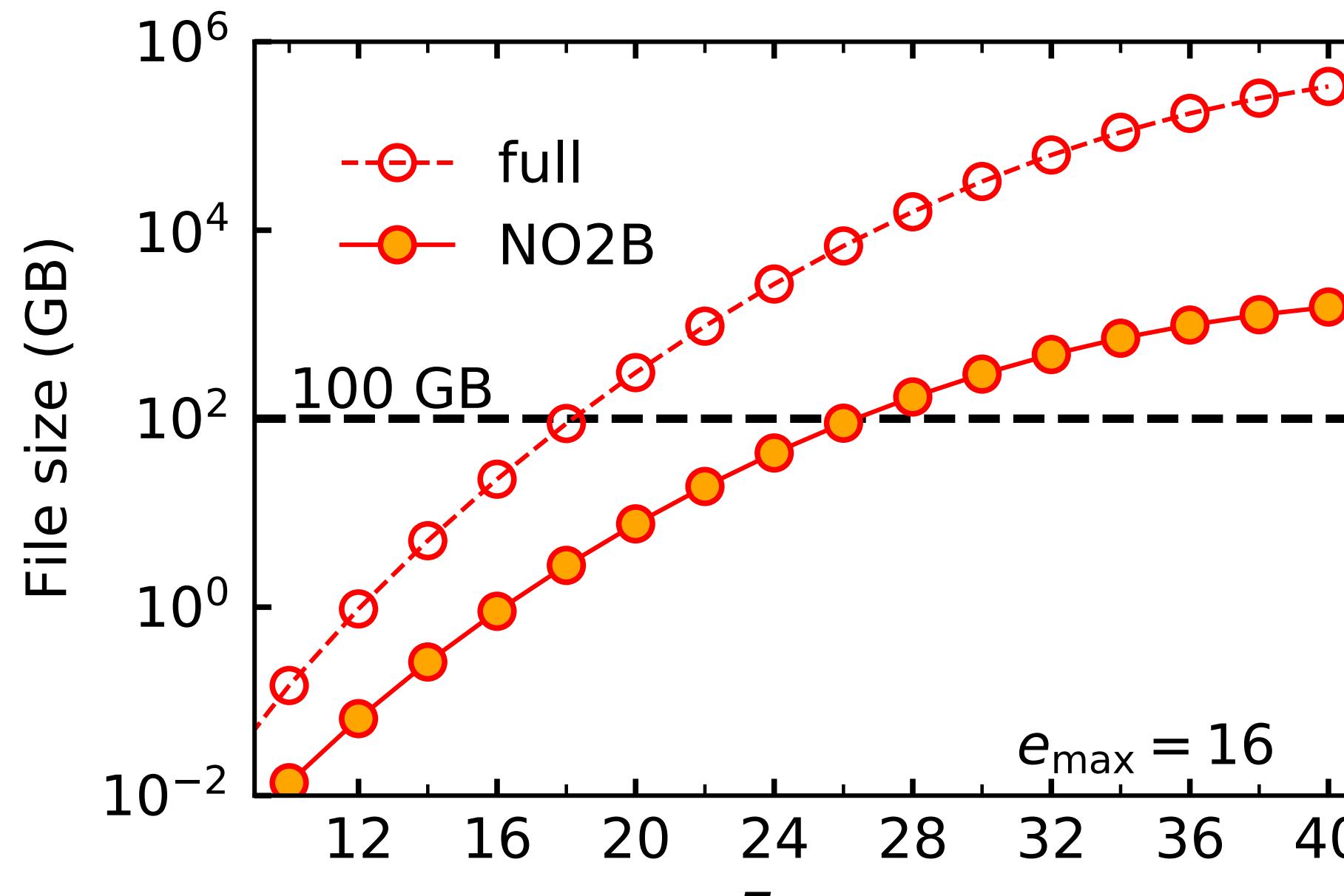
Coupled

- 2B basis:  $| (pq)JM_J \rangle$ , ignore trivial  $M_J$  dependence,  
block sparse in ( $\Pi = \Pi_p + \Pi_q, T_z = t_{z,p} + t_{z,q}, J$ )
- 3B basis:  $| [(pq)J_{pq}r]JM_J \rangle$ , ignore trivial  $M_J$  dependence,  
block sparse in ( $\Pi = \Pi_p + \Pi_q + \Pi_r, T_z = t_{z,p} + t_{z,q} + t_{z,r}, J$ )

Uncoupled

- 2B basis:  $| pq \rangle$ , block sparse in ( $\Pi, T_z, M_J = m_p + m_q$ )
- 3B basis:  $| pqr \rangle$ , block sparse in ( $\Pi, T_z, M_J = m_p + m_q + m_r$ )
- **Typically partition 3B basis blocks into 2B symmetry subblocks**

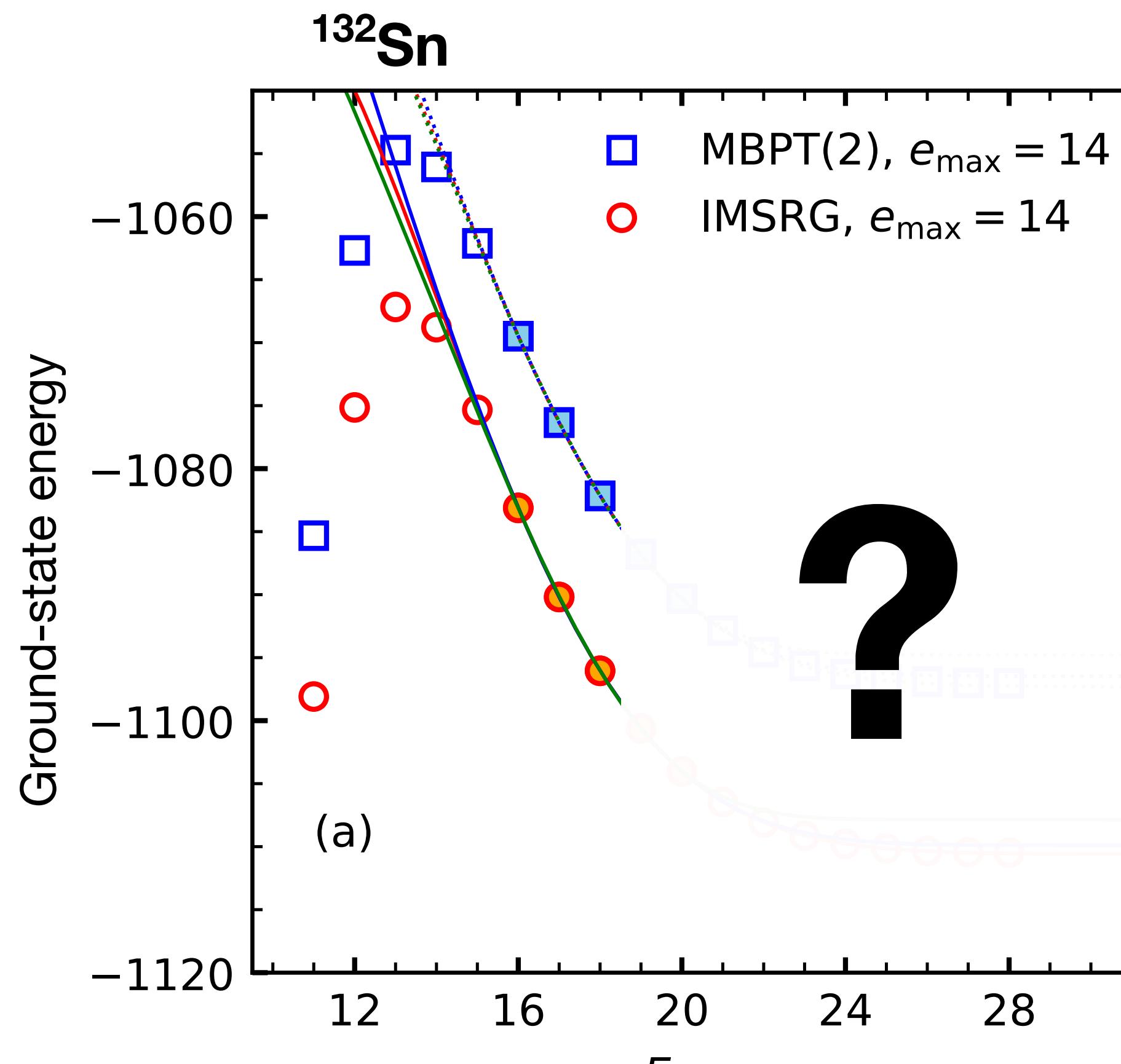
# Converged calculations for heavy nuclei



Miyagi et al., PRC 105 (2022)

- Ab initio calculations of heavy nuclei constrained by three-body force convergence ( $E_{3,\max}$ )
- Novel storage scheme reduces storage costs by 2-3 orders of magnitude
- First converged calculations of Sn-132
- Opens up frontier of heavy nuclei

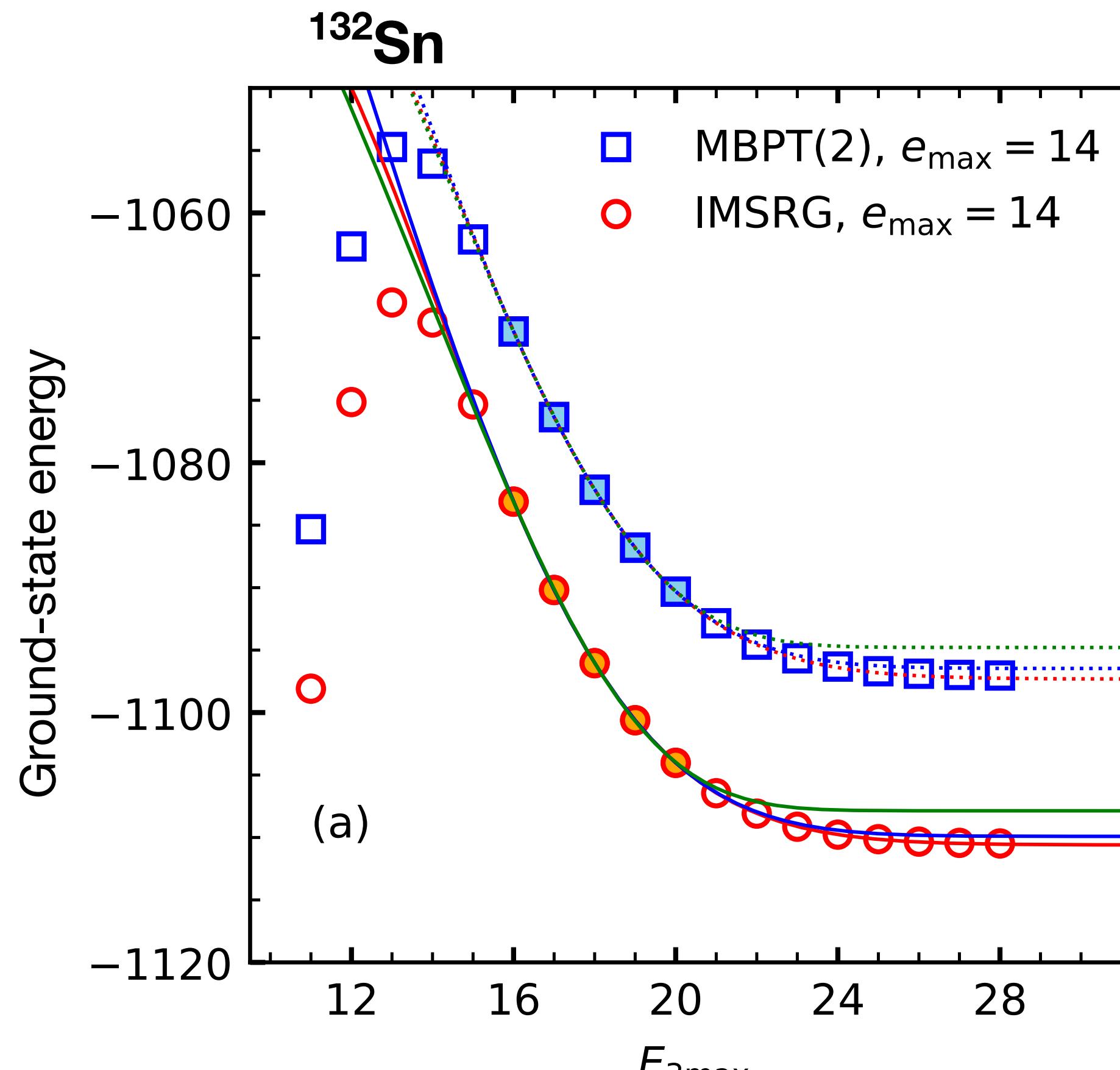
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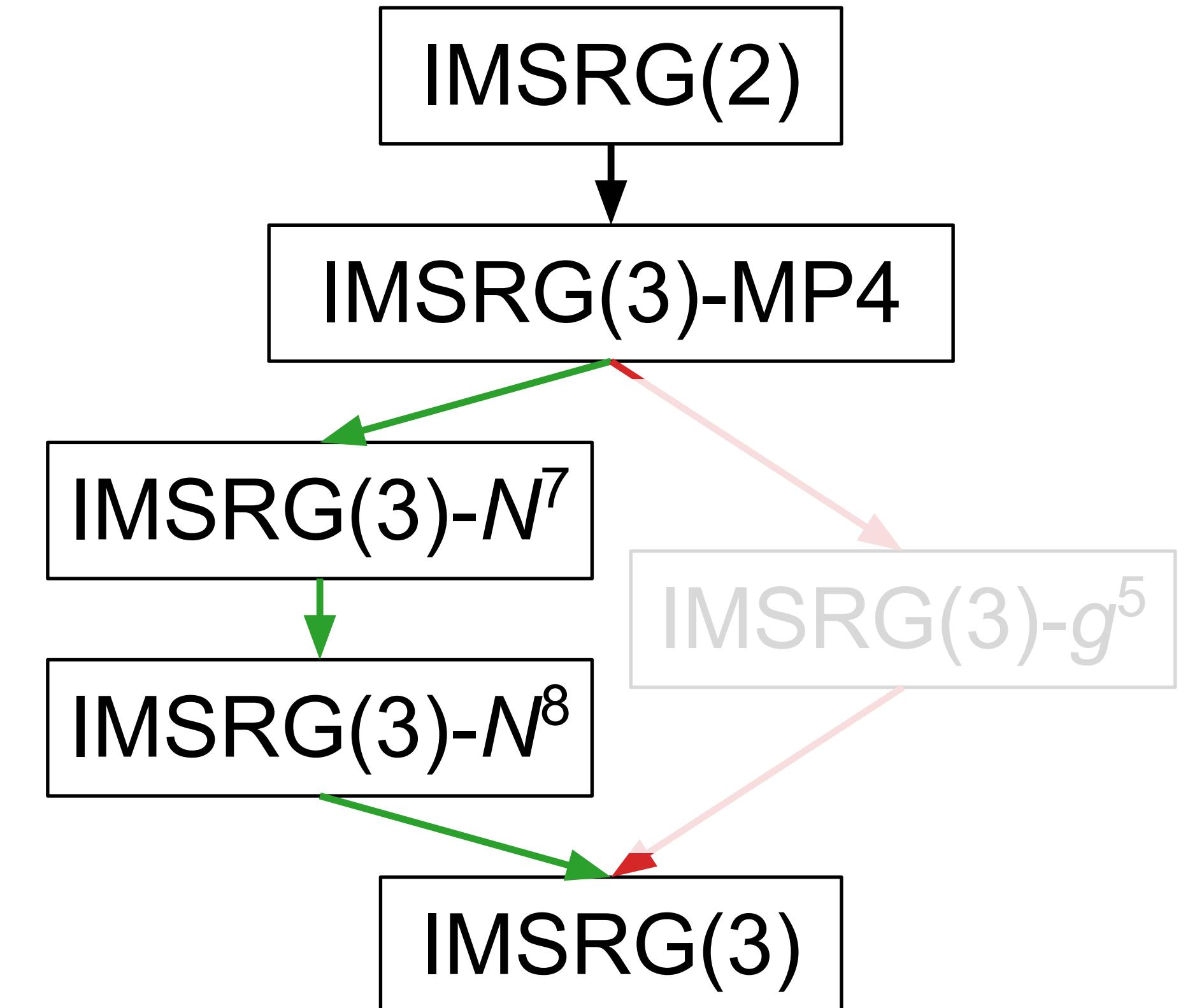
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# Unlocking IMSRG(3) (approximately)

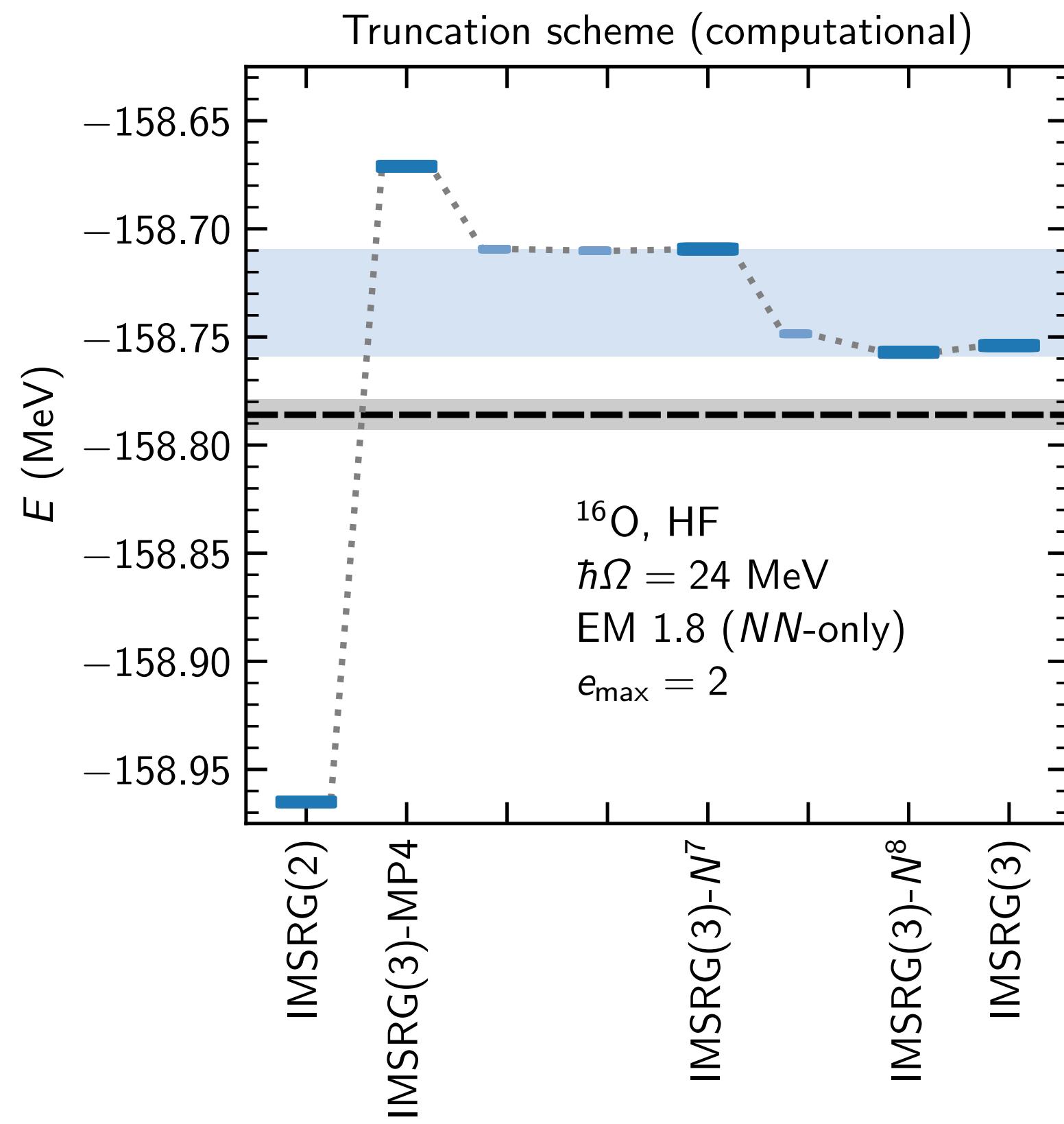
Heinz et al., PRC 103 (2021)

- IMSRG(3) is too expensive
- Systematic approximation based on **computational cost**
- Study in restricted setting to understand many-body convergence



# Unlocking IMSRG(3) (approximately)

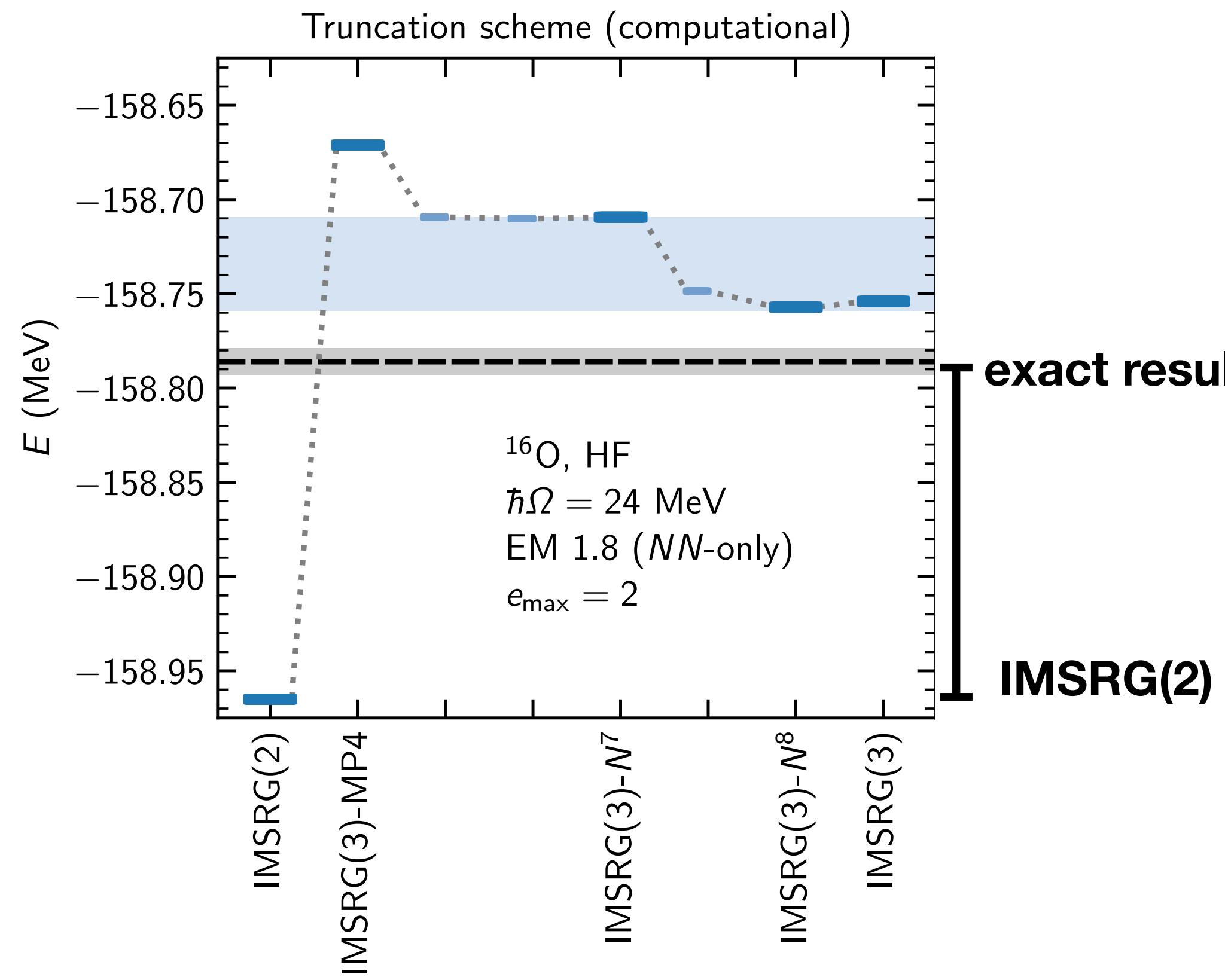
Heinz et al., PRC 103 (2021)



- Systematic convergence to exact result
- IMSRG(2) does very well ...
- ... but IMSRG(3) (and approximations) perform even better
- Benefits of IMSRG(3) largely present in approximations (e.g., IMSRG(3)- $N^7$ )

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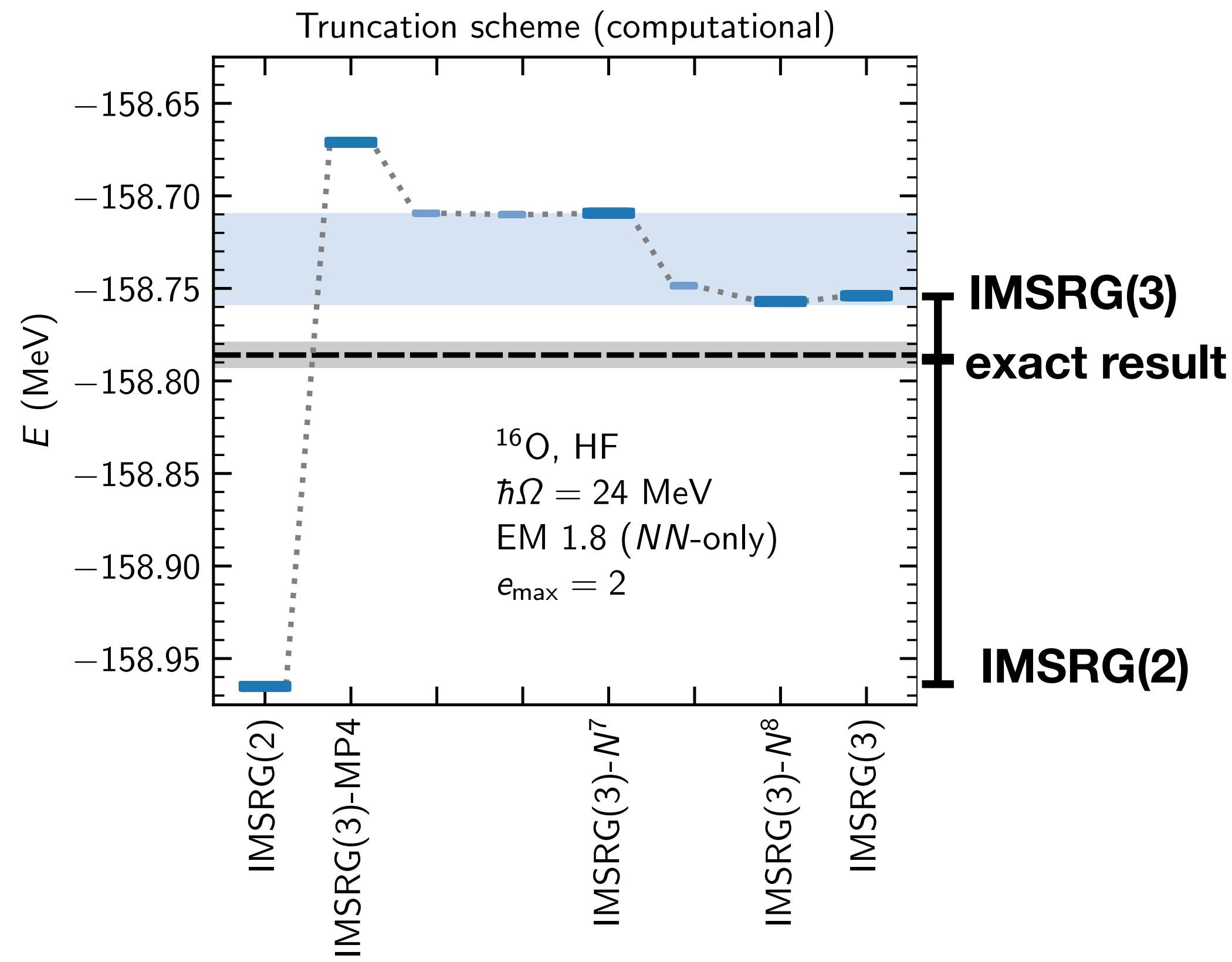
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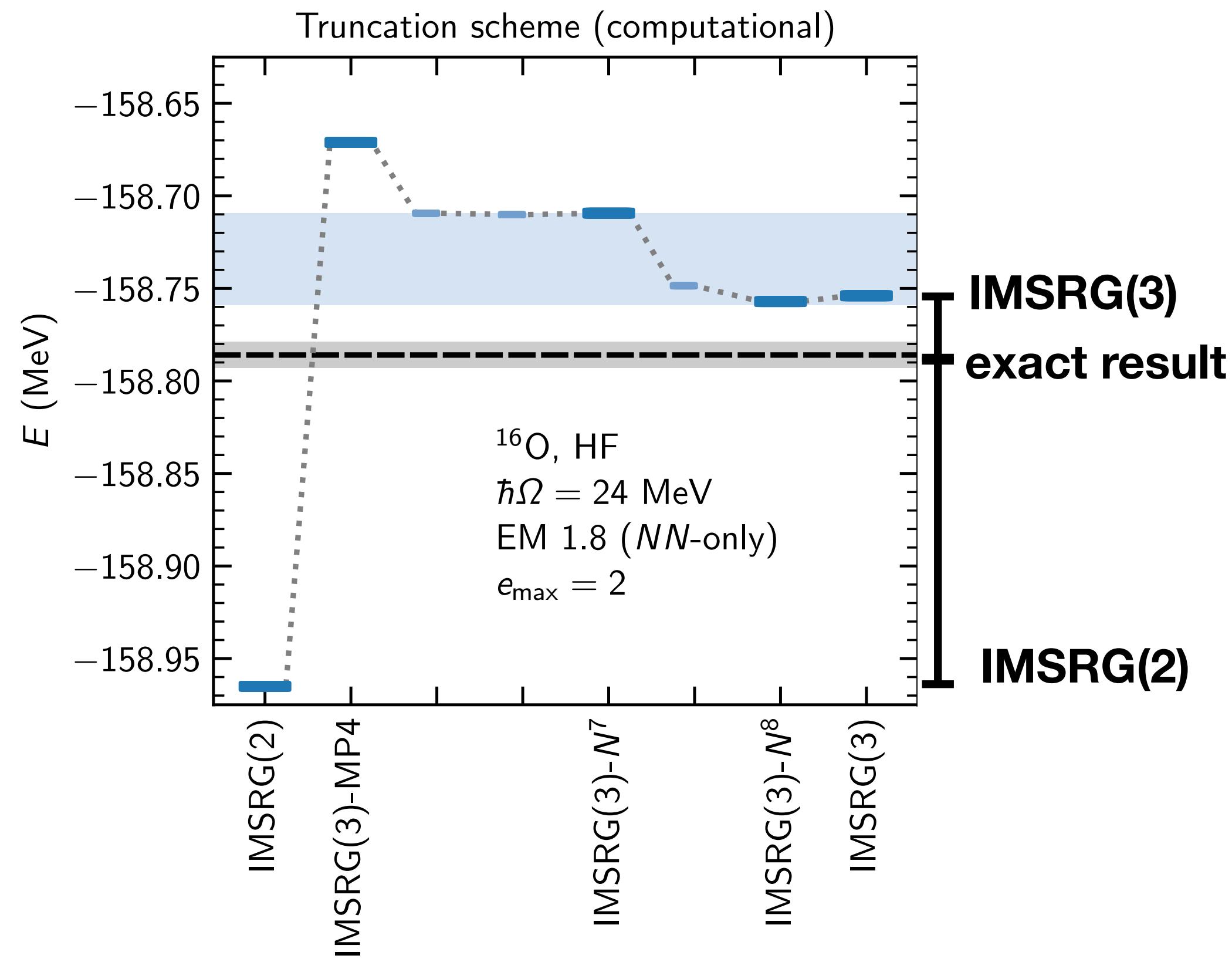
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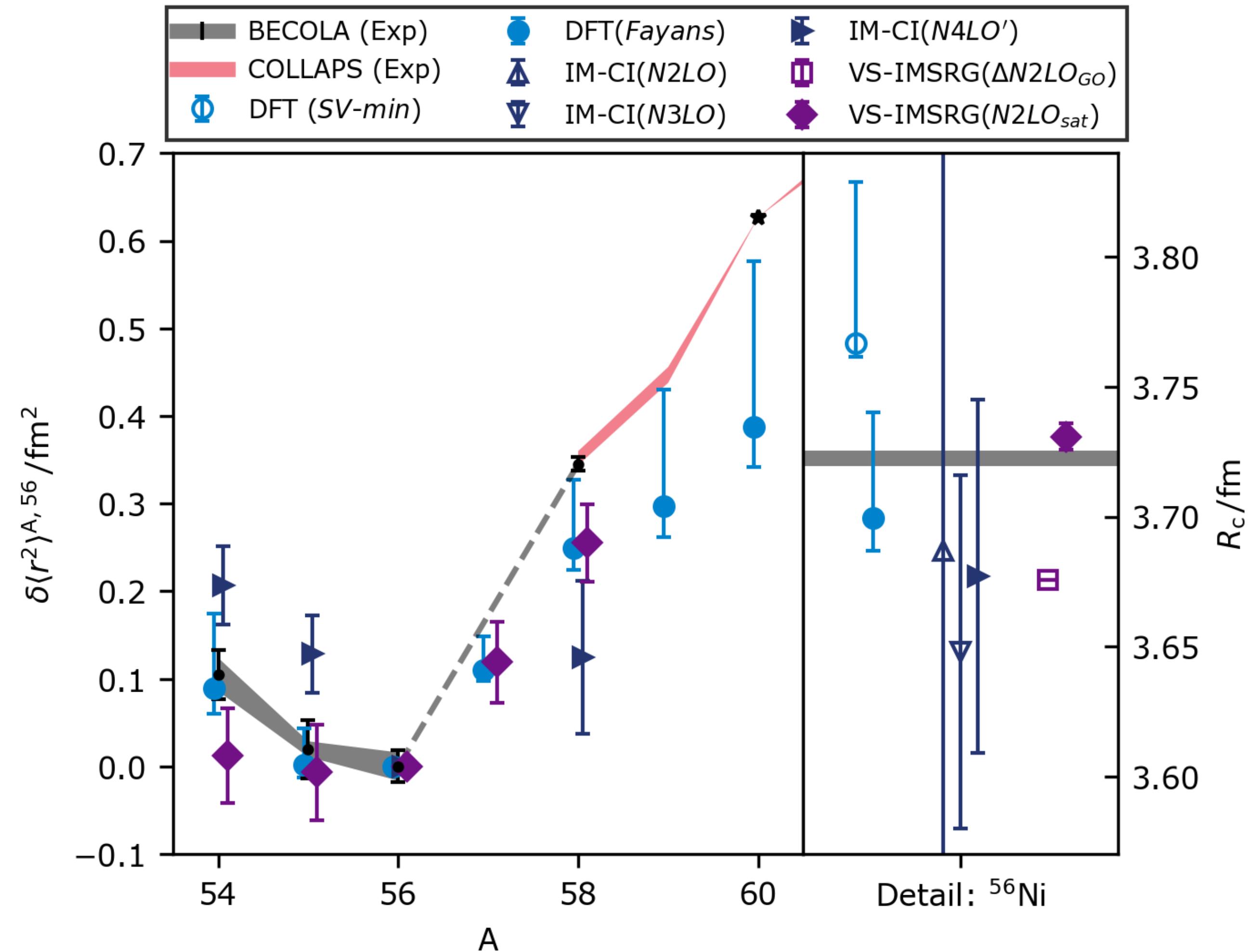
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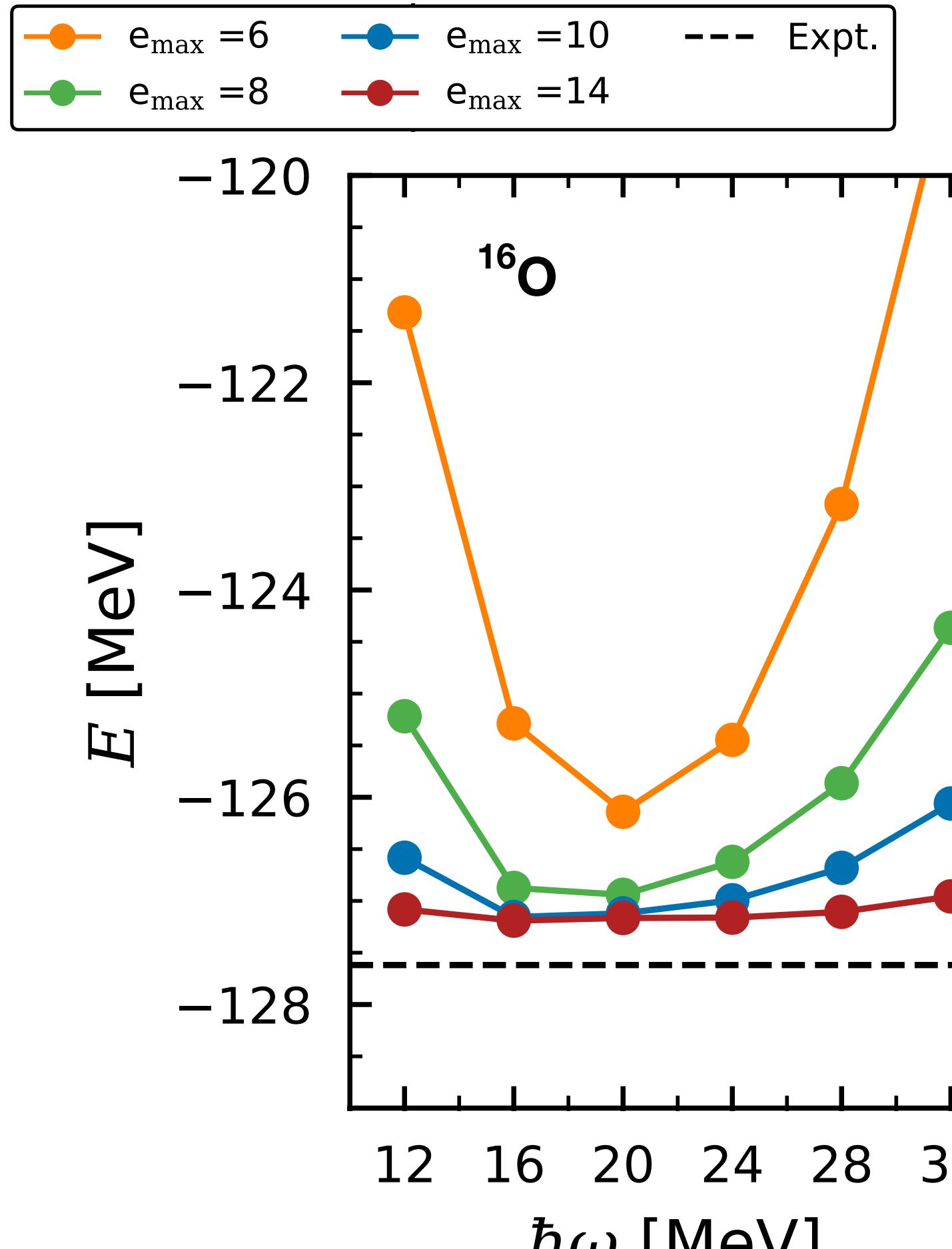
**Goal: Use IMSRG(3) approximations to improve many-body calculations and quantify uncertainties**

# Nickel charge radii



Sommer et al., PRL 129 (2022)

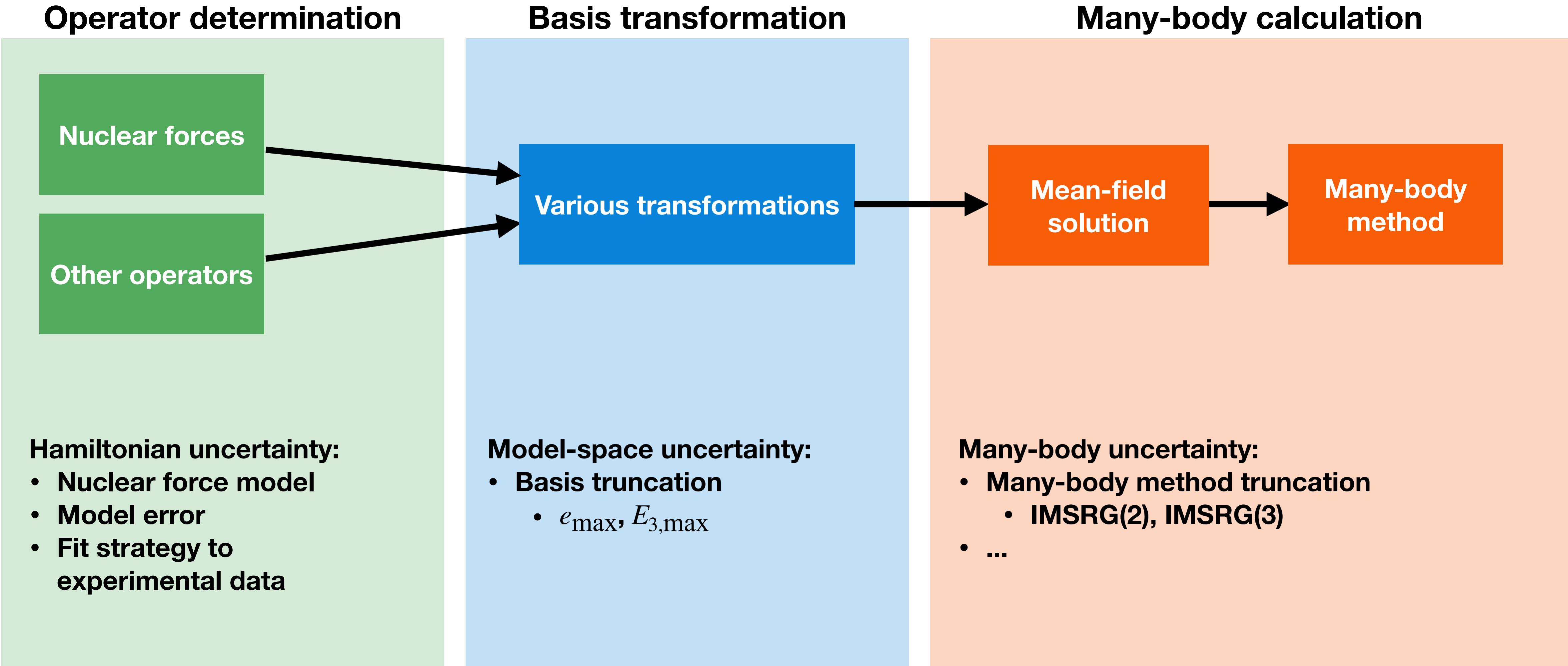
# IMSRG ingredients



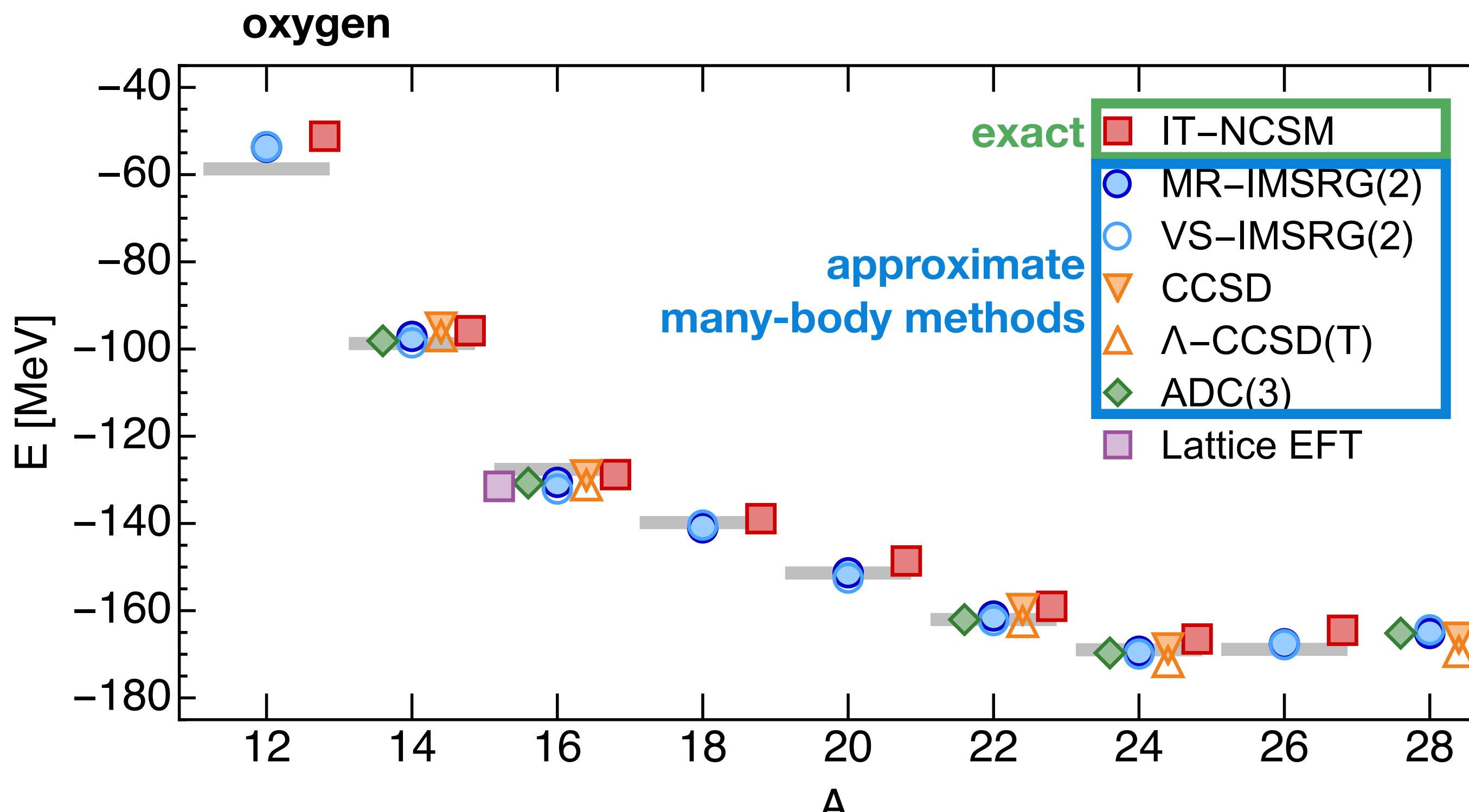
Hoppe, ..., MH et al., PRC 103 (2021)

1. Input Hamiltonian  $H$
2. Solve for mean field (Hartree-Fock)
  - Input dependence:  $H, e_{\max}, E_{3\max}, \hbar\omega$
  - Output: reference state  $|\Phi\rangle$ , basis  $\{\phi_p\}$
3. Solve for many-body correlations [IMSRG(2)]
  - Input dependence:  $\boxed{H, |\Phi\rangle, \{\phi_p\}}$ , other ops ...
  - normal ordering
  - Hebeler, ..., MH et al., PRC 107 (2023)
- Output:  $|\Phi\rangle, E$ , expectation values of ops ...

# Many-body calculation uncertainties



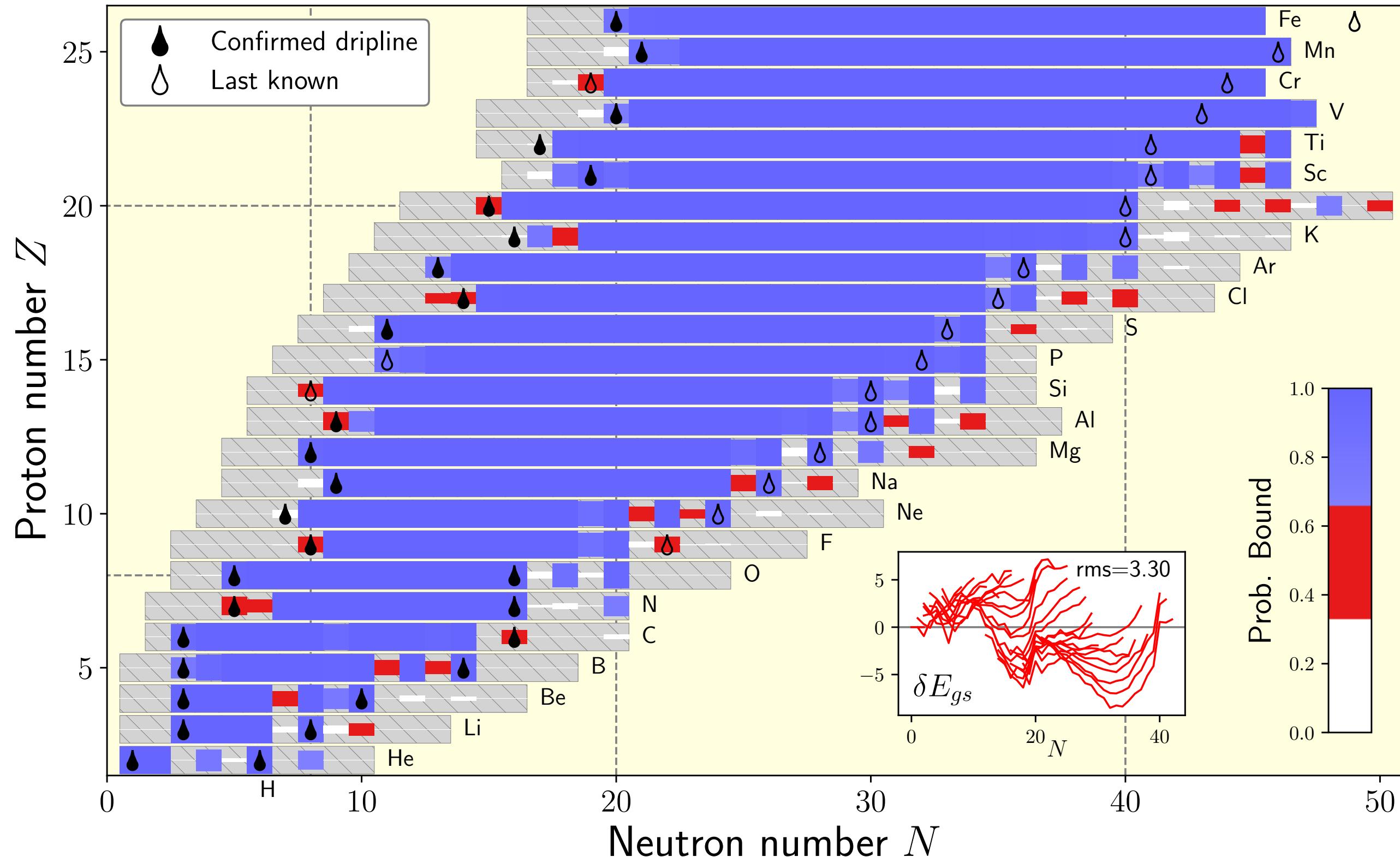
# Many-body consistency in oxygen



Hergert, Front. Phys. 7 (2020)

- For given Hamiltonian, many-body methods are very consistent
- 1-2% discrepancy to **exact result** due to **many-body approximation**

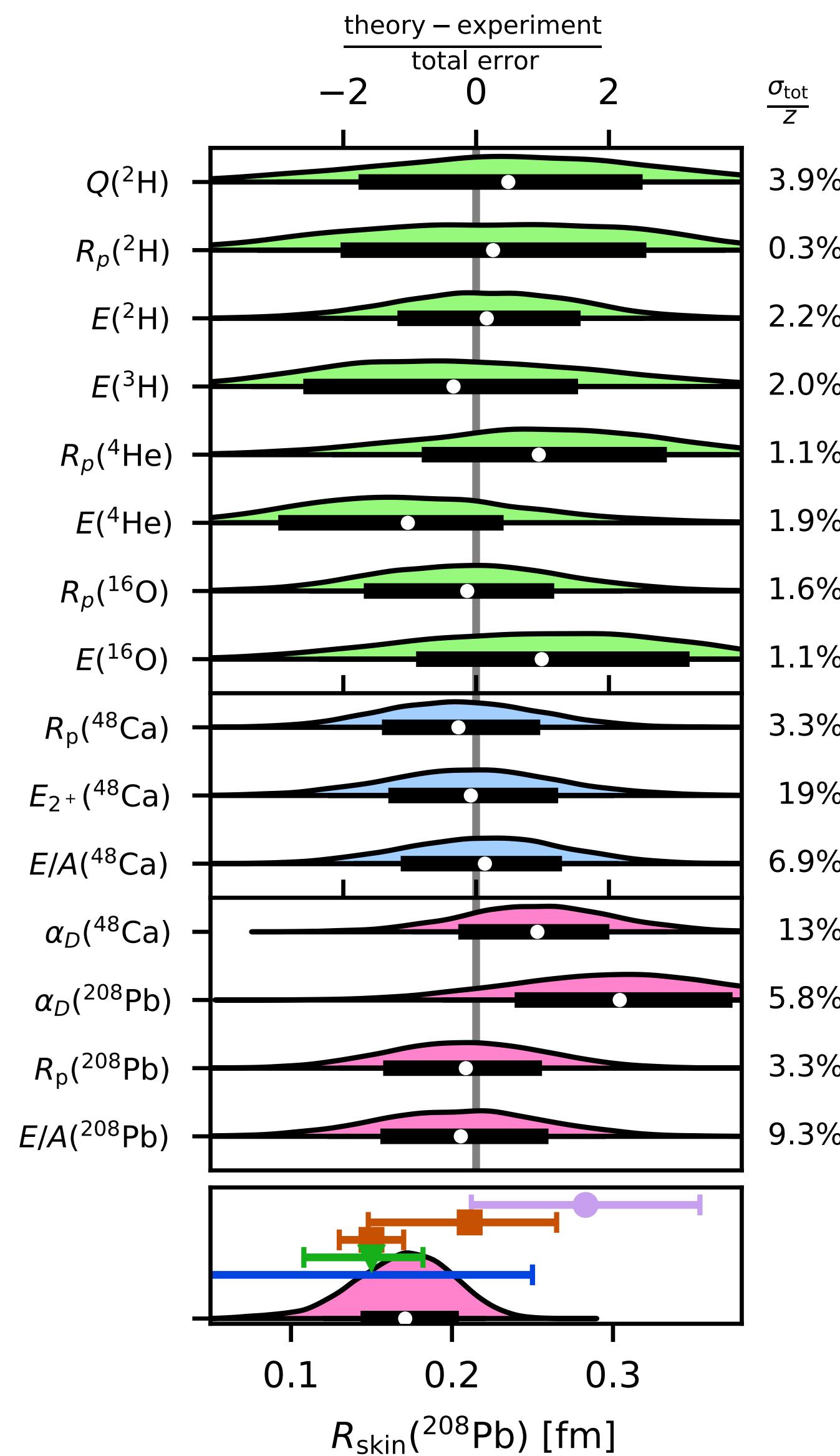
# Drip lines from the VS-IMSRG



agonalization  
ies based on single Hamiltonian  
ors to predict limits of stability

Stroberg et al., PRL 126 (2021)

# First systematic study of $^{208}\text{Pb}$



History matching:

- Start from  $10^9$  Hamiltonians
- Compare predictions with experiment and discard "implausible" Hamiltonians
- Final result: 34 different valid choices

Likelihood calibration:

- Assess quality of valid Hamiltonians in  $^{48}\text{Ca}$
- Assign importance weight based on quality

Uncertainty quantified prediction:

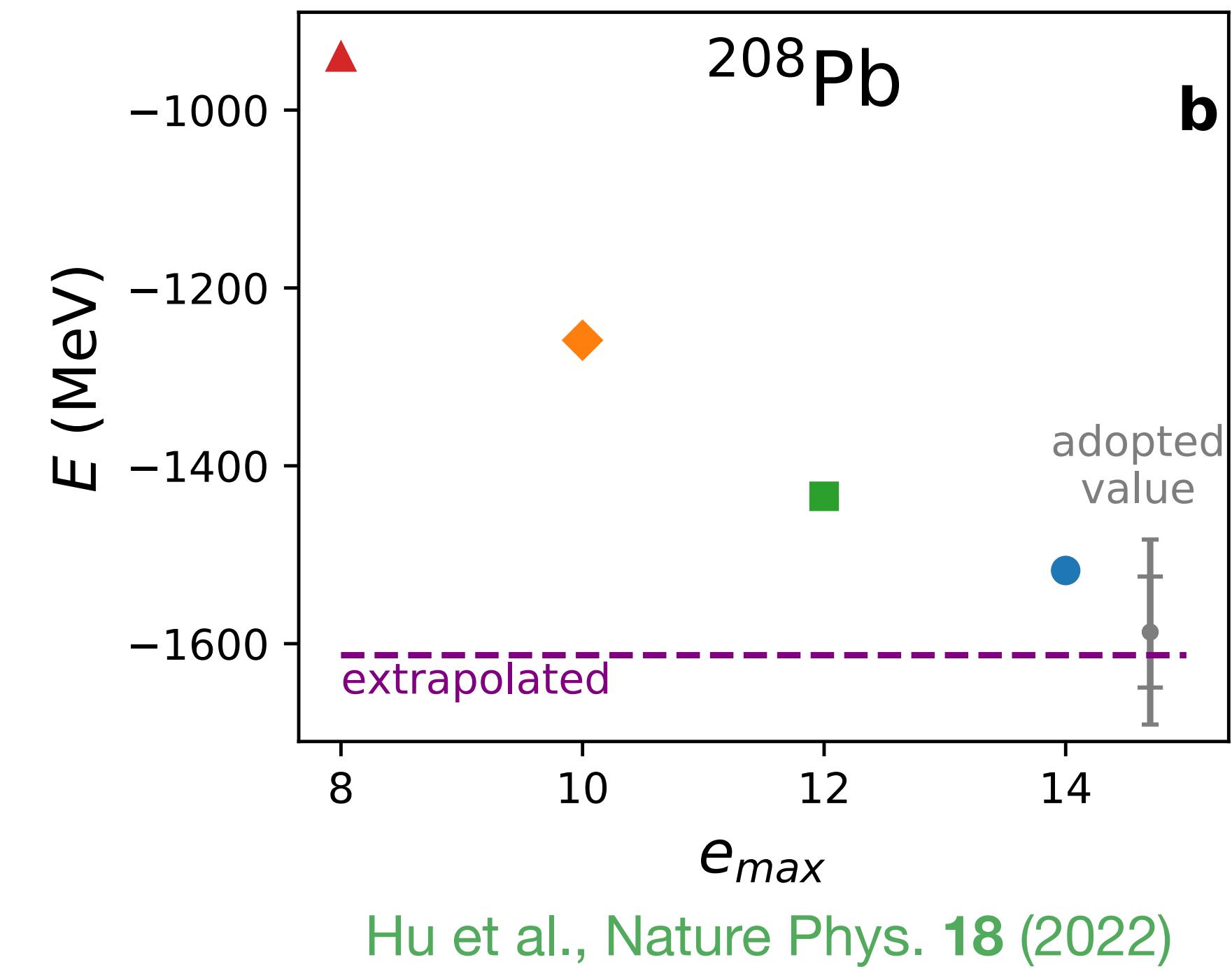
- Make many-body prediction using all valid Hamiltonians

$$R_{\text{skin}} = 0.14\text{--}0.20 \text{ fm}$$

**Hamiltonian, model space, and method uncertainties accounted for at all stages!**

# First systematic study of $^{208}\text{Pb}$

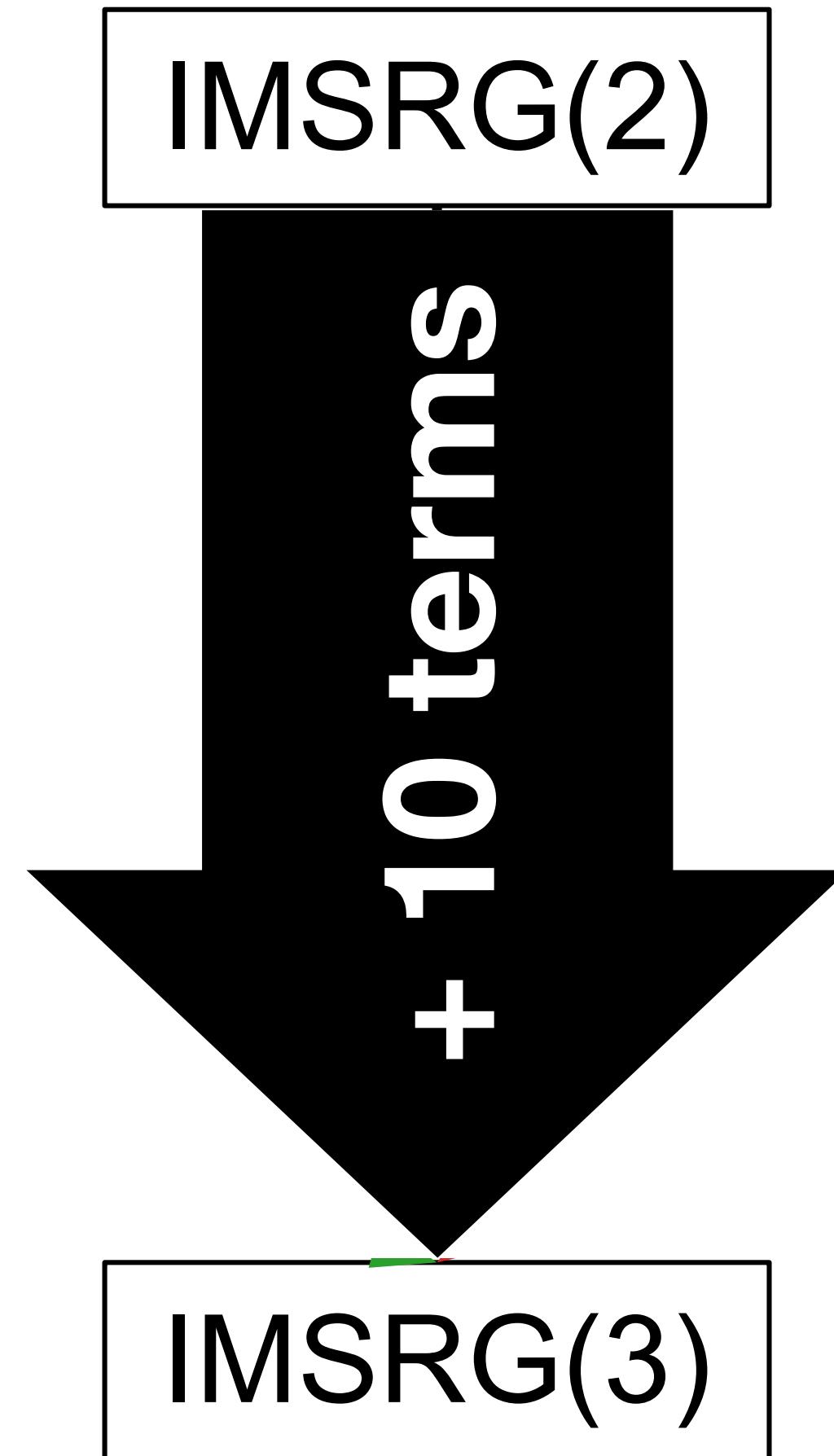
- Convergence of many-body calculation still challenging
- Hamiltonian uncertainties also very large
- Uncertainty quantification methods explicitly take errors into account
- Look at observables where correlated errors cancel?



# Reaching IMSRG(3)

MH et al., PRC 103 (2021)

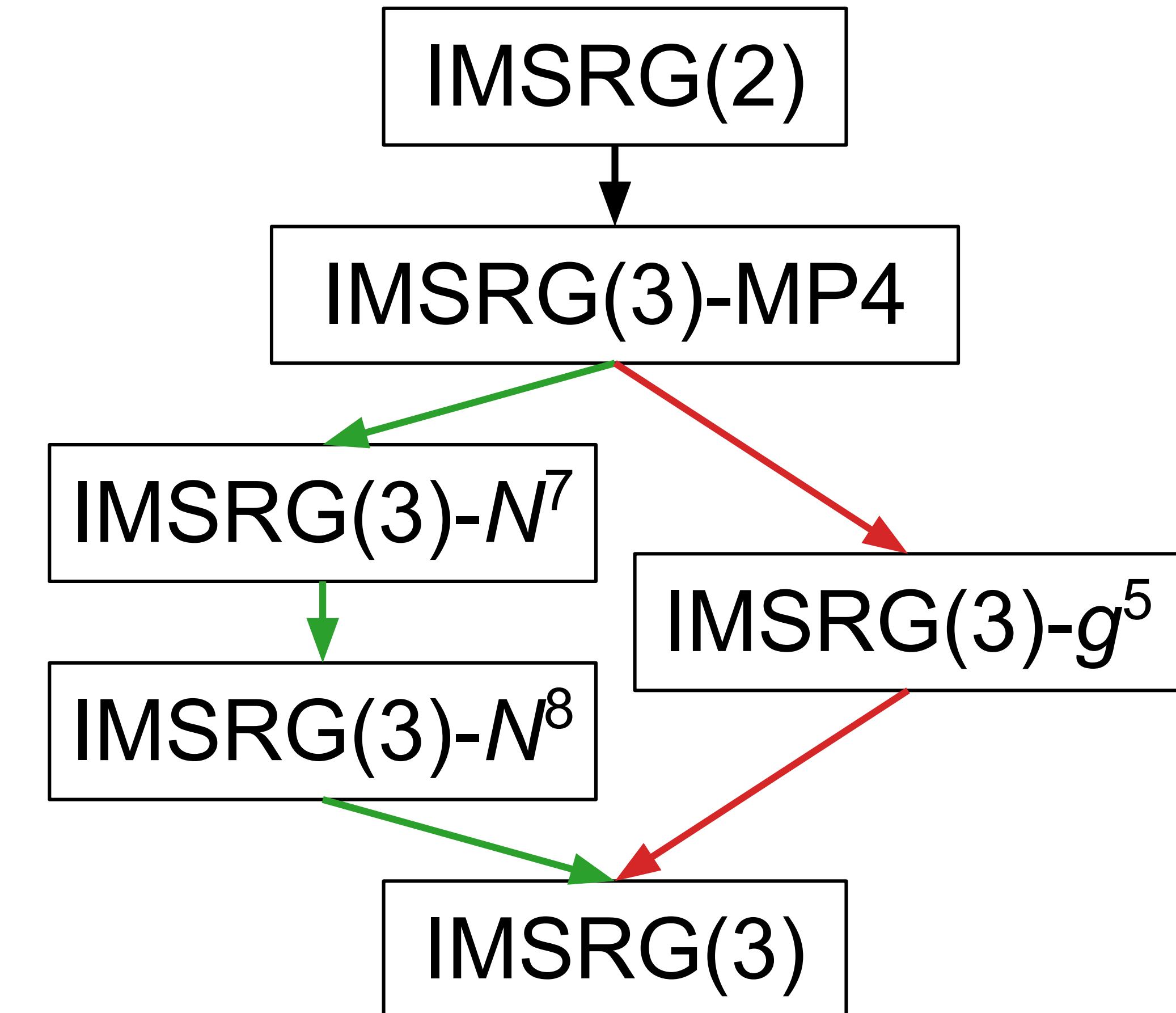
- IMSRG(3) is too expensive
- What can we afford?
- What is actually important?



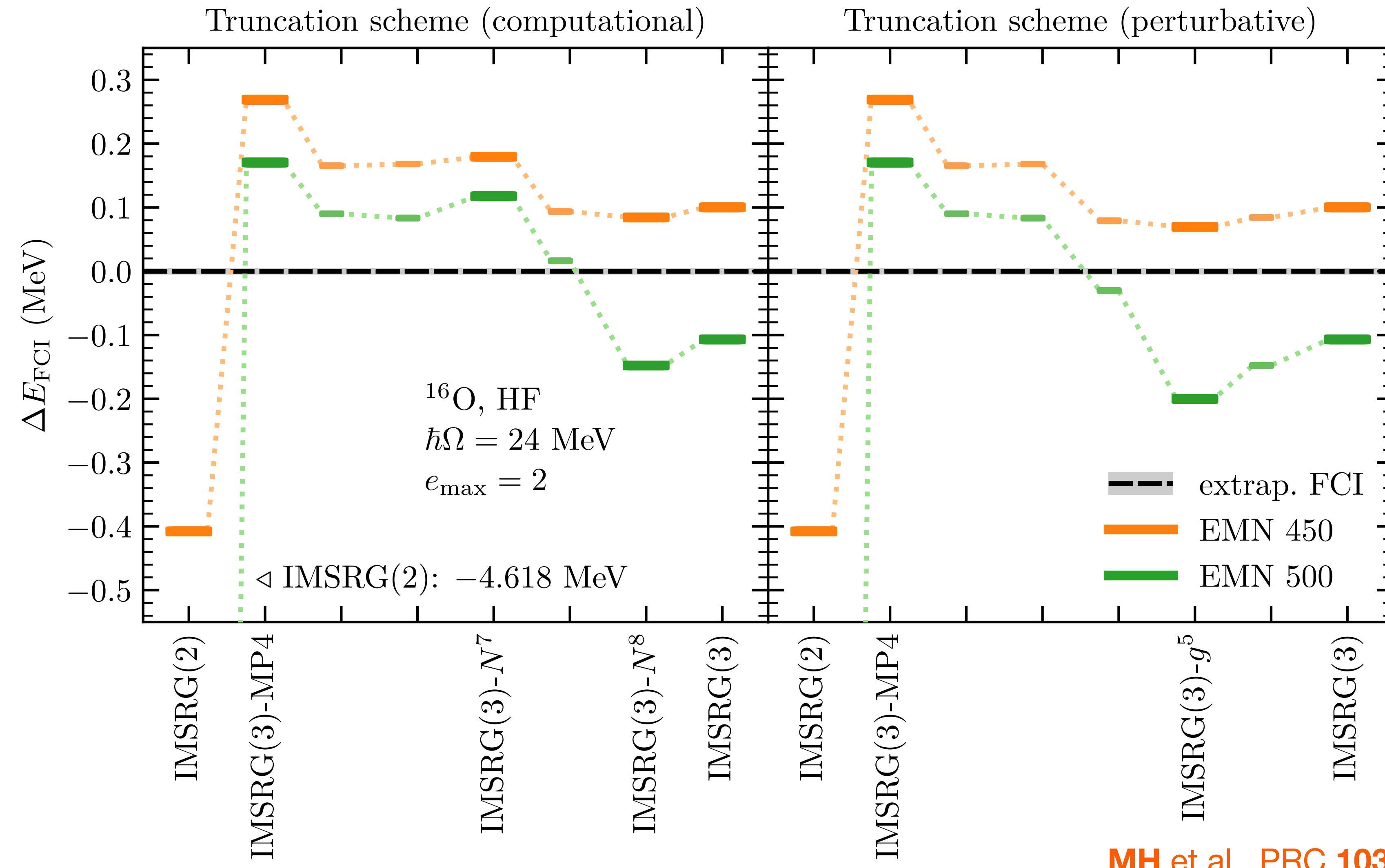
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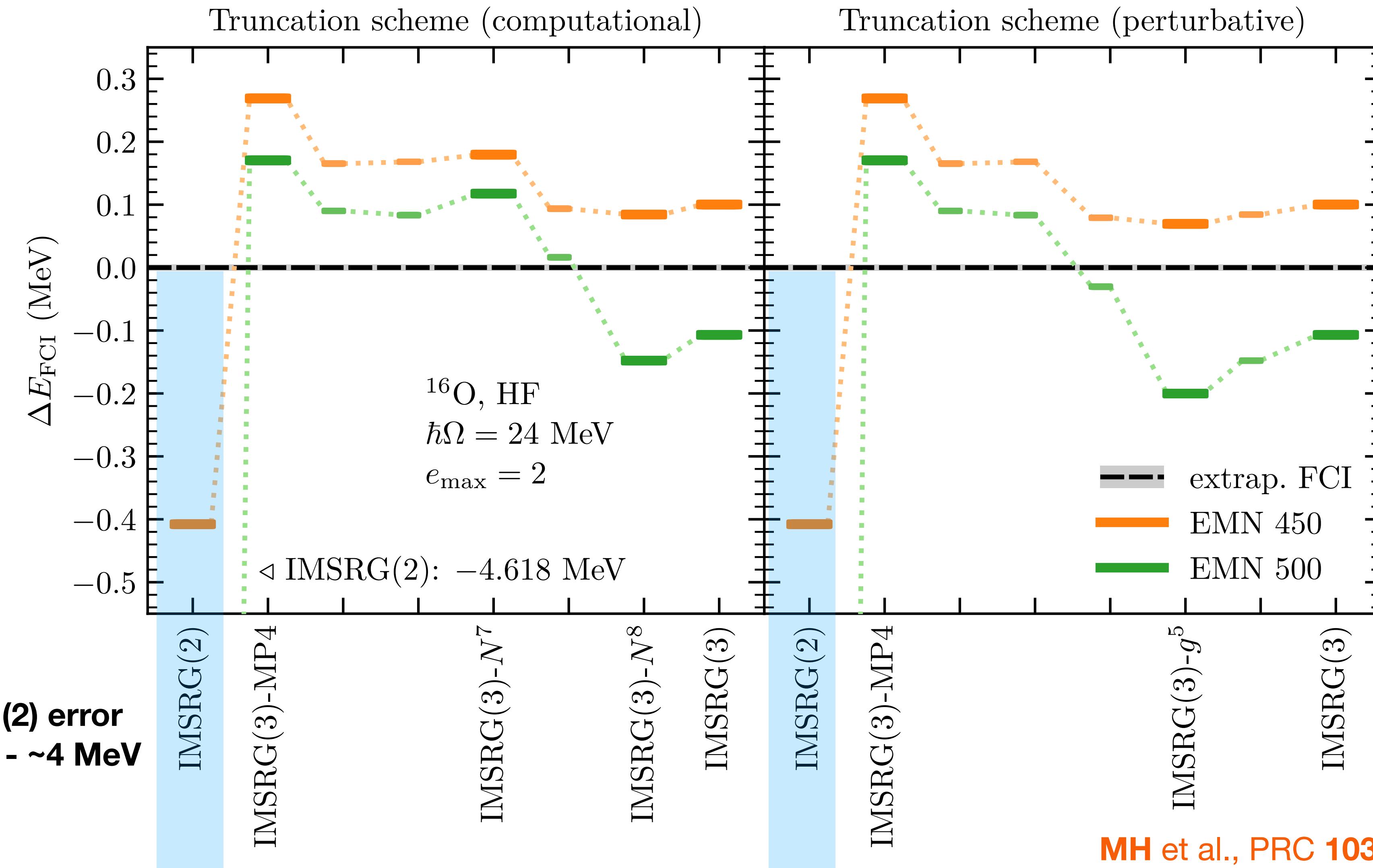
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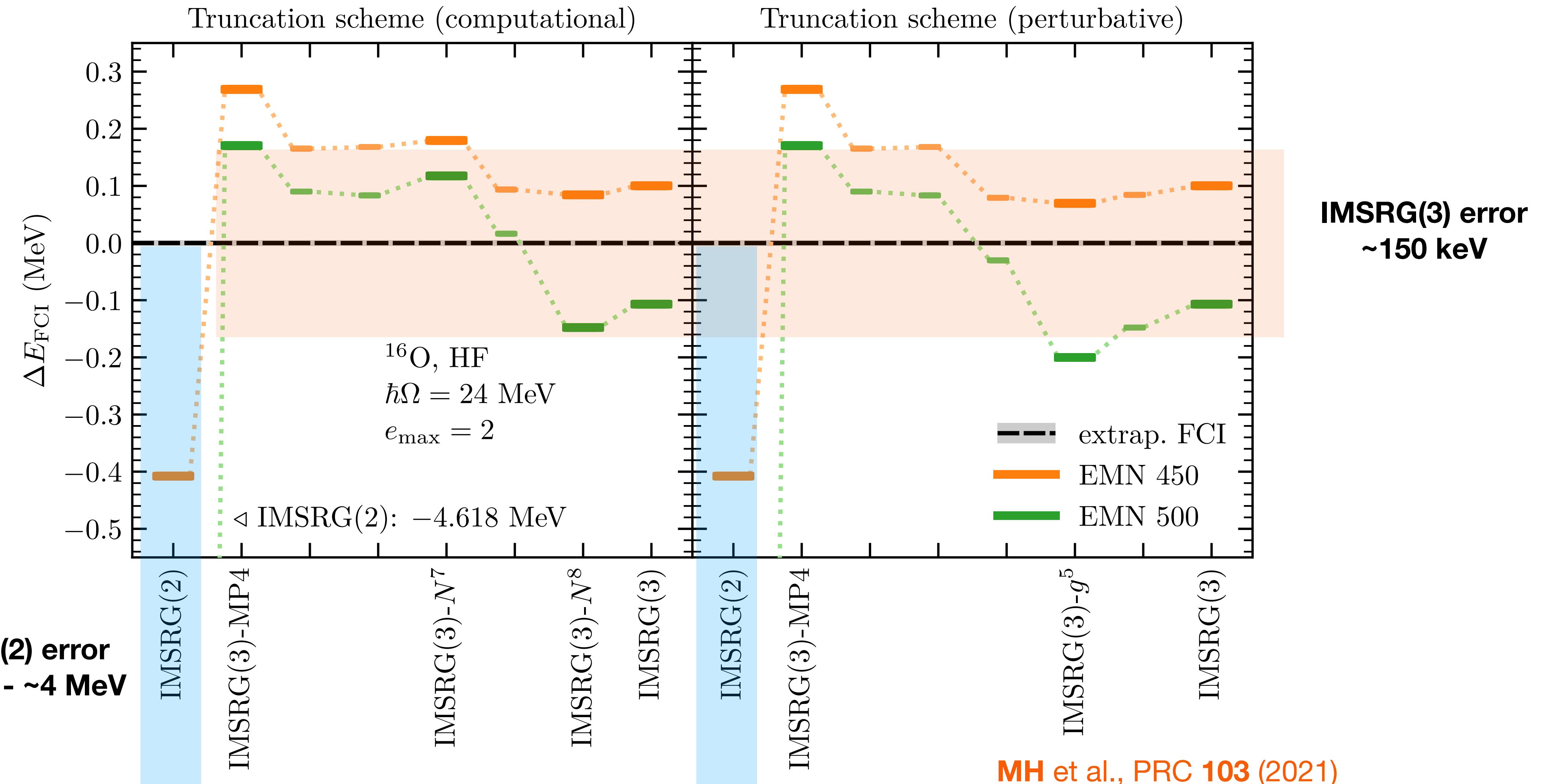
# The IMSRG(3) difference



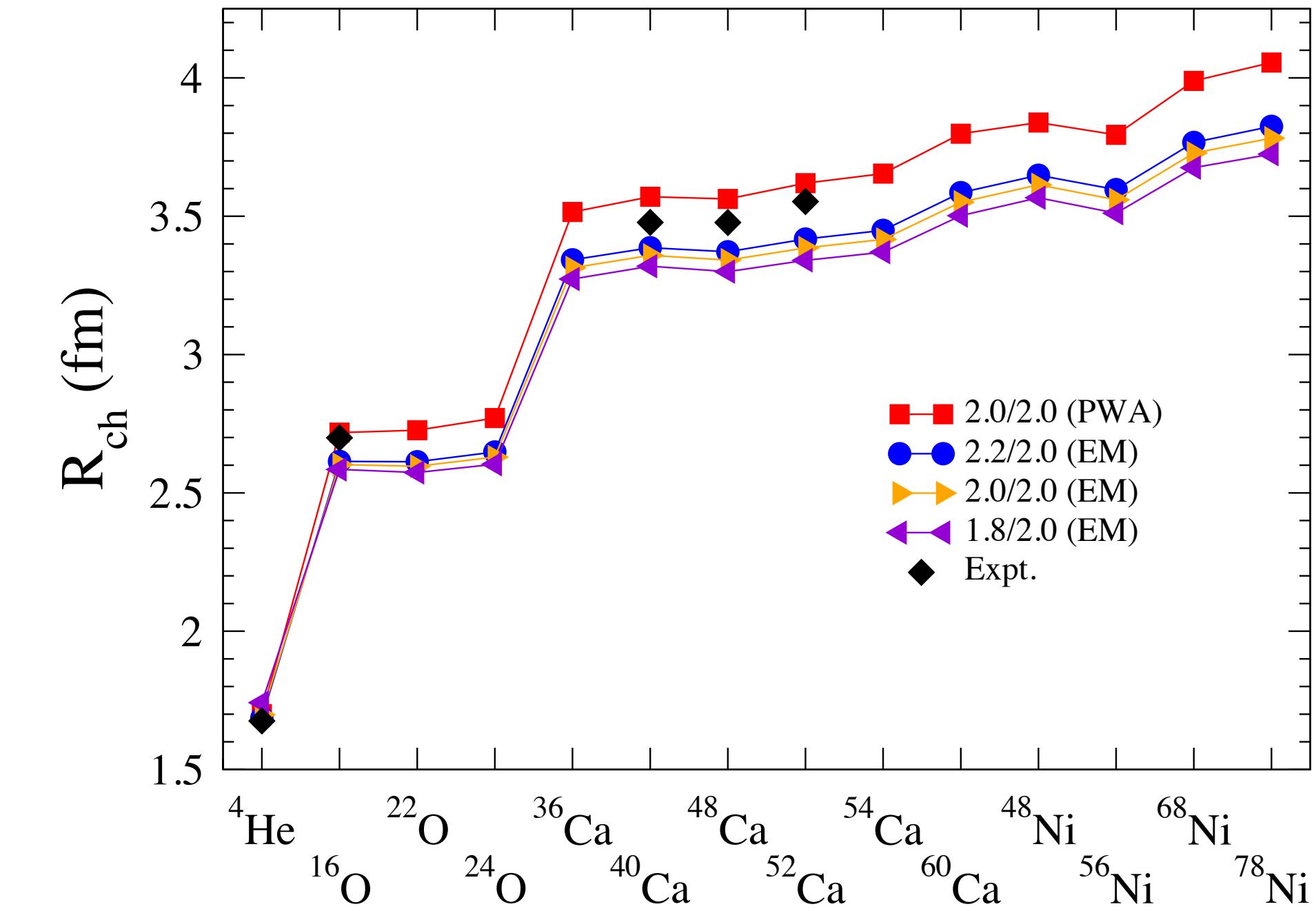
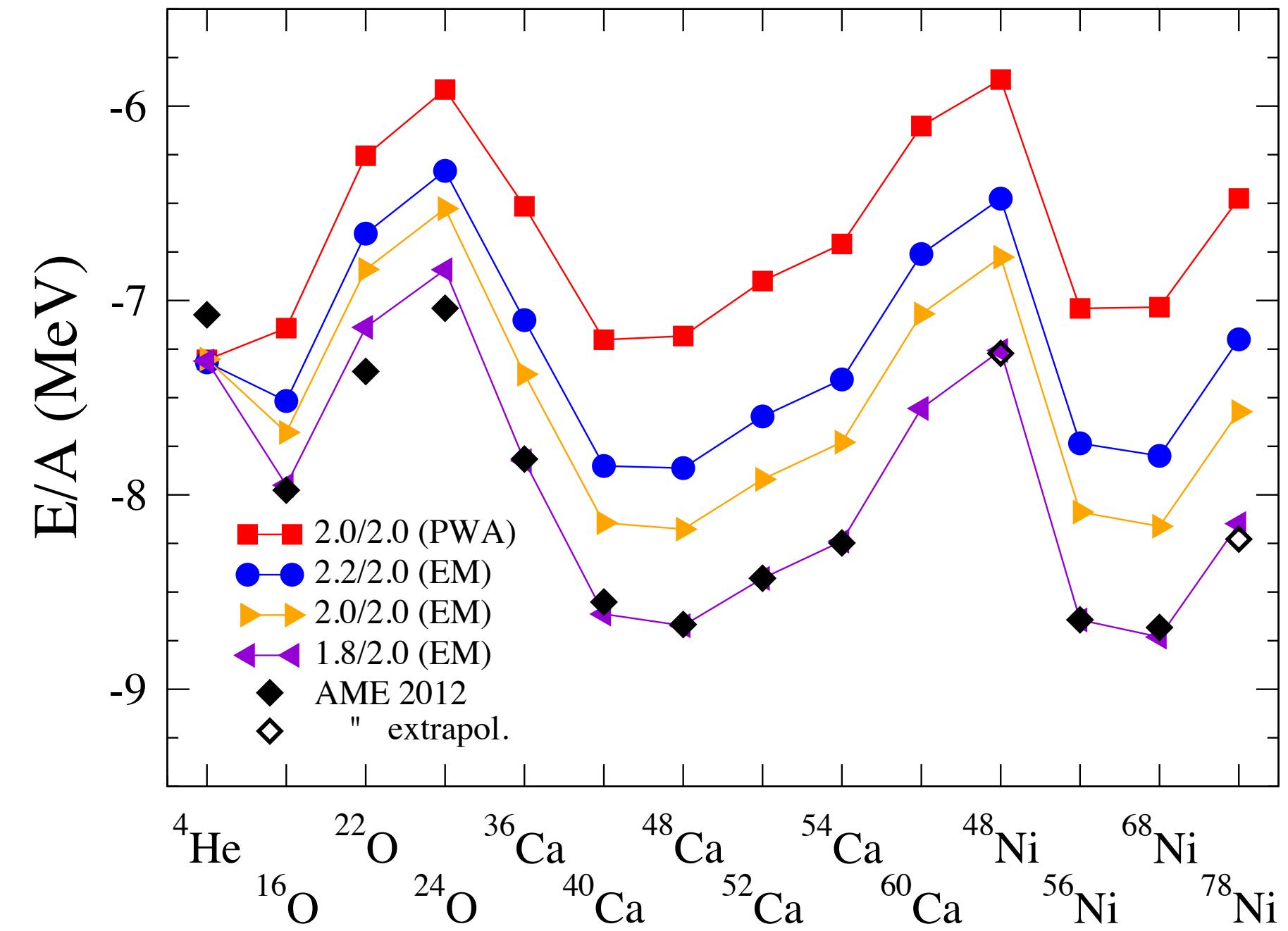
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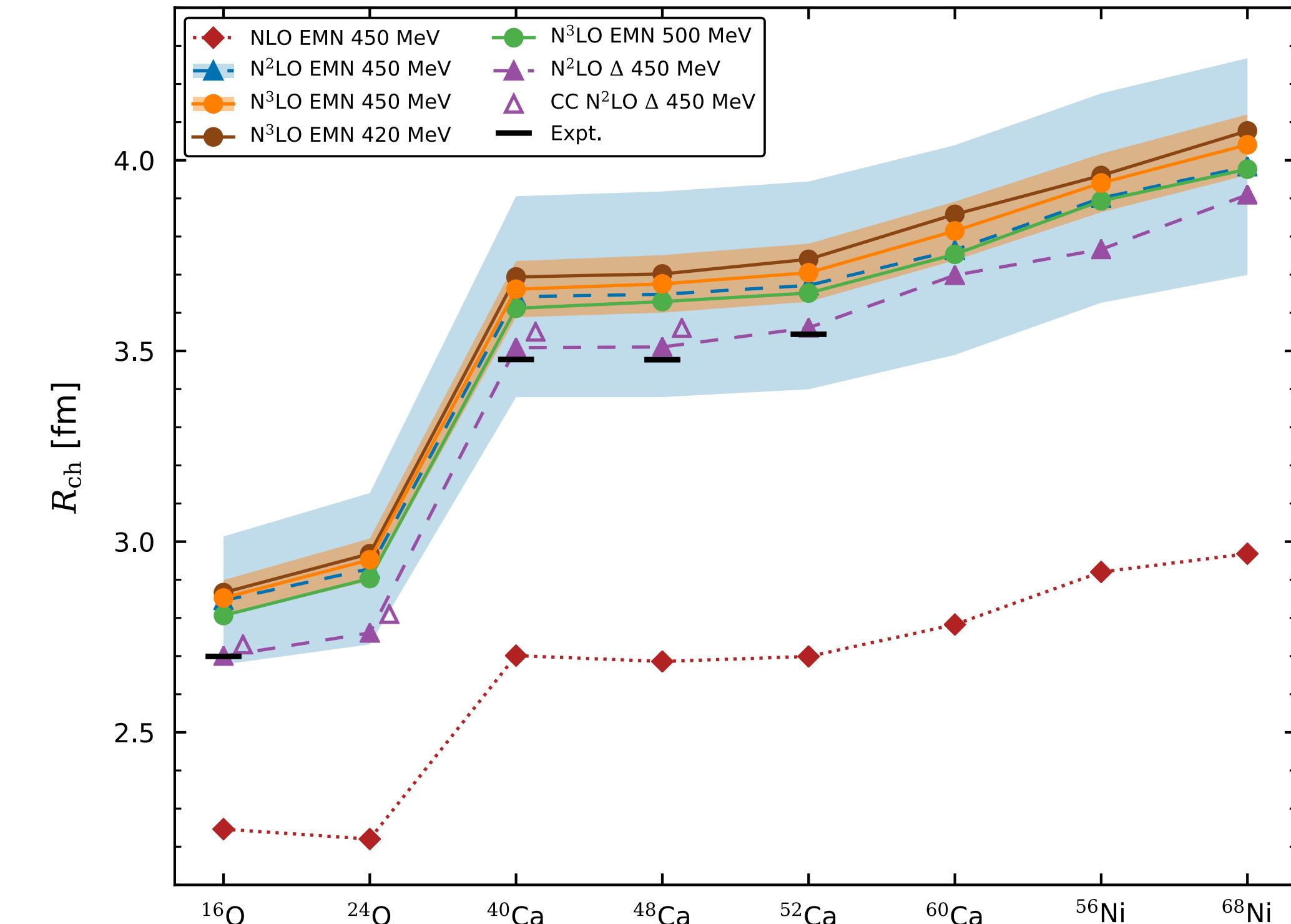
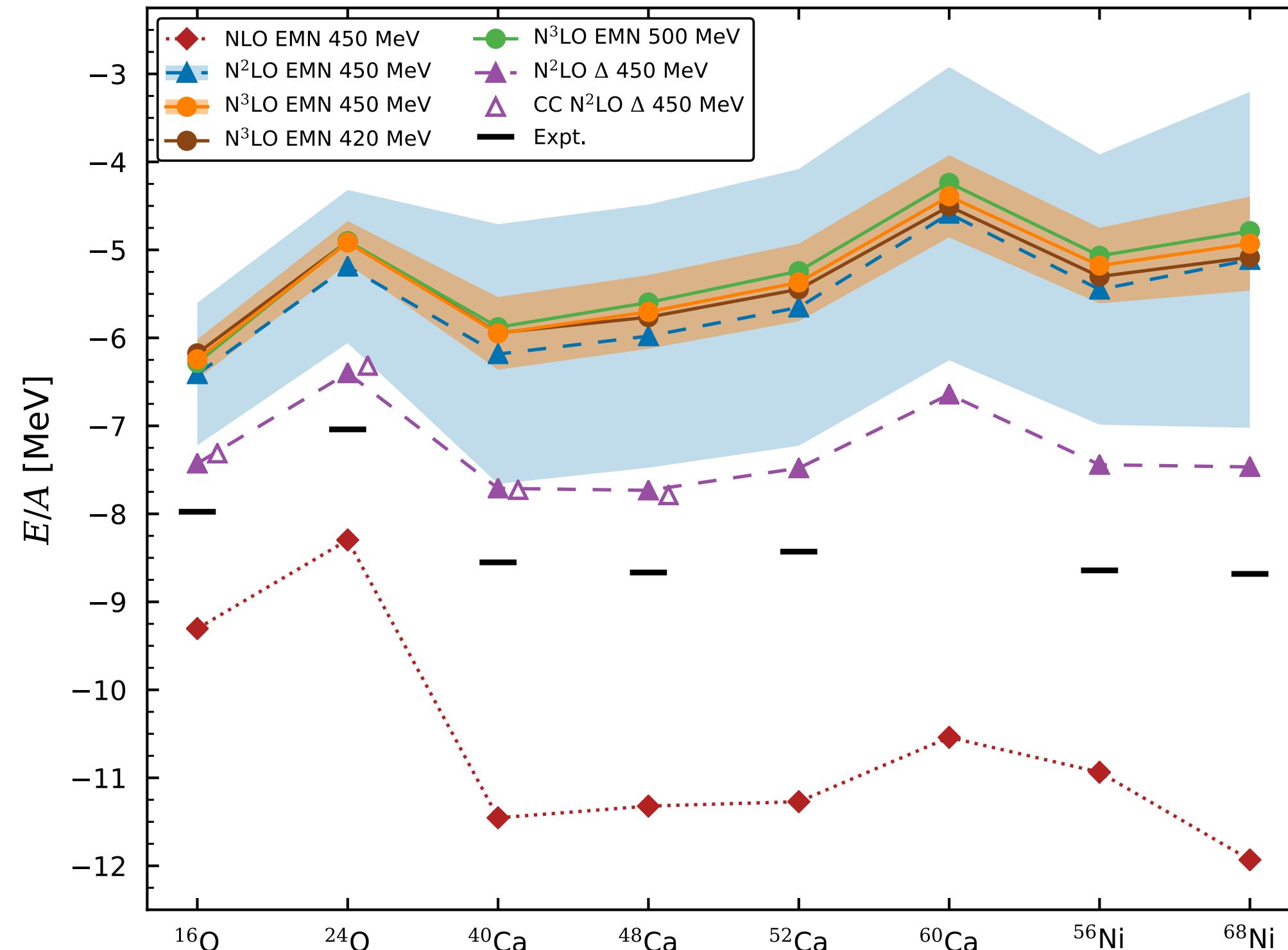
# Challenge: Energies and radii



Simonis et al., PRC 96 (2017)

**Accurate binding energies lead to underpredicted radii!**

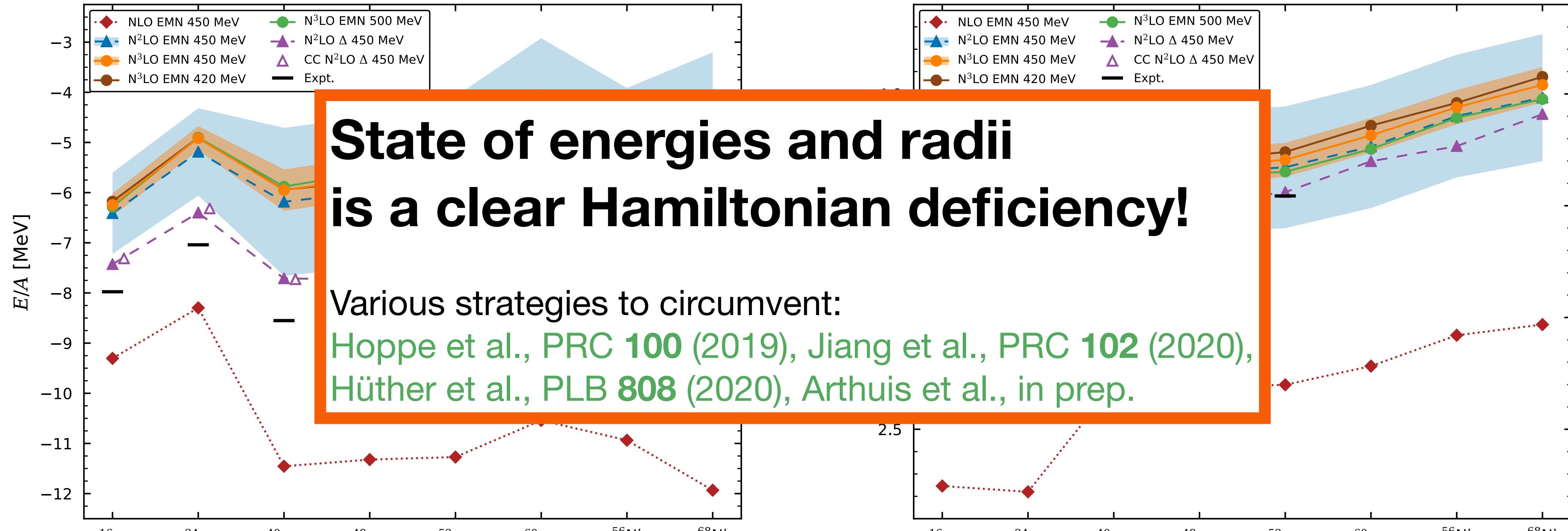
# Challenge: Energies and radii



Hoppe et al., PRC 100 (2019)

**Accurate radii lead to underpredicted binding energies!**

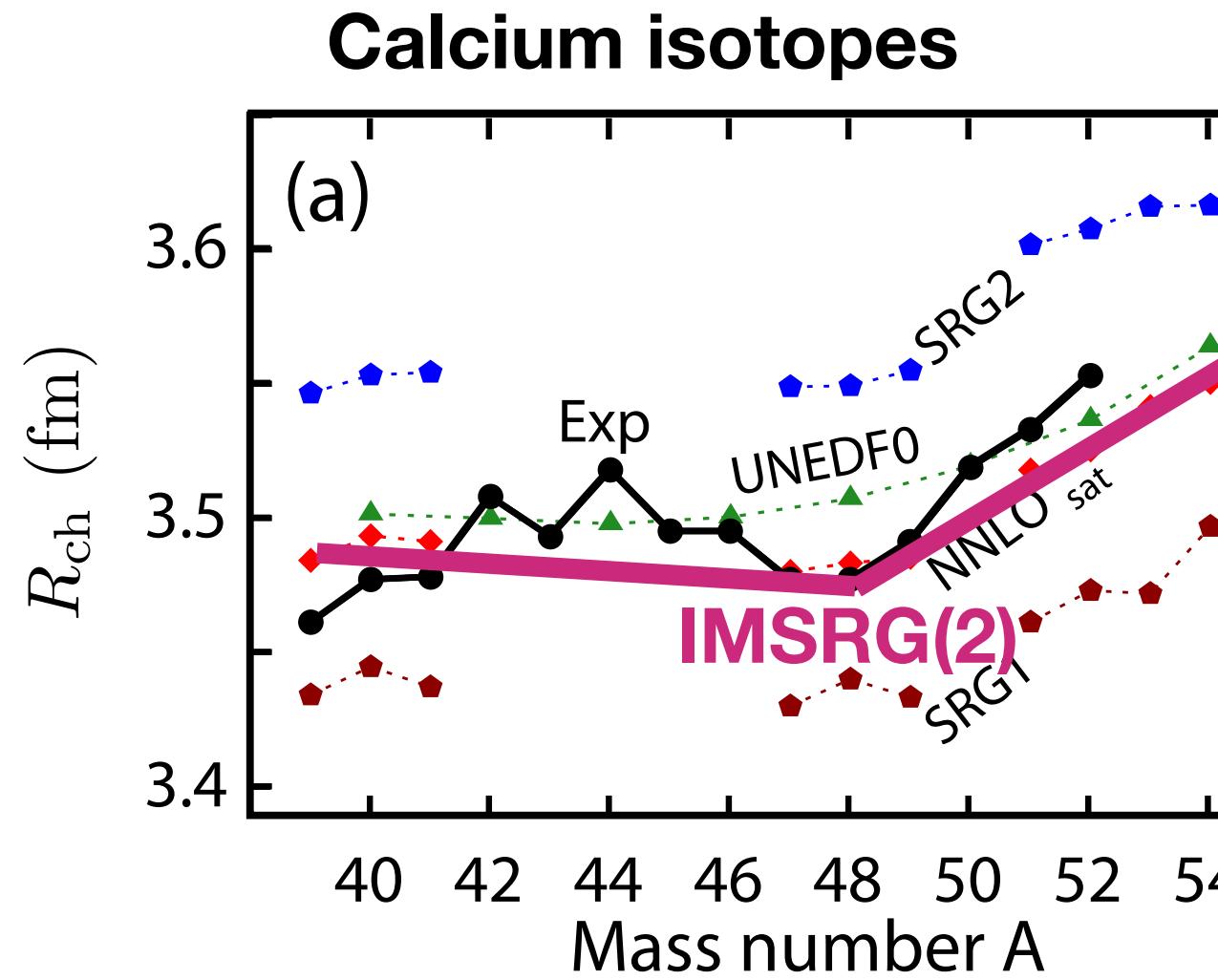
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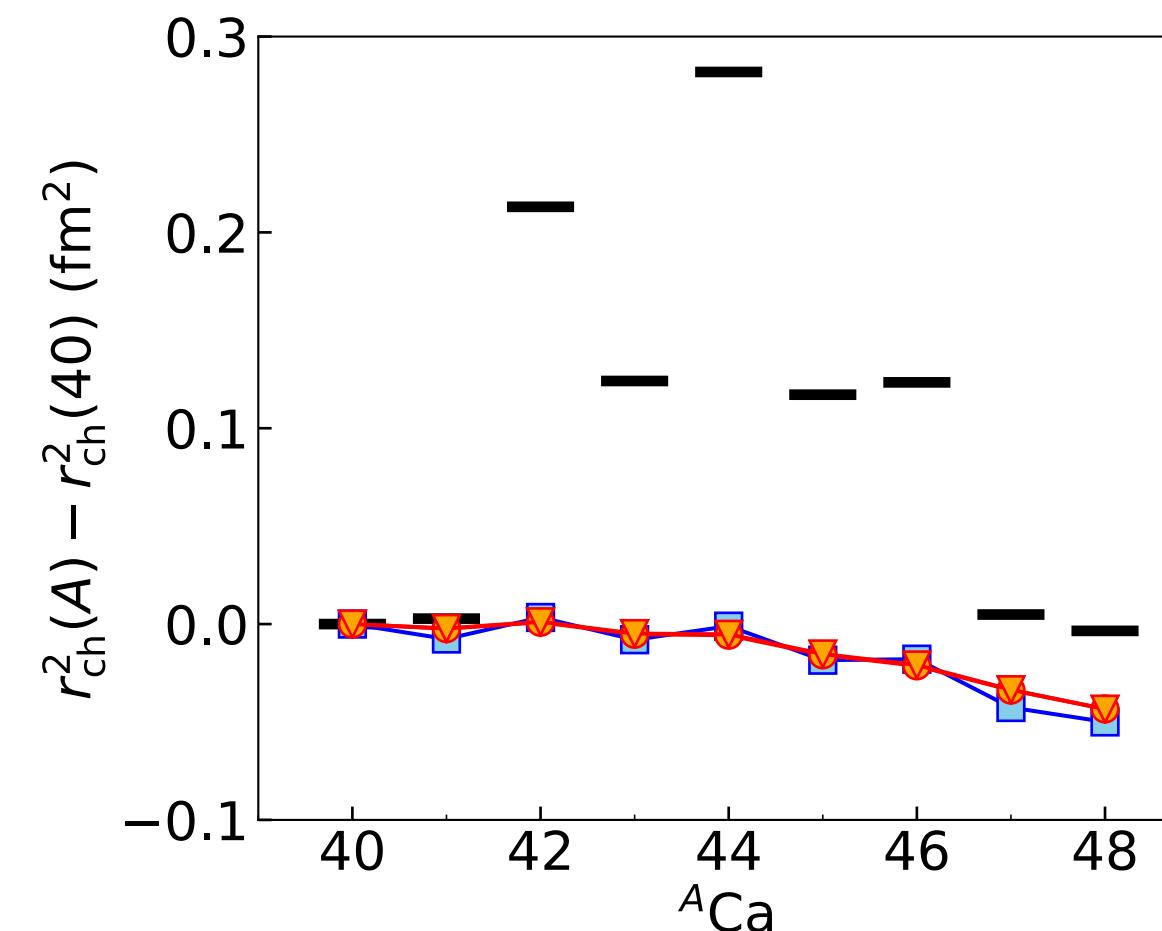
Hoppe et al., PRC 100 (2019)

**Accurate radii lead to underpredicted binding energies!**

# Challenge: Charge radii in Ni and Ca



Garcia Ruiz et al., Nature Physics 12 (2016)



Miyagi et al., PRC 102 (2020)

- Many-body calculations have trouble reproducing
  - $R_{\text{ch}}$  parabola between  $^{40}\text{Ca}$  and  $^{48}\text{Ca}$
  - Steep increase in  $R_{\text{ch}}$  after  $^{48}\text{Ca}$
- Unclear whether Hamiltonian or many-body error
- Similar trend observed in nickel isotopes

Sommer et al., PRL 129 (2022)