

# Emulator of shell-model calculations via eigenvector continuation & a surrogate model for IMSRG

“Eigenvector continuation and related techniques in nuclear structure and reaction theory”  
30 May – 2 June @CEA Saclay

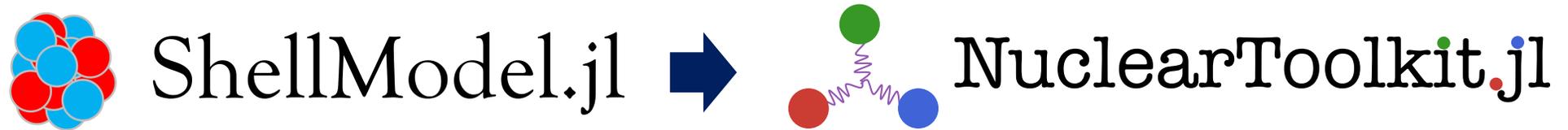
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- ◆ EC + valence shell model
- ◆ A Julia package for nuclear structure calculations



- ◆ IMSRG-Net: a surrogate model for IMSRG
- ◆ Summary

## Constructing approximate shell-model wavefunctions by eigenvector continuation

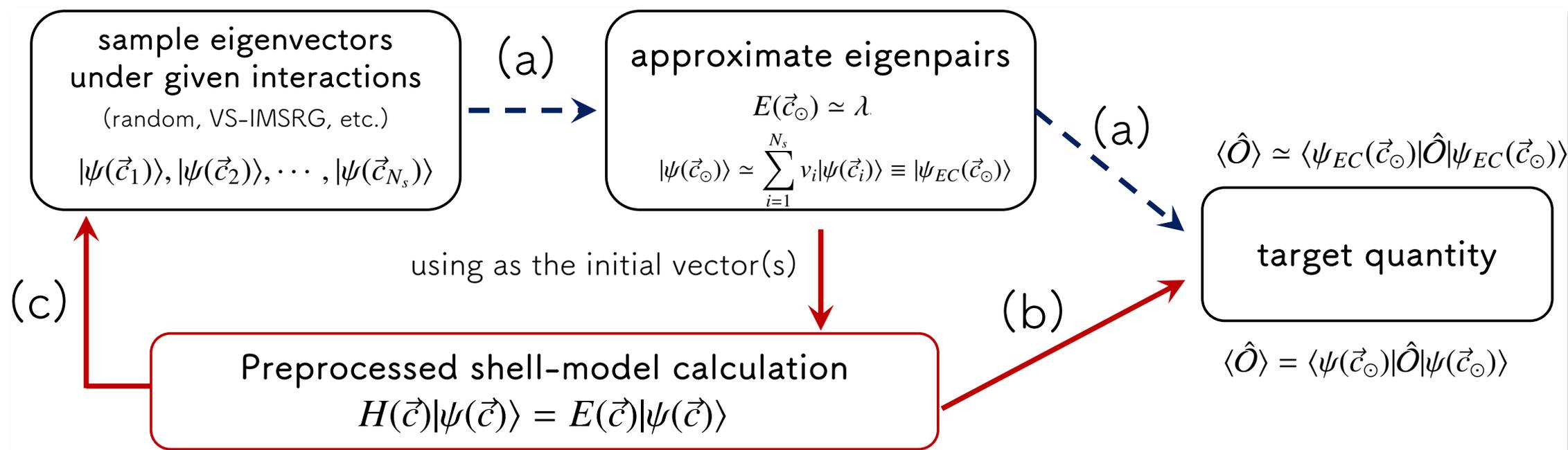


FIG. 1

original problem:

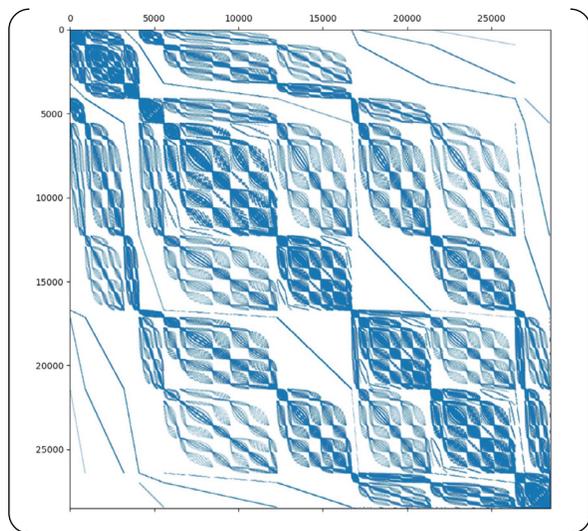
$$H(\vec{c})|\psi(\vec{c})\rangle = E(\vec{c})|\psi(\vec{c})\rangle$$

$$H = H^{(1)} + H^{(2)} = \sum_{ac} h_{ac}^{(1)} c_a^\dagger c_c + \frac{1}{4} \sum_{abcd} h_{abcd}^{(2)} c_a^\dagger c_b^\dagger c_d c_c,$$

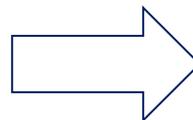
SPEs
TBMEs

large sparse matrix

e.g., <sup>24</sup>Mg M-scheme Dim. = 28,503



row & column:  
many-body configurations

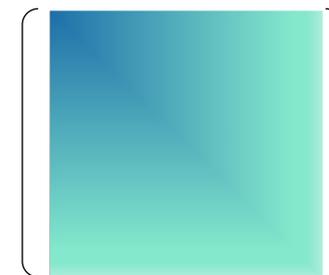


row & column:  
indices for EC samples

EC + shell model:

$$\begin{aligned} \tilde{H}\vec{v} &= \lambda N\vec{v}, \\ \tilde{H}_{i,j} &= \langle \psi(\vec{c}_i) | H(\vec{c}_\odot) | \psi(\vec{c}_j) \rangle, \\ N_{i,j} &= \langle \psi(\vec{c}_i) | \psi(\vec{c}_j) \rangle. \end{aligned}$$

small dense matrix  
Dim. = # of sample ~ 100



N. Shimizu et al., Comp. Phys. Comm. 244 (2019) 372–384

~ sec.

~ msec. (or less?)

Example:

sd-shell ( $^{16}\text{O}$  core + 0d5/0d3/1s1 valence orbits)

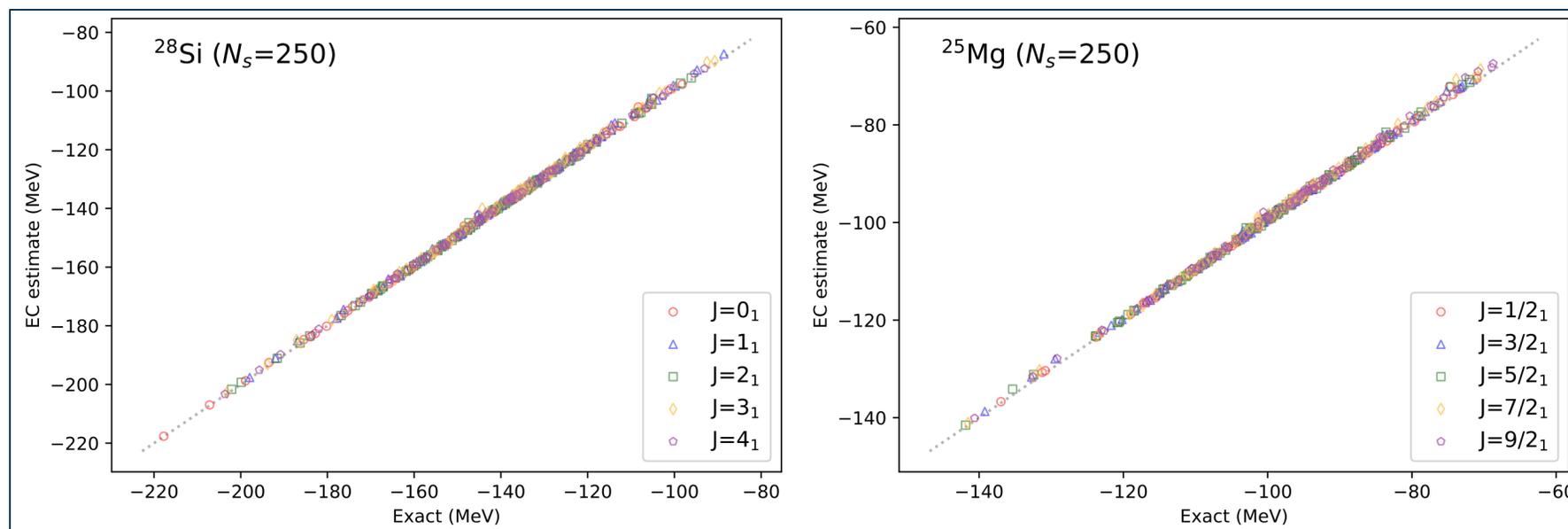
parameters: 66 (3 SPEs & 63 TBMEs, w/ isospin)

target nuclei:  $^{25}\text{Mg}$  (vp=4,vn=5),  $^{28}\text{Si}$  (vp=vn=6, dim.  $\sim 90,000$ )

sampling 5 states for given total J at 50 (random) different points ( $5 \times 50 = 250$  samples) around USDB

$$\begin{aligned} \tilde{H}\vec{v} &= \lambda N\vec{v}, & E(\vec{c}_0) &\simeq \lambda, \\ \tilde{H}_{i,j} &= \langle \psi(\vec{c}_i) | H(\vec{c}_0) | \psi(\vec{c}_j) \rangle, & |\psi(\vec{c}_0)\rangle &\simeq \sum_{i=1}^{N_s} v_i |\psi(\vec{c}_i)\rangle \equiv |\psi_{EC}(\vec{c}_0)\rangle. \\ N_{i,j} &= \langle \psi(\vec{c}_i) | \psi(\vec{c}_j) \rangle. \end{aligned}$$

“validation” for 100 random parameters



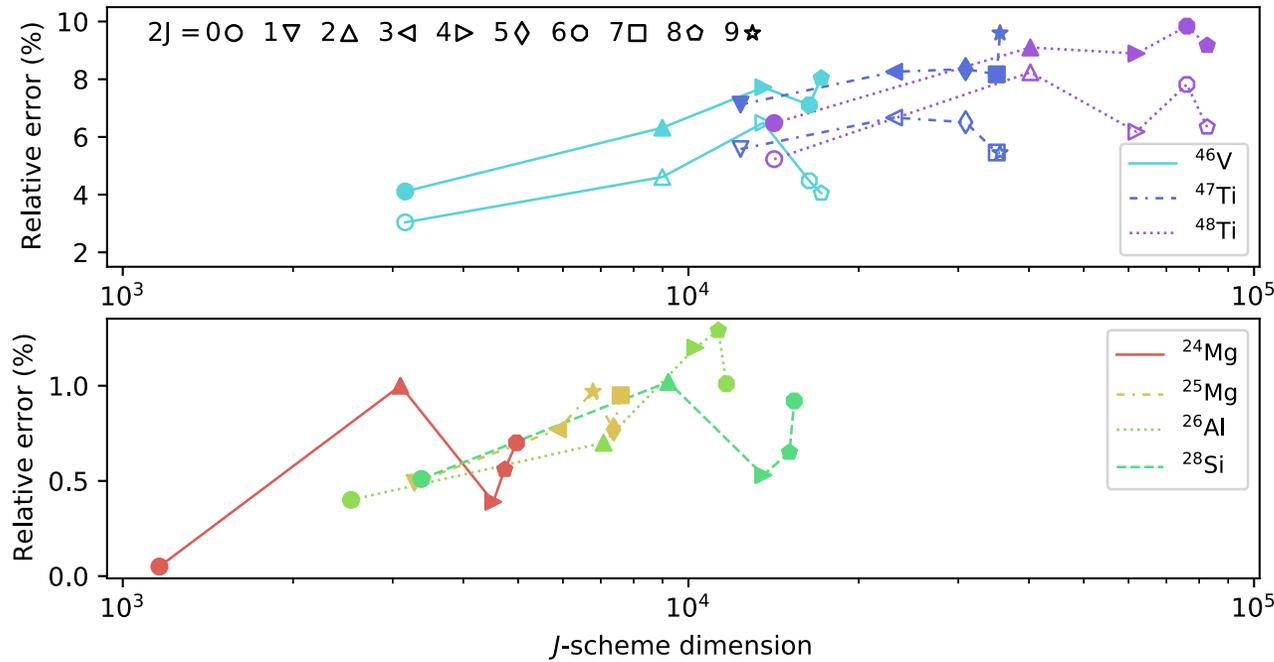
EC approximates energies within a few percent accuracy

**Table 1.** Average sizes of two errors by EC estimates for the five yrast states of the four  $sd$ -shell nuclei: One is the relative error (%) of absolute energies, and the other one is the error of excitation energies. The sample size  $N_s$  means the product of the number of random interactions and the number of excited states used as the *sample* eigenvectors in Eqs. (8)–(9).  $N_s = 250^*$  with  $\sigma_{\text{int.}} = 3$  means that the standard deviation to generate the random interactions is increased from the default value  $\sigma_{\text{int.}} = 1$ , and  $N_s = 250^*$  (LHS,  $L = 2$ ) corresponds to the result using Latin hypercube sampling (LHS).

$N_s$ (# interaction $\times$ # states)	mean of $\text{relative error}(\%) \equiv 100 \left  \frac{E_{\text{exact}} - E_{\text{EC}}}{E_{\text{exact}}} \right $				mean of $\text{ex. error (MeV)} \equiv  E_{\text{exact}}^{\text{ex.}} - E_{\text{EC}}^{\text{ex.}} $			
	$^{28}\text{Si}$	$^{26}\text{Al}$	$^{25}\text{Mg}$	$^{24}\text{Mg}$	$^{28}\text{Si}$	$^{26}\text{Al}$	$^{25}\text{Mg}$	$^{24}\text{Mg}$
50 (50 $\times$ 1)	1.4	2.1	1.8	1.3	0.66	1.22	0.62	0.65
50* (25 $\times$ 2)	1.8	2.3	2.1	1.7	0.82	1.16	0.61	0.97
150 (50 $\times$ 3)	0.9	1.2	1.1	0.7	0.44	0.85	0.42	0.62
250 (50 $\times$ 5)	0.7	0.9	0.8	0.5	0.39	0.70	0.37	0.51
250* (50 $\times$ 5; $\sigma_{\text{int.}} = 3$ )	2.8	3.3	3.1	2.3	1.35	2.35	1.09	1.96
250* (50 $\times$ 5; LHS, $L = 2$ )	0.8	1.0	0.9	0.6	0.47	0.73	0.40	0.57

each row shows different settings (# of samples, way of sampling)

- sample not only g.s. but also excited states if you want to know excited states too
- odd or odd-odd nuclei are more difficult than even (even-even) ones



**Fig. 7.** Relative errors with  $N_s = 250$  samples against  $J$ -scheme dimensions for  $^{24,25}\text{Mg}$ ,  $^{26}\text{Al}$ ,  $^{28}\text{Si}$  (lower panel),  $^{46}\text{V}$ , and  $^{47,48}\text{Ti}$  (upper panel). The filled symbols correspond to the samples generated by varying all the parameters with  $\sigma_{\text{int.}} = 1$  around the reference values (USDB and GXPF1A). The open symbols in the upper panel show the results with samples in which only the 32 parameters related to  $f_{7/2}$  and  $p_{3/2}$  were varied with the same  $\sigma_{\text{int.}}$ .

Dotted lines are results considering only the samples spread over  $f_{7/2}$  &  $p_{3/2}$

it is better to sample over a subspace more relevant to what you want to know

Q. What is a better sampling strategy?

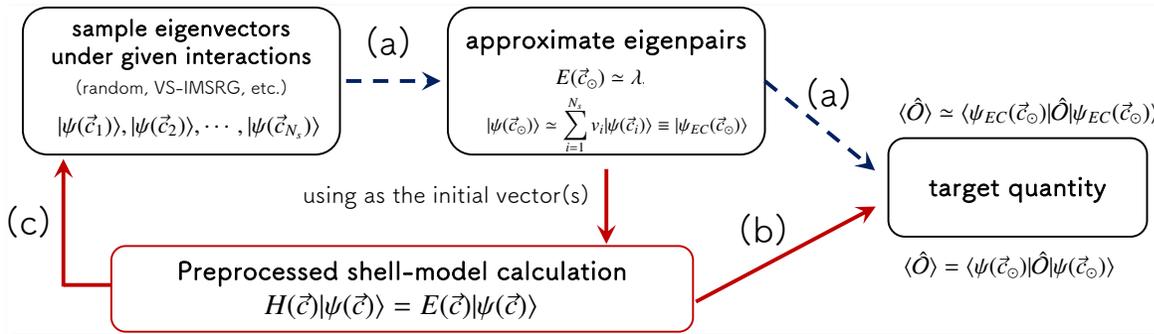
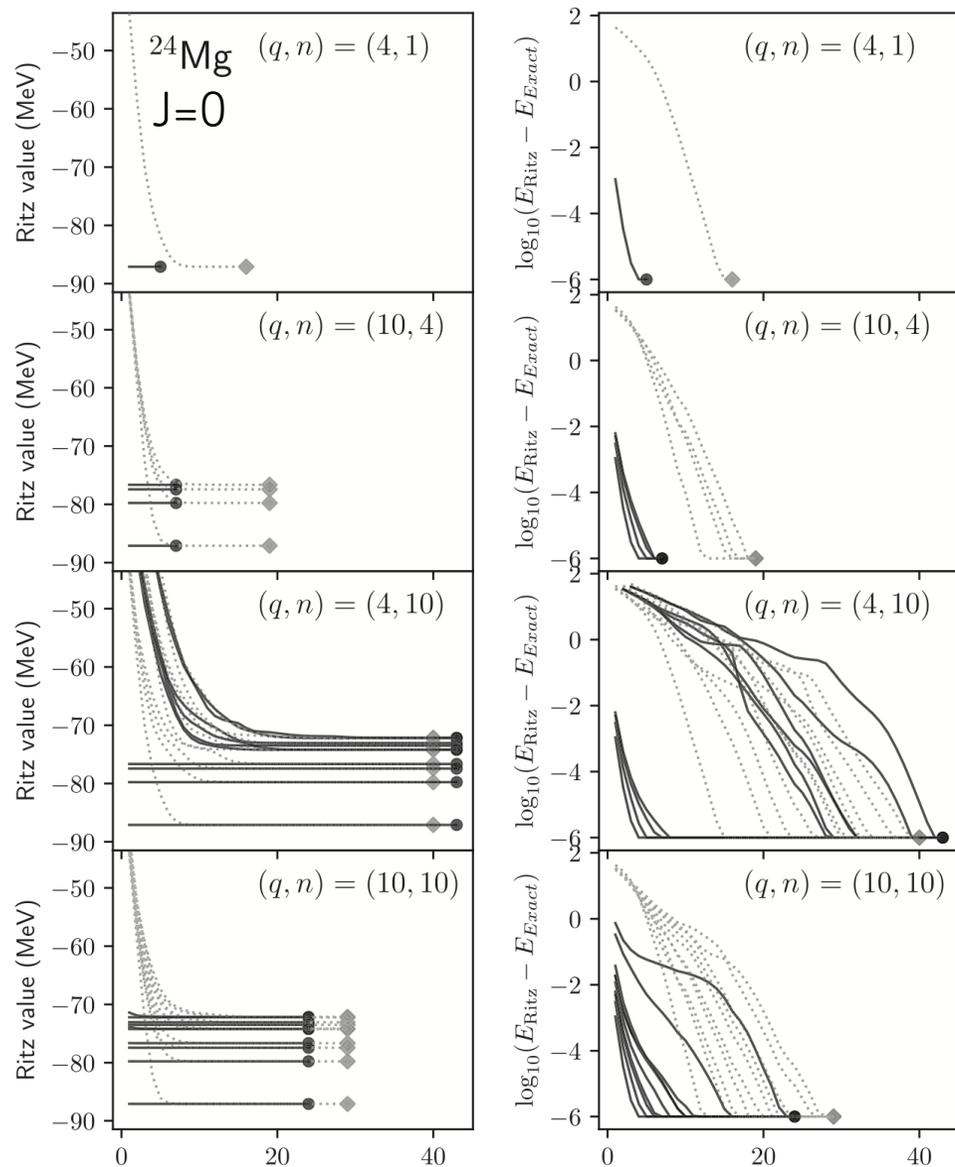
Is there any way to maximize information gain from a next observation?

no answer will be shown in this talk though...

pf-shell => 199 parameters (4 SPEs & 195 TBMEs)

(craziest application of EC method !?)

# As a preprocessor



$q$ : size of initial “block” vector

$n$ : # of excited states of interest

dotted: initialized by random vectors

solid: initialized by EC eigenvectors

Starting from better initial guess,  
# of manipulation could be reduced!!

Exception =>  $(q, n) = (4, 10)$

since the emulator is trained with 5 lowest states,  
such emulator do not have much info. on higher states

Number of H operation during the (block) Lanczos method

converged results are obtained at ●/◆

Sampling itself is not easy ...

$$\tilde{H}\vec{v} = \lambda N\vec{v},$$

$$\tilde{H}_{i,j} = \langle \psi(\vec{c}_i) | H(\vec{c}_\odot) | \psi(\vec{c}_j) \rangle, \quad \leftarrow \text{most time-consuming part}$$

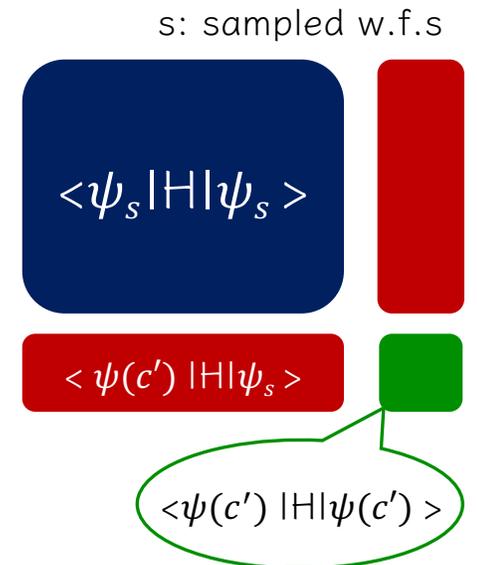
$$N_{i,j} = \langle \psi(\vec{c}_i) | \psi(\vec{c}_j) \rangle.$$

- You don't need to explicitly calculate  $H(\vec{c}_\odot) | \psi(\vec{c}_j) \rangle$  for each parameter  $\vec{c}_\odot$  to evaluate  $\tilde{H}$  above:

all you need is 1&2-body transition densities

$$\tilde{H}_{i,j} = \sum_k \underbrace{h_k^{(1)}}_{\text{SPEs}} \times \underbrace{\overline{\text{OBTD}}_{i,j,k}}_{\text{OBTD}} + \sum_k \underbrace{V_J(abcd)_k}_{\text{TBMEs}} \times \underbrace{\overline{\text{TBTD}}_{i,j,k}}_{\text{TBTD}}$$

- If you want to increase sample number (for better accuracy), prepare new sample (green) and calc. overlap (transition densities) between new w.f. and previous samples (red)

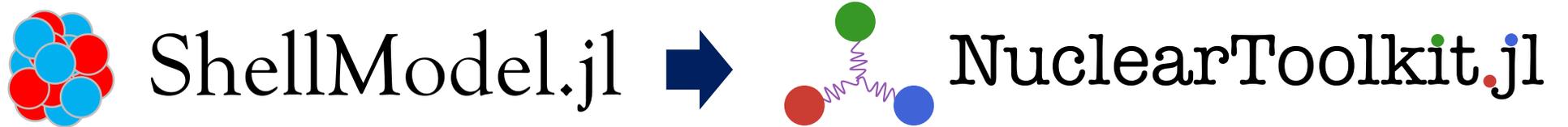


# Outline

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- ◆ EC + valence shell model

- 👉 ◆ A Julia package for nuclear structure calculations



- ◆ IMSRG-Net: a surrogate model for IMSRG

- ◆ Summary

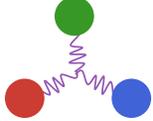
- A single method (code) can be lengthy ~ 100,000 lines  
two language problem (Fortran/C++ & shell/Python)
- Especially in Japan, research methods are “clusterized” (localized)  
to a specific group. This can be an obstacle to  
collective intelligence or co-creation  
seen in e.g. ML community through ML frameworks
- Educations for next generation



secret sauce (source)  
in Prof. XX Group



can circumvent the situation...?



- ChiEFTint ~ 8000 lines
  - ◆ NN potential, Entem-Machleidt(N3LO), EMN(EKMN, N4LO)
  - ◆ SRG in momentum space (NN-only)
  - ◆ effective NN from 3NF
  - ◆ valence NN interaction  $\neq$  effective interaction
  - ◆ input for No-core shell model (in KSHELL fmt)
  - ◆ genuine 3NF (only in Jacobi HO form)
    - ✧ plz use NuHamil (by Takayuki Miyagi@TUDarmstadt)
- HartreeFock ~ 3000 lines
  - ◆ spherical HF (from snt/snt.bin/memory)
  - ◆ HFMBPT Energy=> 3rd order, Scaler operator => 2nd order
  - ◆ Normal ordering w.r.t. target reference, ensemble normal ordering
- IM-SRG ~2700 lines
  - free space
    - ◆ IMSRG(2) calculation => g.s. properties
    - ◆ consistent IMSRG flow of operators with Magnus expansion
  - valence space (VS-IMSRG)
    - ◆ derive effective interaction for a target model space
    - ◆ effective operators (only scaler ones for now)
-  ShellModel.jl ~5000 lines
  - ◆ eigenvector continuation (fast emulator of exact wavefunctions)

~ 20,000 lines  
(including document)

Do you think it's lengthy? 🤔

Docs are automatically generated from docstring (in markdown)

```

"""
    HF_MBPT2(bininfo,modelspace,fp,fn,e1b_p,e1b_n,Chan2b,Gamma)
Calculate 2nd order correction to HF energy
```math
E^{(2)} = \frac{1}{4} \sum_{abij} \frac{\bar{H}_{abij} \bar{H}_{ijab}}{\epsilon_{ij}^{ab}} = \frac{1}{4} \sum_{\bar{a}\bar{b}\bar{i}\bar{j}} \sum_{\{m\}} \sum_{JJ'MM'} \frac{J \bar{H}_{\bar{a}\bar{b}\bar{i}\bar{j}}^{[2]} J \bar{H}_{\bar{i}\bar{j}\bar{a}\bar{b}}^{[2]}}{\epsilon_{ij}^{ab}} (j_a j_b m_a m_b | JM) (j_a j_b m_a m_b | J' M') (j_i j_j)
(j_a j_b m_a m_b | J' M')
(j_i j_j m_i m_j | JM)
(j_i j_j m_i m_j | J' M')
```math
=
\frac{1}{4} \sum_{\bar{a}\bar{b}\bar{i}\bar{j}} \sum_{\{m\}} \sum_{JJ'MM'} \frac{J \bar{H}_{\bar{a}\bar{b}\bar{i}\bar{j}}^{[2]} J \bar{H}_{\bar{i}\bar{j}\bar{a}\bar{b}}^{[2]}}{\epsilon_{ij}^{ab}} (j_a j_b m_a m_b | JM) (j_a j_b m_a m_b | J' M') (j_i j_j)
\delta_{JJ'} \delta_{MM'} = \frac{1}{4} \sum_{\bar{a}\bar{b}\bar{i}\bar{j}} \sum_{JJ'MM'} \frac{J \bar{H}_{\bar{a}\bar{b}\bar{i}\bar{j}}^{[2]} J \bar{H}_{\bar{i}\bar{j}\bar{a}\bar{b}}^{[2]}}{\epsilon_{ij}^{ab}} \delta_{JJ'} \delta_{MM'} = \frac{1}{4} \sum_{\bar{a}\bar{b}\bar{i}\bar{j}} \sum_J (2J + 1) \frac{J \bar{H}_{\bar{a}\bar{b}\bar{i}\bar{j}}^{[2]} J \bar{H}_{\bar{i}\bar{j}\bar{a}\bar{b}}^{[2]}}{\epsilon_{ij}^{ab}}
"""
function HF_MBPT2(bininfo,modelspace,fp,fn,e1b_p,e1b_n,Chan2b,Gamma)
    p_sps = modelspace.p_sps

```

← docstring



NuclearToolkit.HF\_MBPT2 – Method

---

HF\_MBPT2(bininfo,modelspace,fp,fn,e1b\_p,e1b\_n,Chan2b,Gamma)

Calculate 2nd order correction to HF energy

$$E^{(2)} = \frac{1}{4} \sum_{abij} \frac{\bar{H}_{abij} \bar{H}_{ijab}}{\epsilon_{ij}^{ab}} = \frac{1}{4} \sum_{\bar{a}\bar{b}\bar{i}\bar{j}} \sum_{\{m\}} \sum_{JJ'MM'} \frac{J \bar{H}_{\bar{a}\bar{b}\bar{i}\bar{j}}^{[2]} J \bar{H}_{\bar{i}\bar{j}\bar{a}\bar{b}}^{[2]}}{\epsilon_{ij}^{ab}} (j_a j_b m_a m_b | JM) (j_a j_b m_a m_b | J' M') (j_i j_j)$$

$$= \frac{1}{4} \sum_{\bar{a}\bar{b}\bar{i}\bar{j}} \sum_{JJ'MM'} \frac{J \bar{H}_{\bar{a}\bar{b}\bar{i}\bar{j}}^{[2]} J \bar{H}_{\bar{i}\bar{j}\bar{a}\bar{b}}^{[2]}}{\epsilon_{ij}^{ab}} \delta_{JJ'} \delta_{MM'} = \frac{1}{4} \sum_{\bar{a}\bar{b}\bar{i}\bar{j}} \sum_J (2J + 1) \frac{J \bar{H}_{\bar{a}\bar{b}\bar{i}\bar{j}}^{[2]} J \bar{H}_{\bar{i}\bar{j}\bar{a}\bar{b}}^{[2]}}{\epsilon_{ij}^{ab}}$$

NuclearToolkit
⚙️

## NuclearToolkit

Julia Toolkit for nuclear structure calculations

### How to Start

First, prepare Julia environment  $v \geq 1.6$ .

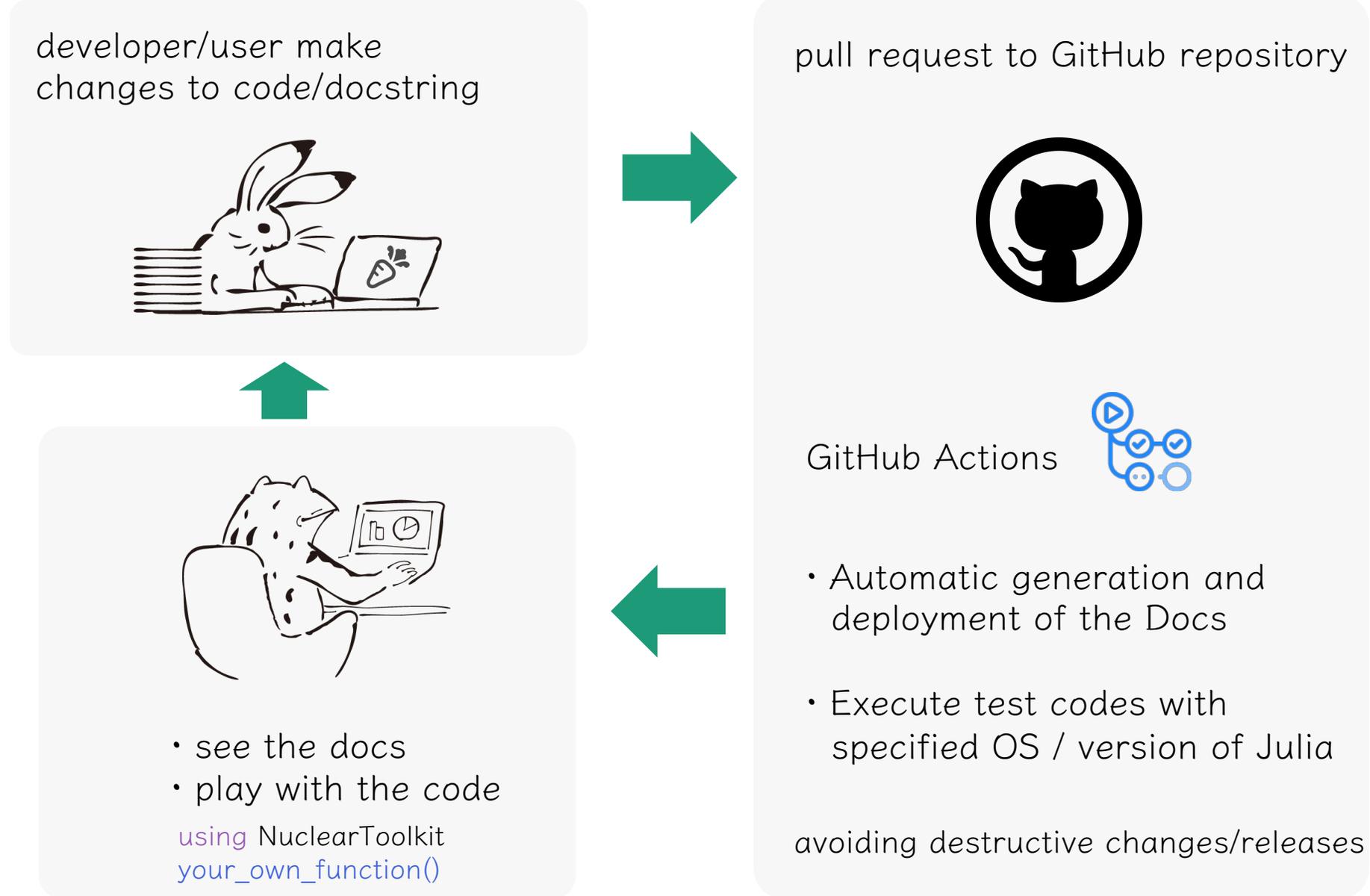
Second, add the package in Pkg mode

```
julia>add NuclearToolkit
```

Note: The above is currently not working, since this package has not yet registered as an official package. The adhoc prescription for now is to execute below

```
julia src/package_install.jl
```

### Package features and building blocks



## ◆ Installation of Julia

Download Julia binary and add to PATH

## ◆ Installation of NuclearToolkit.jl

Download src (recommended)

```
$git clone https://github.com/SotaYoshida/NuclearToolkit.jl
```

Note:

NuclearToolkit.jl was registered as a Julia package.

You can install the package in Julia's REPL

like “pip” in Python

```
Documentation: https://docs.julialang.org
Type "?" for help, "]" for Pkg help.
Version 1.7.3 (2022-05-06)
Official https://julialang.org/ release

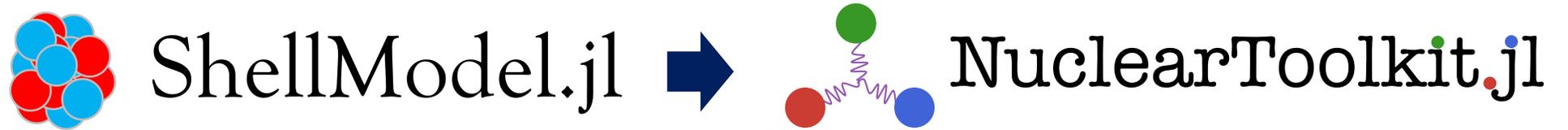
julia> using Pkg

julia> Pkg.add("NuclearToolkit")
Updating registry at `~/.julia/registries/General`
Updating git-repo `https://github.com/JuliaRegistries/General.git`
Resolving package versions...
```

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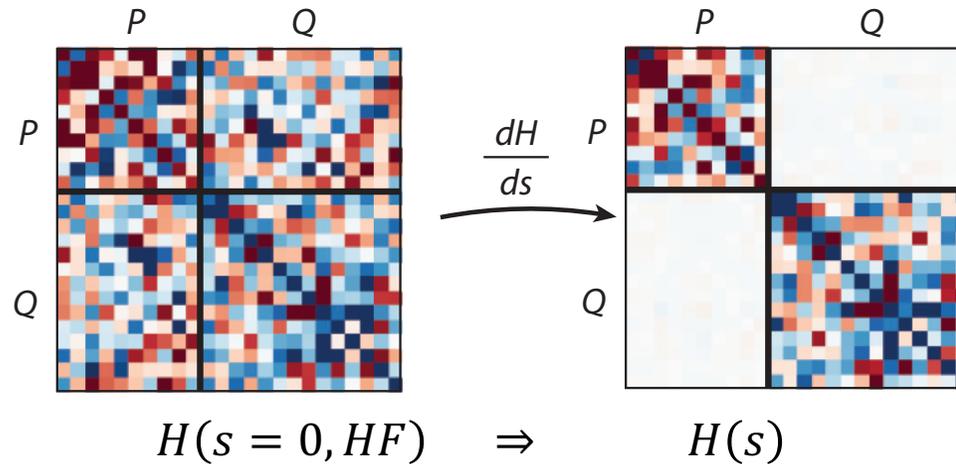
K. Tsukiyama, S. K. Bogner, and A. Schwenk, PRL **106**, 222502 (2011).  
 K. Tsukiyama, S. K. Bogner, and A. Schwenk, PRC **85**, 061304 (2012).

IMSRGflow:

$$\frac{dH(s)}{ds} = [\eta(s), H(s)],$$

$H(s)$  {  
 0-body: scalar,  $E_0(s) = \langle H(s) \rangle$   
 1-body:  $(n,n)$  matrix ( $n = \#$  of sps, block structure)  
 2-body:  $(\mathbf{d}_i, \mathbf{d}_i)$  matrix ( $i$  is label of  $\{J, P, T_z\}$ )

NO2B  $\rightarrow$  IMSRG(2)



under HF reference state,  
 divide single particle states as

P: hole

Q: particle

S.R.Stroberg et al., Annu. Rev. Nucl. Part. Sci. 2019. 69:307–62 (2019)

“off-diagonal” component, P-Q  $\lim_{s \rightarrow \infty} H^{od}(s) = 0$

“diagonal” component, P-P, Q-Q  $\lim_{s \rightarrow \infty} H^d(s) = H_{eff.} \rightarrow$  better  $E_0(s)$

adopting a certain generator  $\eta$  to achieve this “decoupling”

$$\begin{aligned}
 H(s) = U^\dagger(s) H U(s) &\Leftrightarrow \frac{dH(s)}{ds} = [\eta(s), H(s)], \\
 O(s) = U^\dagger(s) O U(s) &\Leftrightarrow \frac{dO(s)}{ds} = [\eta(s), O(s)],
 \end{aligned}$$

formulation with Magnus expansion

$$U(s) = e^{\Omega(s)}$$

$$H(s + ds) = e^{\eta(s)ds} e^{\Omega(s)} H(0) e^{-\Omega(s)} e^{-\eta(s)ds}$$

$$e^{\Omega(s+ds)} \equiv e^{\eta(s)ds} e^{\Omega(s)}$$

$$\Omega(s + ds) = \Omega(s) + \eta(s)ds + \frac{1}{2}[\eta(s), \Omega(s)]ds + \frac{1}{12}[\Omega(s), [\Omega(s), \eta(s)]]ds + \dots$$

explicit calculation of unitary transformation via Magnus operator  $\Omega$

- nested commutator (BCH formula)
- 1-step is enough (Euler method), nice for memory or I/O
- reuse  $U(s) = \exp(\Omega(s)) \rightarrow$  Any operators can be evolved simultaneously

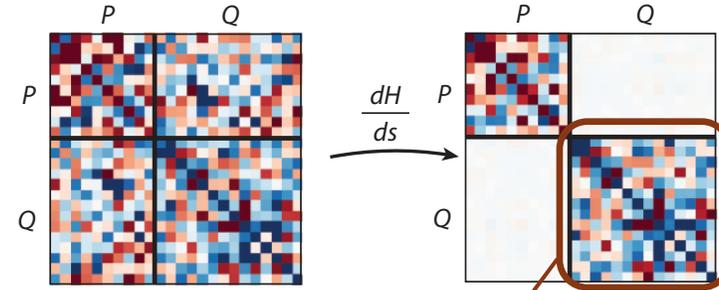
IMSRG: P: hole, Q: particle

“off-diagonal” component, P-Q  $\lim_{s \rightarrow \infty} H^{od}(s) = 0$

“diagonal” component, P-P, Q-Q  $\lim_{s \rightarrow \infty} H^d(s) = H_{\text{eff.}}$

adopting a certain generator  $\eta$  to achieve this “decoupling”

S.R.Stroberg, et al., PRL 118, 032502 (2017)  
 S.R.Stroberg et al., Annu. Rev. Nucl. Part. Sci. 2019. 69:307–62 (2019)

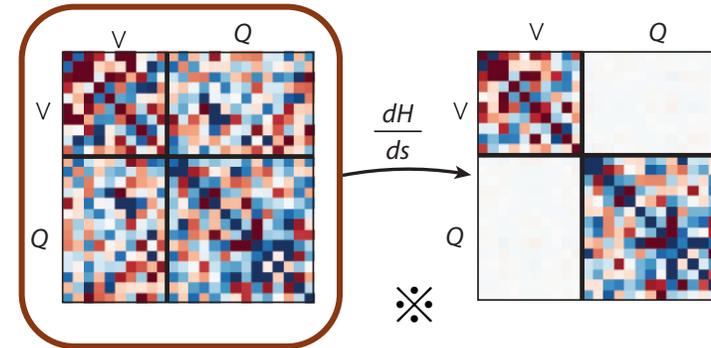


VS-IMSRG: P => valence, Q => q-space

$$\eta_{abij} = \frac{1}{2} \arctan \left( \frac{2\Gamma_{abij}}{f_{aa} + f_{bb} - f_{ii} - f_{jj} + G_{abij} + \Delta} \right),$$

$$G_{abij} = \Gamma_{abab} + \Gamma_{ijij} - (\Gamma_{aiai} + \Gamma_{bjbj} + [a \leftrightarrow b]).$$

denominator Delta: prescription for multi-shell interaction  
 see T. Miyagi et.al., PRC 102, 034320 (2020)



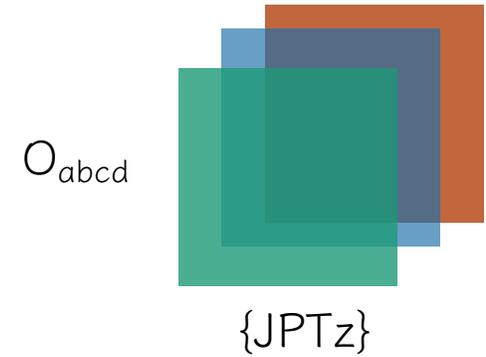
$H_v(s \rightarrow \infty) =$  Effective interactions for a valence space

(※ core is still considered in VS-IMSRG calculations, this figure is edited (by me) to focus on v-space)

One-body Hamiltonians  $\langle a|V|b\rangle$

Two-body Hamiltonians  $\langle ab|V|cd\rangle_{\{JPTz\}}$

Three-body ... (ignored in IMSRG(2) truncation)



$$111 \quad [A^{(1)}, B^{(1)}]^{(1)} = \sum_{ij} \sum_a :a_i^\dagger a_j: (A_{ia}B_{aj} - B_{ia}A_{aj})$$

$$110 \quad [A^{(1)}, B^{(1)}]^{(0)} = \sum_{ij} A_{ij}B_{ji}(n_i - n_j)$$

$$122 \quad [A^{(1)}, B^{(2)}]^{(2)} = \frac{1}{4} \sum_{ijkl} \sum_a :a_i^\dagger a_j^\dagger a_l a_k: \{ (1 - P_{ij})A_{ia}B_{ajkl} - (1 - P_{kl})A_{ak}B_{ijal} \}$$

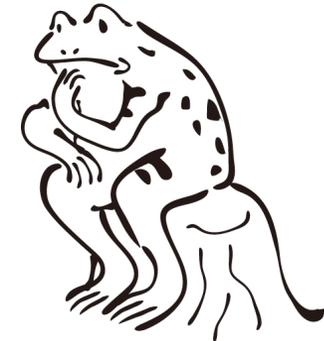
$$121 \quad [A^{(1)}, B^{(2)}]^{(1)} = \sum_{ij} \sum_{ab} :a_i^\dagger a_j: \{ (n_a - n_b)A_{ab}B_{biaj} \}$$

$$222 \quad [A^{(2)}, B^{(2)}]^{(2)} = \frac{1}{4} \sum_{ijkl} \sum_{ab} :a_i^\dagger a_j^\dagger a_l a_k: \left\{ \frac{1}{2} (A_{ijab}B_{abkl} - B_{ijab}A_{abkl})(1 - n_a - n_b) \right. \\ \left. + (n_a - n_b)(1 - P_{ij} - P_{kl} + P_{ij}P_{kl})A_{aibk}B_{bjal} \right\}$$

$$221 \quad [A^{(2)}, B^{(2)}]^{(1)} = \frac{1}{2} \sum_{ij} \sum_{abc} :a_i^\dagger a_j^\dagger: (A_{ciab}B_{abcj} - B_{ciab}A_{abcj})(\bar{n}_a \bar{n}_b n_c + n_a n_b \bar{n}_c)$$

$$220 \quad [A^{(2)}, B^{(2)}]^{(0)} = \frac{1}{4} \sum_{ijkl} n_i n_j \bar{n}_k \bar{n}_l (A_{ijkl}B_{kl ij} - B_{ijkl}A_{kl ij})$$

How can we accelerate...



c.f. talk by

- Jacob Davison, TRIUMF workshop 2023
- H. Hergert, INT Program 21r-1c "Tensor Networks in Many Body and Quantum Field Theory", 2023

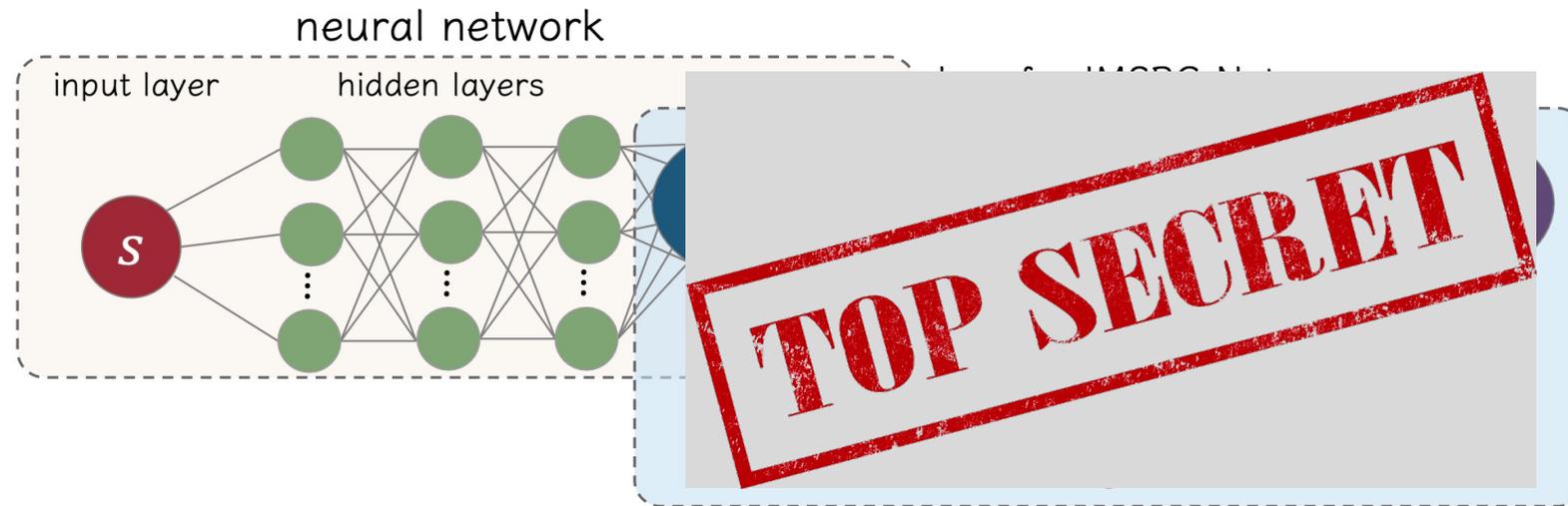
One may expect... “If we feed neural networks many data, they will learn underlying law”.

It is usually not the case. We need some “inductive biases” or more constraints.

## IMSRG-Net: A machine-learning based solver for In-Medium Similarity Renormalization Group

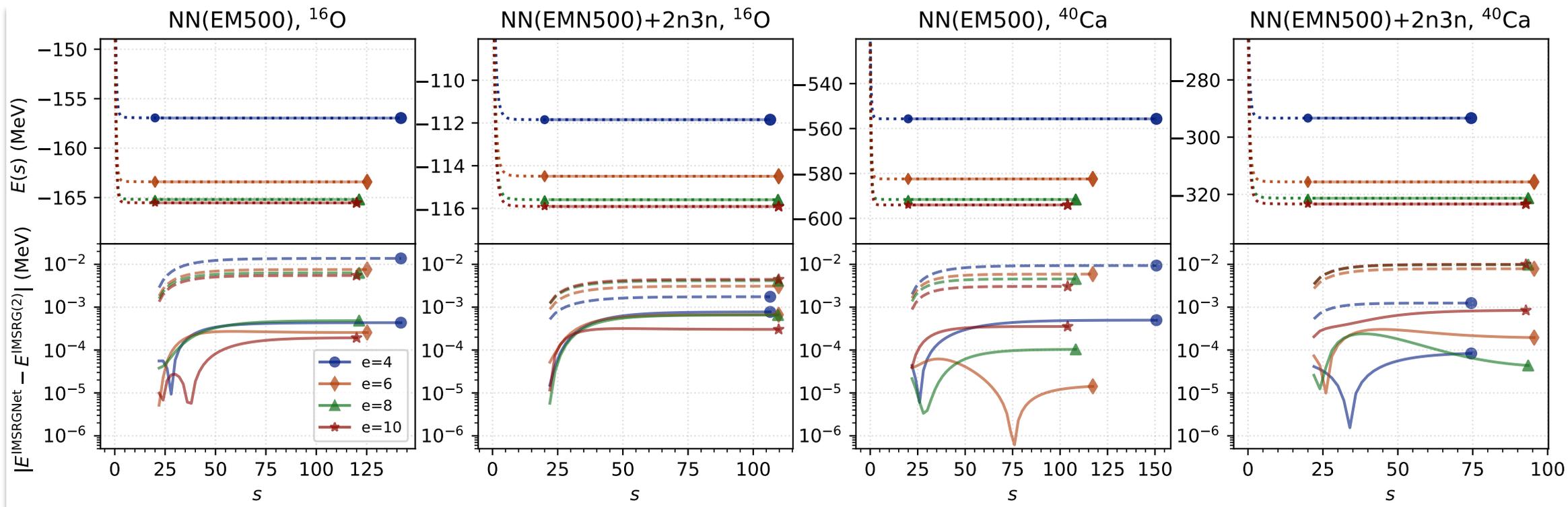
Sota Yoshida<sup>1,\*</sup>

<sup>1</sup>*Institute for Promotion of Higher Academic Education,  
Utsunomiya University, Mine, Utsunomiya, 321-8505, Japan*



hopefully, it will be available on arXiv&GitHub within a couple of ~~days~~ weeks !!

Aim is to predict  $E(s)$  or  $O(s)$  at  $s=\infty$  giving converged IMSRG(2) results



Upper panels: g.s. energy, dotted = IMSRG(2), solid = IMSRG-Net using 10 points around  $s=20$   
 Lower panels: energy diff., solid = IMSRG-Net, dashed = naïve(?) ANN

Symbols:  $s = 20$  or the point giving converged IMSRG(2)

prediction errors are less than 1 keV ( $\sim 0.5$  keV level)

target	interaction	$e_{\max}$	Energy (MeV)				$R_{\text{ch}}$ (fm)			
			$s = 20$		$s = \infty$		$s = 20$		$s = \infty$	
			IMSRG(2)	IMSRG-Net	IMSRG(2)	IMSRG-Net	IMSRG(2)	IMSRG-Net	IMSRG(2)	IMSRG-Net
$^{16}\text{O}$	EM500	4	-156.9474	-156.9474	-156.9611	-156.9607	2.2578	2.2578	2.2612	2.2610
		6	-163.4079	-163.4079	-163.4153	-163.4150	2.2526	2.2526	2.2547	2.2546
		8	-165.1876	-165.1875	-165.1932	-165.1927	2.2482	2.2482	2.2499	2.2497
		10	-165.5309	-165.5309	-165.5359	-165.5357	2.2469	2.2469	2.2485	2.2484
	EMN500+2n3n	4	-111.8453	-111.8453	-111.8470	-111.8462	2.3600	2.3600	2.3607	2.3605
		6	-114.4895	-114.4895	-114.4925	-114.4918	2.3681	2.3681	2.3692	2.3690
		8	-115.5894	-115.5894	-115.5930	-115.5925	2.3735	2.3735	2.3748	2.3746
		10	-115.9040	-115.9040	-115.9079	-115.9082	2.3751	2.3751	2.3765	2.3765
$^{40}\text{Ca}$	EM500	4	-555.6791	-555.6791	-555.6884	-555.6879	2.5947	2.5947	2.5959	2.5958
		6	-582.4293	-582.4293	-582.4350	-582.4351	2.5960	2.5960	2.5967	2.5967
		8	-591.5783	-591.5782	-591.5822	-591.5821	2.5915	2.5915	2.5920	2.5920
		10	-594.0215	-594.0215	-594.0242	-594.0246	2.5890	2.5890	2.5894	2.5895
	EMN500+2n3n	4	-293.3474	-293.3474	-293.3486	-293.3487	2.8579	2.8579	2.8581	2.8581
		6	-315.6334	-315.6334	-315.6411	-315.6409	2.9082	2.9082	2.9089	2.9088
		8	-321.3233	-321.3232	-321.3320	-321.3319	2.9201	2.9200	2.9209	2.9208
		10	-323.3519	-321.3520	-323.3605	-323.3613	2.9252	2.9252	2.9260	2.9260

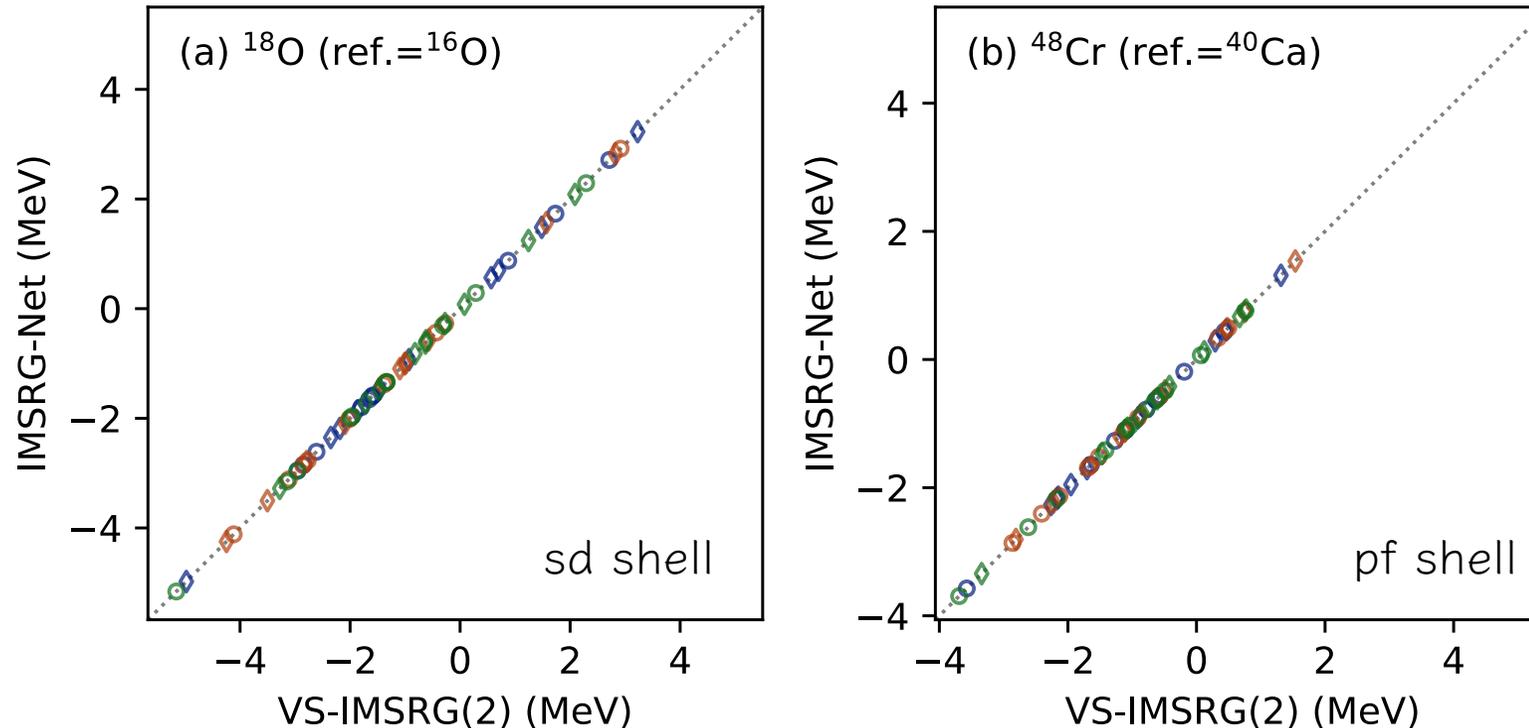
prediction errors of IMSRG-Net are much smaller than the *residuals*

to be gained through the rest IMSRG flow ( $s = \text{from } 20 \text{ to } \infty$ )

80-150

IMSRG-Net for a valence space

- same architecture and training strategy
- trained w/ earlier VS-IMSRG(2) flow



shell-model results agree in typically  $\leq 1$  keV level !!

✂ Some show  $\sim 10$  keV error, attributed not to IMSRG-Net, but numerical instability of VS-IMSRG



# Transition densities

The one-body transition densities (OBTDs) is given as

$$\text{OBTD}(fi; j_a j_b; \lambda) \equiv \frac{1}{\sqrt{2\lambda + 1}} \langle \psi_{J_f M_f} | [c_{j_a}^\dagger \otimes \tilde{c}_{j_b}]^{(\lambda)} | \psi_{J_i M_i} \rangle, \quad (\text{A1})$$

where we introduced  $\tilde{c}_{j_b} \equiv (-1)^{j_b - m_b} c_{j_b}$ , and  $\langle || \cdot || \rangle$  means taking the so-called reduced matrix element, and the notation  $[\cdot \otimes \cdot]^{(\lambda)}$  is for the rank- $\lambda$  irreducible tensor operators. For more details on the tensor algebra, see e.g., [1, 2]. Since we are interested in the  $\lambda = 0$  (scaler in terms of irreducible tensor operator) and the diagonal ( $f = i$ ,  $J_f = J_i$ ,  $M_f = M_i$ ) component, which contributes to  $\tilde{H}$  in Eq. (13) in the main text.  $\overline{\text{OBTD}}$  for the  $k$ -th single particle state is defined as

$$\overline{\text{OBTD}}_k \equiv \sqrt{\frac{2j_k + 1}{2J_i + 1}} \text{OBTD}(ii; j_k j_k; 0) = \langle \psi_{J_i M_i} | N_k | \psi_{J_i M_i} \rangle, \quad (\text{A2})$$

where  $N_k$  is the occupation number of the  $k$ -th orbital, and the factor  $\sqrt{(2j_k + 1)/(2J_i + 1)}$  is introduced to make  $\overline{\text{OBTD}}_k$  identical with the occupation number of  $k$ -th orbital.

The two-body transition densities (TBTDs) are defined as

$$\begin{aligned} & \text{TBTD}(fi; abcd; J_{ab} J_{cd}; \lambda) \\ & \equiv \frac{1}{\sqrt{2\lambda + 1}} \langle \psi_{J_f M_f} | [A^\dagger(ab; J_{ab} M_{ab}) \otimes \tilde{A}(cd; J_{cd} M_{cd})]^{(\lambda)} | \psi_{J_i M_i} \rangle, \end{aligned} \quad (\text{A3})$$

$$\tilde{A}(cd; J_{cd} M_{cd}) \equiv (-1)^{J_{cd} + M_{cd}} A(j_c j_d; J_{cd} - M_{cd}), \quad (\text{A4})$$

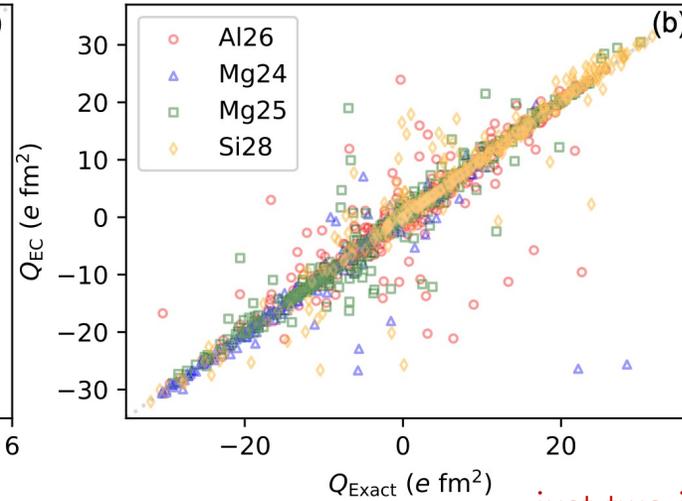
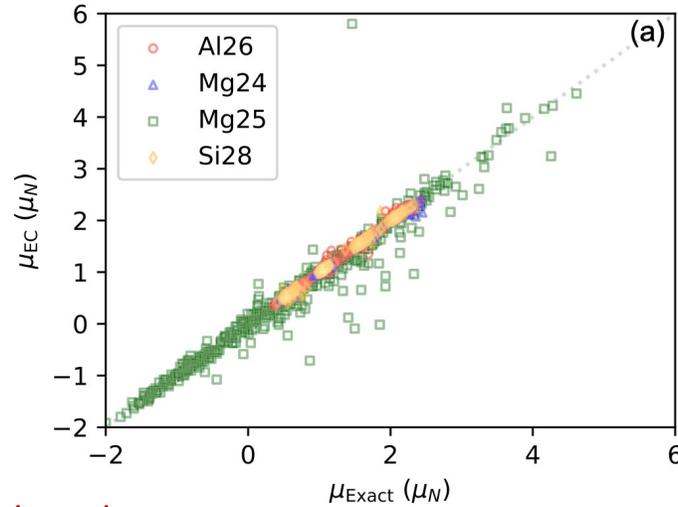
where  $A^\dagger$  and  $A$  are the same as in Eqs. (5-6). For the factorization in Eq. (13). the  $\overline{\text{TBTD}}$  for a two-body interaction  $V_J(abcd)$  is defined as follows

$$\overline{\text{TBTD}} \equiv \sqrt{\frac{2J_{ab} + 1}{2J_i + 1}} \text{TBTD}(fi; abcd; J_{ab} J_{ab}; 0), \quad (\text{A5})$$

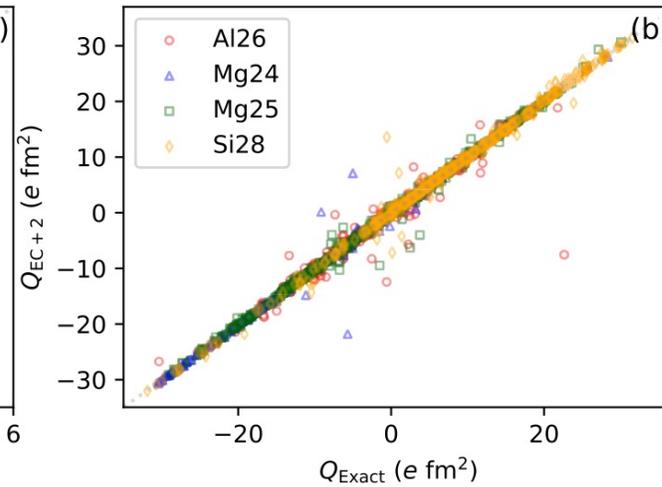
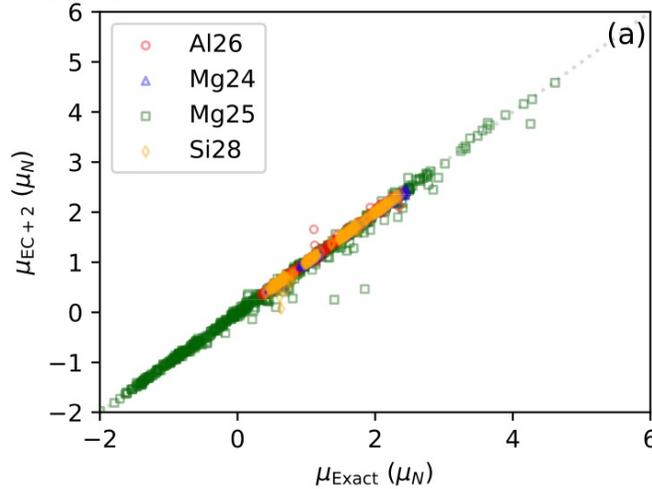
where only the term with  $\lambda = 0$ ,  $J_{cd} = J_{ab}$ ,  $M_{cd} = M_{ab}$ ,  $J_f = J_i$ ,  $M_f = M_i$  is needed due to the symmetry.

# $\mu$ & Q moments

magnetic moments & quadrupole moments for the lowest states



just two iterations



just two iterations

# Why Julia ?

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Since 2012:

Becoming popular in physics, DS, Machine Learning, etc.

- MIT LICENSE
  - Multiple dispatch
  - Dynamically typed
  - JIT(Just-In-Time) compilation by LLVM
  - Fast as C++/Fortran
  - Macros like Lisp
  - Package manager
  - Easy to call Python, C, Fortran, etc.
- High readability and productivity like Python
- High performance like C++/Fortran

If you are “greedy”, you should consider to use Julia 😊

# NuclearToolkit.jl: How to start



## NuclearToolkit.jl

Home

- Installation and example
- Package features
- Building blocks
- Optional parameters
- Issues/Pull requests

HowToUse

FileFormat

Optional parameters

Contributing to NuclearToolkit

References

- ChiEFTint
- HartreeFock
- IMSRG
- ShellModel

Version dev

Home [Edit on GitHub](#)

## NuclearToolkit

Julia Toolkit for nuclear structure calculations

### Installation and example

First, prepare Julia environment  $v \geq 1.7.0$ .

Second, add the package in Pkg mode

```
 julia>]add NuclearToolkit 
```

You can try the package by 1. or 2.:

- clone the repository and run `test/sample_script.jl` in the repository like `$ julia -t 8 sample_script.jl` This performs:
  - calculating NN potential from Chiral EFT
  - HFMBPT(3) and IMSRG/VS-IMSRG(2) calculation with it
  - shell-model calculations with the effective interaction derived by VS-IMSRG

An expected results using the latest dev branch can be found [here](#).

- Try sample codes in [HowToUse](#) page.

Please make sure to use the latest version of the package. Update can be done with

```
 julia>]up NuclearToolkit 
```

In the Julia REPL, you can see the UUIDs and versions of the installed packages

```
 julia>using Pkg
 julia>Pkg.status()
```

### Package features

NuclearToolkit.jl provides a self-contained set of nuclear structure calculation codes covering from nuclear