

# Eigenvector continuation for the anharmonic oscillator and the pairing Hamiltonian

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**European Research Council**

Established by the European Commission



- For Hamiltonians  $\hat{H}(c)$  with continuous parametric dependence

$$\hat{H}(c) = \hat{H}_0 + c \cdot \hat{H}_1$$

- Based on **analytical continuation** of eigenvalues and eigenvectors outside original domain  
*Frame et al., PRL (2018)*

- Predictions obtained by solving the generalized eigenvalue problem

$$HX = \epsilon NX$$

- With matrix elements

$$H_{pq} = \langle \Psi^{(p)} | \hat{H}(c) | \Psi^{(q)} \rangle$$

$$N_{pq} = \langle \Psi^{(p)} | \Psi^{(q)} \rangle$$

- For set of training vectors  $|\Psi^{(p)}\rangle$



This talk: Two different applications

- ▷ As a PT resummation tool
  - ▷ State corrections as EC basis  
e.g. Demol *et al.* PRC (2020)  
Sakar and Lee PRL (2021)
  - ▷ Here for the anharmonic oscillator (AHO) with high-order state corrections
  - ▷ For the pairing Hamiltonian with low-order state corrections
- ▷ As an emulator tool
  - ▷ Solution known for some  $c_i$
  - ▷ Exact solutions used as EC basis  
e.g. König *et al.* PLB (2020)  
Furnstahl *et al.* PLB (2020)  
Drischler *et al.* PLB (2021)
  - ▷ For the pairing Hamiltonian with truncated training vectors

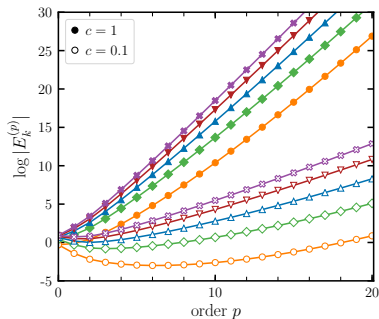
- Hamiltonian of a **harmonic** oscillator with a **quartic** perturbation (in natural units)

$$\hat{H}(c) = \frac{1}{2}\hat{p}^2 + \frac{1}{2}\hat{x}^2 + c\hat{x}^4$$

- Perturbation expansion is well known to **diverge**  
Bender, Wu, PR (1969) , PRD (1973)
- Expression in terms of ladder operators

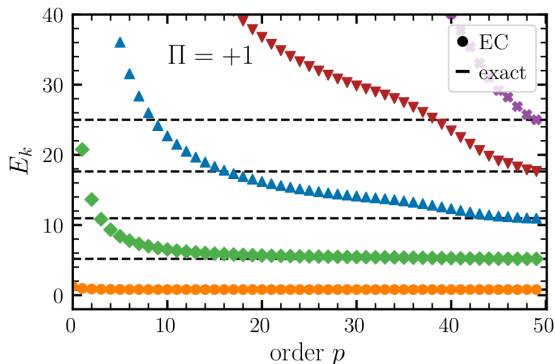
$$\hat{H} = \hat{a}^\dagger \hat{a} + \frac{1}{2} + c \left( \frac{\hat{a}^\dagger + \hat{a}}{\sqrt{2}} \right)^4$$

- Pentadiagonal matrix



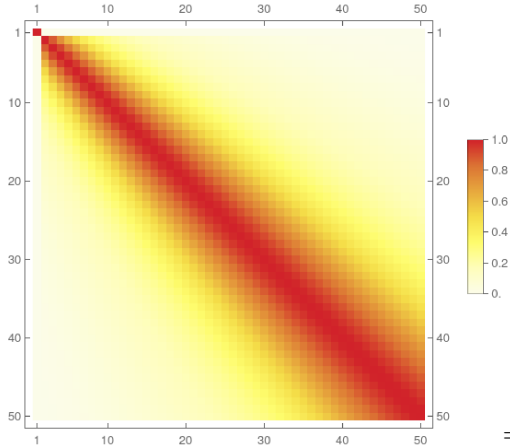
Companys Franzke et al., PLB (2022)

- EC basis generated by PT expansion
  - ▷ Reference state: ground state of HO
- Most rapid convergence for ground state
- Goal: Extend EC framework for excited states
- Also sequences for **excited states**
  - ▷ Solving the **same** generalized eigenvalue problem
  - ▷ Only eigenstates of the same symmetries can be accessed

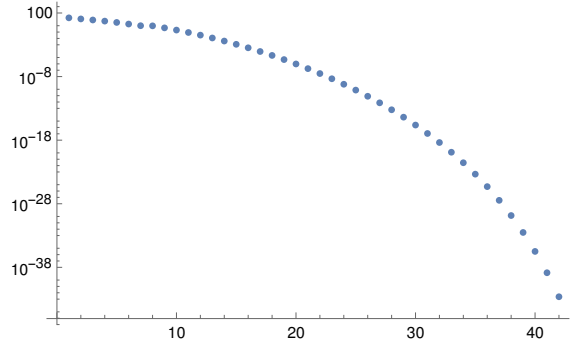


Companys Franzke et al., PLB (2022)

# Computational difficulties



Singular Value



⇒ Norm matrix becomes singular for high orders



- EC solves generalized eigenvalue problem  $\mathbf{H}^V X = \epsilon \mathbf{N}^V X$  with EC basis  $V$  and

$$\mathbf{H}^V = V^T H V$$

$$\mathbf{N}^V = V^T V$$

- Equivalent to eigenvalue problem  $(\mathbf{N}^V)^{-1} \mathbf{H}^V X = \epsilon X$



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- Can be shown to be equivalent to problem in orthogonal basis  $W$





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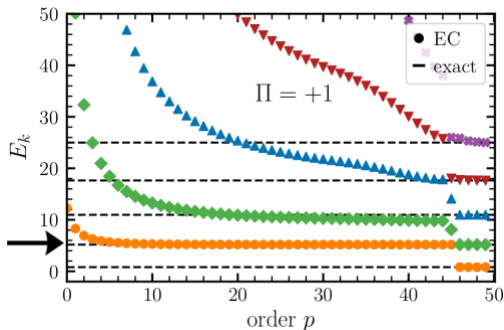
- Can be shown to be equivalent to problem in orthogonal basis  $W$

- Let  $T$  and  $T^{-1}$  be the transformation matrices, such that  $V = W \cdot T$ .

$$\implies (\mathbf{N}^V)^{-1} \mathbf{H}^V = T^{-1} \mathbf{H}^W T$$

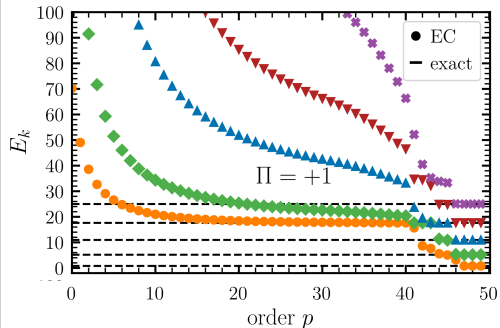
- Solve  $\mathbf{H}^W Y = \epsilon Y$  for  $Y = T X$  instead

- EC basis generated by PT expansion
  - ▷ reference state: **excited** state of HO
- Reference state sequence converges quickly
- Other sequences do not converge to exact eigenvalues until high order
  - ▷ Jump at high orders
  - ▷ Variational principle



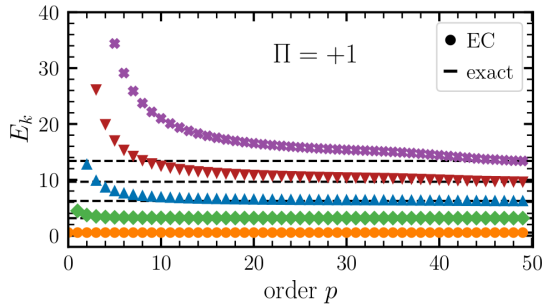
Comanys Franzke et al., PLB (2022)

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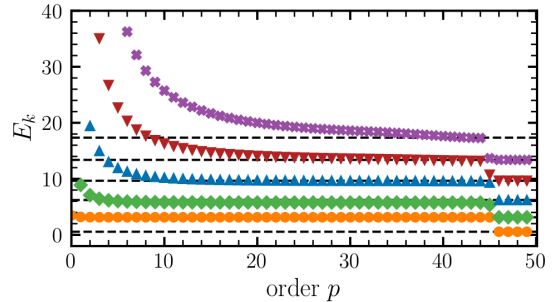


# EC for coupling strength ( $c = 0.1$ )

Reference state: ground state



Reference state: first excited state



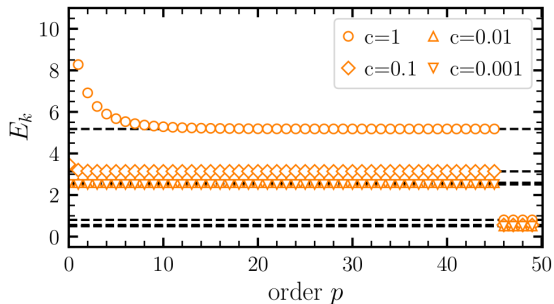
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Same general trend as for  $c = 1$

Converges faster

# Variation of the coupling strength for EC

- EC basis from PT for first excited state for different  $c$ 
  - ▷ Only lowest sequence shown
- Jump **independent of  $c$** 
  - ▷ EC spans the same subspace

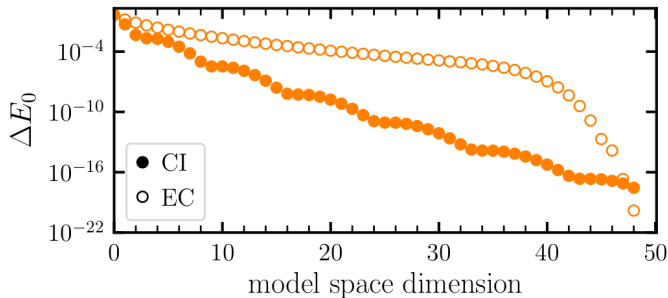


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# EC compared to configuration interaction

- Configuration interaction (CI) basis consists of the eigenstates of the HO
- Diagonalization in truncated HO basis
- PT strongly divergent
  - ▷ EC performs worse than CI
  - ▷ EC still converges

⇒ PT chooses poor basis



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# Summary AHO

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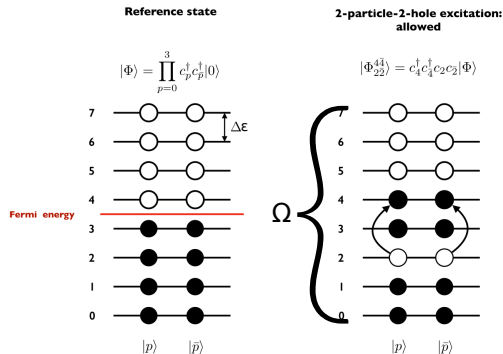
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- EC sequence converges quickly for the reference state
  - EC for the ground-state reference state also produces sequences converging to excited states
  - Reference state EC sequence converges faster than same state in ground-state EC
    - ▷ Simultaneous assessment of multiple excited states challenging
  - The EC order where the jump occurs is independent of  $c$
- ⇒ EC provides a robust framework to extract ground and excited states

- Pairing Hamiltonian for model-space size  $\Omega$  and pair states  $p$  and  $\bar{p}$

$$\hat{H}_{\text{pairing}} \equiv \sum_p^{\Omega} \epsilon_p (c_p^\dagger c_p + c_{\bar{p}}^\dagger c_{\bar{p}}) - g \sum_{pq} c_p^\dagger c_{\bar{p}}^\dagger c_{\bar{q}} c_q,$$

⇒ **exactly solvable** due to Richardson  
without large-scale diagonalization  
e.g. Richardson *et al.* PL (1964)



Courtesy of Alexander Tichai

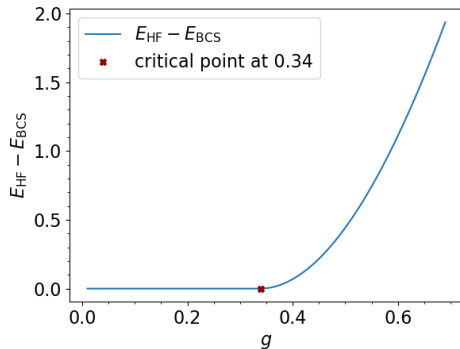


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- Phase transition to superfluid state for  $g > g_{\text{crit}}$



- PT:  $\hat{H}_0 = \hat{H}^{(0)} + \hat{H}^{(1)}$  and  $\hat{H}_1 = \hat{H}^{(2)}$
- $\hat{H}^{(k)}$  normal ordered  $k$ -body part of the Hamiltonian

$$E^{(0)}|\Phi\rangle = H_0|\Phi\rangle = (2 \sum_i \epsilon_i - gN_{\text{occ}})|\Phi\rangle$$

$$|\Psi^{(1)}\rangle = \frac{1}{2} \sum_{ai} \frac{g}{f_i - f_a} |\Phi_{ii}^{a\bar{a}}\rangle$$

$$E^{(2)} = -\frac{1}{2} \sum_{ai} \frac{g^2}{f_i - f_a}$$

- Where  $f_p = \epsilon_p - n_p g$

- Reference state:

$$|\Phi\rangle \equiv \prod_{i=1}^{N_{\text{occ}}} c_i^\dagger c_i^\dagger |0\rangle.$$

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- Reference state:

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- EC:

$$\mathbf{H} = \begin{pmatrix} \langle \Psi^{(0)} | \hat{H} | \Psi^{(0)} \rangle & \langle \Psi^{(0)} | \hat{H} | \Psi^{(1)} \rangle \\ \langle \Psi^{(1)} | \hat{H} | \Psi^{(0)} \rangle & \langle \Psi^{(1)} | \hat{H} | \Psi^{(1)} \rangle \end{pmatrix} = \begin{pmatrix} E^{(0)} & E^{(2)} \\ E^{(2)} & \mathbf{H}_{11} \end{pmatrix}$$

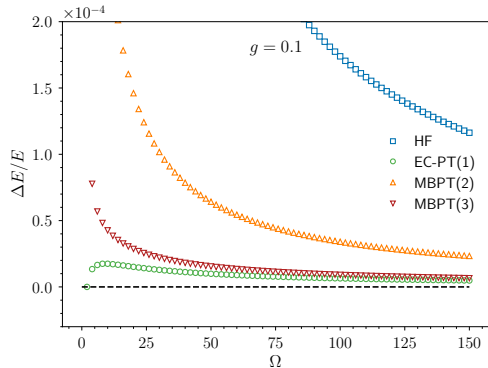
$$\mathbf{N} = \begin{pmatrix} \langle \Psi^{(0)} | \Psi^{(0)} \rangle & \langle \Psi^{(0)} | \Psi^{(1)} \rangle \\ \langle \Psi^{(1)} | \Psi^{(0)} \rangle & \langle \Psi^{(1)} | \Psi^{(1)} \rangle \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & \mathbf{N}_{11} \end{pmatrix}.$$

# EC as a resummation method of PT (2/2)

$$\mathbf{N}_{11} = \frac{1}{4} \sum_{abij} \frac{g^2 \langle \Phi_{ii}^{a\bar{a}} | \Phi_{ii}^{a\bar{a}} \rangle}{(f_i - f_a)(f_j - f_b)} = \frac{1}{4} \sum_{ai} \frac{g^2}{(f_i - f_a)^2}$$

$$\mathbf{H}_{11} = \frac{1}{4} \sum_{ai} \frac{g^2 (2 \sum_k \epsilon_k - g N_{\text{occ}})}{(f_i - f_a)^2} - \frac{1}{2} \sum_{ai} \frac{g^2}{f_i - f_a} - \frac{1}{4} \left( \sum_{abi} \frac{g^3}{(f_i - f_a)(f_i - f_b)} + \sum_{aij} \frac{g^3}{(f_i - f_a)(f_j - f_a)} \right)$$

- Scales polynomially with system size
- Similar in spirit to **Ekström and Hagen PRL (2019)**



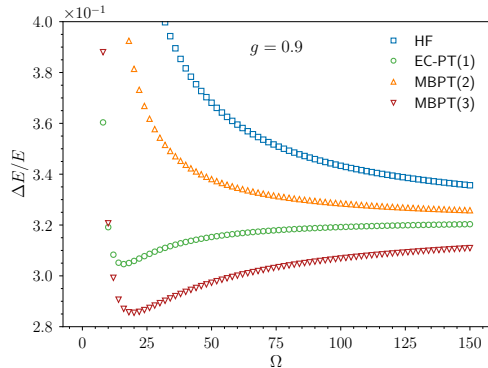
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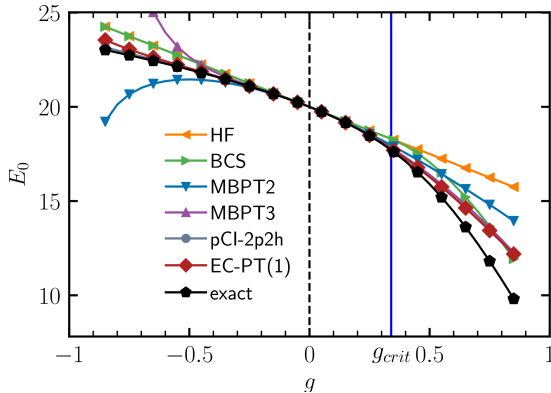
$$\mathbf{H}_{11} = \frac{1}{4} \sum_{ai} \frac{g^2 (2 \sum_k \epsilon_k - g N_{\text{occ}})}{(f_i - f_a)^2} - \frac{1}{2} \sum_{ai} \frac{g^2}{f_i - f_a} - \frac{1}{4} \left( \sum_{abi} \frac{g^3}{(f_i - f_a)(f_i - f_b)} + \sum_{aij} \frac{g^3}{(f_i - f_a)(f_j - f_a)} \right)$$

- Scales polynomially with system size
- Similar in spirit to **Ekström and Hagen PRL (2019)**
- Superfluidity cannot be captured by MBPT around normal state



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# EC compared to other methods

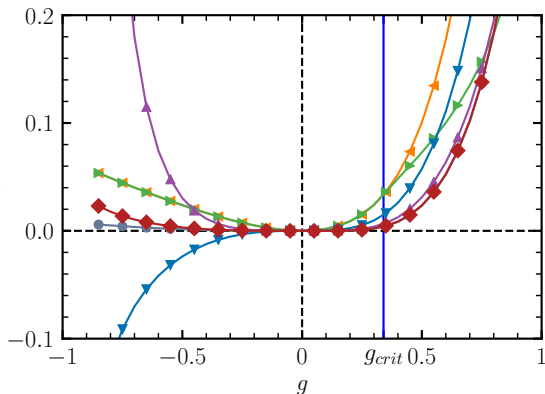
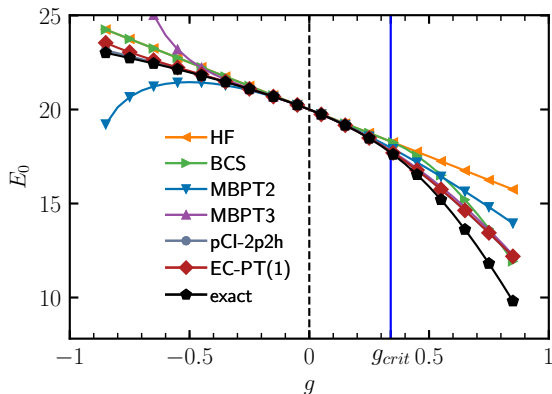


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$$E^{(2)} = -\frac{1}{2} \sum_{ai} \frac{g^2}{f_i - f_a} \text{ with } f_p = \epsilon_p - n_p g$$

- Singularity at  $g = -\Delta\epsilon = -1$
- pCI-2p2h and EC-PT(1) both are diagonalizations on 2p2h-spaces

# EC compared to other methods



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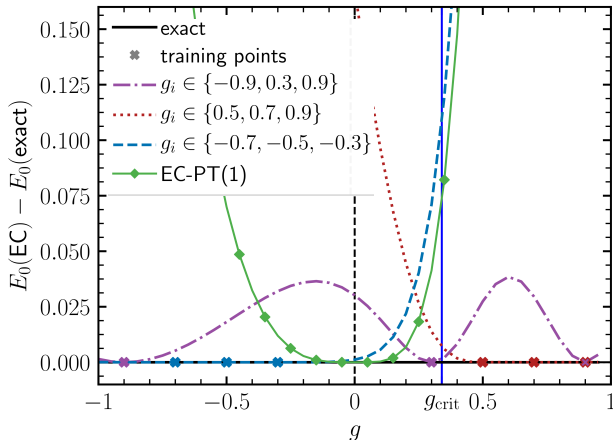
⇒ EC gives good approximations for large coupling range, although HF and MBPT(2) do not

- EC can be used to emulate Hamiltonians from training data
  - e.g. Frame *et al.* PRL (2018)
  - Baran and Nichita PRB (2023)
- Can be used to
  - ▷ Interpolate with training points from both sides of  $g_{crit}$
  - ▷ Extrapolate from only normal or superfluid training points
- Either exact or approximated training data can be used
  - ▷ Here the CI truncation of the training states is varied

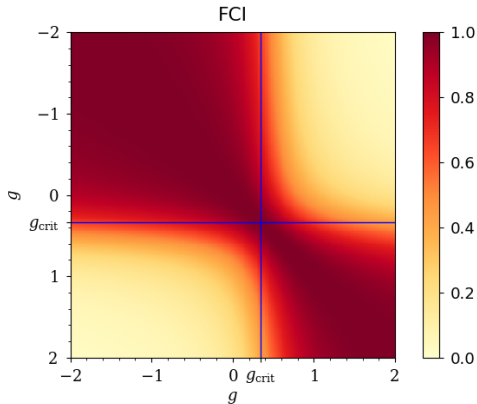


# Different EC applications for the pairing Hamiltonian

- EC from PT state corrections only good around  $g = 0$
- One-sided training points only approximate the same side well
- Training points from both sides of the interval give good results



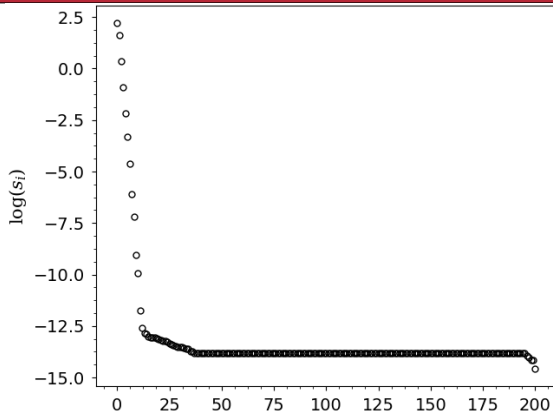
# Norm matrix for the pairing Hamiltonian



- Block matrix structure
- Dependent on critical  $g$
- Strongly suppressed overlap between normal and superfluid ground state

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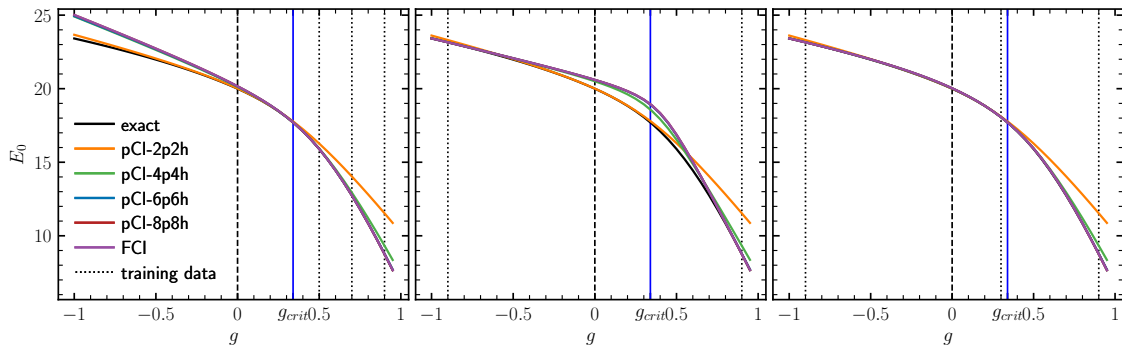
# Norm matrix for the pairing Hamiltonian



Companys Franzke, M.Sc. thesis (2023)

- Block matrix structure
- Dependent on critical  $g$
- Strongly suppressed overlap between normal and superfluid ground state
- Most singular values zero

# Variation of the CI truncation of the EC training points

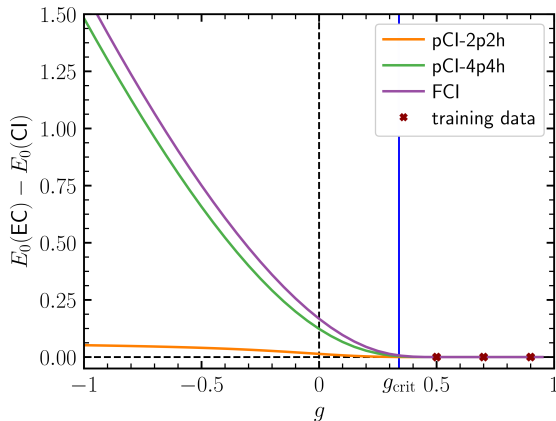


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⇒ 6p6h truncated CI gives basically the same results as full CI

# Eigenvector continuation compared to configuration interaction

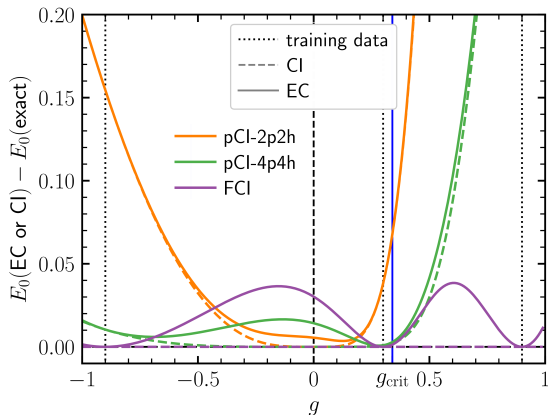
- EC approximates CI with lower truncation better



Companys Franzke et al., arXiv:2302.08373 (2023)

# Eigenvector continuation compared to configuration interaction

- EC approximates CI with lower truncation better
- Higher CI truncations give more accurate training vectors



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# Summary Pairing Hamiltonian

arXiv:2302.08373



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- Pairing Hamiltonian changes from normal to superfluid ground state with increasing coupling
- Eigenvector continuation with MBPT state correction only approximates solutions around  $g = 0$  well
- One-sided training points can only approximate coupling values from the same side well
- Overlap between eigenvectors from superfluid couplings with eigenvectors from normal couplings is small
- Eigenvector continuation approximates lower CI truncations better



- Design of many-body emulator for chiral Hamiltonians following the work of Ekström and Hagen PRL (2019)
- Hartree-Fock states as EC basis





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# Thank you for your attention!