Overview of nuclear parton distribution functions from LHC to EIC

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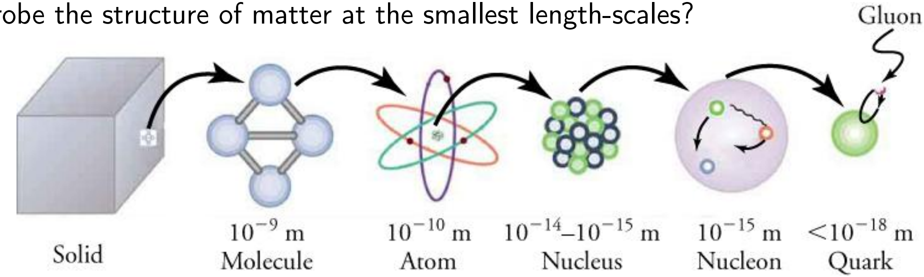
ESNT workshop

“Deciphering nuclear phenomenology across energy scales”

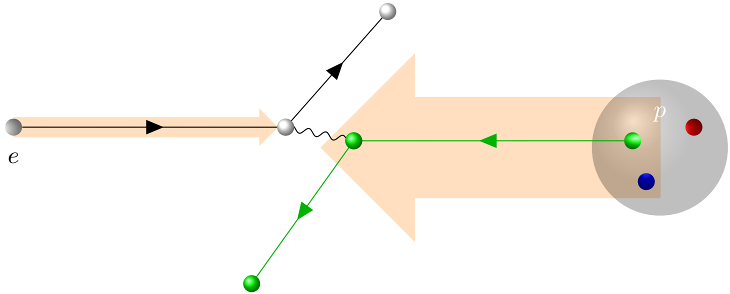
23 September 2022

1. What, how and why? / The age-old questions
2. Where are we now? / Recent developments
3. What we would like to know better? / More A and towards 3D

Q: How to probe the structure of matter at the smallest length-scales?



A: Hit it as hard as you can and see what comes out



But what is it that we are probing at sub-nucleon level? ➔ Need help from theory!

At the heart of it all: Collinear factorisation of QCD

The cross section for producing an inclusive final state $k + X$ can be described as a convolution of...

...**C**oefficient **F**unctions $\mathrm{d}\hat{\sigma}^{ij\to k+X'}$ which are calculable from perturbative QCD...

$$\mathrm{d}\sigma^{AB\to k+X}(Q^2) \overset{Q\gg\Lambda_{\mathrm{QCD}}}{=} \sum_{i,j,X'} f_i^A(Q^2) \otimes \mathrm{d}\hat{\sigma}^{ij\to k+X'}(Q^2) \otimes f_j^B(Q^2) + \mathcal{O}(1/Q^2)$$

...and **P**arton **D**istribution **F**unctions f_i^A, f_j^B which contain long-range physics and cannot be obtained by perturbative means...

...plus "**H**igher **T**wist" corrections which are suppressed at high enough momentum scale $Q \gg \Lambda_{\mathrm{QCD}}$

The PDFs $f_i^A(x, Q^2)$ are *universal*, process independent,

parent hadron or nucleus

factorisation scale

parton flavour

fraction of momentum carried by the parton

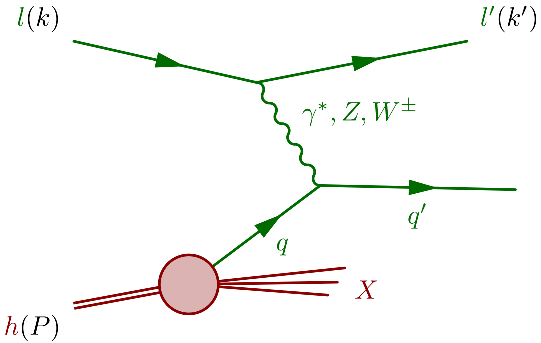
and obey the DGLAP equations $Q^2 \dfrac{\partial f_i^A}{\partial Q^2} = \sum_j P_{ij} \otimes f_j^A$

splitting functions

Mellin conv.

...this is the framework which every PDF analysis and application relies on and tests!

Deep inelastic scattering (DIS)



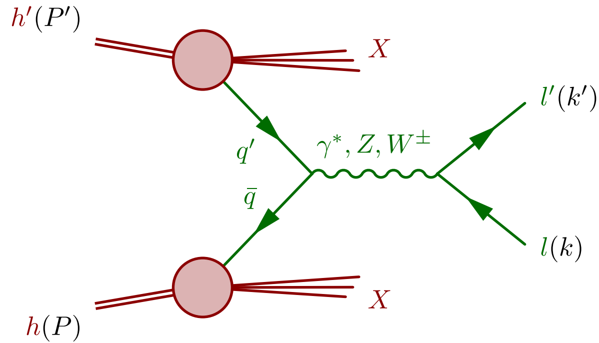
For the photon-mediated case:

$$\frac{\mathrm{d}^2\sigma^{\mathrm{DIS}}}{\mathrm{d}x\mathrm{d}Q^2} = \frac{\mathrm{d}^2\hat{\sigma}}{\mathrm{d}x\mathrm{d}Q^2}\sum_{i\in\{q,\bar{q}\}} e_i^2 f_i^h(x,Q^2) + {}_{\text{corrections}}^{\text{NLO}}$$

$$\left.\begin{aligned} Q^2 &= -(k-k')^2 \\ x &= \frac{Q^2}{2P\cdot(k-k')} \end{aligned}\right\}$$

← access scale and momentum-fraction dependence through external kinematics

Drell-Yan (DY)



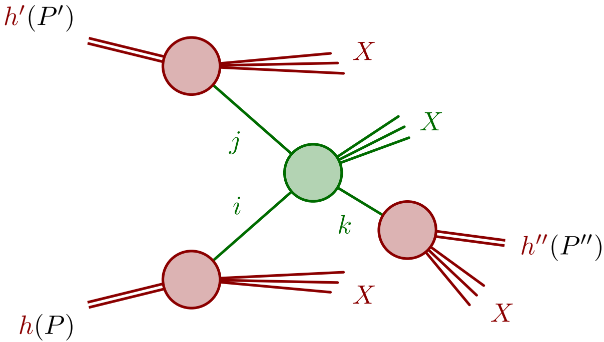
The photon-mediated case:

$$\frac{\mathrm{d}^2\sigma^{\mathrm{DY}}}{\mathrm{d}y\mathrm{d}M^2} = \frac{4\pi\alpha_{\mathrm{e.m.}}^2}{9M^4}\sum_{i\in\{q,\bar{q}\}} e_i^2 x_1 x_2 f_i^h(x_1,M^2) f_{\bar{i}}^{h'}(x_2,M^2) + {}_{\text{corrections}}^{\text{NLO}}$$

$$M^2 = (k+k')^2 = x_1 x_2 s$$

$$y = \frac{1}{2}\log\frac{(k_0+k_0')+(k_3+k_3')}{(k_0+k_0')-(k_3+k_3')} = \frac{1}{2}\log\frac{x_1}{x_2}$$

Hadron-production

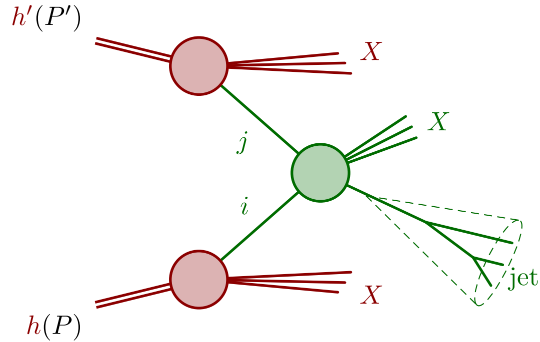


$$\sigma^{h+h'\to h''+X} = \sum_{i,j,k\in\{q,\bar{q},g\}} f_i^h \otimes f_j^{h'} \otimes \hat{\sigma}^{ij\to k+X} \otimes D_k^{h''}$$

Account for the hadronization effects with the *parton to hadron fragmentation functions* $D_k^{h''}$

➔ a source of uncertainty for PDF fits

Jet-production



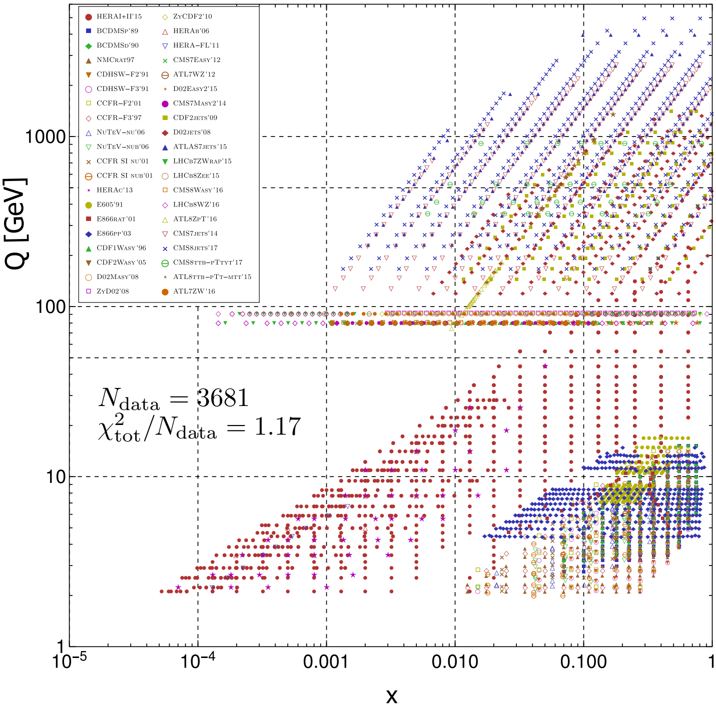
$$\sigma^{h+h'\to \text{jet}+X} = \sum_{i,j\in\{q,\bar{q},g\}} f_i^h \otimes f_j^{h'} \otimes \hat{\sigma}^{ij\to \text{jet}+X}$$

Additional complications:

- need an IR-safe definition of a jet
- non-perturbative corrections

Global analysis – a multi-experiment–multi-observable fit

Experimental data in CT18 PDF analysis



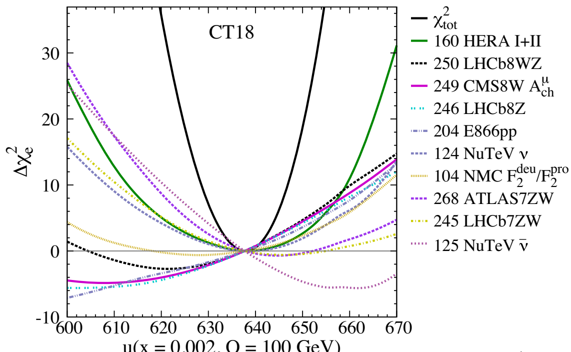
Multi-observable fit needed to constrain individual flavours, minimise:

Sum over data sets

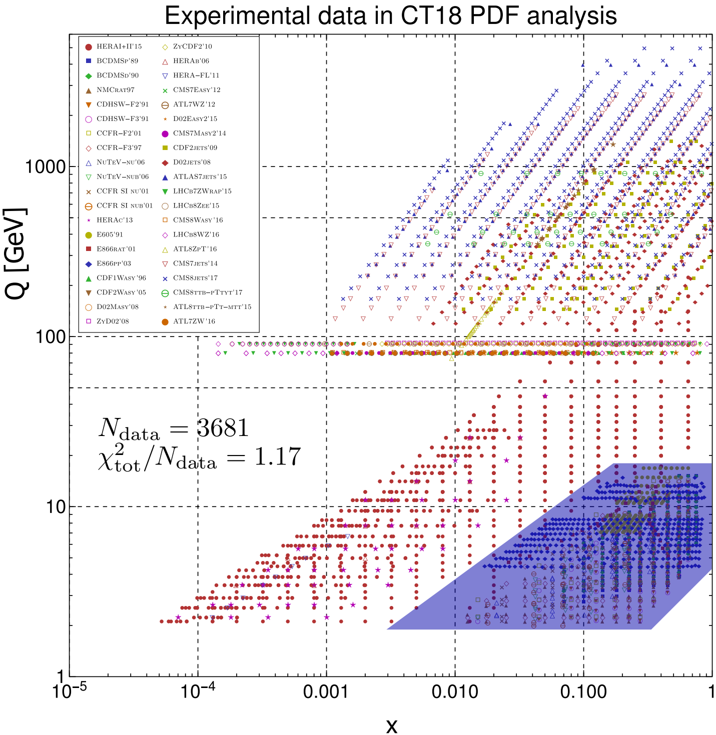
$$\chi^2_{\rm tot} = \sum_k (\underbrace{D_k}_{\text{data}} - \underbrace{T_k}_{\text{theory}})^T \underbrace{C_k^{-1}}_{\text{cov.}} (D_k - T_k)$$

Correlations important!

figs. from Hou et al., Phys. Rev. D 103 (2021) 014013



Global analysis – a multi-experiment–multi-observable fit

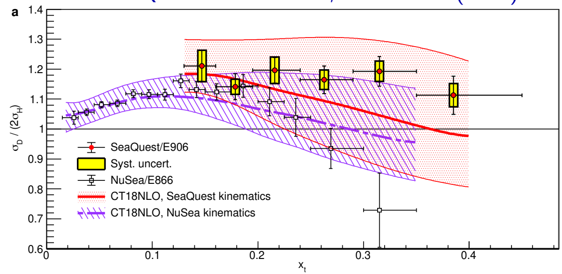


Fixed-target DIS and DY important in setting the large-x quark distributions (valence/sea and flavour separation)

New data still coming from Fermilab & JLab!

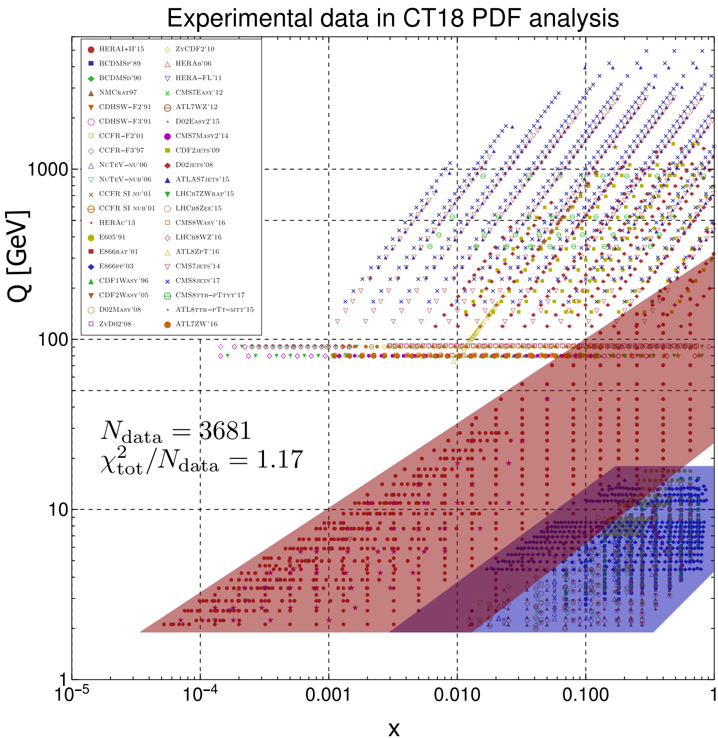
new!

SeaQuest Collaboration, Nature 590 (2021) 561



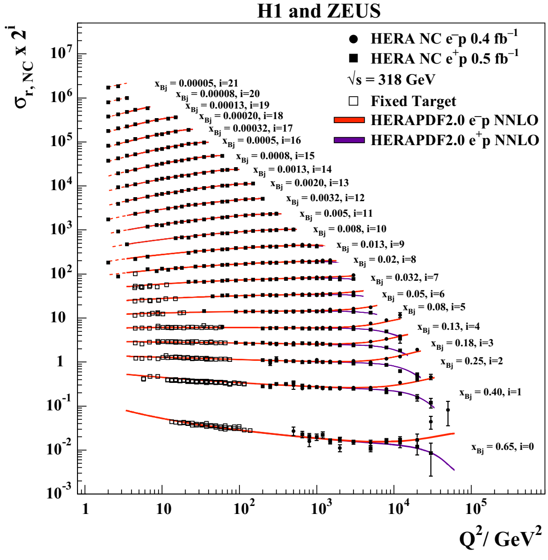
$$\left.\frac{\sigma^{pd}}{2\sigma^{pp}}\right|_{x_{\rm beam} \gg x_{\rm target}} \approx \frac{1}{2}\left[1 + \frac{\bar{d}\,(x_{\rm target})}{\bar{u}\,(x_{\rm target})}\right]$$

Global analysis – a multi-experiment–multi-observable fit

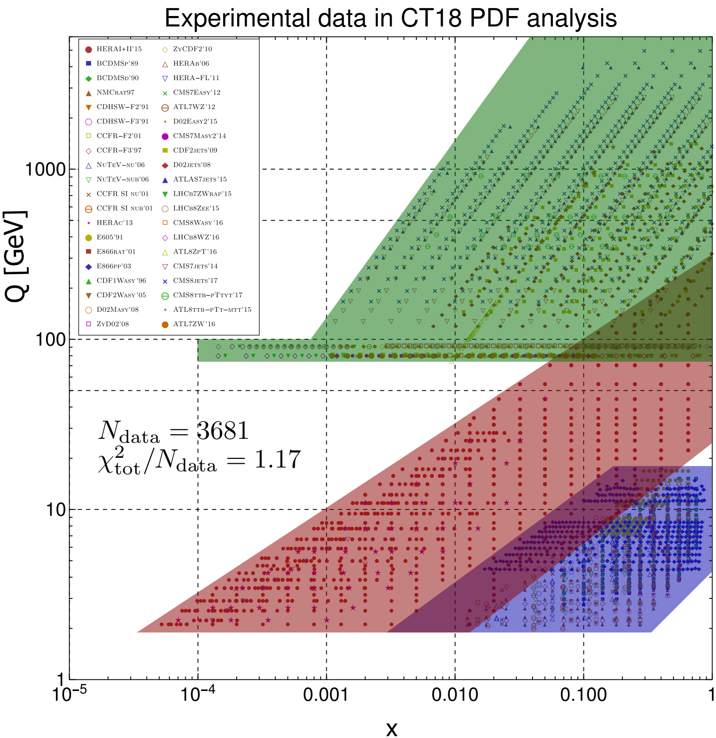


Collider DIS from HERA with large x, Q^2 lever arm ➔ gluons through DGLAP

H1 and ZEUS, Eur. Phys. J. C 75 (2015) 580

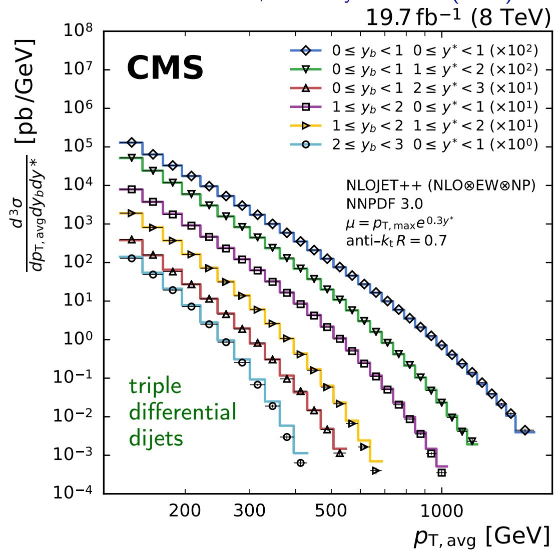


Global analysis – a multi-experiment–multi-observable fit

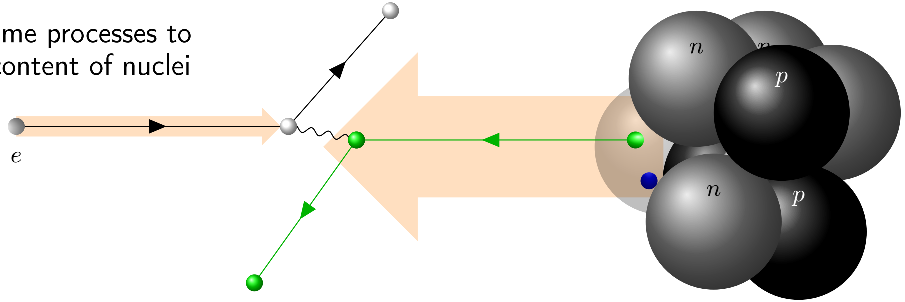


Hadron colliders give access to new processes → $W^\pm, Z$, jets, $t\bar{t}\ldots$

CMS Collaboration, Eur. Phys. J. C 77 (2017) 746



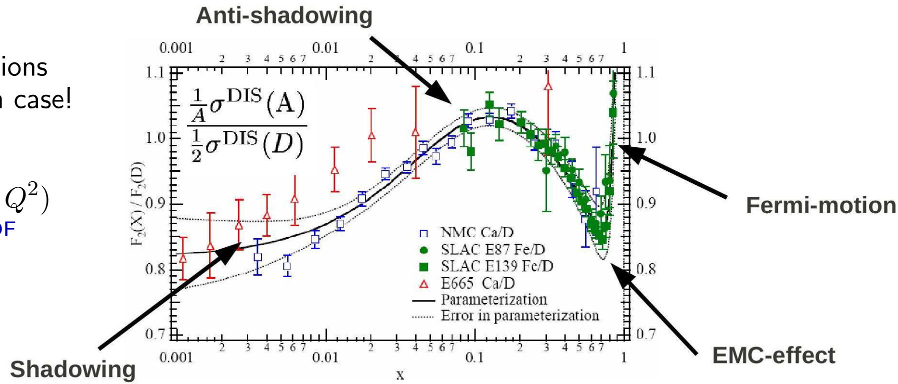
We can use the same processes to study the parton content of nuclei



Observe *significant* modifications compared to the free-nucleon case!

$$f_i^{p/A}(x, Q^2) \neq f_i^{p}(x, Q^2)$$

bound-proton PDF　　free-proton PDF



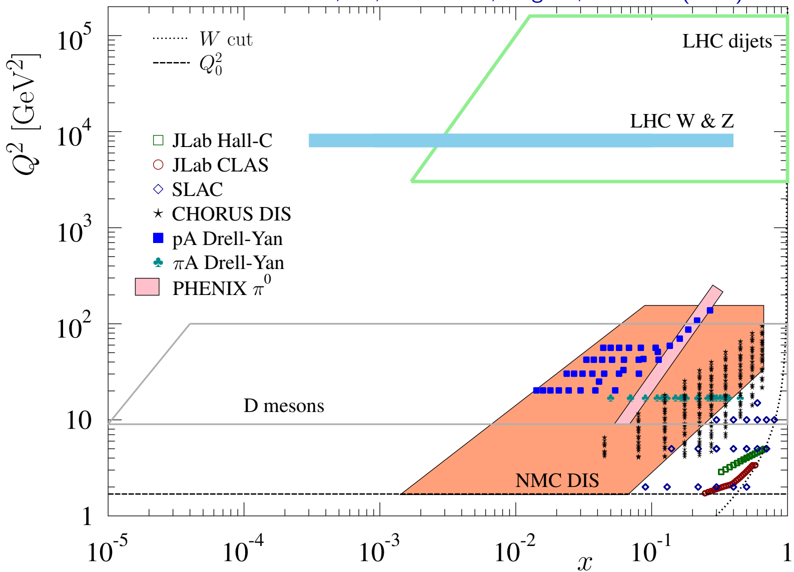
Nuclear PDFs from global analyses

Nuclear PDFs (nPDFs) are fitted with similar global analyses as their free-proton counterparts

- ➔ rely only to the QCD collinear factorisation
- ➔ model-agnostic way to study the nuclear effects

LHC is extending the x, Q^2 reach by orders of magnitude

Eskola, PP, Paukkunen, Salgado, EPJC 82 (2022) 413



Recent nPDF global fits

	KSASG20	TUJU21	EPPS21	nNNPDF3.0	nCTEQ15HQ*
Order in α_s	NLO & NNLO	NLO & NNLO	NLO	NLO	NLO
lA NC DIS	✓	✓	✓	✓	✓
νA CC DIS	✓	✓	✓	✓	
pA DY	✓		✓	✓	✓
πA DY			✓		
RHIC dAu π^0,$\pi^\pm$			✓		✓
LHC pPb π^0,$\pi^\pm$,$K^\pm$					✓
LHC pPb dijets			✓	✓	
LHC pPb HQ			✓ GMVFN	✓ FO+PS	✓ ME fitting
LHC pPb W,Z		✓	✓	✓	✓
LHC pPb γ				✓	
Q, W cut in DIS	1.3, 0.0 GeV	1.87, 3.5 GeV	1.3, 1.8 GeV	1.87, 3.5 GeV	2.0, 3.5 GeV
p_{T} cut in HQ,inc.-h	N/A	N/A	3.0 GeV	0.0 GeV	3.0 GeV
Data points	4353	2410	2077	2188	1496
Free parameters	9	16	24	256	19
Error analysis	Hessian	Hessian	Hessian	Monte Carlo	Hessian
Free-proton PDFs	CT18	own fit	CT18A	~NNPDF4.0	~CTEQ6M
Free-proton corr.	no	no	yes	yes	no
HQ treatment	FONLL	FONLL	S-ACOT	FONLL	S-ACOT
Indep. flavours	3	4	6	6	5
Reference	PRD 104, 034010	arXiv:2112.11904	EPJC 82, 413	arXiv:2201.12363	arXiv:2204.09982

*see also PRD 103 (2021) 114015 & arXiv:2204.13157

Example parametrization: EPPS21

- Define nuclear PDFs in terms of

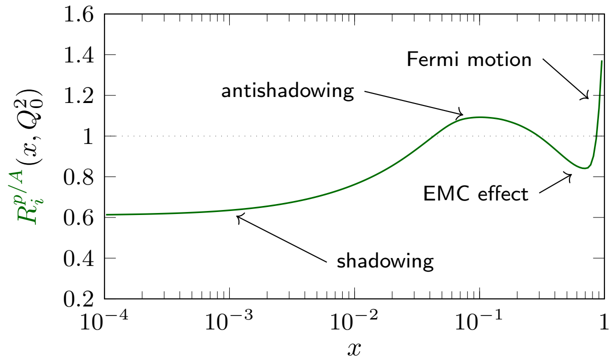
$$\underbrace{f_i^{p/A}}_{\text{bound-proton PDF}}(x,Q^2) = \overbrace{R_i^{p/A}}^{\text{nuclear modification}}(x,Q^2)\ \underbrace{f_i^{p}}_{\text{free-proton PDF}}(x,Q^2)$$

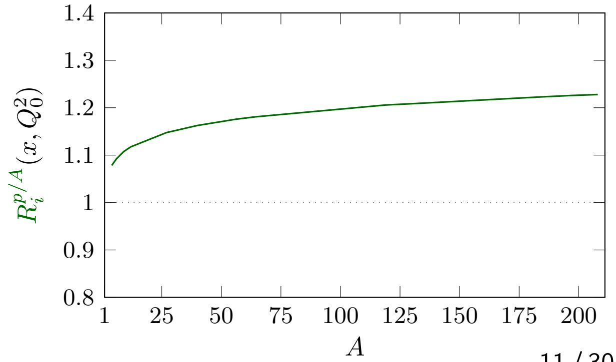
- PDFs of the full nucleus are then constructed with

$$f_i^{A}(x,Q^2) = Z f_i^{p/A}(x,Q^2) + N f_i^{n/A}(x,Q^2),$$

and assuming $f_i^{p/A} \overset{\text{isospin}}{\longleftrightarrow} f_j^{n/A}$

- Parametrize the x and A dependence of $R_i^{p/A}(x,Q_0^2)$ at $Q_0 = m_{\text{charm}} = 1.3$ GeV
  - Use a phenomenologically motivated piecewise function in x
  - Use a power-law type function in A





data from: Seely et al., PRL 103 (2009) 202301
CLAS Collaboration, Nature 566 (2019) 354

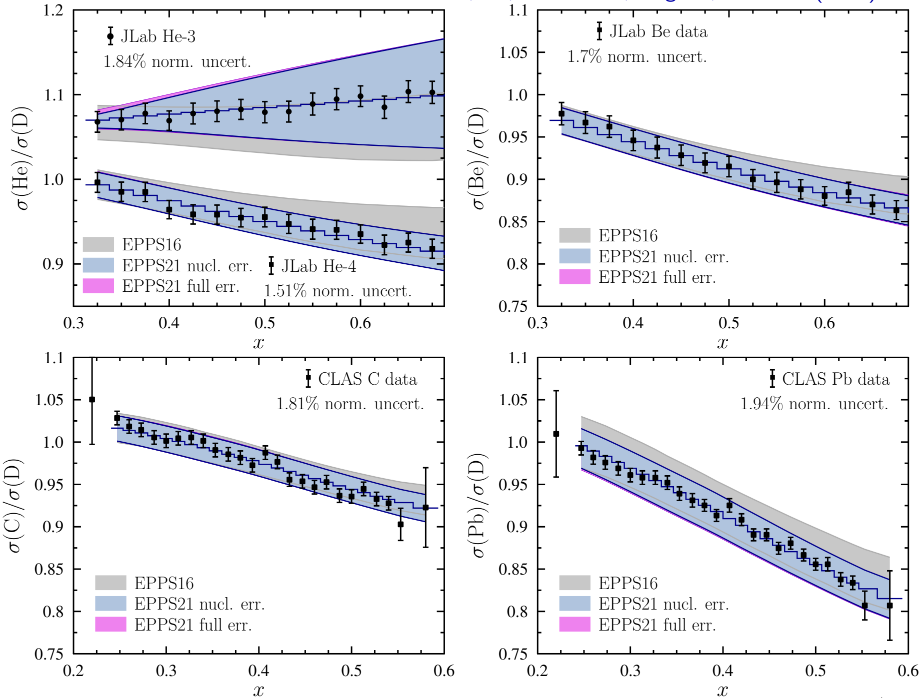
Eskola, PP, Paukkunen, Salgado, EPJC 82 (2022) 413

DIS in the “transition region” $W \gtrsim 1.7$ GeV just above the resonance-dominated one

Target-mass corrections important!

N.B. A-dependence not necessarily smooth for light nuclei ➔ need to scale the nuclear modifications for He-3 and Li-6 by factors

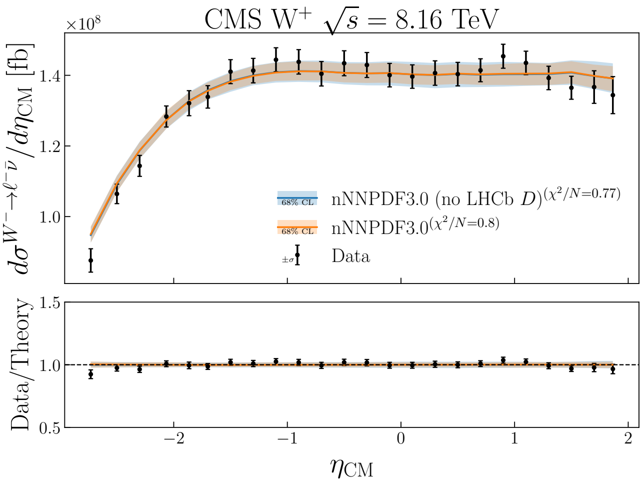
$$f_3 = 0.291, \quad f_6 = 0.495$$



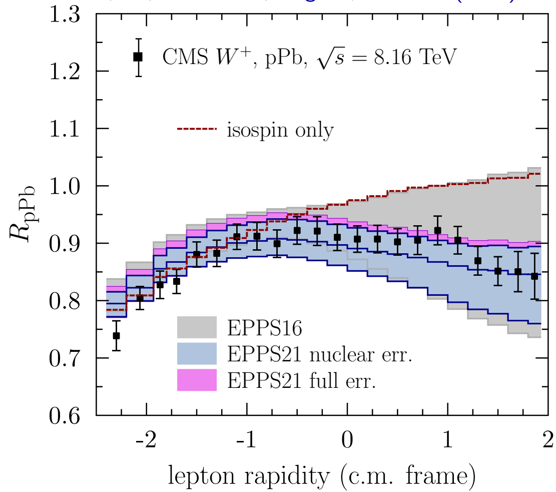
W bosons in pPb at 8.16 TeV

data from: CMS Collaboration, PLB 800 (2020) 135048
pp baseline: CMS Collaboration, EPJC 76 (2016) 469

Abdul Khalek et al., arXiv:2201.12363



Eskola, PP, Paukkunen, Salgado, EPJC 82 (2022) 413



EW bosons important probes of flavour separation

- $u\bar{d}$ ($c\bar{s}$) $\to W^+$
- $\bar{u}d$ ($\bar{c}s$) $\to W^-$

Small-x, high-Q^2 quarks and gluons correlated by DGLAP evolution ➔ sensitivity to gluons

nCTEQ15WZSIH, TUJU21 and nNNPDF3.0 fit to absolute cross sections

EPPS21 uses nuclear-modification ratios to cancel proton-PDF uncertainties
Eskola, PP, Paukkunen, Salgado, EPJC 82 (2022) 271

Mitigating free-proton PDF uncertainty

data from: CMS Collaboration, PLB 800 (2020) 135048
pp baseline: CMS Collaboration, EPJC 76 (2016) 469

Eskola, PP, Paukkunen, Salgado, EPJC 82 (2022) 271

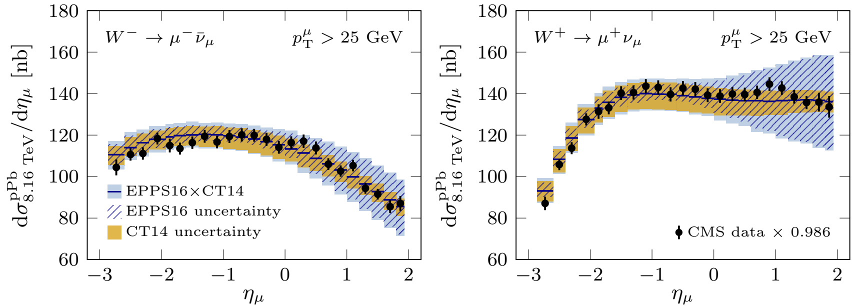
Absolute pPb cross sections sensitive to proton-PDF uncertainties!

Difficult to disentangle nuclear modifications from free-proton d.o.f.s

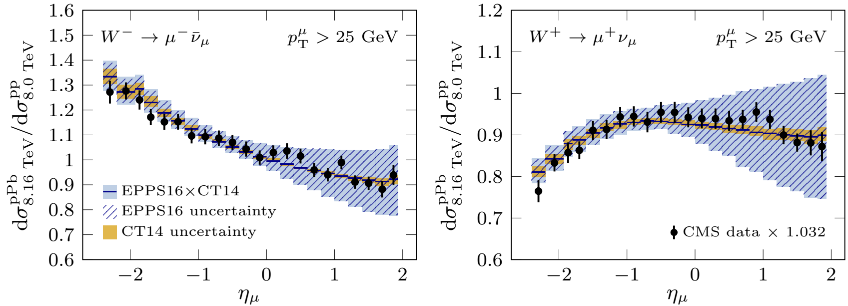
nCTEQ15WZSIH, TUJU21 and nNNPDF3.0 fit to absolute cross sections

Wherever possible, EPPS21 uses nuclear modification ratios to cancel the free-proton-PDF uncertainties

- ➔ can still become relevant with LHC Run 3 statistics
  Eskola, PP, Paukkunen, Salgado, EPJC 82 (2022) 271

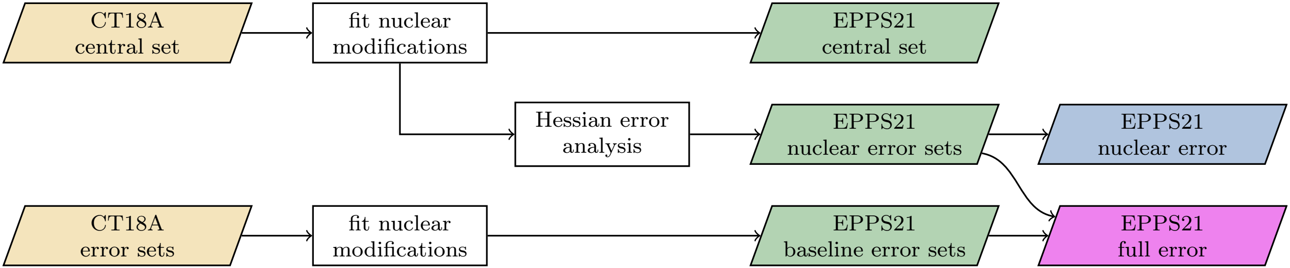


↓ Cancel proton-PDF uncertainty ↓



Propagating free-proton PDF uncertainty

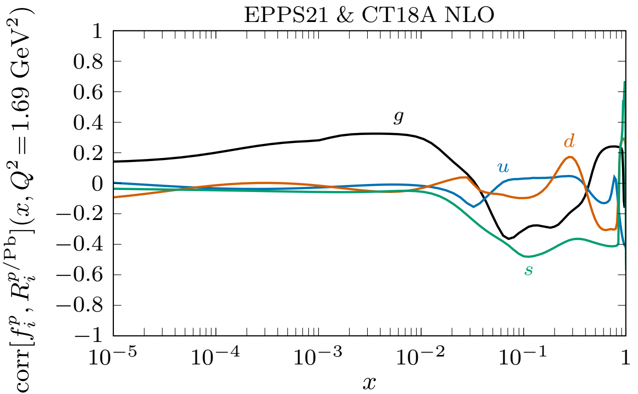
Eskola, PP, Paukkunen, Salgado, EPJC 82 (2022) 413



EPPS21: fit nuclear modifications for each CT18A error set separately

nNNPDF3.0 uses similar approach in Monte Carlo framework

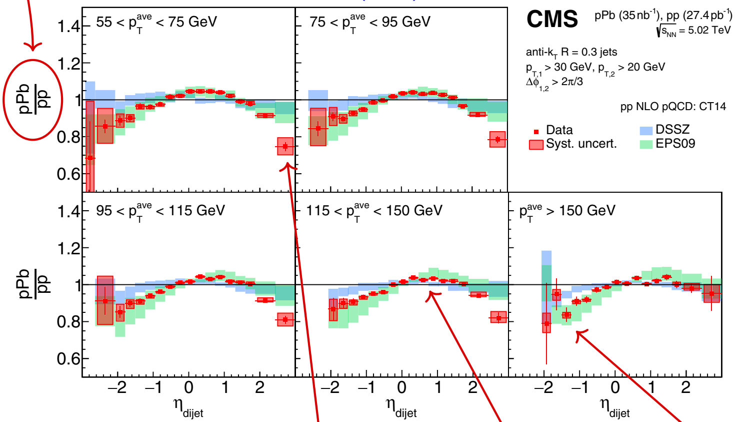
Note: nuclear modifications and proton PDFs become correlated!



Dijets in pPb at 5.02 TeV

Ratio of ratios: $R^{\rm norm.}_{\rm pPb} = \frac{{\rm d}^2\sigma^{\rm pPb}/{\rm d}p_{\rm T}^{\rm ave}{\rm d}\eta_{\rm dijet}}{{\rm d}\sigma^{\rm pPb}/{\rm d}p_{\rm T}^{\rm ave}} \Big/ \frac{{\rm d}^2\sigma^{\rm pp}/{\rm d}p_{\rm T}^{\rm ave}{\rm d}\eta_{\rm dijet}}{{\rm d}\sigma^{\rm pp}/{\rm d}p_{\rm T}^{\rm ave}}$

CMS Collaboration, PRL 121 (2018) 062002

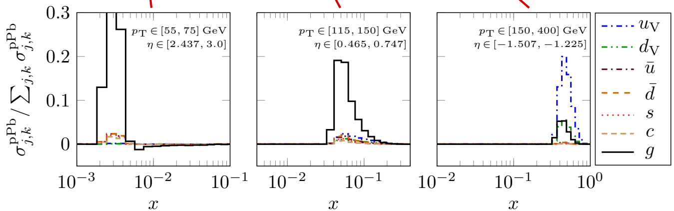


Double ratio convenient for:

- Cancellation of hadronization and luminosity uncertainties separately for pPb and pp
  - → do not expect strong final-state effects
- Cancellation of free-proton-PDF and scale uncertainties in pPb/pp
  - → direct access to nuclear modifications

Eskola, PP, Paukkunen, EPJC 79 (2019) 511

NLO pQCD:

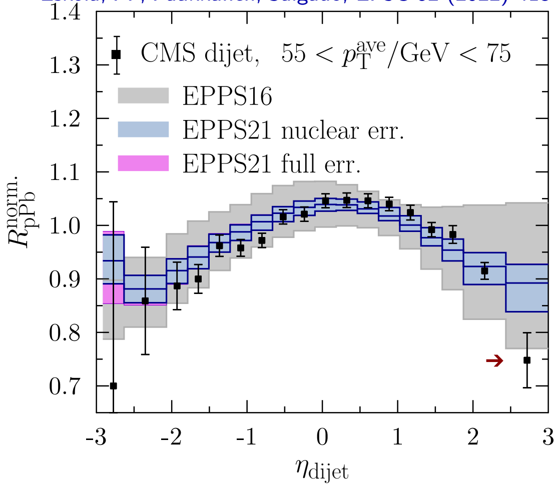


Good resolution to gluon nuclear modifications for $10^{-3} < x < 0.5$

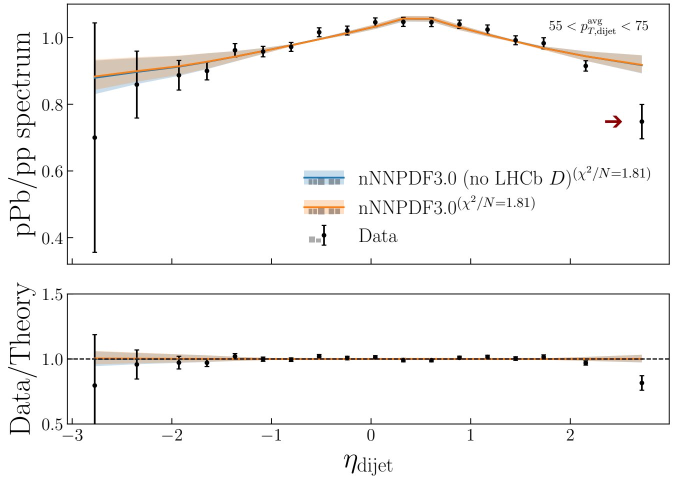
Dijets in EPPS21 and nNNPDF3.0

data from: CMS Collaboration, PRL 121 (2018) 062002

Eskola, PP, Paukkunen, Salgado, EPJC 82 (2022) 413



Abdul Khalek et al., arXiv:2201.12363



Drastic reduction in the nPDF uncertainties!

➔ Important constraints for the nuclear gluons!

Eskola, PP, Paukkunen, EPJC 79 (2019) 511
Eskola, PP, Paukkunen, Salgado, EPJC 82 (2022) 413
Abdul Khalek et al., arXiv:2201.12363

Both EPPS21 and nNNPDF3.0 find difficulties in reproducing the most forward data points

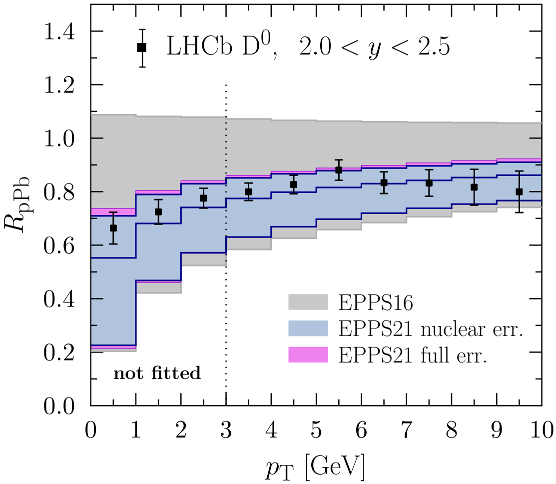
➔ missing data correlations important?

➔ NNLO? non-pert. effects?

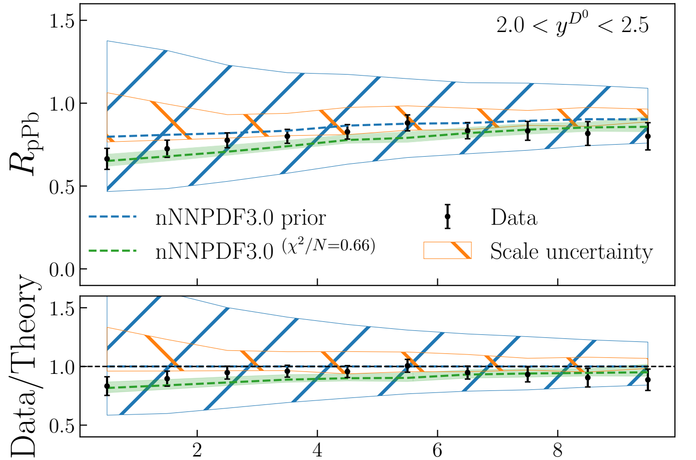
D^0s in EPPS21 and nNNPDF3.0

data from: LHCb Collaboration, JHEP 10 (2017) 090

Eskola, PP, Paukkunen, Salgado, EPJC 82 (2022) 413



Abdul Khalek et al., arXiv:2201.12363



Drastic reduction in the nPDF uncertainties!

➔ Important constraints for the nuclear gluons!

Kusina et al., PRL 121 (2018) 052004
Eskola, Helenius, PP, Paukkunen, JHEP 05 (2020) 037
Eskola, PP, Paukkunen, Salgado, EPJC 82 (2022) 413
Abdul Khalek et al., arXiv:2201.12363

nNNPDF3.0 with POWHEG+PYTHIA finds a large scale uncertainty ➔ fit only forward data

not seen in the S-ACOT-m_T GM-VFNS used in EPPS21

Helenius & Paukkunen, JHEP 05 (2018) 196
Eskola, Helenius, PP, Paukkunen, JHEP 05 (2020) 037

FFNS

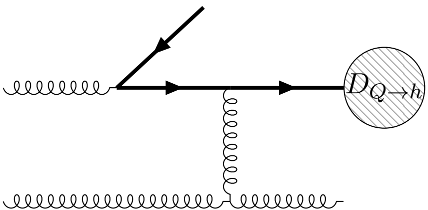
In *fixed flavour number scheme*, valid at small p_T, heavy quarks are produced only at the matrix element level

Contains $\log(p_\mathrm{T}/m)$ and $\mathcal{O}(m)$ terms

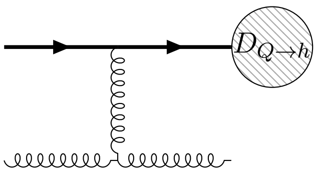
ZM-VFNS

In *zero-mass variable flavour number scheme*, valid at large p_T, heavy quarks are treated as massless particles produced also in ISR/FSR

Resums $\log(p_\mathrm{T}/m)$ but ignores $\mathcal{O}(m)$ terms



$-$ subtraction term $+$



GM-VFNS

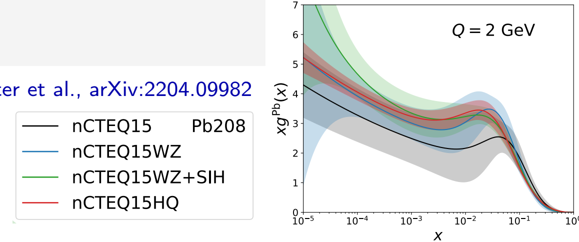
A *general-mass variable flavour number scheme* combines the two by supplementing subtraction terms to prevent double counting of the resummed splittings, valid at all p_T

Resums $\log(p_\mathrm{T}/m)$ and includes $\mathcal{O}(m)$ terms in the FFNS matrix elements

Important: includes also gluon-to-HF fragmentation – large contribution to the cross section!

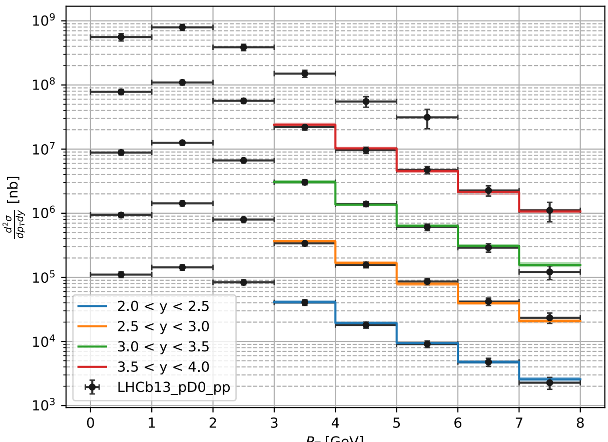
A data-driven approach – nCTEQ15HQ

Duwentäster et al., arXiv:2204.09982

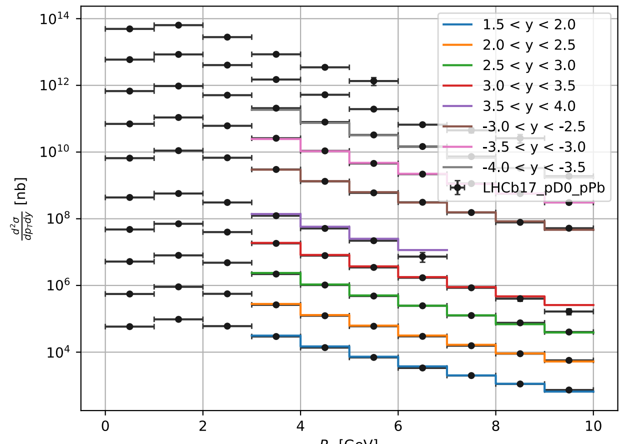


nCTEQ15HQ uses a data-driven approach
Lansberg & Shao, EPJC 77 (2017) 1
Kusina et al., PRL 121 (2018) 052004
to fit the D^0 and J/ψ data:

1. Fit the matrix elements to pp data...
   (assume $2 \to 2$ kinematics, neglect IS quarks)



2. ...use the fitted matrix elements to fit nuclear PDFs with pPb data



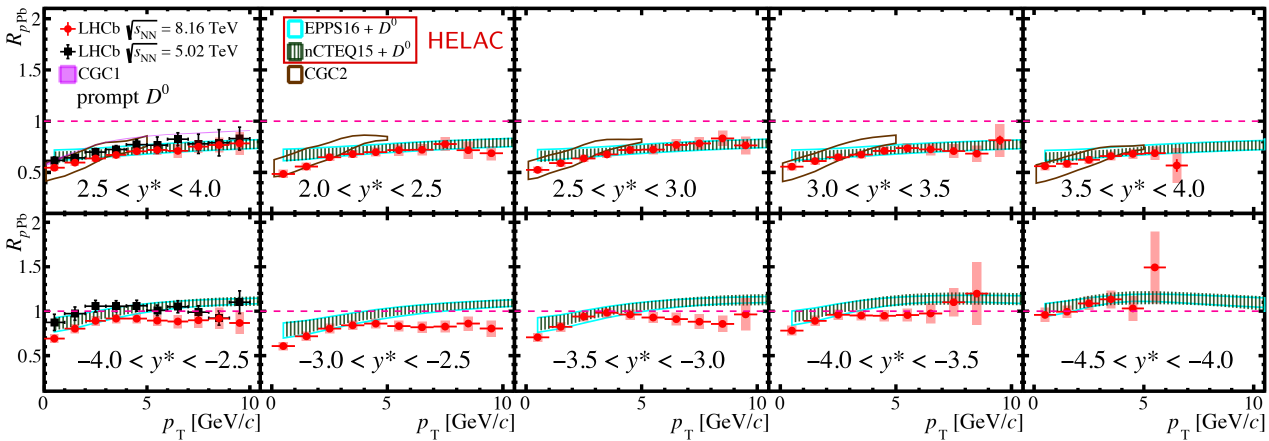
New LHCb measurement at 8.16 TeV
(not included in the nPDF analyses yet)

pp reference interpolated from 5 and 13 TeV measurements

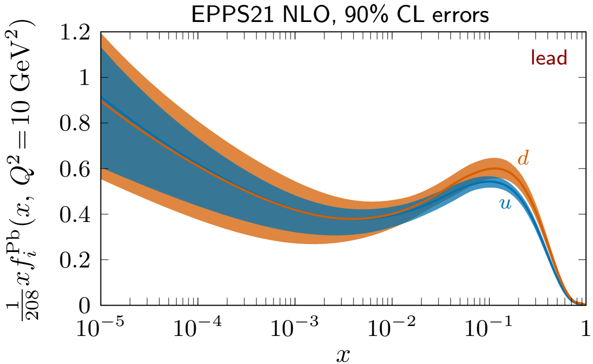
So far compared only against the HELAC matrix-element-fitting results
Kusina et al., PRL 121 (2018) 052004

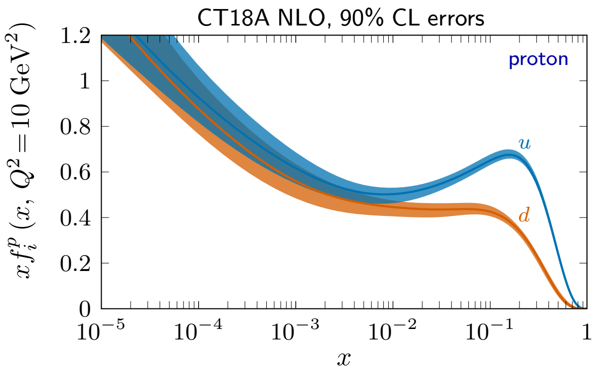
➔ to be scrutinised with the direct pQCD calculations

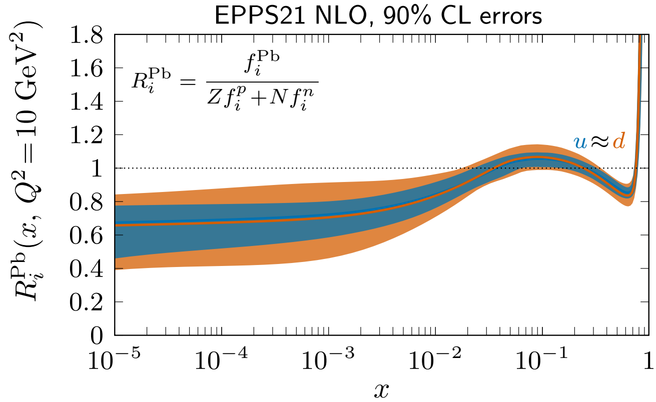
LHCb Collaboration, arXiv:2205.03936



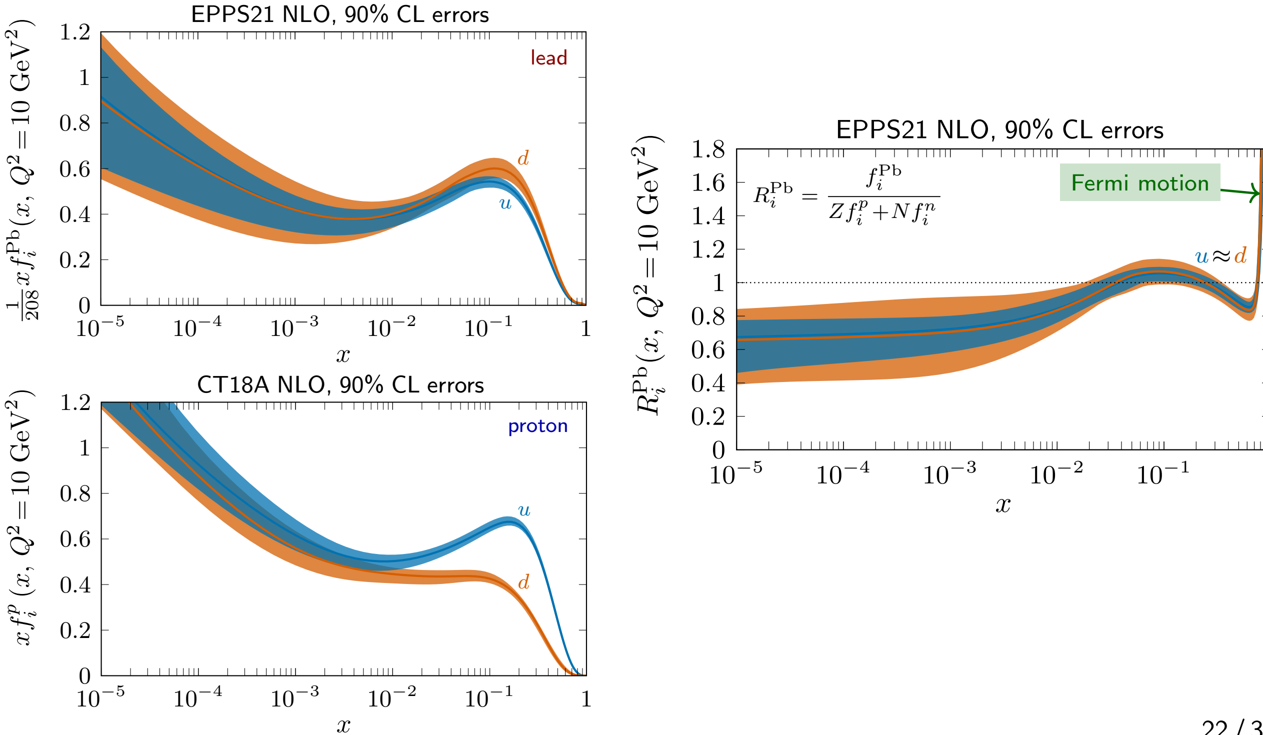
Comparing nuclear and proton PDFs – u and d



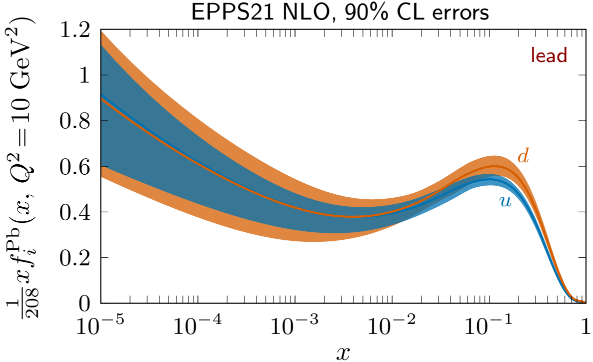


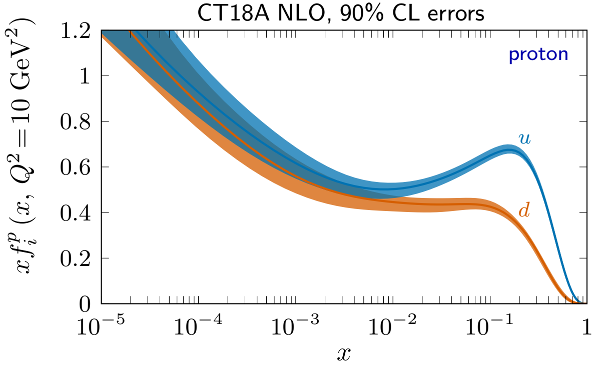


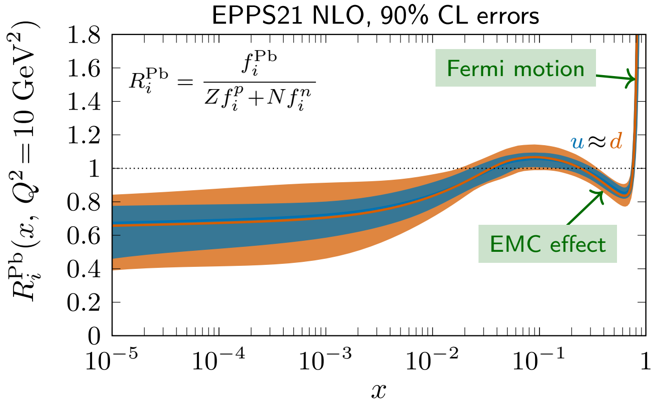
Comparing nuclear and proton PDFs – u and d



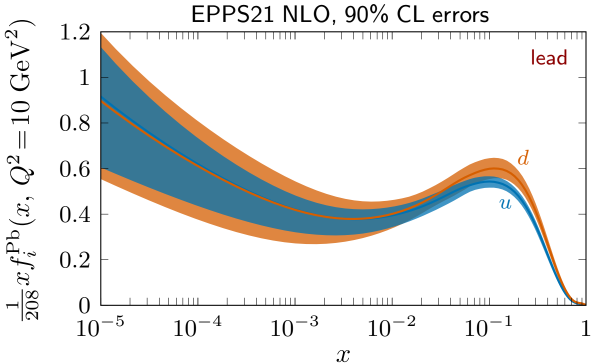
Comparing nuclear and proton PDFs – u and d

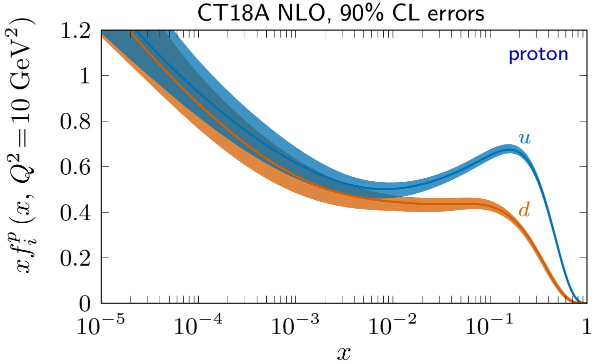


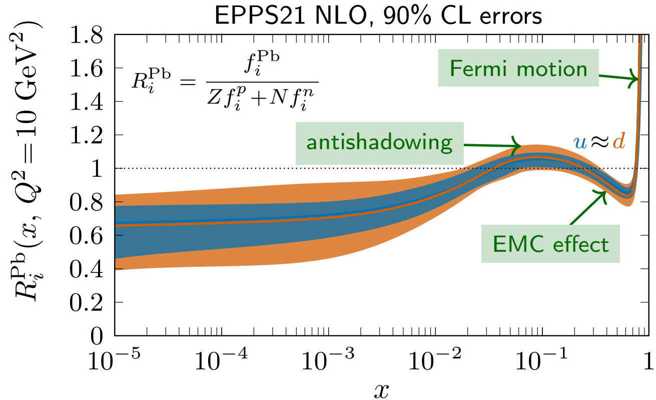




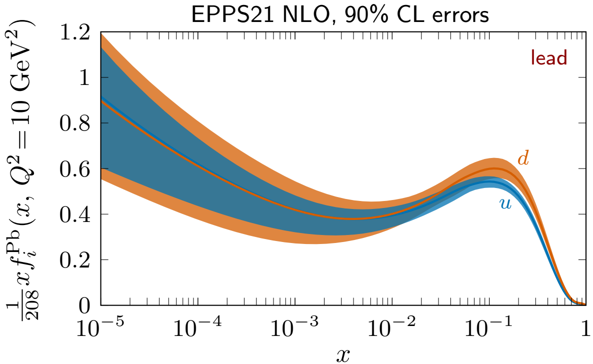
Comparing nuclear and proton PDFs – u and d

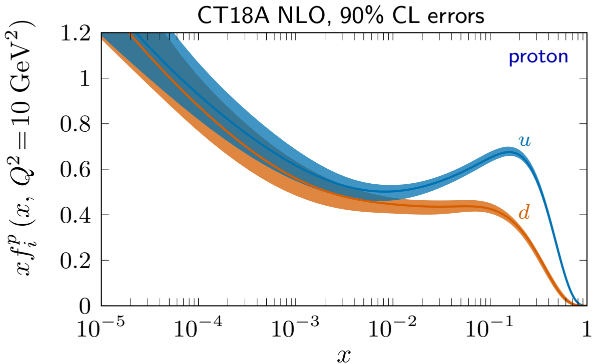


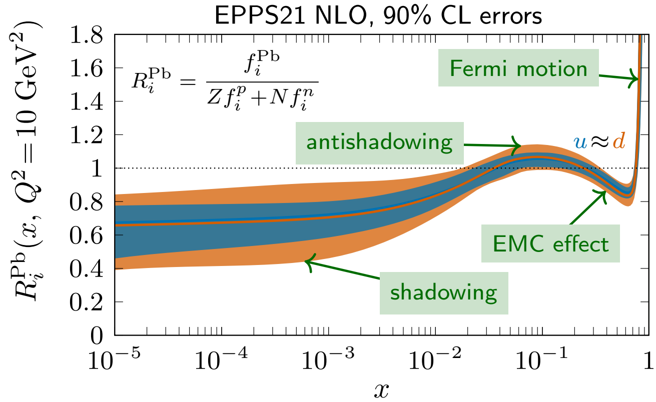




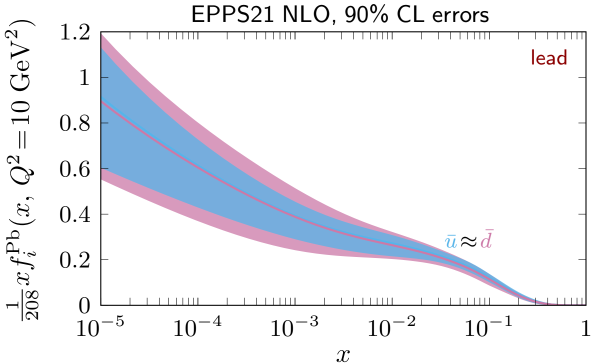
Comparing nuclear and proton PDFs – u and d

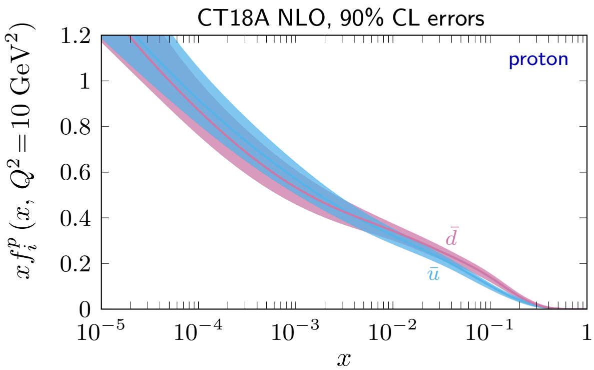


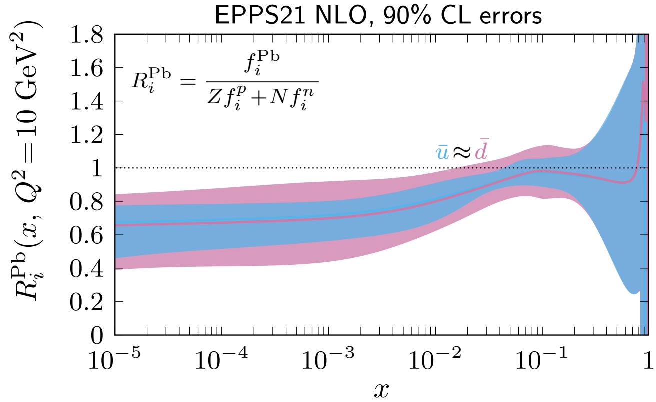




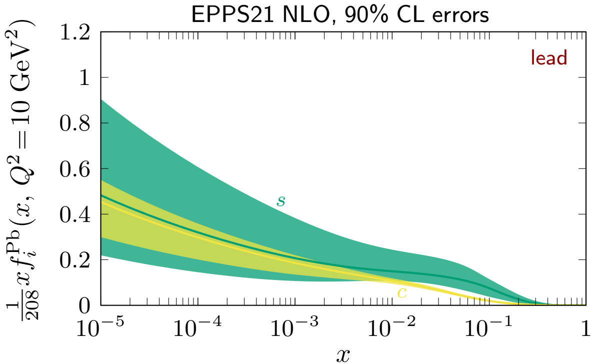
Comparing nuclear and proton PDFs – $\bar{u}$ and $\bar{d}$

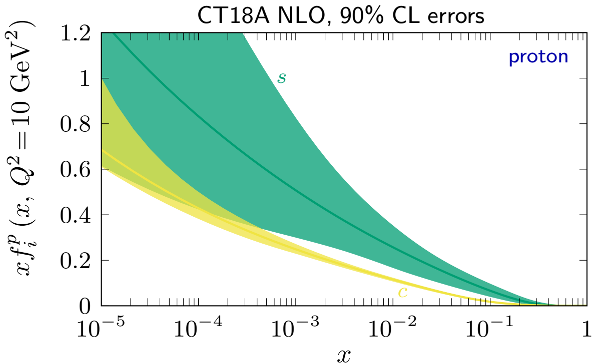


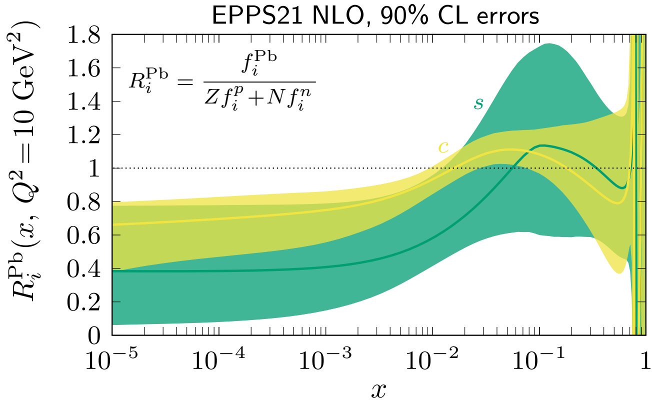




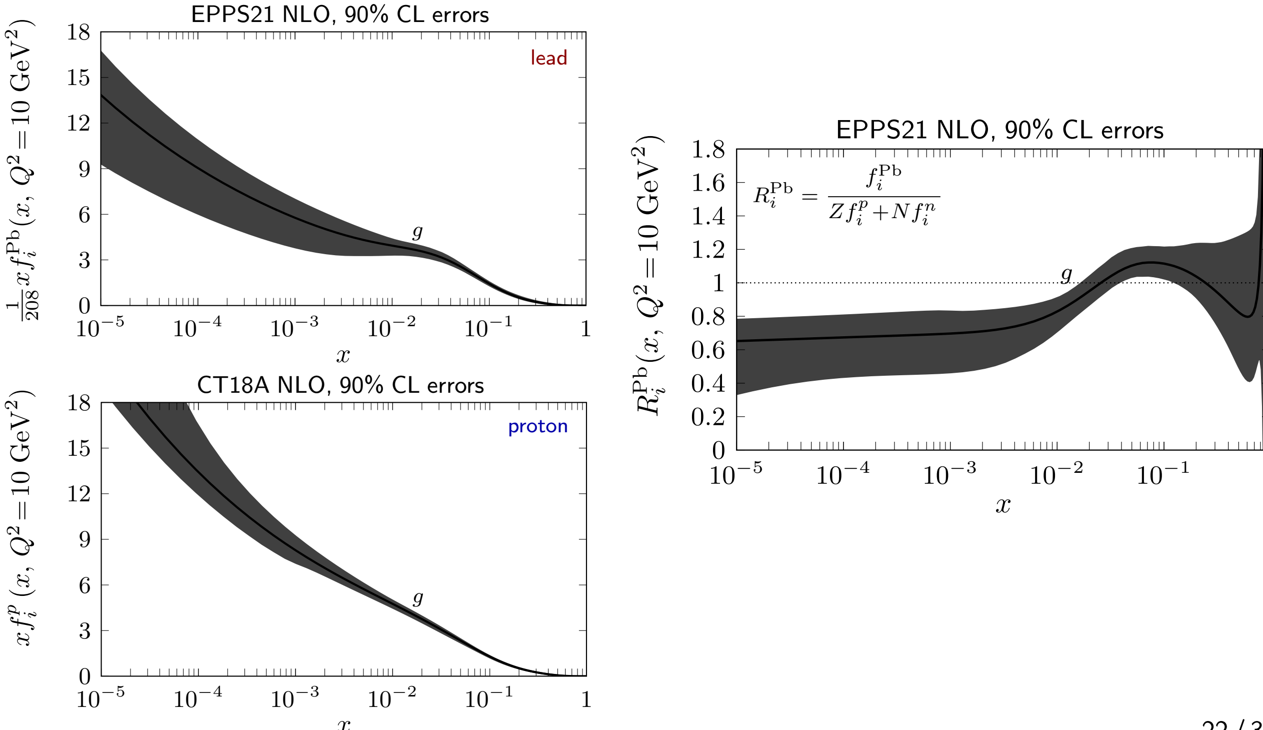
Comparing nuclear and proton PDFs – s and c



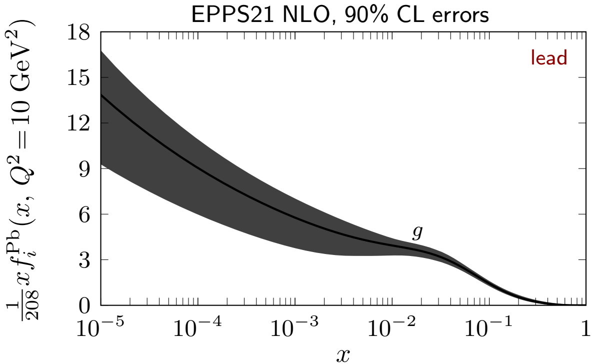


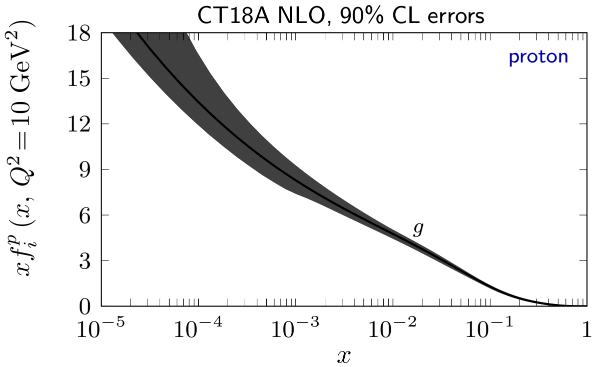


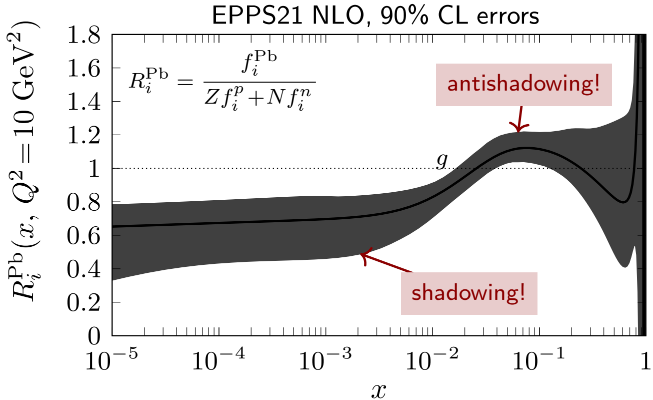
Comparing nuclear and proton PDFs – *glue*



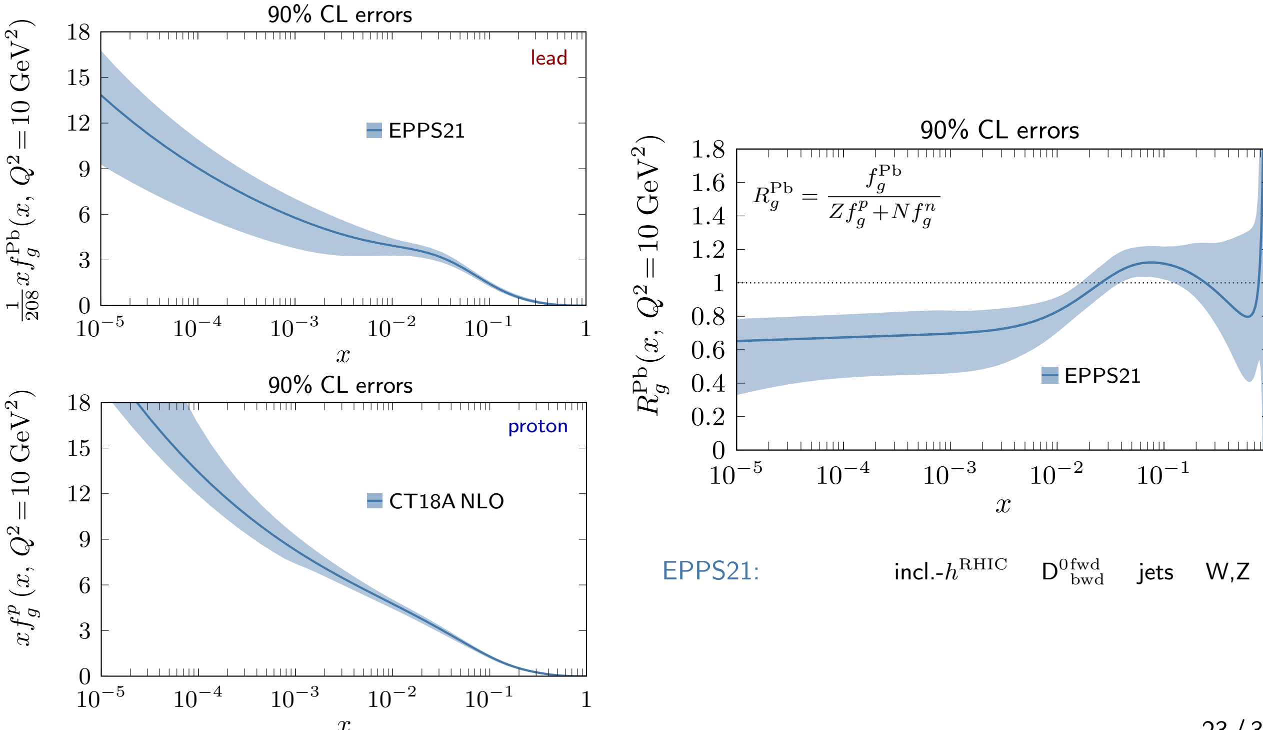
Comparing nuclear and proton PDFs – *glue*





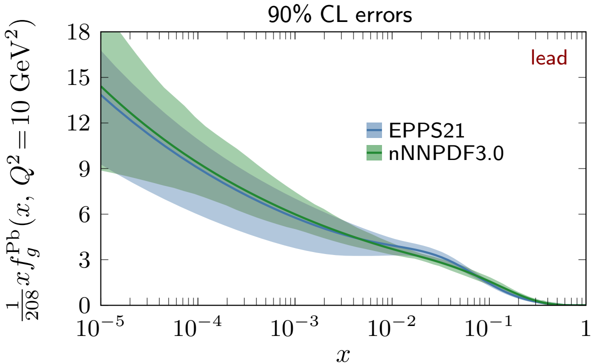


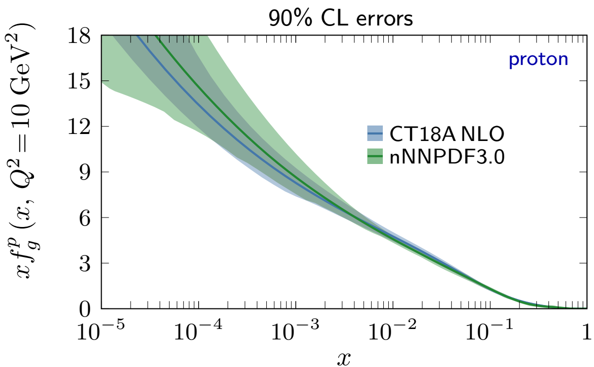
nPDF comparison – *glue*

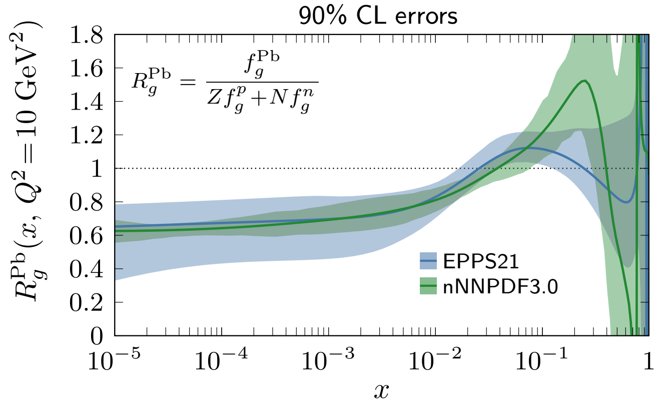


EPPS21: incl.-$h^{\rm RHIC}$ $D^{0\,{\rm fwd}}_{\rm bwd}$ jets W,Z

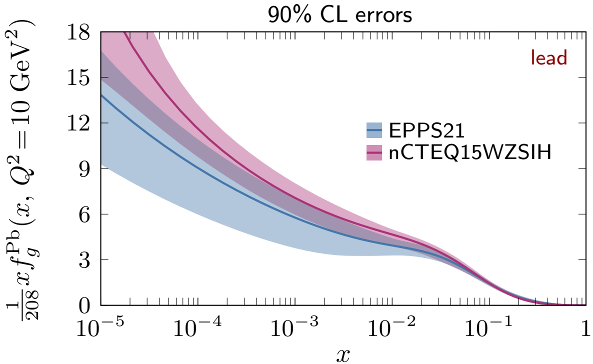
nPDF comparison – *glue*

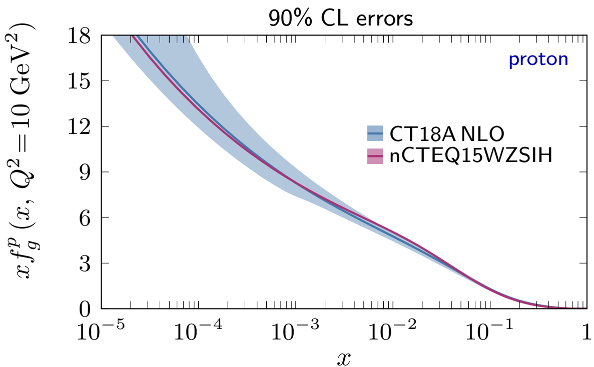


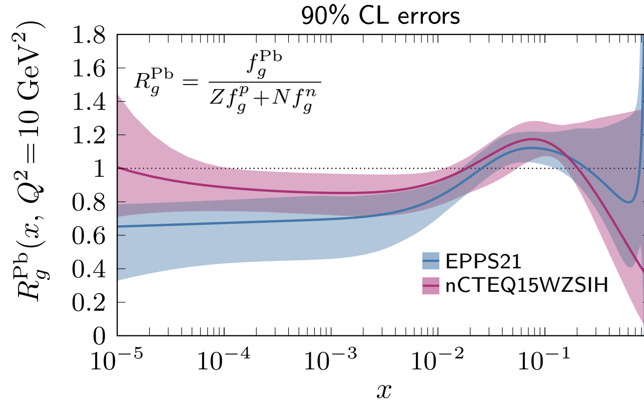




EPPS21:	incl.-h^{RHIC}	$\mathrm{D}^{0\,\mathrm{fwd}}_{\mathrm{bwd}}$	jets	W,Z
nNNPDF3.0:		$\mathrm{D}^{0\,\mathrm{fwd}}$	jets	W,Z

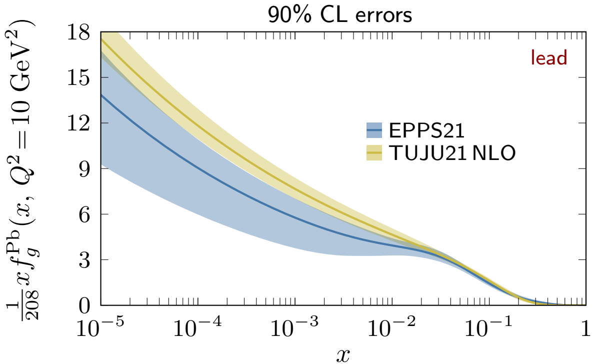


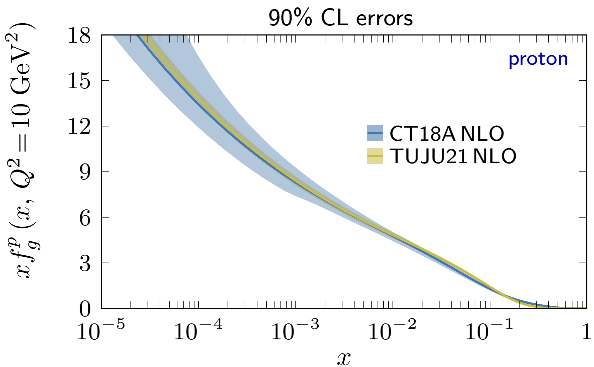


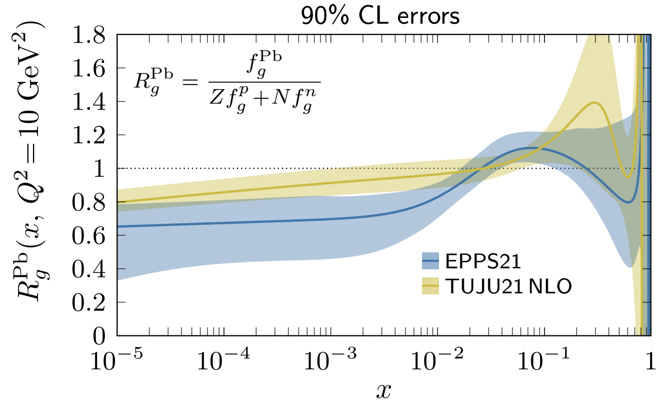


EPPS21:	incl.-h^{RHIC}	$\mathrm{D}^{0\,\mathrm{fwd}}_{\mathrm{bwd}}$	jets	W,Z
nNNPDF3.0:		$\mathrm{D}^{0\,\mathrm{fwd}}$	jets	W,Z
nCTEQ15WZSIH:	incl.-$h^{\mathrm{RHIC}}_{\mathrm{LHC}}$			W,Z

nPDF comparison – *glue*

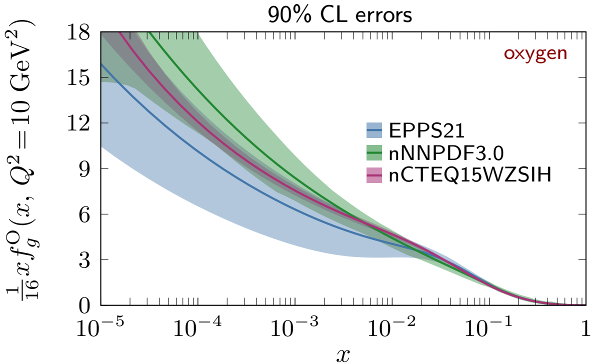


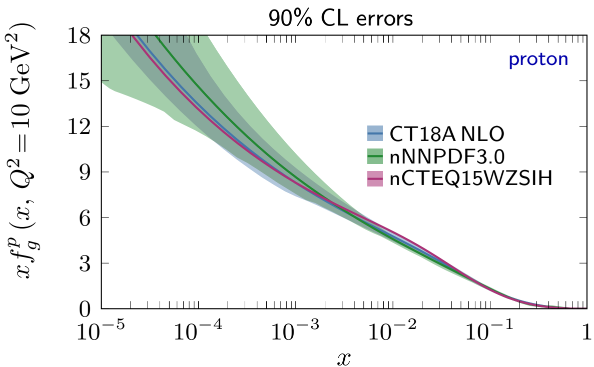




EPPS21:	incl.-$h^{\rm RHIC}$	$D^{0\,\rm fwd}_{\rm bwd}$	jets	W,Z
nNNPDF3.0:		$D^{0\,\rm fwd}$	jets	W,Z
nCTEQ15WZSIH:	incl.-$h^{\rm RHIC}_{\rm LHC}$			W,Z
TUJU21:				W,Z

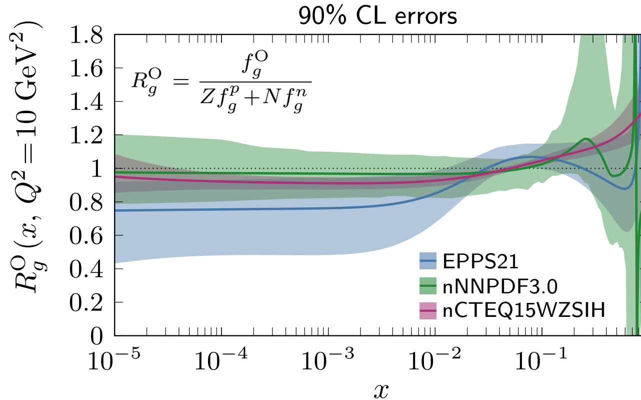
nPDF comparison – *glue* in oxygen





nPDFs a major source of uncertainty in testing small-system energy loss with OO

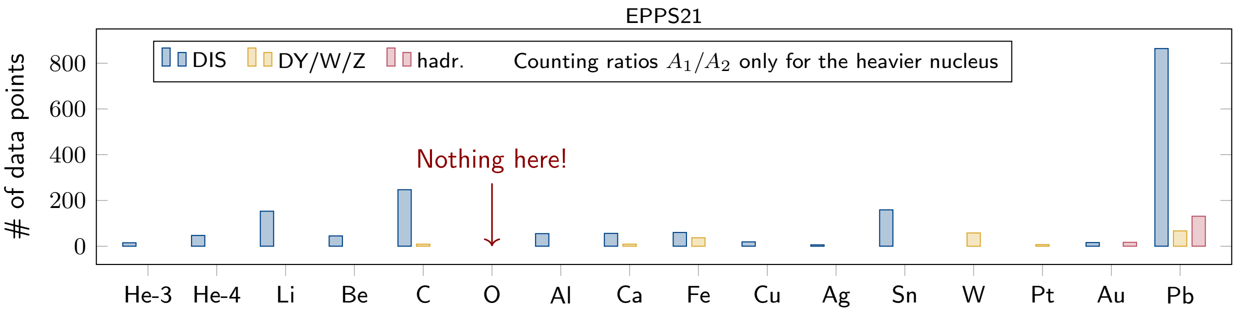
Huss et al., PRL 126 (2021) 192301
Brewer et al., PRD 105 (2022) 074040



EPPS21:	incl.-h^{RHIC}	$\mathrm{D}^{0\,\mathrm{fwd}}_{\mathrm{bwd}}$	jets	W,Z
nNNPDF3.0:		$\mathrm{D}^{0\,\mathrm{fwd}}$	jets	W,Z
nCTEQ15WZSIH:	incl.-$h^{\mathrm{RHIC}}_{\mathrm{LHC}}$			W,Z

only dAu and pPb!

Data availability w.r.t. A



$\sim 50\%$ of the data points are for Pb!

- ☺ Good coverage of DIS measurements for different A (but only fixed target!)
- 😐 DY data more scarce, but OK A coverage
- ☹ Hadronic observables available only for heavy nuclei!

Light-ion runs at LHC could:

- Complement other light-nuclei DY data with W and Z production (strangeness!)
- Give first direct constraints (e.g. dijets, D-mesons) on light-nuclei small-x gluon distributions!

PP, PRD 105 (2022) L031504

Similar setup as in CMS 5.02 TeV pPb measurement

Total integrated pO cross section of $81\ \mu\text{b}$

- Compare with $\sim 330\ \mu\text{b}$ in pPb at 5.02 TeV
- Sufficient to give reasonable statistics even at relatively low luminosities

  16000 events at $0.2\ \text{nb}^{-1}$

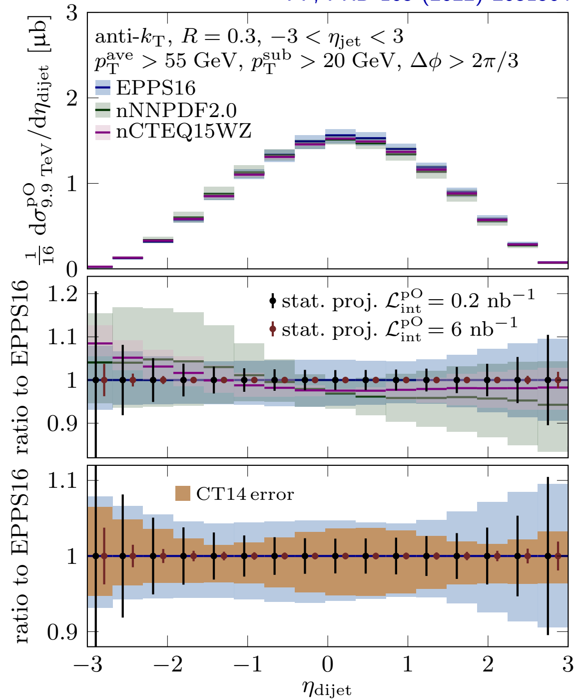
  486000 events at $6\ \text{nb}^{-1}$

Problem: absolute cross sections very sensitive to the used free-proton PDFs

- Difficult to disentangle nuclear modifications from the free-proton d.o.f.s

Problem: We do not expect pp reference at 9.9 TeV

- Could we use a mixed energy ratio pO(9.9 TeV)/pp(8.8 TeV)?



Dijet R_{pO} in pO at 9.9 TeV

PP, PRD 105 (2022) L031504

Problem: We do not expect pp reference at 9.9 TeV

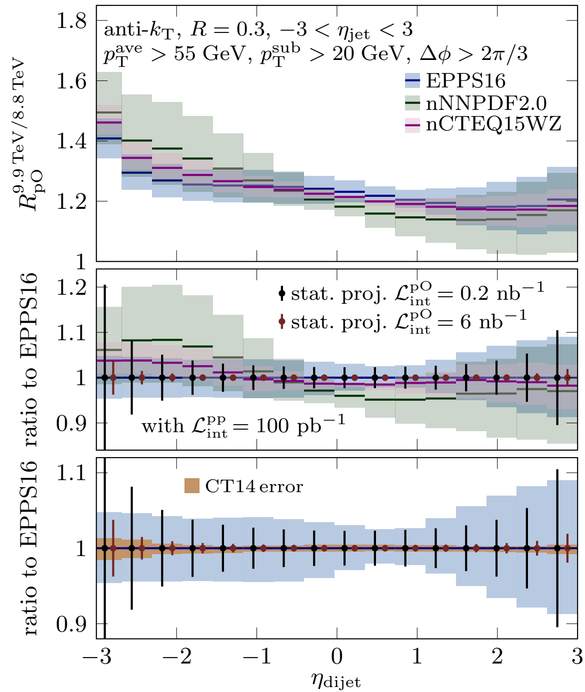
- Could we use a mixed energy ratio pO(9.9 TeV)/pp(8.8 TeV)? Yes!

Excellent cancellation of free-proton PDFs

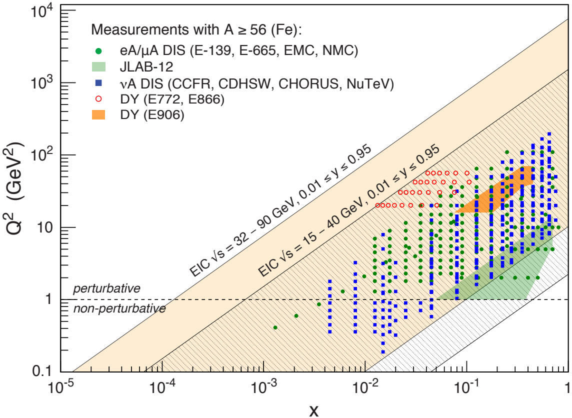
→ Direct access to nuclear modifications

Already few nb^{-1} can be expected to be enough to put new constraints on nPDFs (if we have sufficient statistics for the pp reference)

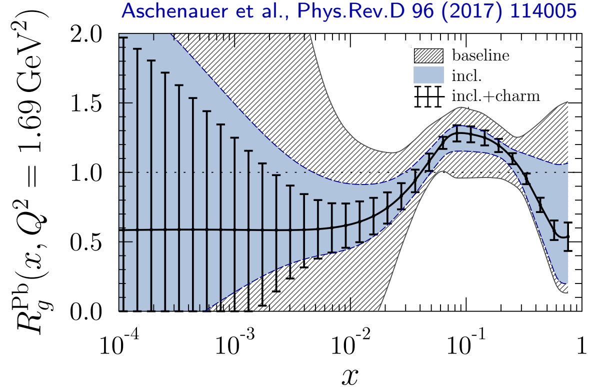
→ Can resolve different nPDF parametrisations!



Gluon constraints from EIC



Aschenauer et al., Phys.Rev.D 96 (2017) 114005



EIC will significantly widen the kinematic range of DIS constraints for nPDFs

- Comparing with LHC measurements will put collinear factorization with nuclei to a stringent test

With the F_{L} extraction cabability, EIC provides a clean probe to study small-x gluons

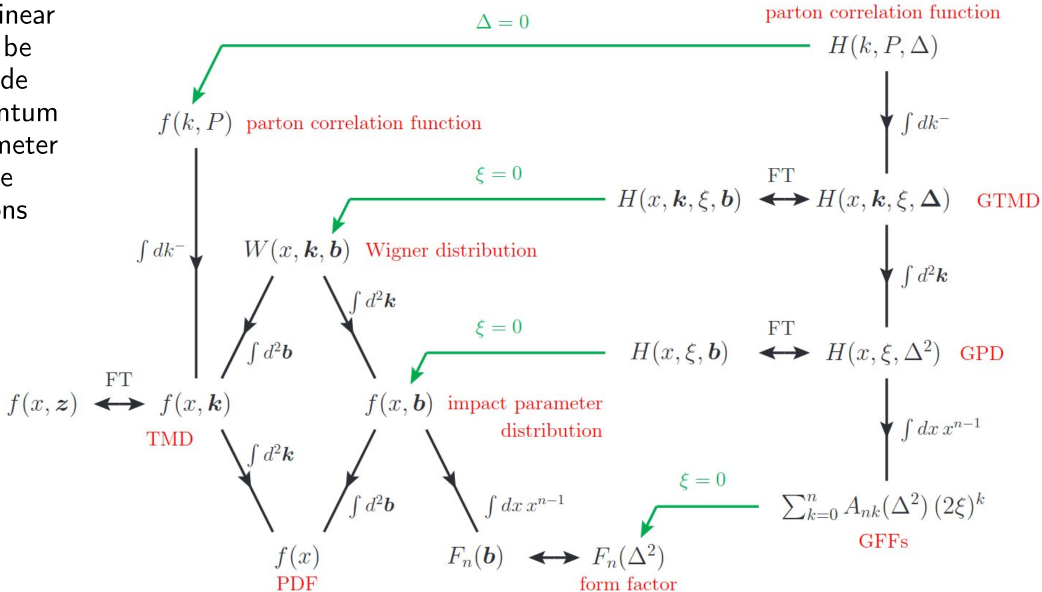
- Good constraining power to well down to 10^{-2} in a high-energy scenario

Charm-tagged cross-section measurement can vastly reduce high-x gluon uncertainty

see also: Kelsey et al., Phys.Rev.D 104 (2021) 054002

The "simple" collinear factorization can be extended to include transverse-momentum and impact-parameter dependence of the parton distributions

Diehl, Eur. Phys. J. A 52 (2016) 149



...each of these have their own region of applicability in terms of experimental observables, outside of which the factorization is not quaranted to work!

Impact-parameter dependent nPDFs EPS09s and EKS98s

Helenius et al., JHEP 07 (2012) 073

Generalize the nPDFs by defining

nuclear modification — spatially dependent nuclear modification

$$R_i^{p/A}(x,Q^2) = \frac{1}{A}\int d^2\mathbf{s}\ T_A(\mathbf{s})\ r_i^A(x,Q^2,\mathbf{s})$$

nuclear thickness function

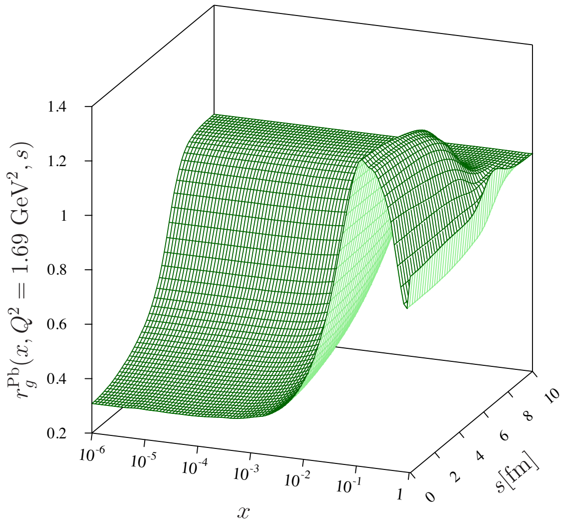
such that

$$dN^{AB\to k+X}(\mathbf{b}) = \sum_{i,j,X'}\sum_{N_A,N_B}\int d^2\mathbf{s}_1 T_A(\mathbf{s}_1) r_i^A(x,Q^2,\mathbf{s}_1) f_i^{N_A}(x,Q^2)\otimes \int d^2\mathbf{s}_2 T_B(\mathbf{s}_2) r_i^B(x,Q^2,\mathbf{s}_2) f_i^{N_B}(x,Q^2)\otimes d\hat{\sigma}^{ij\to k+X'}\delta(\mathbf{s}_2-\mathbf{s}_1-\mathbf{b})$$

reduces to the min. bias cross section when integrated over the impact parameter b

<u>To my knowledge</u> not proven to coincide with the GPD definition of impact-parameter distribution!

See Wu, JHEP 07 (2021) 002 for factorization derivation with similar objects in soft-collinear effective theory



Thank you!

Z bosons in pPb at 8.16 TeV

data from: CMS Collaboration, JHEP 05 (2021) 182
pp baseline: CMS Collaboration, EPJC 75 (2015) 147

New Run 2 data from CMS

CMS Collaboration, JHEP 05 (2021) 182

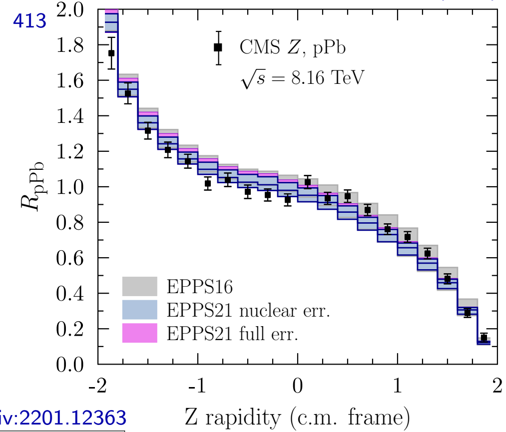
- nNNPDF3.0 include both low-mass and on-peak data
- R_{pPb} studied in EPPS21 → not included in the final fit

Both EPPS21 and nNNPDF3.0 observe some tension between the data and fit

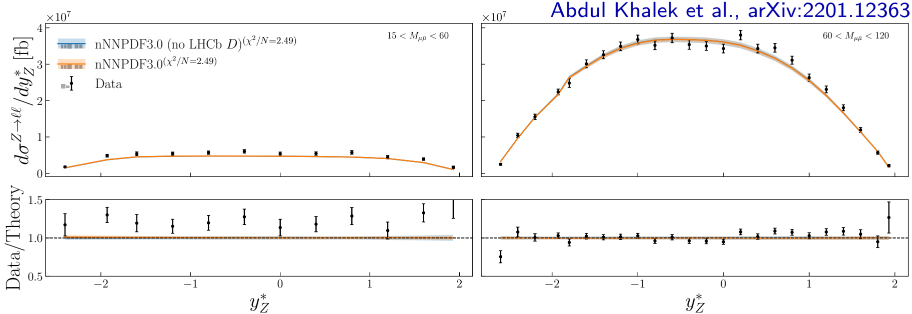
- abrupt change in the shape at midrapidity
- NNLO to cure for the low-mass data?

Abdul Khalek et al., arXiv:2201.12363

Eskola, PP, Paukkunen, Salgado, EPJC 82 (2022) 413



Abdul Khalek et al., arXiv:2201.12363



	χ^2/N_{data}
EPPS21	2.1
nNNPDF3.0	2.49

PHENIX pion production small-system scan

New mid-rapidity π^0 data from PHENIX
PHENIX Collaboration, arXiv:2111.05756

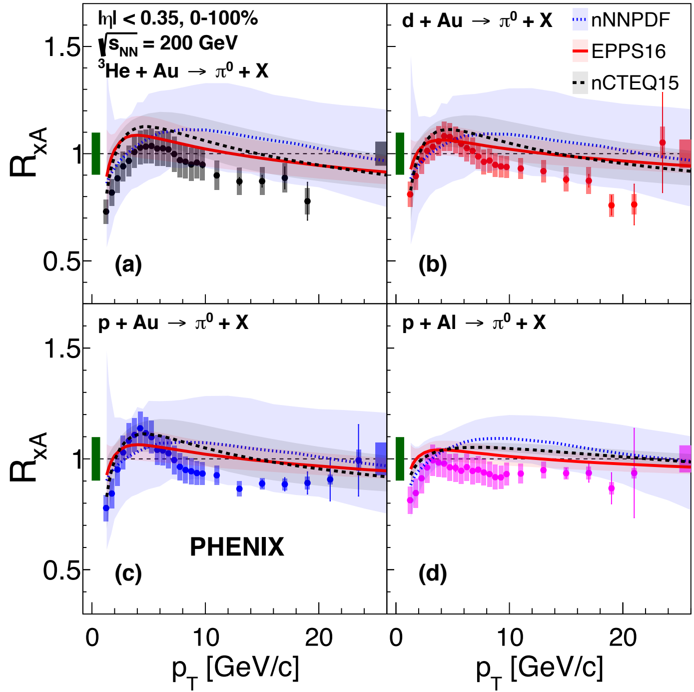
- improved precision
- higher p_T → larger x

Contrary to nPDF expectations, measured "Cronin peak" size follows the ordering ${}^3\mathrm{He}+\mathrm{Au} < d+\mathrm{Au} < p+\mathrm{Au}$

- higher-twist (multiple-scattering)?
- flow-like component?

At high p_T the nPDF predictions overshoot the data, but mind the large normalisation uncertainties

PHENIX Collaboration, arXiv:2111.05756



UPCs in collinear factorisation

First phenomenological implementation of the exclusive J/ψ photoproduction NLO corrections

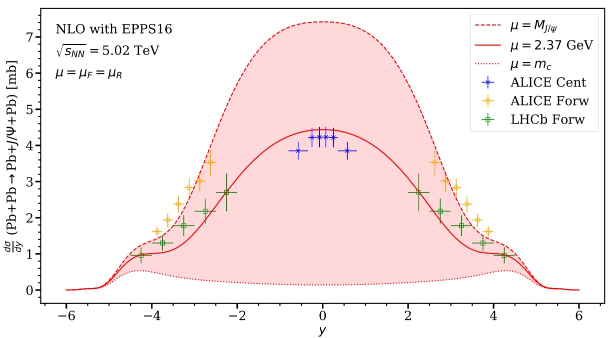
Ivanov et al., EPJC 34 (2004) 297
Jones et al., J. Phys. G 43 (2016) 035002

in ultrapheripheral Pb+Pb

Large scale uncertainty

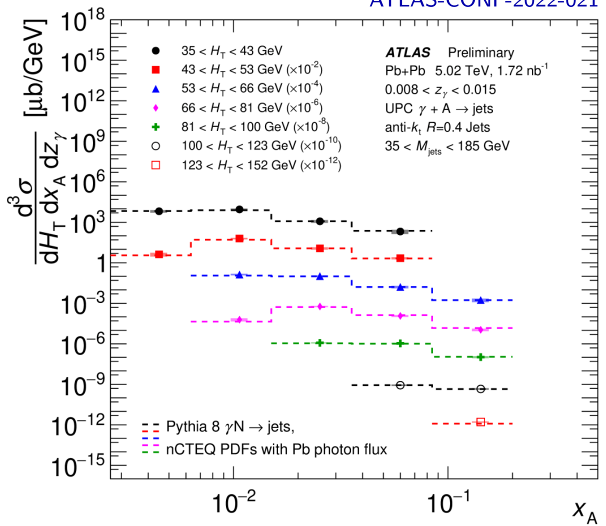
- ➔ perturbative convergence?
- ➔ cancel with nuclear ratios?

Eskola et al., arXiv:2203.11613

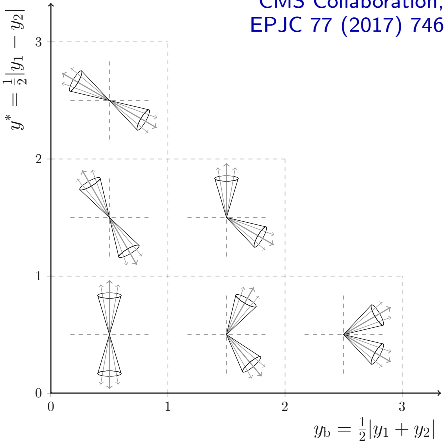


ATLAS inclusive dijet photoproduction measurement now fully unfolded

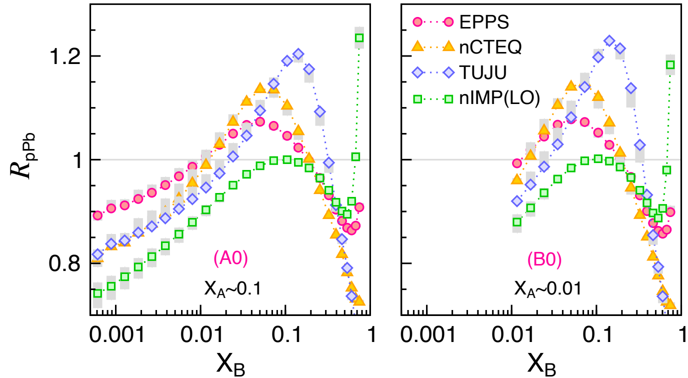
ATLAS-CONF-2022-021



Triple-differential dijets in pPb?



CMS Collaboration, EPJC 77 (2017) 746



Shen et al., arXiv:2112.11819

Triple-differential measurement fixes partonic kinematics at LO

➔ powerful test of factorisation and PDFs

Measured in pp at 8 TeV

CMS Collaboration, EPJC 77 (2017) 746

Should be feasible in pPb with Run 2/3 statistics?

Various observable choices possible, e.g. X_A, X_B, y^*

measurable!

$$X_B = \sum_{n \in \text{dijet}} \frac{E_{\text{T}n}}{\sqrt{s}} e^{-y_n} \overset{\text{LO}}{=} x_{\text{Pb}}$$

momentum fraction from the lead side

Ellis & Soper, PRL 74 (1995) 5182

Proton strangeness from νA DIS vs. LHC EW data

$$K_s = \frac{\int_0^1 dx\,x[s(x,Q^2)+\bar{s}(x,Q^2)]}{\int_0^1 dx\,x[\bar{u}(x,Q^2)+\bar{d}(x,Q^2)]}$$

table and fig. from Feng et al., "The Forward Physics Facility at the High-Luminosity LHC", arXiv:2203.05090

			Proton PDF sets				Nuclear PDF sets			
Data set		Ref.	ABMP16	CT18	MSHT20	NNPDF4.0	EPPS21	nCTEQ15	nNNPDF3.0	TUJU21
CHORUS $\sigma_{CC}^{\nu,\bar{\nu}}$	Pb	[1238]	✗	✗	✓	✓	✓	✗	✓	✓
CHORUS	Pb	[1239]	✓	✗	✗	✗	✗	✗	✗	✗
NOMAD $\mathcal{R}_{\mu\mu}$	Fe	[1195]	✓	✗	✗	(✓)	✗	✗	✗	✗
CCFR xF_3^p	Fe	[1240]	✗	✓	✗	✗	✗	✗	✗	✗
CCFR F_2^p	Fe	[1241]	✗	✓	✗	✗	✗	✗	✗	✗
CDSHW F_2^p, xF_3^p	Fe	[1242]	✗	✓	✗	✗	✗	✗	✗	✓
NuTeV $\sigma_{CC}^{\nu,\bar{\nu}}$	Fe	[1196]	✓	✓	✓	✓	✗	✗	✓	✗
NuTeV F_2, F_3	Fe	[1194]	✗	✗	✓	✗	✗	✗	✗	✗

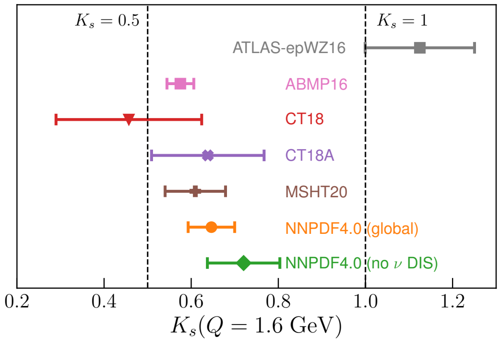
Proton-PDF fits traditionally include neutrino-*nucleus* DIS for improved strange-quark constraints ➔ suppressed strangeness

Complementary data from ATLAS EW-boson production confronts this view with preference for unsuppressed strange

ATLAS Collaboration, PRL 109 (2012) 012001
EPJC 77 (2017) 367

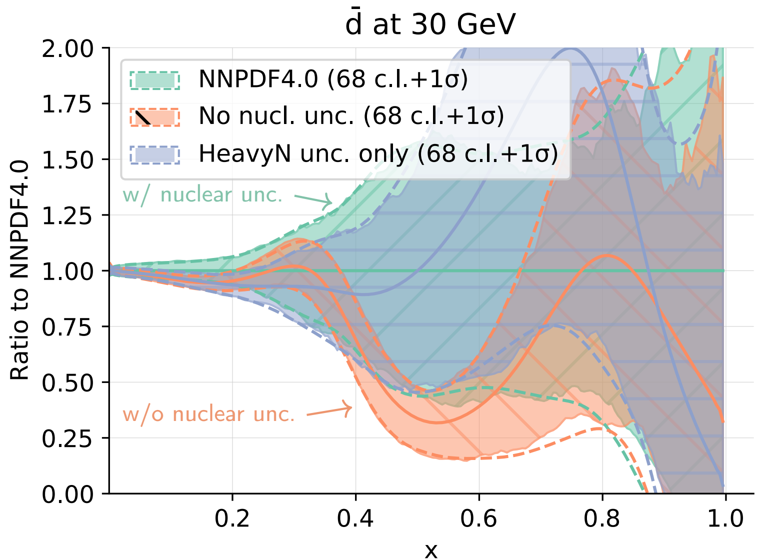
Simultaneous fit feasible w/ NNLO c-quark mass corrections

Faura et al., EPJC 80 (2020) 1168
Bailey et al., EPJC 81 (2021) 341
Ball et al., arXiv:2109.02653



Nuclear uncertainties in proton-PDF fits

Ball et al., arXiv:2109.02653



Nuclear effects can impact the proton-PDF fits!

NNPDF4.0:

- Different large-x sea-quark behaviour depending on whether the uncertainties from nNNPDF2.0 nuclear PDFs were included or not
- Nuclear data found to constrain the proton PDFs even with nuclear uncertainties included

Ball et al., EPJC 79 (2019) 282
Ball et al., arXiv:2109.02653

MSHT20: take nuclear corrections from DSSZ + additional 3-param. fit
Bailey et al., EPJC 81 (2021) 341

CT18: does not report on any use of nuclear corrections