



DE LA RECHERCHE À L'INDUSTRIE

# PES topology and interplay with shell effects

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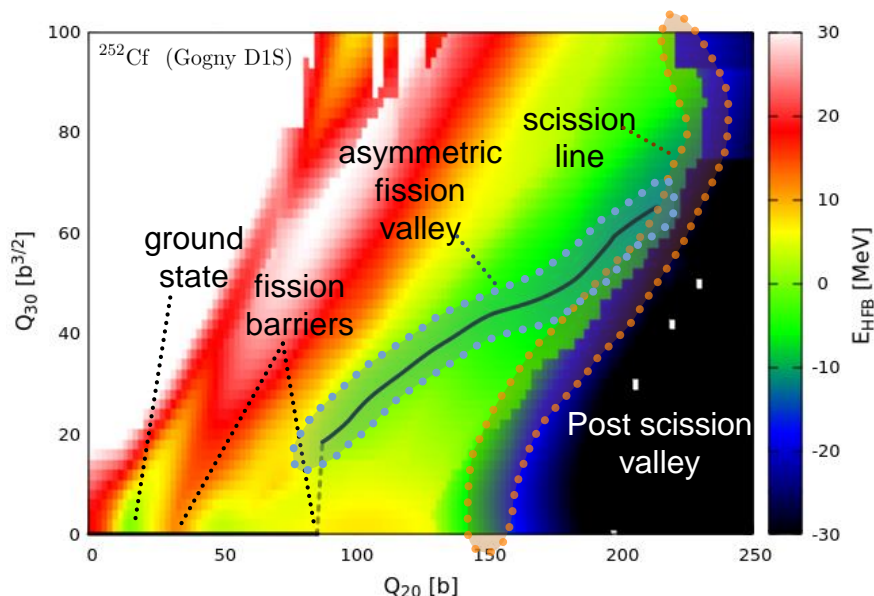
Rémi Bernard

CEA, Cadarache

***IRISNE | DER | SPRC | LEPH***

- Smoothing discontinuities from Potential Energy Surfaces
  - Motivations
  - The Dynamic Programming Method in 1 D
  - The Dynamic Programming Method in 2 D
  
- Shell effects in fission
  - Level densities at the Fermi surfaces
  - Connection to (pre)fragments
  - 180Hg case

Hartree-Fock-Bogoliubov (or any other mean field) under constraints calculations to generate the PES



Adiabatic hypothesis : the PES is built from mean field state minima for each deformation.

- It implies discontinuities in unconstrained variables.
- Missing the physics at the discontinuity

How to deal with discontinuities before the scission line ?

- Adding a dimension in the calculations ?
- Linear interpolation ?
- Other ?

## The Time-Dependent Generator Coordinate Method

$$|\Psi_{\text{GCM}}(t)\rangle = \int d\mathbf{q} f(\mathbf{q}, t) |\Phi(\mathbf{q})\rangle$$



Hill-Wheeler equation

$$\int d\mathbf{q}' \left( \underbrace{\mathcal{H}(\mathbf{q}, \mathbf{q}')}_{\text{Hamiltonian kernel}} - i\hbar \underbrace{\mathcal{N}(\mathbf{q}, \mathbf{q}') \frac{d}{dt}}_{\text{overlap kernel}} \right) f(\mathbf{q}', t) = 0$$

$$\mathcal{H}(\mathbf{q}, \mathbf{q}') = \langle \Phi(\mathbf{q}) | \hat{H} | \Phi(\mathbf{q}') \rangle$$

Hamiltonian kernel

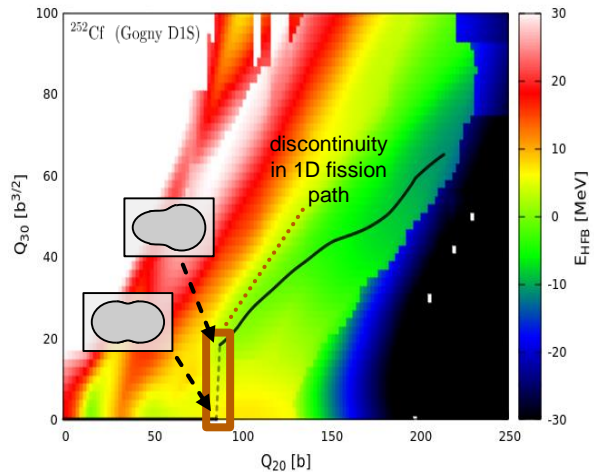
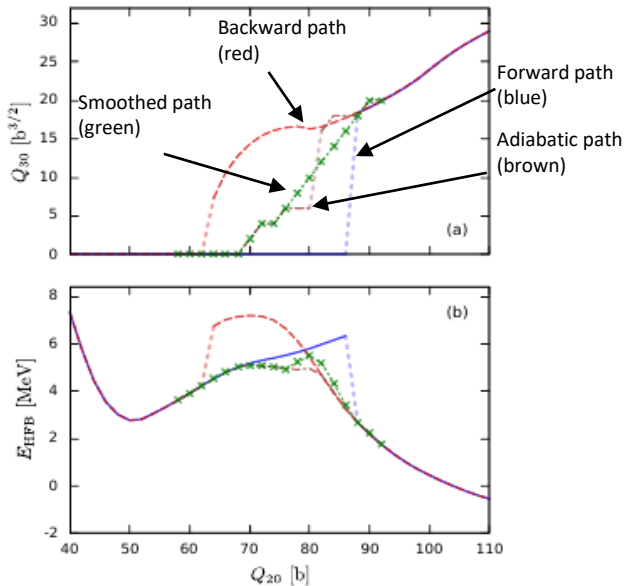
$$\mathcal{N}(\mathbf{q}, \mathbf{q}') = \langle \Phi(\mathbf{q}) | \Phi(\mathbf{q}') \rangle$$

overlap kernel

## Overlap and Hamiltonian kernels

- Inversion of the square root of the overlap kernel usually done to solve the TDGCM equation
- Any inertia formula is based on wrong mean field states locally

→ issue for the TDGCM wave function evolution



*N.-W. T. Lau, et al, Phys. Rev. C 105, 034617 (2022)*

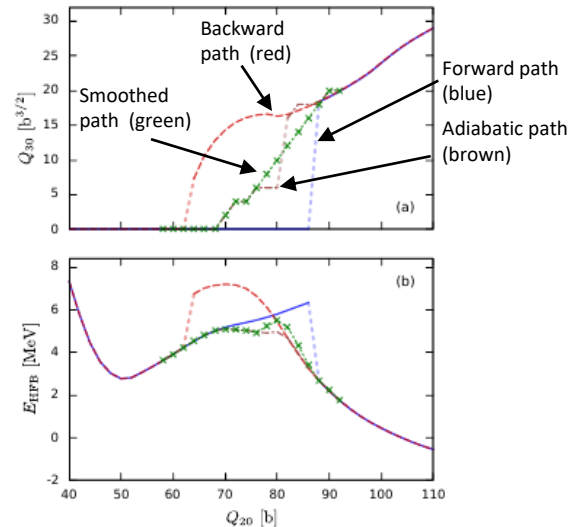
Useful tools to detect discontinuities :

- Overlaps (Onishi Formula) between adjacent states  $\mathcal{N}(\mathbf{q}, \mathbf{q}') = \langle \Phi(\mathbf{q}) | \Phi(\mathbf{q}') \rangle$
- Density distance

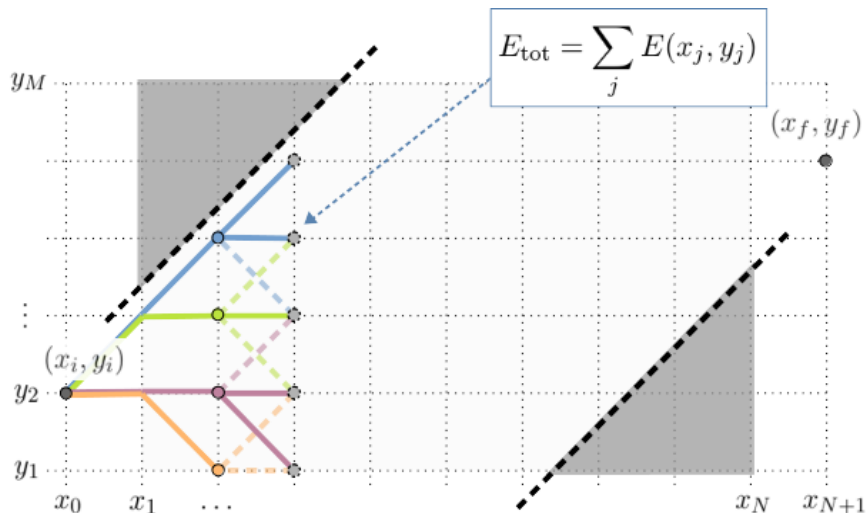
*N.-W. T. Lau, et al, Phys. Rev. C 105, 034617 (2022)*

The method to clean the discontinuities is built from the following criteria :

- Minimizing the total energy
- Remaining smooth in energy
- The 1D/2D PES must be continuous in terms of overlap or density distance
- Method must be efficient to facilitate its application to many discontinuities and PESs at a reasonable cost



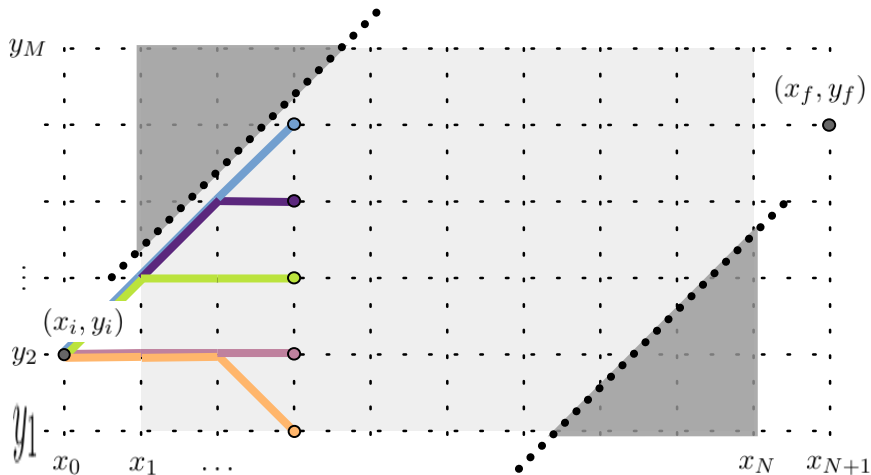
*The solution is not unique, one may start from other requirements or add some extra criteria...*



Conventional tree search:  
 $M^N$  paths

Dynamic Programming  
 Method (DPM):  
 $MN$  trajectories by keeping  
 the best path for each point

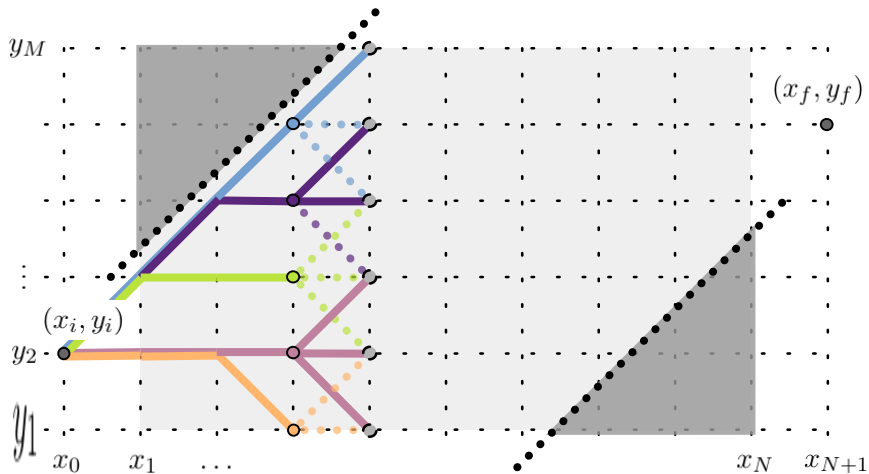
*A. Baran et al. 1981 Nuclear Physics A361 83-101*  
*J. Sadhukhan et al. 2013 Phys. Rev. C 88 064314*



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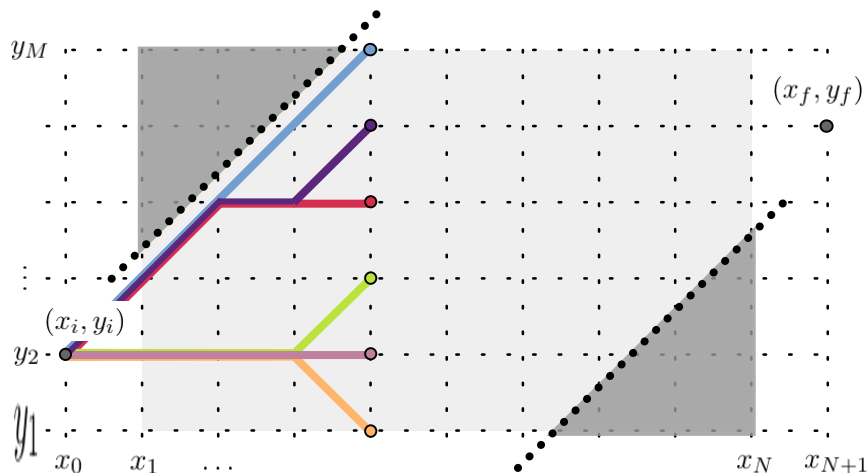
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**1D DPM** : for  $M = 9$  and  $N = 19$ ; an extremely short run (of few seconds on the CDC-CYBER 7316 computer) allows us to choose the minimal trajectory from among a huge number,  $919 \times 1018$ , of all trial paths. (Baran NPA 81)

*A. Baran et al. 1981 Nuclear Physics A361 83-101*  
*J. Sadhukhan et al. 2013 Phys. Rev. C 88 064314*

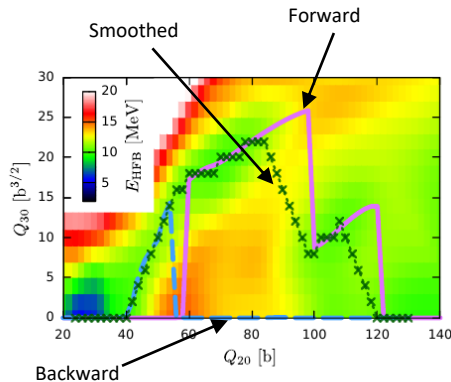
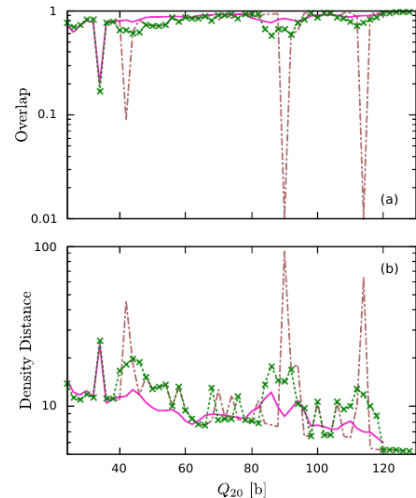
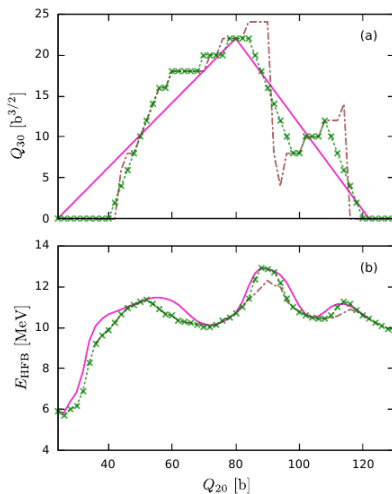


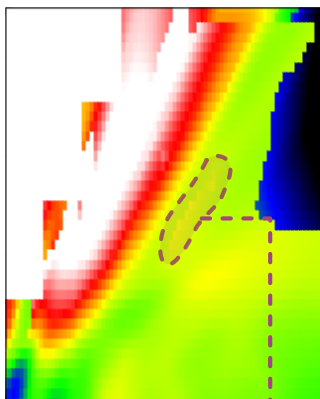
FIG. 8. A subsection of the 2D PES in  $Q_{20}$  and  $Q_{30}$  for  $^{222}\text{Th}$ , generated with  $N_{\perp} = 15$ ,  $N_z = 22$  shells and fixed oscillator lengths  $b_{\perp} = 2.0$ ,  $b_z = 3.0$ . Energies are shown relative to the ground state. The HFB paths propagated in the  $+Q_{20}$  and  $-Q_{20}$  directions are drawn with solid pink and dashed blue respectively, while the smoothed path generated using DPM is depicted in dotted green with crosses.



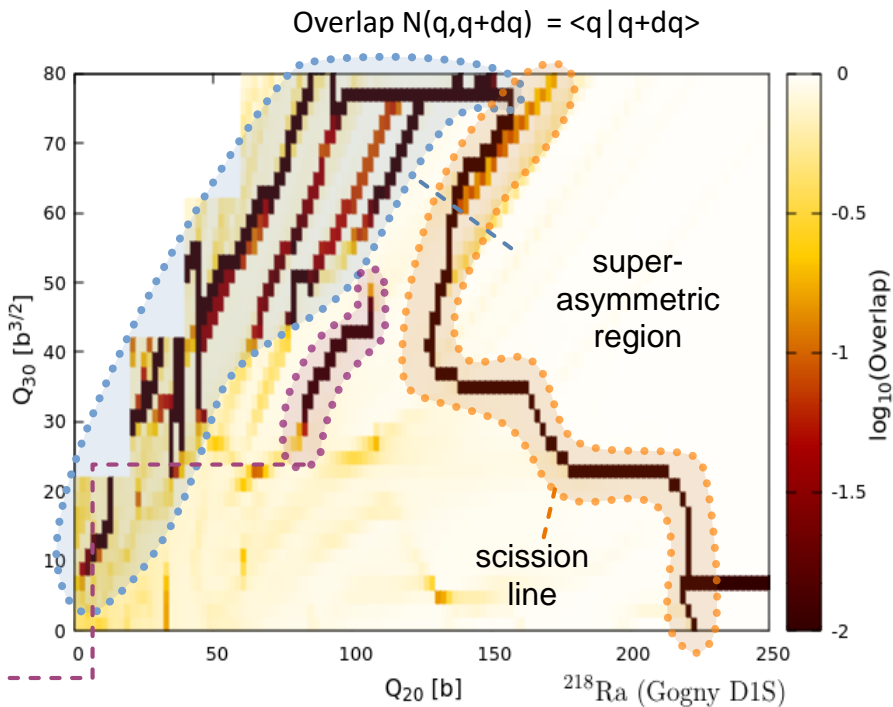
Pink : Linear interpolation  
Brown : Adiabatic path  
Green : Smoothed

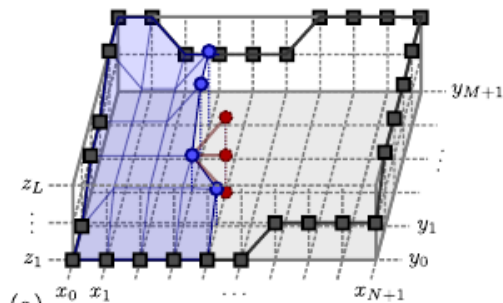
Identification : overlap or density distance

<sup>214</sup>Ra Potential Energy Surface



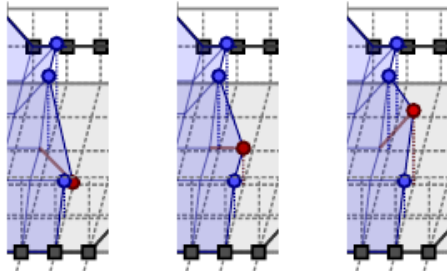
2D  
discontinuity





(a)

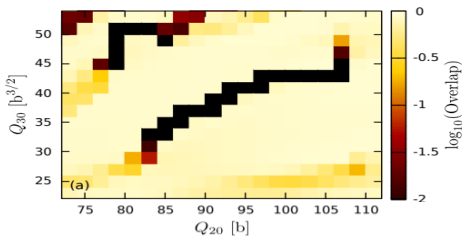
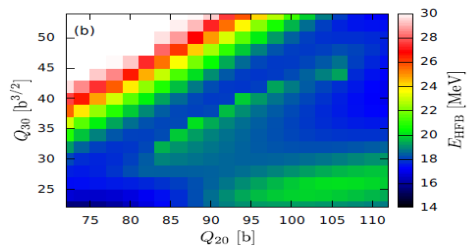
(b)



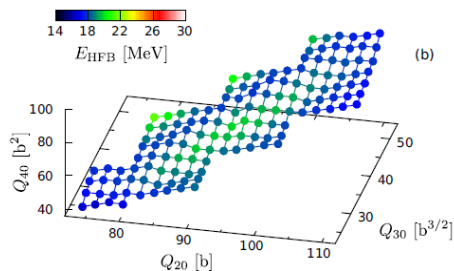
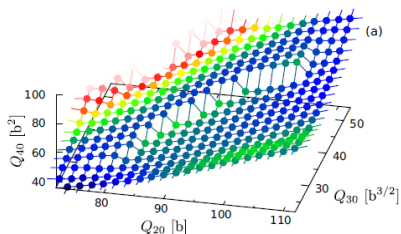
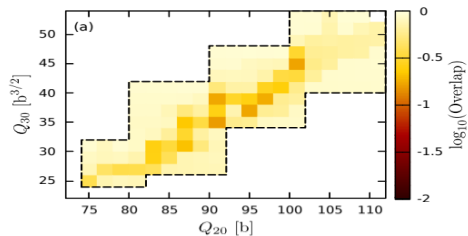
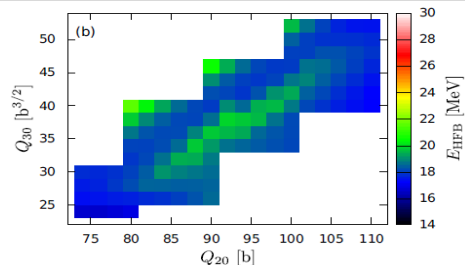
Number of possible surfaces =  $L^{M \cdot N}$

- (1) The initial set of partial surfaces contains only one surface consisting of the boundary points.
- (2) Choose the next point in the  $(x, y)$  plane to add to the surface
- (3) For each possible  $z$ -value at this point, generate a successor surface from each existing surface by the point. Keep the surface that lower the total energy
- (4) Repeat from step 2 until the entire surface within the boundaries is filled.

Number of calculated surfaces in the 2D DPM  
 $= M \cdot N \cdot L^{M+1}$



“Frontier  
DPM”



Avoid the addition of another full dimension → keep the computational resources for other purposes

No need to get the adiabatic PES first

- Shell effects have been studied by Strutinsky corrections to explain asymmetric fission on the way to fission
- Recent studies based on microscopic theories are focused on the 1D path scission point
- What are the compound nucleus shell effects along the asymmetric fission path?
- What are the role of fragment spherical and deformed shell effect in fission ?

Quantum shells stabilizing fission fragments with various shapes have been invoked as a factor determining the distribution of nucleons between the fragments at scission.

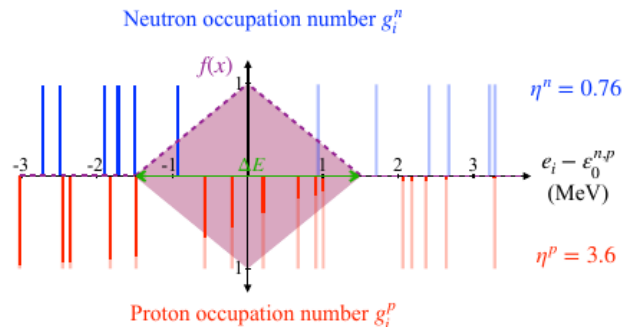
Spherical :  $Z = 50$ ,  $N = 82$

Deformed :  $Z = 52, 54, 56$   $N = 88$

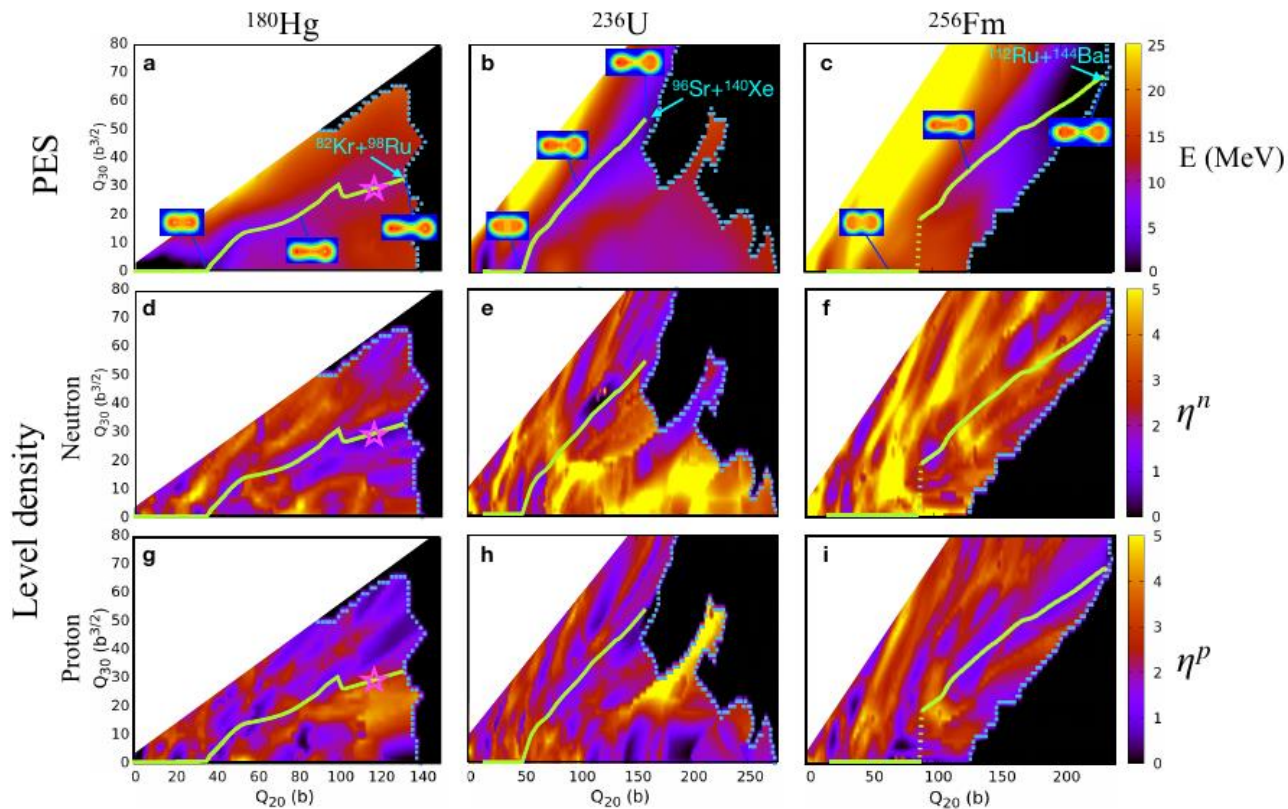
Other shells are also expected in the fissioning nucleus.

Idea : looking at the shell effects only at the Fermi level within all the PES

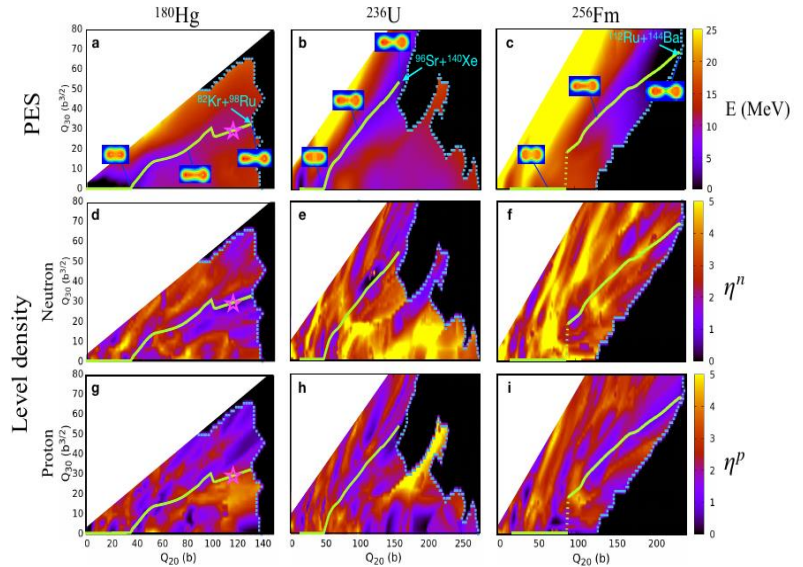
- counting the number of single particle energy states in an energy window around the Fermi level
- applying a smoothing function  $f$



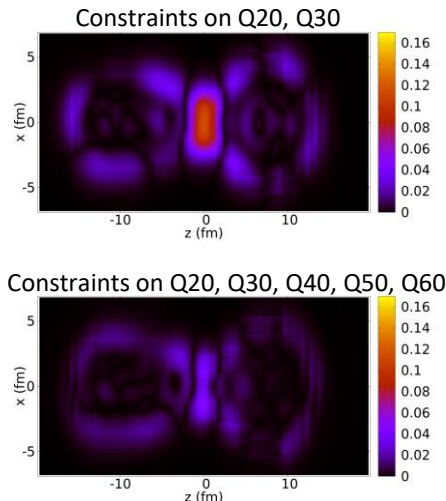
- Diagonalisation of the  $h$  part of the Bogoliubov matrix to get Single Particle Energy
- Fermi level as the middle of the gap
- Energy window : a few MeV around the Fermi level



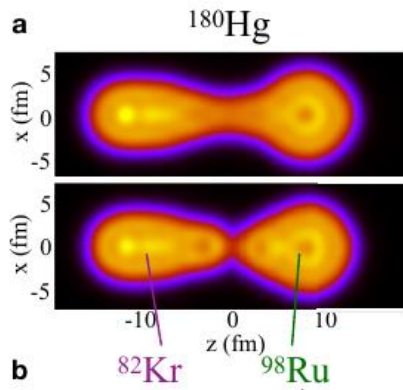
- Asymmetric paths follow Fermi shell effects
- 2 to 3 proton/neutron shell effects are on the path of asymmetric fission
- some 1D path discontinuities may be understood looking at the competition of shell effects



Total density difference  $\left| \rho_{comp}(\vec{r}) - \rho_{frag}(\vec{r}) \right|$



Total density



- Localizing the density minimum in the neck
- Integrating the space densities to get the fragment properties
- HFB calculation for (pre)fragments with 5 constraints on multipole moments from Q20 to Q60 to lower  $\int d\vec{r} \left| \rho_{comp}(\vec{r}) - \rho_{frag}(\vec{r}) \right|$  from 3,76 (2 constraints) to 1,90 (5 constraints)

Prefragments remains the 82Kr and 98Ru within the last 180 neutron shell effect

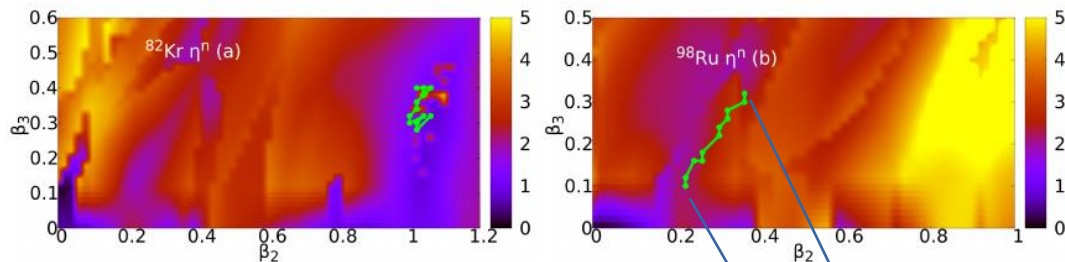
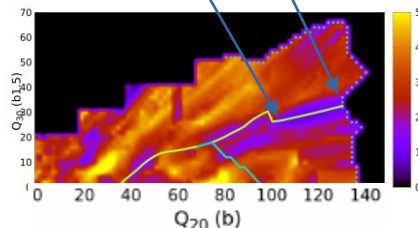


FIG. 2: Neutron sld of  $^{82}\text{Kr}$  in panel (a) and  $^{98}\text{Ru}$  in panel (b). The green dotted lines give the prefragment deformation during the final stages of the asymmetric  $^{180}\text{Hg}$  fission. The energy window around the Fermi level is fixed to 5 MeV for medium mass nuclei. Constraints are done on  $Q_{20}$  and  $Q_{30}$  and higher multipole moments are let free.

- No low sld in the proton side
- *Andreyev PRL 2010* : « This is a new type of asymmetric fission, not caused by large shell effects related to fragment magic proton and neutron numbers, as observed in the actinide region. »
- Conclusions : N=46 driver at the last stages



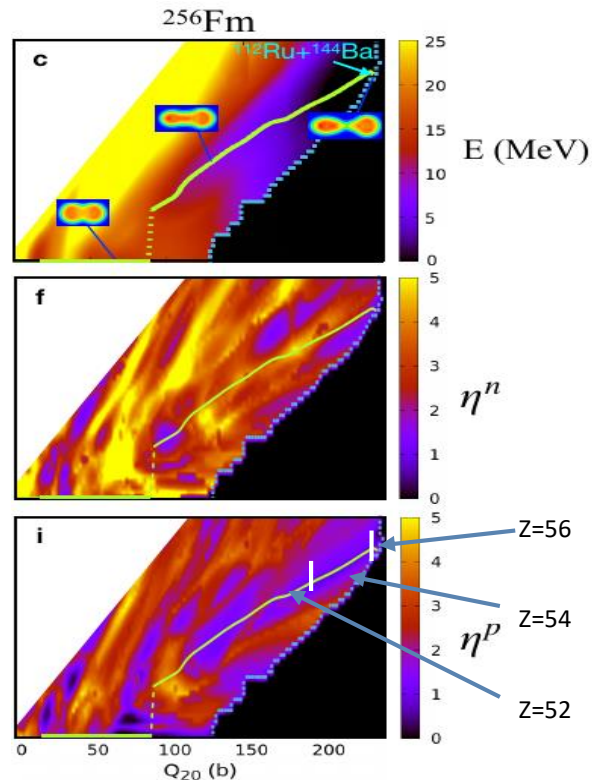
Are the 1D path  $Z_{\text{frag}}$  and  $N_{\text{frag}}$  due :

- Always to fragment shell effects ?  
(configuration at scission)
- To other Prefragment shell effects ?  
(configurations before scission)

→ stable prefragments before scission (82Kr and 98Ru)

→ successive shell effects from several prefragments during the last shell effect before scission

- 256Fm – Spherical  $Z=52$   
 – Deformed  $Z=54, 56$
- 236U – Spherical  $Z=52, N=82$   
 – Deformed  $Z=38, 40, 54, N=58, 62$



+ Local PES topology may be convex at the asymmetric valley around fission (ie 236U, ...)

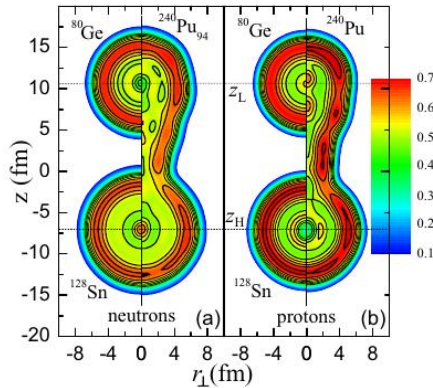
**Discontinuities :**

- The cleaning of the Potential Energy Surfaces from discontinuities is required for the TDGCM wave function evolution
- Detection of discontinuities done with simple tools : mean field states overlap or density distance operator
- an adapted Dynamic Programming Method used to clean discontinuities in 1 or 2 D involving a local  $N+1$  dimension calculation

**Shell effects :**

- The 1 D paths in PES favored by low sld at the Fermi surfaces
- Several proton and neutron shell effects guiding the asymmetric path
- Final shell effects correlated to (pre)fragment ones : several deformed and/or spherical fragment shell effects to lead the 1D path

## Localization function



Sadhukhan et al, PRC 2017

## Quantum loca. of the fragments

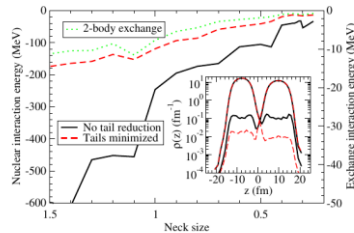


FIG. 1 (color online). Interaction energies plotted as a function of neck size ( $Q_N$ ). The solid black and red dashed curves are the nuclear interaction energies before and after localization, respectively, (energy scale on the left y axis), and the dotted green curve is the exchange part of the 2-body component of the interaction energy (energy scale on the right y axis). The inset shows the effect of the localization on the densities of the fragments at scission.

Younes et al, PRL 2011

