

Laser Spectroscopy of Highly Charged Ions

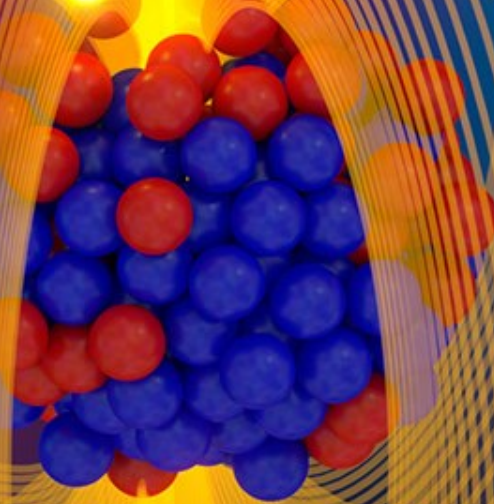
Wilfried Nörtershäuser



TECHNISCHE
UNIVERSITÄT
DARMSTADT



L A S E R
SpHERE



DE LA RECHERCHE À L'INDUSTRIE

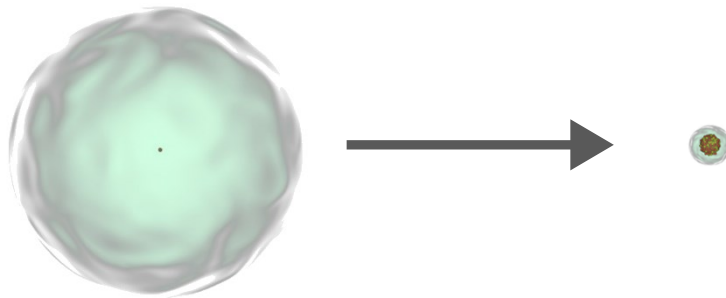
cea

ESNT

Espace de Structure Nucléaire Théorique

Workshop
“Laser spectroscopy as a
tool for nuclear theories”

Interest in Highly Charged Ions (HCI)



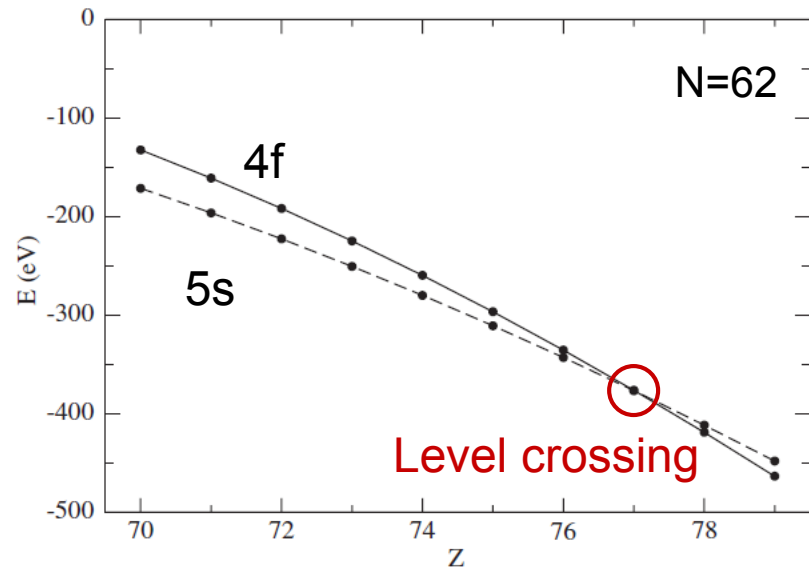
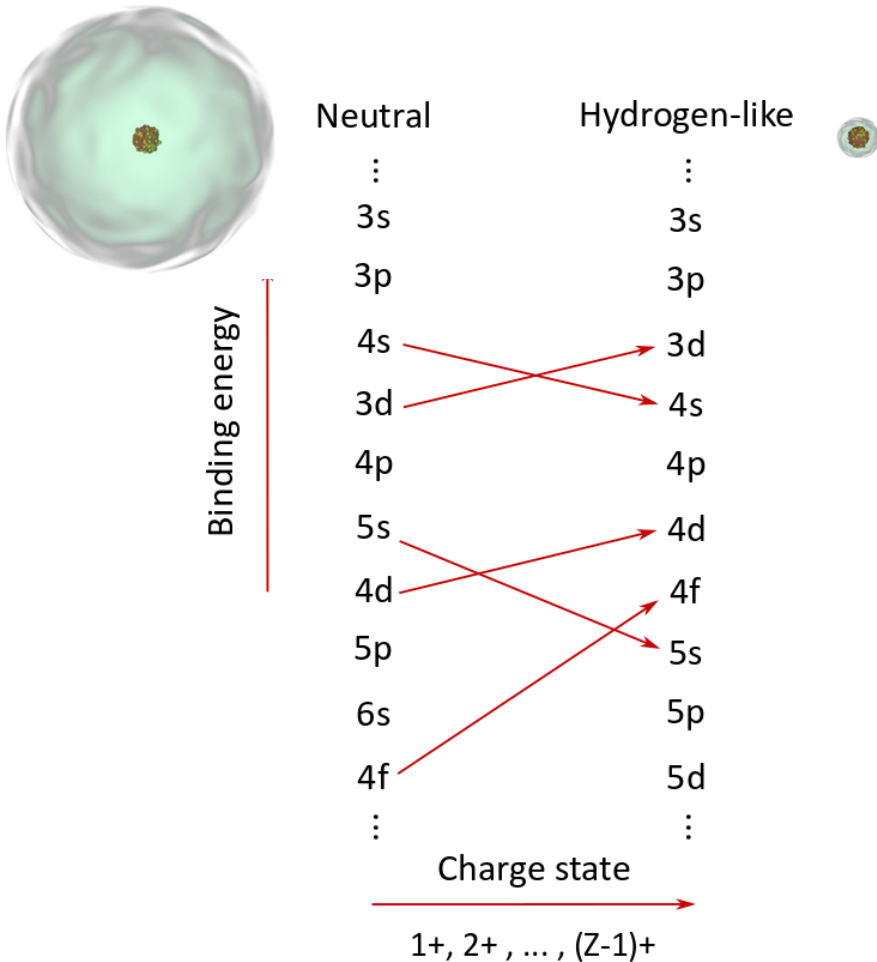
Radius	Z^{-1}
Polarizability	Z^{-1}
Gross structure	Z^2
Fine structure splitting	Z^4
Hyperfine structure	Z^3
QED effects	Z^4

- Relativistic, QED, nuclear effects
- Forbidden optical transitions sensitive to α , μ , violation of local Lorentz invariance

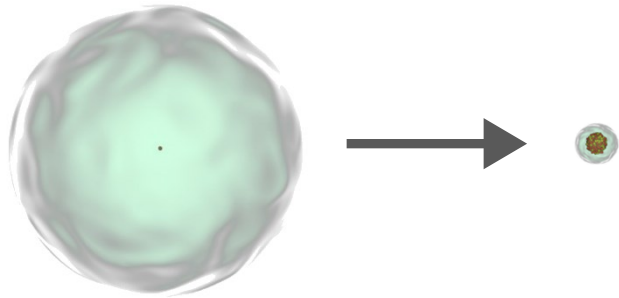
Kozlov et al., Rev. Mod. Phys. **90**, 045005 (2018) , arXiv:180306532



Optical Transitions in HCI



Advantages of HCl for High-Precision Experiments and Optical Clocks



Linear Stark shift	Z^{-1}
Second order Stark shift	Z^{-4}
Linear Zeeman shift	Z^0
Second order Zeeman shift	$Z^{-3\dots-4}$
Electric quadrupole shift	Z^{-2}

Less sensitive to external perturbations



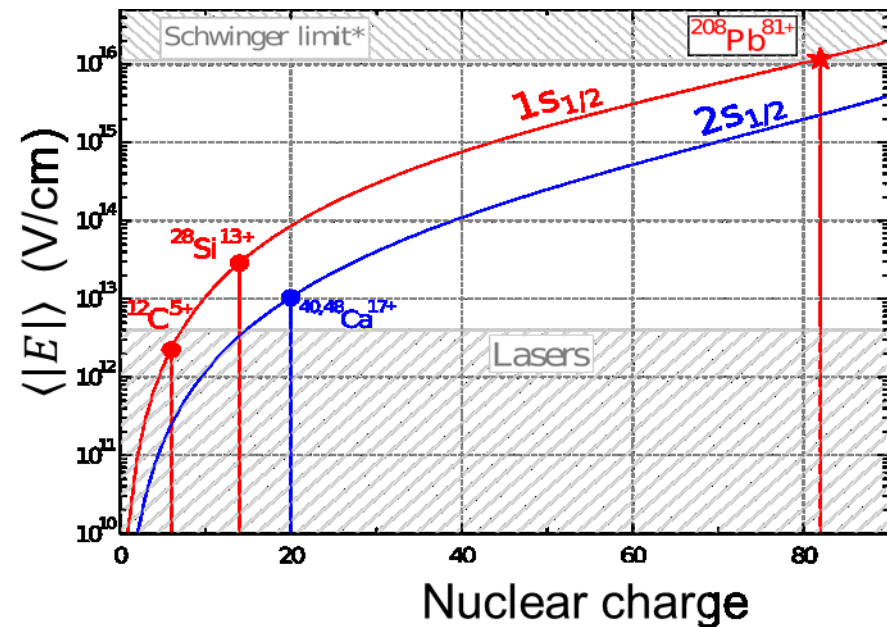
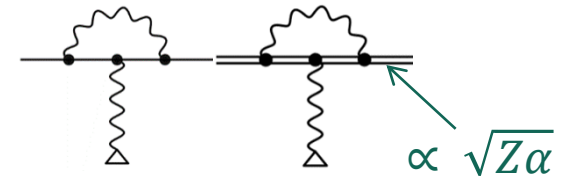
QED in Highly Charged Ions (HCI)

- Impressive theory predictions and most precise experimental results, e.g. electron magnetic moments: $g-2$, g -factor in HCI [1]
 - Most stringently tested theory in weak fields

- Validity of QED in strong fields?
Test of bound state QED (BS-QED) under extreme conditions in high electric and magnetic fields of heavy HCI

- g -factor measurements in H or Li-like ions
- Fine structure and hyperfine structure spectroscopy

→ Combine strongest field and highest precision



[1] H. Häfner, et al., Phys. Rev. Lett. 85, 5308 (2000)

J.Verdu, et al., Phys. Rev. Lett. 92, 093002 (2004)

[2] S. Sturm, et al., Nature 506, 7489 (2014)

[3] S. Sturm, et al., Phys. Rev. Lett. 107, 023002 (2011)

[4] F. Köhler, et al. Nat. Comm. 7, 10246 (2016)



Outline

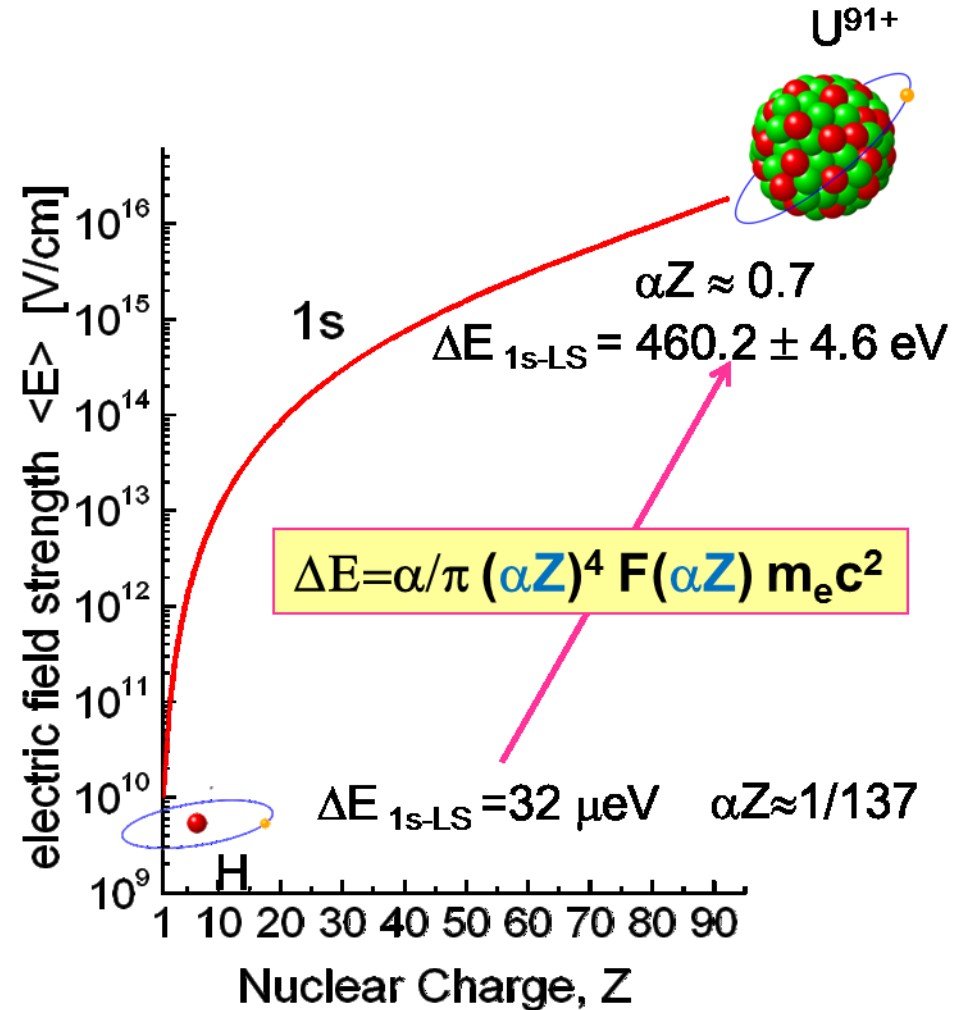
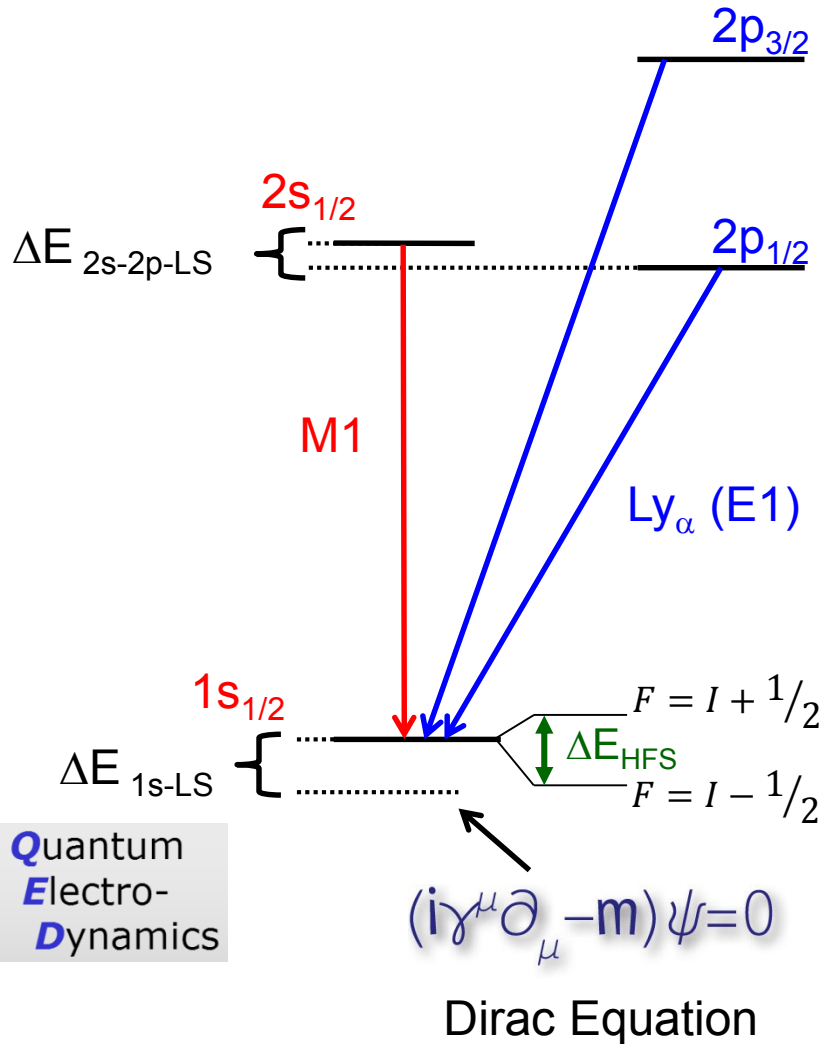
Laser Spectroscopy in Heavy Highly Charged Ions and the
„Hyperfine Puzzle of strong-field bound-state QED“

Nuclear Magnetic Resonance and its contribution to a solution of
the „Hyperfine Puzzle“

Into the future: Experimental developments towards high-precision
spectroscopy in HCI

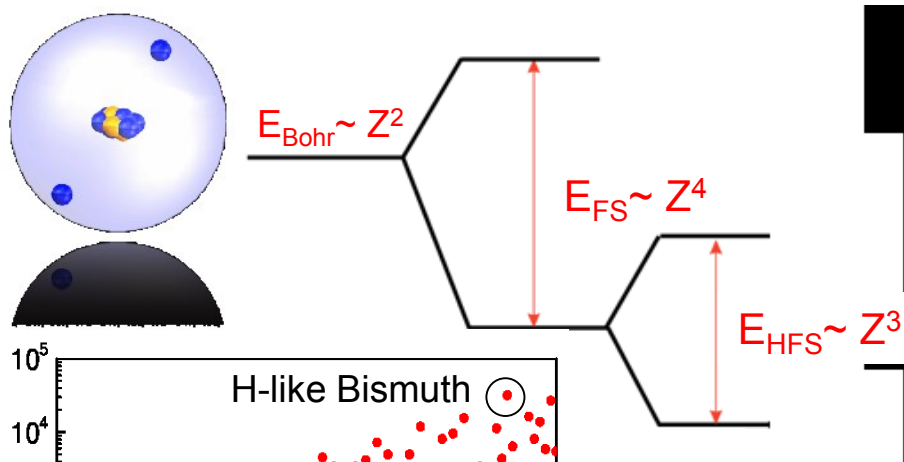


Structure of Heavy H-Like Ions

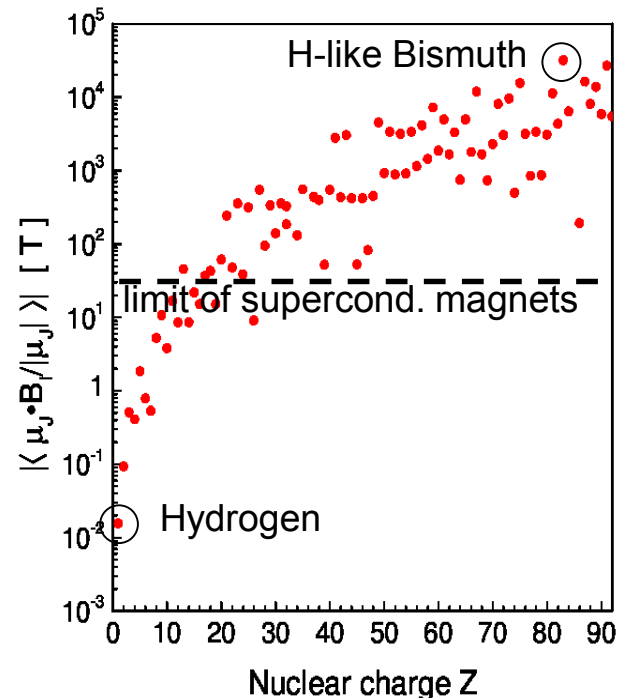
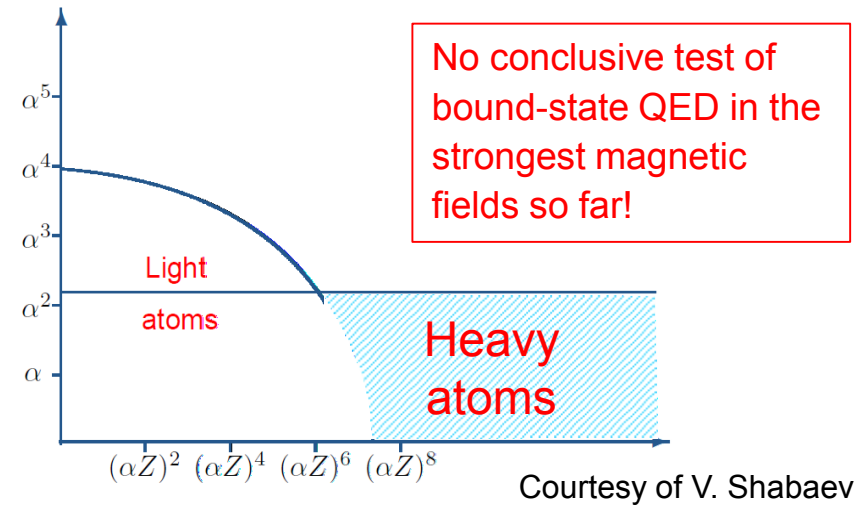


Test of BS-QED in Highly Charged Ions

atomic structure:



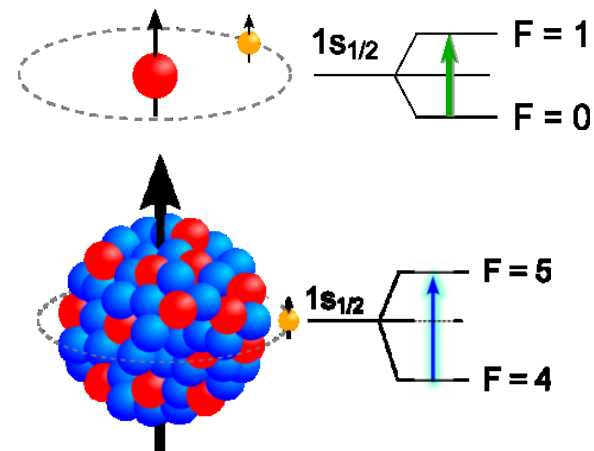
Tests of QED to lowest orders in α and to all orders in αZ



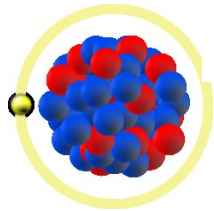
H: $\lambda = 21 \text{ cm}$
 $\tau = 11 \text{ Ma}$

$^{209}\text{Bi}^{82+}$: $\lambda = 244 \text{ nm}$
 $\tau = 0.39 \text{ ms}$

$^{209}\text{Bi}^{80+}$: $\lambda = 1555 \text{ nm}$
 $\tau = 83.3 \text{ ms}$



HFS-Measurements in Heavy H-Like Ions



$$\Delta E^{(1s)} = \underbrace{\frac{4}{3} \alpha (\alpha Z)^3 \frac{\mu_I}{\mu_N} \frac{m}{m_p} \frac{2I+1}{2I} mc^2}_{\text{Dirac}} \times \underbrace{(A(\alpha Z)^{1s} (1 - \epsilon_{BR}^{1s}) (1 - \epsilon_{BW}^{1s}) + x_{red}^{1s})}_{\text{nuclear contributions}} + \underbrace{x_{red}^{1s}}_{\text{QED}}$$

Bi⁸²⁺

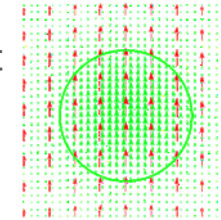
Dirac

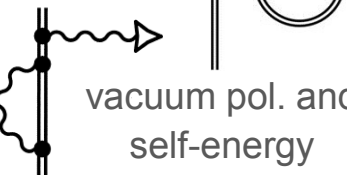
nuclear contributions

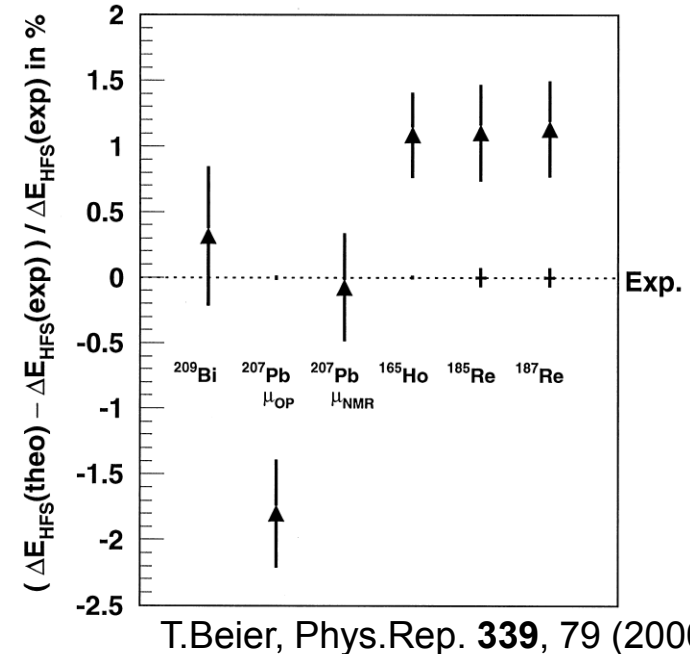
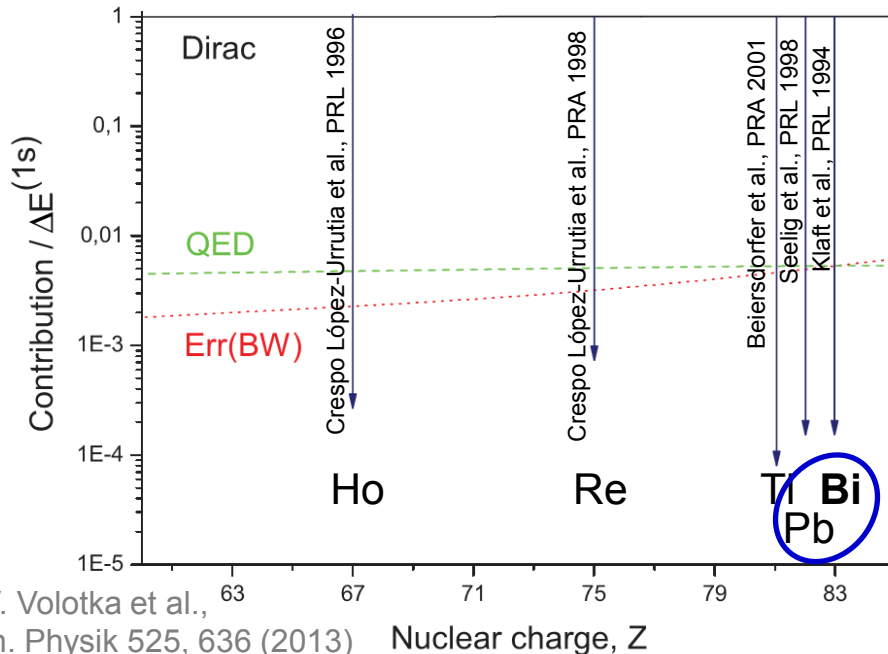
QED

known (?) 

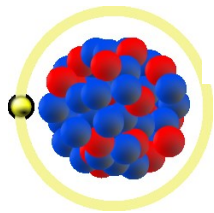
Bohr-Weisskopf effect
Breit-Rosenthal effect
nuclear polarization



vacuum pol. and self-energy 

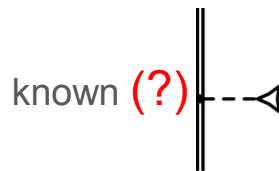


Theory: Trouble & Solution

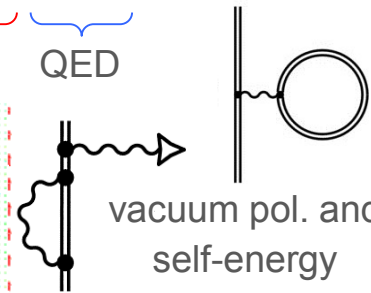
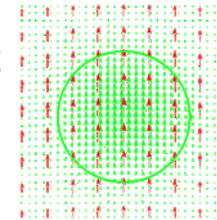


$$\Delta E^{(1s)} = \underbrace{\frac{4}{3} \alpha (\alpha Z)^3 \frac{\mu_I}{\mu_N} \frac{m}{m_p} \frac{2I+1}{2I} mc^2}_{\text{Dirac}} \times \underbrace{(A(\alpha Z)^{1s} (1 - \epsilon_{\text{BR}}^{1s}) (1 - \epsilon_{\text{BW}}^{1s}) + x_{\text{red}}^{1s})}_{\text{nuclear contributions}} \underbrace{+ x_{\text{red}}^{1s}}_{\text{QED}}$$

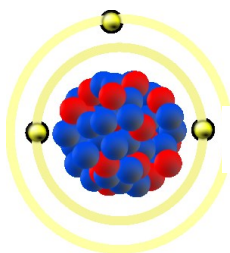
Bi⁸²⁺



Bohr-Weisskopf effect
Breit-Rosenthal effect
nuclear polarization



uncertainty
Bohr-Weisskopf > QED contribution



Bi⁸⁰⁺

$$\Delta E^{(2s)} = \frac{1}{6} \alpha (\alpha Z)^3 \frac{\mu_I}{\mu_N} \frac{m}{m_p} \frac{2I+1}{2I} mc^2 \times (A(\alpha Z)^{2s} (1 - \epsilon_{\text{BR}}^{2s}) (1 - \epsilon_{\text{BW}}^{2s}) + x_{\text{red}}^{2s} + B_Z + C_Z)$$

specific difference (V. M. Shabaev et al. PRL 86 (2001))

$$\Delta'E = \Delta E^{(2s)} - \xi \Delta E^{(1s)}$$

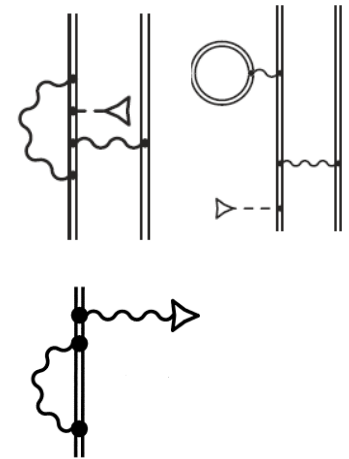
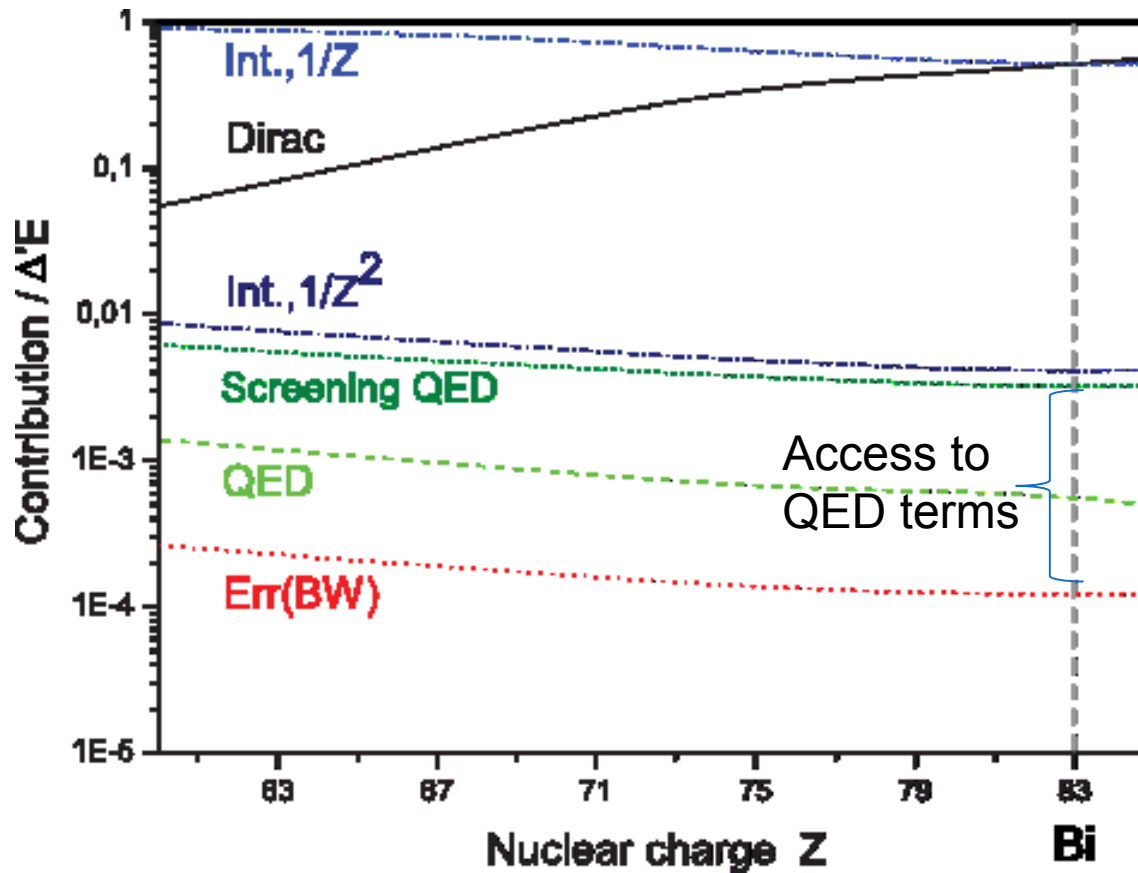
nuclear contributions removed



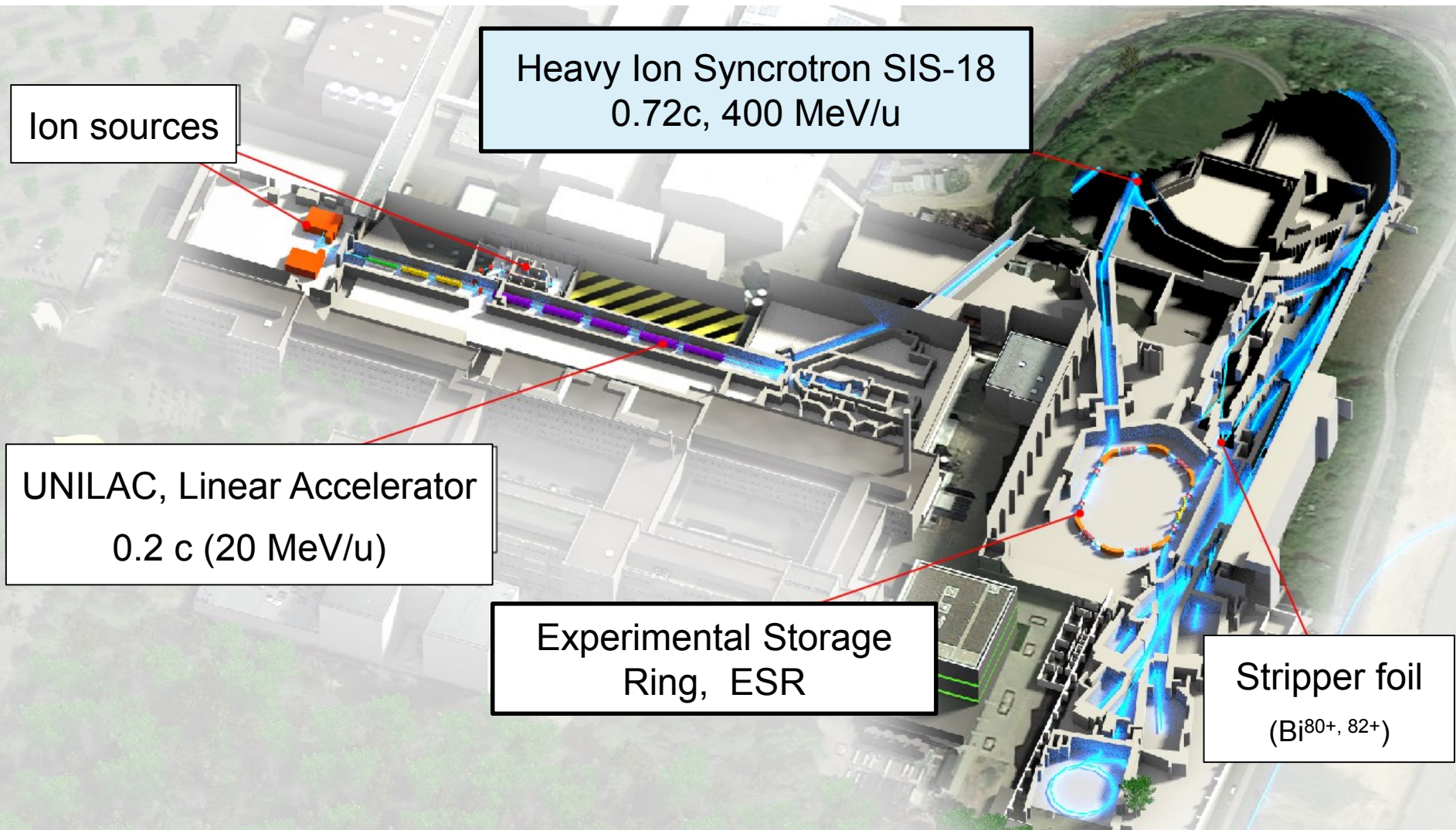
Precision Test of QED: The Specific Difference

specific difference (V. M. Shabaev et al. PRL 86 (2001))

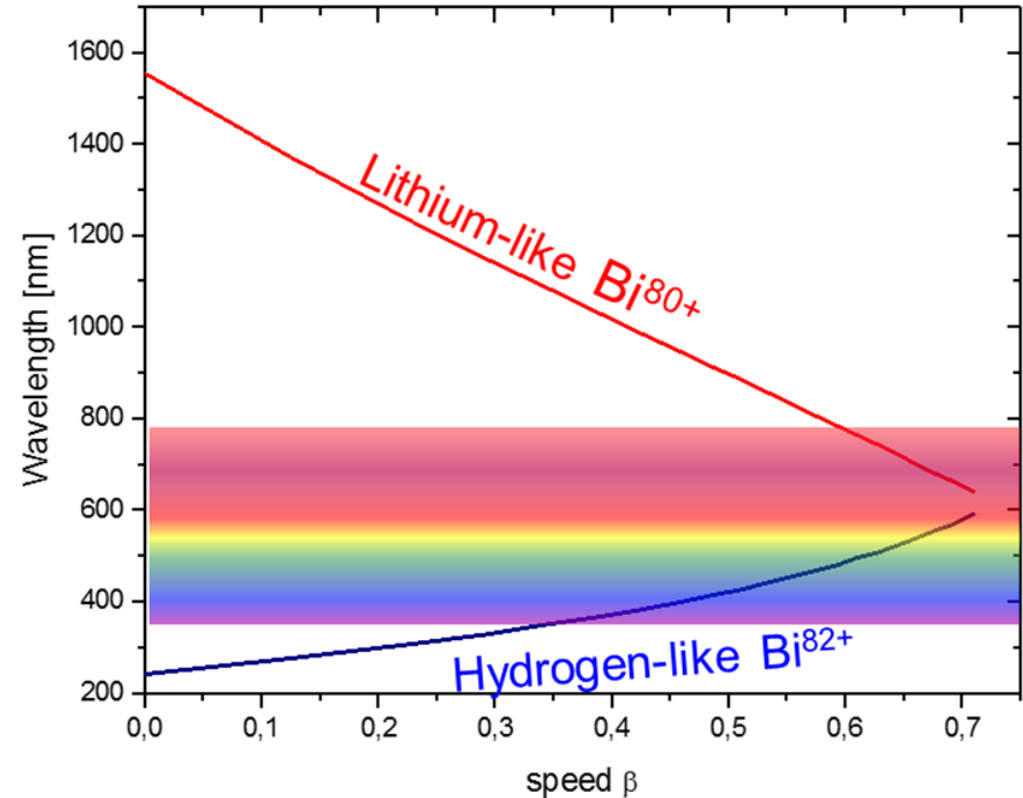
$$\Delta'E = \Delta E^{(2s)} - \xi \Delta E^{(1s)}$$



Production of $^{209}\text{Bi}^{80,82+}$



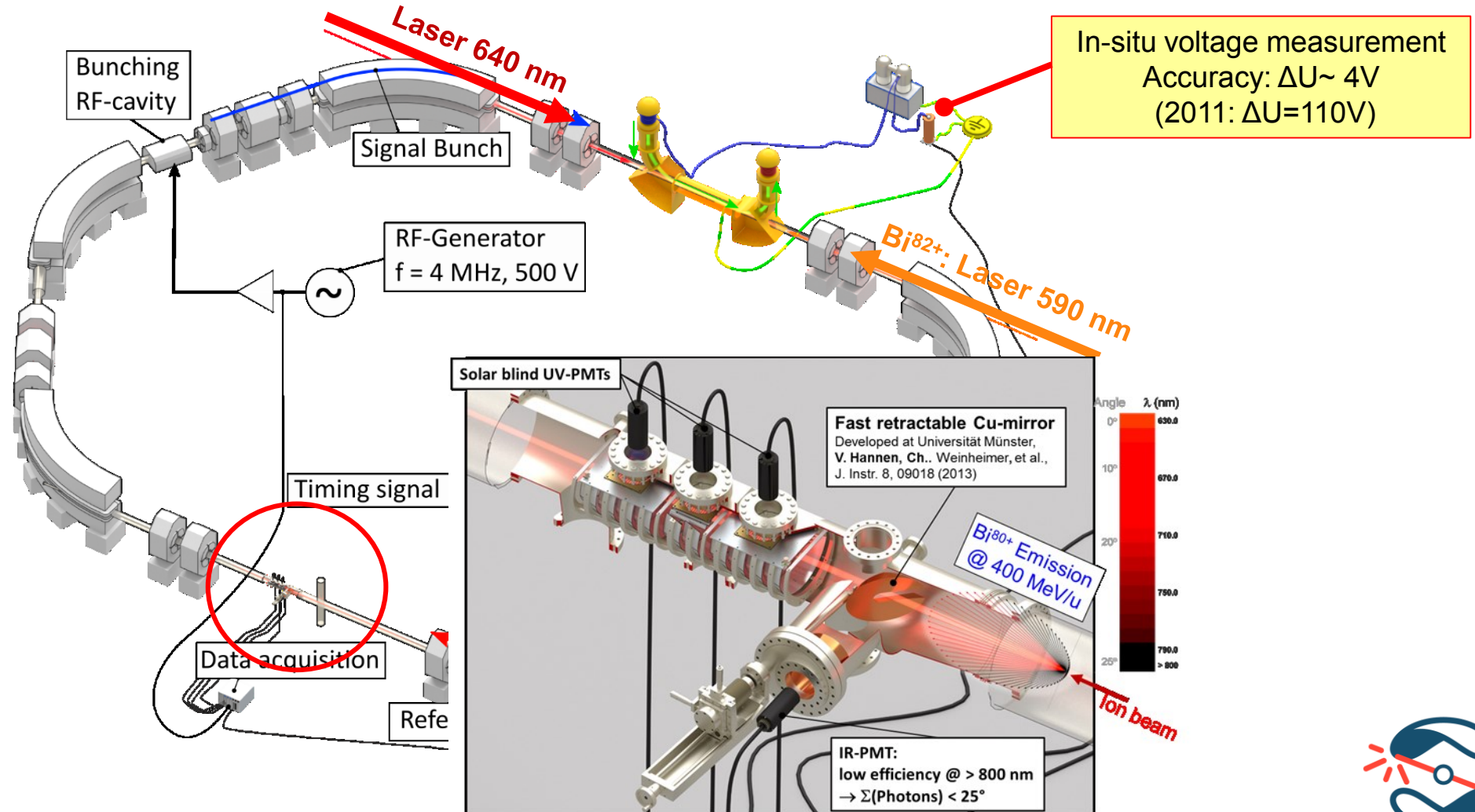
Reminder: Optical Doppler Effect



... for Dummies !



First High Accuracy Measurement of $\Delta'E$



The Experimental Storage Ring ESR



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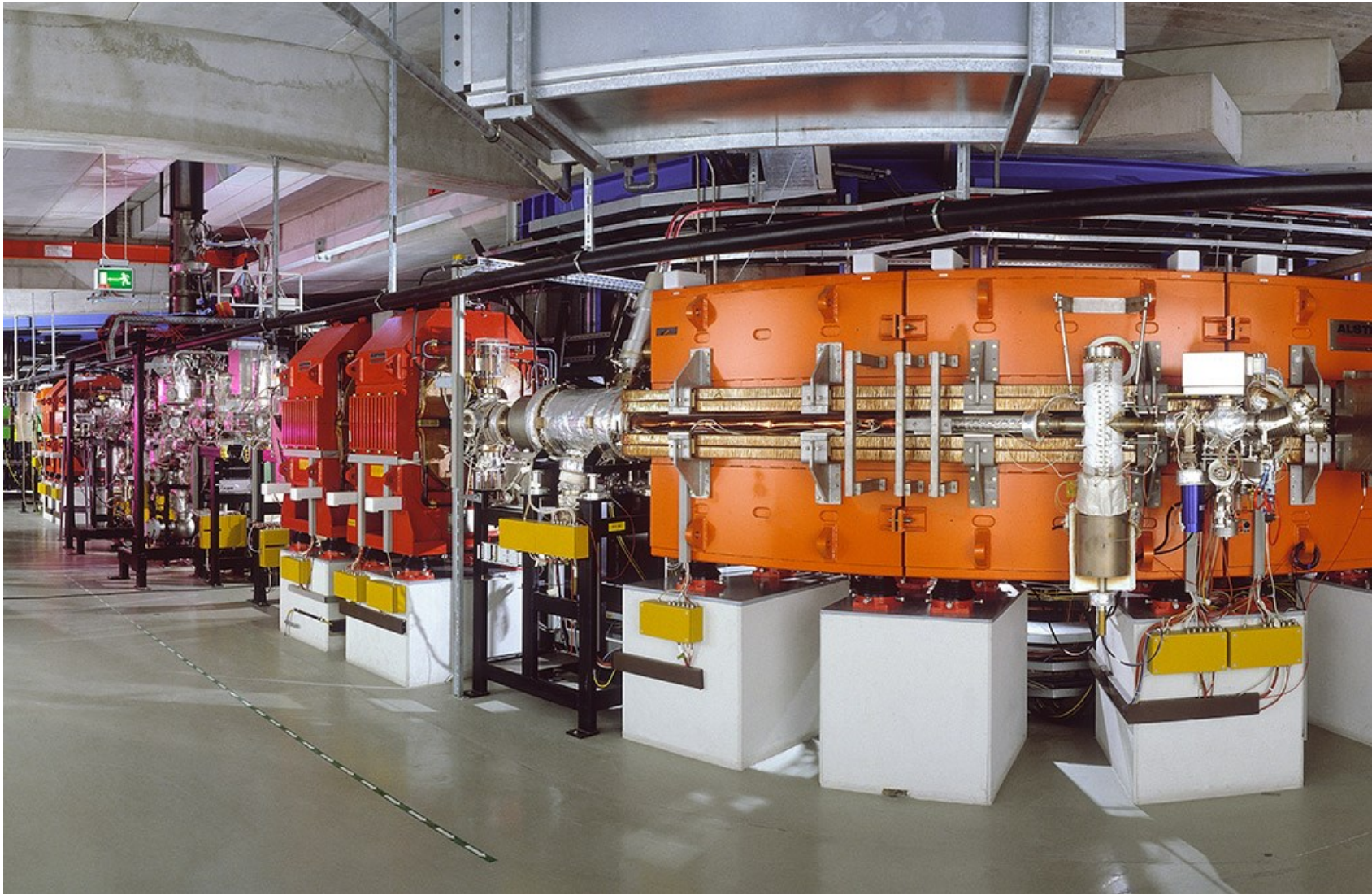


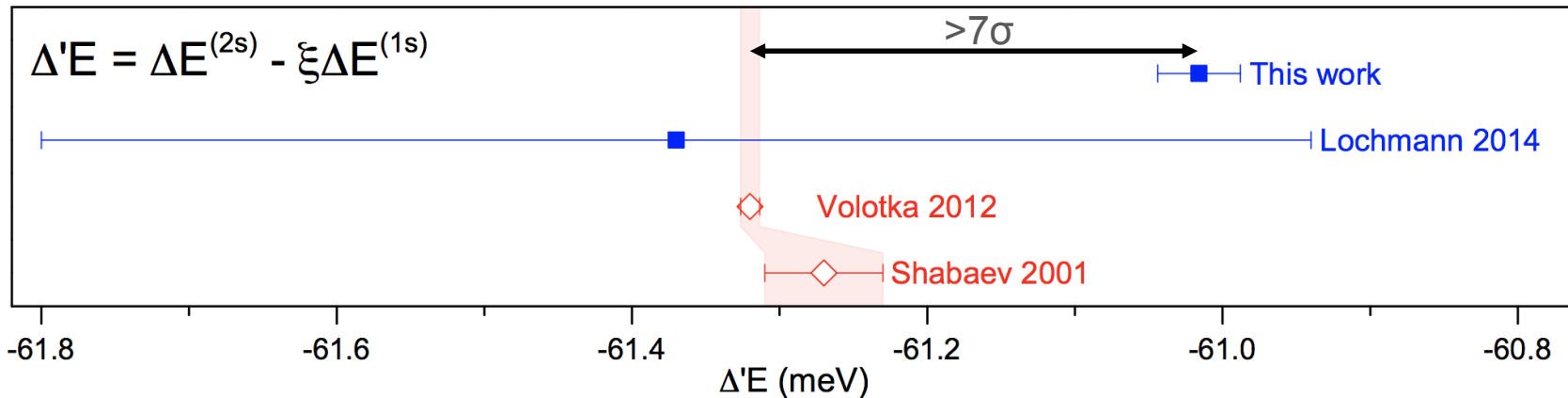
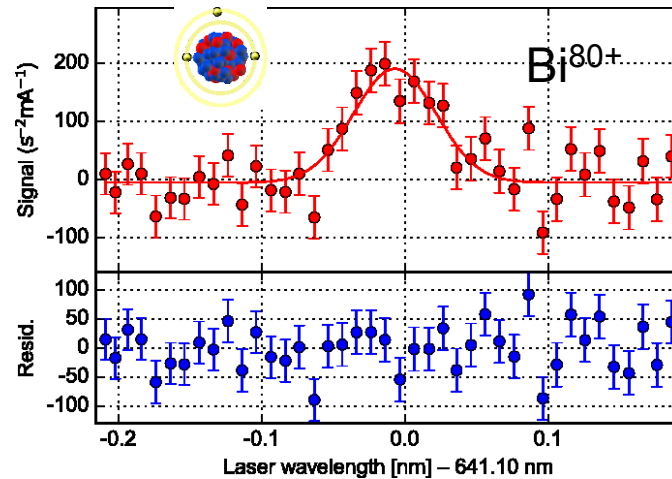
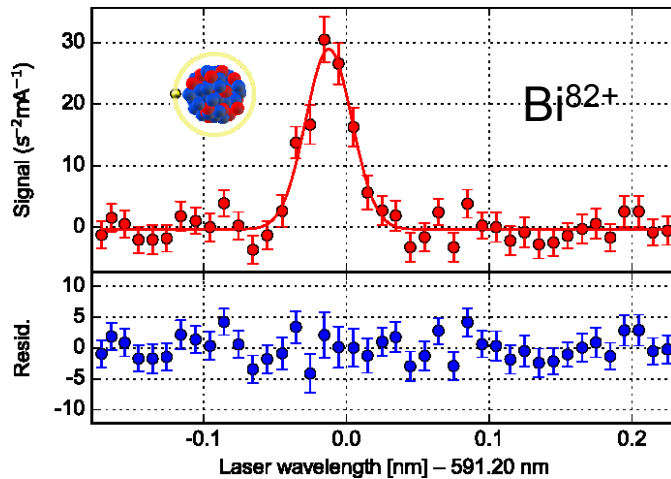
Foto:
J. Mai,
GSI



Results: Hyperfine Transition Wavelength



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Explanations for the Discrepancy ?

Experiment

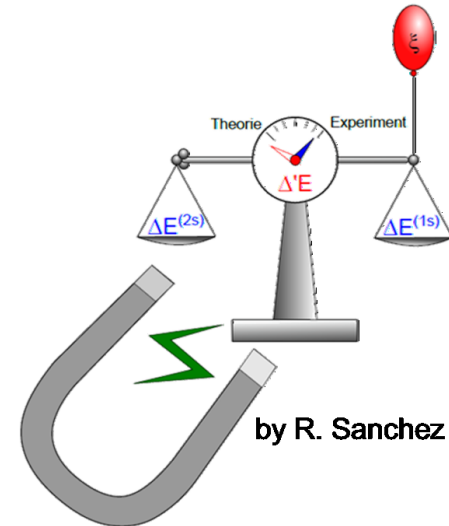
- ~~Our experimental value is wrong?! unlikely~~
- The literature value of the nuclear magnetic moment of ^{209}Bi is wrong?!
(Dirac term of the ground state hfs is proportional to the experimental value)

Theory

- The specific difference does not work as expected?!

$$\Delta'E = \Delta E^{(2s)} - \xi \Delta E^{(1s)}$$

- ~~BS Strong-Field QED is wrong?! First exclude other possibilities !~~



Challenge: Magnetic Moments Required

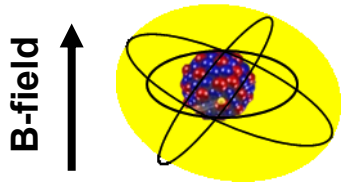
Author	$(1 - \sigma_{\text{dia}})^{-1}$	σ_{chem}	$\mu(^{209}\text{Bi}) [\mu_N]$
Proctor & Yu ¹ [27]			4,0400(7) ³
Ting & Williams ¹ [28]	1,0104		4,0810(4)
Flynn ¹ [31]			4,0391(2) ³
Raghavan ² [32]	1,0177		4,1106(2)
Baštuĝ et al. ² [33]	1,01757(6)		4,1103(5)
Gustavsson & Mårtensson-Pendrill ³ [30]		0,000(1)	4,110(4)

Bi(NO₃)₃ · 5H₂O solved in HNO₃ (nitric acid)

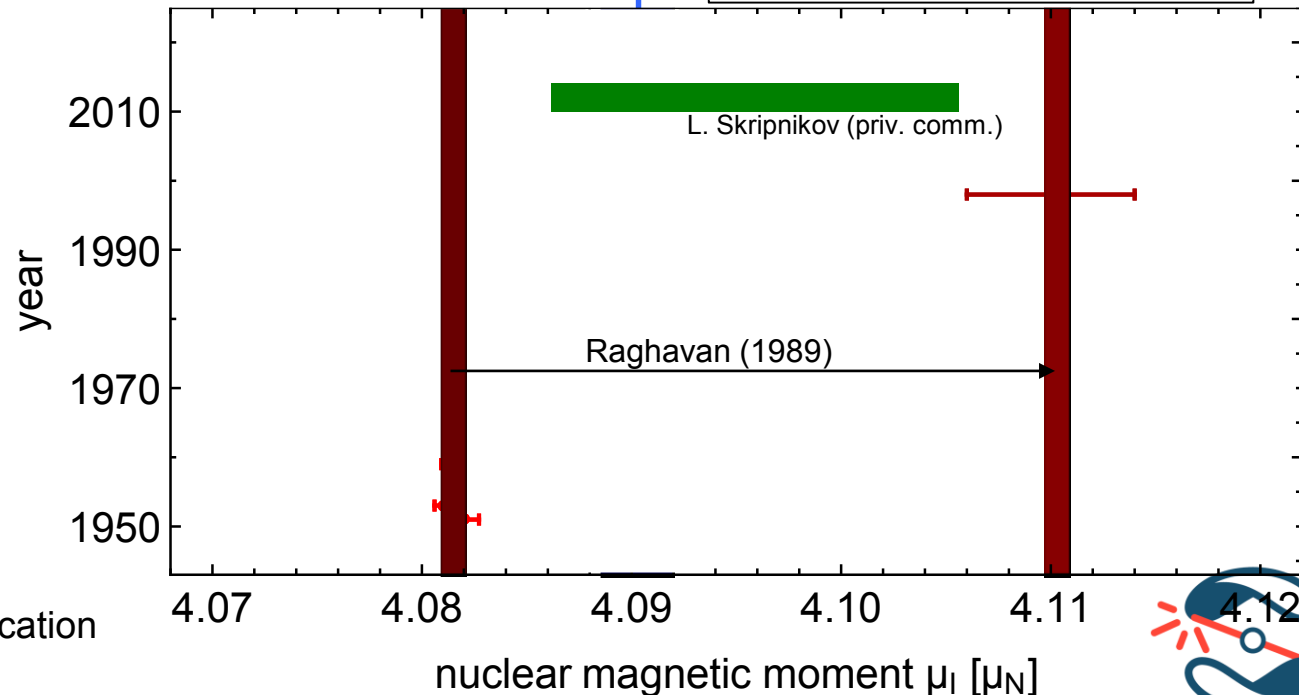
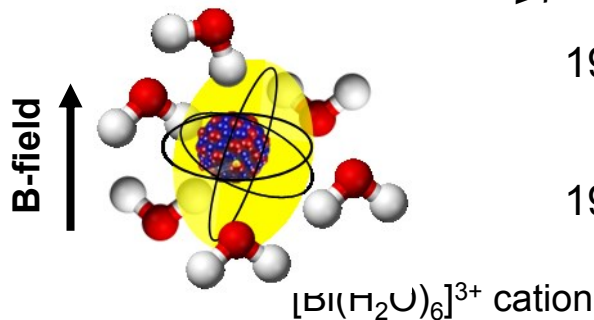
$$\mu_{I,\text{Exp}} = \frac{\Delta' E_{\text{Exp}}}{\Delta' E_{\text{Theo}}} \cdot \mu_{I,\text{Lit}}$$

$$\mu_{\text{Probe}} = \mu'_{\text{Probe}} / [1 - (\sigma_{\text{dia}} + \sigma_{\text{chem}})]$$

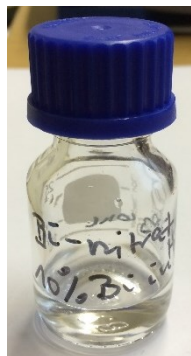
diamagnetic shielding



chemical shift



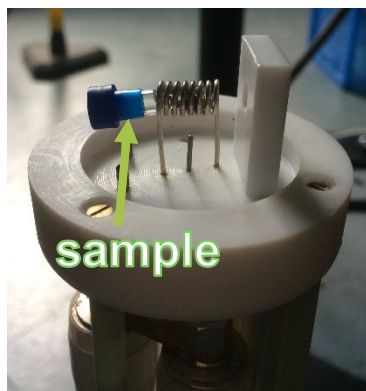
NMR Measurements @ TU Darmstadt



sample

$\text{Bi}(\text{NO}_3)_3$ solved
 HNO_3 (nitric acid)

samples prepared
at Institut für Kernchemie
(Uni Mainz)



Cooperation: Prof. Michael Vogel

TU Darmstadt, Institut für Festkörperphysik



Alexei Privalov

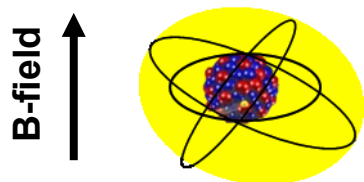
Benjamin Kresse



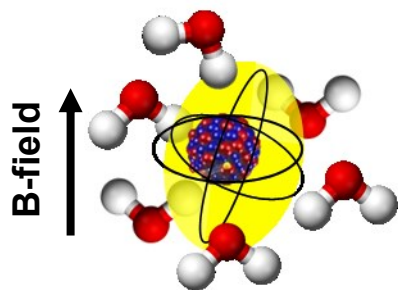
Chemical Shift in NMR Spectra of $\text{Bi}(\text{NO}_3)_3$ (aq)

$$\mu_{\text{Probe}} = \mu'_{\text{Probe}} / [1 - (\sigma_{\text{dia}} + \sigma_{\text{chem}})]$$

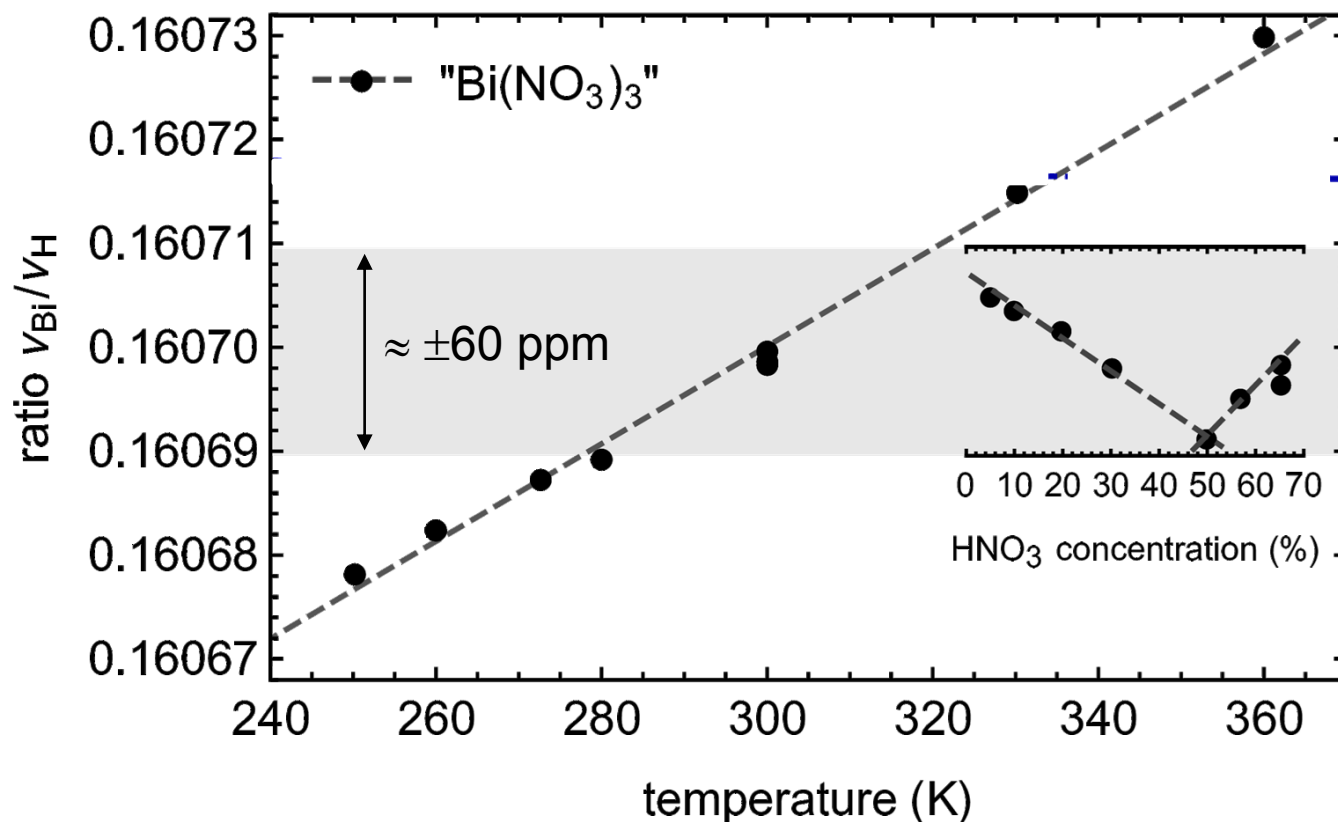
diamagnetic shielding



chemical shift



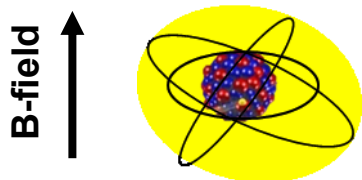
$[\text{Bi}(\text{H}_2\text{O})_6]^{3+}$ cation



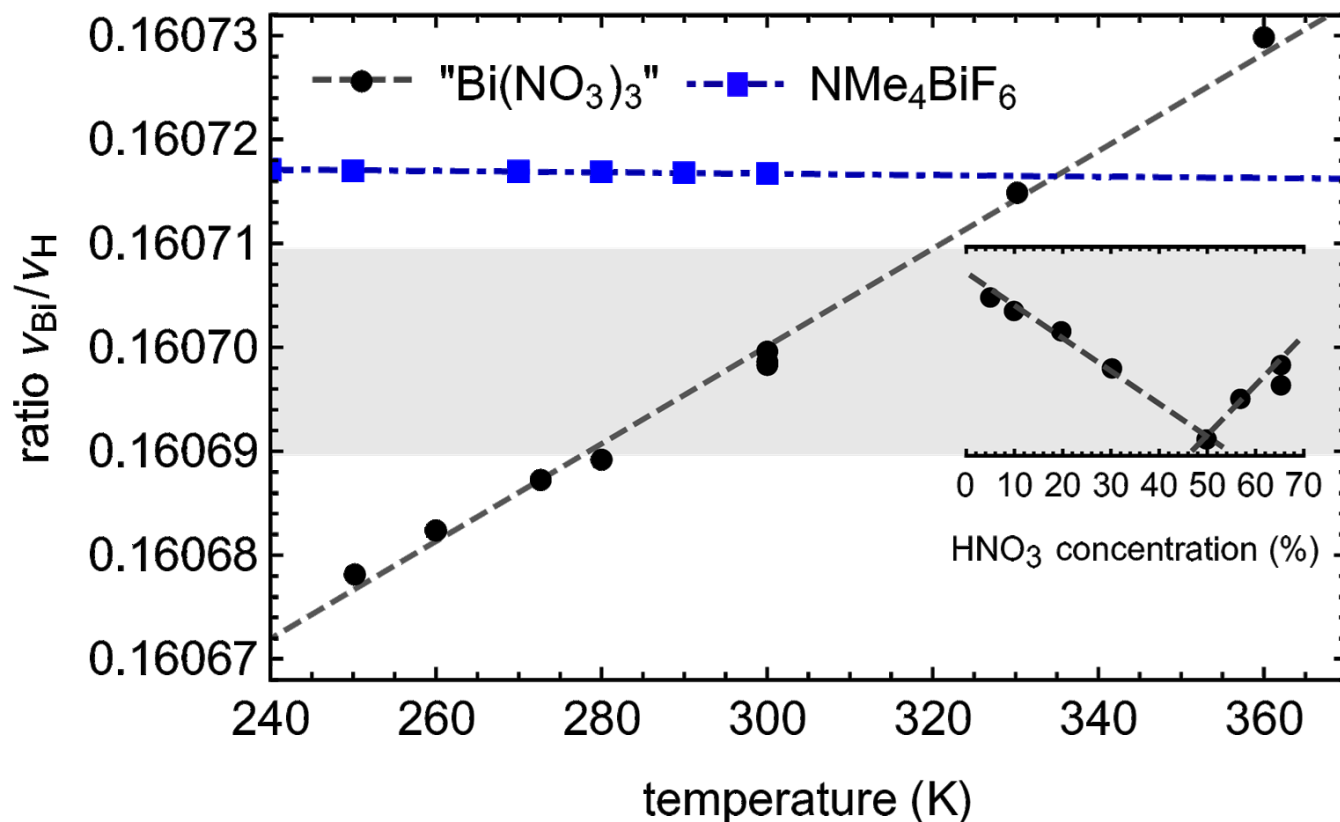
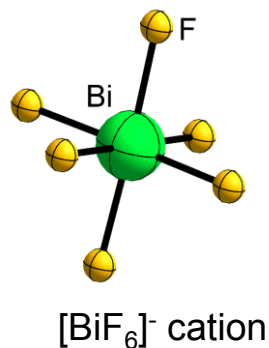
Chemical Shift in NMR Spectra of $\text{Bi}(\text{NO}_3)_3$ (aq)

$$\mu_{\text{Probe}} = \mu'_{\text{Probe}} / [1 - (\sigma_{\text{dia}} + \sigma_{\text{chem}})]$$

diamagnetic shielding



chemical shift

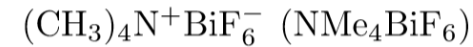
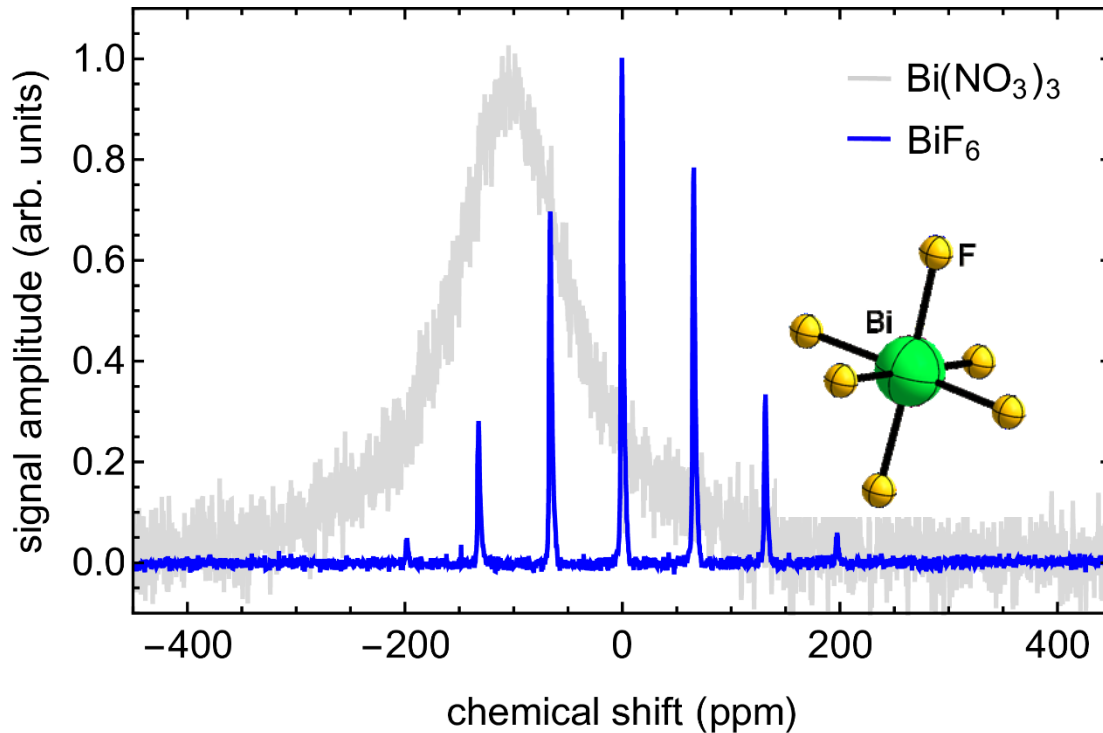


[Fedotov, M. A.; et al.
Neorg. Khim. 1998, 43, 307–310]

Previous measurement of BiF_6^- : Morgan et al., J. Magn. Res. **52**, 139 (1969)



NMR Signals



Old:

$\mu(^{209}\text{Bi}) = 4.1106(2) \mu_N$
Raghavan, At. Data Nucl. Data
Tables **42**, 189 (1989)

New:

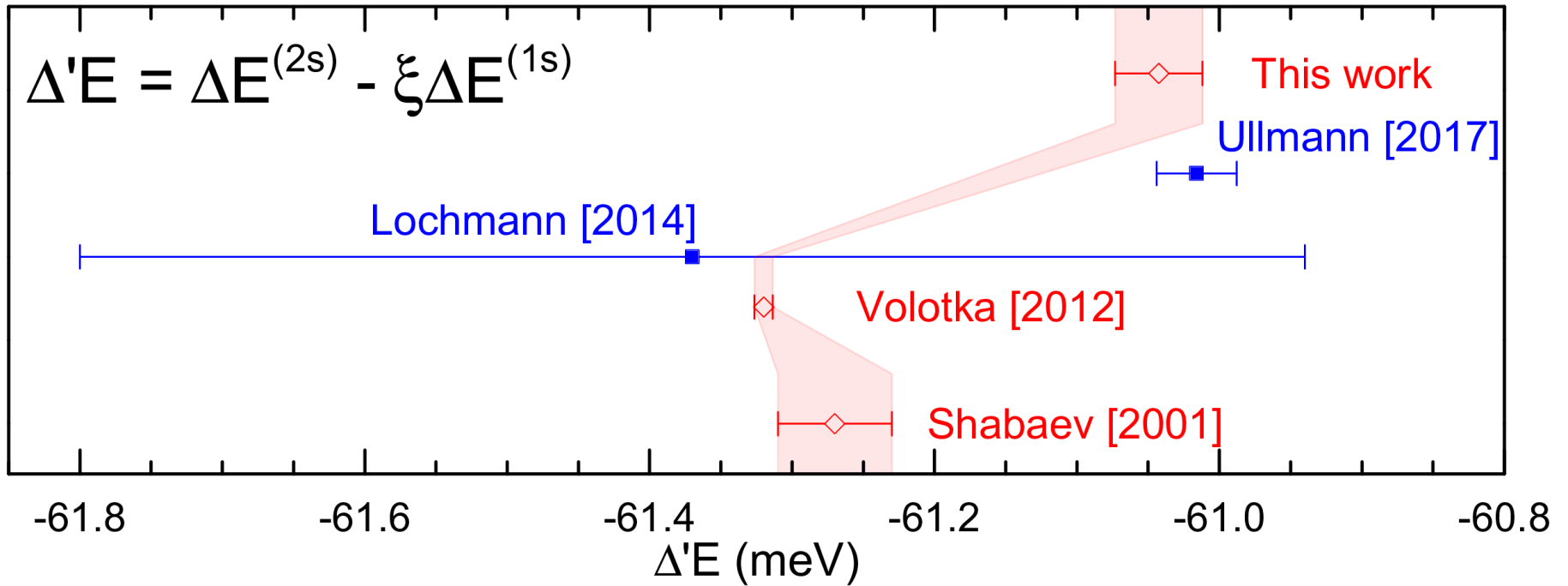
$\mu(^{209}\text{Bi}) = 4.092(2) \mu_N$
uncertainty dominated by theory

Diff. $> 8\sigma_{\text{comb}}$

	Theory		Experiment
	$\mu_I(\text{old})$	$\mu_I(\text{new})$	
$\Delta E^{(1s)}$	5112(-5/+20)	5089(-5/+20)(2)	5085.03(2)(9)
$\Delta E^{(2s)}$	801.9(-9/+34)	798.3(-9/+34)(4)	797.645(4)(14)



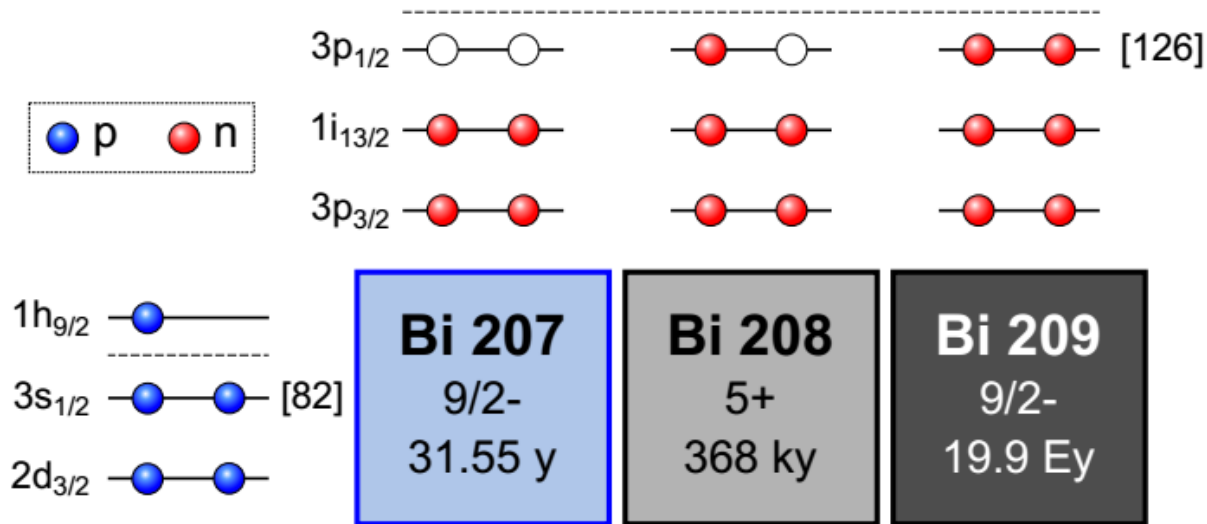
Result for $\Delta'E$



L. Skripnikov *et al.*, Phys. Rev. Lett. **120** 093001 (2018)



Bi Isotopes of Interest



$\Delta E^{(1s)}$ (eV)	5.085	5.661	5.08503(9)
$\lambda^{(1s)}$ (nm)	244	219	243.8221(44)
$\Delta E^{(2s)}$ (eV)	0.7976	0.8894	0.797645(14)
$\lambda^{(2s)}$ (nm)	1554	1394	1554.377(29)

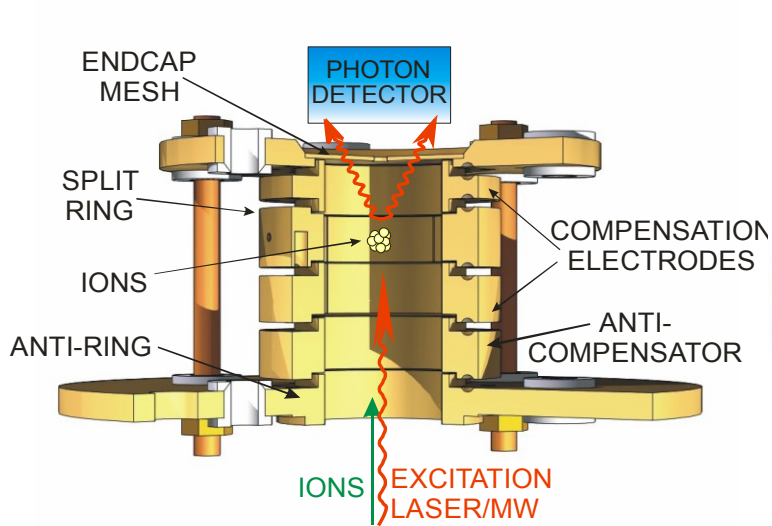


Towards Laser Spectroscopy of HCl in Ion Traps

- EBIT (MPIK)
Laser Fluorescence Spectroscopy (high T, pulsed lasers)
Mäckel *et al.*, PRL **107**, 143002 (2011)
- SPECTRAP (GSI)
Fluorescence Spectroscopy at 4K, Doppler limited
Andelkovich *et al.*, Phys. Rev. A **87**, 033423 (2013)
- ARTEMIS (GSI)
RF-Laser Fluorescence Double-Resonance Spectroscopy
Quint *et al.*, Phys. Rev. A **78** 032517 (2008)
- α TRAP (MPIK)
Spin-flip Detection with Continuous Stern-Gerlach Effekt (≈ 1 K)
Egl *et al.*, Phys. Rev. Lett. **123**, 123001 (2019)
- Cryogenic Paul Trap (PTB / MPIK)
Quantum Logic Spectroscopy (highest resolution)
Schmöger *et al.*, Science **347**, 1233 (2015)
Micke *et al.*, submitted to Nature (2019)



Long-Term Perspective at GSI



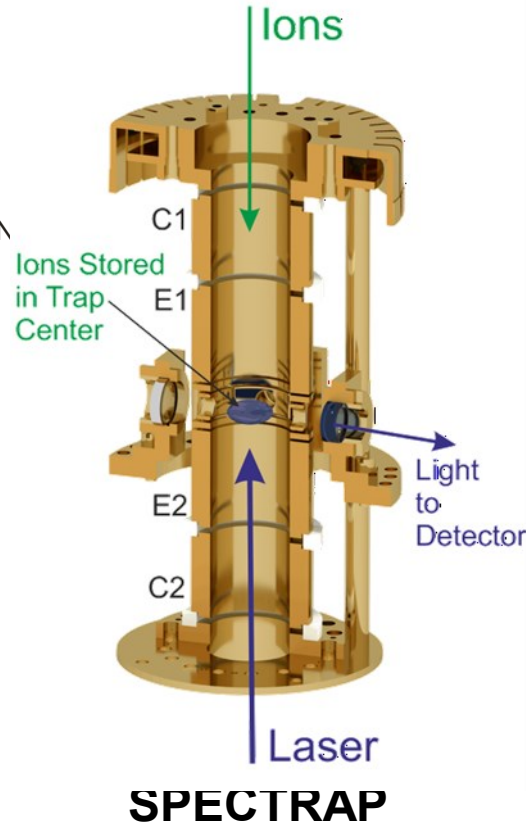
ARTEMIS

W. Quint, M. Vogel, G. Birkel, et al.

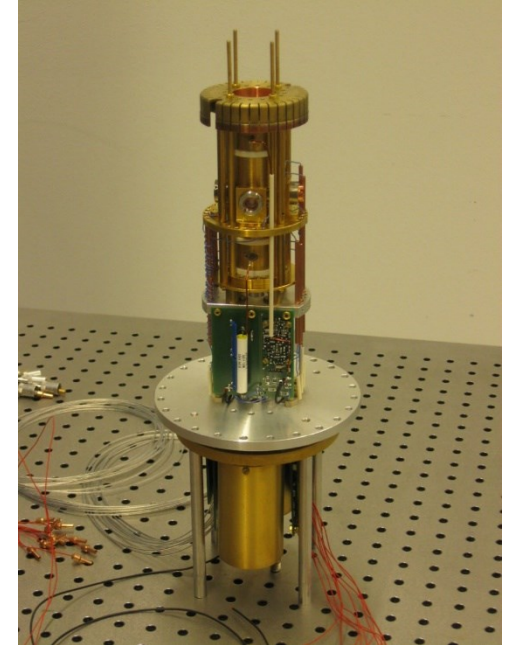
Technique:

Laser-microwave double resonance spectroscopy

- g-factor of bound electrons
- **nuclear g factor**
(free of diamagnetic correction)
- atomic lifetimes
- higher-order Zeeman shifts



→ Goal: Improved measurement of $E(1s)$, $E(2s)$ and $\Delta'E$
(gain: 2-3 orders of magnitude)

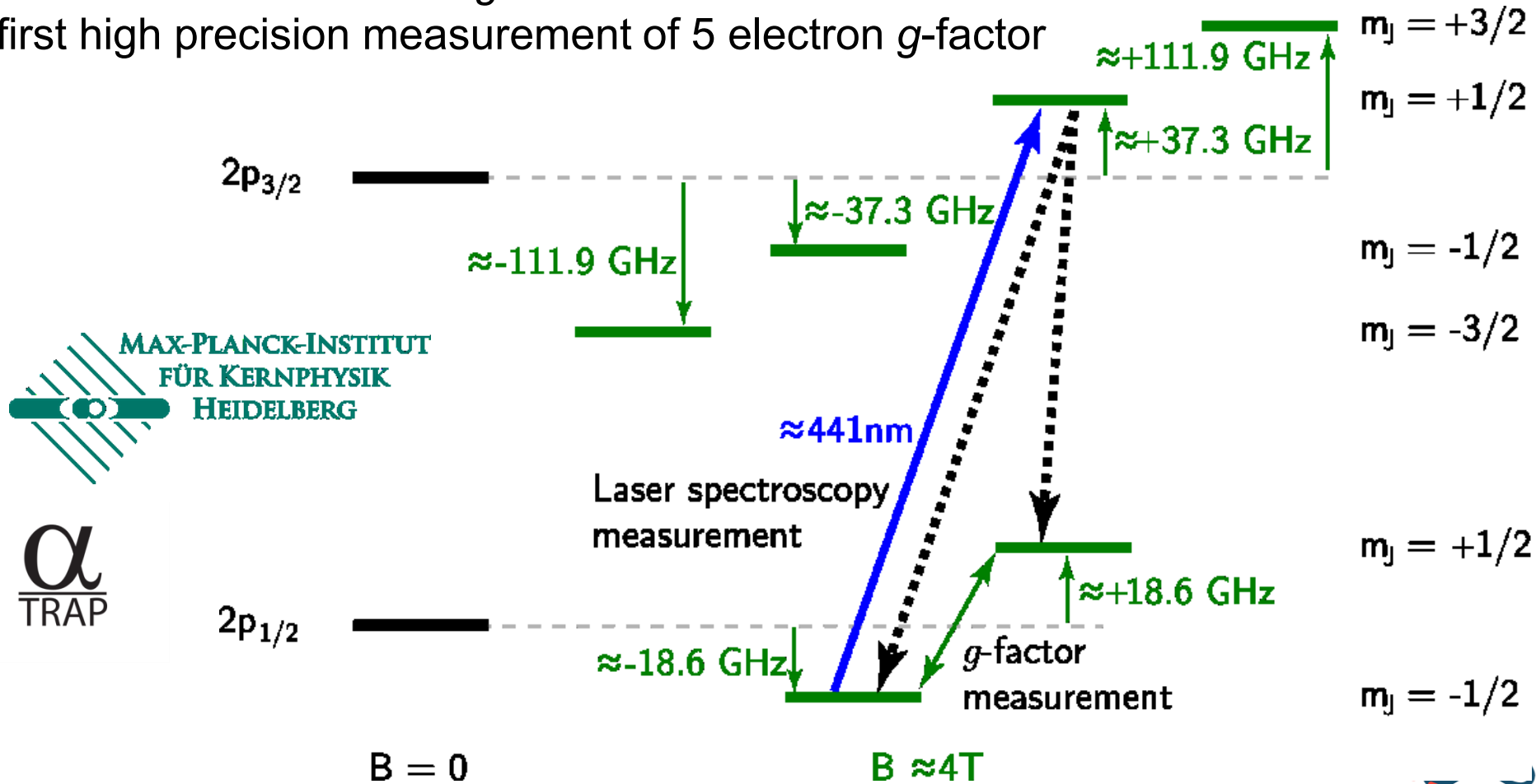


Z. Andelkovich, G. Birkel, M. Vogel,
W. Nörtershäuser et al.



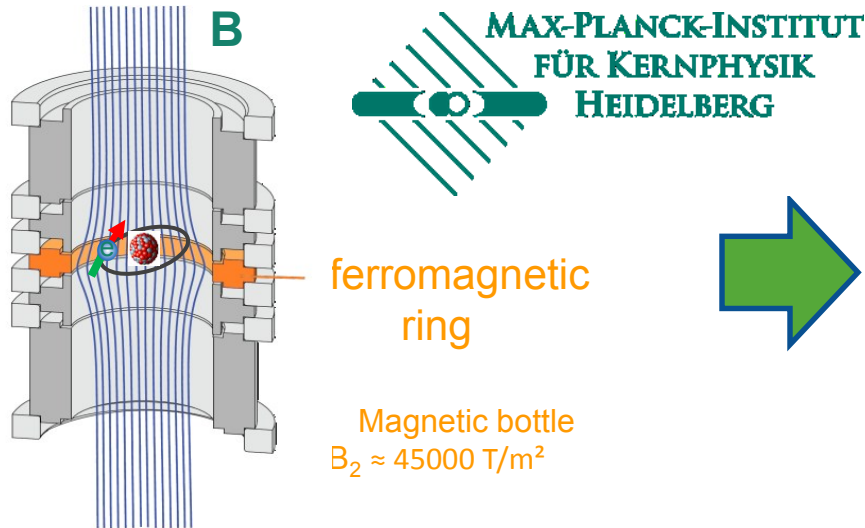
ALPHATRAP: Measurements on $^{40}\text{Ar}^{13+}$

$^{40}\text{Ar}^{13+}$ measurement of g -factor:
first high precision measurement of 5 electron g -factor



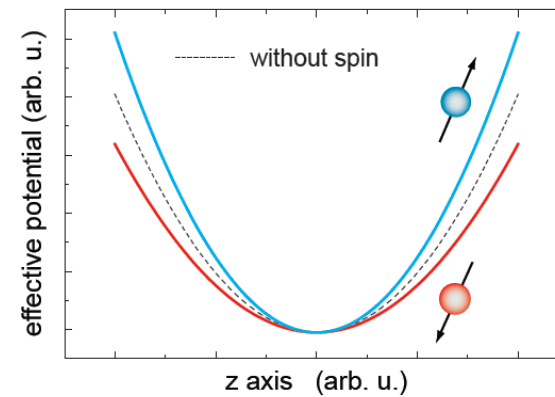
ALPHATRAP – Fluorescence-Free Detection Continuous Stern-Gerlach Effect in a Penning Trap

Introducing magnetic bottle inhomogeneity



additional potential:

$$\Phi_z^{\text{mag}} = \pm \mu_z \left[B_0 + B_2 \left(z^2 - \frac{\rho^2}{2} \right) \right]$$



α
TRAP

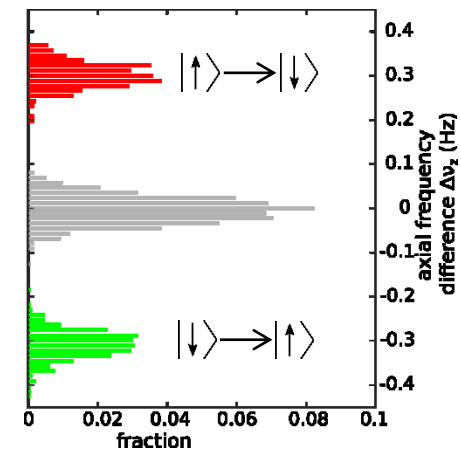
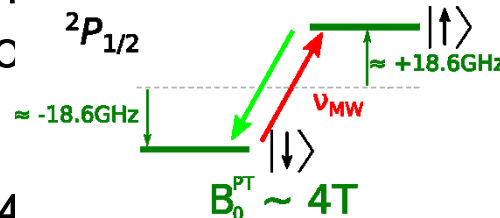
Courtesy of Alexander Egl

Continuous Stern-Gerlach effect:

axial frequency offset between
“up” and “down” spin orientatic

$$\Delta \nu_z \approx \frac{B_2 g \mu_B}{4\pi^2 m_{ion} \nu_z}$$

@ $\nu_z \approx 334 \text{ MHz}$

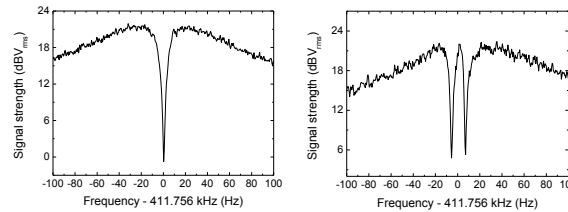


$^{12}\text{C}^{5+}$ $^{28}\text{Si}^{13+}$ $^{40}\text{Ar}^{13+}$ $^{208}\text{Pb}^{81+}$

$\Delta \nu_z$ 3.1 Hz 1.3 Hz 312 mHz 156 mHz

Measurement Cycle

PT: Measurement of motional frequencies and probe with spectroscopy frequency

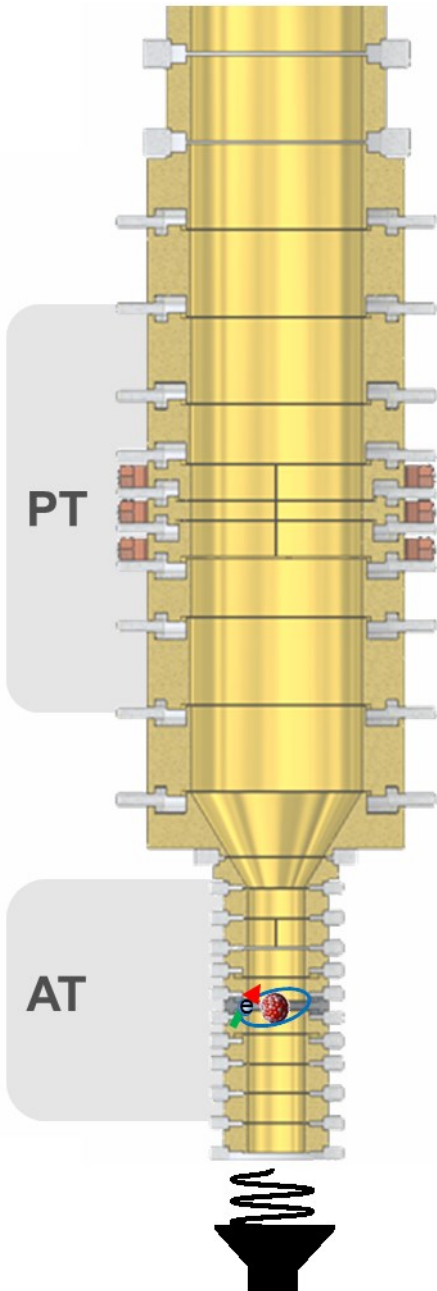
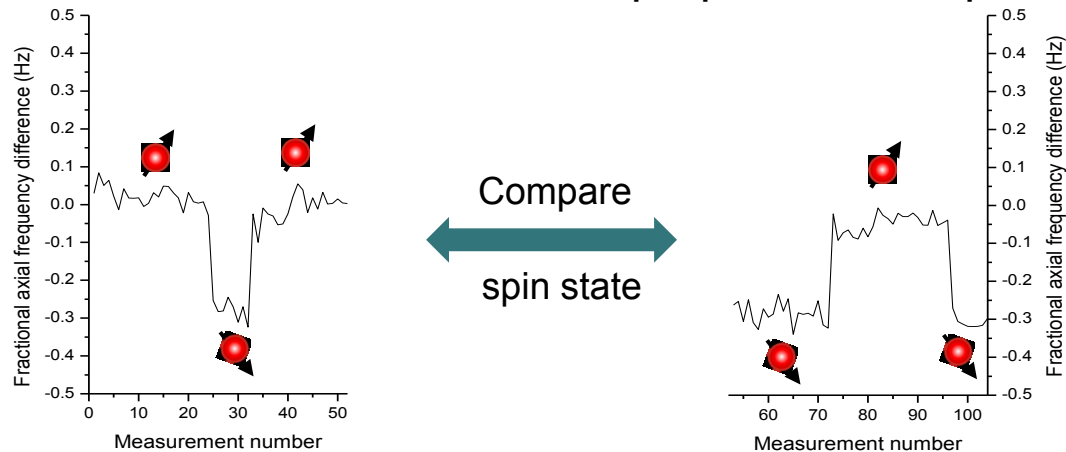


α
TRAP



$$v_Z^2 + v_+^2 + v_-^2 = v_C^2$$

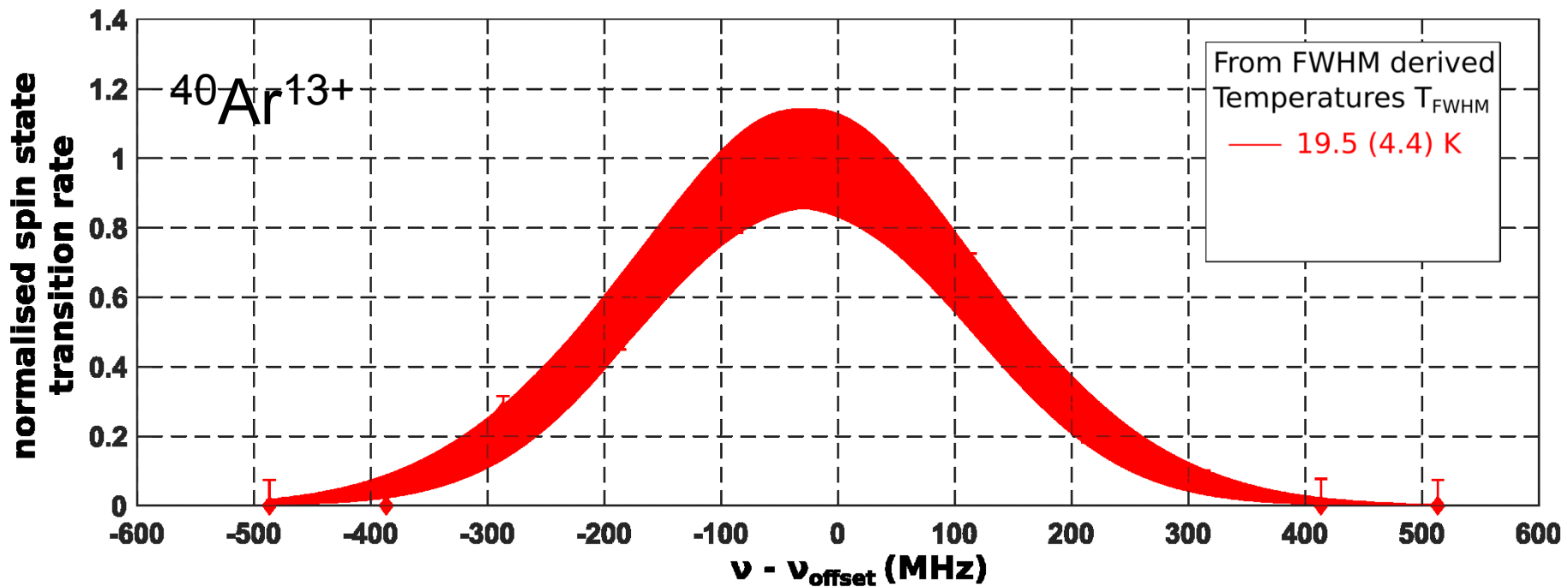
AT: Detection and preparation of spin orientation



Laser Spectroscopy Applying Continuous Stern-Gerlach Effect

Results – Resonance for $|1/2, -1/2\rangle \leftrightarrow |3/2, +1/2\rangle$ ($= |J, m_j\rangle$)
A. Egl, *et al.*, PRL **123**, 123001 (2019)

- Derived temperature from FWHM $\rightarrow 19.5 \pm 4.4$ K



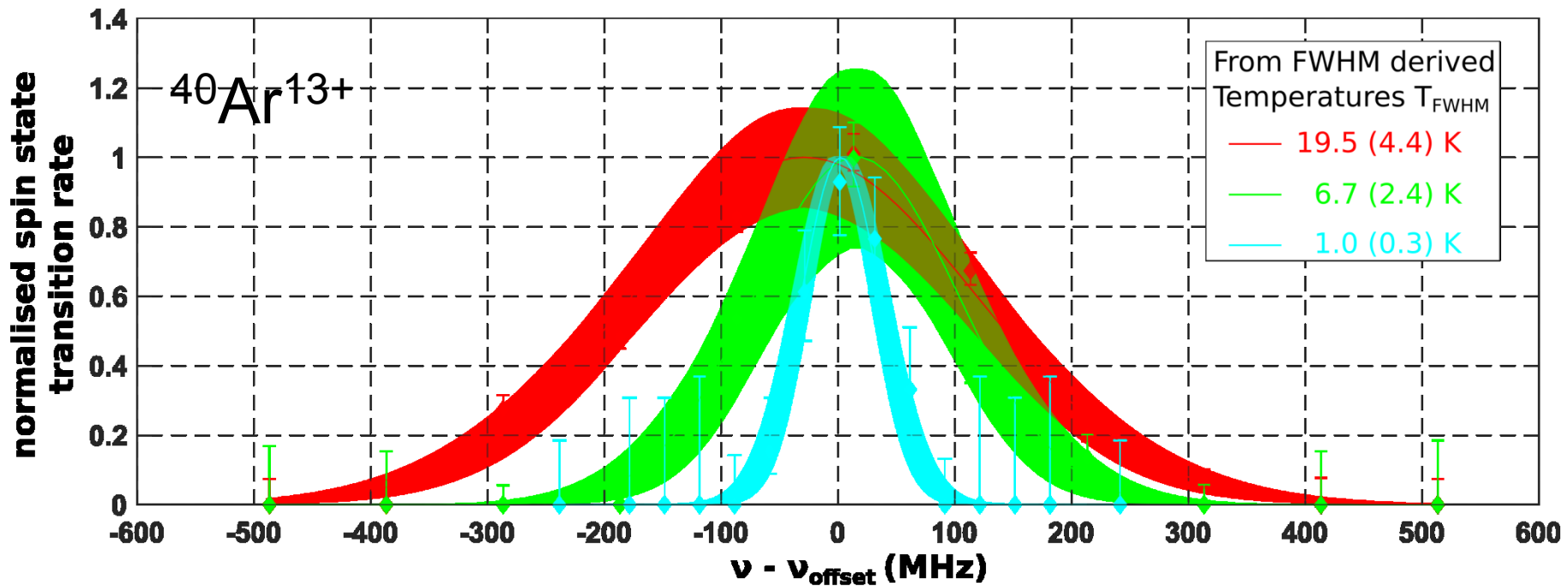
Laser Spectroscopy Applying Continuous Stern-Gerlach Effect

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α
TRAP

- Negative electronic feedback applied, expected lower temperature of focus
- Adiabatic cooling by lowering the trapping potential depth by a factor 3.8
→ lowered temperature by $\approx 3 \times 3.8 = 11.4$



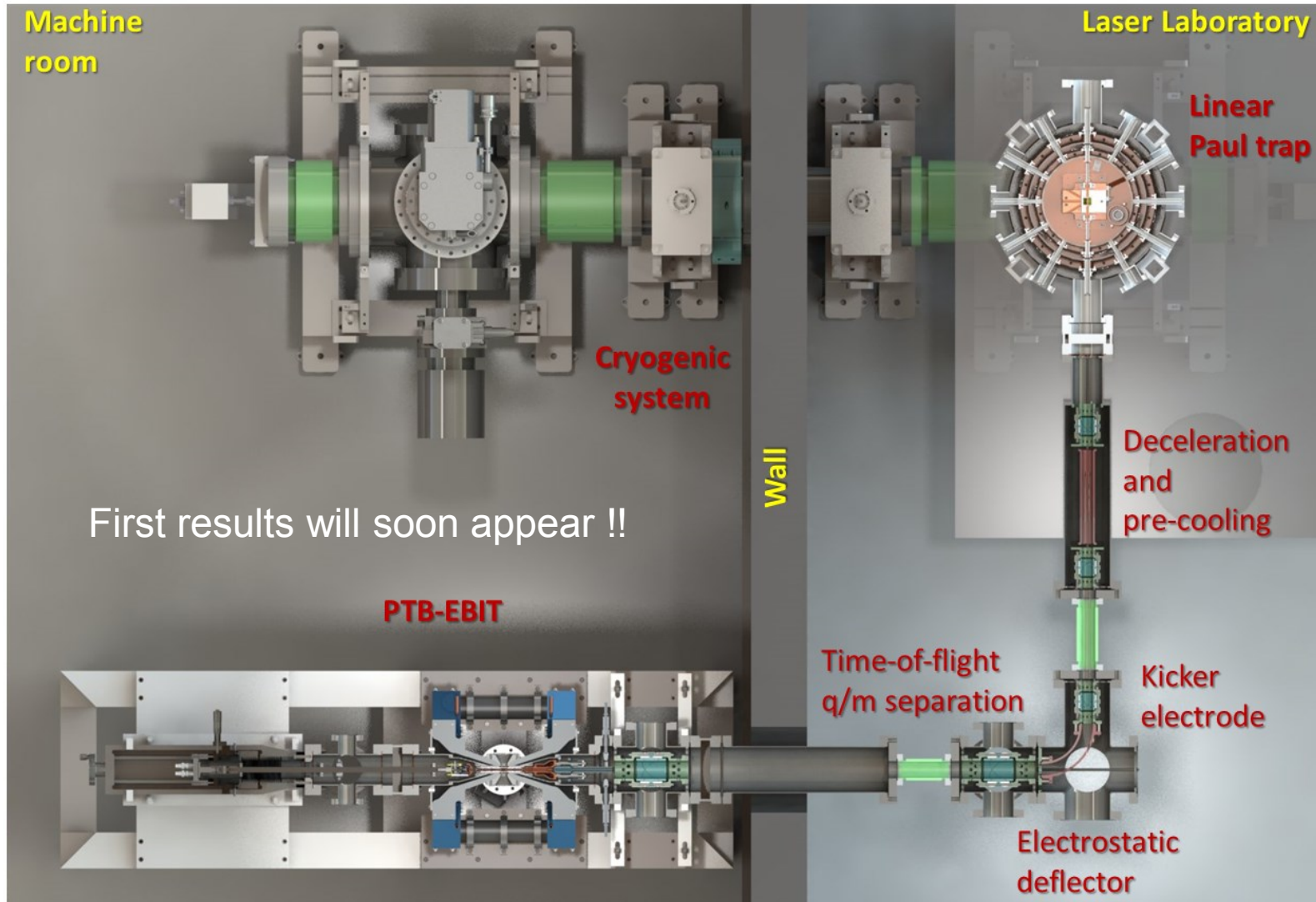
$$\nu = 679.216464 (4)_{\text{stat}}(5)_{\text{syst}} \text{ THz}, \Delta\nu/\nu = 9 \times 10^{-9}$$

Courtesy of Alexander Egl

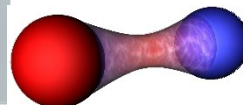
PTB-Experiment on cold HCl in a Paul-Trap



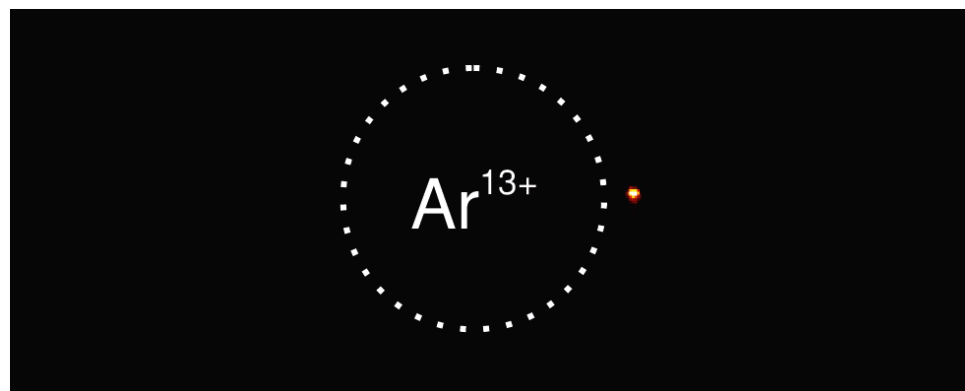
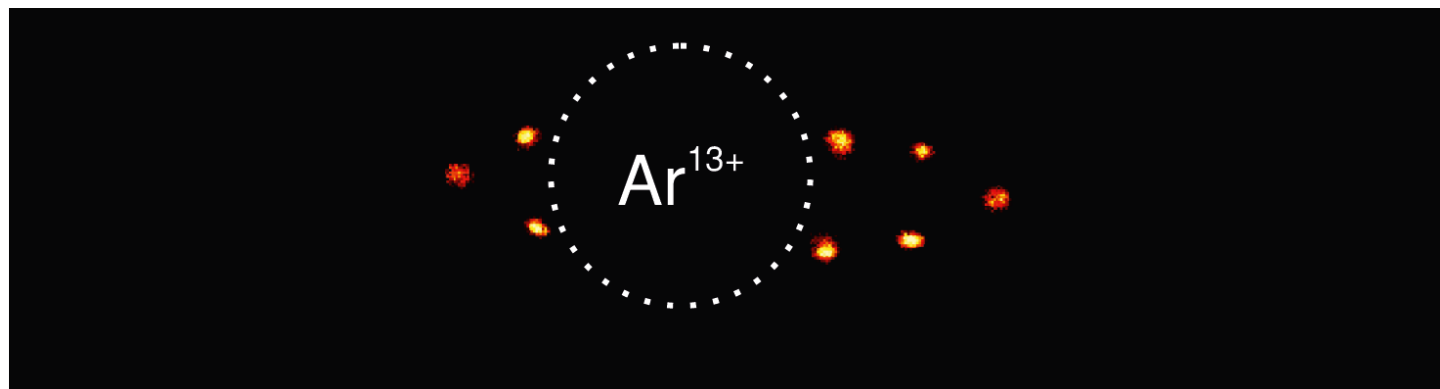
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Courtesy of Peter Micke



Two-Ion Crystal Preparation



Summary

- The measurement of $\Delta'E$ at the ESR strongly disagreed with QED calculations („HFS-Puzzle“).
- NMR Measurements and new chemical shift calculations in BiF_6^- were performed and proved that the Hyperfine Puzzle was caused by a wrong nuclear magnetic moment of ^{209}Bi .
- The measurements represents now a first (still inaccurate) test of bound-state QED in strong magnetic fields.

The road ahead of us:

- Determination of the nuclear magnetic moment of ^{208}Bi
- Prediction for the specific difference in ^{208}Bi
- Measurement of the hfs splitting in $^{208}\text{Bi}^{80+,82+}$ at the ESR storage ring
→ show independence of $\Delta'E$ from Bohr-Weisskopf effect
- TRAP measurements will provide higher accuracy (HITRAP → ARTEMIS & SPECTRAP, α TRAP, PTB-TRAP)



Partners and Funding



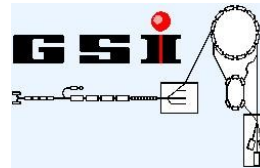
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