



THE AUSTRALIAN NATIONAL UNIVERSITY

# What we can learn from E0 transitions and how?

T. Kibèdi (ANU)



## Outline

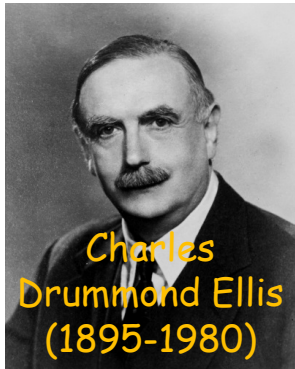
- Emission of E0 transitions
- New tabulation of  $\Omega_{CE}(E0)$
- e- $\gamma$  angular correlations of E0+M1+E2 transitions
- E0 transitions and the evolution of shape co-existence
  - ❑ W-Os-Pt-Hg (Z=82, N=126)
  - ❑ Fe (Z=N=28)
  - ❑  $^{12}\text{C}$  (Z=N=8)
- Monopole transitions in atomic nuclei - new review

# E0 transitions - 75 years on

*The Absolute Intensities and Internal Conversion Coefficients of the  $\gamma$ -Rays of Radium B and Radium C.*

By C. D. ELLIS, F.R.S., Fellow of Trinity College, Cambridge, and G. H. ASTON, M.A., B.Sc., Trinity Hall, Cambridge.

Proc. Roy. Soc. (London) **129** (1931) 180-207



Highly converted K, L1, M1 lines corresponding to 1.426 MeV transition

1.4155 MeV E0 in  $^{214}\text{Po}$   
 $T_{1/2} = 99(3)$  ps;  $\rho^2 = 0.0013(2)$

6.05 MeV E0 pairs in  $^{16}\text{O}$   
 Fowler & Lauritsen, Phys. Rev. 56 (1939) 840

6.05 MeV E0 ( $^{16}\text{O}$ ) E1+M1 double photon  
 Gorodetzky et al., Phys. Rev. Lett. 7 (1961) 170

CONSTITUTION OF  
 ATOMIC NUCLEI  
 AND  
 RADIOACTIVITY

BY  
 G. G A M O W



OXFORD  
 AT THE CLARENDON PRESS  
 1931

1931:  
 $0 \rightarrow 0$  transition:  
 quantum transition  
 forbidden

# Formation of E0 transitions

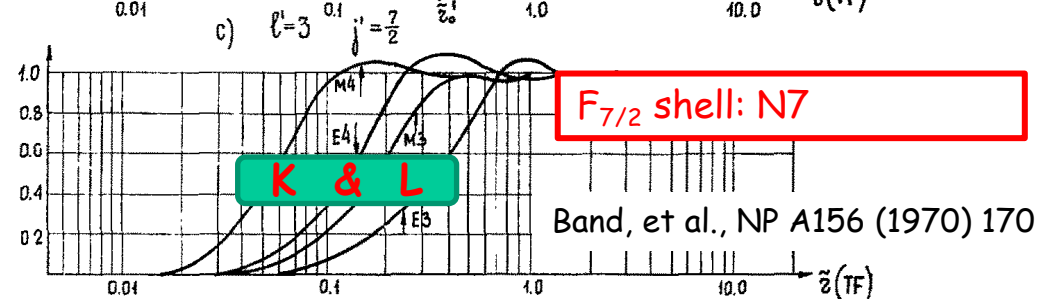
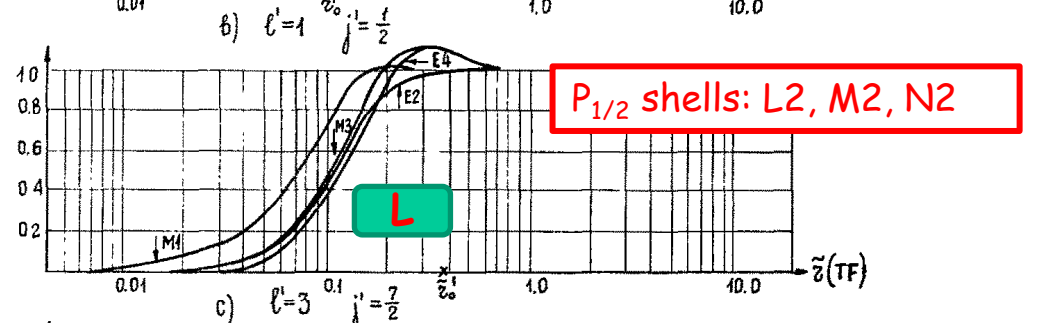
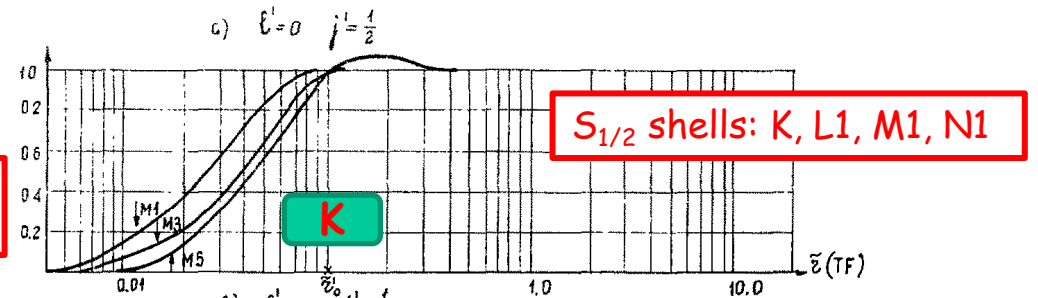
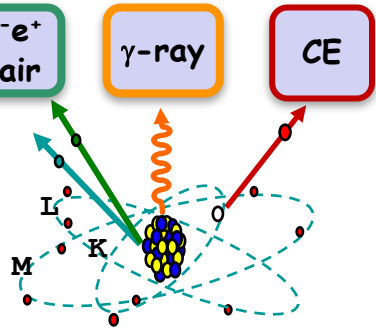
➤ "Normal conversion" ( $L > 0$ ):

- ❑ Small contribution from inside the nucleus
- ❑ Inner part of the atom (K-L-M shells)

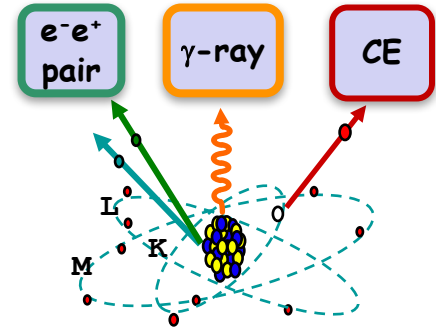
$$\alpha^{\lambda L}(r) / \alpha^{\lambda L}$$

➤ E0 conversion electrons:

- ❑ Monopole potential localised inside the nucleus
- ❑ Point nucleus approx.:  $W_{E0} \Rightarrow$  vanishes
- ❑ Purely penetration effect



Band, et al., NP A156 (1970) 170



*Transition probability*

$$W_T = W_\gamma + W_{CE} + W_\pi$$

*Conversion coefficient*

$$\alpha_{ce,\pi} = W_{ce,\pi} / W_\gamma$$

**CE &  $\pi$**

$$W_{ce,\pi} = W_\gamma \times \alpha_{ce,\pi}$$

**E0**

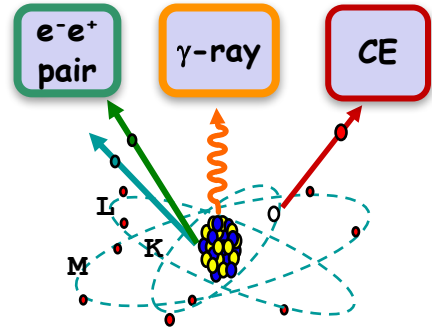
$$W_{ce,\pi} = \rho^2(0^+ \rightarrow 0^+) \times \Omega_{ce,\pi}(Z, \kappa)$$

**Monopole strength parameter**

$$\rho = \frac{\langle 0_f^+ | \sum e_j r_j^2 | 0_i^+ \rangle}{eR^2} = \frac{\langle 0_f^+ | m(E0) | 0_i^+ \rangle}{eR^2}$$

**Reduced transition probability**

$$B(E0) = e^2 R^4 \rho^2$$



*Transition probability*

$$W_T = W_\gamma + W_{CE} + W_\pi$$

*Conversion Coefficient*

$$\alpha_{ce,\pi} = W_{ce,\pi} / W_\gamma$$

**CE &  $\pi$**

$$W_{ce,\pi} = W_\gamma \times \alpha_{ce,\pi}$$

**E0**

$$W_{ce,\pi} = \rho^2(0^+ \rightarrow 0^+) \times \Omega_{ce,\pi}(Z, \kappa)$$

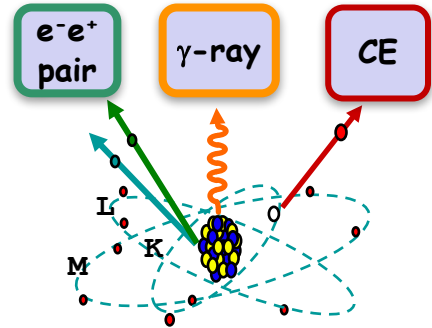
**Separation is less complete!**

**nuclear**

**atomic**

Piet Van Isacker (GANIL) talk on Monday

$$\rho = \sum_{k \in \text{protons}}^Z \left\langle f \left( \left( \frac{r_k}{R} \right)^2 - \sigma \left( \frac{r_k}{R} \right)^4 + \dots \right) \middle| i \right\rangle \quad \left( R = r_0 A^{1/3}, r_0 = 1.2 \text{ fm} \right)$$



*Transition probability*

$$W_T = W_\gamma + W_{CE} + W_\pi$$

*Conversion Coefficient*

$$\alpha_{ce,\pi} = W_{ce,\pi} / W_\gamma$$

**CE &  $\pi$**

$$W_{ce,\pi} = W_\gamma \times \alpha_{ce,\pi}$$

**E0**

$$W_{ce,\pi} = \rho^2(0^+ \rightarrow 0^+) \times \Omega_{ce,\pi}(Z, \kappa)$$

**nuclear**                      **atomic**

$L \neq 0$ :  $W_\gamma$  and  $W_{ce,\pi}$   
 $L = 0$  i.e. E0: only  $W_{ce,\pi}$ !  
**Need to measure electrons and/or electron-positron pairs**

- $\Omega_{\kappa,\pi}(E0) \sim I_{\kappa,\pi}(E0)$
- Only ratios of  $\Omega_{\kappa,\pi}(E0)$  could be measured
  - $K/\pi$
  - $K/L, K/LM, K/MN$
  - $L/M$

$S_{1/2}, P_{1/2}$  and  $P_{3/2}$  shells only

New  $\Omega_{CE}$  calculations

**Takahe** (Jackson Dowie, ANU)

- Modified version of the CATAR code (Puli & Raff, 1975)
- Relativistic Hartree-Fock-Slater atomic model
- Sliv's surface current model
- Directional & polarization particle parameters
- Penetration parameters
- $\Omega_{CE}$  electronic factors
- $Z$  up to 126

$\Omega_{CE}(E0), Mo (Z=42) 1000 \text{ keV}$			
	Takahe (2017)	Bell et al. (1972)	Hager & Seltzer (1968)
K	1.454E+9	1.459E+9	1.438E+9
L1	1.629E+8	1.574E+8	1.611E+8
L2	1.358E+6	1.152E+6	1.346E+6
M1	2.913E+7		
M2	2.566E+5		
N1	4.983E+6		
N2	3.788E+4		
O1	3.807E+5		



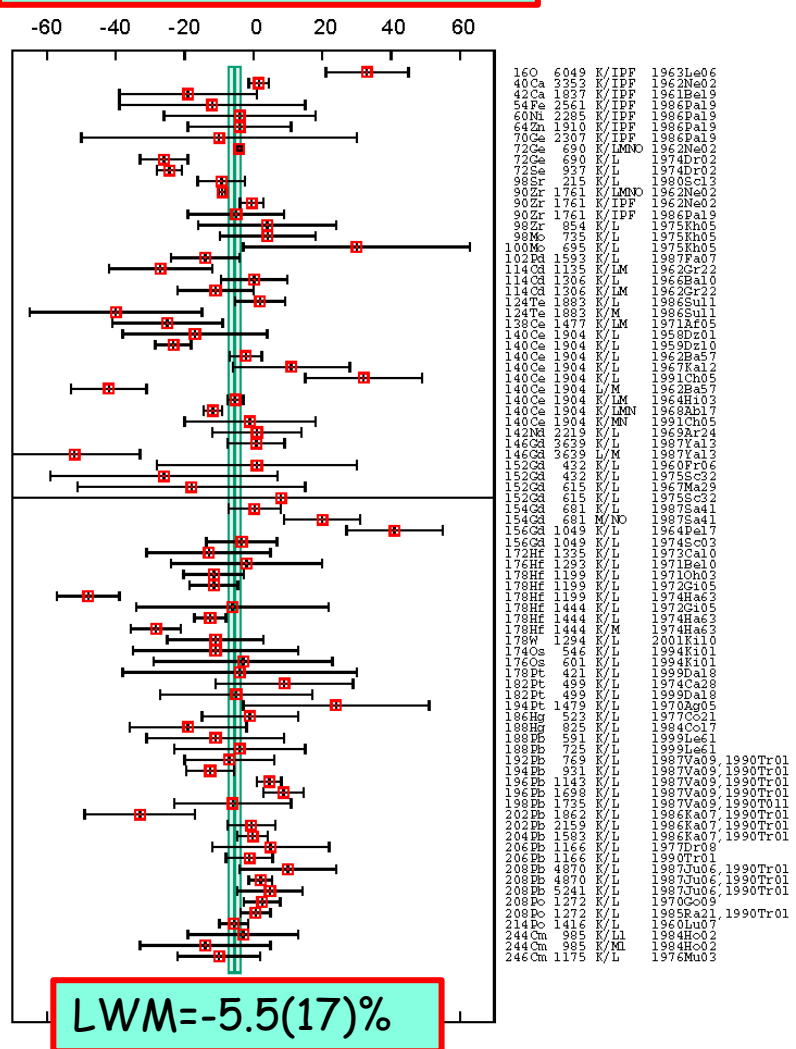
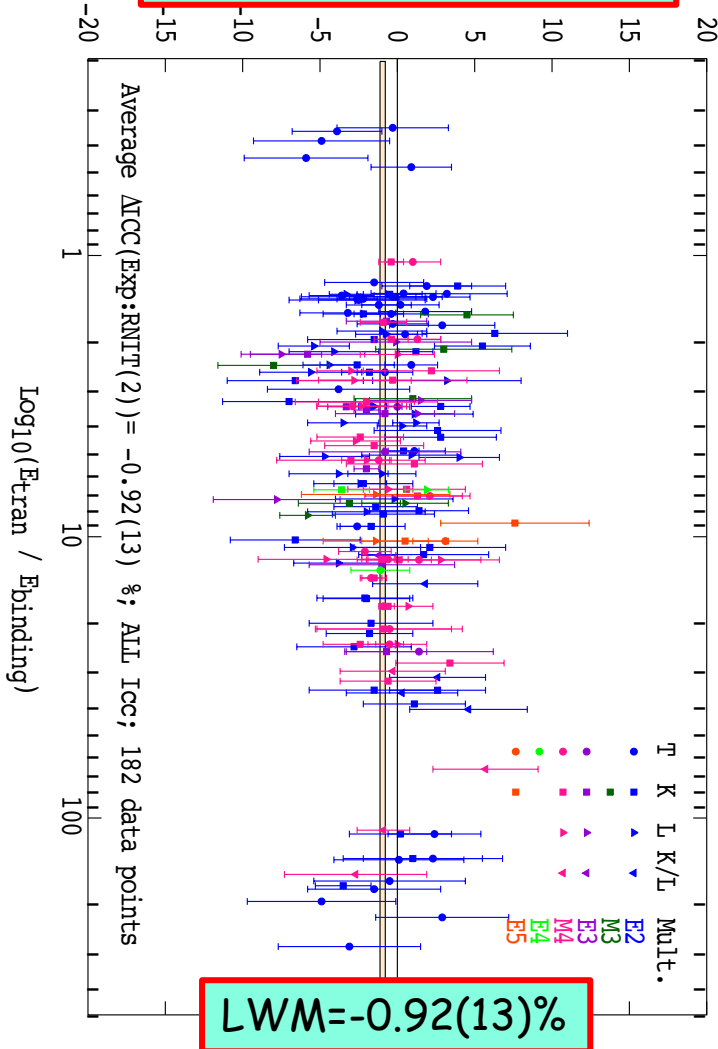
# $\Omega(E0)$ - how good are they?

BrICC

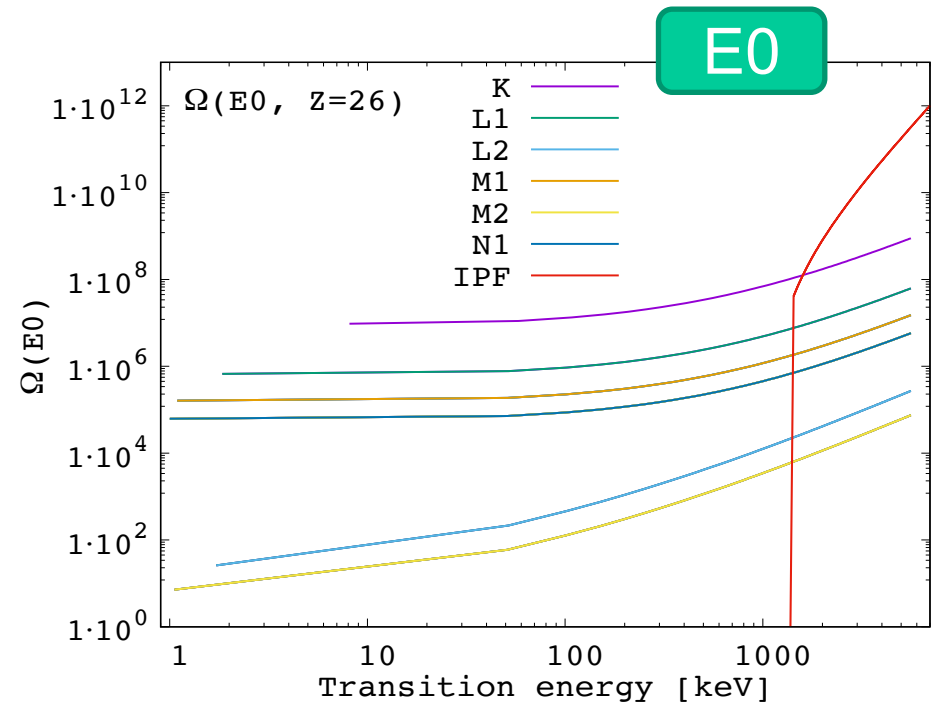
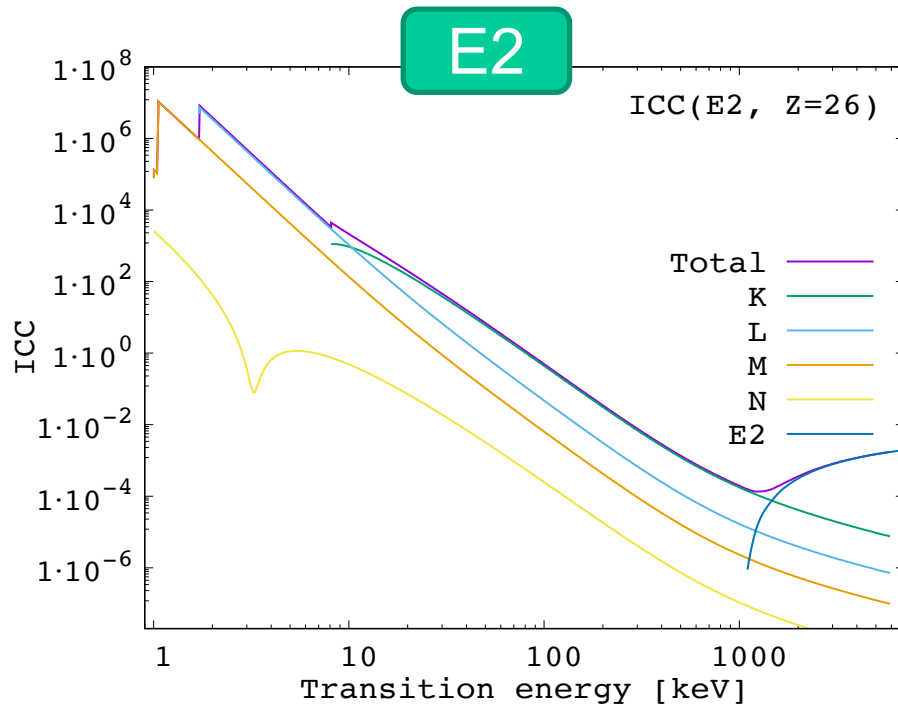
Experiment Vs. Theory  
precision ICC

Experiment Vs. Theory  
Ratios of  $\Omega_{CE}$  &  $\Omega_{\pi}$

CATAR



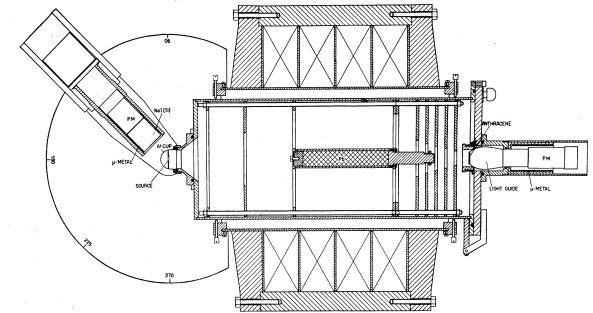
# E0 vs E2 transitions



- Z and atomic shell dependent
- Increases with energy
- K/L weak dependence on Z (4 to 9)
- Pair conversion dominant at low Z

Church, Rose and Weneser (1958)

- E0 can proceed in competition of E2 & M1
- e- $\gamma$  angular correlations: a sensitive test
- First observation: 334.0-356.5 cascade in  $^{196}\text{Pt}$  (Gerholm & Pettersson 1958)



Mixing ratios:

$$\delta^2 = \frac{N_{\gamma}^{E2}}{N_{\gamma}^{M1}}$$

$$p^2 = \frac{\alpha_K^{M1}}{\alpha_K^{E2}} \delta^2$$

$$q^2 = \frac{N_K^{E0}}{N_K^{E2}}$$

$$W(\gamma\gamma, M1 + E2) = P_0 + \frac{1}{1+\delta^2} [A_2^e + 2\delta A_2 + \delta^2 A_2^m] P_2 + \frac{1}{1+\delta^2} [A_4^e] P_4$$

$$W(e\gamma, M1 + E2) = P_0 + \frac{1}{1+p^2} [b_2^e A_2^e + 2p b_2 A_2 + p^2 b_2^m A_2^m] P_2 + \frac{1}{1+p^2} [b_4^e A_4^e] P_4$$

$$W(e\gamma, E0 + M1 + E2) = \frac{1+p^2}{1+p^2+q^2} W(e\gamma, M1 + E2) + \frac{q^2}{1+p^2+q^2} P_0 + \frac{q}{1+p^2+q^2} b_0 P_2$$

Multi-detector electron- $\gamma$  arrays  
Need to evaluate numerically the e-e/e- $\gamma$  correlation!

## Mixed E0+M1+E2 conversion coefficients

$$\alpha_K^{exp}(E0 + M1 + E2) = \frac{1}{1 + \delta^2} [\delta^2 [1 + q^2] \alpha_K(E2) + \alpha_K(M1)]$$

Mixing ratios:  $\delta^2 = \frac{N_\gamma^{E2}}{N_\gamma^{M1}}$        $q^2 = \frac{N_K^{E0}}{N_K^{E2}}$

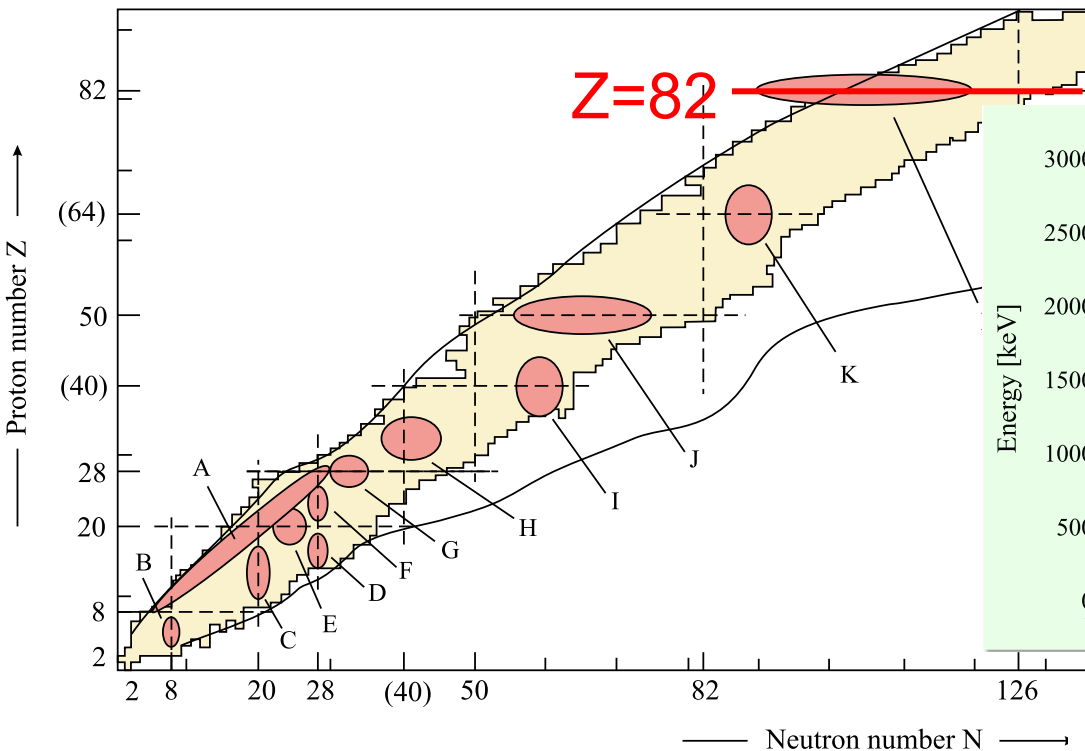
## Reduced E0 matrix element

$$\rho^2(E0) = \frac{1}{\tau(E0) \times \Omega(E0)} = q^2 \frac{\alpha_K(E2) W_\gamma(E2)}{\Omega(E0)}$$

## B(E0)/B(E2) ratio (Rasmussen 1960)

$$X(E0/E2) = \frac{\rho^2(E0) e^2 R_0^4}{B(E2)} = \frac{2.54 \times 10^9 \times A^{4/3} E_\gamma^5 q^2 \alpha_K(E2)}{\Omega_K(E0)}$$

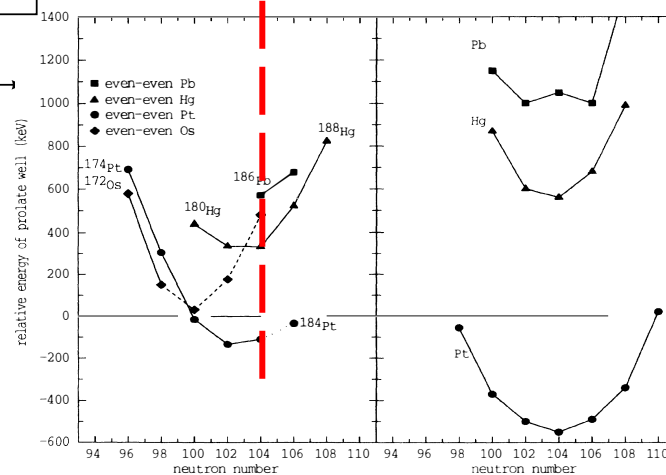
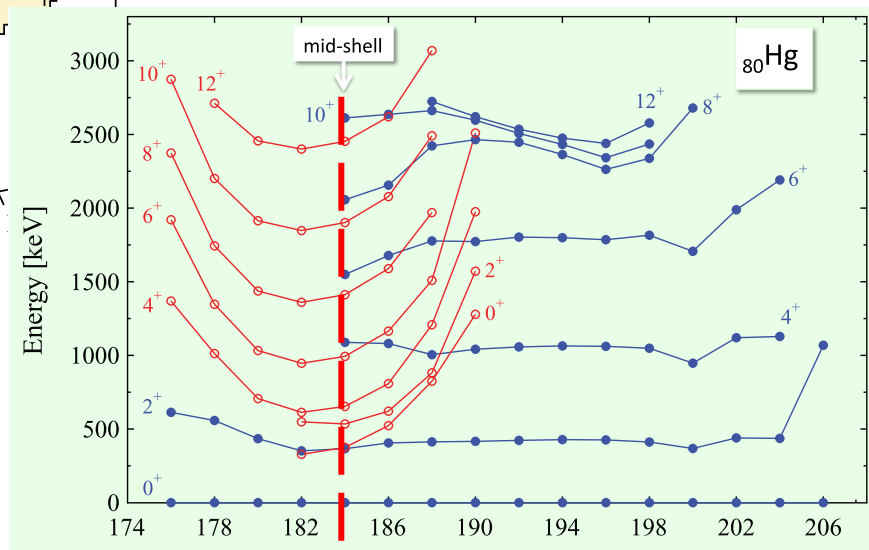
# Evolution of shape co-existence



Heyde & Wood, Rev. Mod. Phys. 83 (2011) 1467

Textbook example around Pb(Z=82):  
Hg(Z=80), Po(Z=84)

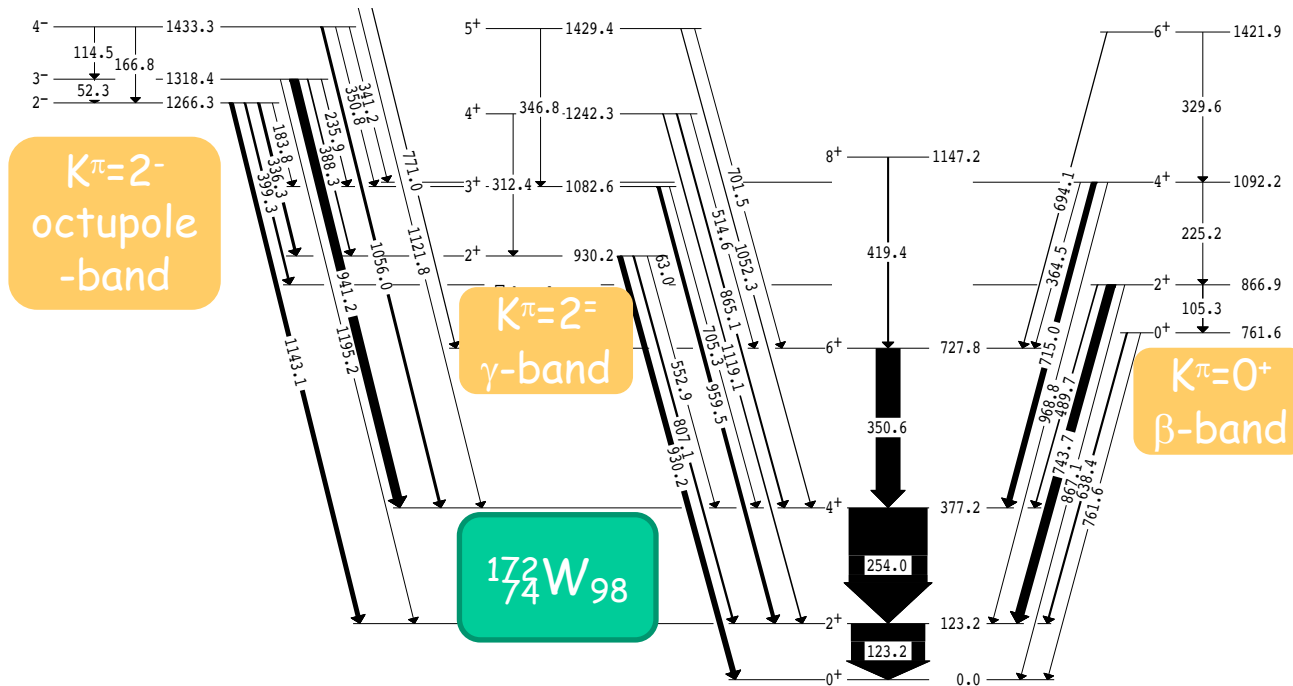
Heyde & Wood, Phys. Scr. 91 (2016) 083008



Dracoulis PRC 49 (1994) 3324

E0 Workshop, CEA 2017

# Evolution of shape co-existence W-Os-Pt



From radioactive decay:  $\gamma$ ,  $\gamma$ -g, CE

- $^{174}\text{Pt}$ : Dracoulis, et al., PRC 44, R1246 (1991)
- $^{176}\text{Pt}$ : Dracoulis, et al., J. Phys. G 12, L97 (1986)
- $^{172}\text{Os}$ : Davidson, et al., Nucl. Phys. A568, 90 (1994)
- $^{174-182}\text{Os}$ : Kibedi, et al., Nucl. Phys. A567, 183 (1994)
- $^{172-178}\text{W}$ : Kibedi, et al., Nucl. Phys. A 688 669 (2001)

172-178W: Kibedi, et al.,

Nucl. Phys. A 688 669 (2001)

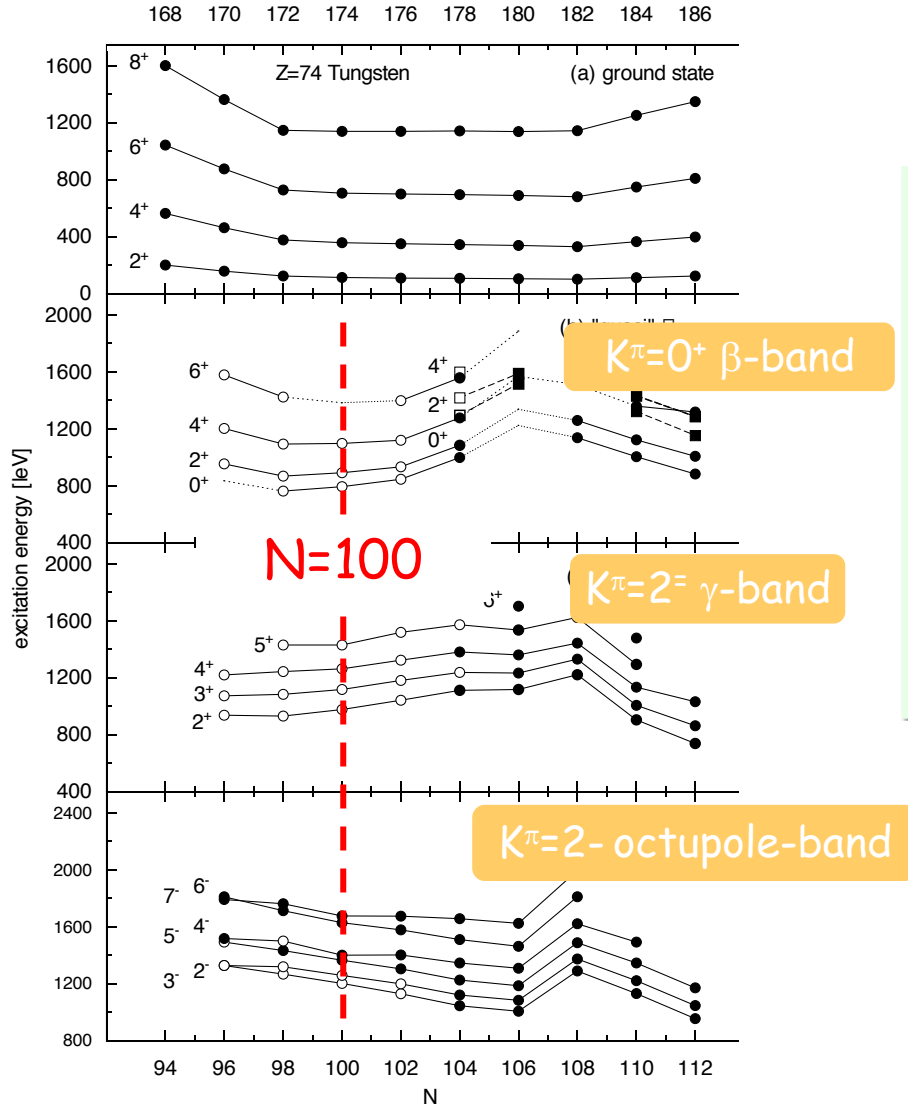
- 0<sup>+</sup> - 0<sup>+</sup> E0 transitions (4)
- J<sup>+</sup> - J<sup>+</sup> (J=2,4,6) E0+M1+E2 transitions (19)
- No T1/2 - only X=B(E0)/B(E2)

Table 6

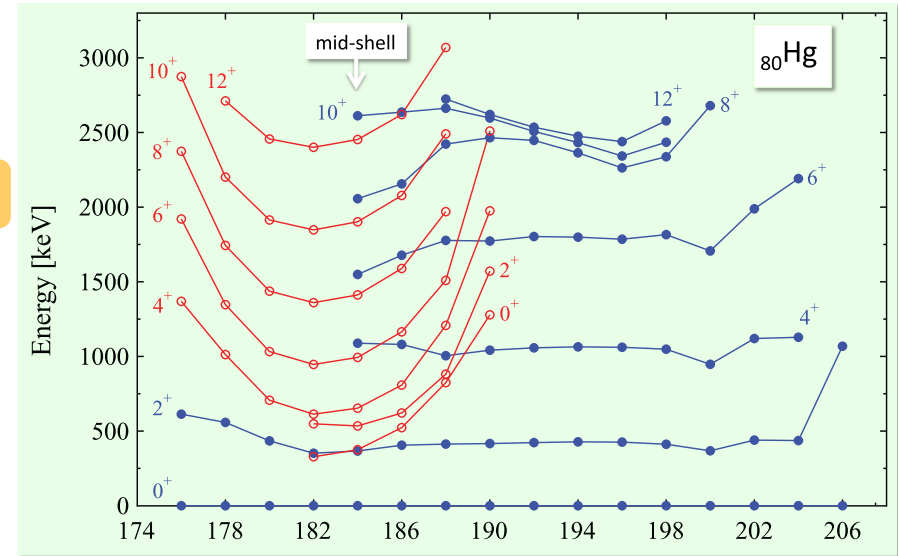
E0 component of  $J_i^+ \rightarrow J_f^+$  transitions in  $A = 170$  to  $178$  tungsten isotopes (only for  $K$ -conversion electron lines have been considered)

$J_i^+ \rightarrow J_f^+$	A	$E_\gamma$ [keV]	$\alpha_K$ (exp) × 100	$\alpha_K$ (E2) × 100	$\alpha_K$ (M1) × 100	$\delta$ (E2/M1)	$q^2$ (E0/E2)	X(E0/E2)
$0_\beta^+ \rightarrow 0_g^+$	172	761.6	E0				3.6(11)	0.060(18)
	174	792.2	E0				2.1(4)	0.041(2)
	176	844.0	E0				2.5(6)	0.058(15)
$0_3^+ \rightarrow 0_g^+$	178	1294.4	E0				26(2)	1.73(12)
$2_\beta^+ \rightarrow 2_g^+$	172	743.7	4.52(20)	0.677	1.78	$-10.3^{+3.0}_{-7.0}$	5.7(3)	0.153(8)
	174	777.0	8.4(9)	0.618	16.0	$-4.5^{+0.9}_{-1.3}$	13.1(15)	0.39(5)
	176	822.2	5.6(4)	0.551	1.38	$-2.7^{+0.4}_{-0.5}$	$10.2^{+1.3}_{-1.1}$	$0.346^{+0.042}_{-0.038}$
	178	976.5	0.76(5)	0.392	0.902	$-12.3^{+2.8}_{-6.4}$	$0.94^{+0.13}_{-0.13}$	$0.045^{+0.006}_{-0.006}$
$2_\gamma^+ \rightarrow 2_g^+$	172	807.1	1.14(15)	0.572	1.45	$+7.6^{+2.5}_{-7.4}$	7.4(19)	0.17(4)
	176	932.4	0.83(16)	0.429	1.01	$+3.0^{+1.0}_{-0.7}$	$0.89^{+0.42}_{-0.47}$	$0.039^{+0.018}_{-0.018}$
	178	1004.6	0.90(8)	0.370	0.840	> +2	$1.5^{+0.3}_{-0.3}$	$0.076^{+0.014}_{-0.013}$
$2_4^+ \rightarrow 2_g^+$	178	1311.5	1.49(13)	0.223	0.436	> 2  <sup>a</sup>	$6.9^{+0.7}_{-1.7}$	$0.61^{+0.07}_{-0.15}$
$4_\beta^+ \rightarrow 4_g^+$	170	739.8	4.6(11)	0.684	1.81	$-3.3^{+1.6}$	6.1(18)	0.16(5)
	172	715.0	7.0(7)	0.735	1.97	$-4.1^{+3.6}_{-1.9}$	8.9(11)	0.221(26)
	174	739.4	5.8(9)	0.685	1.81	$-4.2^{+0.7}_{-1.1}$	7.8(14)	0.21(4)
	176	768.7	6.6(7)	0.632	1.64	$-2.2^{+0.6}_{-1.2}$	$11.1^{+3.0}_{-2.2}$	$0.33^{+0.09}_{-0.06}$
	178	932.4	1.76(11)	0.429	1.01	$-6.6^{+1.5}_{-3.0}$	$3.15^{+0.30}_{-0.28}$	$0.140^{+0.013}_{-0.013}$
$4_\gamma^+ \rightarrow 4_g^+$	172	865.1	1.4(4)	0.497	1.22	$+4.2^{+1.1}_{-2.6}$	1.8(9)	0.007(3)
	178	1037.4	0.51(5)	0.497	1.22	$-1.9^{+0.7}_{-1.2}$	0.7(5)	0.014(18)
$4_4^+ \rightarrow 4_g^+$	178	1255.1	1.01(8)	0.242	0.485	> 2  <sup>a</sup>	$3.7^{+0.4}_{-0.9}$	$0.30^{+0.03}_{-0.07}$
$6_\beta^+ \rightarrow 6_g^+$	170	702.8	6.7(15)	0.762	2.06	$-1.7^{+0.8}_{-2.5}$	10(3)	0.23(7)
	172	694.1	6.4(14)	0.782	2.12	$-5.0^{+3.2}$	10(3)	0.23(7)
	176	696.6	4.9(9)	0.776	2.10	> 2  <sup>a</sup>	$6.2^{+1.5}_{-2.1}$	$0.15^{+0.04}_{-0.05}$

# Evolution of shape co-existence W-Os-Pt

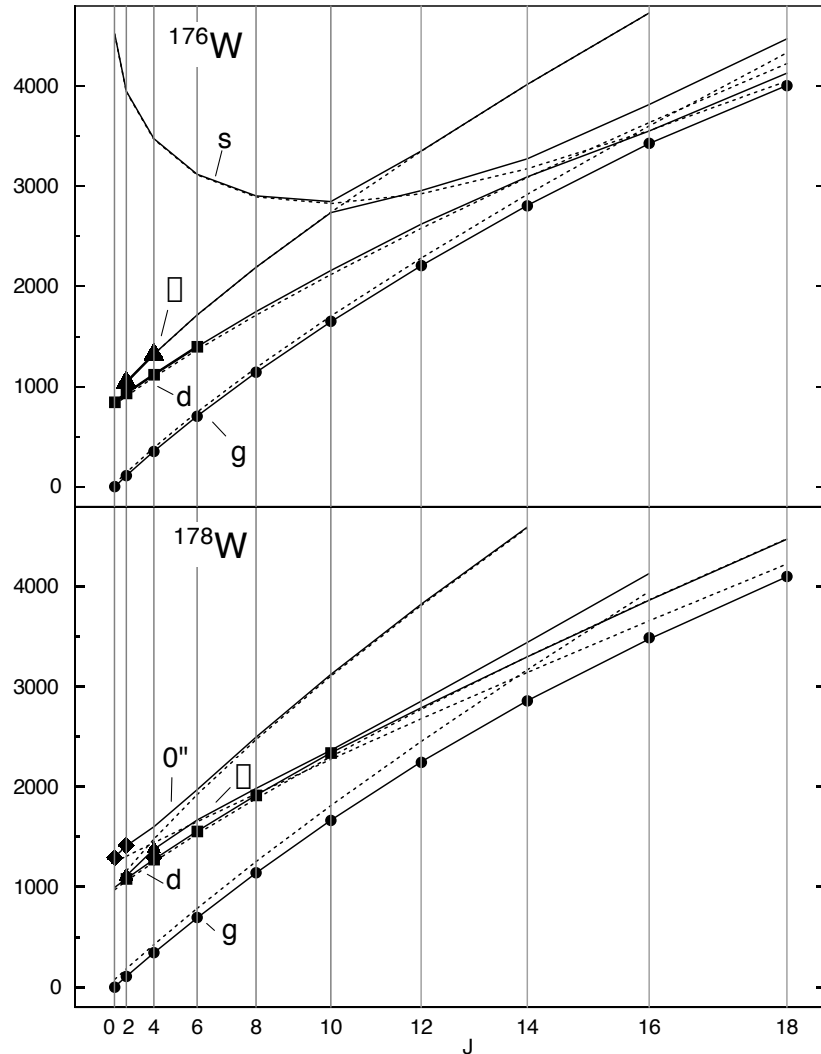


## Level systematics



Kibedi, et al., Nucl. Phys. A 688 669 (2001)





Kibedi, et al., Nucl. Phys. A 688 669 (2001)

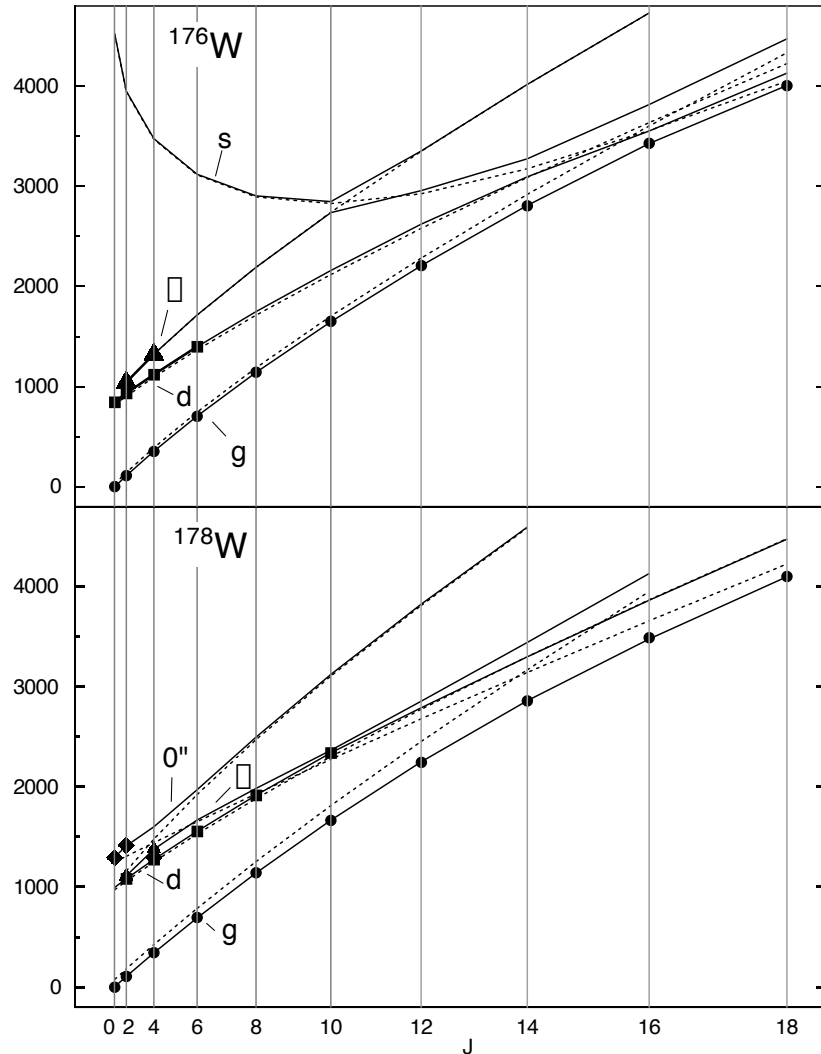
Tibor Kibedi, Dep. of Nuclear Physics, Australian National University

### 4-band mixing calculations

- K=0 g.s. unperturbed ground-state rotation
- K=0 "deformed" band
- K=2 g-band
- K=0 "s"-band, unperturbed rotation-aligned band, back banding observed in (HI,xn)

### ➤ Parameters to fit excitation energies:

- Moment of inertia (VMI)
- Unperturbed band-head energies
- Spin-independent interactions
- aligned angular momentum of the s-band



### 4-band mixing calculations

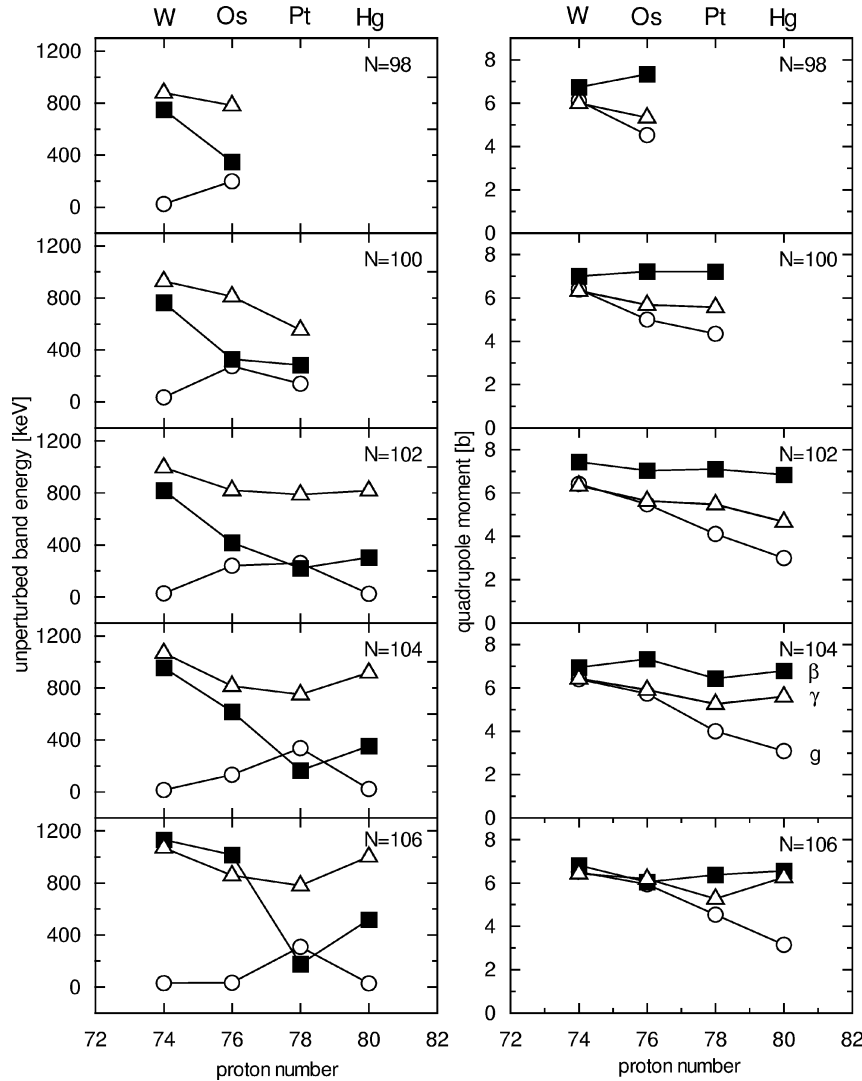
- K=0 g.s. unperturbed ground-state rotation
- K=0 "deformed" band
- K=2 g-band
- K=0 "s"-band, unperturbed rotation-aligned band, back banding observed in (HI,xn)

### ➤ Interactions:

- 150 keV (g-d)
- 30 keV (g- $\gamma$ )
- 5 keV (d- $\gamma$ )

Kibedi, et al., Nucl. Phys. A 688 669 (2001)

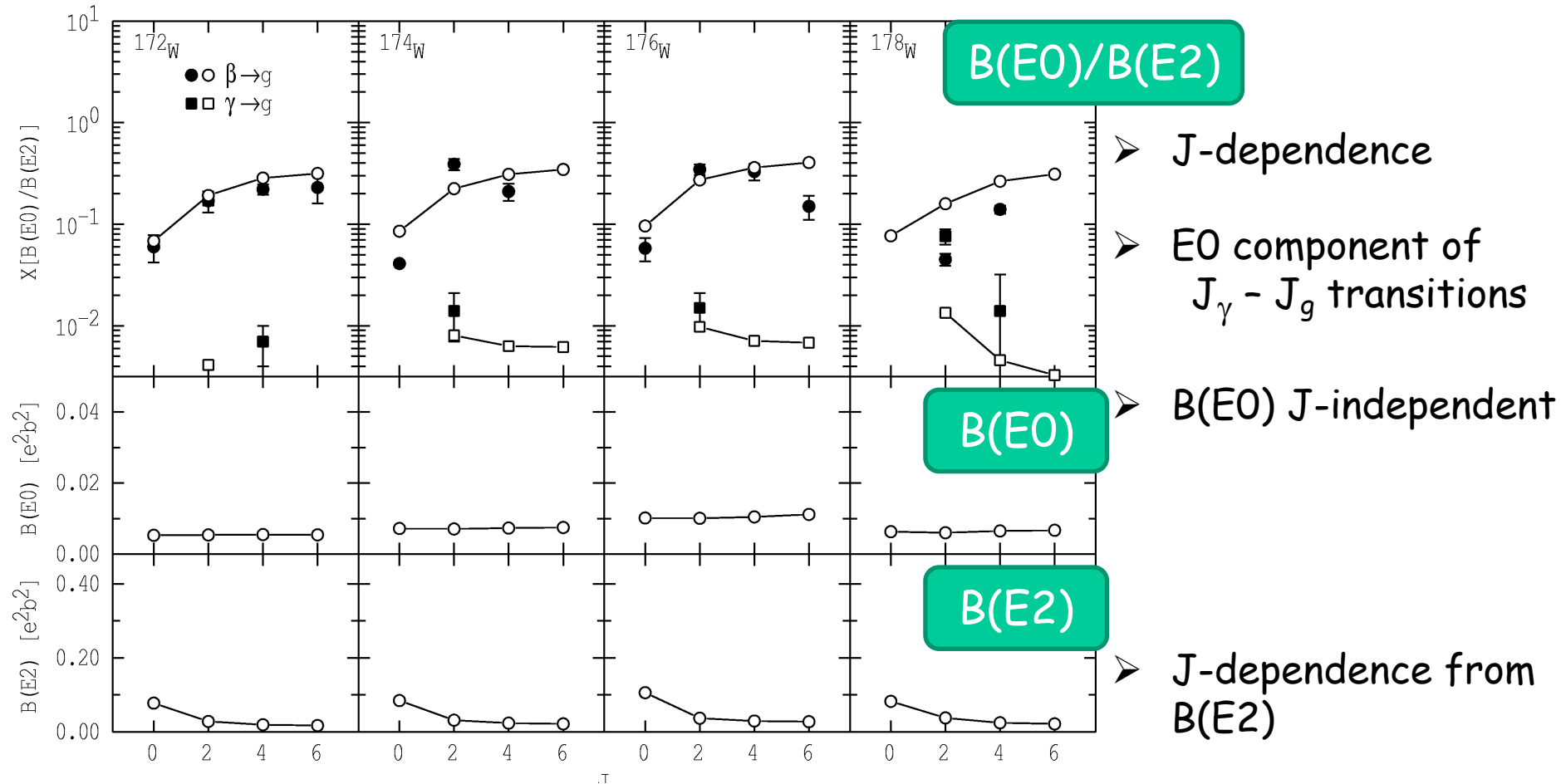
# Evolution of shape co-existence W-Os-Pt



- Smooth evolution across  $Z=80$  to  $74$  and  $N=98$  to  $106$
- Differences in deformation
- Unperturbed  $\gamma$  and  $d$  bands shifted down in energy as  $Z \gg 82$

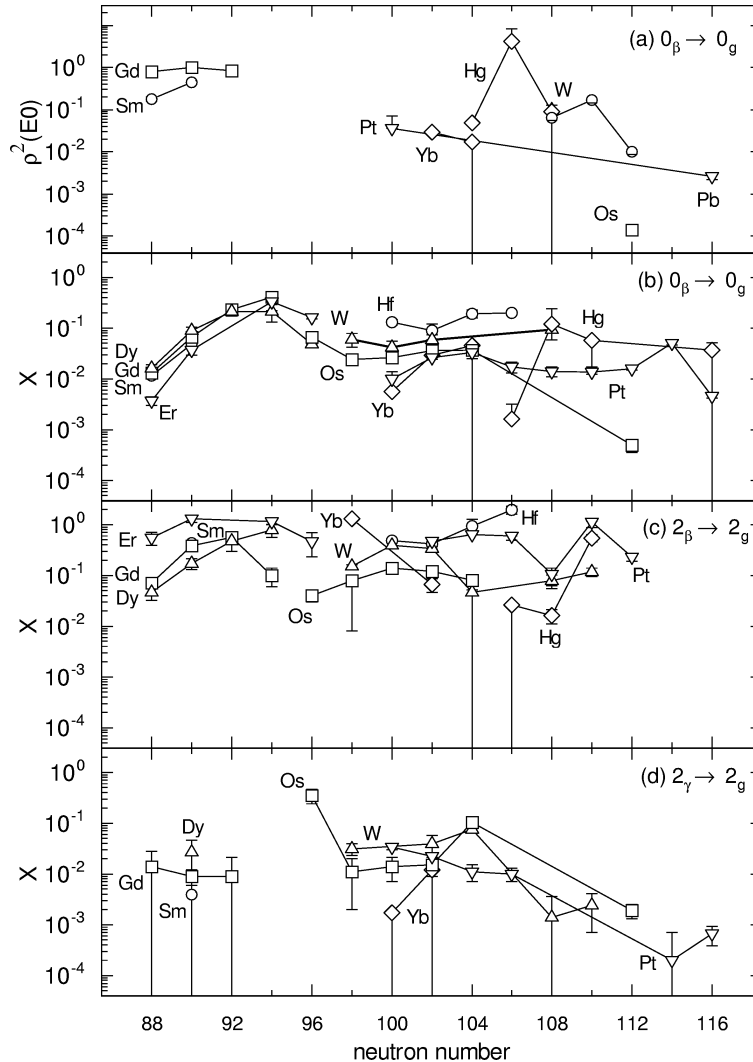
Kibedi, et al., Nucl. Phys. A 688 669 (2001)

# Evolution of shape co-existence W-Os-Pt



Kibedi, et al., Nucl. Phys. A 688 669 (2001)

# Evolution of shape co-existence x(E0/E2) systematics



$\rho^2(E0)$   
 $0_2 - 0_1$

15 data

$B(E0)/B(E2)$   
 $0_2 - 0_1$

42 data

$B(E0)/B(E2)$   
 $2_2 - 2_1$

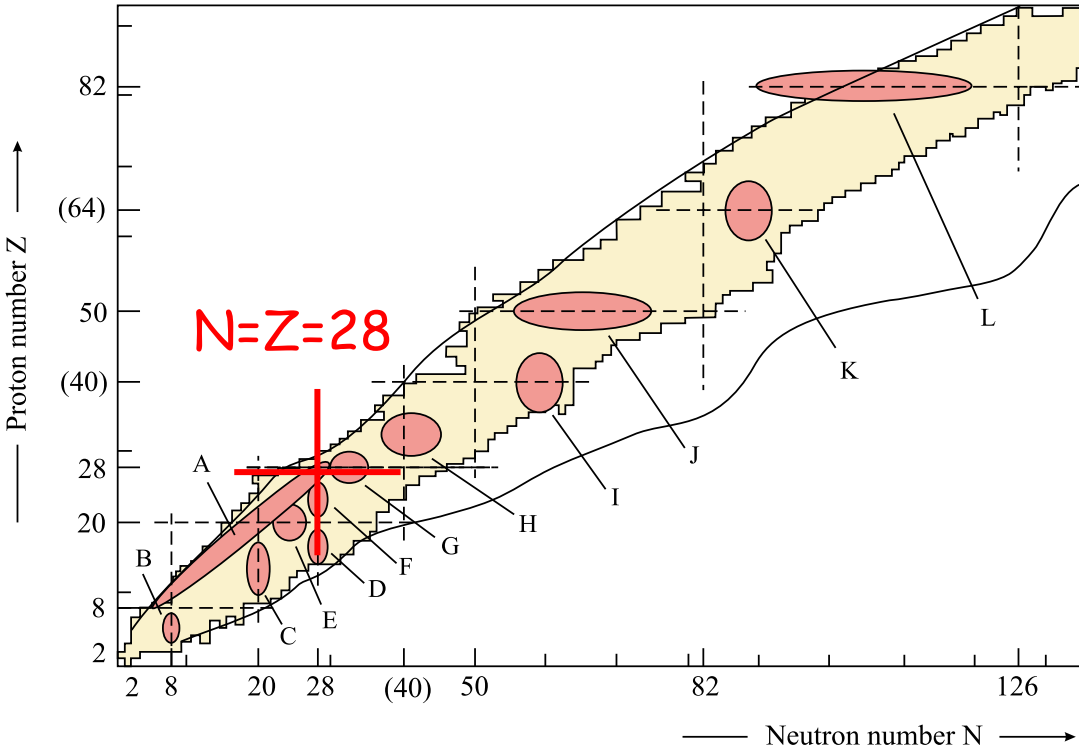
39 data

$B(E0)/B(E2)$   
 $2_{\gamma} - 2_1$

24 data

Kibedi, et al., Nucl. Phys. A 688 669 (2001)

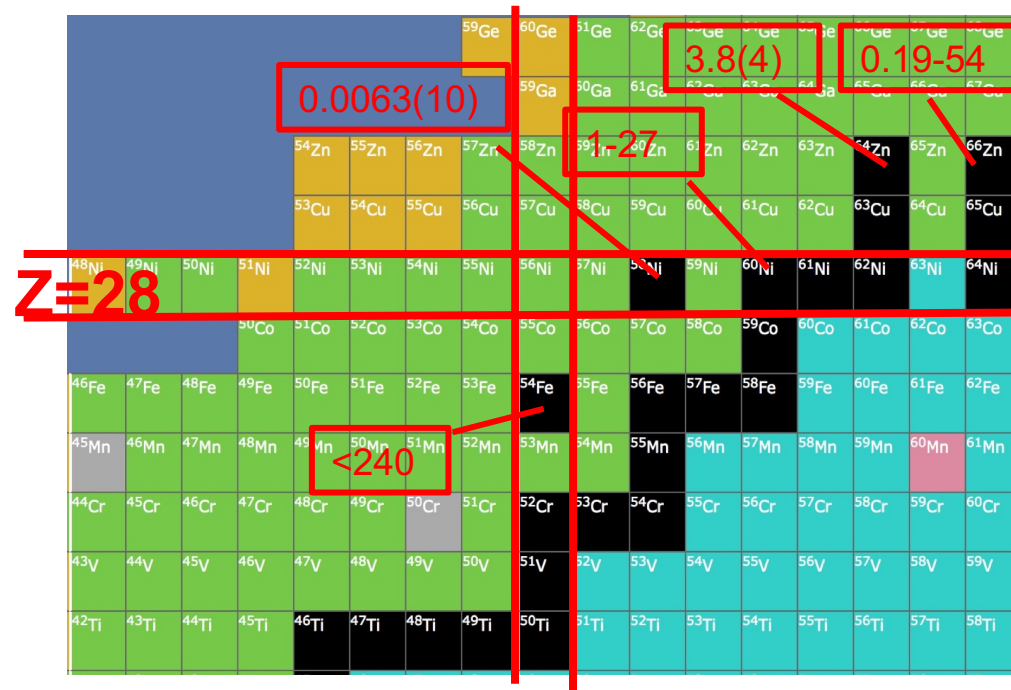
# Shape co-existence around $N=Z=28$



Heyde & Wood, Rev. Mod. Phys. 83 (2011) 1467

$N=Z=28$  double magic  
Shape co-existence "could emerge"

- Excited  $0^+$  around  $N=Z=28$ : from mp-mh excitations from the  $1f_{7/2}$  to the  $1f_{5/2}$ ,  $2p_{1/2}$  and  $2p_{3/2}$  orbits
- E0 transitions: not very well known
- Aim:** Characterise E0 transitions in  $Z=26$ ,  $N=28,30,32$
- E0's in Ni isotopes talk by Adam Garnsworthy

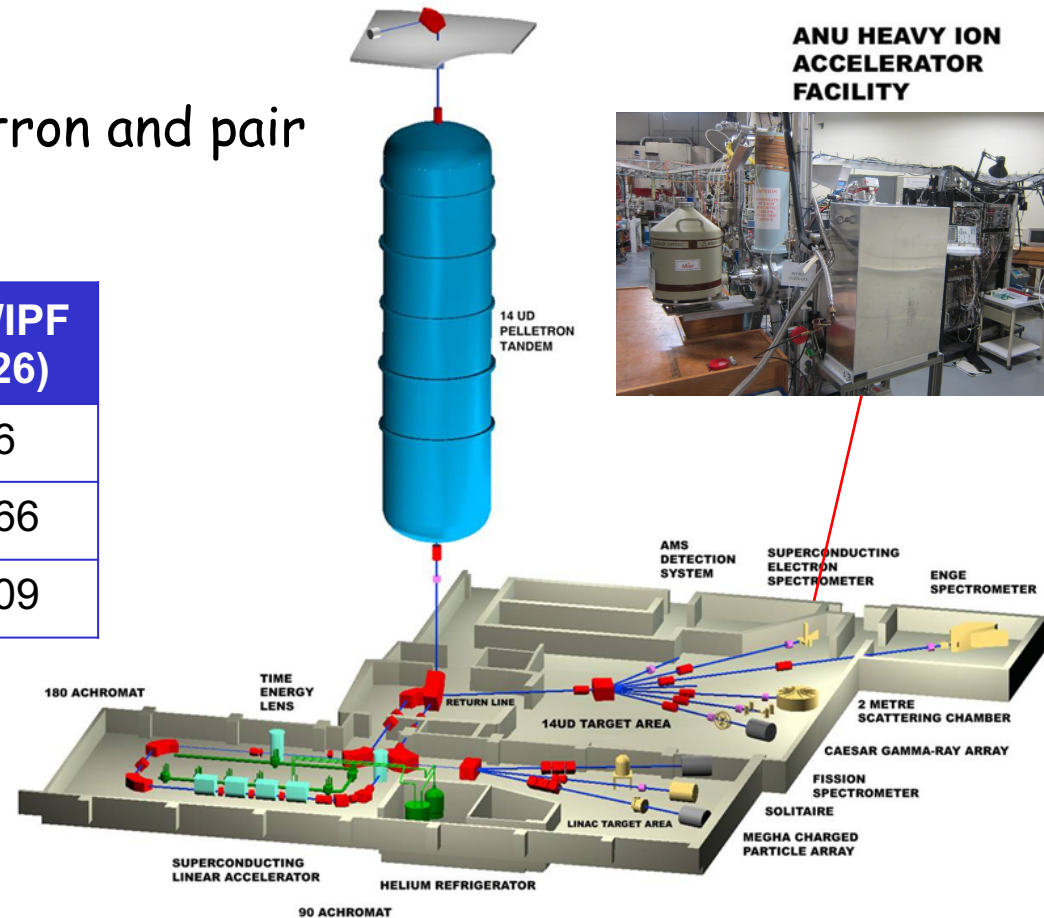
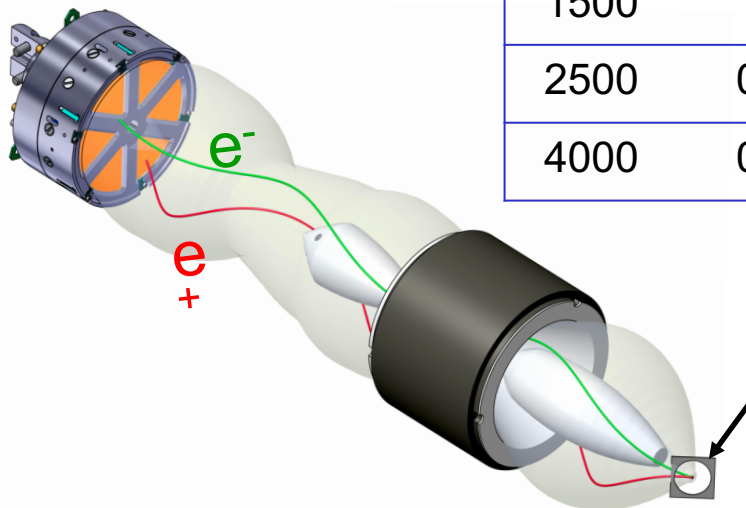


- (p,p') reaction,  $E_p=6.7-7$  MeV DC beam
- 1-2 mg/cm<sup>2</sup> <sup>54,56,58</sup>Fe targets
- Singles gamma, conversion electron and pair conversion (Super-e)

$$\varepsilon_{CE} = 0.1-0.5\%$$

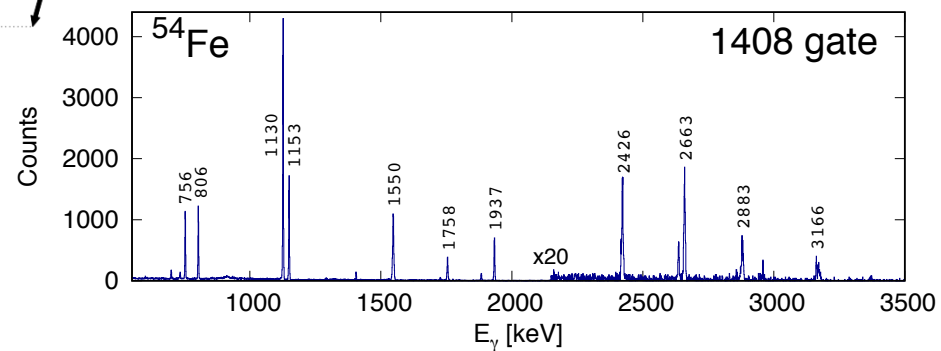
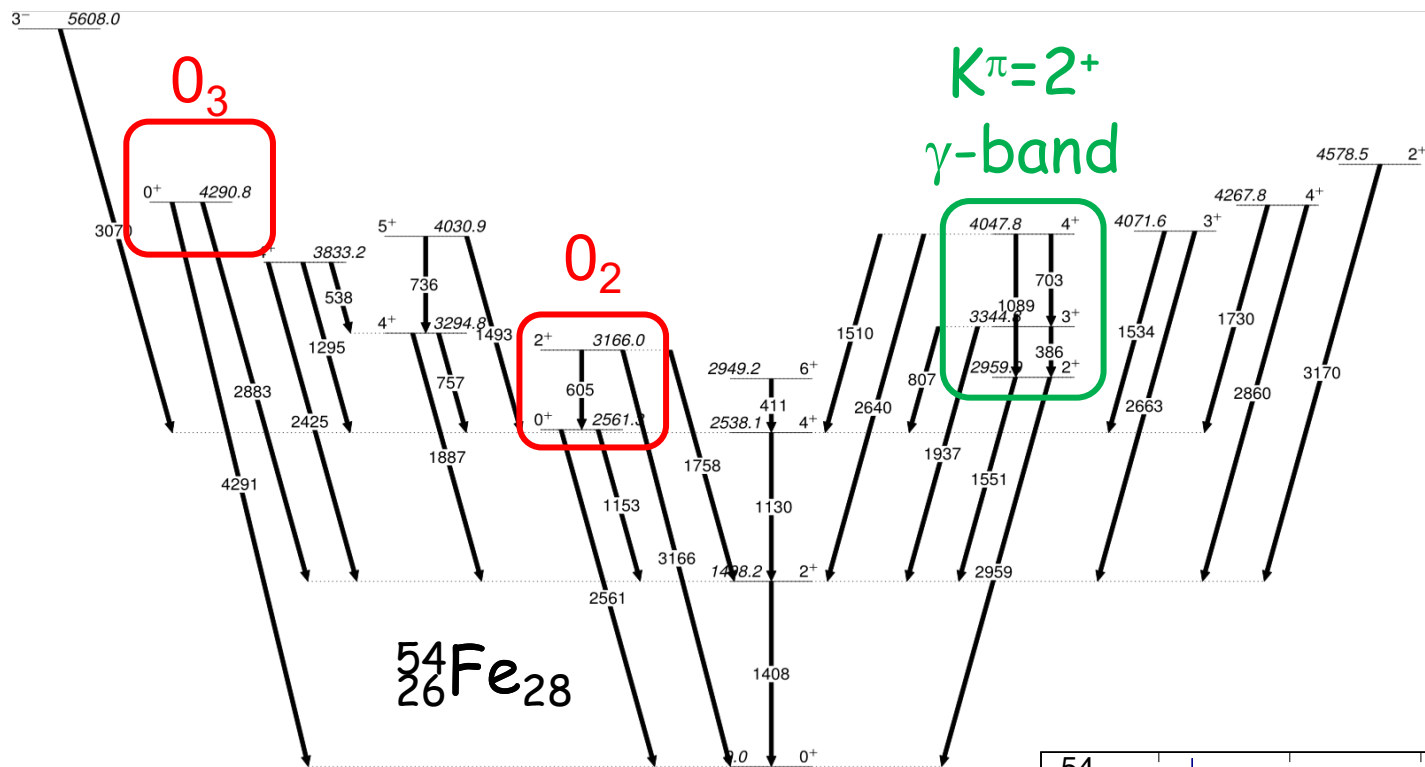
$$\varepsilon_{IPF} = 0.01\%$$

E [keV]	CEK/IPF (Z=26)
1500	1.6
2500	0.066
4000	0.009

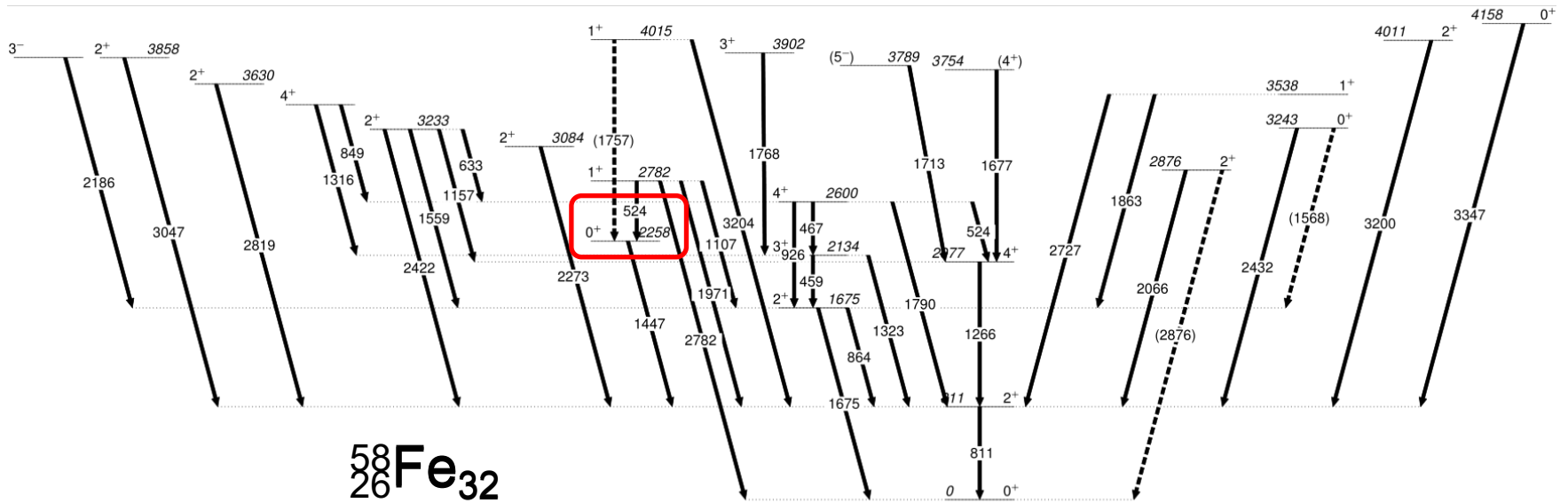




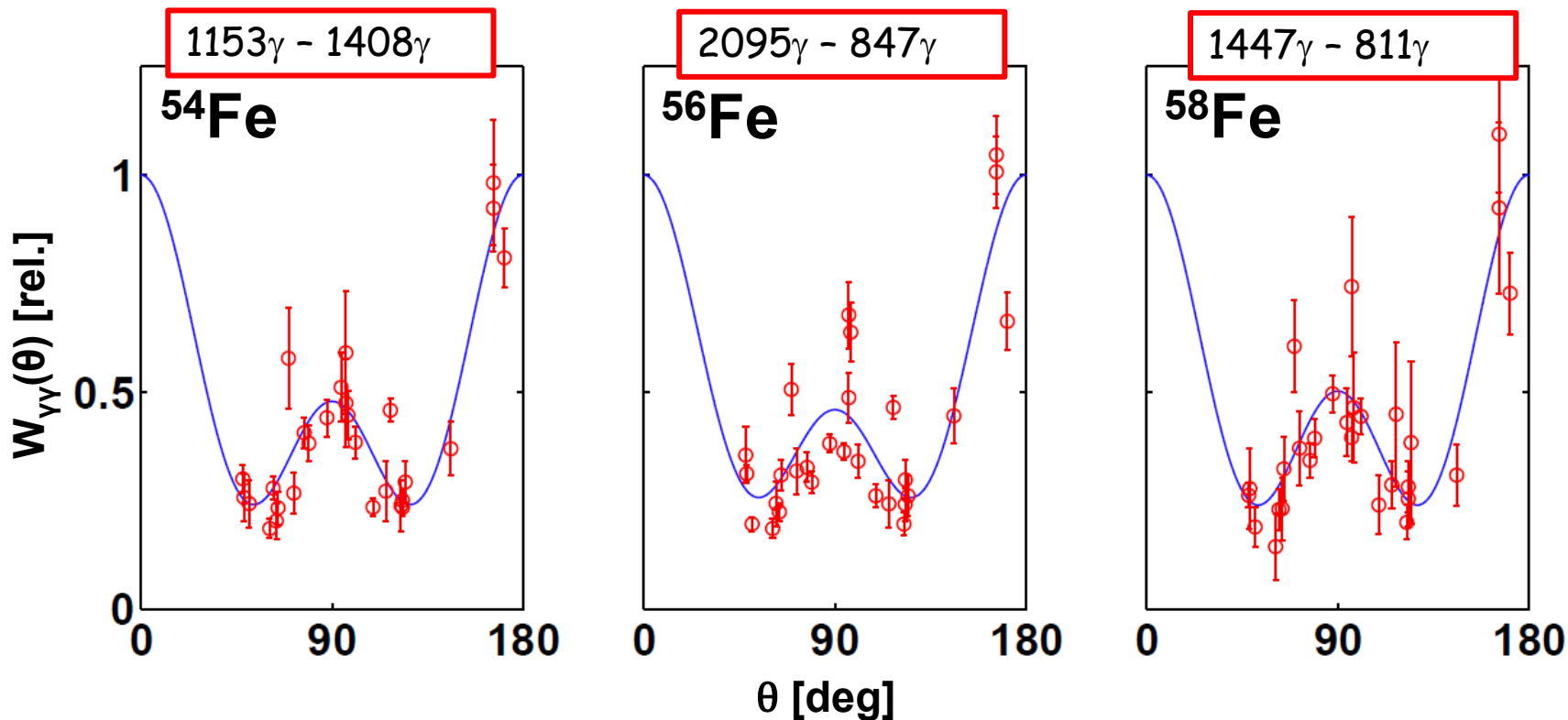
# Extended level schemes



# Extended level schemes



# Angular correlation of 0-2-0 cascades

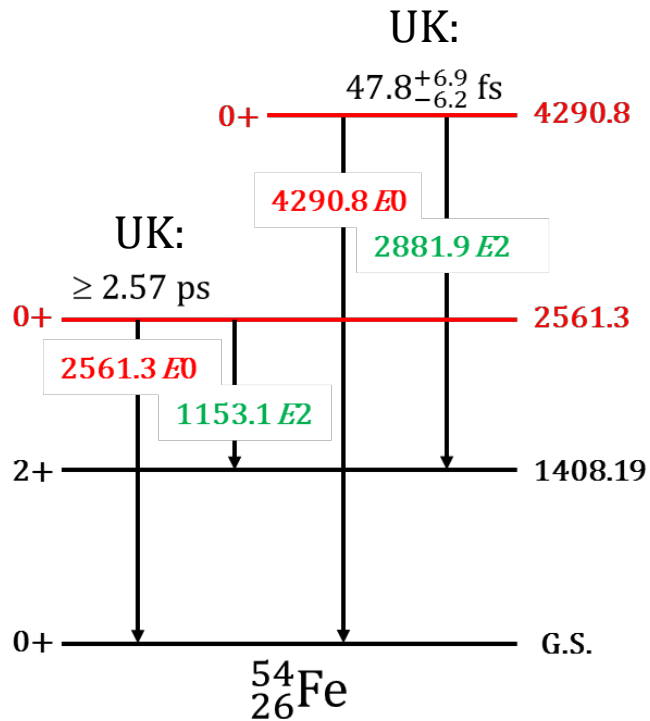


$$W_{YY}(\theta) = N(1 + a_2 \cos^2(\theta) + a_4 \cos^4(\theta))$$

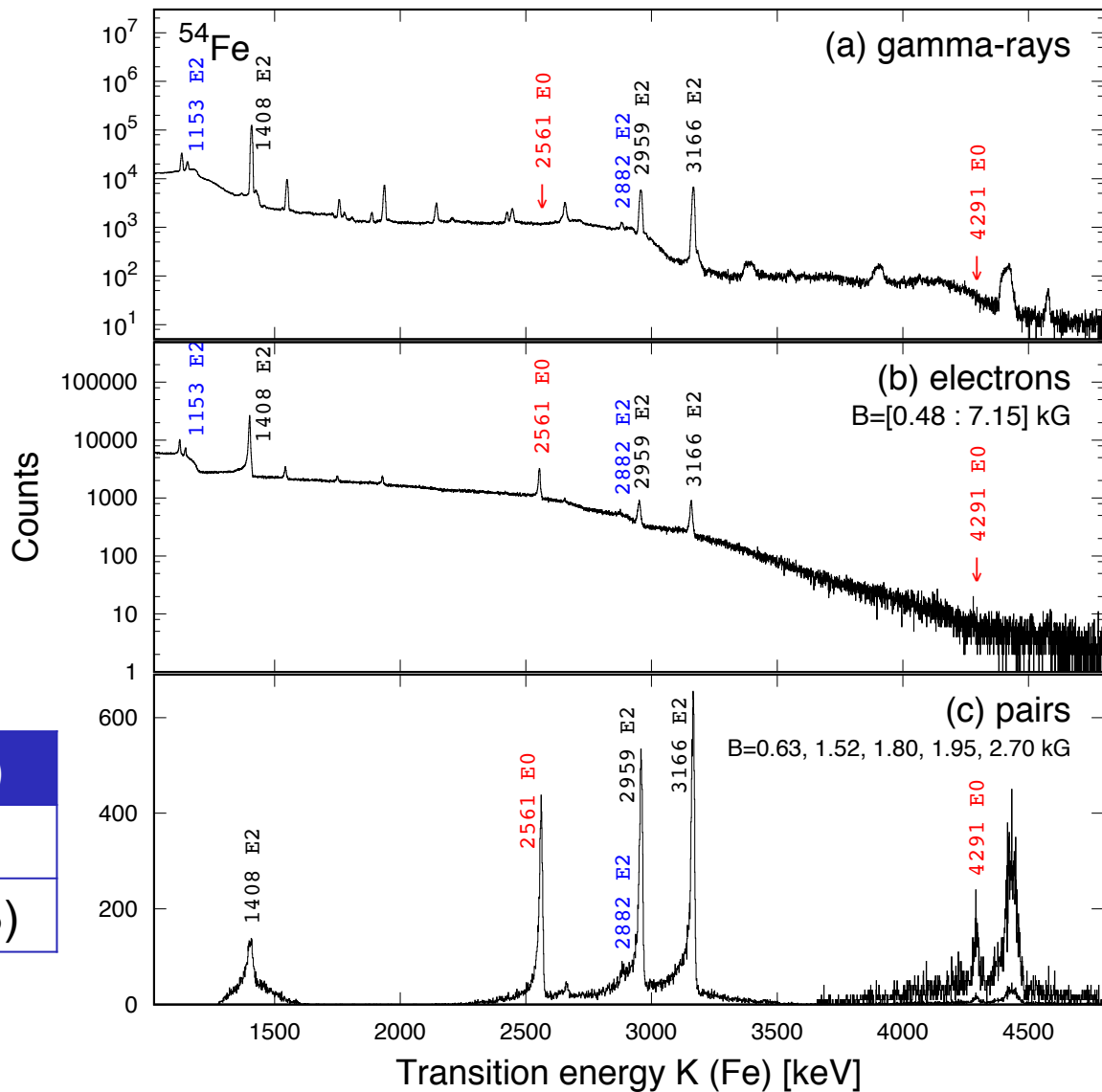
$$0^+ \rightarrow 2^+ \rightarrow 0^+$$

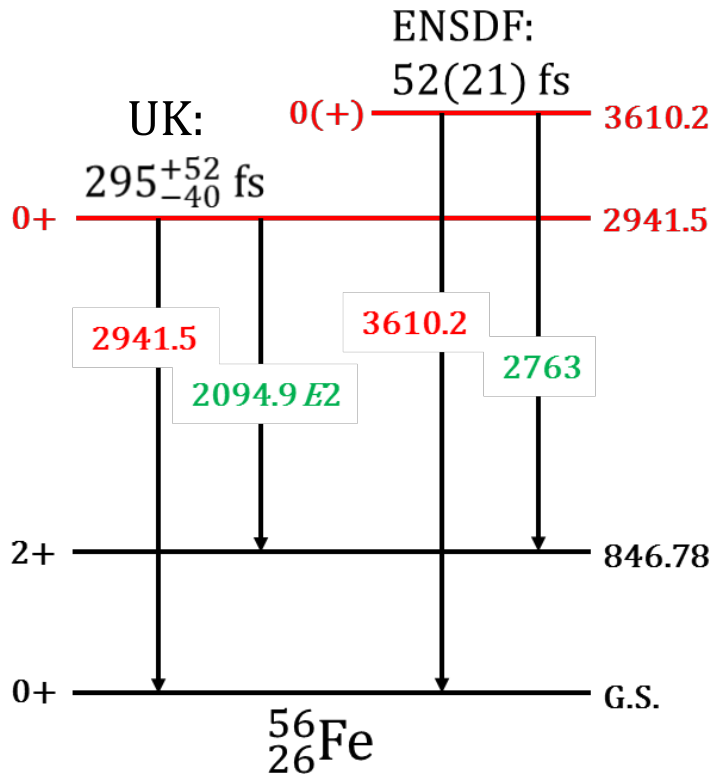
$$a_2 = -3, a_4 = 4$$

	$a_2$	$a_4$
$^{54}\text{Fe}$	$-2.8(6)$	$3.9(7)$
$^{56}\text{Fe}$	$-2.6(9)$	$3.7(10)$
$^{58}\text{Fe}$	$-2.8(8)$	$3.8(9)$

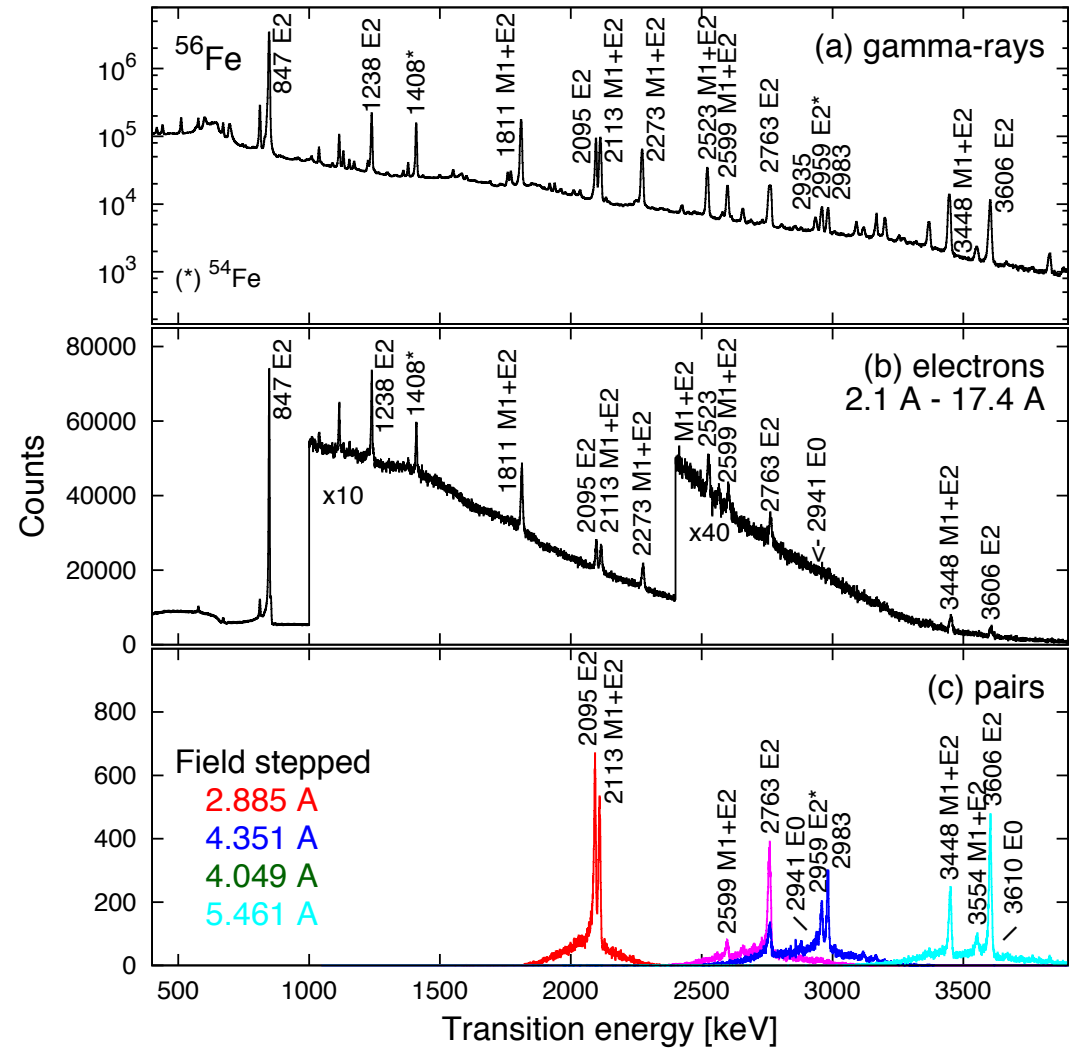


	$10^3 \rho^2(E0)$
2561.0 E0	<80
4290.8 E0	92(+22-23)





	$10^3 \rho^2(E0)$
2941.5 E0	3.2(11)
3610.2 E0	?



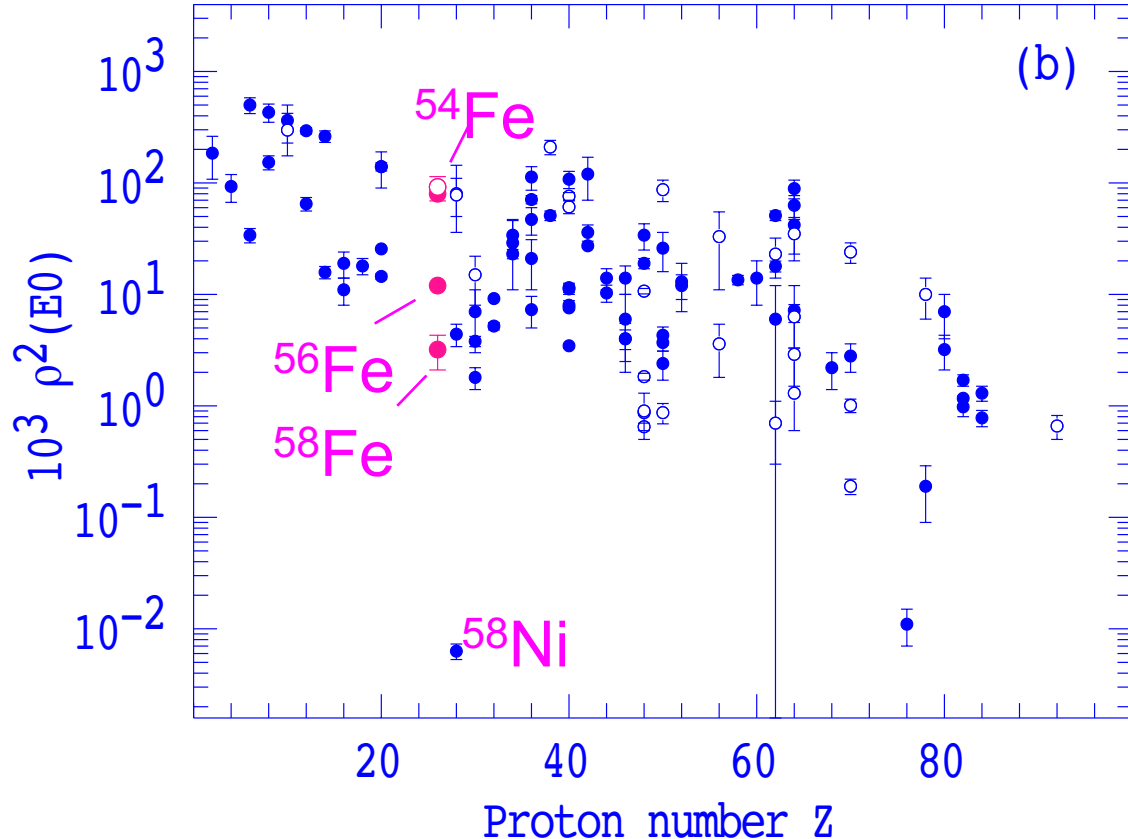
Experimental monopole strength

Monopole strength and band mixing

$$\rho^2(E0) = \frac{1}{\Omega(E0) \times \tau(E0)}$$

$$\rho^2(E0) = \frac{Z^2}{R_0^4} a^2 (1 - a^2) [\Delta\langle r^2 \rangle]^2$$

		$10^3 r^2(E0)$
$^{54}\text{Fe}$	2561 ( $0_2-0_1$ )	<80
	4291 ( $0_3-0_1$ )	92(+22-23)
$^{56}\text{Fe}$	2942 ( $0_2-0_1$ )	3.2(11)
$^{58}\text{Fe}$	2258 ( $0_2-0_1$ )	<12



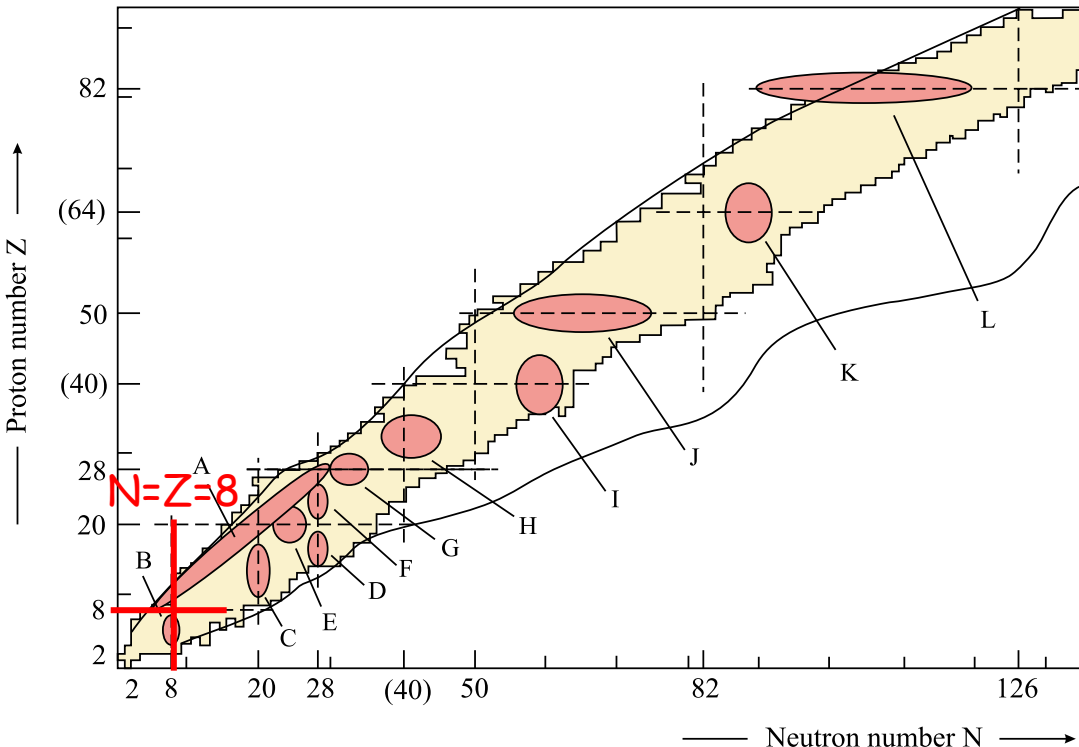
$^{54}\text{Fe}$  1758 ( $2_2-2_1$ )  
 $\alpha_K = 4.9\text{E-}5(3)$  M1+E2  
 Small/negligible E0

# Shape co-existence around N=Z=28

## Summary

Z	20	22	24	26	28	30	32	34	36
Ge	32								
Zn	30							3.8 <sub>4</sub>	0.19-54
Ni	28					0.0063 <sub>10</sub>	1-27		
Fe	26				<80	3.2 <sub>11</sub>	<12		
Cr	24								
Ti	22								
Ca	20	25.6 <sub>7</sub>	140 <sub>12</sub>	140 <sub>50</sub>	14.5 <sub>9</sub>				

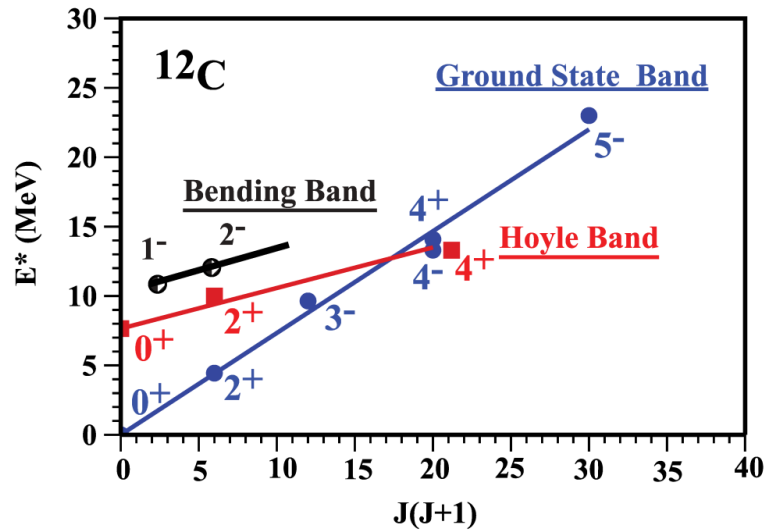
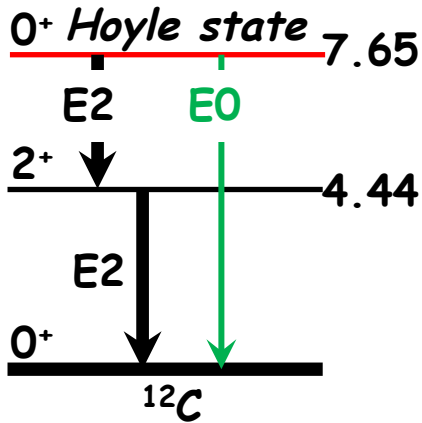
- New results in <sup>54,56,58</sup>Fe: extended level schemes, new T1/2,  $\delta(E2/M1)$ , E0 transitions,  $\rho^2(E0)$
- Future: look for E0s between J>0 states to characterise bands built on excited 0<sup>+</sup> states
- Interpretation within the bandmixing approach
- E0s in Cr (Z=24) and Ti(Z=22) to explore N=28 isotones



Heyde & Wood, Rev. Mod. Phys. 83 (2011) 1467

Hoyle state: not a typical excited  $0^+$  at south-west from the double magic  $^{16}\text{O}$

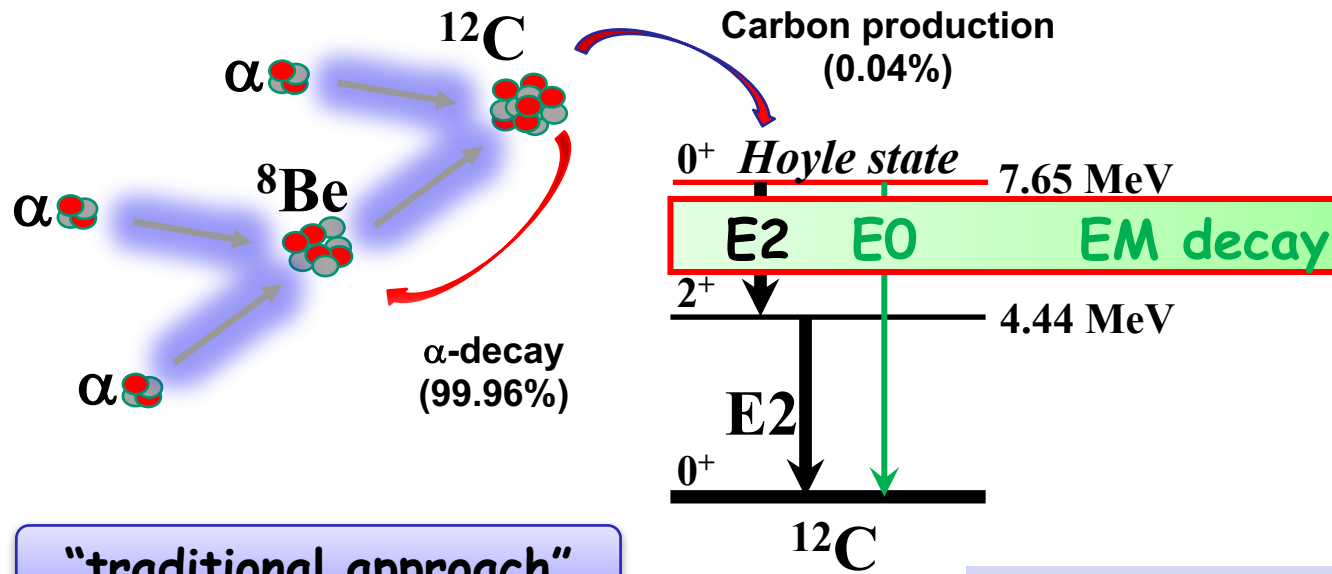




D.J. Marín-Lámbbari, et al., PRL 113 (2014) 012502

- $\Gamma(0_2)=9.3(9)$  eV;  $T_{1/2}(0_2)=3.5(3) \times 10^{-17}$  s
- “Extended object” (Brink 1966)  
RMS=2.89(4) fm = 1.2\* RMS(g.s.)  
PRC 80 (2009) 054603
- $\rho^2(E0)=500(81)$  ADNDT 89 (2005) 77
- $2^+$  at 9.8 MeV  
Nucl. Phys. A738, (2004) 268; Phys. Rev. C 84 (2011) 054308; 80 (2009) 041303(R); 84 (2011) 027304; 86 (2012) 034320; PRL 113 (2014) 012502
- $4^+$  at 13.3 MeV  
Phys. Rev. C83 (2011) 034314
- $0_3^+$  at 10.3 MeV;  $\Gamma(0_3)=2.7$  MeV  
Nucl. Phys. A738, (2004) 268
- Microscopic  $\alpha$ -cluster model /exp  
 $E(0_2)-E_{3\alpha}=0.23 / 0.38$  MeV  
 $\Gamma(0_2)=7.6 / 9.3(9)$   $\mu\text{eV}$   
 $M(E0)=6.3 / 5.4(2)$   $\text{fm}^2$   
Yasuro Funaki, Phys. Rev. C 94 (2016) 024344

# The radiative width of the Hoyle state



“traditional approach”

$$r_{3\alpha} \propto [\Gamma_{\text{rad}}] \exp(-[Q_{3\alpha}] / kT)$$

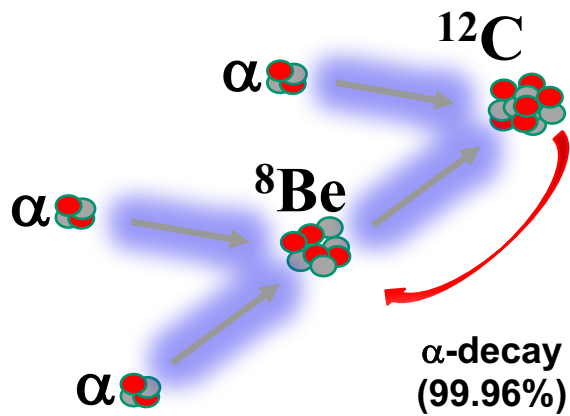
$$\Gamma_{\text{rad}} = \left[ \frac{\Gamma_{\text{rad}}}{\Gamma} \right] \times \left[ \frac{\Gamma}{\Gamma_{\pi}(E0)} \right] \times [\Gamma_{\pi}(E0)]$$

$$\Gamma_{\text{rad}} = 3.9(4) \times 10^{-3} \text{ eV EM}$$

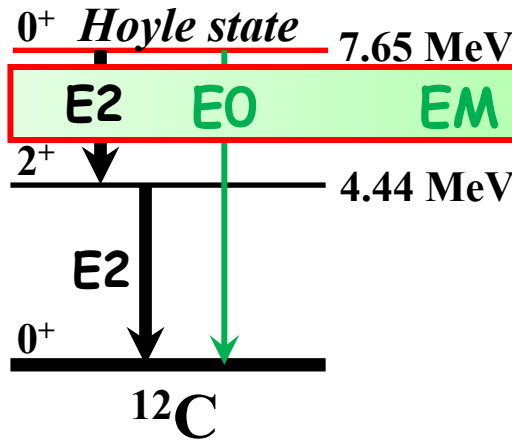
branching: 0.042%

		Unc. [%]	# exp
$\Gamma_{\text{rad}}/\Gamma (\times 10^{-4})$	4.19(11)	2.7	8
$\Gamma_{\pi}(E0)/\Gamma (\times 10^{-6})$	6.7(6)	8.9	4
$\Gamma_{\pi}(E0) (\mu\text{eV})$	62.3(20)	3.2	8

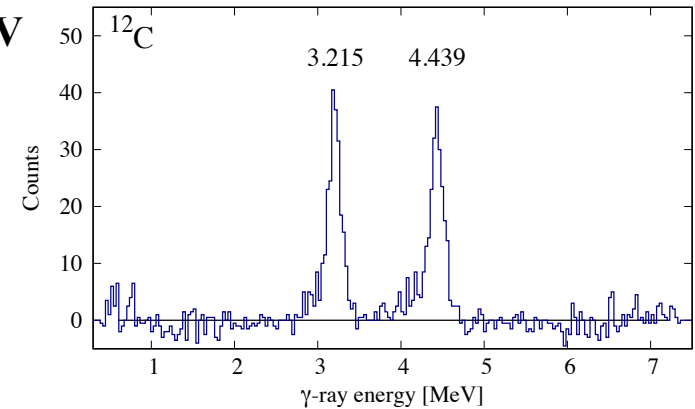
# The radiative width from p- $\gamma$ - $\gamma$



Carbon production (0.04%)



Univ. Oslo Cyclotron  
CACTUS + SiRi



"traditional approach"

$$r_{3\alpha} \propto [\Gamma_{rad}] \exp(-[Q_{3\alpha}] / kT)$$

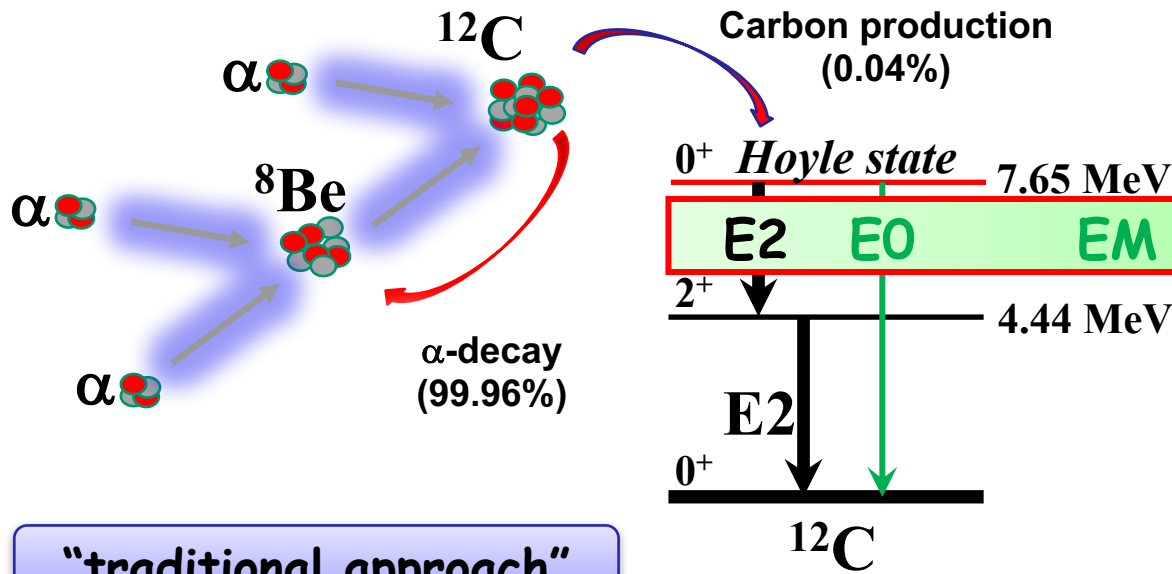
$$\Gamma_{rad} = \left[ \frac{\Gamma_{rad}}{\Gamma} \right] \times \left[ \frac{\Gamma}{\Gamma_{\pi}(E0)} \right] \times [\Gamma_{\pi}(E0)]$$

Events in 12 days:

- Total: 6.0E+9
- p-singles (7.65): 2.72E+8
- p- $\gamma$ - $\gamma$  (7.65): 260(16)

$$\left[ \frac{\Gamma_{rad}}{\Gamma} \right] = 2.76(21)\text{E-4}$$

# The radiative width from pair conversion measurements



ANU Super-e pair spectrometer

"traditional approach"

Out of the ~0.04%

$\Gamma_\gamma(E2) = 98.5\%$

$\Gamma_\pi(E0) = 1.5\%$

$\Gamma_\pi(E2) = 0.088\%$

$\Gamma_{CE}(E2), \Gamma_{CE}(E0) \sim 0$

Internal pair conversion

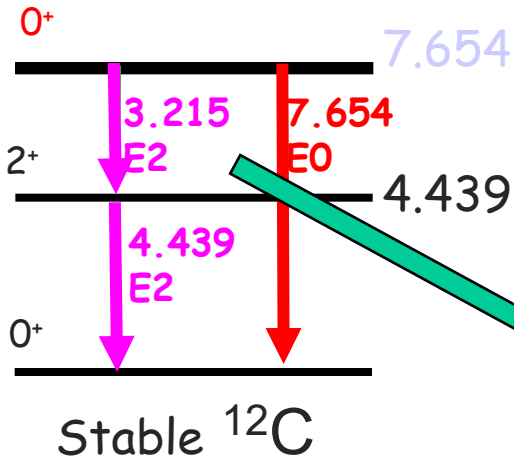
$$r_{3\alpha} \propto [\Gamma_{rad}] \exp(-[Q_{3\alpha}] / kT)$$

$$\Gamma_{rad} = \left[ \frac{\Gamma_{rad}}{\Gamma} \right] \times \left[ \frac{\Gamma}{\Gamma_\pi(E0)} \right] \times [\Gamma_\pi(E0)]$$

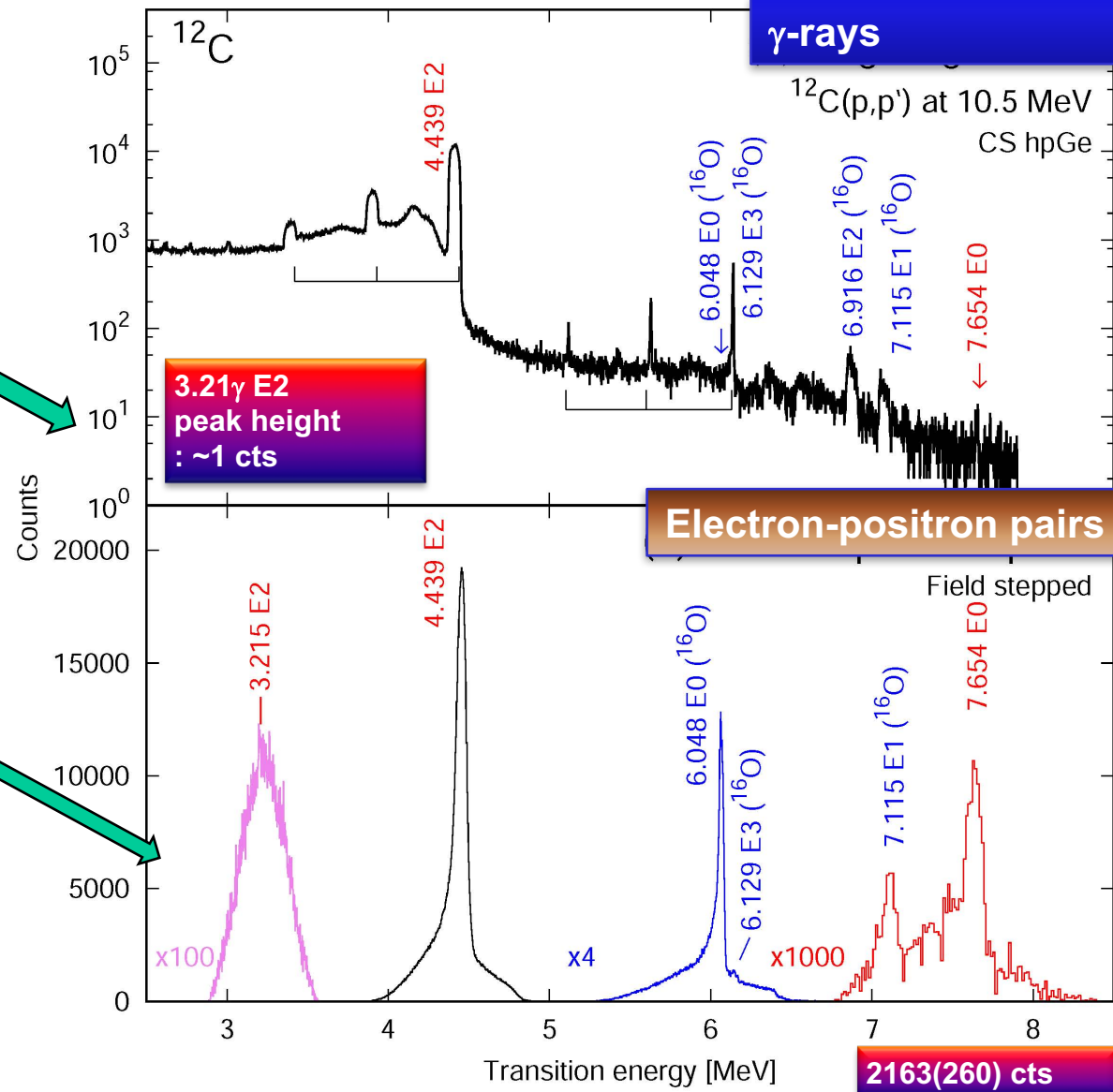
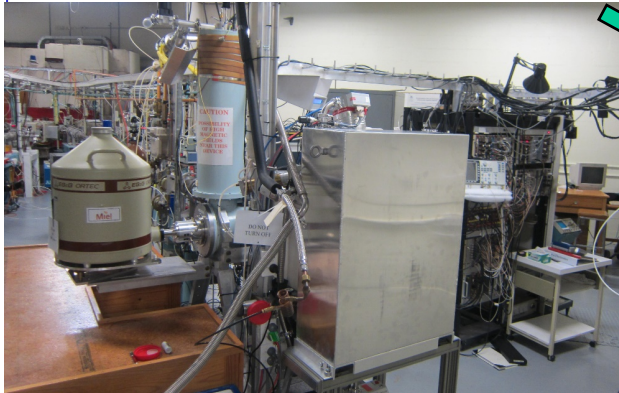
$$\Gamma_{rad} = \left[ \frac{\Gamma_\pi(E2)}{\Gamma_\pi(E0)} \right] \times \left( 1 + \frac{1}{[\alpha_\pi(E2)]} \right) + 1 \times [\Gamma_\pi(E0)]$$

$$\alpha_\pi(E2) = 8.765 \times 10^{-4}$$

# The radiative width from pair conversion measurements

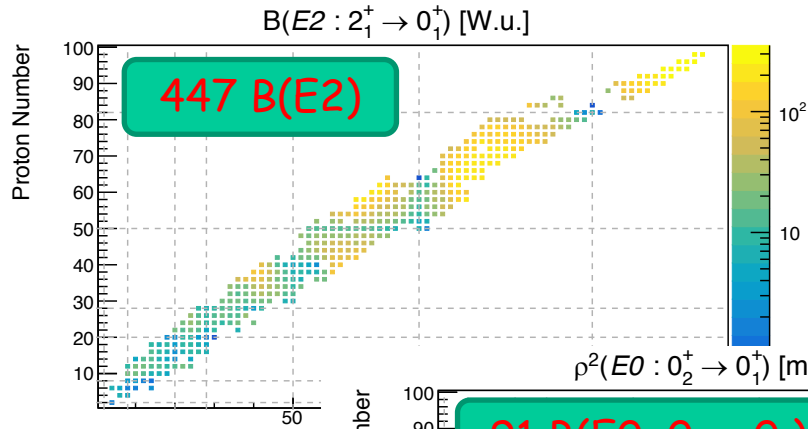


	gamma	pairs
4.439 E2	= 1	1.32E-3
3.215 E2	1.1E-4	9.2E-8
7.654 E0	--	1.6E-6

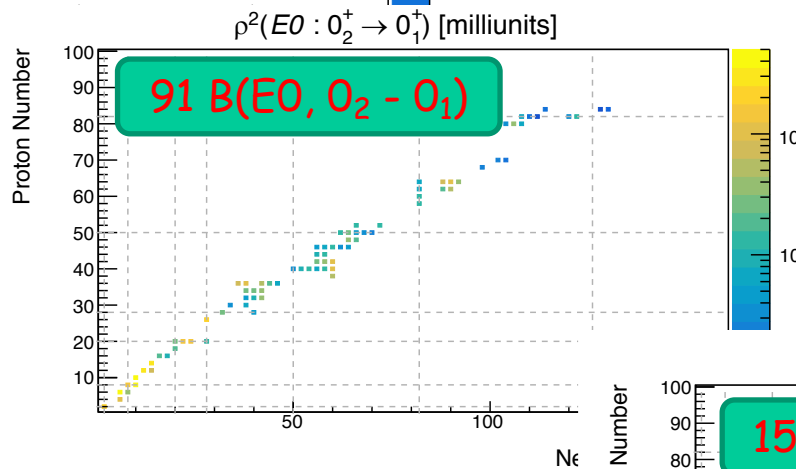


# Properties of Monopole transitions in atomic nuclei

With J.L. Wood and A. Garnsworthy

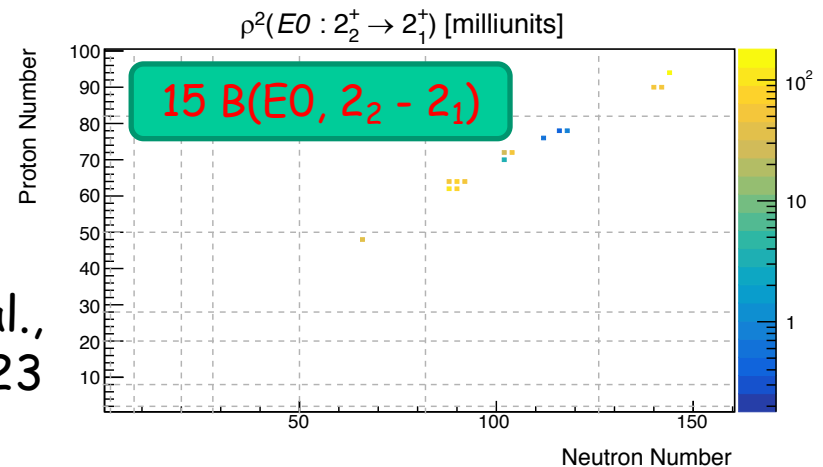


Pritychenko et al.,  
ADNDT 107 (2016) 1



Kibèdi & Spear,  
ADNDT 89 (2005) 77

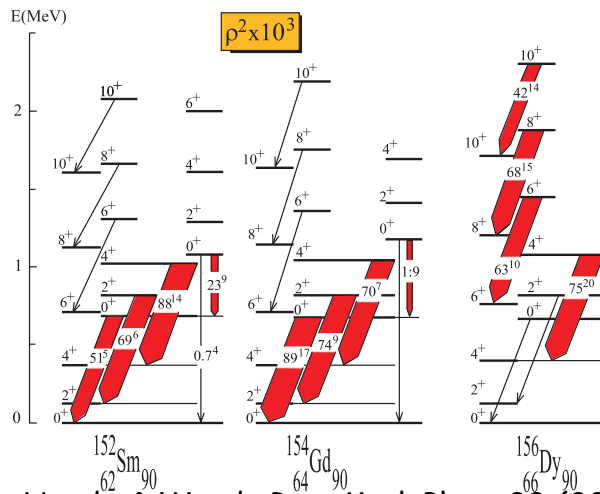
Wood et al.,  
Nucl. Phys. **A651** (1999) 323



# Properties of Monopole transitions in atomic nuclei

## Motivation:

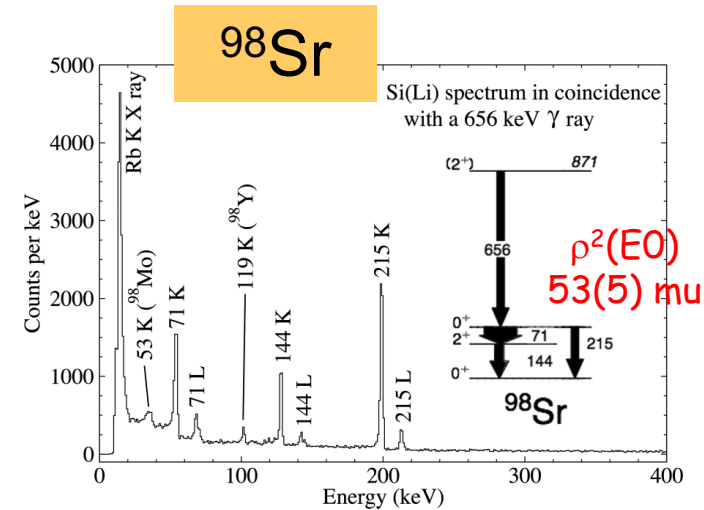
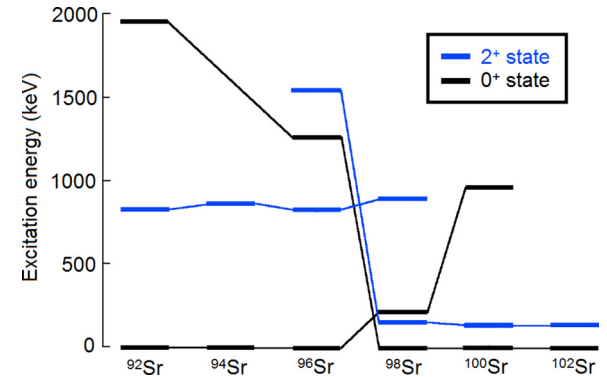
- New data since the last evaluations
- First pass of ENSDF: 174 j-j (j>0) transitions
- $E_g = [41.9:1877] \text{ keV}$
- $J_i = 1[8], 2[109], 3[4], 4[32], 5[3], 6[6], 8[2], 9[1], 10[1], 16[1]$



Heyde & Wood., Rev. Mod. Phys. 83 (2011) 1467

E0 strength model independent probe of co-existing structures -spin dependence

With J.L. Wood and A. Garnsworthy



Park et al., PRC 93, 014315 (2016)

Strongly deformed g.s., weakly deformed  $0^+$  at 215 keV

# Properties of Monopole transitions in atomic nuclei

With J.L. Wood and A. Garnsworthy

## Approach

- Consistent treatment of the data known
  - ❑ Combine data on  $0^+ \rightarrow 0^+$  and  $J^+ \rightarrow J^+$  ( $J > 0$ ) transitions
  - ❑ Collect and adopt  $T_{1/2}$ , EM branching ratios, multipolarities, mixing ratios, experimental conversion coefficients from original references
  - ❑ Accept data if:  $T_{1/2}$ , ICC and mixing ratio are known
- New conversion coefficients,  $\Omega(E0)$  electronic factors

$^{238}\text{U}$  966.1 keV  $2^+$

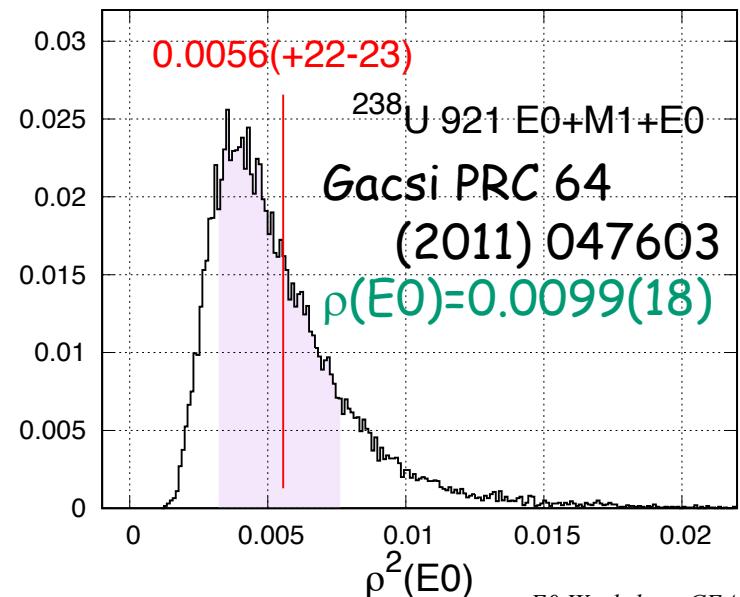
$T_{1/2} = 2.4(+17-7)$  ps

921.2 keV  $E0 + M1 + E2$

$\alpha_K = 0.191(30)$

$\alpha_K(M1) = 0.0390$

$\alpha_K(E2) = 0.00966$





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