

# Review of shell-model results in $^{100}\text{Sn}$ region

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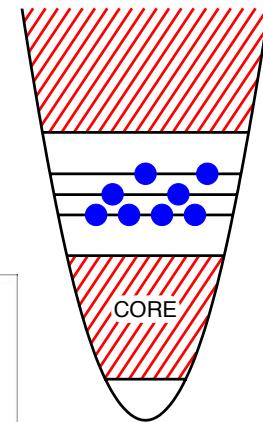
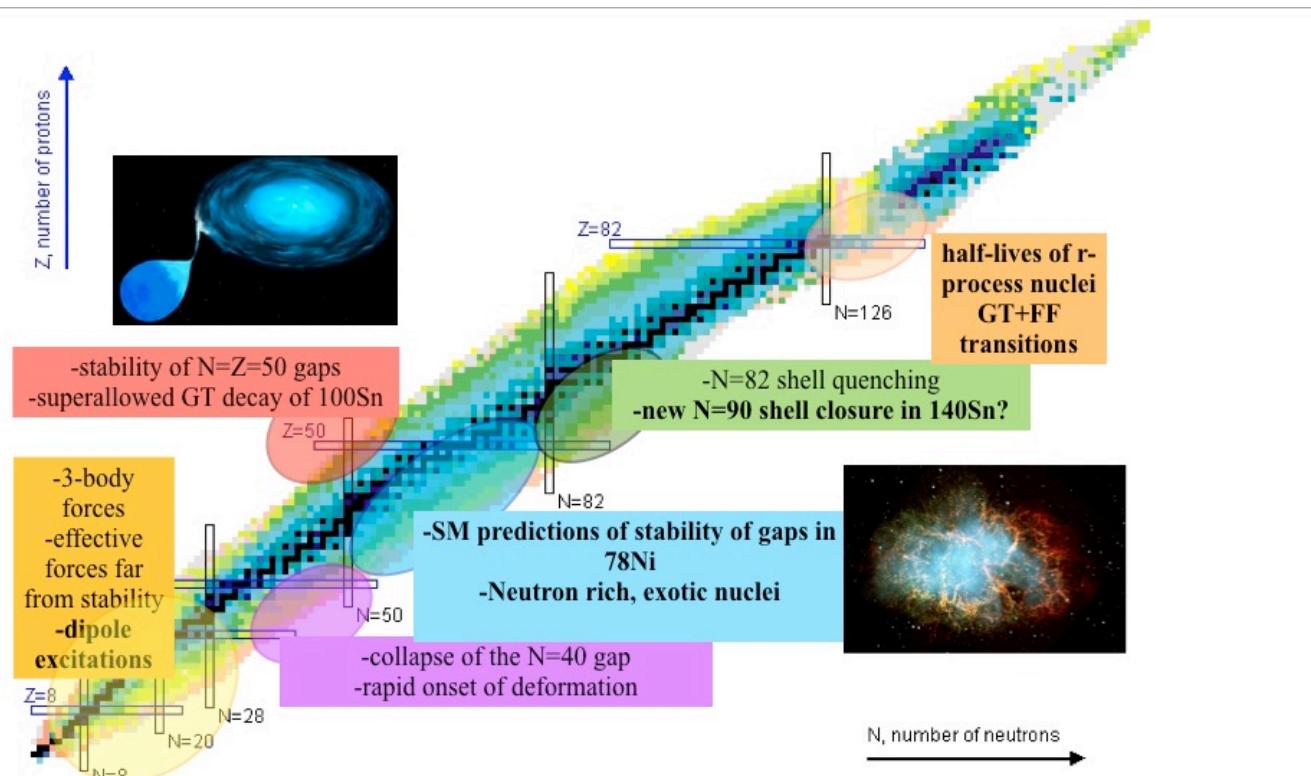
27-31.03.2017

## **Several shell-model studies in $^{100}\text{Sn}$ region:**

- Stability of shell gaps  $N=Z=50$  and single-particle structure of  $^{101}\text{Sn}$ 
  1. Core-excited states, spin-gap isomers
  2. The puzzle of  $B(E2)$  transitions
  3. Superallowed GT decay of  $^{100}\text{Sn}$
  
- Proton-neutron pair correlations in  $N\sim Z$  systems
  1. Collectivity in light Xe isotopes

# SM studies: generalities

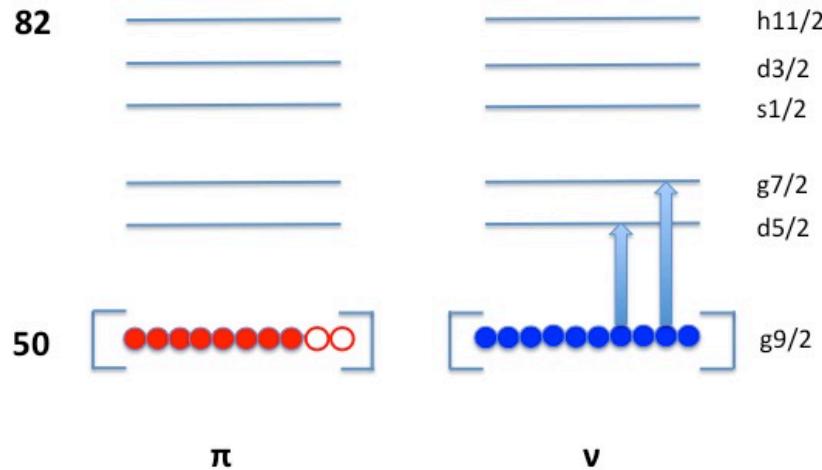
- SM relies on the possibility of diagonalizing the Hamiltonian matrix and constraining (empirically) the effective interaction



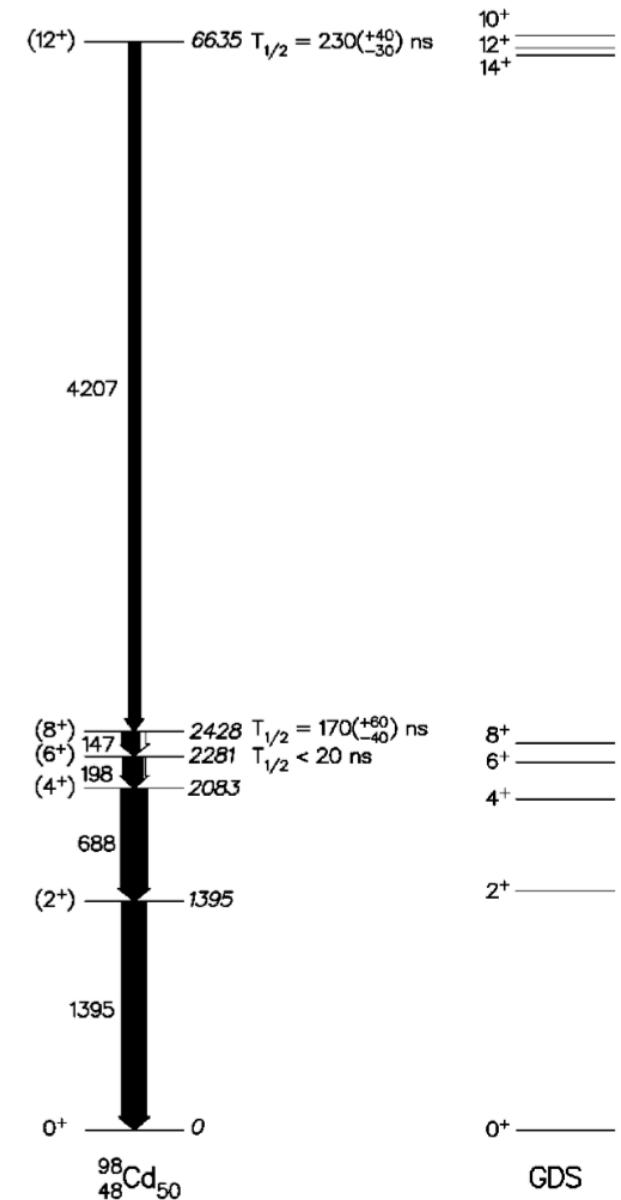
$$H_{\text{eff}} \Psi_{\text{eff}} = E \Psi_{\text{eff}}$$

# Core-excited states: $^{98}\text{Cd}$

A. Blazhev et al., Phys. Rev. C69, 064304, 2004

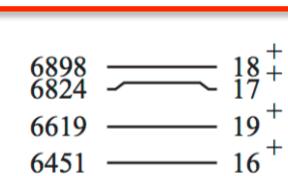
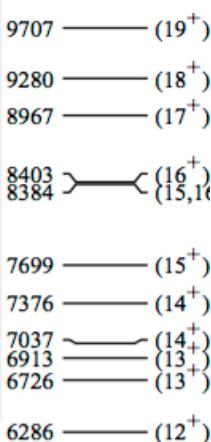
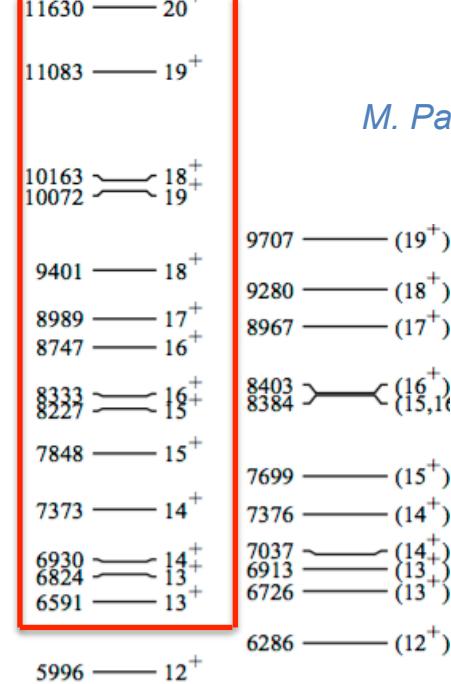


- Shell model calculations in  $sdg$  model space, up to 4p-4h excitations across the N=Z=50 gaps
- Clean constraint on the N=50 shell gap: **6.5MeV** estimated from SM calculations

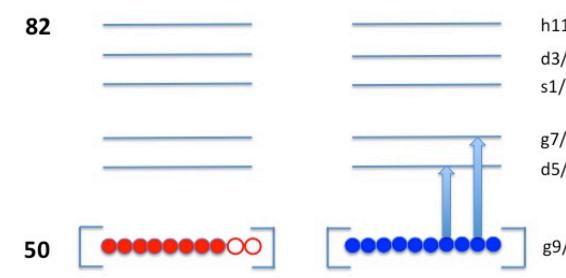
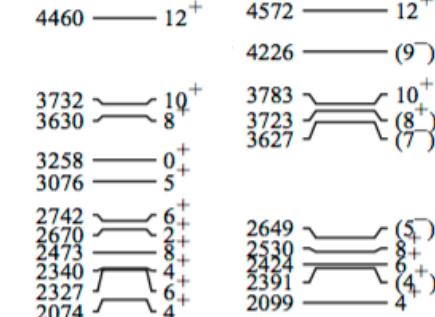


# Core excited states: $^{96}\text{Pd}$ , $^{96}\text{Ag}$

M. Palacz, J. Nyberg, H. Grawe, K. Sieja et al., Phys. Rev. C88, 014318, 2012  
 P. Boutachkov et al., Phys. Rev. C84, 044311, 2011



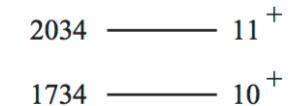
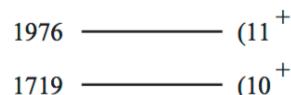
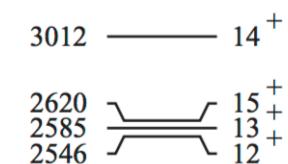
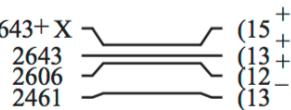
- High-spin states dominated by neutron-core excitations consistent with 6.5-6.7 MeV N=50 gap



$^0\text{LSSM}$

$^0\text{EXP}$

Z=46, N=50



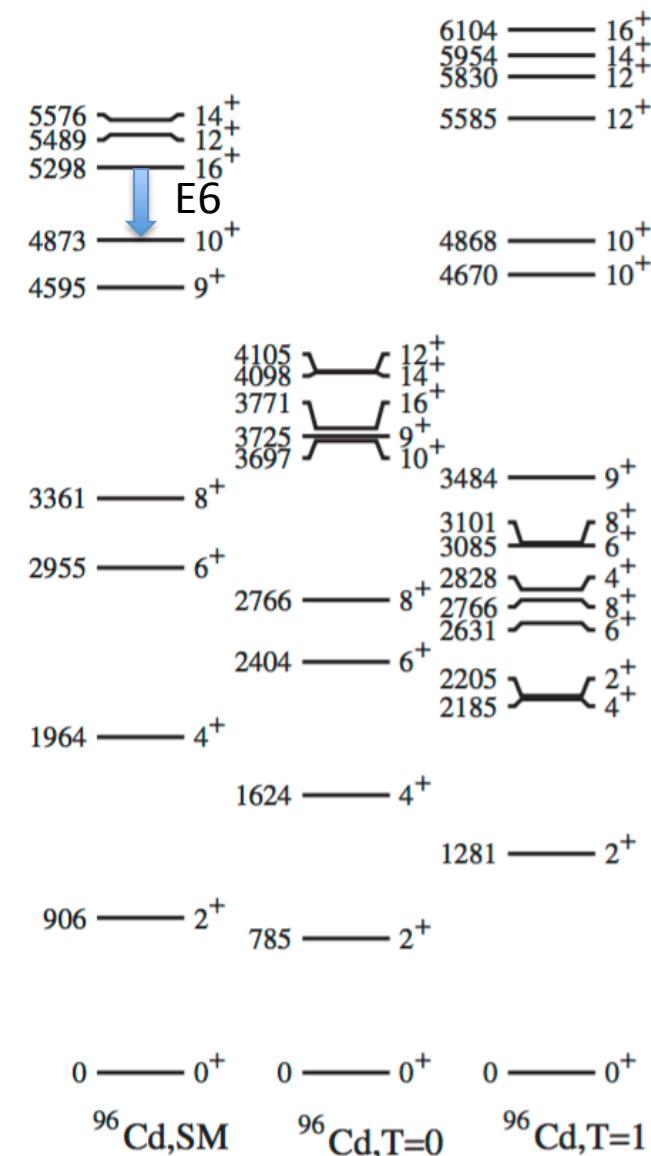
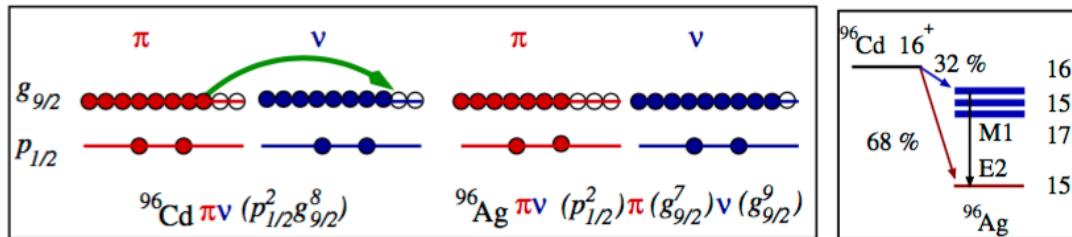
$^0\text{EXP}$

$^0\text{GDS t=5}$

Z=47, N=49

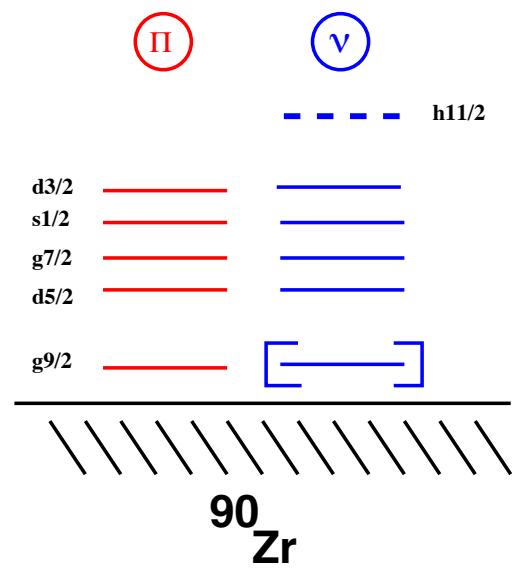
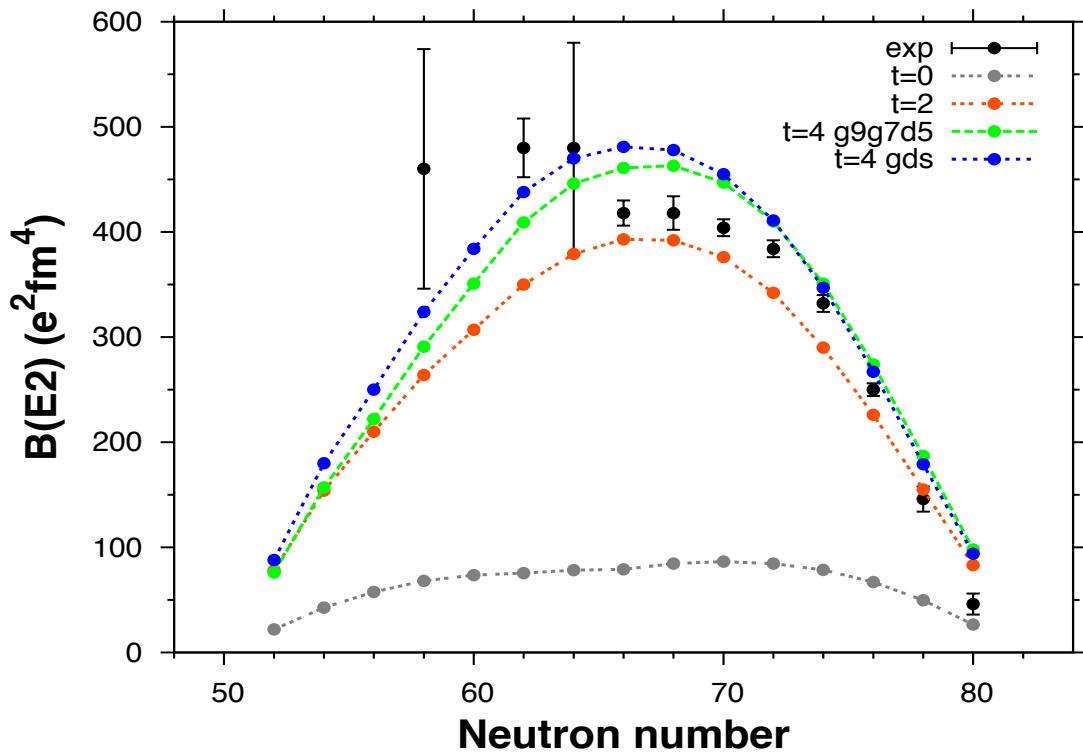
# Spin-gap 16<sup>+</sup> isomer in <sup>96</sup>Cd and T=0 interaction

B. S. Nara Singh et al. Phys. Rev. Lett. 107, 172502, 2011



- Long-standing shell-model prediction of the spin-gap isomer *Phys. Rev. C28, 958, 1983* confirmed in experiment and in modern large-scale shell-model calculations with core excitations
- Important role of the T=0 p-n interaction in the creation of isomer

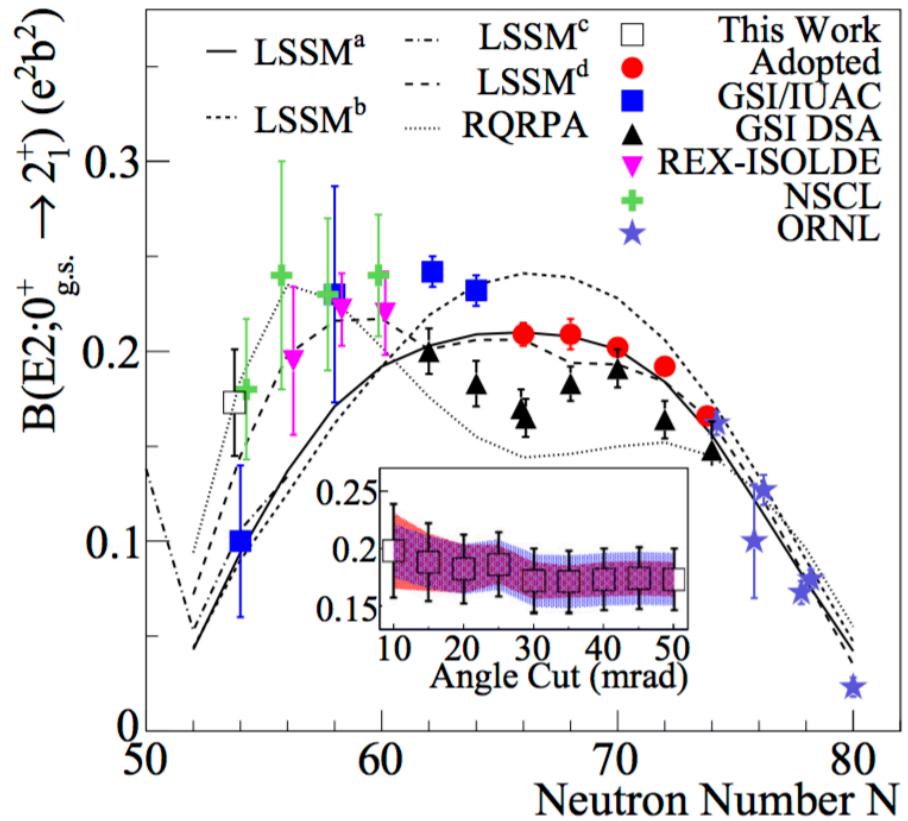
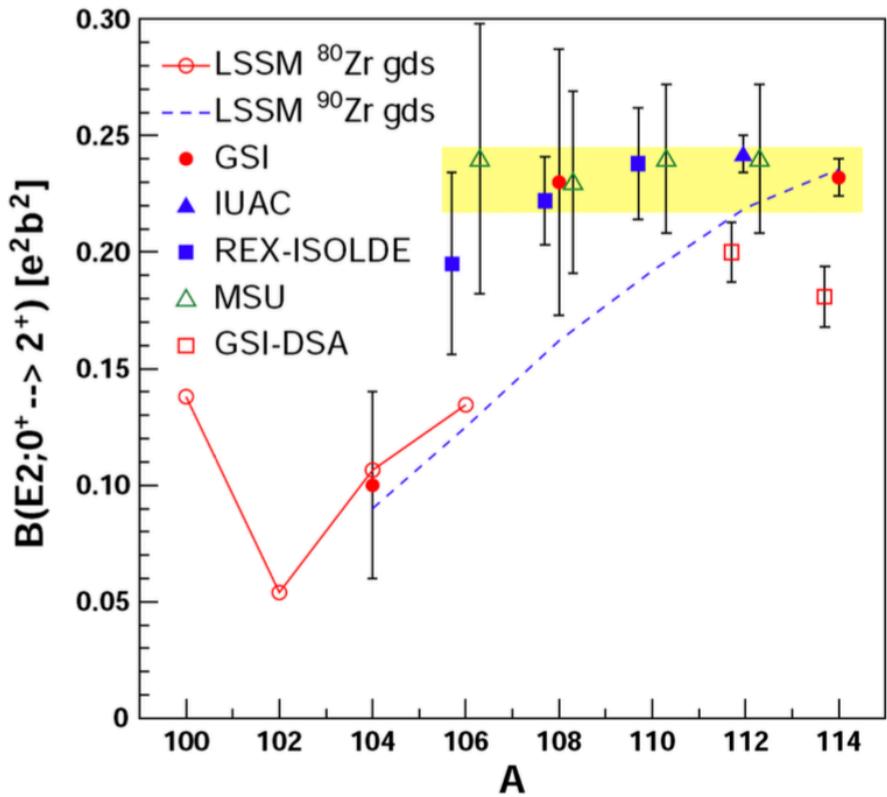
# B(E2) values of tins: reminder



- Full space diagonalization possible
- Not applicable for light tins

A. Banu et al., Phys. Rev. C72, 06135(R), 2005

# B(E2) transitions in light tins



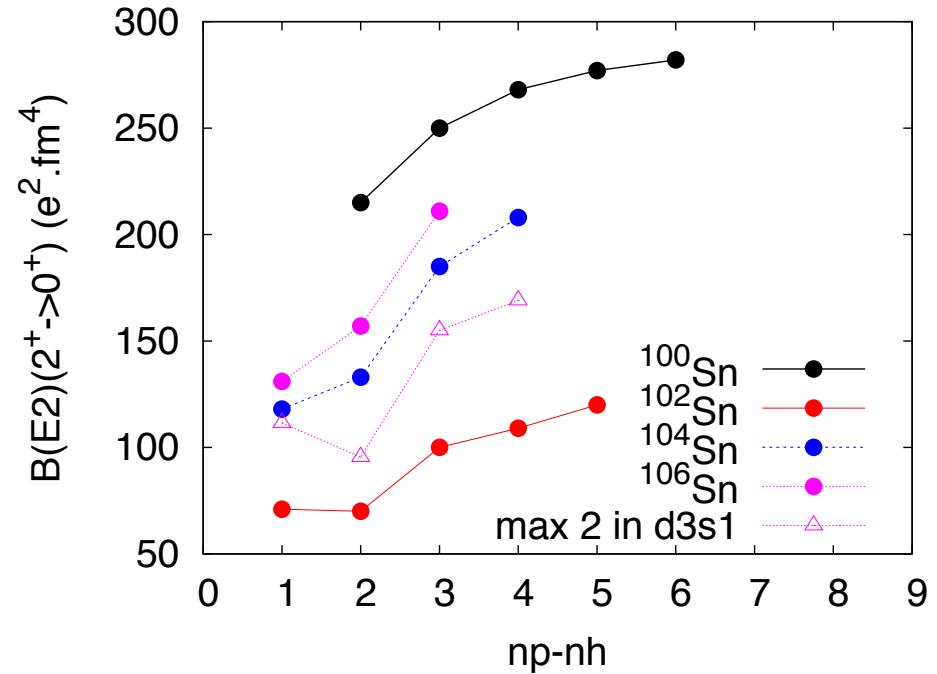
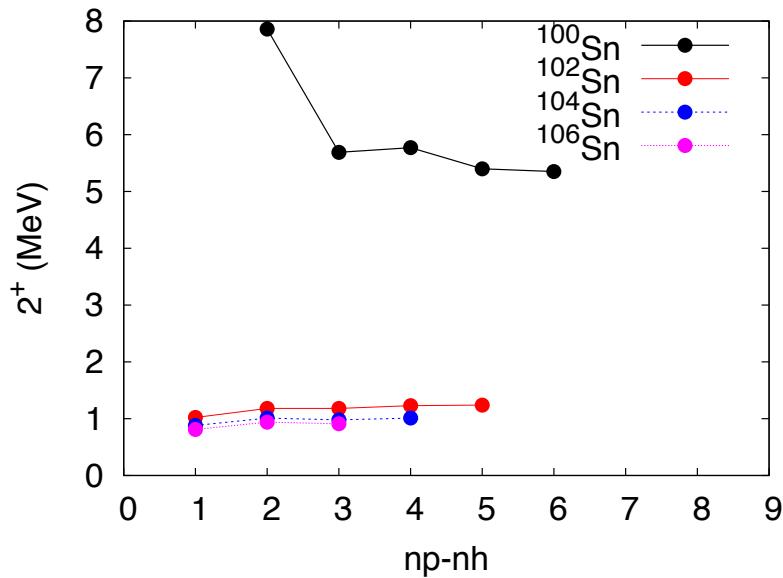
G. Guastalla et al., Phys. Rev. Lett. 110,  
172501, 2013

P. Doornenbal et al., Phys. Rev. C90,  
063102(R), 2014

*What is the true pattern of transitions in tins ?*

# B(E2) transitions in light tins

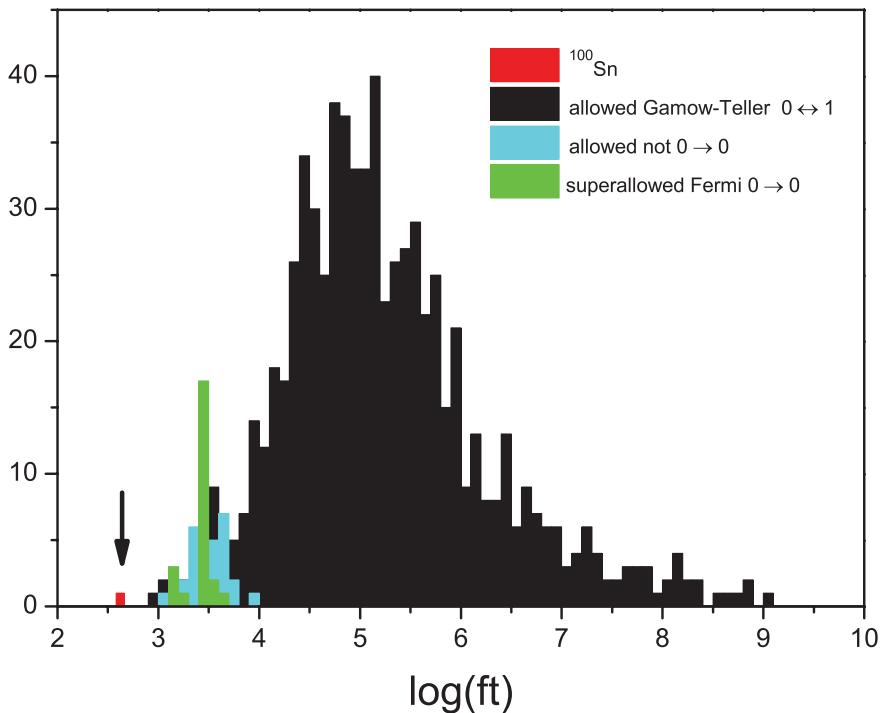
- What is sure: Core-excitations are necessary for light tins
- SM calculations are suffering from computing limitations for the lightest tins



*Where is the  $2^+$  of  $^{100}\text{Sn}$  located ?*

# Superallowed Gamow-Teller decay of $^{100}\text{Sn}$

C. Hinke et al., Nature 341, 486, 2012



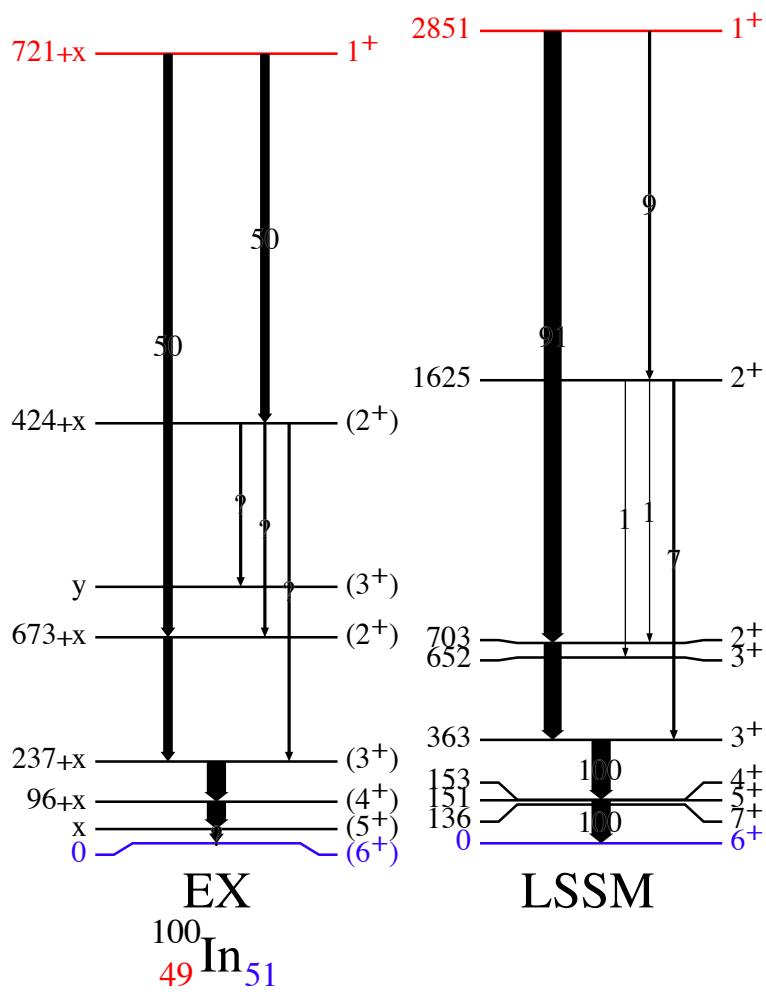
- ❑ Fastest GT decay ever observed
- ❑  $B(\text{GT})=9.1(3.0)$
- ❑  $\text{Log (ft)}=2.62(0.13)$
- ❑ Single-particle limit:  
 $B(\text{GT})=10$  (with 0.75 quenching)

$$B_{\text{GT}}^{\text{ESPM}} = \frac{4\ell}{2\ell+1} \cdot \left(1 - \frac{N_{\nu g7/2}}{8}\right) \cdot N_{\pi g9/2}$$

Experiment: RISING GSI

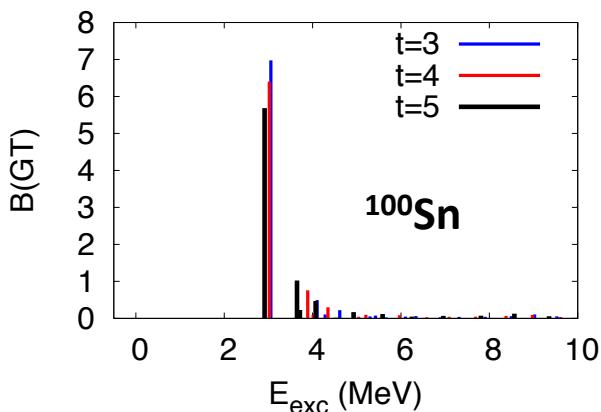
# Superallowed Gamow-Teller of $^{100}\text{Sn}$

3479 —————  $1^+$

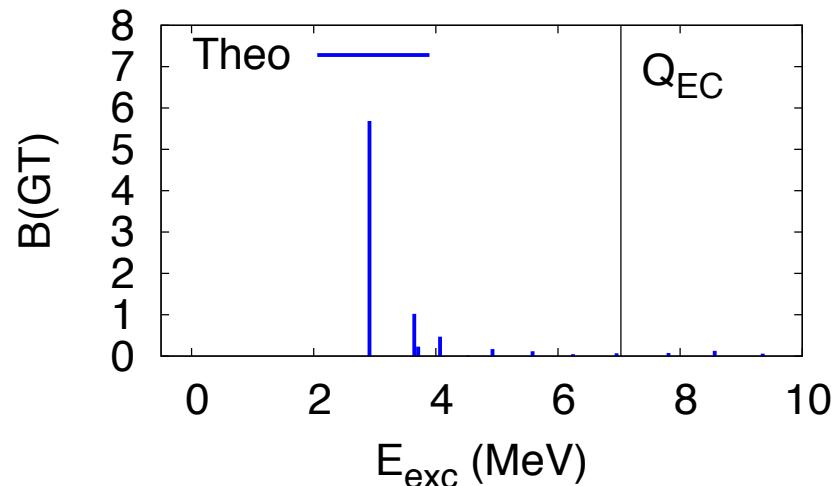


- Shell model calculations in  $sdg$  model space, up to 5p-5h excitations across the  $N=Z=50$  gaps
- Shell model strength in the Q-window:  
 $B(\text{GT})=7.82$
- To the first  $1^+$   
 $B(\text{GT})=5.68$  (69% of the strength)
- Pure wave-functions:
  - $\pi g_{9/2}^{10} \otimes \nu g_{7/2}^0$  (82%)
  - $\pi g_{9/2}^9 \otimes \nu g_{7/2}^1$  (54%)

# Superallowed Gamow-Teller of $^{100}\text{Sn}$

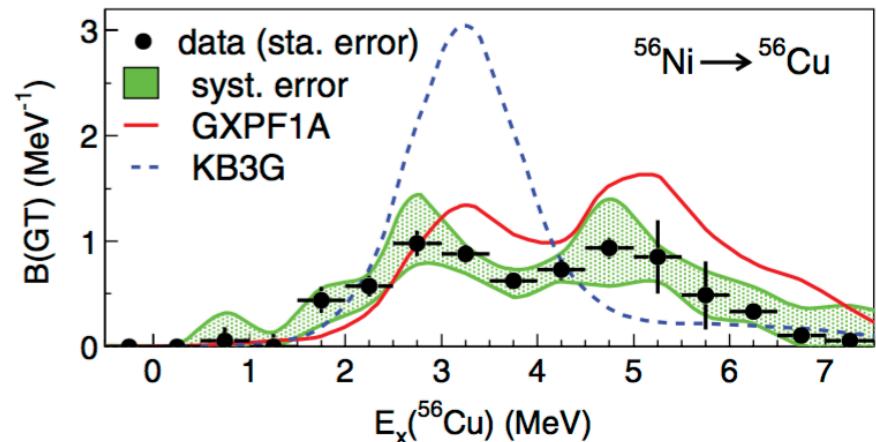


$\beta$  decay of  $^{100}\text{Sn}$



- Correlations reduce the  $B(\text{GT})$  sum rule

	$^{94}\text{Ru}$	$^{96}\text{Pd}$	$^{98}\text{Cd}$	$^{102}\text{Sn}$
0p-0h	7.11	10.66	14.22	16.02
2p-2h	4.36	7.68	11.30	12.15
4p-4h	3.44	6.64	10.36	10.95
5p-5h	2.21	5.44	10.18	
6p-6h	1.77			



- 82% of closed shell in  $^{100}\text{Sn}$  predicted in SM vs 50-60% in  $^{56}\text{Ni}$

# Superallowed Gamow-Teller of $^{100}\text{Sn}$

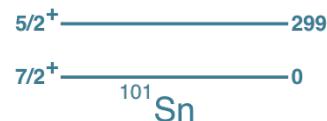
- The gaps N=Z=50 were constrained from earlier studies (e.g.  $^{98}\text{Cd}$ )
- SM reproduced well  $^{100}\text{In}$  structure
- GT appears sensitive to the  $d_{5/2}g_{7/2}$ - the rest of the shell splitting

*What is the single-particle structure of  $^{101}\text{Sn}$  ?*

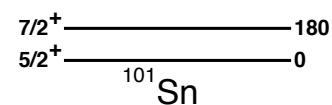
B(GT)=5.01



B(GT)=5.02



B(GT)=5.68



# Collectivity in light Xe isotopes

PRL 99, 022501 (2007)

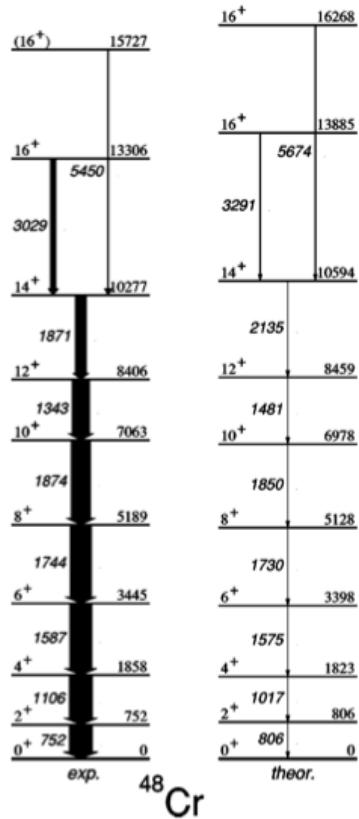
PHYSICAL REVIEW LETTERS

week ending  
13 JULY 2007



## Identification of Excited States in the $T_z = 1$ Nucleus $^{110}\text{Xe}$ : Evidence for Enhanced Collectivity near the $N = Z = 50$ Double Shell Closure

M. Sandzelius,<sup>1</sup> B. Hadinia,<sup>1</sup> B. Cederwall,<sup>1,\*</sup> K. Andgren,<sup>1</sup> E. Ganioglu,<sup>2</sup> I. G. Darby,<sup>3</sup> M. R. Dimmock,<sup>3</sup> S. Eeckhaudt,<sup>4</sup> T. Grahn,<sup>4,†</sup> P. T. Greenlees,<sup>4</sup> E. Ideguchi,<sup>5</sup> P. M. Jones,<sup>4</sup> D. T. Joss,<sup>3</sup> R. Julin,<sup>4</sup> S. Juutinen,<sup>4</sup> A. Khaplanov,<sup>1</sup> M. Leino,<sup>4</sup> L. Nelson,<sup>3</sup> M. Niikura,<sup>5</sup> M. Nyman,<sup>4</sup> R. D. Page,<sup>3</sup> J. Pakarinen,<sup>4,†</sup> E. S. Paul,<sup>3</sup> M. Petri,<sup>3</sup> P. Rahkila,<sup>4</sup> J. Sarén,<sup>4</sup> C. Scholey,<sup>4</sup> J. Sorri,<sup>4</sup> J. Uusitalo,<sup>4</sup> R. Wadsworth,<sup>6</sup> and R. Wyss<sup>1</sup>



Known rotors in the proximity  
of the closed shells:  
*sd*-shell  $^{24}\text{Mg}$   $N=Z=4$   
*pf*-shell  $^{48}\text{Cr}$   $N=Z=4$   
*sdg*-shell  $^{108}\text{Xe}$   $N=Z=4$  ?

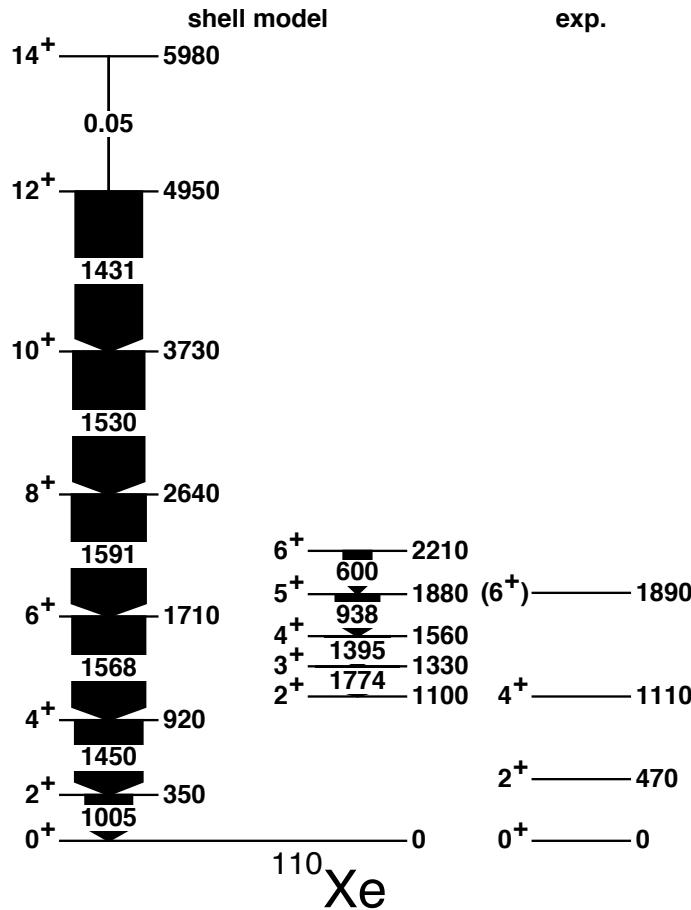
The proposed mechanism  
behind the light cases was a  
realizations of the approximate  
 $SU(3)$  symmetry scheme

« This unusual  
feature is  
suggested to be  
an effect of  
enhanced  
collectivity,  
possibly arising  
from isoscalar  $n$ - $p$   
interactions  
becoming  
increasingly  
important close to  
the  $N = Z$  line »

# Collectivity in light Xenon isotopes

E. Caurier, F. Nowacki, A. Poves and K. Sieja, Phys. Rev. C82, 064304, 2010

- Shell Model calculations in  $g7d5s1d3h11$  model space with GCN5082 interaction



$J$	$E^*$	$E_\gamma$	$B(E2)$	$Q_{\text{spec}}$	$Q_0$ from $B(E2)$	$Q_0$ from $Q_{\text{spec}}$	$\beta$
2 <sup>+</sup>	0.35	0.35	1005	-62	225	217	0.17
4 <sup>+</sup>	0.92	0.57	1450	-78	226	215	0.17
6 <sup>+</sup>	1.71	0.79	1568	-83	224	208	0.17
8 <sup>+</sup>	2.64	0.93	1591	-87	220	207	0.17
10 <sup>+</sup>	3.73	1.09	1530	-86	213	198	0.17
12 <sup>+</sup>	4.95	1.22	1431	-85	204	191	0.16
14 <sup>+</sup>	5.98	1.03	0.05	-126			
16 <sup>+</sup>	6.63	0.65	111	-125			
18 <sup>+</sup>	7.51	0.88	1184	-130			
20 <sup>+</sup>	8.51	1.00	1043	-134			
<b>110Xe</b>							
$J$	$E^*$	$E_\gamma$	$B(E2)$	$Q_{\text{spec}}$	$Q_0$ from $B(E2)$	$Q_0$ from $Q_{\text{spec}}$	$\beta$
2 <sub>2</sub> <sup>+</sup>	1.10			+61		214	0.17
3 <sup>+</sup>	1.33	0.23	1774	-1.3			
4 <sub>2</sub> <sup>+</sup>	1.56	0.23	1395	-38	219	261	0.18
5 <sup>+</sup>	1.88	0.32	938	-54	217	234	0.17
6 <sub>2</sub> <sup>+</sup>	2.21	0.33	600	-74	209	259	0.18

# Collectivity in light Xe isotopes

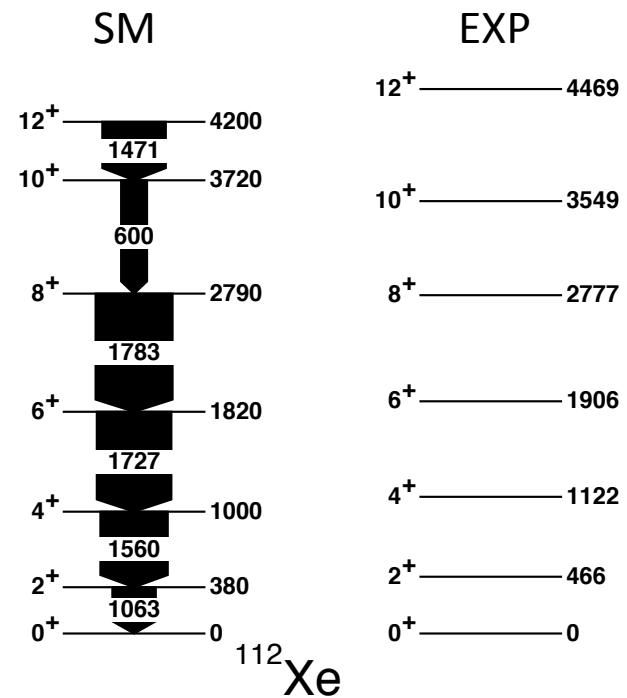
## Large-scale SM calculations:

- ❑  $^{112}\text{Xe}$ : dim. 9.324.751.339
- in the  $g7d5s1d3h11$  model space
- ❑ 2p-2h in  $^{110}\text{Xe}$  5.10<sup>9</sup>
- ❑ 2p-2h in  $^{112}\text{Xe}$  17.888.383.122  
(out of reach)

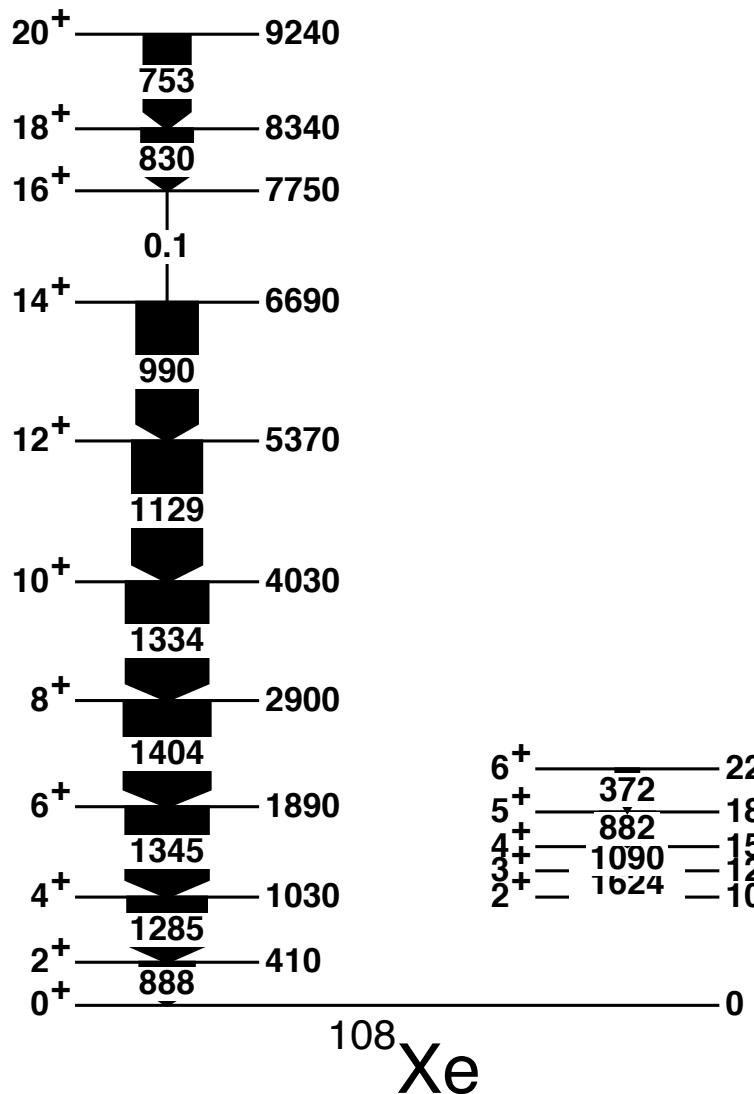
$J$	$E^*$	$E_{\text{exp}}$	$B(E2)$	$Q_{\text{spec}}$	$Q_0$ from $B(E2)$	$Q_0$ from $Q_{\text{spec}}$	$\beta$
2 <sup>+</sup>	0.38	0.46	1063	-62	217	231	0.17
4 <sup>+</sup>	1.00	1.12	1560	-75	206	236	0.17
6 <sup>+</sup>	1.82	1.91	1727	-76	190	236	0.17
8 <sup>+</sup>	2.79	2.78	1783	-74	176	232	0.17
10 <sup>+</sup>	3.72	3.55	600	-97			
12 <sup>+</sup>	4.20	4.47	1471	-118			
<b><math>^{112}\text{Xe}</math></b>							

Is  $^{100}\text{Sn}$  a good closed core?

	Number of promoted particles	Effect on $Q_2$
$^{108}\text{Xe}$		
1p-1h	0.25	15%
2p-2h	0.5	19%
$^{110}\text{Xe}$		
1p-1h	0.22	17%
2p-2h	0.5	21%
$^{112}\text{Xe}$		
1p-1h	0.25	20%



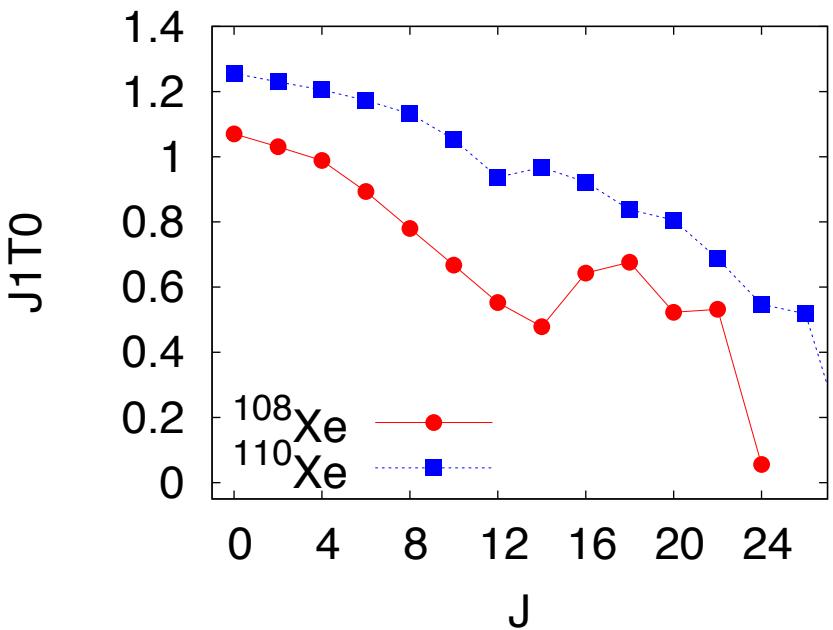
# Collectivity in Xe isotopes: the N=Z case



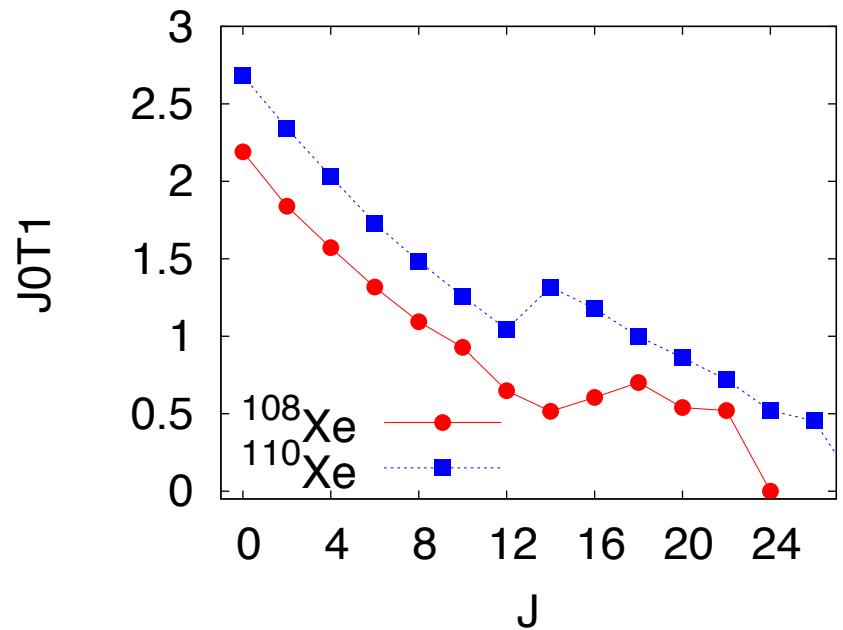
$J$	$E^*$	$E_\gamma$	$B(E2)$	$Q_{\text{spec}}$	$Q_0$ from $B(E2)$	$Q_0$ from $Q_{\text{spec}}$	$\beta$
$2^+$	0.41	0.41	888	-57	200	211	0.16
$4^+$	1.03	0.62	1285	-71	195	210	0.16
$6^+$	1.89	0.86	1345	-65	163	208	0.16
$8^+$	2.90	1.01	1404	-64	154	206	0.16
$10^+$	4.03	1.13	1334	-67	160	198	0.15
$12^+$	5.37	1.34	1129	-71	175	182	0.15
$14^+$	6.69	1.32	990	-79	176	168	0.14
$16^+$	7.75	1.06	0.1	-137			
$18^+$	8.34	0.59	830	-140			
$20^+$	9.24	0.90	753	-143			

$J$	$E^*$	$E_\gamma$	$B(E2)$	$Q_{\text{spec}}$	$Q_0$ from $Q_{\text{spec}}$	$\beta$
$2_2^+$	1.03			+59	196	0.16
$3^+$	1.28	0.25	1624	-1.3		
$4_2^+$	1.51	0.23	1090	-38	265	0.18
$5^+$	1.84	0.33	882	-51	220	0.17
$6_2^+$	2.25	0.41	372	-83	290	0.19

# T=0 vs T=1 pairing in Xe isotopes



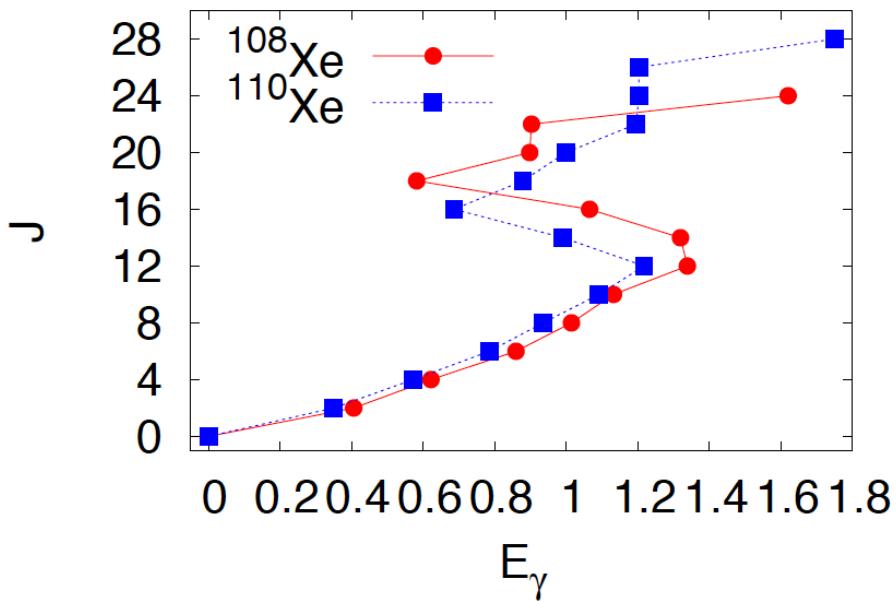
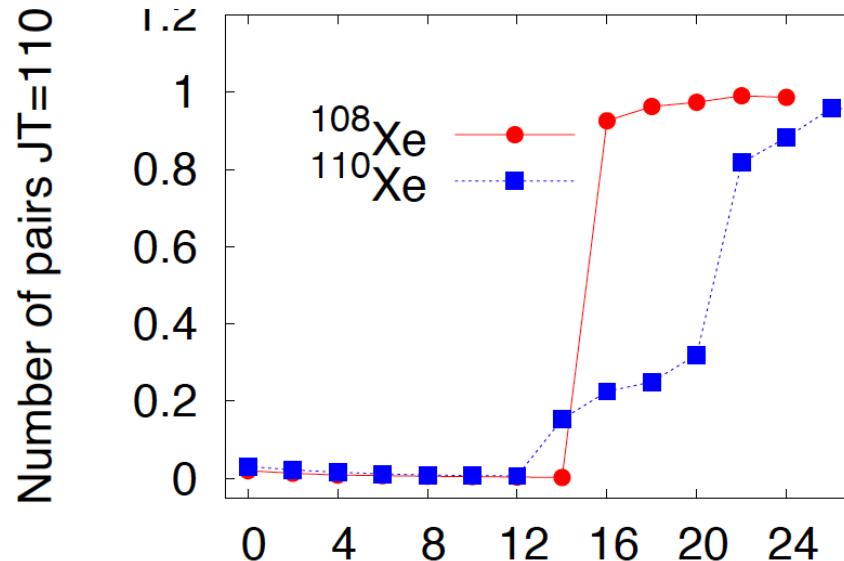
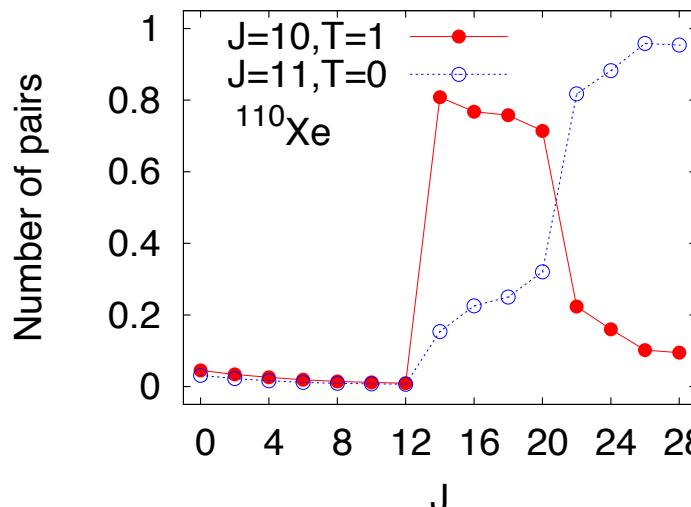
Number of p-n pairs ( $J=1$ ,  $T=0$ )



Number of p-p, p-n, n-n pairs ( $J=0$ ,  $T=1$ )

- No evidence for g.s. condensates
- Pair content decreasing with spin

# Backbending and pair alignment



- Backbend takes place at  $J=14$  in both nuclei
- Two-step process in  $^{110}\text{Xe}$

## Conclusions

New, more precise data are welcome to revive the theoretical progress and interest in the  $^{100}\text{Sn}$  region. The questions remain to be addressed experimentally:

- What is the true pattern of transitions in tins ?*
- Where is the  $2^+$  of  $^{100}\text{Sn}$  located ?*
- What is the single-particle structure of  $^{101}\text{Sn}$  ?*