

Few atoms in atomic waveguides





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Saclay - 03 February 2017



• Introduction:

Few-body problems: a common interest in nuclear and ultracold physics

• Context:

- Atomic gases near s-wave magnetic Feshbach resonances
- Three-body resonant problem for bosons \rightarrow Efimov physics
- Atomic gas in low dimensions
 - \rightarrow the 1D Bose gas and the Lieb Liniger model

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• 2- and 3-body problem in 1D atomic waveguides

- Two-channel model
- Confined induced dimers
- Some results of the Skorniakov Ter Martirosian equation
- Summary & perspectives

Deep links between nuclear physics and ultracold atoms

Two different energy scales ($\sim 10^{-13}~\text{ev}\equiv 1~\text{nK}\to\sim 10~\text{Mev}$) ... but a lot of things to share:

- common interest on the few-body problem and the limit toward many-body states (ex: few n \rightarrow neutron star)
- importance of resonant phenomena
- search for exotic quantum states (ex: 4n, Efimov ...)



Example : contact models & universality in few-body physics

Introduction

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Some key features of ultracold-atoms not present in nuclear physics

- Tunable effective interactions
- Tunable dimensionality 3D \leftrightarrow 2D, 3D \leftrightarrow 1D
- Dilute limit $nb^3 \ll 1$ & possible large correlations $na^3 \gtrsim 1$ (unitary limit)
- Two-body microscopic potentials well known and sufficient for quantitative calculations

Conclusi

Few-body properties : at the heart of ultracold gases experiments

The *N*-body ground state at T=0K is a solid



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 $\iff \mathsf{many} \ \mathsf{deep} \ \mathsf{bound} \ \mathsf{states} \ \mathsf{for} \ 2, 3, \dots \textit{N}\text{-}\mathsf{atoms}$ Three-body recombination processes



• Bose gas : large losses for $na^3 > 1$

 \implies No thermodynamical equilibrium at resonance!

- Two-component Fermi gas : stability at resonance ! (Petrov-Shlyapnikov-Salomon 2004)
- \Rightarrow Experimental study of the BEC-BCS crossover & unitary Fermi gas

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Order of magnitude in ultracold physics

- Typical trap size: few μ m
- Number of atoms: 10⁵
- Atomic density: $10^{13} \lesssim n \lesssim 10^{15}$ atoms/cm³
- Temperature: 1 nK \lesssim T \lesssim μ K
- Range of the interaction: $b \sim R_{\rm vdW} = \frac{1}{2} \left(\frac{mC_6}{\hbar^2} \right)^{1/4} \sim 10$ nm

Scale separation:

$$\frac{\hbar^2 n^{2/3}}{m}, k_B T, \mu \ldots \ll \frac{\hbar^2}{m h^2}$$

• "Low" energy: collective modes, bound states near scattering resonances . . .

"High" energy: usual molecules, clusters ...

-channel mod

STM equation

Results

Conclusions

Resonance and Feshbach mechanism

$$\begin{array}{c} & \bigoplus \\ & \bigoplus \\ & \bigoplus \\ & \bullet \end{array} \end{array}^{+ \text{ external magnetic field } \mathcal{B} \\ & \mathcal{E}_{\text{mol}} \sim \mathcal{E}_{\text{mol}}^{0} + \delta \mathcal{M} \times \mathcal{B} \quad \Longrightarrow \quad a = a_{\text{bg}} \left(1 - \frac{\Delta \mathcal{B}}{\mathcal{B} - \mathcal{B}_{0}} \right) \end{array}$$

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Conclusion

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s-wave scattering length near the resonance



- Feshbach resonance at : $B = B_0$
- Shape resonance: $|a_{
 m bg}| \gg R_{
 m vdW}$
- Away from a shape resonance: $a_{\mathrm{bg}} = \mathcal{O}(R_{\mathrm{vdW}})$

Efimov spectrum



Discrete scaling symmetry $E(M') = \lambda^2 \times E(M)$ where $\lambda = \exp(\pi/|s|)$.

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Evidence of the Efimov effect in ultracold atoms Studies in the vicinity of magnetic Feshbach resonances Three-bosons

- Group at Innsbrück university: experiments using ¹³³Cs
- Group at Firenze university: experiments using $^{\rm 39}{\rm K}$
- Group at Bar-Ilan University: experiments using ⁷Li
- Group at Rice university: experiments using ⁷Li

Three-spin-component fermions ⁶Li

- Group of the university of Penn-State
- Group of the university of Heidelberg
- Group of the university of Tokyo

Hetero-nuclear systems

Group of Firenze : KKRb and KRbRb Efimov trimers.

See review : F. Ferlaino, R. Grimm, Physics 3, 9 (2010)

Introduction

STM equation

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Atoms in Low dimensions

Possible to tune the external potential

 \implies studies of atoms in quasi-1D or quasi-2D geometries



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• Atoms in a 2D isotropic harmonic trap



 \longrightarrow 1D atomic wave guide $E_{1-\text{particle}} = \frac{\hbar^2 k^2}{2m} + \hbar \omega_{\perp} (2n + |m| + 1)$

Low energy monomode regime

Characteristic energies : $k_B T$, $\frac{\hbar^2}{m} \rho^{2/3} \cdots \ll \hbar \omega_{\perp}$

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 $\langle z|\psi_{k_0}
angle \sim e^{ik_0z} + f_{\rm 1D}(k_0)e^{ik_0|z|} \qquad f_{\rm 1D}(k_0) = rac{-1}{1+ik_{\rm 1D}a_{\rm 1D}}$

• Lieb-Liniger model : δ interaction \rightarrow reproduce exaclty f_{1D} at the 2-body level

- $a_{1D} \rightarrow 0^-$: Tonks Girardeau regime
- $a_{1D} > 0$: 1D N-body bound states \rightarrow solitonic regime:

Energy :
$$\frac{-\hbar^2 N(N^2-1)}{6ma_{1D}^2}$$

3-body problem: heart of the integrability in 1D



rij & tij : reflexion and transmission coefficients

$$r_{ij} = f_{1\mathrm{D}}\left(rac{k_i - k_j}{2}
ight)$$
 & $t_{ij} = 1 + f_{1\mathrm{D}}\left(rac{k_i - k_j}{2}
ight)$

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Purpose of the present study

Three-body bound states for bosons in a 1D waveguide

From the Efimov law $-E_0 \exp(-2n\pi/s_0)$... to the trimer of the Lieb-Liniger model $-\frac{4\hbar^2}{ma_{1D}^2}$

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Two-channel modeling of the Feshbach resonance

Atom :
$$a_{\mathbf{k}}^{\dagger}|0
angle$$
 ; Molecule : $b_{\mathbf{k}}^{\dagger}|0
angle$

$$H = \int \frac{d^{3}k}{(2\pi)^{3}} \left[E_{\mathbf{k}} a_{\mathbf{k}}^{\dagger} a_{\mathbf{k}} + \left(\frac{E_{\mathbf{k}}}{2} + E_{\mathrm{mol}} \right) b_{\mathbf{k}}^{\dagger} b_{\mathbf{k}} \right]$$

(Kinetic term & $E_{\mathrm{mol}} = E_{\mathrm{mol}}^{0} + \delta \mathcal{M} \mathcal{B}$)

$$+\frac{g}{2}\int\frac{d^{3}kd^{3}Kd^{3}k'}{(2\pi)^{9}}\langle k'|\delta_{\epsilon}\rangle\langle\delta_{\epsilon}|k\rangle a^{\dagger}_{\frac{K}{2}-\mathbf{k}'}a^{\dagger}_{\frac{K}{2}+\mathbf{k}'}a_{\frac{K}{2}+\mathbf{k}}a_{\frac{K}{2}-\mathbf{k}}.$$
(atom-atom interaction)

$$E_{\mathbf{k}} = \frac{\hbar^2 k^2}{2m} \& \langle k | \delta_{\epsilon} \rangle = \exp\left(-\frac{k^2 \epsilon^2}{4}\right) \text{ cut-off function}$$

$$\epsilon : \text{ short range parameter } \epsilon \equiv \mathcal{O}(b)$$

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Parameters obtained from the 2-body properties at low E

Transition operator: $\hat{T}(E) = \frac{|\delta_{\epsilon}\rangle\langle\delta_{\epsilon}|}{\left(g + \frac{\Lambda^2}{E - E_{\text{mol}}}\right)^{-1} - \langle\delta_{\epsilon}|\hat{G}_0^{\text{rel}}(E)|\delta_{\epsilon}\rangle},$ $\hat{G}_0^{\text{rel}}(E)$: resolvent for the relative particle

- background scattering length: $a_{\mathrm{bg}} \leftrightarrow g$
- scattering length: $a = a_{bg} \left(1 \frac{\Delta B}{B B_0} \right)$
- width parameter (Petrov 2004) :

$$R^{\star} = rac{\dot{h^2}}{m a_{
m bg} \delta \mathcal{M} \Delta \mathcal{B}} \propto rac{1}{\Lambda^2}.$$

If $R^* \gg R_{\rm vdw} \equiv$ natural cut-off for low energy properties

Confinement Induced Dimers in the 1D waveguide

- Zero-range model Olshanii (1998)
- Two-channel model:



Study of trimers in the waveguide

- $E = E_{3body} E_{Com} < 0$
- s-wave sector
- Skorniakov Ter Martirosian equation for trimers in the 1D waveguide :

$$D(E^{\mathrm{rel}})f(\underline{n},k) - 2\sum_{\underline{n}'=0}^{\infty} \int \frac{dk'}{2\pi} \langle \underline{n}, k | \mathcal{K}(E) | \underline{n}', k' \rangle f(\underline{n}',k') = 0$$

- Explore the case $a_{\perp}/\epsilon = 20$ for several resonances
- n_{\max} from 100 to 400 $\Longrightarrow (2n_{\max}+1)^3$ values of $d^j_{m,m'}(\theta)$.
- E such that the integral equation has a zero eigenvalue (typical matrix sizes \sim 30 000)

Example of a narrow resonance

 39 K at $B_0=752$ G ; $R^{\star}=36.4R_{
m vdw}$ $a_{\perp}=20\epsilon~(\omega=2\pi imes55.4$ kHz)



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Example of a narrow resonance 39 K at $B_0 = 752$ G ; $R^{\star} = 36.4 R_{vdw}$ $a_{\perp} = 20\epsilon \ (\omega = 2\pi \times 55.4 \text{ kHz})$



Introduction

A direct link with the Efimov 3-body parameter $a_{\rm 1D} \sim \frac{-a_{\perp}}{2} \left[\zeta \left(\frac{1}{2} \right) + \frac{a_{\perp}}{a_{\rm 3D}} \right].$ (Olshanii PRL-1998)

Threshold for the shallow 1D trimer on the 1D dimer line for:

$$a_{
m 1D}^{
m at-d}
ightarrow \infty \quad \Longleftrightarrow \quad a_{
m 3D}^{
m at-d}
ightarrow 0^+$$

Regime of interest for the dimer binding energy :

- Structureless 3D dimer in the waveguide $\implies |E_{\text{dimer}}| \gg \hbar\omega$
- a_{3D}^{at-d} function of the three-body parameter

$$\Rightarrow |E_{\text{dimer}}| \ll \frac{\hbar^2}{m\epsilon^2}$$

A way to measure the three-body parameter for different values ${\sf Energy}/{\sf magnetic}$ field than in 3D ?



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Example of a broad resonance

⁷Li at $B_0 = 737$ G, $R^{\star} = 1.2 R_{\rm vdw}$ $a_{\perp} = 20\epsilon \ (\omega = 2\pi \times 2.7 \times 10^6 \text{ Hz})$



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Feshbach resonance in the vicinity of a shape resonance





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Results

Comparison with a recent 2D study

2D waveguide & single-channel separable model



Jesper Levinsen, Pietro Massignan, and Meera M. Parish, Phys. Rev. X 4, 031020 (2014).

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Summary & perspectives

• Model including the Feshbach coupling for the 3-body problem in wave guides



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Summary & perspectives

- Model including the Feshbach coupling for the 3-body problem in wave guides
- Regime where a_{1D} is large and positive:
- 1D Lieb-Liniger trimer continuously connected with the dimer branche

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other quasi-degenerate trimers exists continuously connected with Efimov trimers for large a_{\perp}/ϵ

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 - other quasi-degenerate trimers exists continuously connected with Efimov trimers for large a_{\perp}/ϵ
 - Trimer threshold at the Atom-dimer spectrum Systematic shift from $a_{3D}^{\mathrm{at-dim}} \to \infty$ to $a_{3D}^{\mathrm{at-dim}} \to 0$ for the deepest trimer branches & large a_{\perp}/ϵ

Summary & perspectives

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- Near future :
 - Explore larger aspect ratio

 \Longleftrightarrow More insights in the confined induced trimers

• 3-body collisions in 1D waveguides

⇒ violation of integrability