

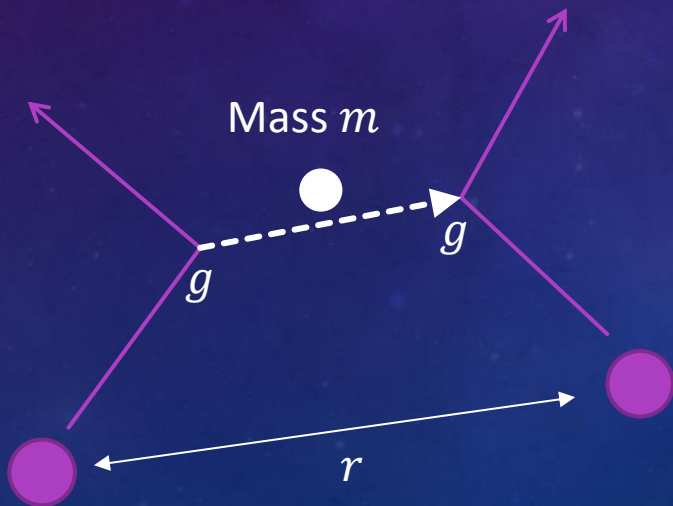
INTERACTING BOSE POLARONS

FROM THE YUKAWA TO THE EFIMOV ATTRACTION

Pascal Naidon, RIKEN

THE YUKAWA POTENTIAL

- In the 1930s, Hideki Yukawa showed that the exchange of a massive boson field between two particles induces a Coulomb potential exponentially screened by the boson mass.



Mediated potential

$$V(r) = -g^2 \frac{e^{-kmr}}{r}$$

THE YUKAWA POTENTIAL

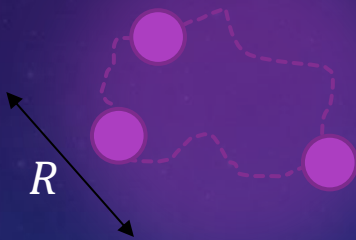
- Coulomb potential of electromagnetism: exchange of photon ($m = 0$) between charged particles
- Tail of the nuclear force: exchange of mesons ($m \geq 135$ MeV) between nucleons

THE EFIMOV POTENTIAL

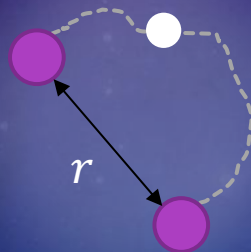
- In the 1970s, Vitaly Efimov showed that a universal three-body force arises between three resonantly-interacting particles.



3 particles of mass M



2 particles of mass M and a boson of mass m



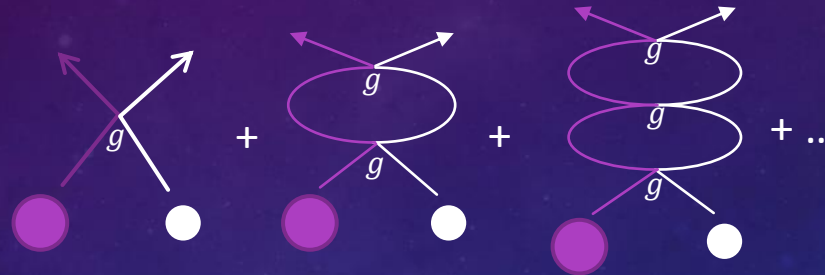
$$V(R) = -\frac{\hbar^2}{2M} \frac{1.006^2}{R^2}$$

$$V(r) = -\frac{\hbar^2}{2m} \frac{0.414^2}{r^2}$$

May also be viewed as a two-body mediated potential

RESONANT INTERACTION

Diagrammatic point of view



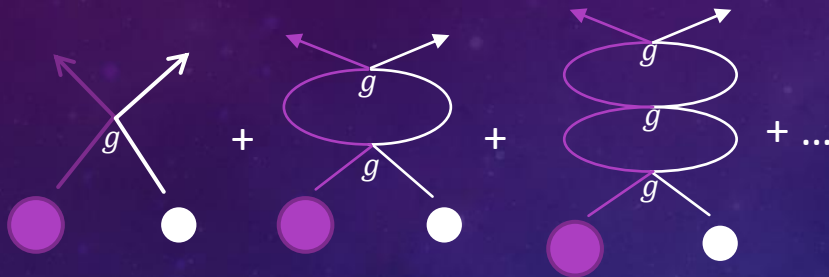
$$\underbrace{\frac{1}{g}}_{< 0} + \underbrace{\sum_k^\Lambda \frac{1}{\frac{\hbar^2 k^2}{2\mu}}}_{> 0} = \frac{2\mu}{4\pi\hbar^2} \frac{1}{a}$$

Scattering length a

$$\text{Reduced mass } \mu = \left(\frac{1}{m} + \frac{1}{M}\right)^{-1}$$

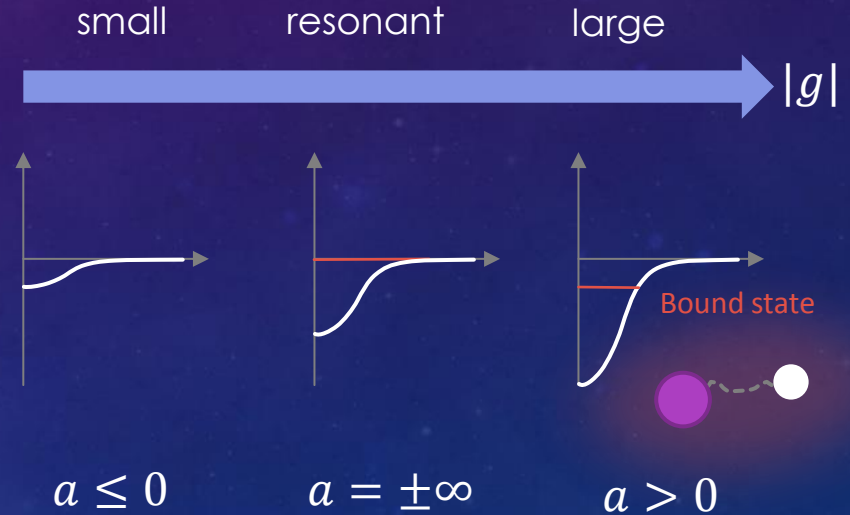
RESONANT INTERACTION

Diagrammatic point of view



$$\underbrace{\frac{1}{g}}_{< 0} + \underbrace{\sum_k^{\Lambda} \frac{1}{\frac{\hbar^2 k^2}{2\mu}}}_{> 0} = \frac{2\mu}{4\pi\hbar^2} \frac{1}{a}$$

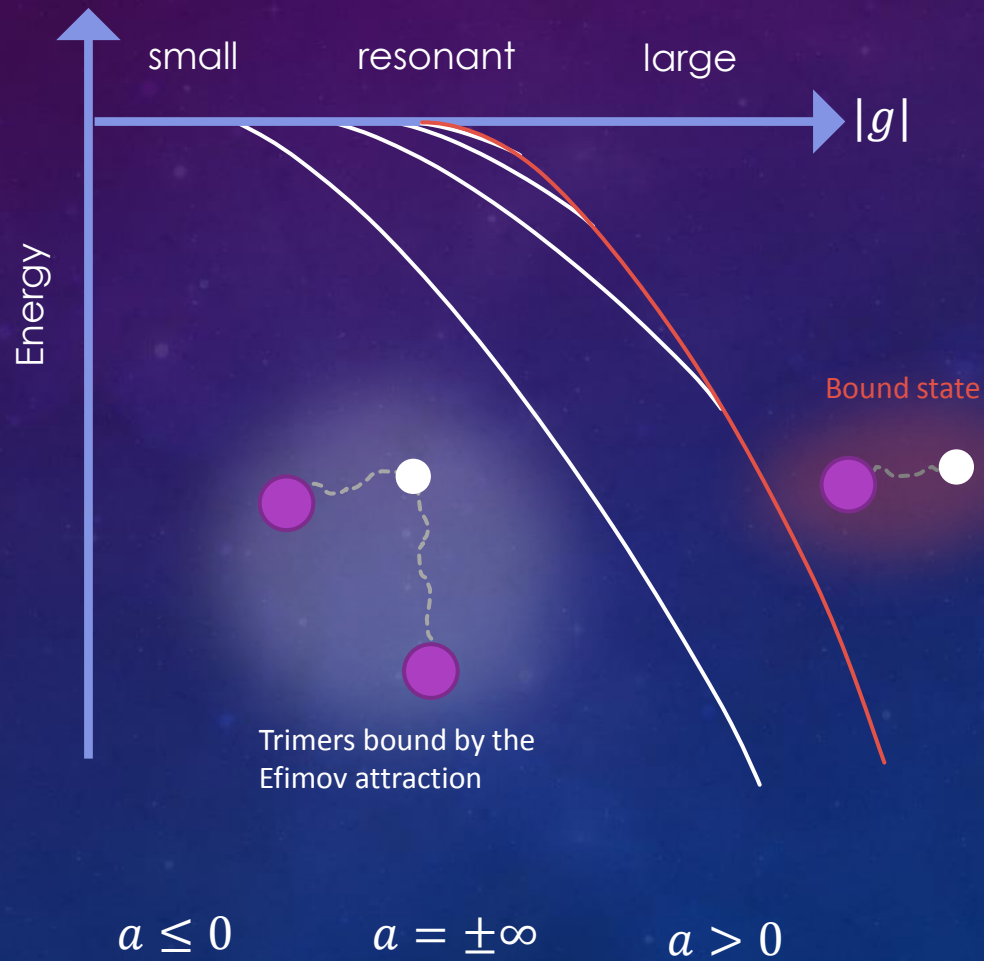
“Potential” point of view



Scattering length a

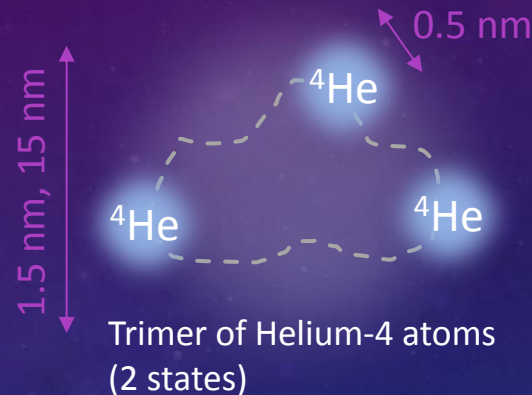
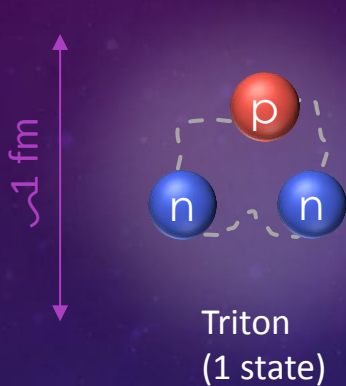
$$\text{Reduced mass } \mu = \left(\frac{1}{m} + \frac{1}{M}\right)^{-1}$$

RESONANT INTERACTION



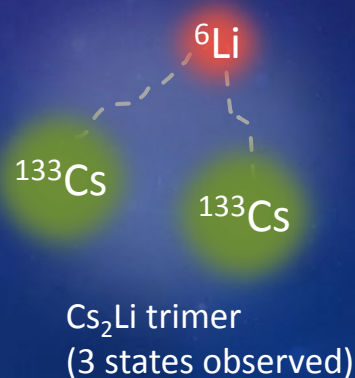
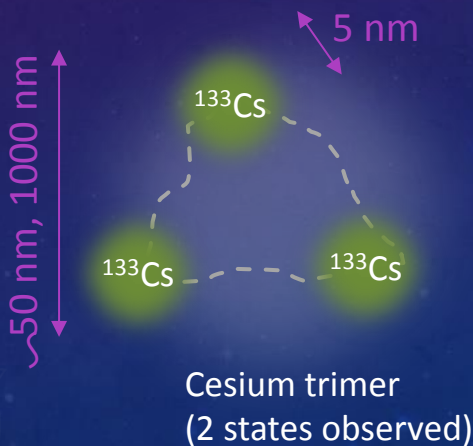
SYSTEMS EXHIBITING THE EFIMOV ATTRACTION

- In Nature : The value of g happens to be close to resonance



Recently observed
Science, 348, 551 (2015)

- Ultra-cold atoms : The value of g can be adjusted by applying a magnetic field



Recently observed
PRL 112, 250404 (2014)
PRL 113, 240402 (2014)

SYSTEMS EXHIBITING THE EFIMOV ATTRACTION

arXiv:1610.09805

Efimov Physics: a review

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November 1, 2016

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Abstract

This article reviews theoretical and experimental advances in Efimov physics, an array of quantum few-body and many-body phenomena arising for particles interacting via short-range resonant interactions, that is based on the appearance of a scale-invariant three-body attraction theoretically discovered by Vitaly Efimov in 1970. This three-body effect was originally proposed to explain the binding of nuclei such as the triton and the Hoyle state of carbon-12, and later considered as a simple explanation for the existence of some halo nuclei. It was subsequently evidenced in trapped ultra-cold atomic clouds and in diffracted molecular beams of gaseous helium. These experiments revealed that the previously undetermined three-body parameter introduced in the Efimov theory to stabilise the three-body attraction typically scales with the range of atomic interactions. The few- and many-body consequences of the Efimov attraction have been since investigated theoretically, and are expected to be observed in a broader spectrum of physical systems.

Contents

I Introduction

1 What is Efimov physics?

2 Why is it important? For which systems?

3 A short history of Efimov physics

4

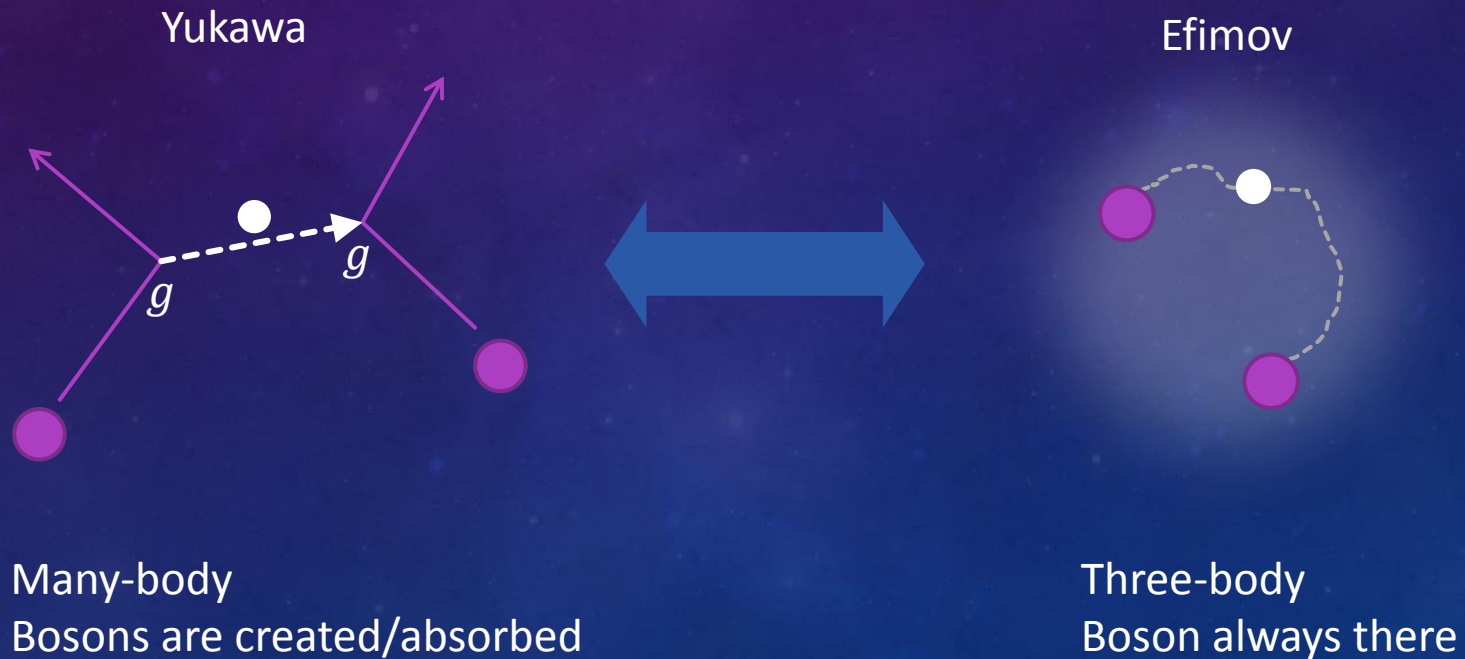
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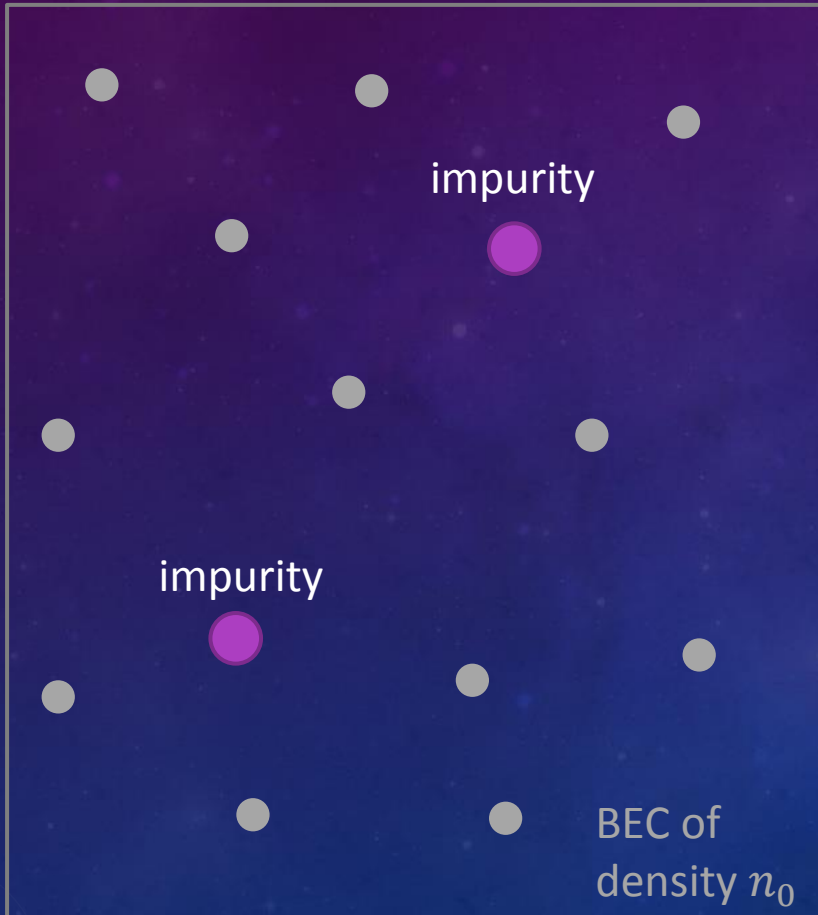
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MOTIVATION

Connection between Yukawa and Efimov mediated potentials?



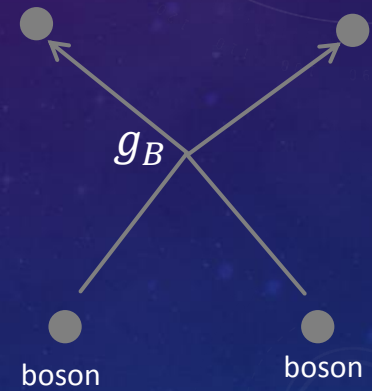
IMPURITIES IN A BOSE-EINSTEIN CONDENSATE



Interactions



$g < 0$ can be large
(attraction)

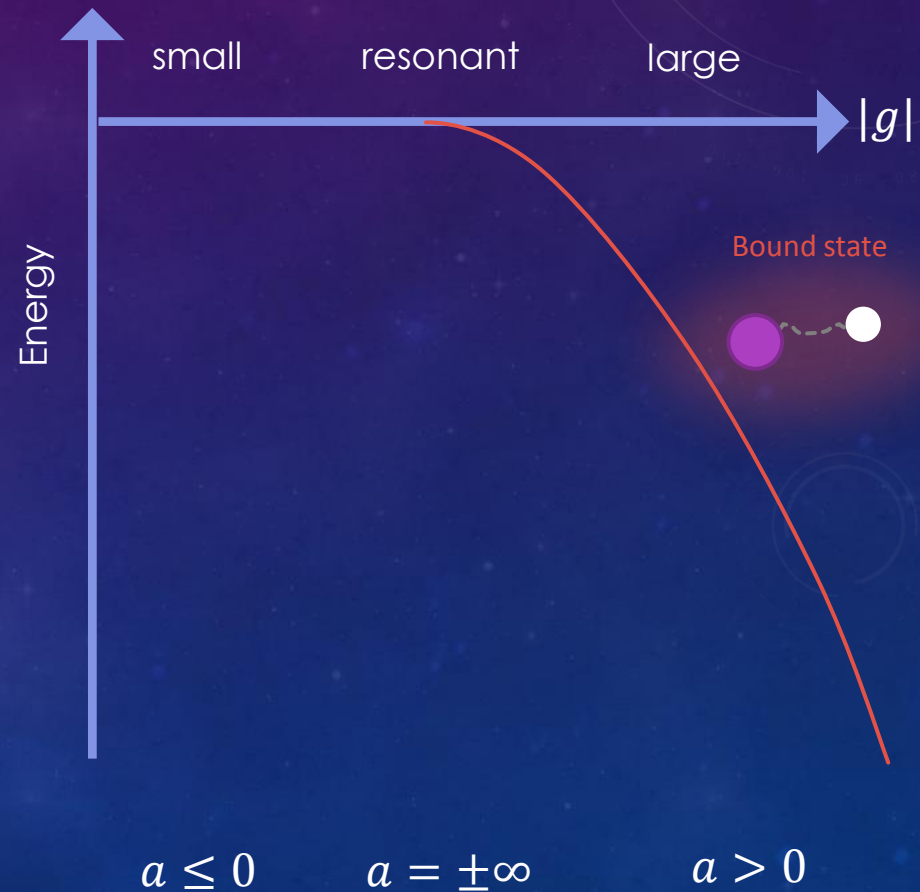
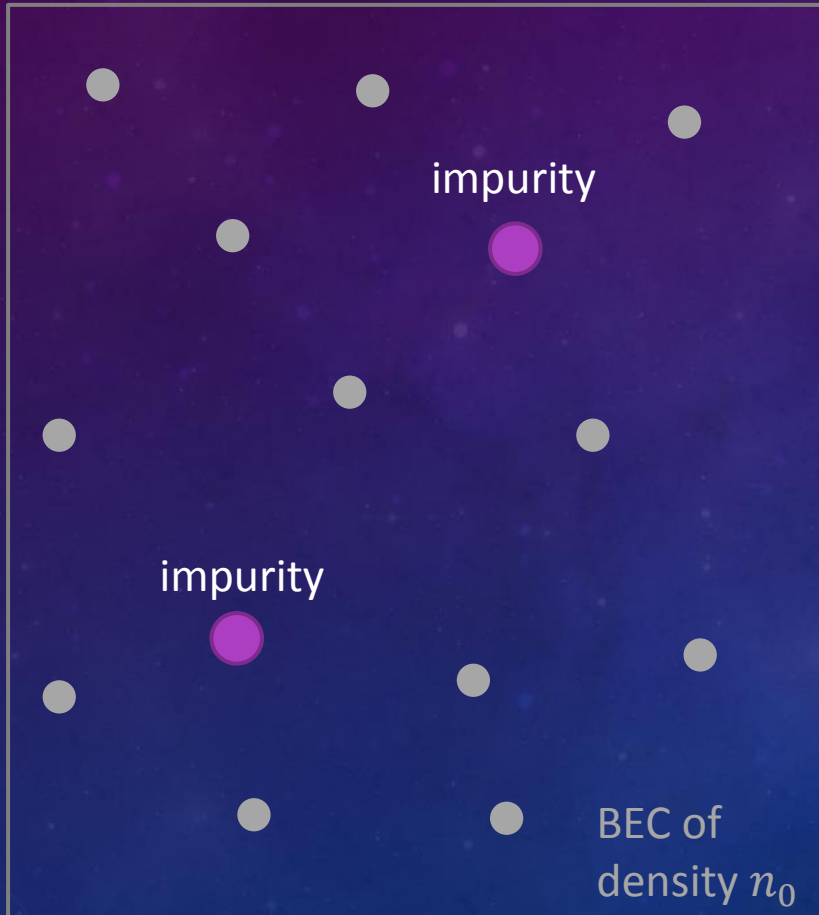


$g_B > 0$ is small
(weak repulsion)

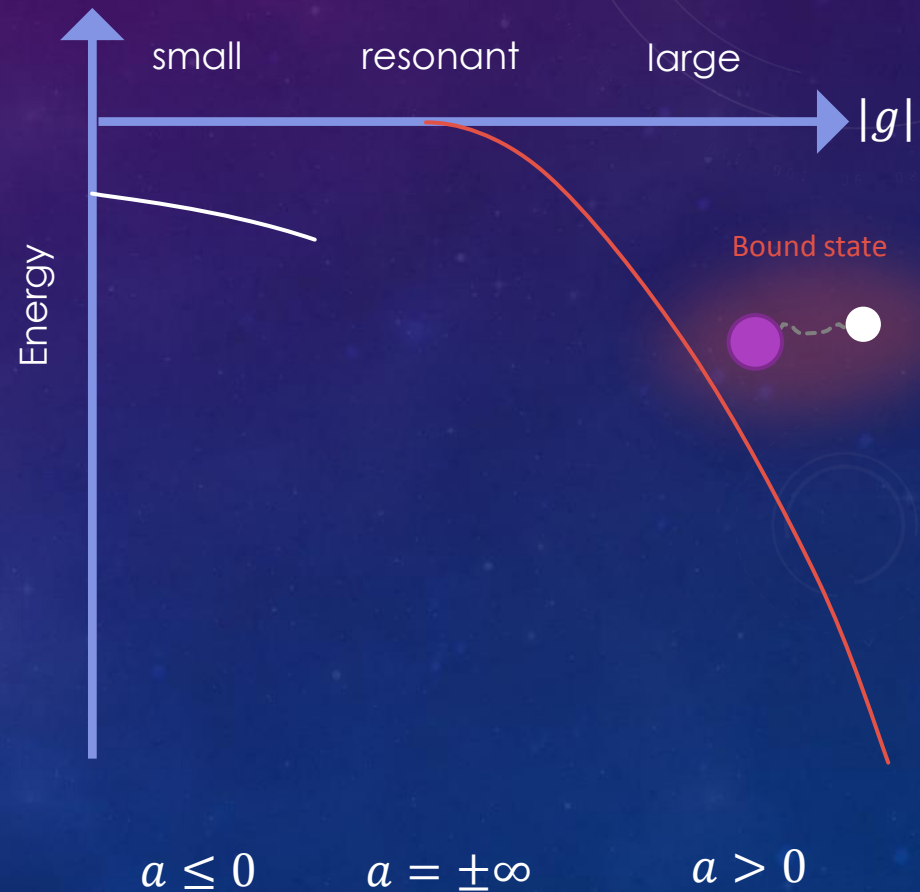
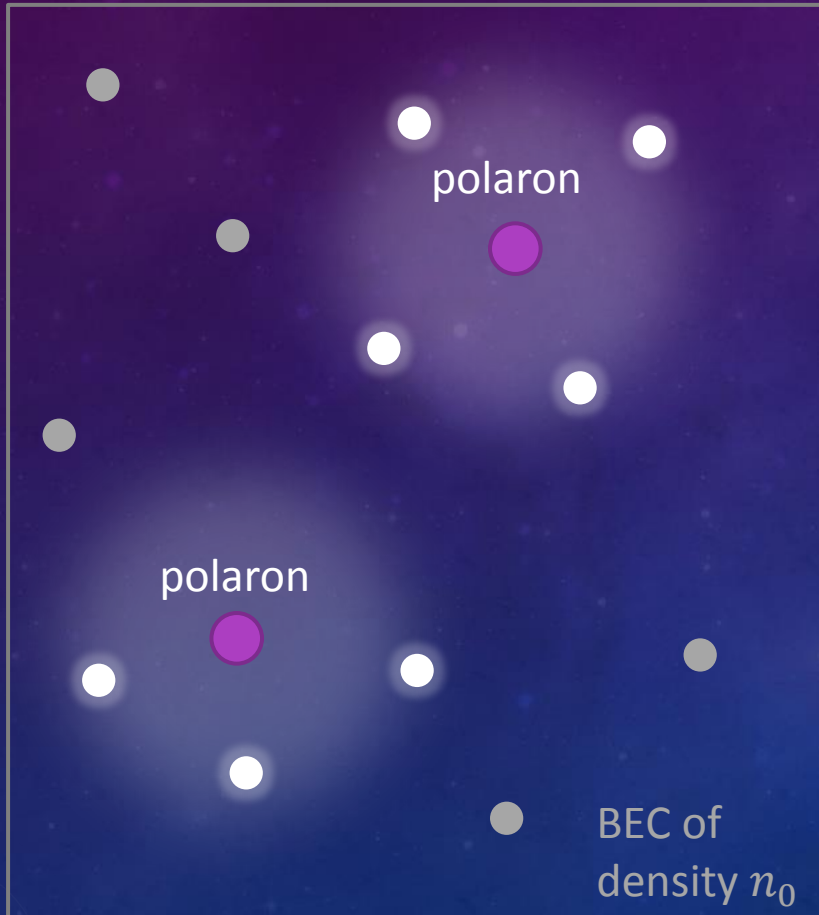
$$n_0 a_B^3 \ll 1$$

Neglect direct interactions between impurities

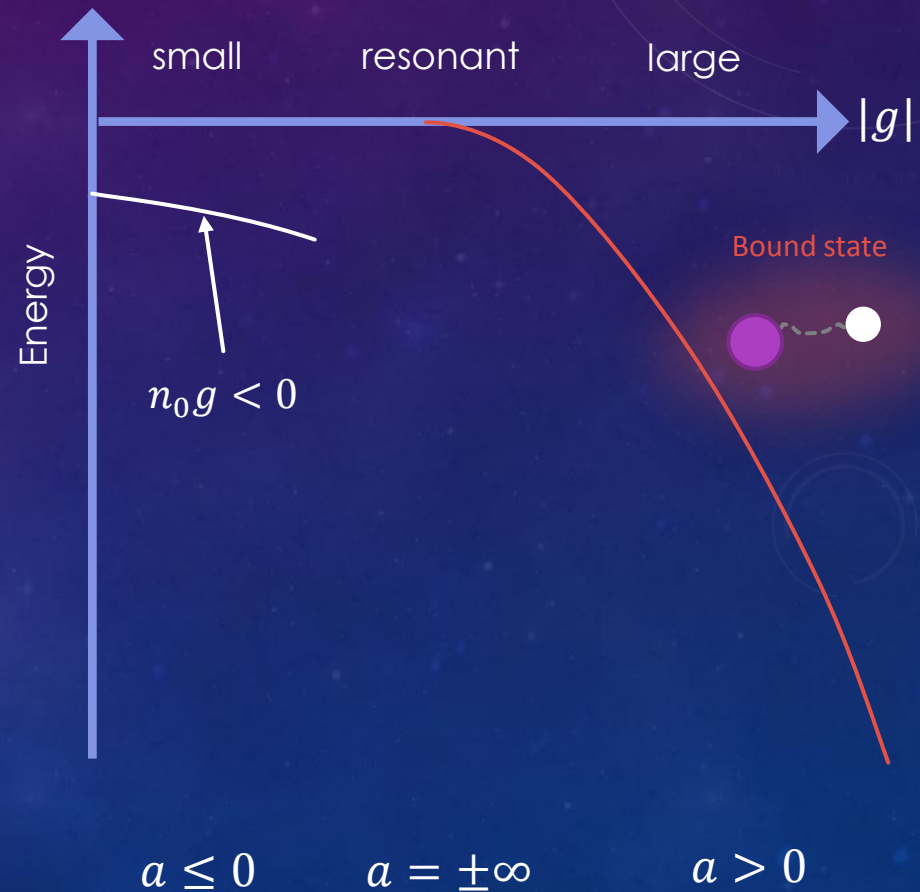
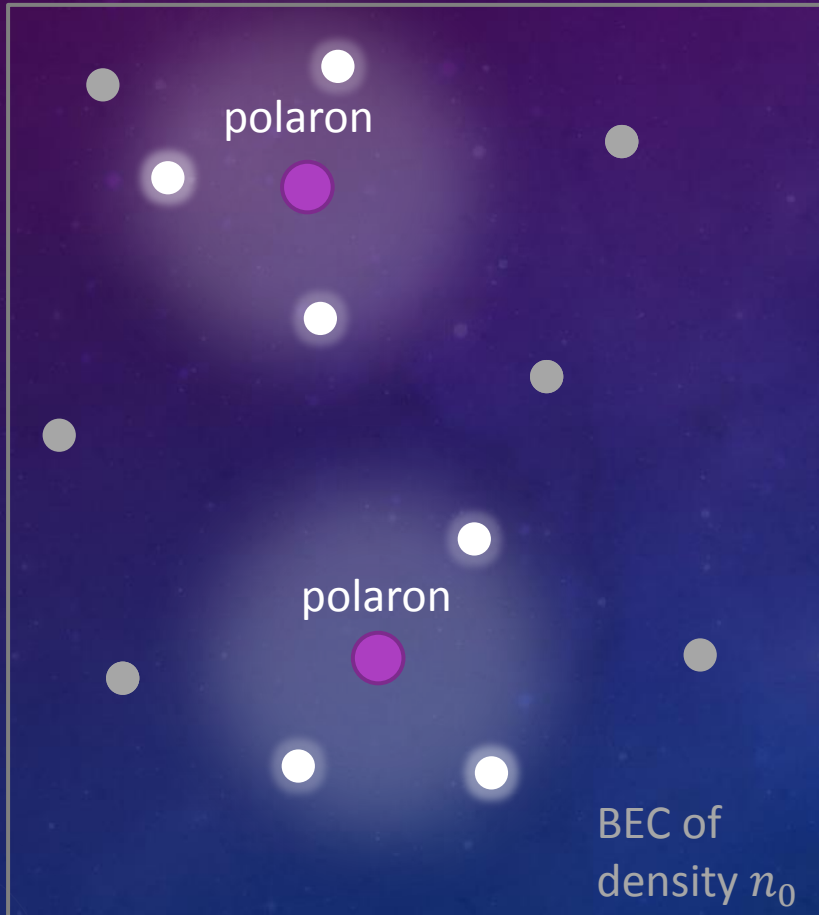
IMPURITIES IN A BOSE-EINSTEIN CONDENSATE



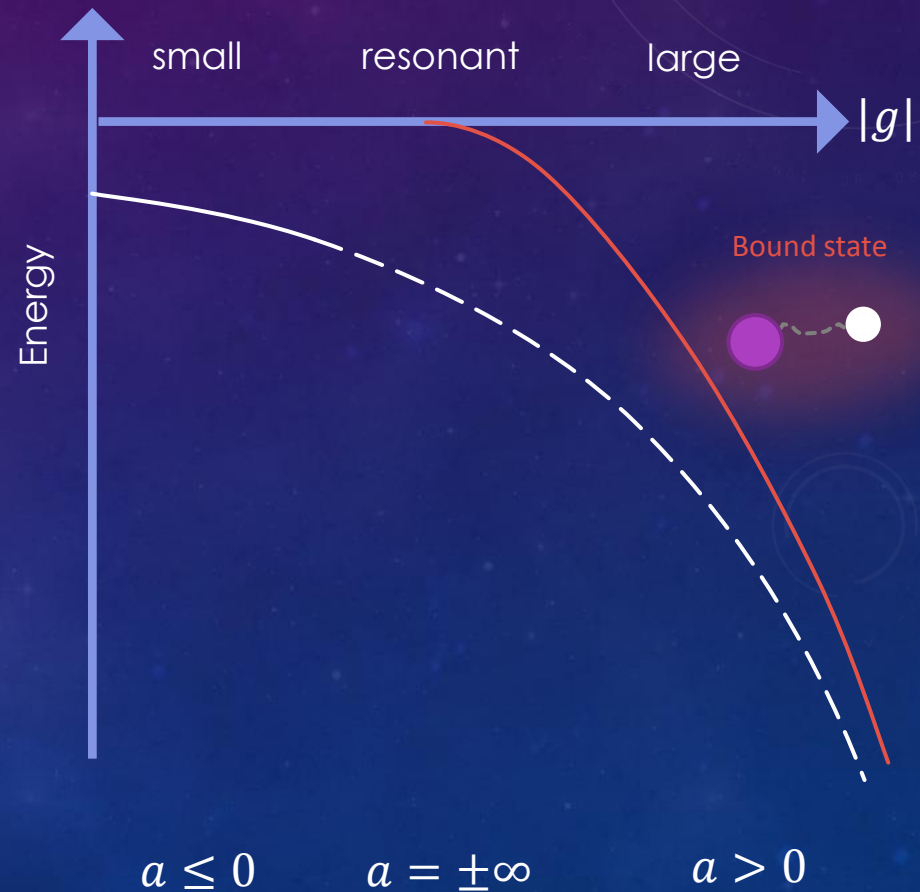
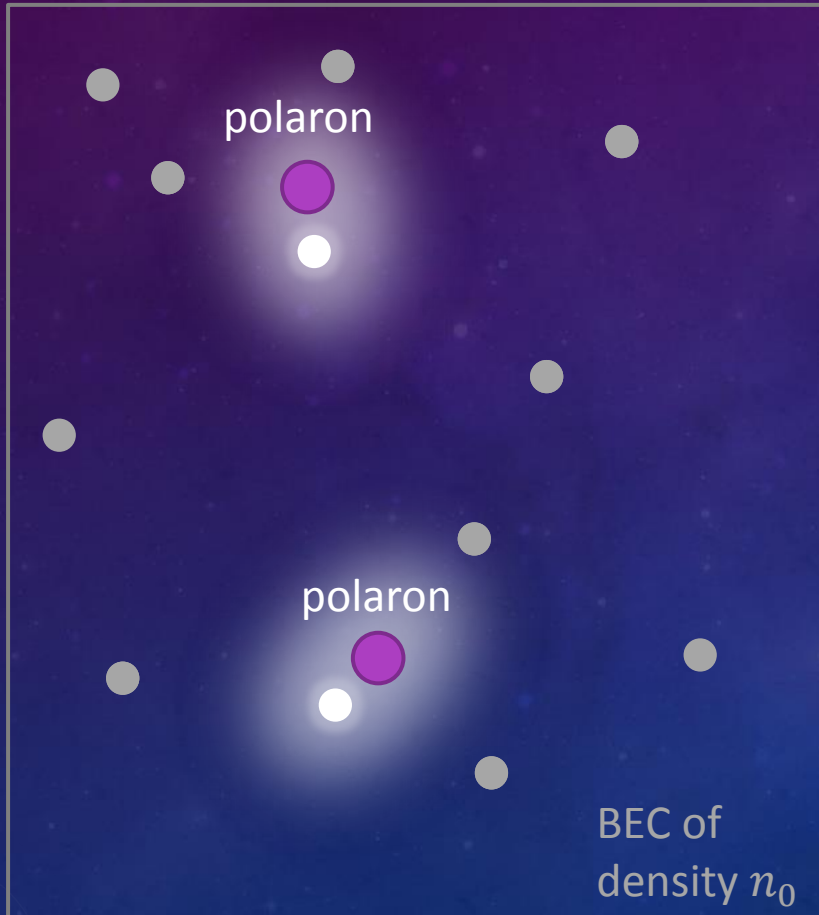
IMPURITIES IN A BOSE-EINSTEIN CONDENSATE



IMPURITIES IN A BOSE-EINSTEIN CONDENSATE

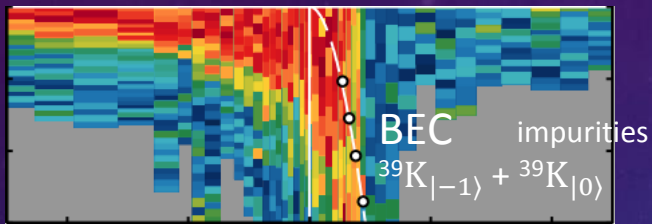


IMPURITIES IN A BOSE-EINSTEIN CONDENSATE

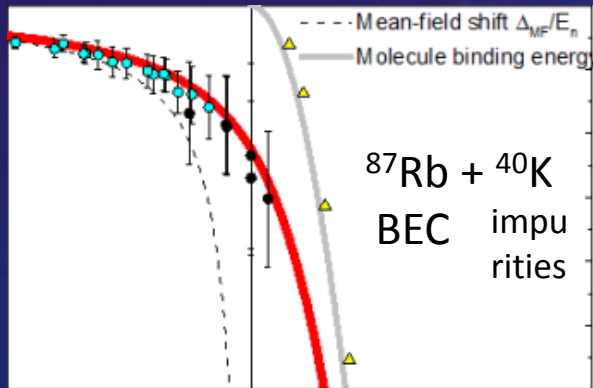


IMPURITIES IN A BOSE-EINSTEIN CONDENSATE

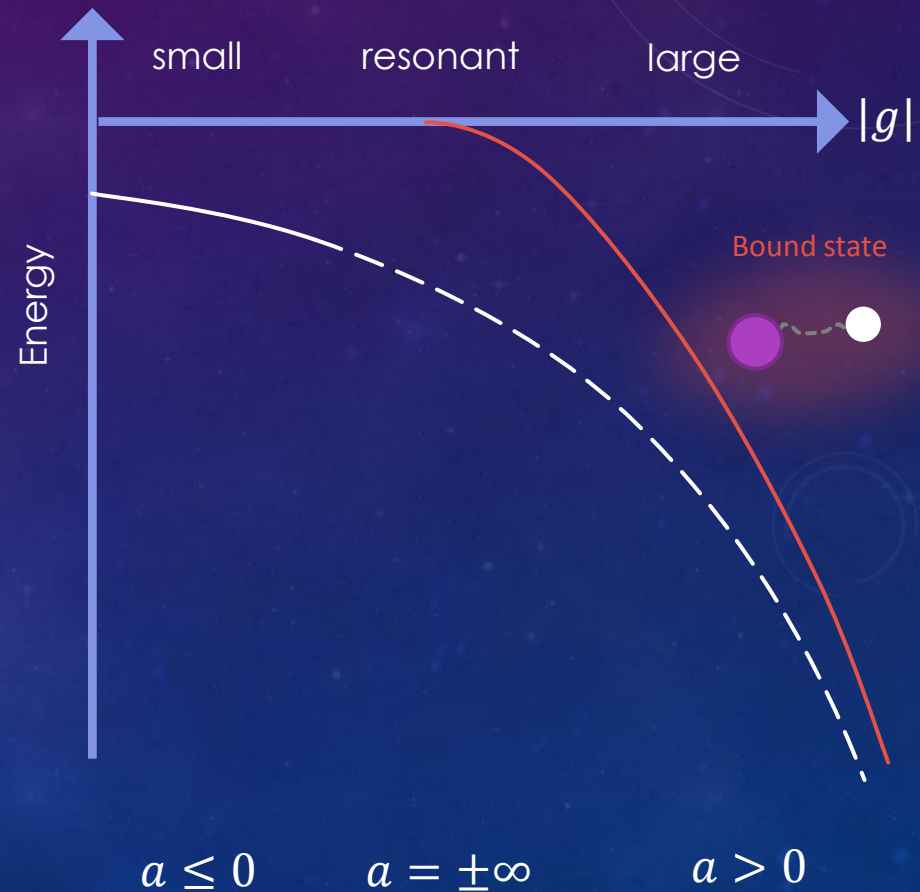
Bose polaron recently observed



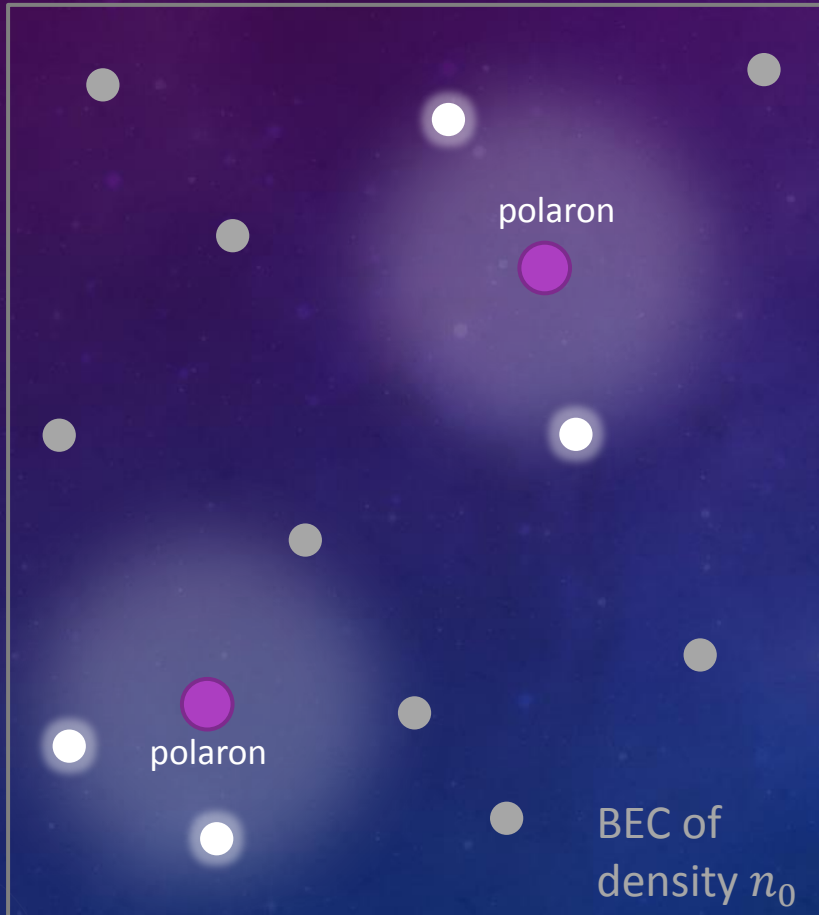
Jørgensen et al, PRL 117, 055302 (2016)



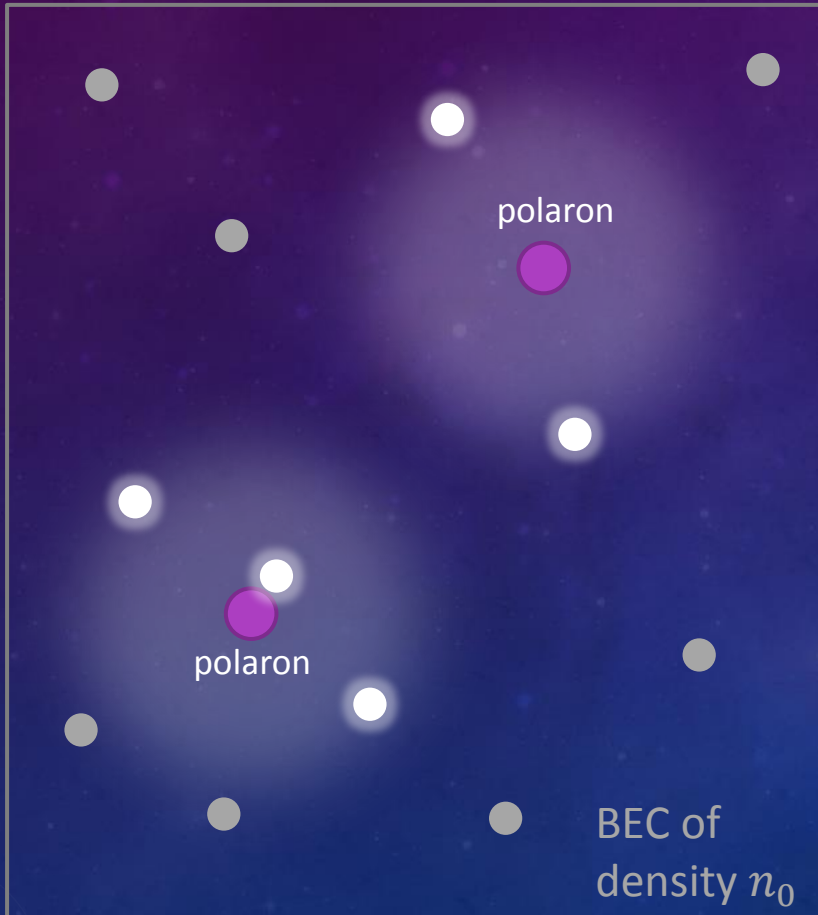
Ming-Guang Hu et al, PRL 117, 055301 (2016)



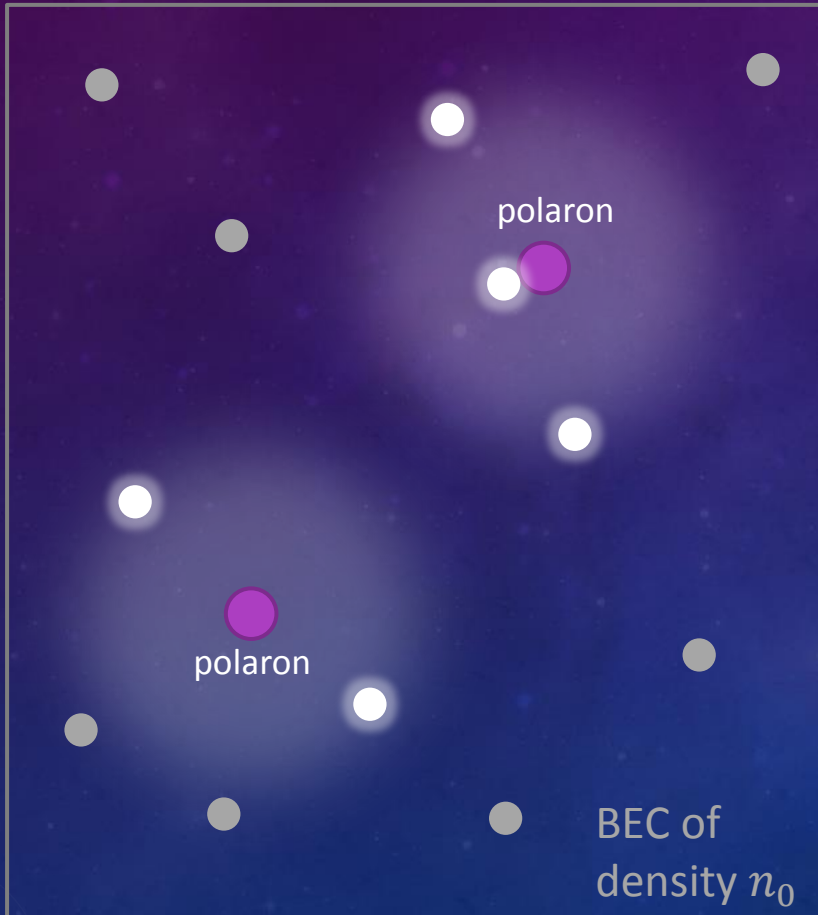
POLARONIC INTERACTION



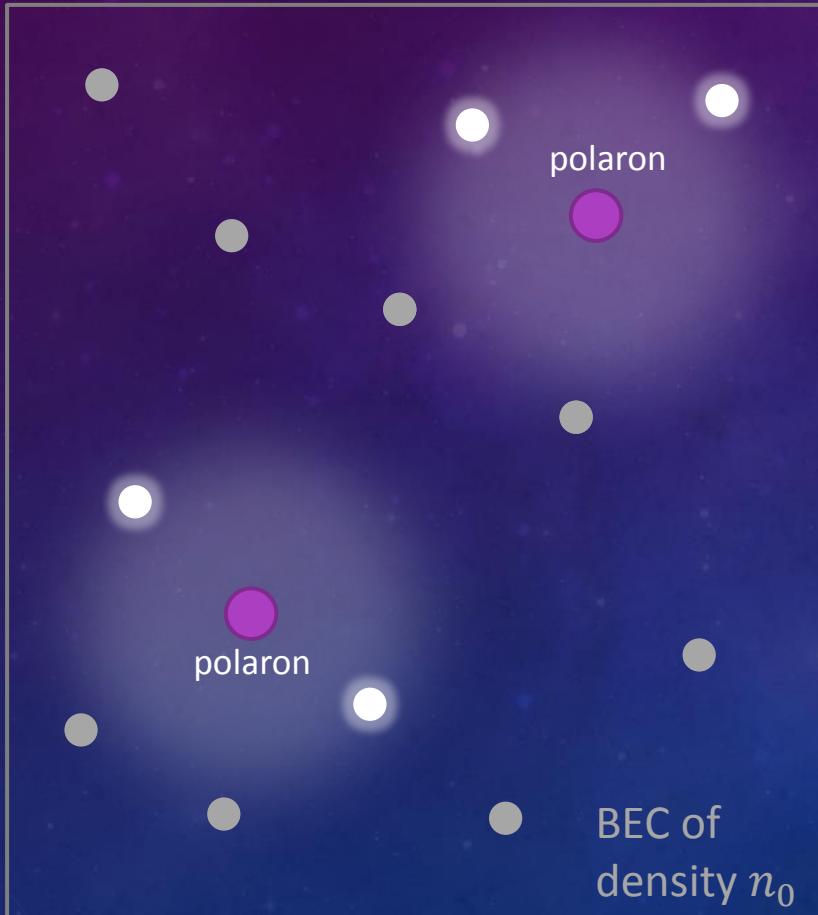
POLARONIC INTERACTION



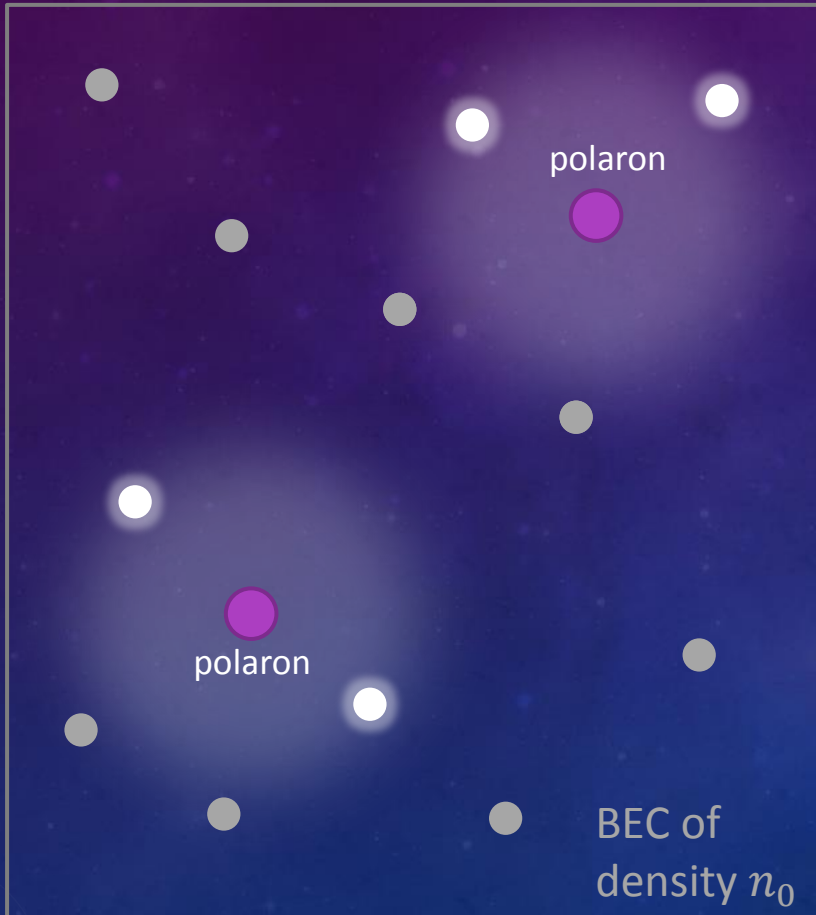
POLARONIC INTERACTION



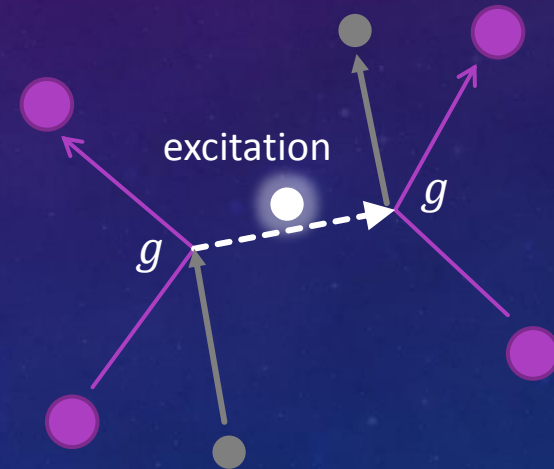
POLARONIC INTERACTION



POLARONIC INTERACTION



The (Bogoliubov) quasi-particles excitations of the BEC can mediate a Yukawa interaction



To second-order in perturbation theory:

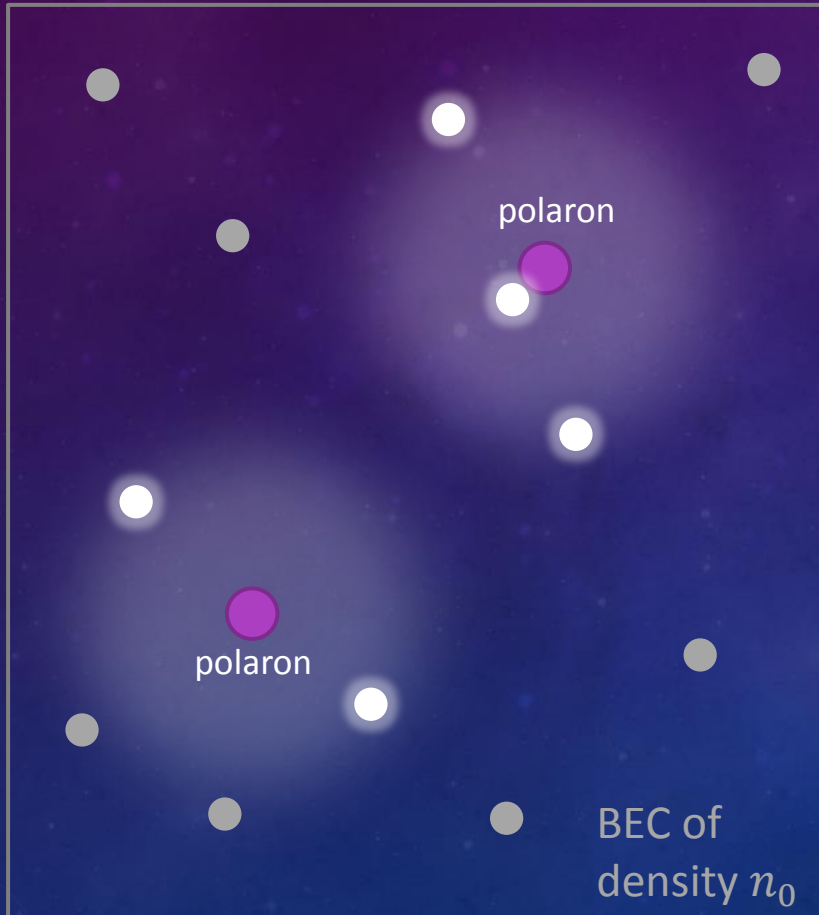
$$V(r) \propto -g^2 n_0 \frac{e^{-\sqrt{2}r/\xi}}{r}$$

BEC coherence length

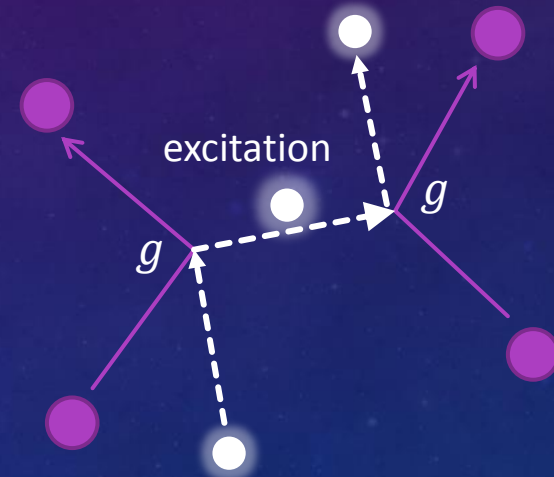
$$\xi = \frac{1}{\sqrt{8\pi n_0 a_B}}$$

Ex: Helium-3 impurities in Helium-4
Phys Rev 156, 207 (1967)

POLARONIC INTERACTION



The (Bogoliubov) quasi-particles excitations of the BEC can also mediate an Efimov interaction



Non-perturbative

NON-PERTURBATIVE METHOD: TRUNCATED BASIS

$$H = \underbrace{\sum_k \epsilon_k b_k^\dagger b_k + \frac{g_B}{2V} \sum_{k,k',p} b_{k'-p}^\dagger b_{k+p}^\dagger b_k b_{k'}}_{\text{Bosons}} + \underbrace{\sum_k \epsilon_k c_k^\dagger c_k}_{\text{Impurities}} + \underbrace{\frac{g}{V} \sum_{k,k',p} b_{k'-p}^\dagger c_{k+p}^\dagger c_k b_{k'}}_{\text{Impurity-boson}}$$

$\epsilon_k = \frac{\hbar^2 k^2}{2m}$
 $\epsilon_k = \frac{\hbar^2 k^2}{2M}$

Bogoliubov approximation
 $b_0 = \sqrt{N_0}$ condensate
 $b_k = u_k \beta_k - v_k \beta_k^\dagger$ Bogoliubov excitation

$$H = \underbrace{E_0 + \sum_k E_k \beta_k^\dagger \beta_k}_{\text{Bogoliubov excitation energy}} + \underbrace{\sum_k (\epsilon_k + g n_0) c_k^\dagger c_k}_{\text{Impurities}} + \underbrace{\sqrt{N_0} \frac{g}{V} \sum_{k,p} (u_p \beta_{-p}^\dagger - v_p \beta_p) c_{k+p}^\dagger c_k + h.c.}_{\text{Fröhlich}}$$

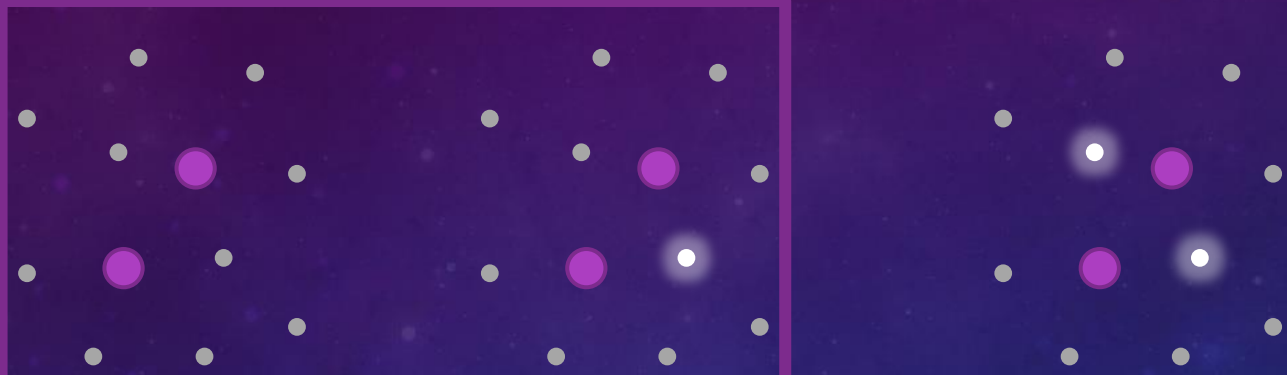
$$+ \frac{g}{V} \sum_{k,k',p} (u_{k'-p} u_{k'} \beta_{k'-p}^\dagger \beta_{k'} + v_{k'-p} v_{k'} \beta_{p-k'} \beta_{-k'}^\dagger) c_{k+p}^\dagger c_k$$

$$+ \frac{g}{V} \sum_{k,k',p} (u_{k'-p} v_{k'} \beta_{k'-p}^\dagger \beta_{-k'}^\dagger + v_{k'-p} u_{k'} \beta_{p-k'} \beta_{k'}) c_{k+p}^\dagger c_k$$

Scattering

Bogoliubov excitation energy
 $E_k = \sqrt{\epsilon_k(\epsilon_k + 2g_B n_0)}$

NON-PERTURBATIVE METHOD: TRUNCATED BASIS



The diagram shows a 2D lattice of particles. Small grey dots represent the BEC ground state. Larger purple circles represent impurities. A white dot represents an excitation. The left lattice is enclosed in a red box, indicating the truncated basis. The right lattice shows the full state with more particles.

$$|\Psi\rangle = \left(\sum_q \alpha_q c_q^\dagger c_{-q}^\dagger + \sum_{q,q'} \alpha_{q,q'} c_q^\dagger c_{q'}^\dagger \beta_{-q-q'}^\dagger + \sum_{q,q',q''} \alpha_{q,q',q''} c_q^\dagger c_{q'}^\dagger \beta_{q''}^\dagger \beta_{-q-q'-q''}^\dagger + \dots \right) |\Phi\rangle$$

Impurity creation operator

Excitation creation operator

BEC ground state

EQUATIONS

Coupled equations for α_q and to $\alpha_{q,q'}$

Zero-range limit (cutoff $\Lambda \rightarrow \infty$)

Three-body Integral equation

$$\frac{1}{T_q(E)} F_q + \frac{1}{V} \sum_k \frac{u_k^2 F_{k-q}}{E_k + \varepsilon_{|k-q|} + \varepsilon_q - E} = \frac{2n_0}{2\varepsilon_q - E} F_q$$

$$F_q = g \frac{1}{V} \sum_k u_k \alpha_{q,k-q}$$

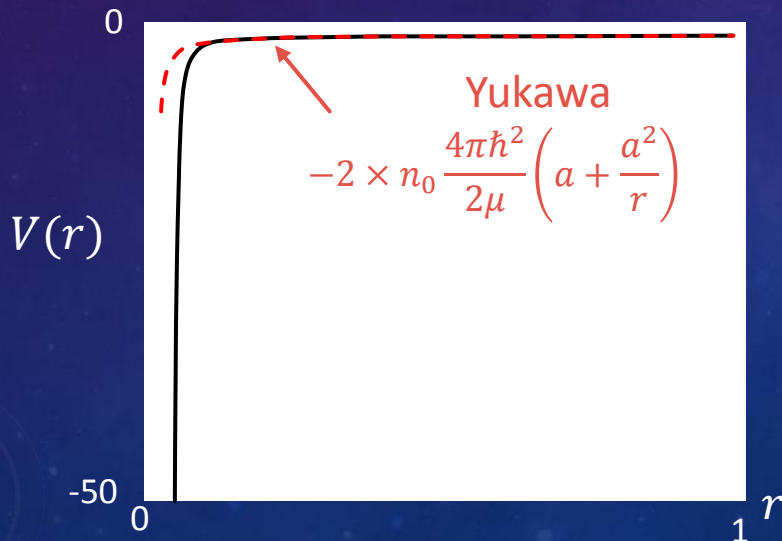
$$\frac{1}{T_q(E)} = \frac{2\mu}{4\pi\hbar^2} \frac{1}{a} + \frac{1}{V} \sum_k \left(\frac{u_k^2}{E_k + \varepsilon_{|k-q|} + \varepsilon_q - E} - \frac{1}{\varepsilon_k + \varepsilon_k} \right)$$

RESULT: POLARONIC POTENTIAL

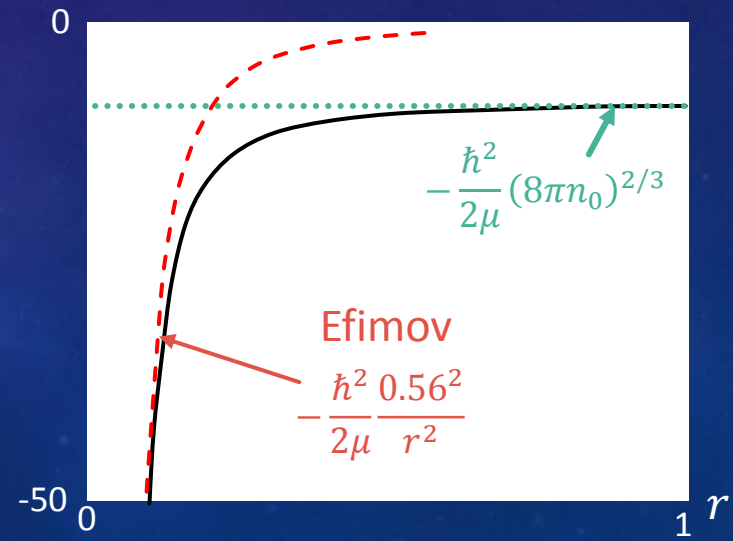
Effective potential (Born-Oppenheimer) between polarons:

$$V(r) = \frac{\hbar^2 \kappa^2}{2\mu} \left(\frac{1}{a} - \kappa + \frac{1}{r} e^{-\kappa r} + \frac{8\pi n_0}{\kappa^2} \right) = 0 \quad (a_B \rightarrow 0)$$

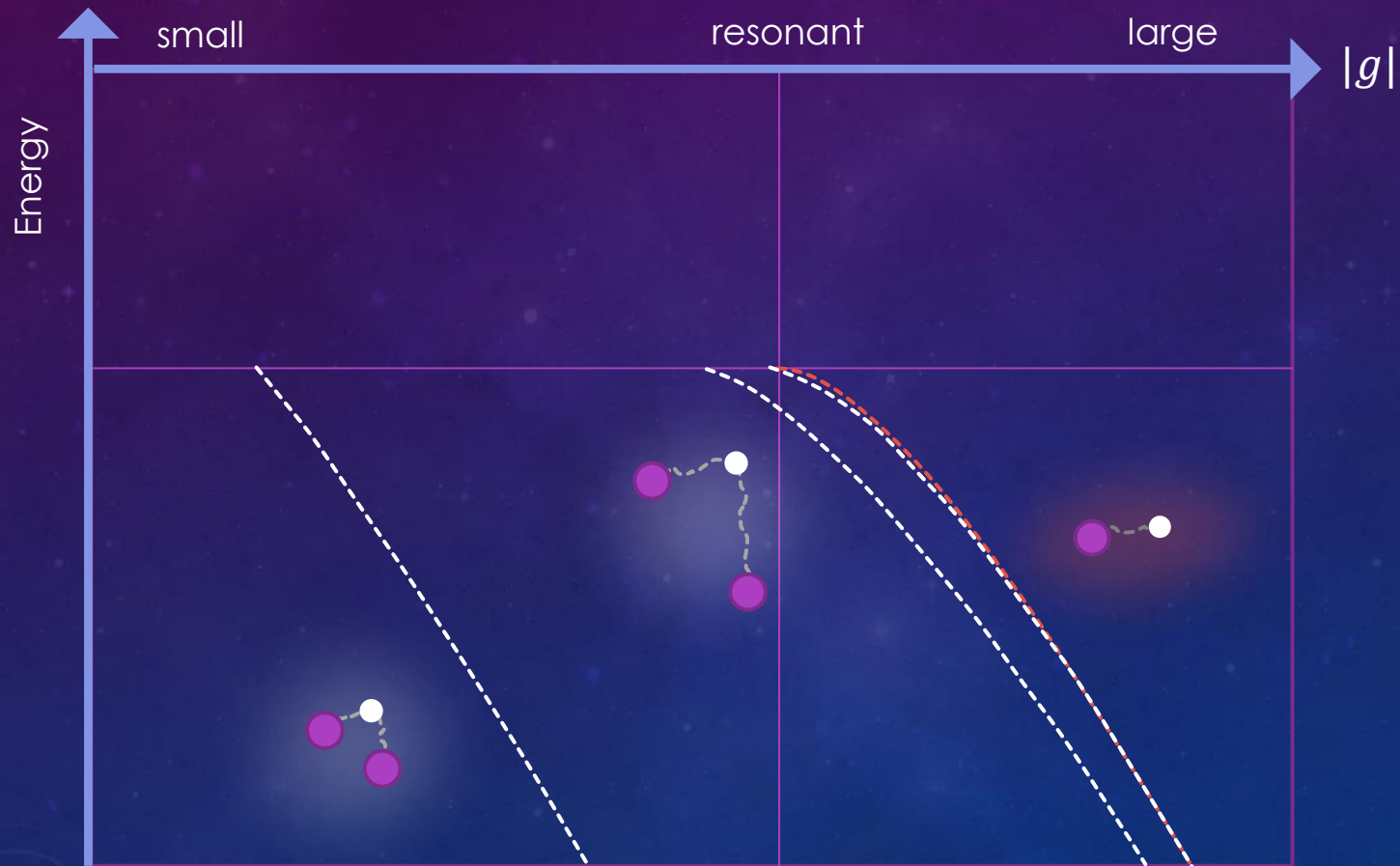
For small $a \lesssim 0$



At resonance $a = \pm\infty$



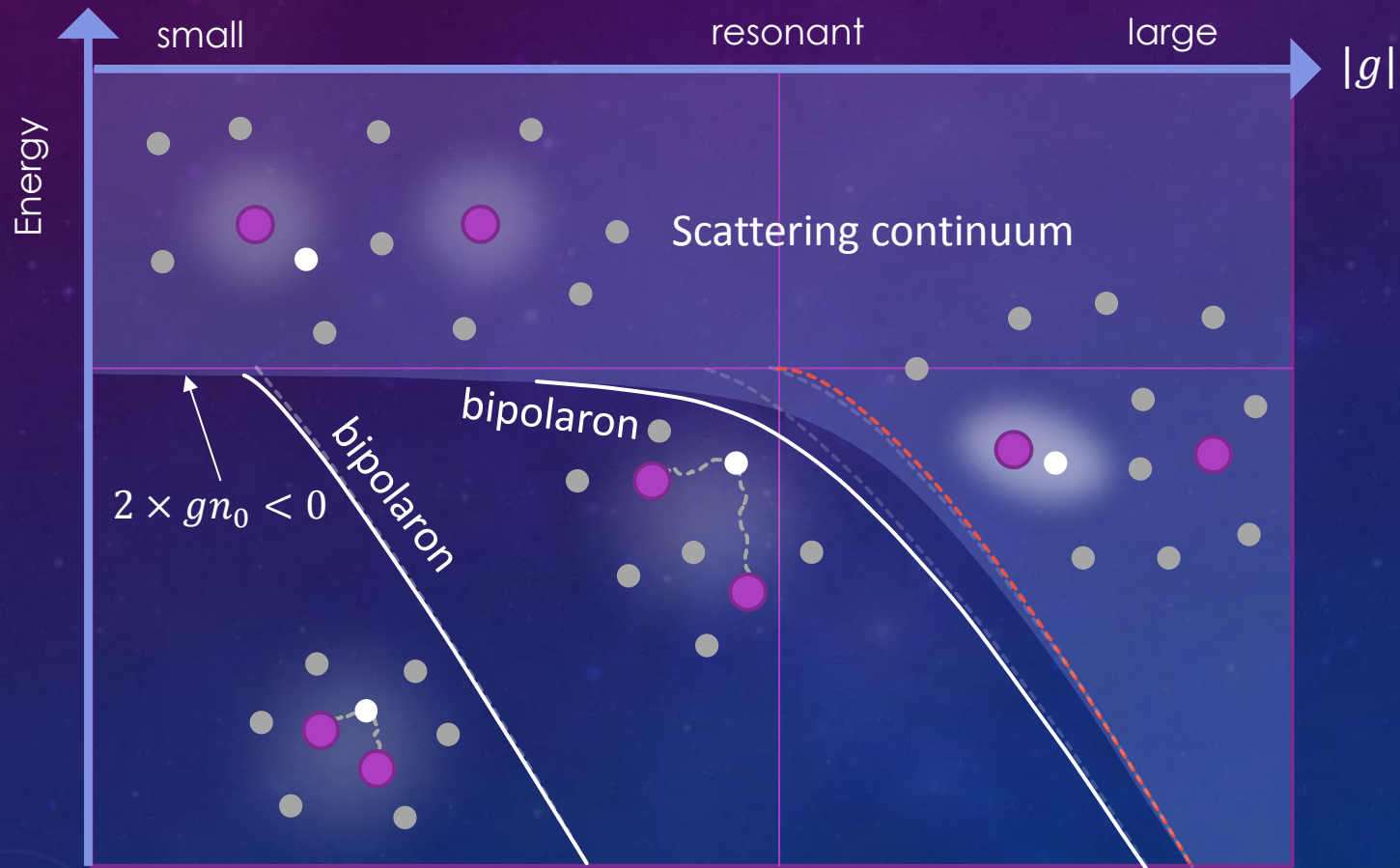
RESULT: ENERGY SPECTRUM



$$\frac{M}{m} = 30$$

$$n_0 a_B^3 = 0.0005$$

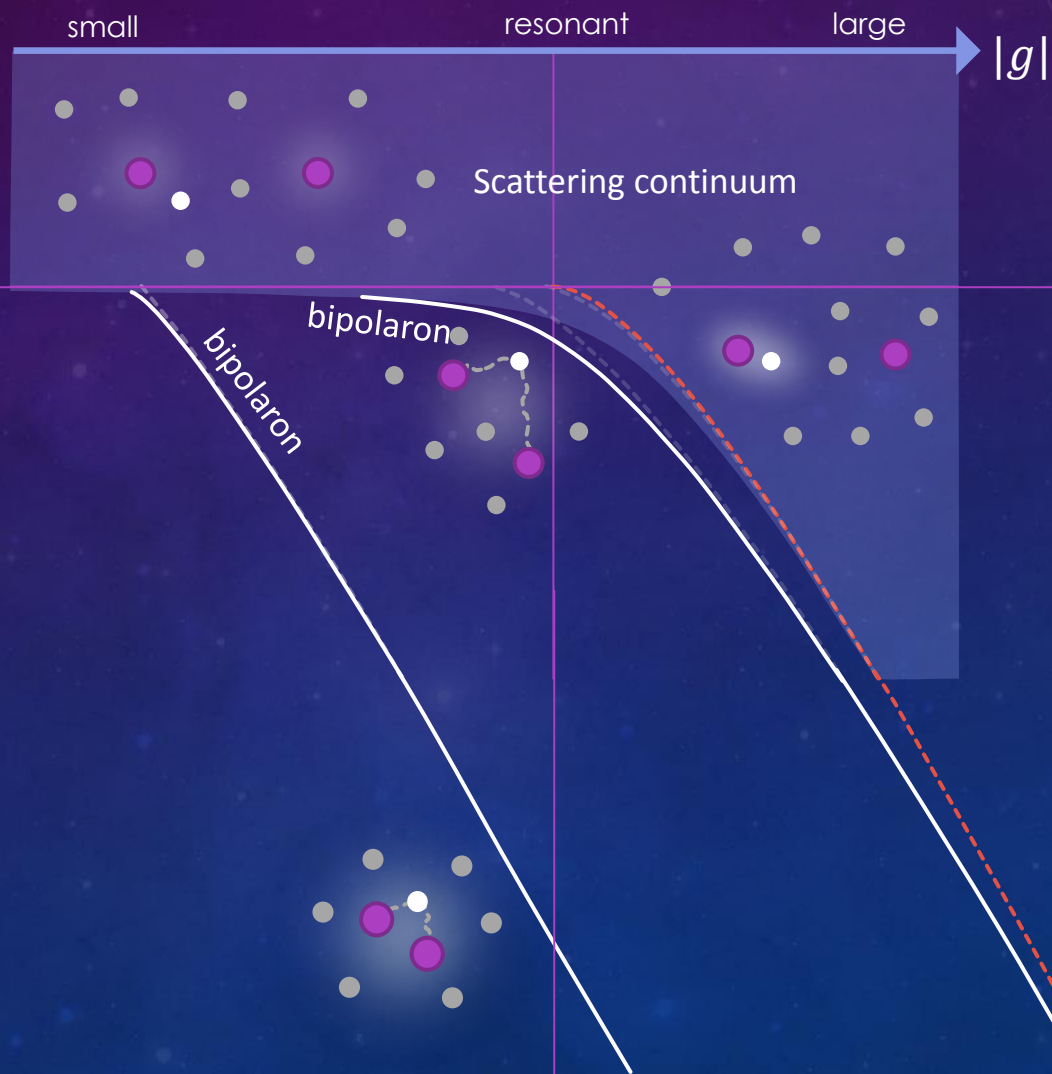
RESULT: ENERGY SPECTRUM



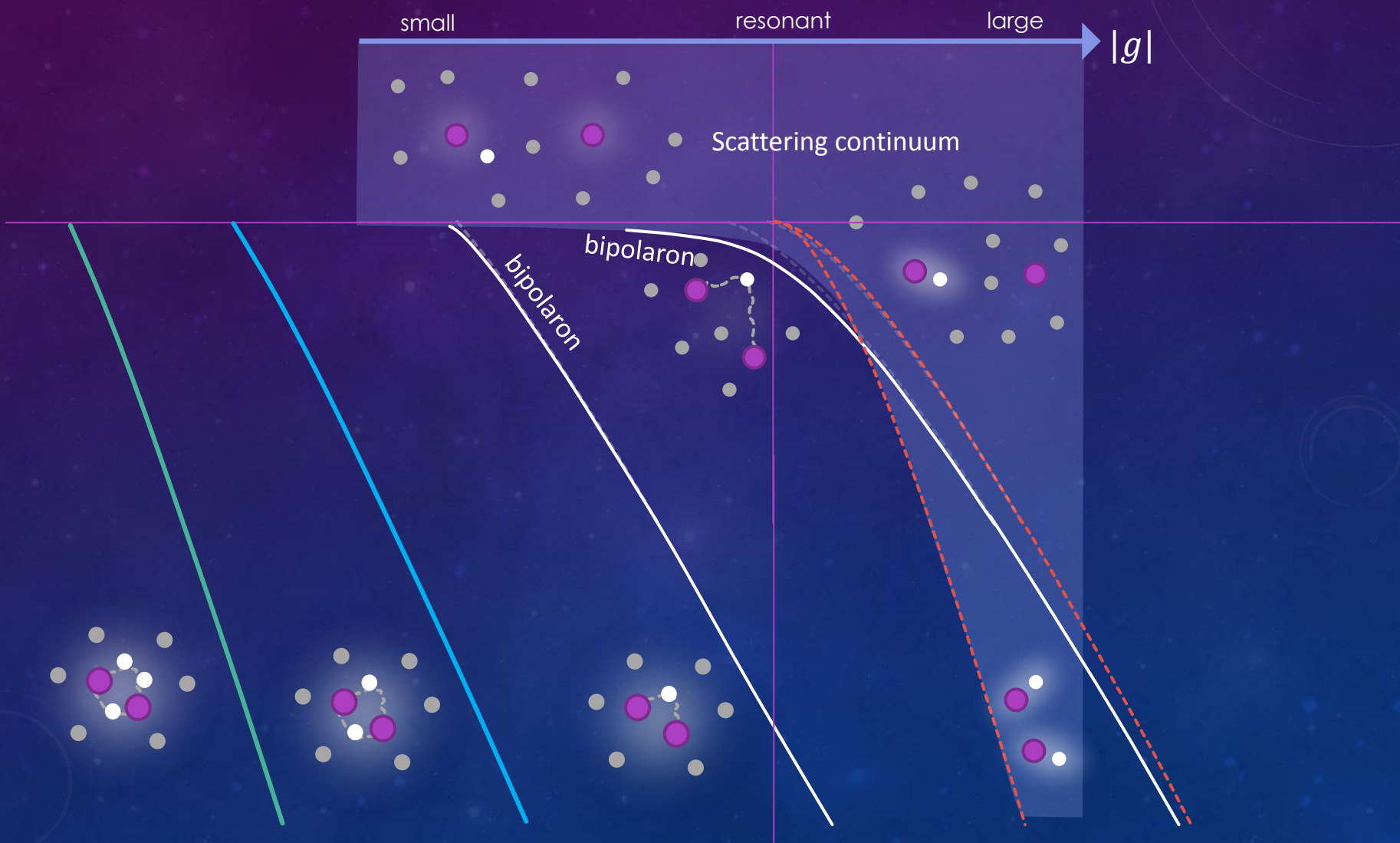
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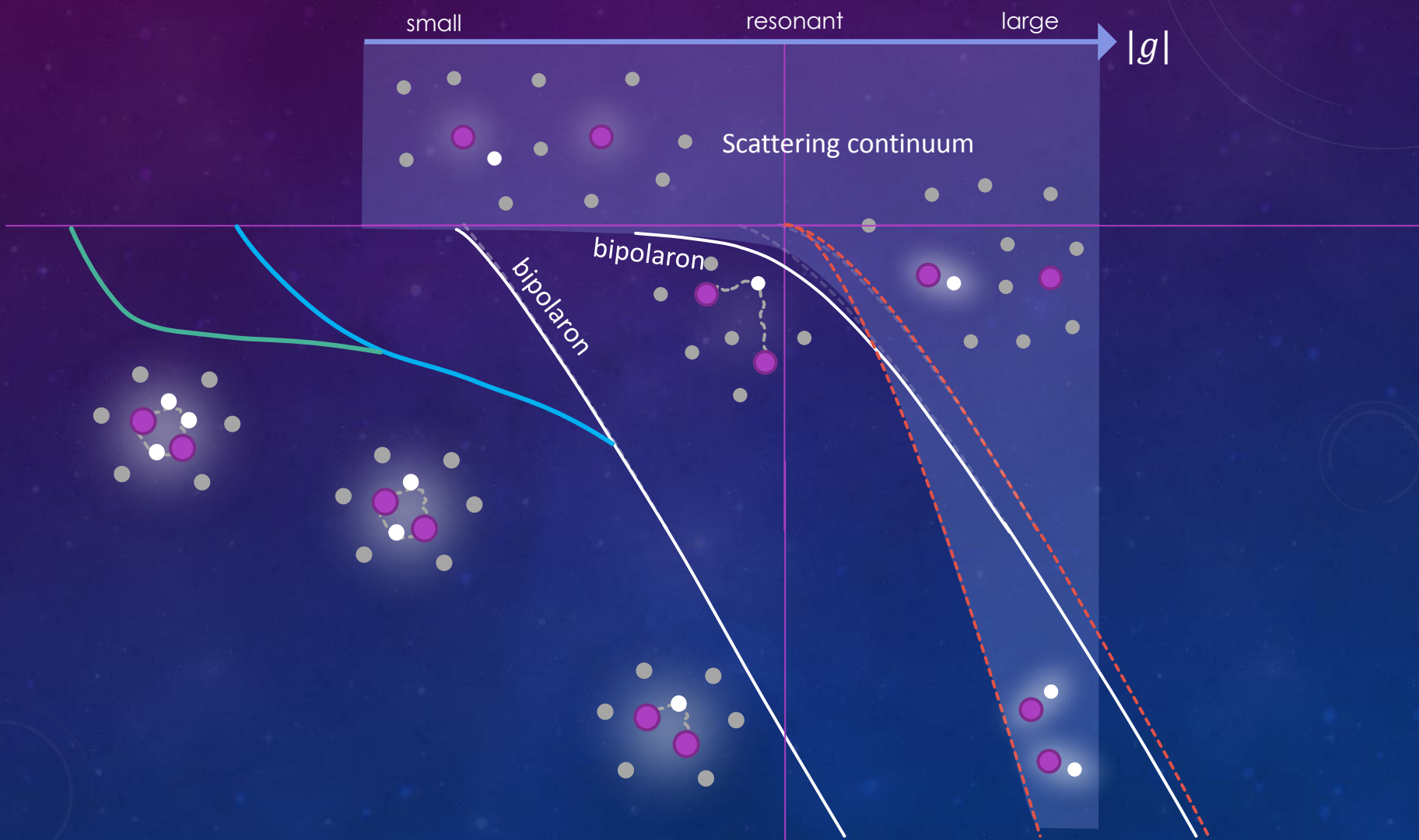
RESULT: ENERGY SPECTRUM



RESULT: ENERGY SPECTRUM



RESULT: ENERGY SPECTRUM



POSSIBLE EXPERIMENTAL OBSERVATIONS

Heavy impurities (e.g. ^{133}Cs) in a condensate of light bosons (e.g. ^7Li)

- Polaron RF spectroscopy: mean-field shift due to the polaron interaction
- Loss by recombination: shift of the loss peak with the light boson density

CONCLUSIONS

- Bose polarons formed by impurities in a Bose-Einstein interact through a Yukawa attraction that turns into an **Efimov attraction** for resonant boson-impurity attraction.
- This attraction can bind the two polarons into one or several **bipolarons**, that asymptote to **Efimov trimers** of two impurities and a boson.

arXiv:1607.04507