Coupling of collective and single-particle degrees of freedom in symmetry-restored GCM

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incomplete work in progress begun while having been at the

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Effective interaction used throughout this talk: SLyMR0

$$\begin{split} \hat{\nu} &= t_0 \left(1 + x_0 \hat{P}_{\sigma} \right) \hat{\delta}_{r_1 r_2} \\ &+ \frac{t_1}{2} \left(1 + x_1 \hat{P}_{\sigma} \right) \left(\hat{\mathbf{k}}_{12}^{\,\prime 2} \hat{\delta}_{r_1 r_2} + \hat{\delta}_{r_1 r_2} \hat{\mathbf{k}}_{12}^{\,2} \right) \\ &+ t_2 \left(1 + x_2 \hat{P}_{\sigma} \right) \hat{\mathbf{k}}_{12}^{\,\prime} \cdot \hat{\delta}_{r_1 r_2} \hat{\mathbf{k}}_{12} \\ &+ \mathrm{i} \, W_0 \left(\hat{\sigma}_1 + \hat{\sigma}_2 \right) \cdot \hat{\mathbf{k}}_{12}^{\,\prime} \times \hat{\delta}_{r_1 r_2} \hat{\mathbf{k}}_{12} \\ &+ u_0 \left(\hat{\delta}_{r_1 r_3} \hat{\delta}_{r_2 r_3} + \hat{\delta}_{r_3 r_2} \hat{\delta}_{r_1 r_2} + \hat{\delta}_{r_2 r_1} \hat{\delta}_{r_3 r_1} \right) \\ &+ v_0 \left(\hat{\delta}_{r_1 r_3} \hat{\delta}_{r_2 r_3} \hat{\delta}_{r_3 r_4} + \hat{\delta}_{r_1 r_2} \hat{\delta}_{r_3 r_2} \hat{\delta}_{r_2 r_4} + \cdots \right) \end{split}$$

J. Sadoudi, M. Bender, K. Bennaceur, D. Davesne, R. Jodon, and T. Duguet, Physica Scripta T154 (2013) 014013



- it is impossible to fullfil the usual nuclear matter constraints, to have stable interactions and attractive pairing
- no "best fit" possible
- very bad performance compared to standard general functionals

J. Sadoudi, M. Bender, K. Bennaceur, D. Davesne, R. Jodon, and T. Duguet, Physica Scripta T154 (2013) 014013



Rotational band in ²⁴Mg

J, K decomposition of the norm J. K decomposition of the energy -190 0.12 $J=0, K_1=K_2=0$ $J=2,K_1=K_2=0$ - J=3,K₁=K₂ = 1 -- J=4,K₁=K₂ = 0 -1950.1-- J=5,K₁=K₂ = 1 · $J=6, K_1=K_2=0$ -- J=5,K₁=K₂ = 1 $J=8,K_1=K_2=0$ $\langle \Phi | \hat{P}^J_{KK} \hat{P}^N \hat{P}^Z | \Phi \rangle$ 0.08 -- J=8,K₁=K₂ = 1 --200 --- J=8,K₁=K₂ = 2 ----- $J=8, K_1=K_2=3$ E (MeV) $J=8, K_1=K_2 = 4$ $J=8, K_1=K_2=5$ -0.06 $J = 8, K_1 = K_2 = 6$ -205 $J=8, K_1=K_2=7$ $J=8, K_1=K_2=8$ 0.04 -2100.02-2150.0 2 3 5 0 4 $\hbar\omega$ $\hbar\omega$ Bender, Avez, Bally, Heenen, to be published

Rotational band in ²⁴Mg

-190

J, K decomposition of the energy





Rotational band in ²⁴Mg





Left: Non-projected total energy of the HFB vacua (without LN correction) relative to the spherical configuration. Middle: N = 26, Z = 20 projected total energy of the HFB vacua relative to the spherical configuration. Right: Energy of the projected N = 26, Z = 20, J = 0 HFB vacua.

Bender & Heenen, to be published



Top row: Right: Energy of the J = 0 HFB vacua. Middle: Energy of the lowest K-mixed J = 2 projected state . Right: Energy of the second K-mixed J = 2 state . Bottom row: Right: Energy of the J = 3 state. Middle: Energy of the lowest K-mixed J = 4 projected state. Right: Energy of the second K-mixed J = 4 state. The total energy is relative to the minimum of the J = 0 energy surface. All states are projected on N = 26, Z = 20,

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æ $12.0 \\ 11.0$ ^{46}Ca 10.0 9.0 0.9 · 6.0 0,8 3.0 4.0 3.0 2.0 1.0 0.0 0.2 20 0.0 0,5 90 0,0 0,2 0.2 0.7 0.0 0.1 $0.2 \ 0.3 \ 0.4$ 0.5 0.6 0.7 0.8 0.9 1.0 1.2 в





Nilsson diagram along the path indicated by cyan dots. Vertical bars indicate the deformation of the minima. Nilsson diagram for a closed path through indicated by yellow dots.

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- There is a sequence of "seniority-2" states with J^π = 2⁺, 4⁺, 6⁺ that in the shell-model is easily obtained by coupling two neutron holes in the 1f_{7/2}shell to these angular momenta.
- These are non-collective; hence, cannot be described by "traditional" GCM.

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M. Bender, IPN Lyon

Coupling of collective and single-particle degrees of freedom



Construction of new forms of effective interactions

Skyrme-type interactions with higher-order terms in derivatives

(not aiming at true Hamiltonians so far, though) Carlsson, Dobaczewski, Kortelainen, PRC 78 (2008) 044326 Raimondi, Carlsson, Dobaczewski, PRC 83 (2011) 054311 Davesne, Pastore, Navarro, JPG 40 (2013) 095104 Becker, Davesne, Meyer, Pastore, Navarro, JPG 42 (2015) 034001

Skyrme-type interactions with explicit three-body interactions

Sadoudi, thèse, Université de Paris-Sud XI (2011)

Sadoudi, Bender, Bennaceur, Davesne, Jodon, Duguet, Phys Scr T154 (2013) 014013

Sadoudi, Duguet, Meyer, Bender, PRC 88 (2013) 064326

regularised contact interactions (replacing the delta function in SKyrme with Gaussians)

Raimondi, Bennaceur, Dobaczewski, JPG 41 (2014) 055112

Bennaceur, Idini, J. Dobaczewski, P. Dobaczewski, Kortelainen, Raimondi, JPG44 (2017) 045106

non-local three-body forces simulating density dependences

Gezerlis, Bertsch, PRL 105 (2010) 212501

Lacroix, Bennaceur, PRC 91 (2015) 011302(R)

or try a different strategy: explicit in-medium correlations from MBPT

Duguet, Bender, Ebran, Lesinski, Somà, EPJA 51 (2015) 162

the most general central Skyrme-type 3-body force up to 2nd order in gradients has been constructed by J. Sadoudi with a dedicated formal algebra code

$$\begin{split} \hat{\nu}_{123} &= u_0 \left(\hat{\delta}_{r_1 r_3} \hat{\delta}_{r_2 r_3} + \hat{\delta}_{r_3 r_2} \hat{\delta}_{r_1 r_2} + \hat{\delta}_{r_2 r_1} \hat{\delta}_{r_3 r_1} \right) \\ &+ \frac{u_1}{2} \left[1 + y_1 P_{12}^{\sigma} \right] \left(\hat{\mathbf{k}}_{12} \cdot \hat{\mathbf{k}}_{12} + \hat{\mathbf{k}}_{12}' \cdot \hat{\mathbf{k}}_{12}' \right) \hat{\delta}_{r_1 r_3} \hat{\delta}_{r_2 r_3} \\ &+ \frac{u_1}{2} \left[1 + y_1 P_{31}^{\sigma} \right] \left(\hat{\mathbf{k}}_{31} \cdot \hat{\mathbf{k}}_{31} + \hat{\mathbf{k}}_{31}' \cdot \hat{\mathbf{k}}_{31}' \right) \hat{\delta}_{r_3 r_2} \hat{\delta}_{r_1 r_2} \\ &+ \frac{u_1}{2} \left[1 + y_1 P_{23}^{\sigma} \right] \left(\hat{\mathbf{k}}_{23} \cdot \hat{\mathbf{k}}_{23} + \hat{\mathbf{k}}_{23}' \cdot \hat{\mathbf{k}}_{23}' \right) \hat{\delta}_{r_2 r_1} \hat{\delta}_{r_3 r_1} \\ &+ u_2 \left[1 + y_{21} P_{12}^{\sigma} + y_{22} (P_{13}^{\sigma} + P_{23}^{\sigma}) \right] \left(\hat{\mathbf{k}}_{12} \cdot \hat{\mathbf{k}}_{12}' \right) \hat{\delta}_{r_1 r_3} \hat{\delta}_{r_2 r_3} \\ &+ u_2 \left[1 + y_{21} P_{31}^{\sigma} + y_{22} (P_{32}^{\sigma} + P_{12}^{\sigma}) \right] \left(\hat{\mathbf{k}}_{31} \cdot \hat{\mathbf{k}}_{31}' \right) \hat{\delta}_{r_3 r_2} \hat{\delta}_{r_1 r_2} \\ &+ u_2 \left[1 + y_{21} P_{23}^{\sigma} + y_{22} (P_{21}^{\sigma} + P_{31}^{\sigma}) \right] \left(\hat{\mathbf{k}}_{23} \cdot \hat{\mathbf{k}}_{23}' \right) \hat{\delta}_{r_2 r_1} \hat{\delta}_{r_3 r_1} \end{split}$$

Sadoudi, Duguet, Meyer, Bender, PRC 88 (2013) 064326

degrees of freedom considered and level of modeling at which they enter: quadrupole and other deformations (SR & MR), pairing correlations (SR & MR), intrinsic angular momentum (SR & MR), quasi-particle excitations (SR & MR).

Breaking of time-reversal invariance in the SR states as prerequisite of coupling single-particle degrees of freedom to collective motion.

phenomena these can be expected to be (particularly) relevant for: Low-lying nuclear structure.

 their treatment as discussed can be (easily) combined with the one of other degrees of freedom?
Formally: yes (higher-order deformations, higher-order quasiparticle excitations). Computationally: not necessarily.

can they be expected to be independent/orthogonal to other degrees of freedom discussed during the workshop: Not necessarily.

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The work presented here would have been impossible without my collaborators

founding fathers Paul Bonche Hubert Flocard Paul-Henri Heenen	SPhT, CEA Saclay CSNSM Orsay Université Libre de Bruxelles
formal aspects of the big picture Thomas Duguet Denis Lacroix	Irfu/CEA Saclay & KU Leuven & NSCL/MSU IPN Orsay
design and implementation of code	extensions
Benoît Avez	CEN Bordeaux Gradignan
Benjamin Bally	CEN Bordeaux Gradignan, now SPhN, CEA Saclay
Veerle Hellemans	Université Libre de Bruxelles
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development and benchmarking of	new functionals
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Robin Jodon	IPN Lyon
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Jeremy Sadoudi Irf	Tu/CEA Saclay first, then CEN Bordeaux Gradignan
Kouhei Washiyama	Université Libre de Bruxelles

color code: active (past) member of the collaboration

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