

Two- and Four- quasiparticle states in
 $Z=92 - 108$ even-even nuclei
In HFB-D1S calculations

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Introduction

- Microscopic method : the HFB approximation with effective force gives the binding energy of the nucleus at the mean field level.
- Effective D1S force
- Time reversal symmetry broken (2 and 4 QP intrinsic excitations)
- Actinides and transactinides $Z= 92$ to 108
- HFB solutions by expanding on 15 major shells

Theory 1

HFB under constraints

Variational principle :

$$\delta[\langle \Phi | \mathbf{H} - \lambda_z \mathbf{Z} - \lambda_n \mathbf{N} - \sum_i \mu_i \mathbf{Q}_i | \Phi \rangle] = 0$$

with $\mathbf{H} = \sum_i \mathbf{T}_i + 1/2 \sum_{i \neq j} \mathbf{V}_{ij}$

where V_{ij} is the nucleon-nucleon effective interaction D1S of GOGNY

Constraint on

- Particle number: $\langle \Phi | \mathbf{Z} \text{ or } \mathbf{N} | \Phi \rangle = Z \text{ or } N$
- Multipole moments: $\langle \Phi | \mathbf{Q}_i | \Phi \rangle = q_i$

Theory 2

The two-quasiparticle excitations are sought for using blocking calculations.

A trial state is defined as $|\Phi'_{ij}\rangle = \eta_i^+ \eta_j^+ |\Phi_{ij}\rangle$
in which η_i^+ is a quasiparticle creation operator.

Then the equation
$$\delta\langle\Phi'_{ij}|\hat{H} - \lambda_N \hat{N} - \lambda_Z \hat{Z}|\Phi'_{ij}\rangle = 0$$

is solved, and the two-quasiparticle excitation energy is

$$E_{2qp}^{ij} = \langle\Phi'_{ij}|\hat{H}|\Phi'_{ij}\rangle - \langle\Phi|\hat{H}|\Phi\rangle.$$

Breaking time-reversal symmetry, the signature partner pairs with angular momentum projections on the z axis,

$$K_- = K_1 - K_2 \text{ and } K_+ = K_1 + K_2,$$

and parity $\pi = \pi_1 \cdot \pi_2$ are no longer degenerate in energy.

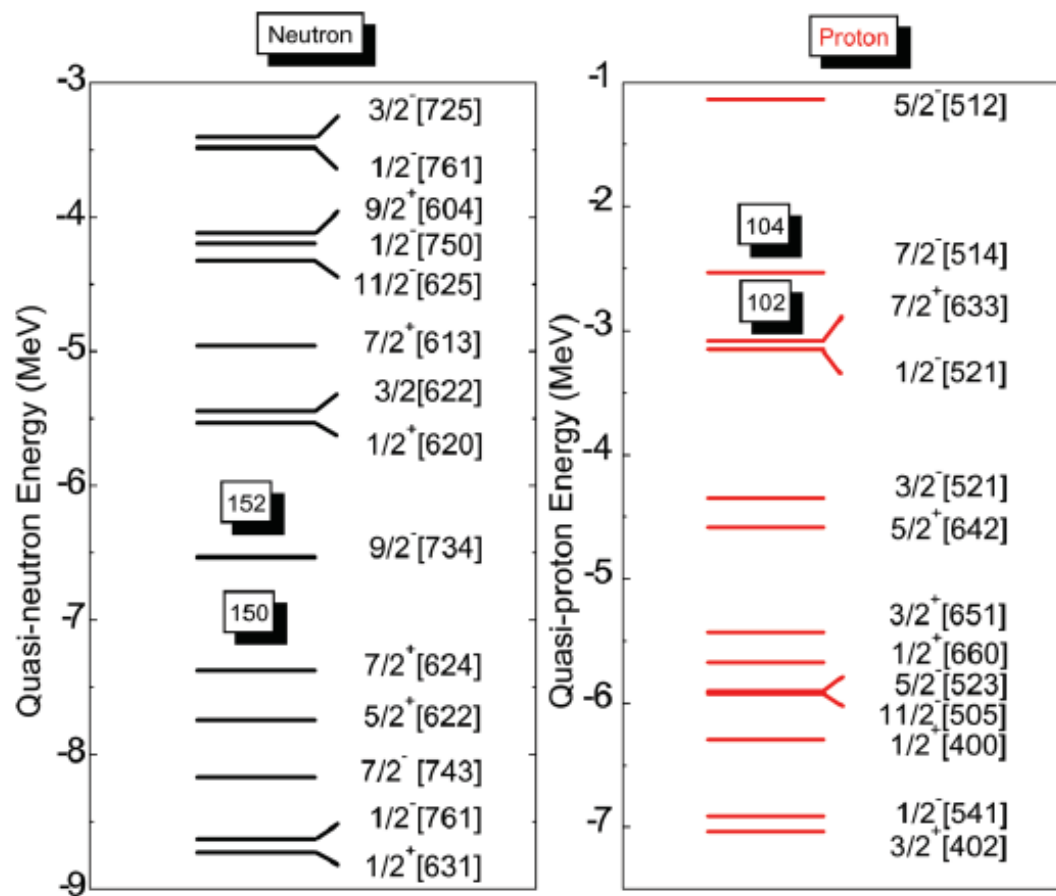
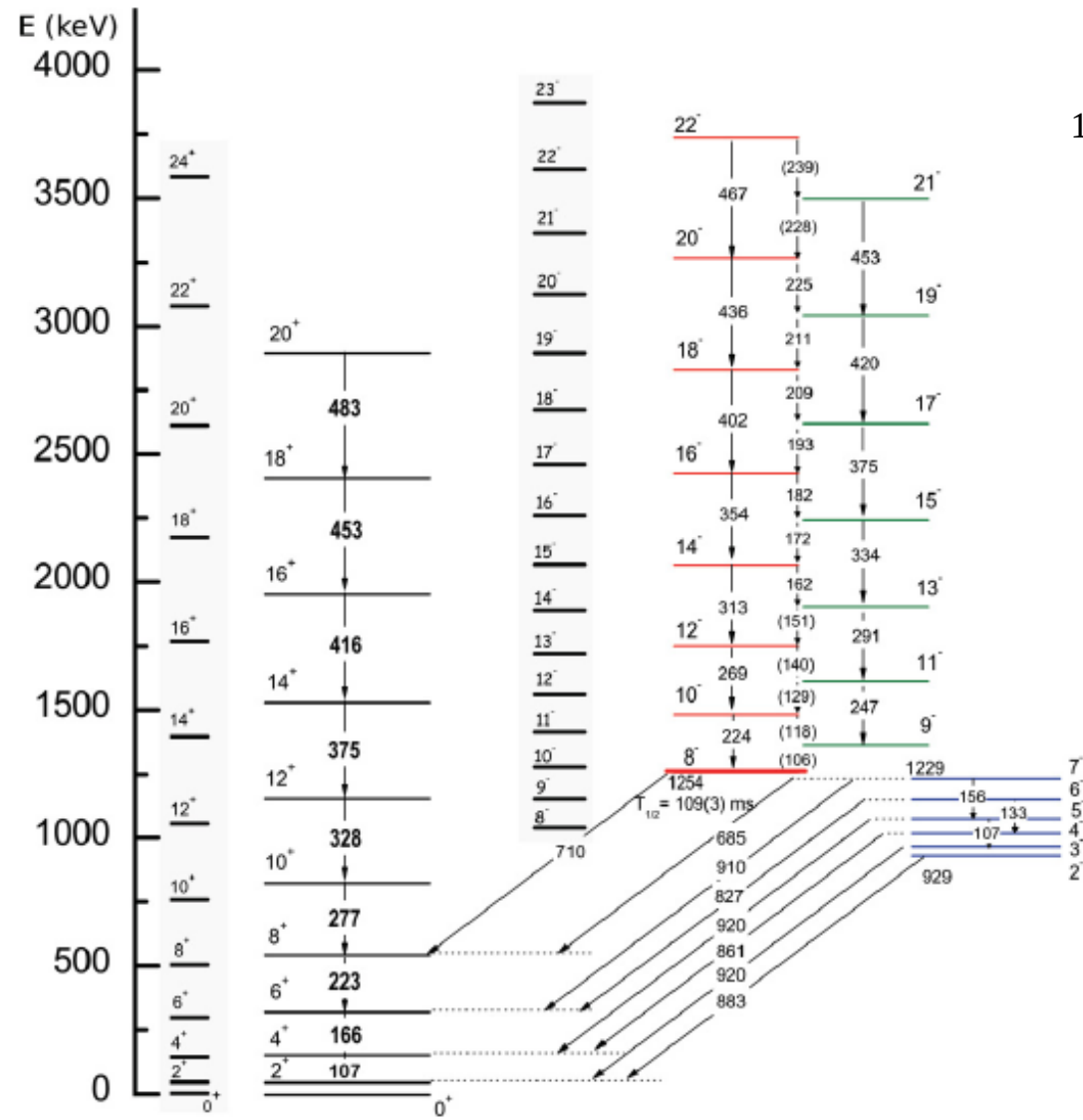


FIG. 9. (Color online) Single-particle energies for neutron and proton states in ^{252}No at axial equilibrium deformation of the HFB energy, where the charge quadrupole moment is $Q_0 = 13.75 e b$. The labels $[Nn_z\Lambda]$ are assigned by analogy with a Nilsson diagram.



^{152}No

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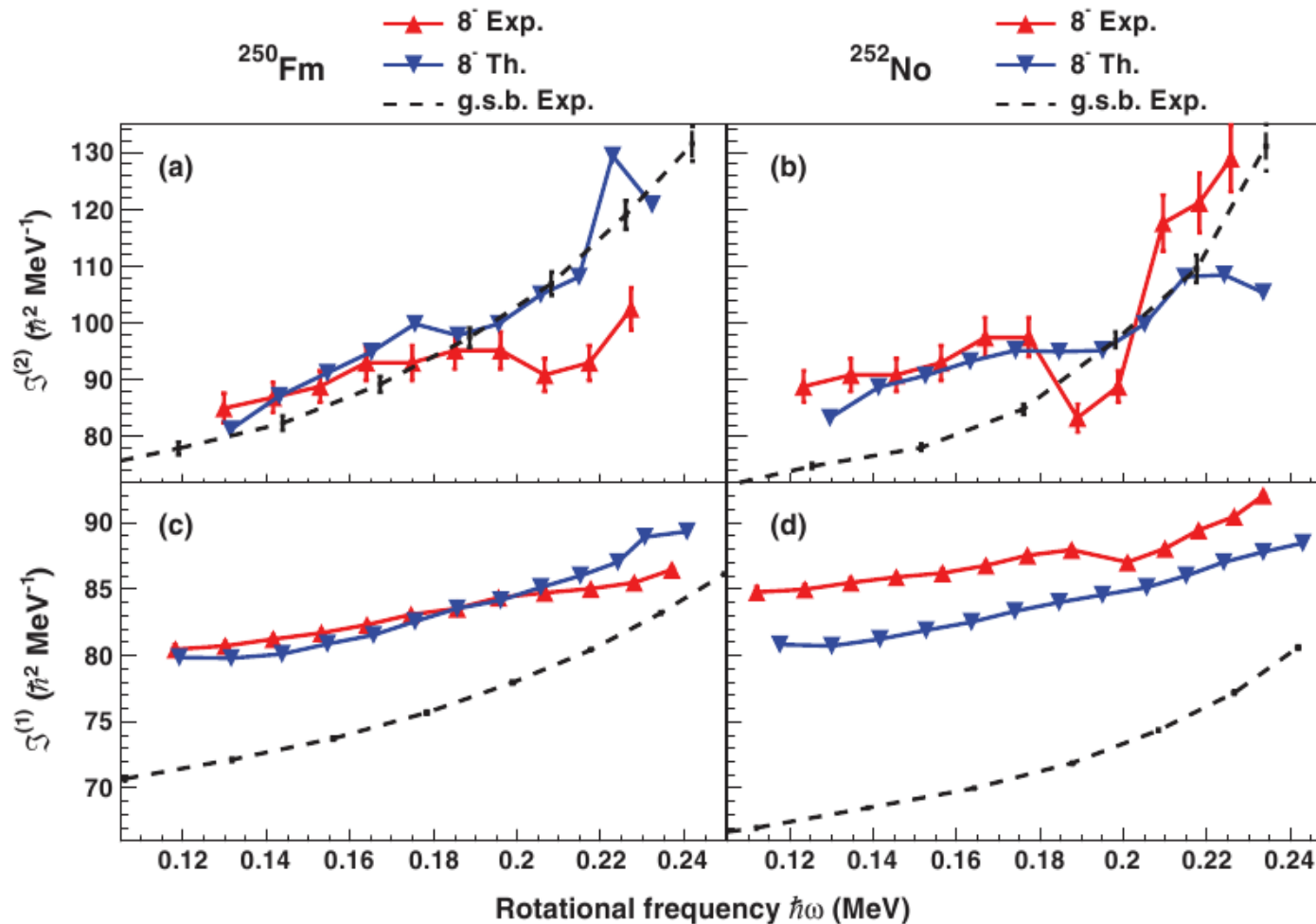
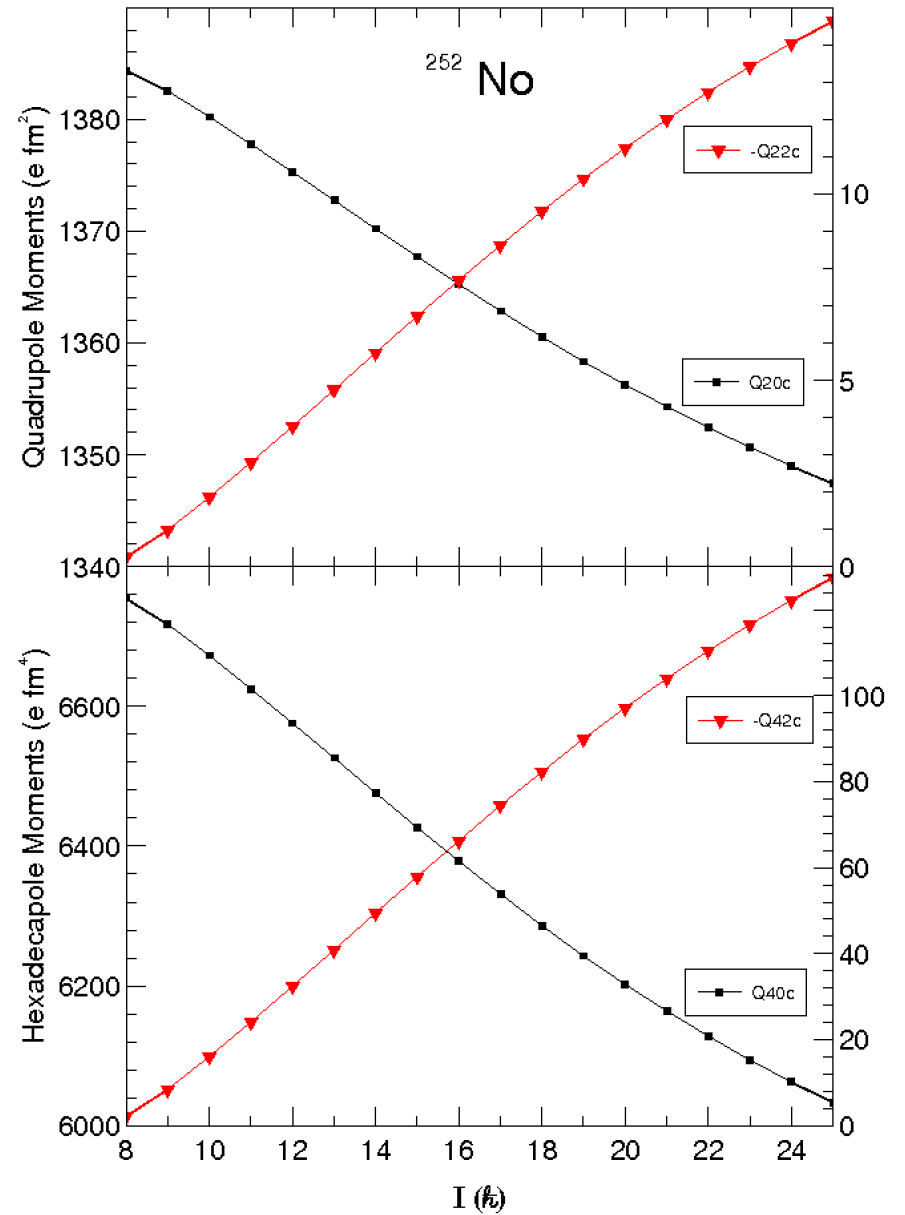
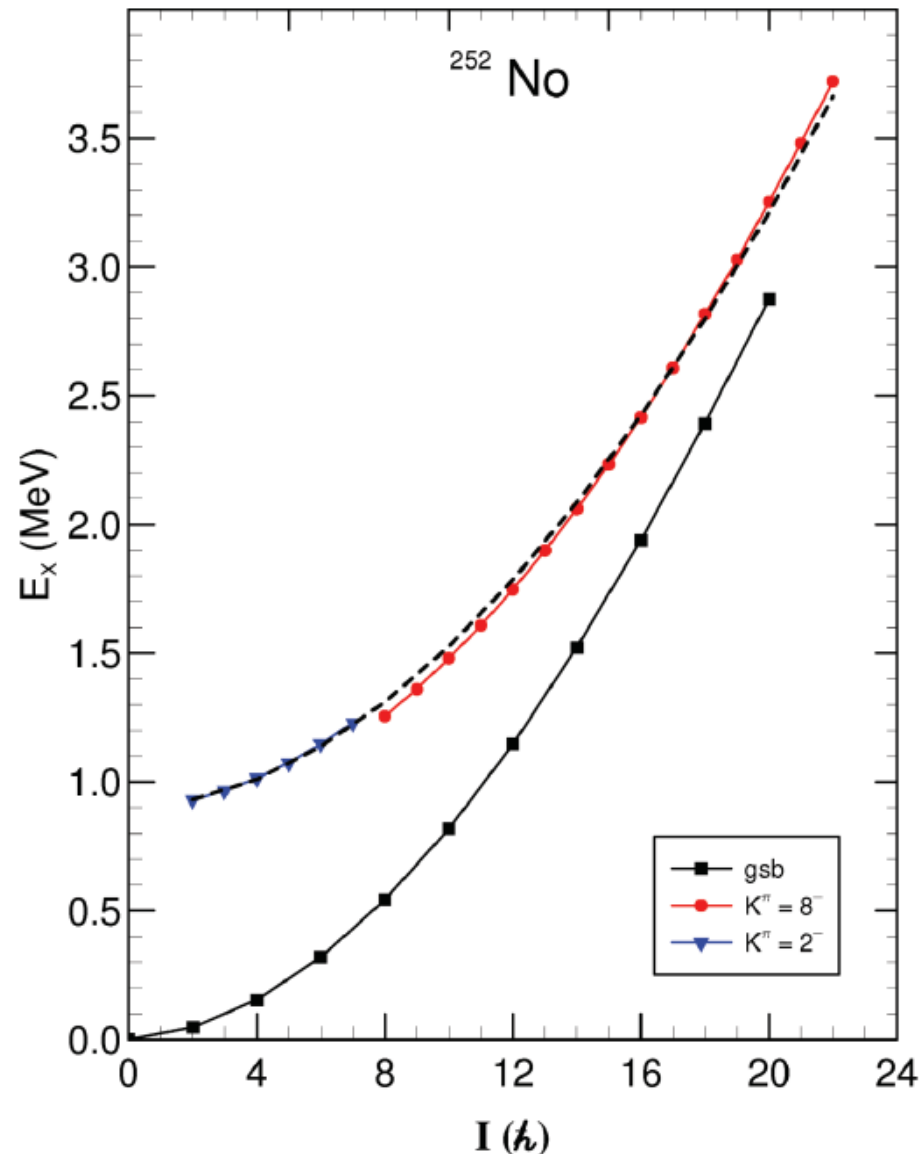


FIG. 6. (Color online) (a) and (b) Dynamic moment of inertia $\mathcal{J}^{(2)}$ vs $\hbar\omega$ for the ground-state rotational band (dashed black lines) and for the isomeric $K^\pi = 8^-$ band, for the ^{250}Fm and ^{252}No (red triangles) isotones. (c) and (d) Same as (a) and (b) but for the kinematic moments of inertia. Triangular blue marks represent theoretical calculations using the DIS Gogny force. See text for details.

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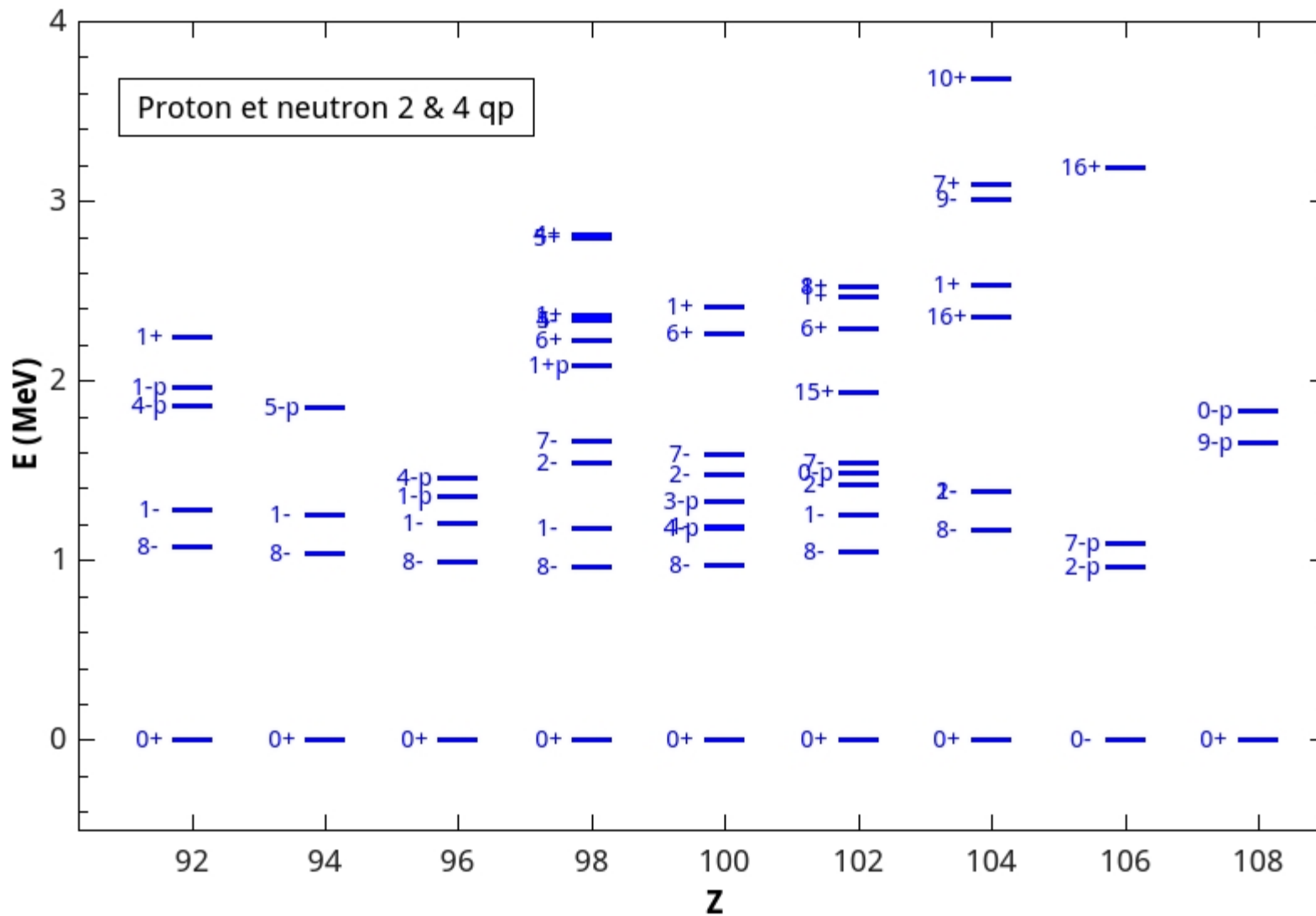
Charge multipole moments in the 2QP 8^- rotating nucleus



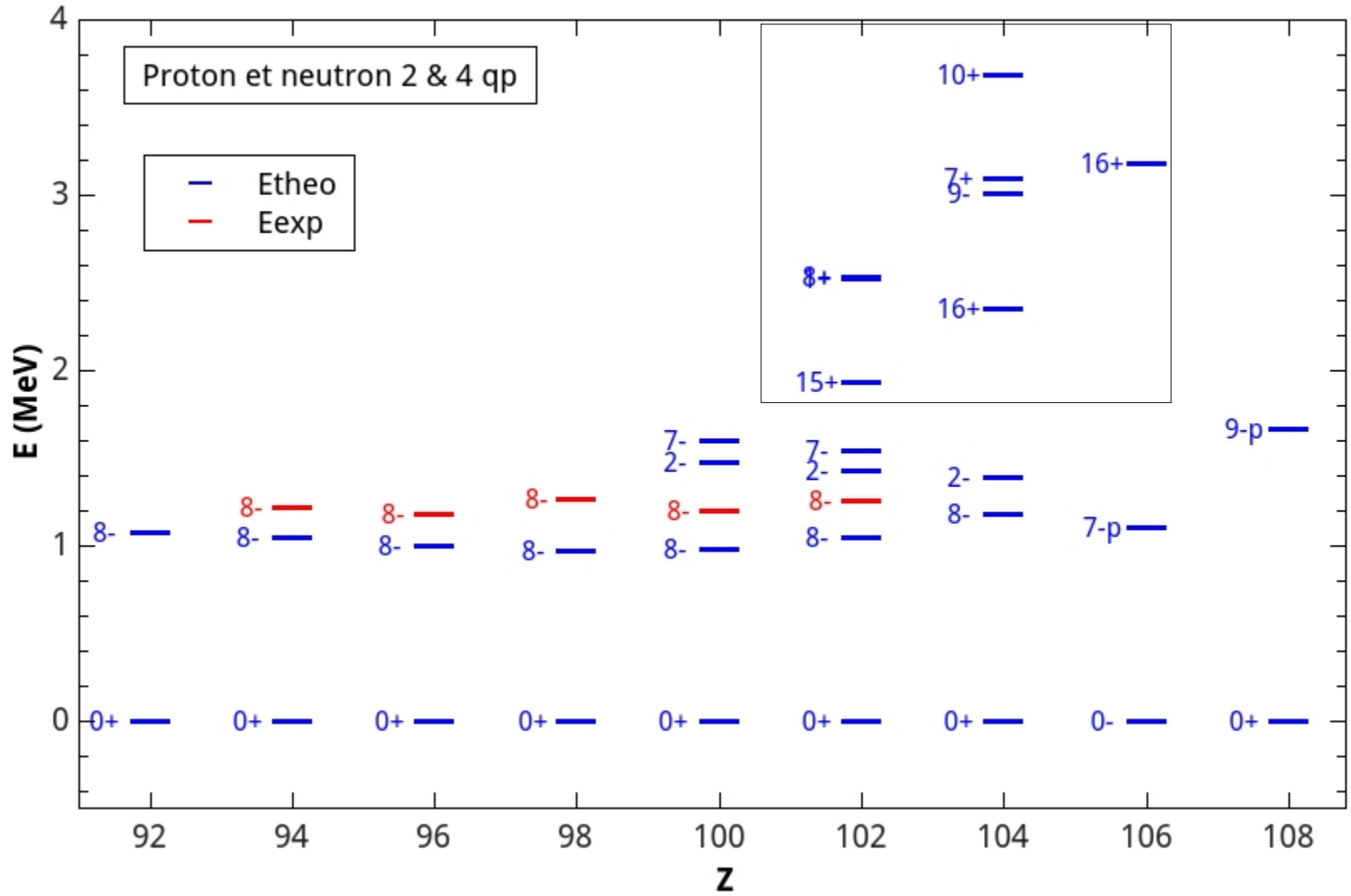


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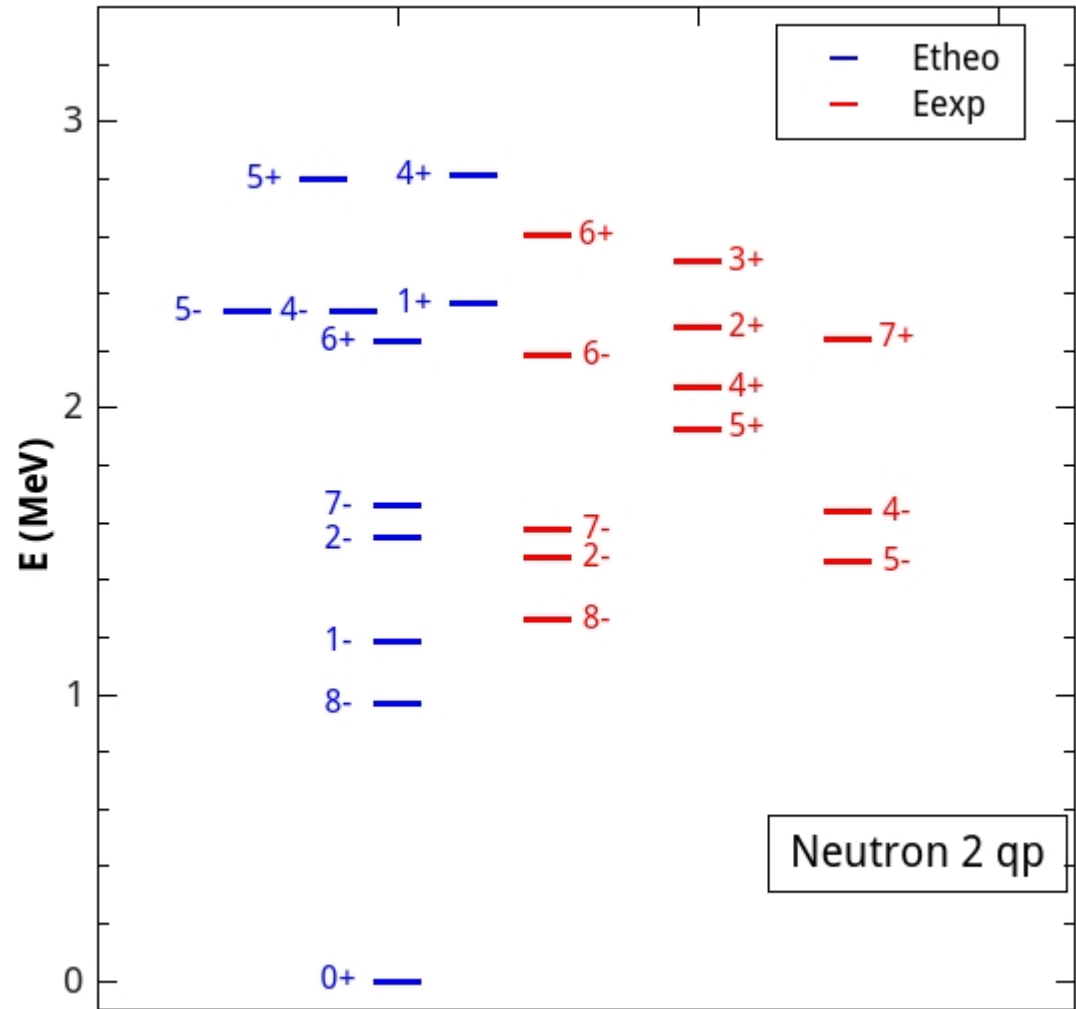
N=150



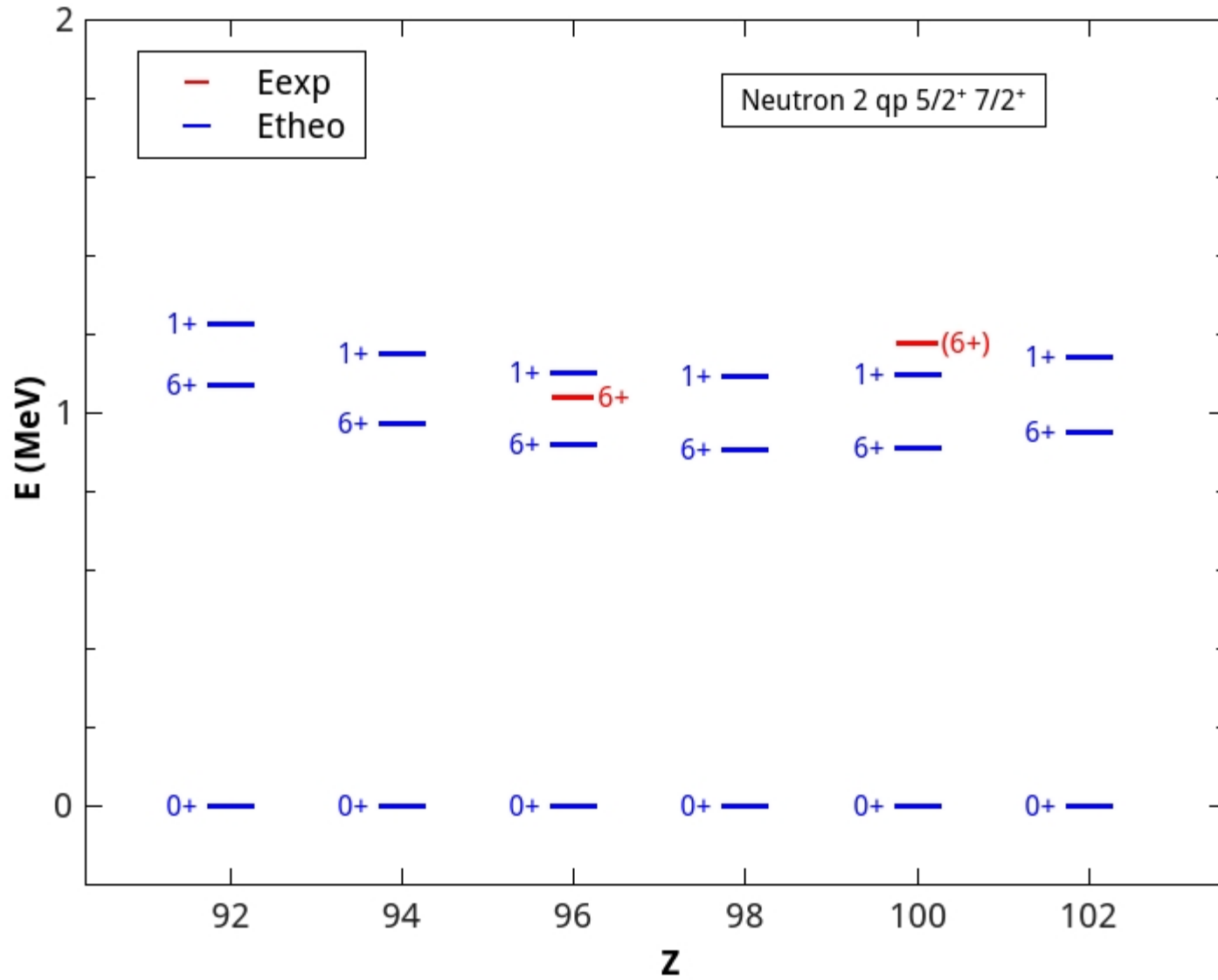
N=150



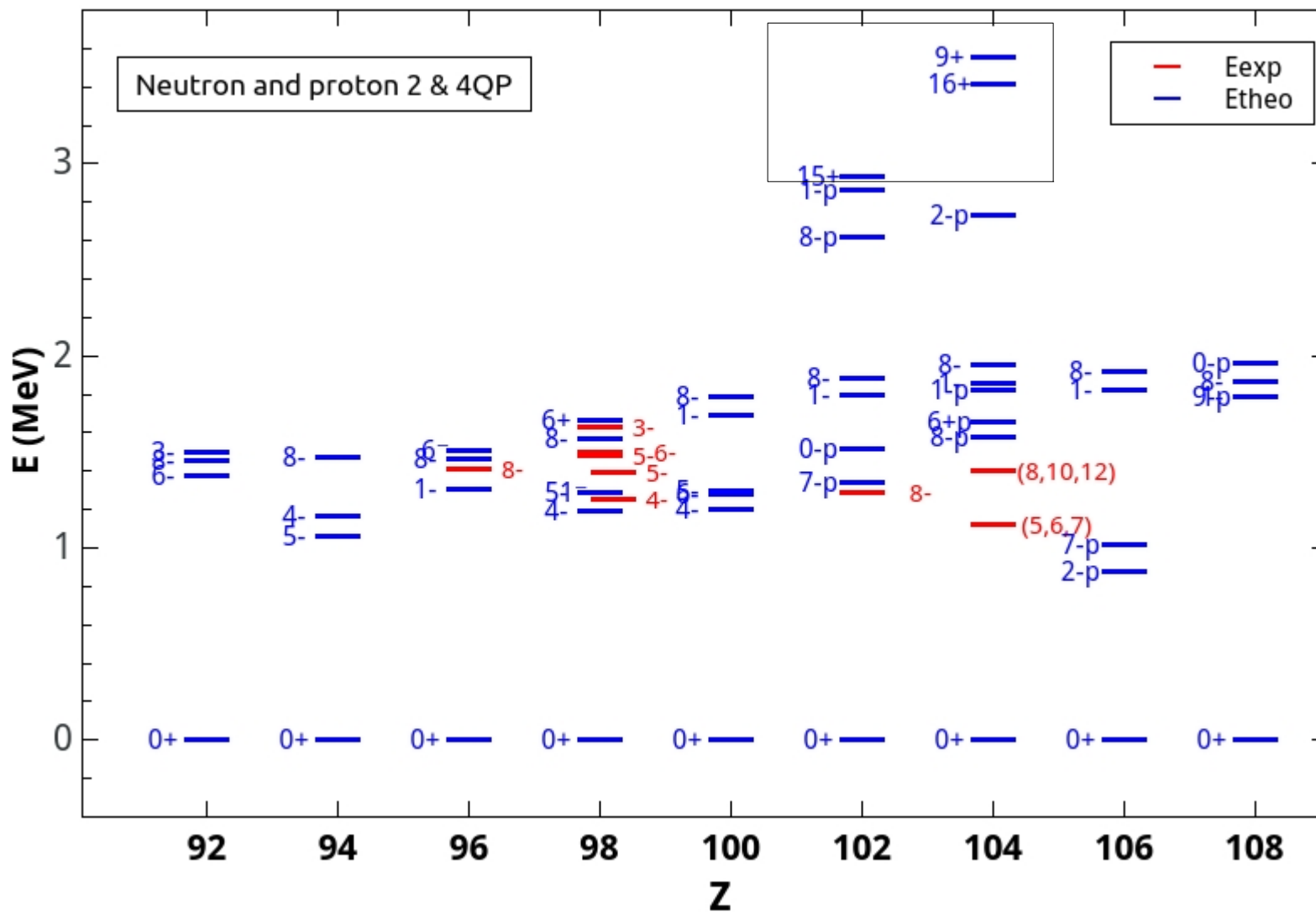
$^{248}_{98}\text{Cf}_{150}$



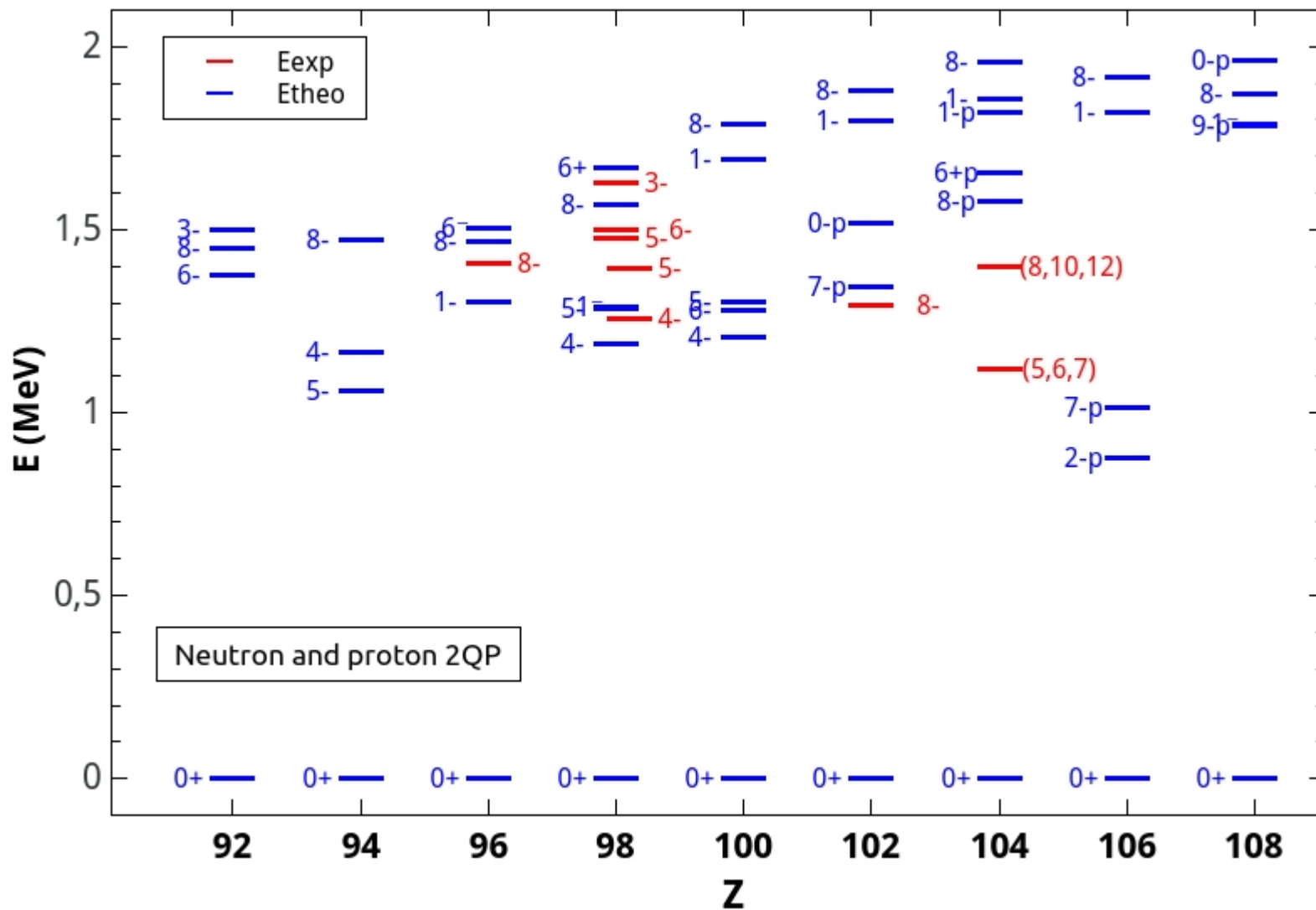
N=148



N=152



N=152

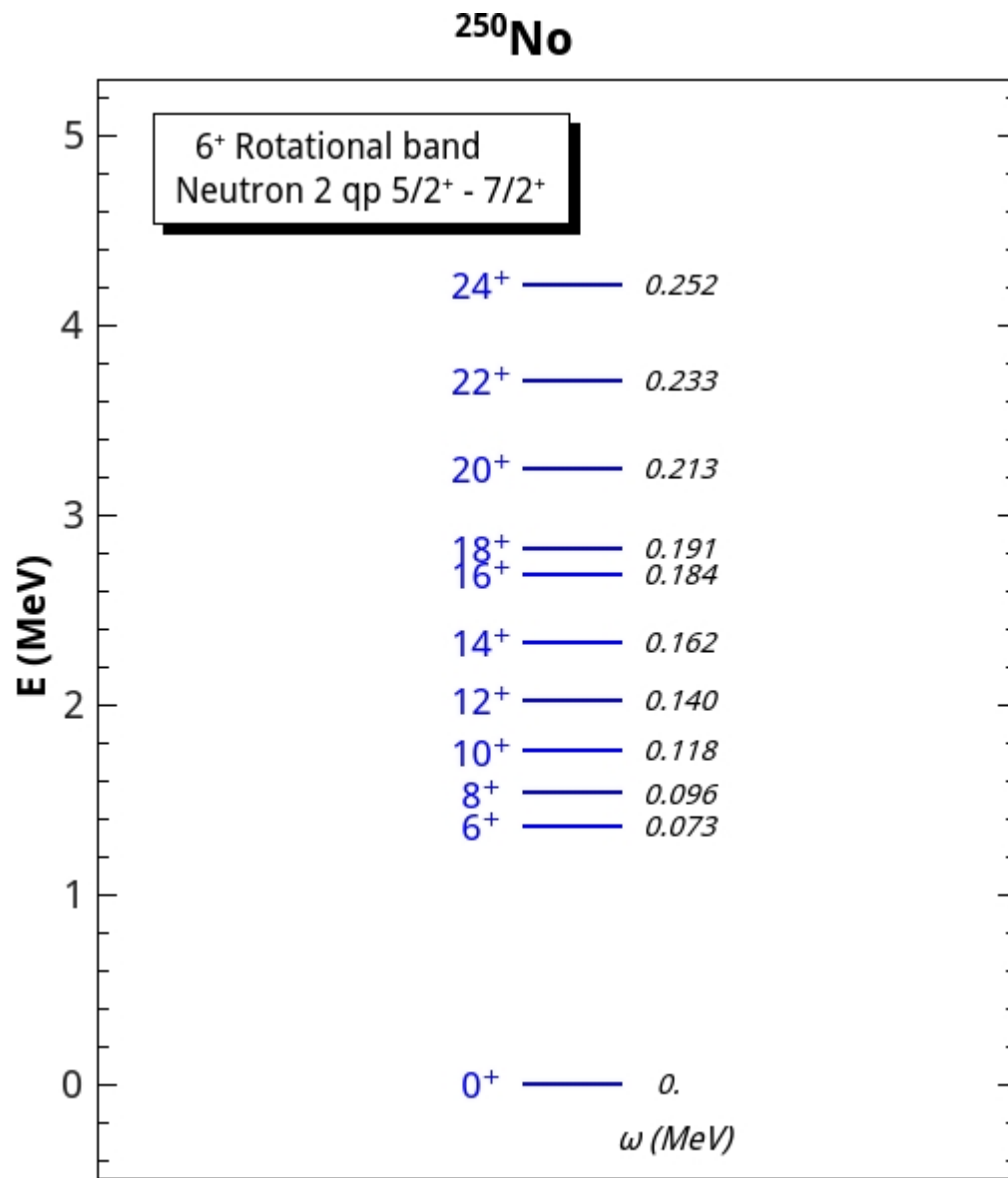


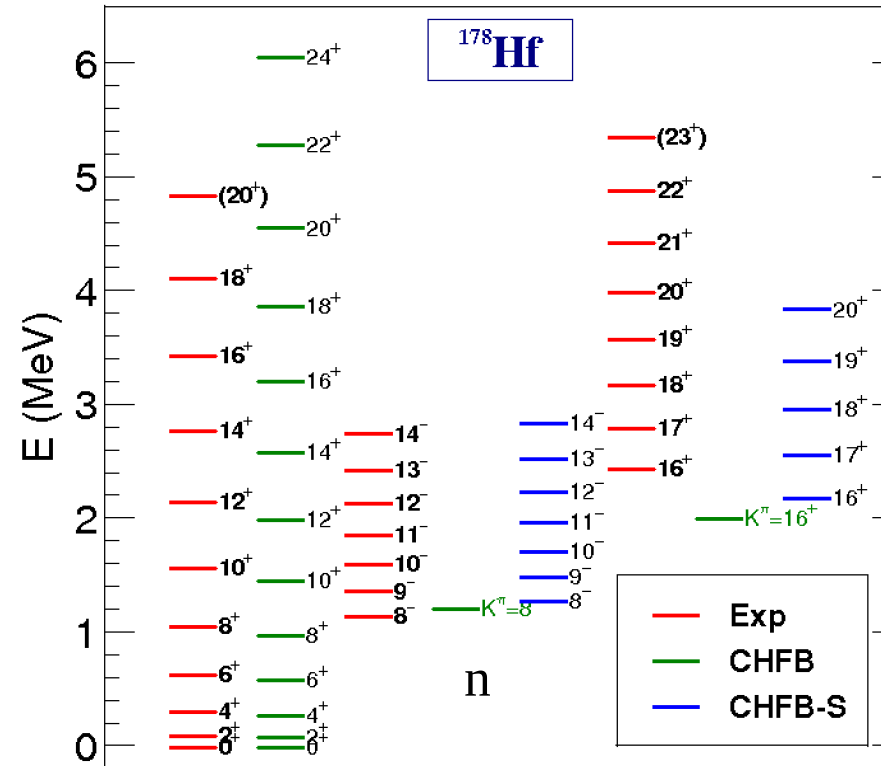
Theoretical rotational band

Quasiparticle excitations

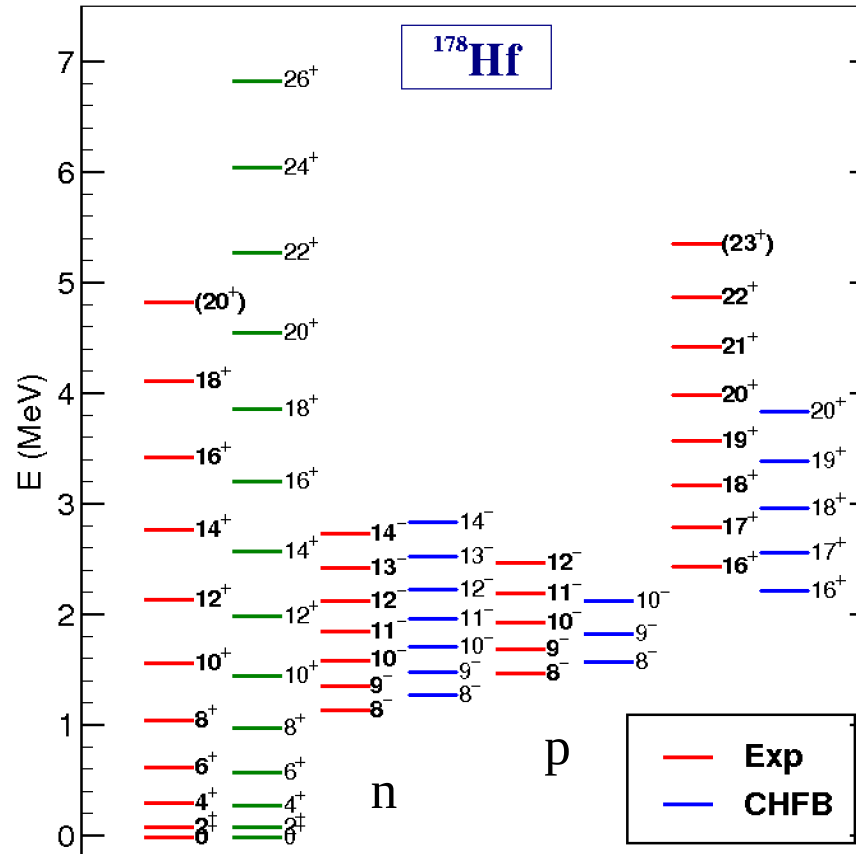
Rotation along the x-axis does break $S_z = ie^{-i\pi J_z}$, the z-signature symmetry

The z-signature is broken because J_x and S_z do not commute.





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Ph.D. thesis



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Conclusions and perspectives

- ♦ 2QP and 4QP calculated at the mean-field level;
- ♦ Fair agreement with data; $N=148, 150, 152$;
- ♦ 2QP and 4QP rotation: underway;
- ♦ Galagher rule not always fulfilled...
- ♦ HFB-D1S blocking 2QP \longrightarrow QRPA-D1S
- ♦ Breaking parity : underway;
- ♦ Including particle-vibration coupling will change “single particle” level schemes (see R.V.Afanasjev and E. Litvinova, PRC 92, 044317 (2015));
- ♦ Are blocking calculations still pertinent?