

Green's functions for nuclear matter

Correlations, equation of state and pairing gaps

+A. Carbone



TECHNISCHE
UNIVERSITÄT
DARMSTADT Alexander von Humboldt
Stiftung/Foundation



+D. Ding, W. H. Dickhoff, H. Dussan
arXiv:1502.05673



+A. Polls



UNIVERSITAT DE BARCELONA



+C. Barbieri

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Department of Physics
University of Surrey

Saclay, 2 April 2015

➔ **Motivation**

➔ *Nuclear matter: Equation of state with 3NFs*

➔ *Neutron matter: beyond-BCS pairing*

Nuclei

- Finite size $\Rightarrow (N, Z)$
- Surface effects
- Single-particle wavefunction?

Nuclear matter

- Infinite matter $\Rightarrow \rho$
- No surface
- Plane waves
- Thermodynamic limit

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Nuclei



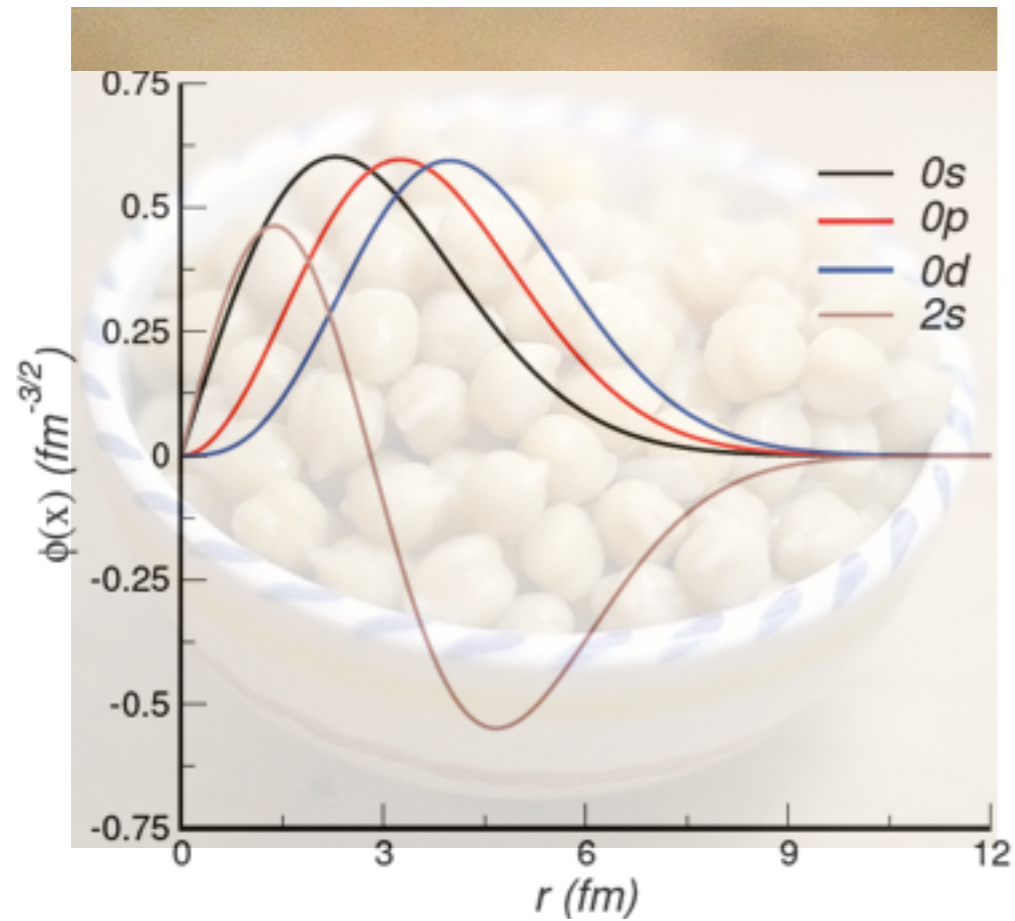
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Nuclear matter



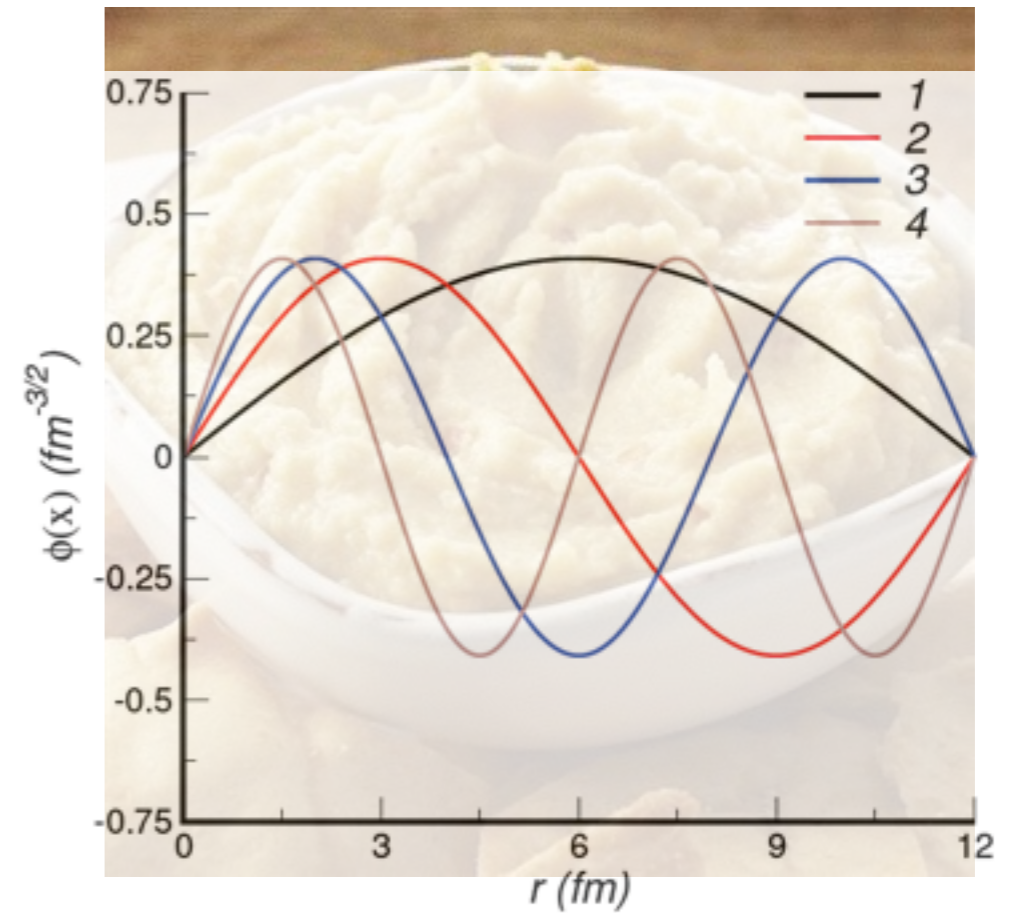
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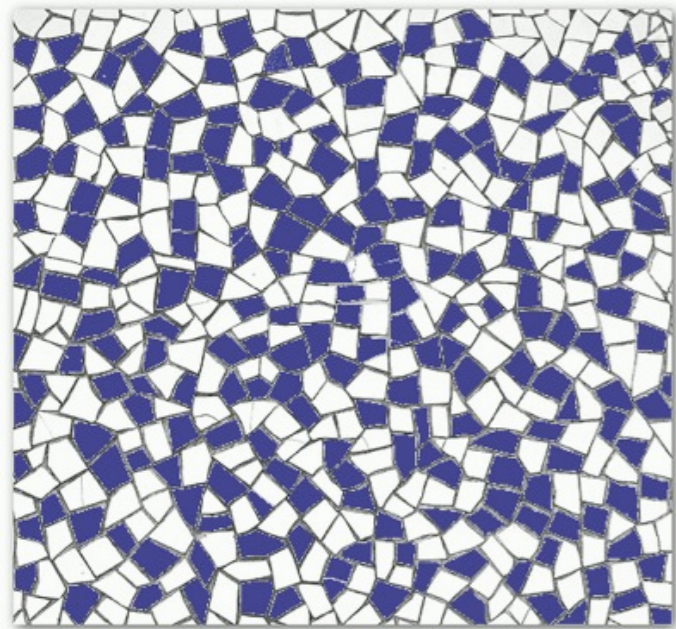


- Infinite matter $\Rightarrow \rho$
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Nuclear “trencadís”

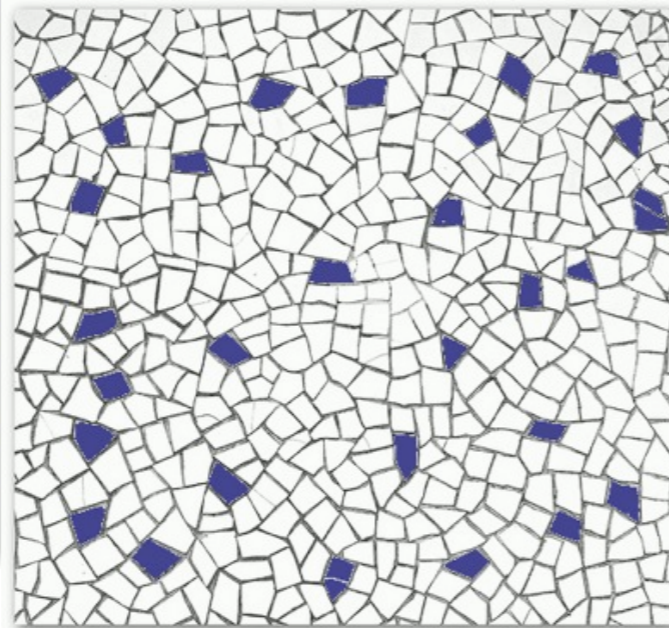
$N=Z, \beta=0$

Symmetric matter



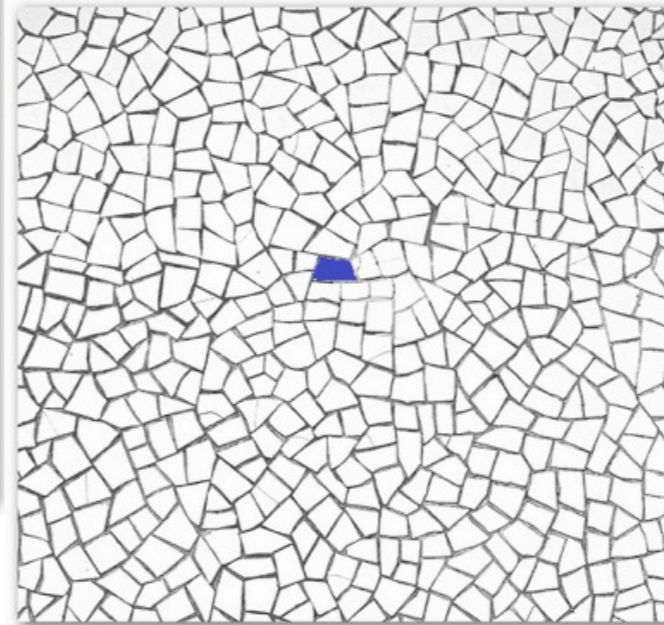
$\beta \neq 0$

Asymmetric nuclei



$\beta \approx 1$

Polaron

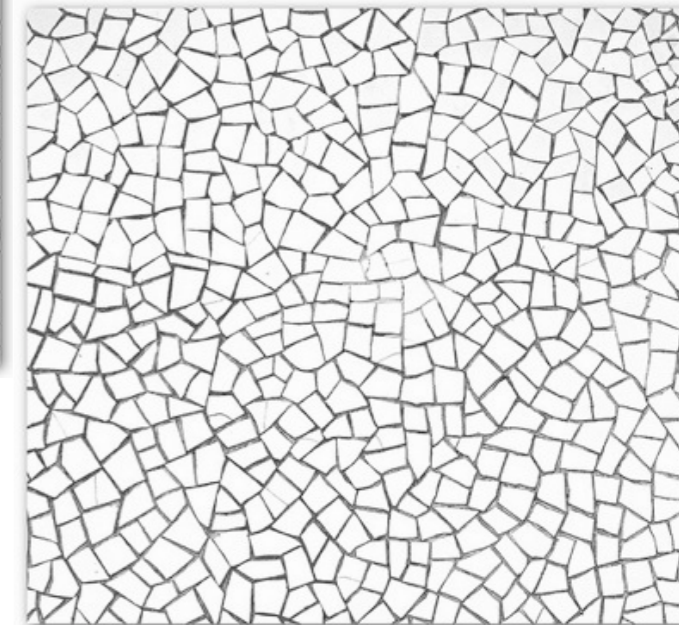


Neutron stars

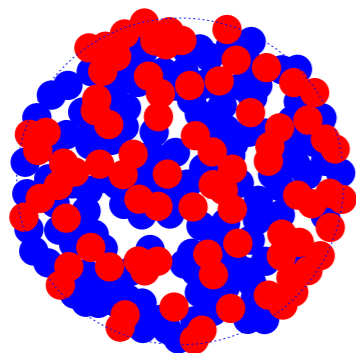


$\beta=1$

Neutron matter



Nuclei



Lead 208

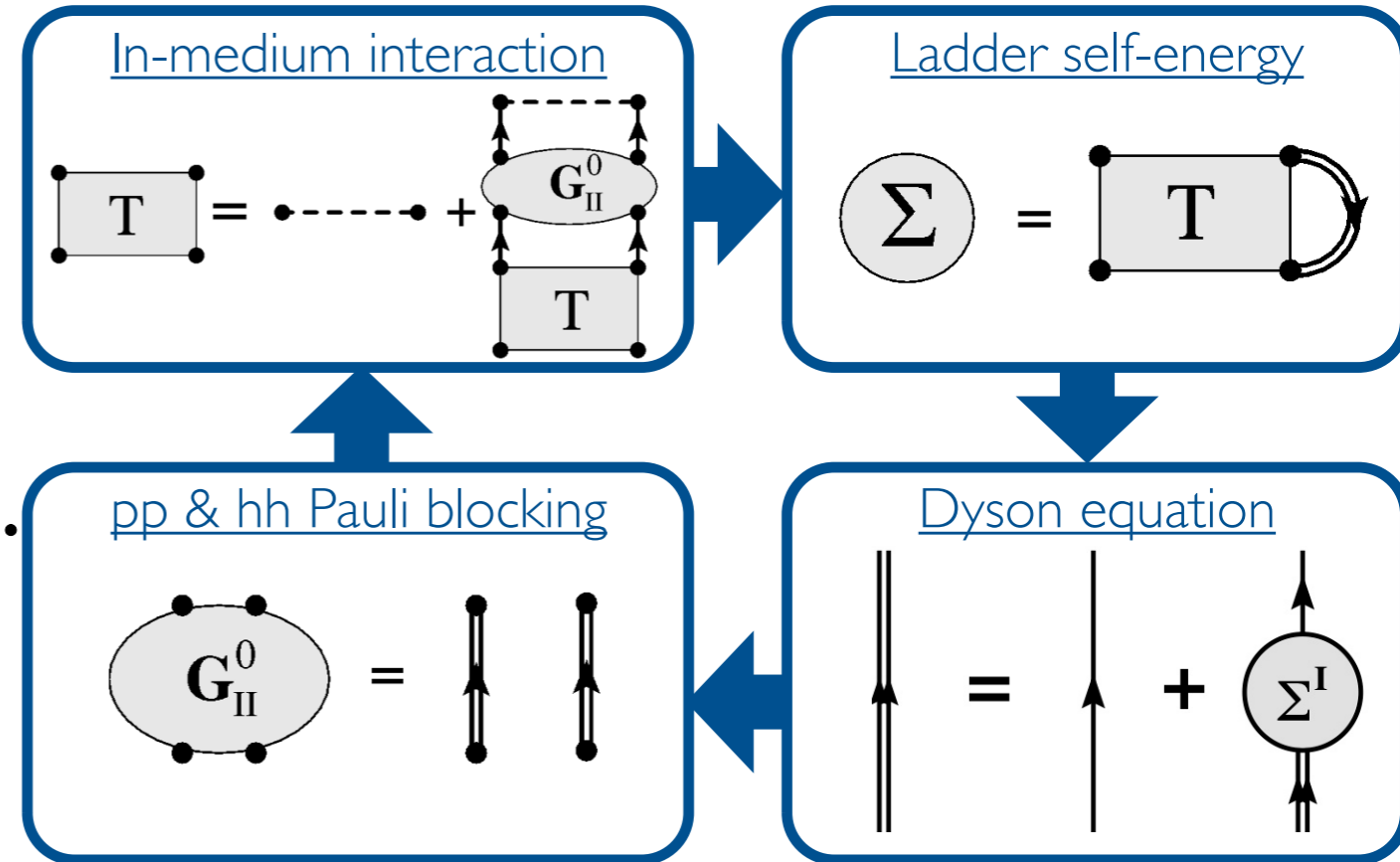
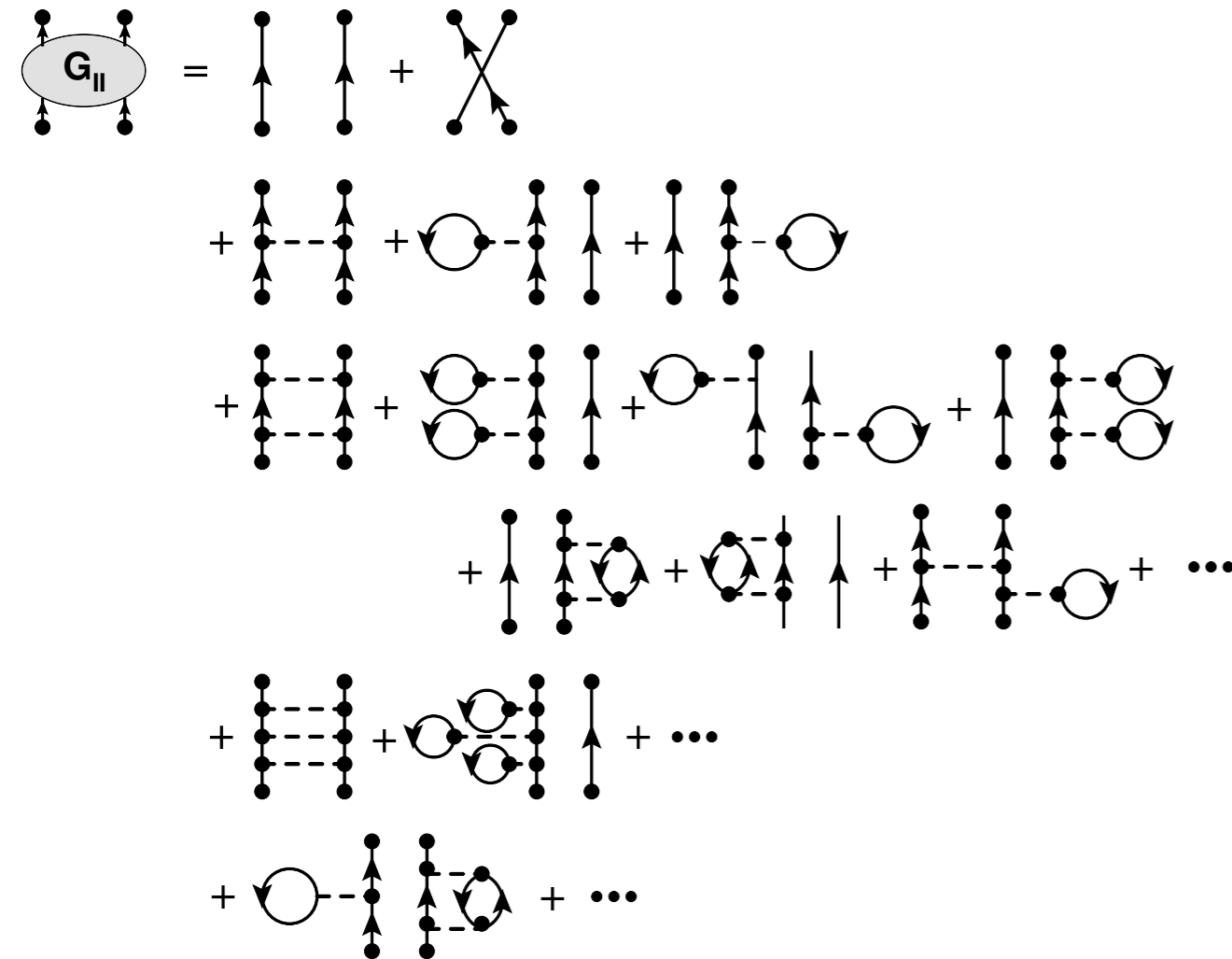
$$\beta = (126 - 82) / 208 = 0.2$$

$$\beta = \frac{N - Z}{N + Z}$$

Experimentally unknown!

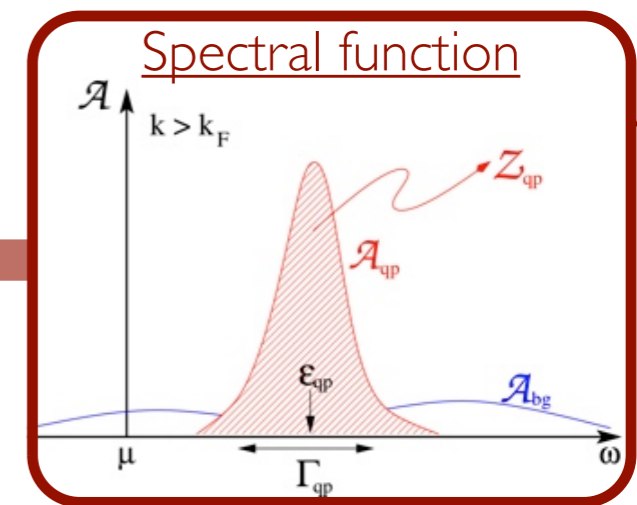
Self-consistent Green's functions

Ladder approximation within SCGF



Ramos, Polls & Dickhoff, NPA **503** 1 (1989)
 Alm *et al.*, PRC **53** 2181 (1996)
 Dewulf *et al.*, PRL **90** 152501 (2003)
 Frick & Muther, PRC **68** 034310 (2003)
 Rios, PhD Thesis, U. Barcelona (2007)
 Soma & Bozek, PRC **78** 054003 (2008)

One-body properties
 Momentum distribution
 Thermodynamics & EoS
Transport

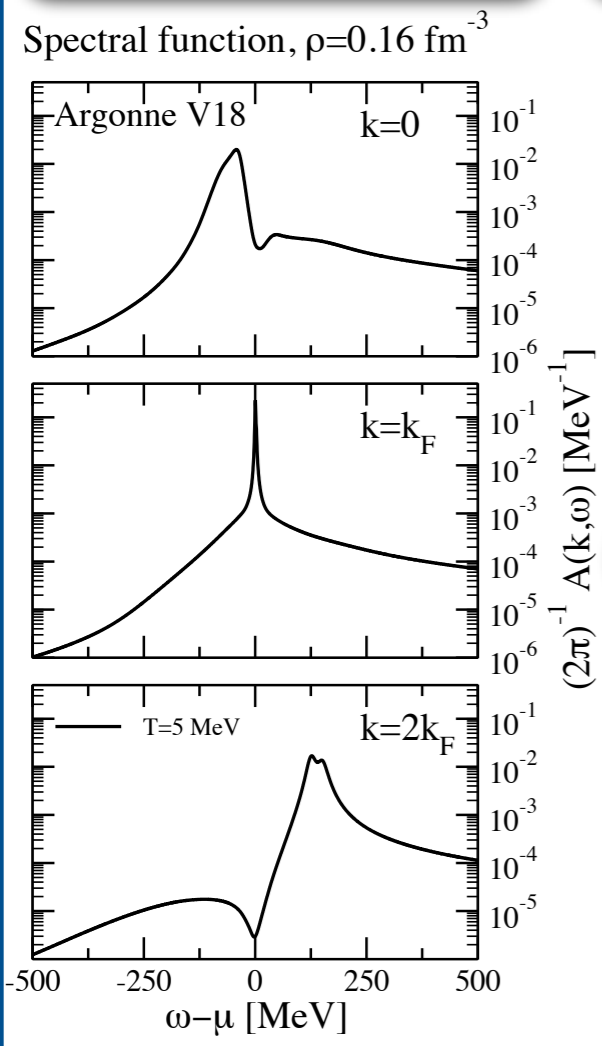


Self-consistency, pp+hh & full off-shell effects
 Finite temperature

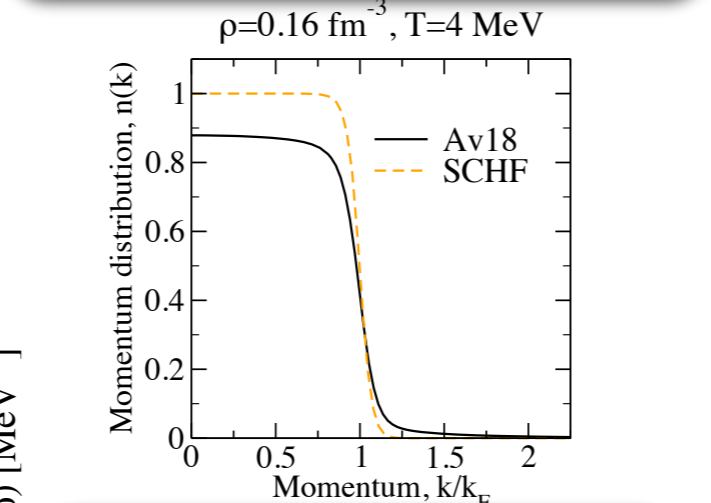
Microscopic properties

Bulk properties

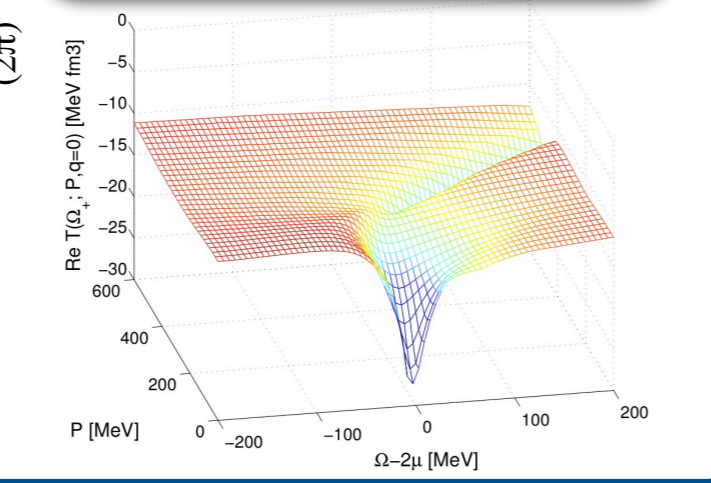
Spectral function



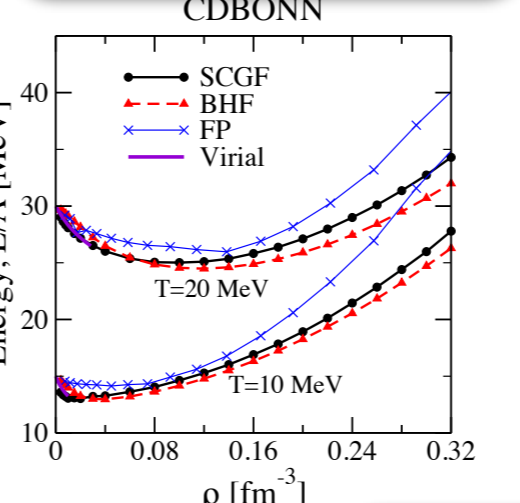
Momentum distribution



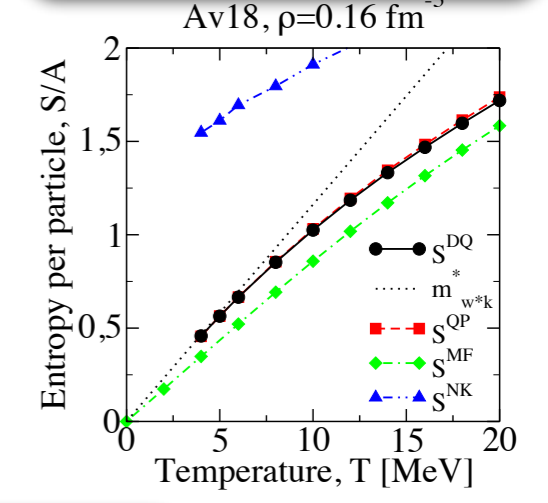
In-medium interaction



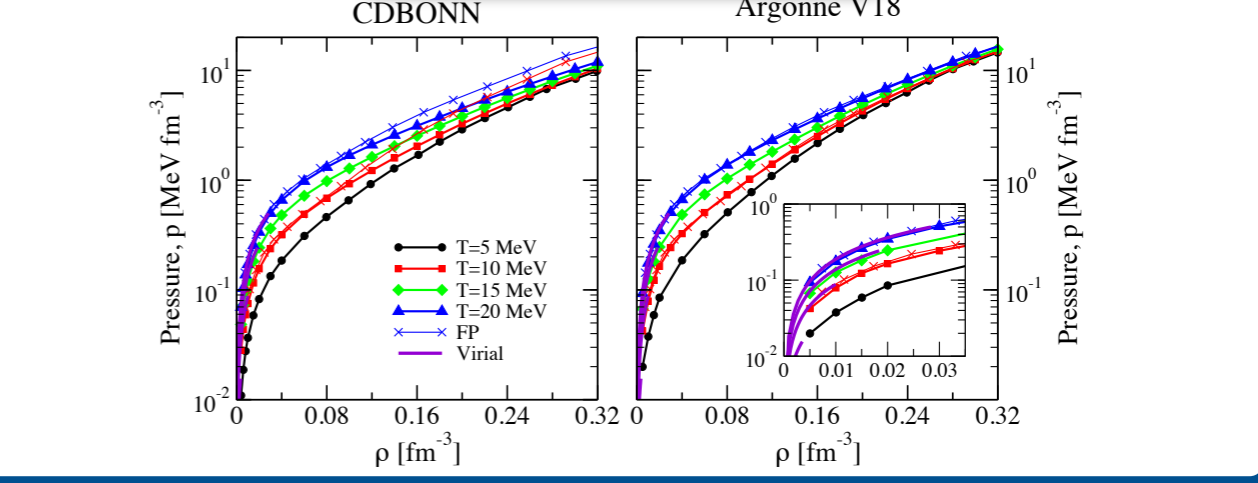
Total Energy



Entropy



Equation of State



+ Transport

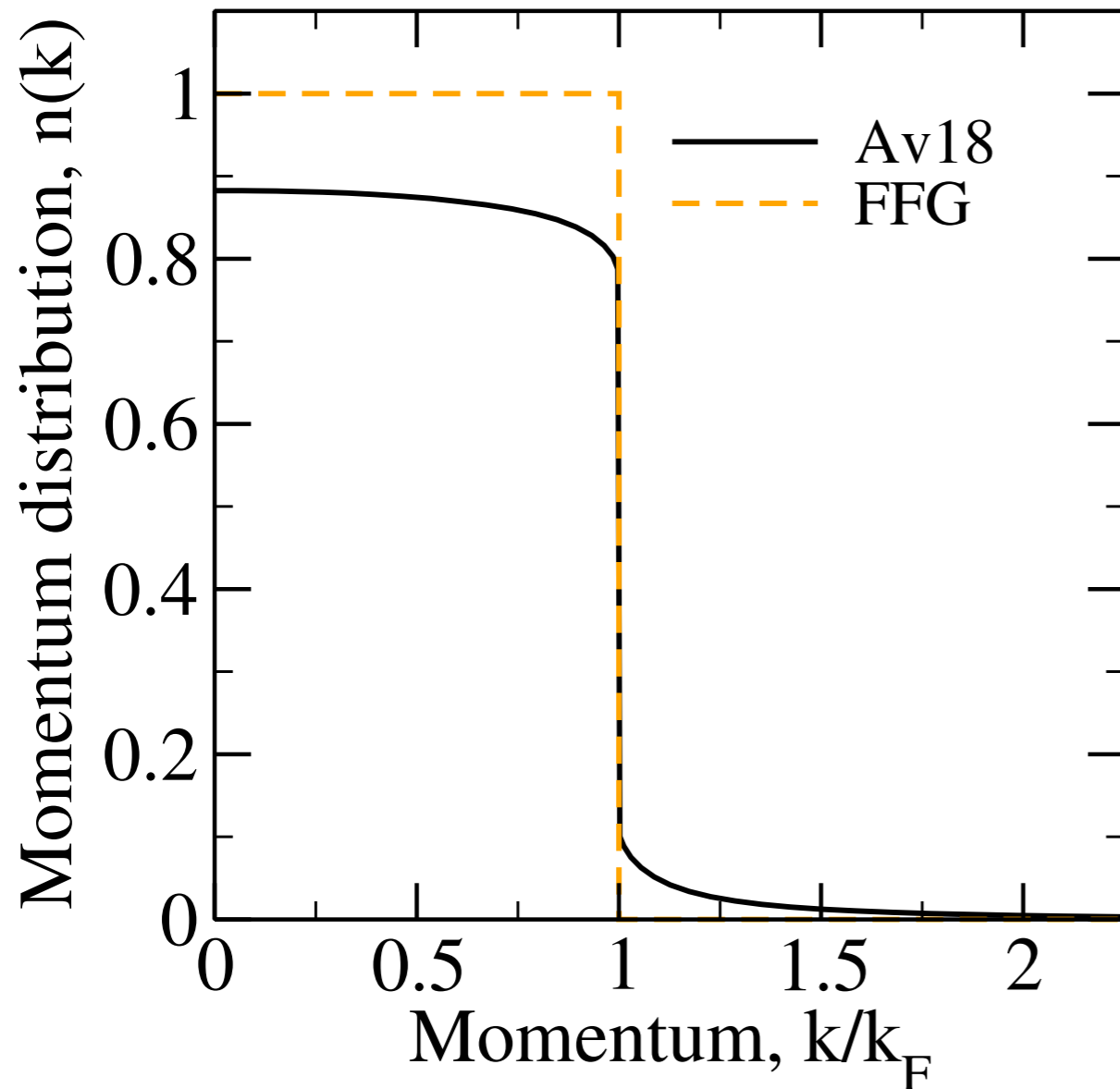
Momentum distribution in SNM

Single-particle occupation

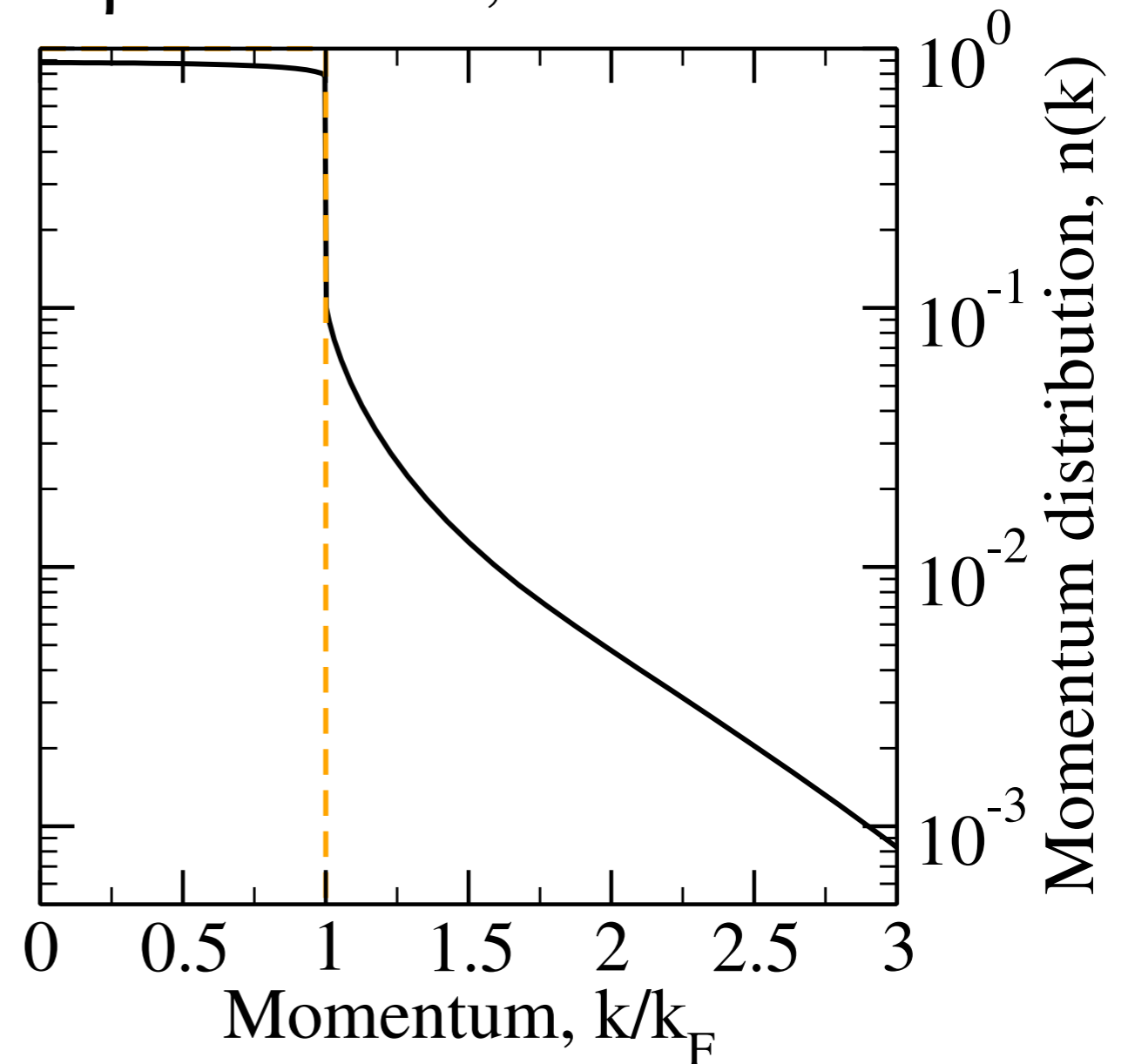
$$n(k) = \langle a_k^\dagger a_k \rangle$$

$$\nu \int \frac{d^3 k}{(2\pi)^3} n(k) = \rho$$

$\rho = 0.16 \text{ fm}^{-3}$, $T = 0 \text{ MeV}$



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- 11-13% depletion at low k , population at high k

Momentum distribution in SNM

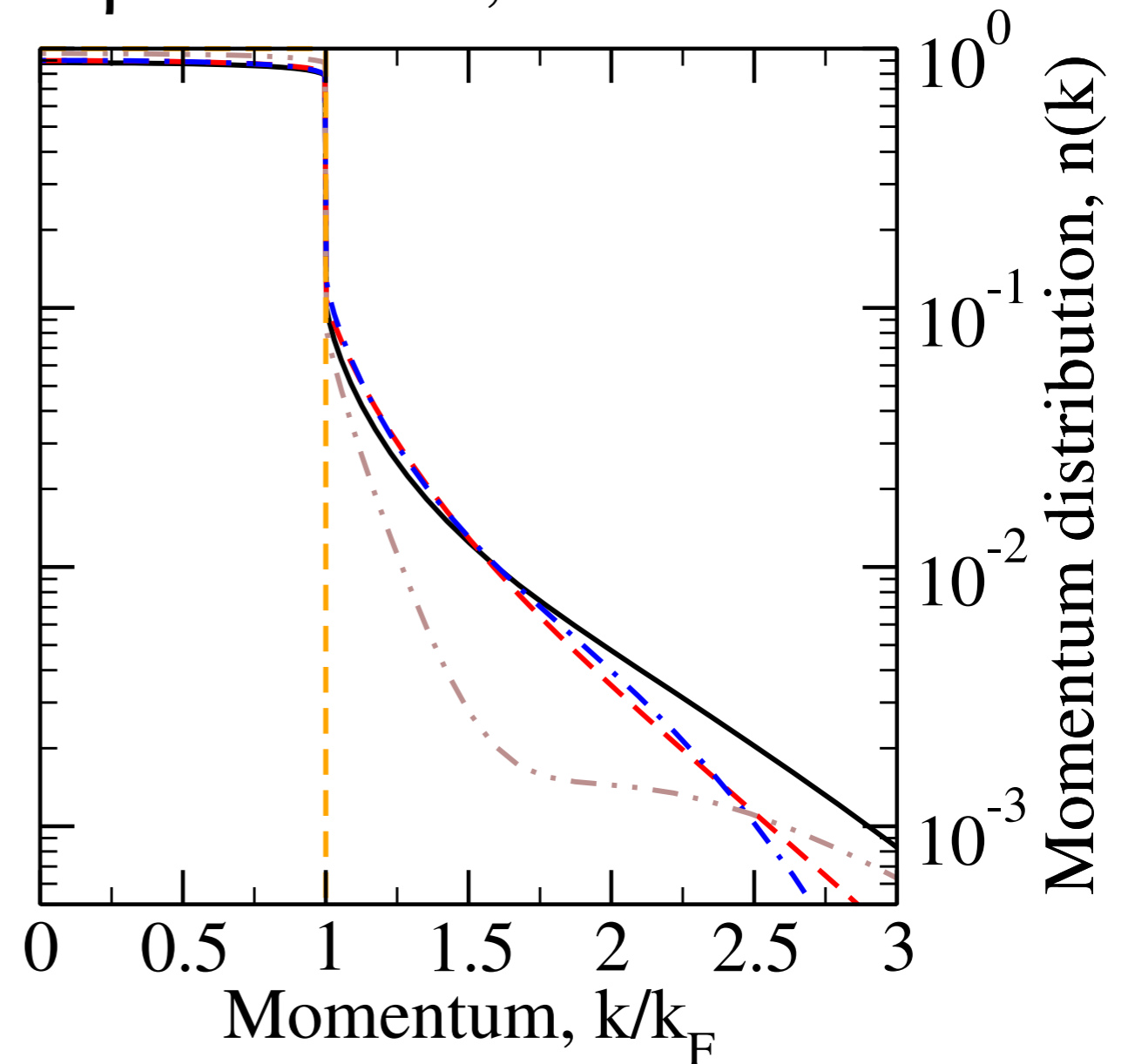
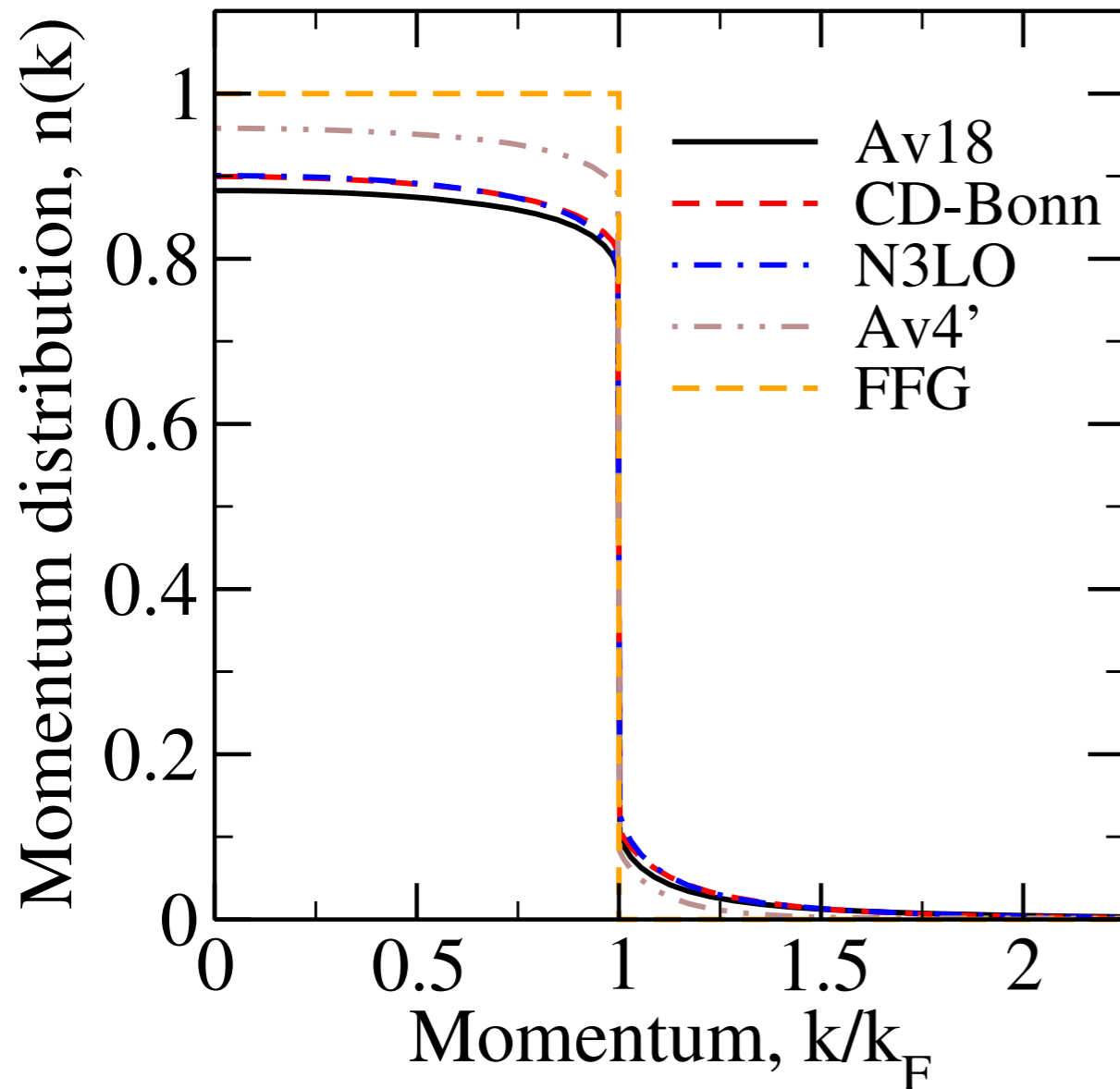
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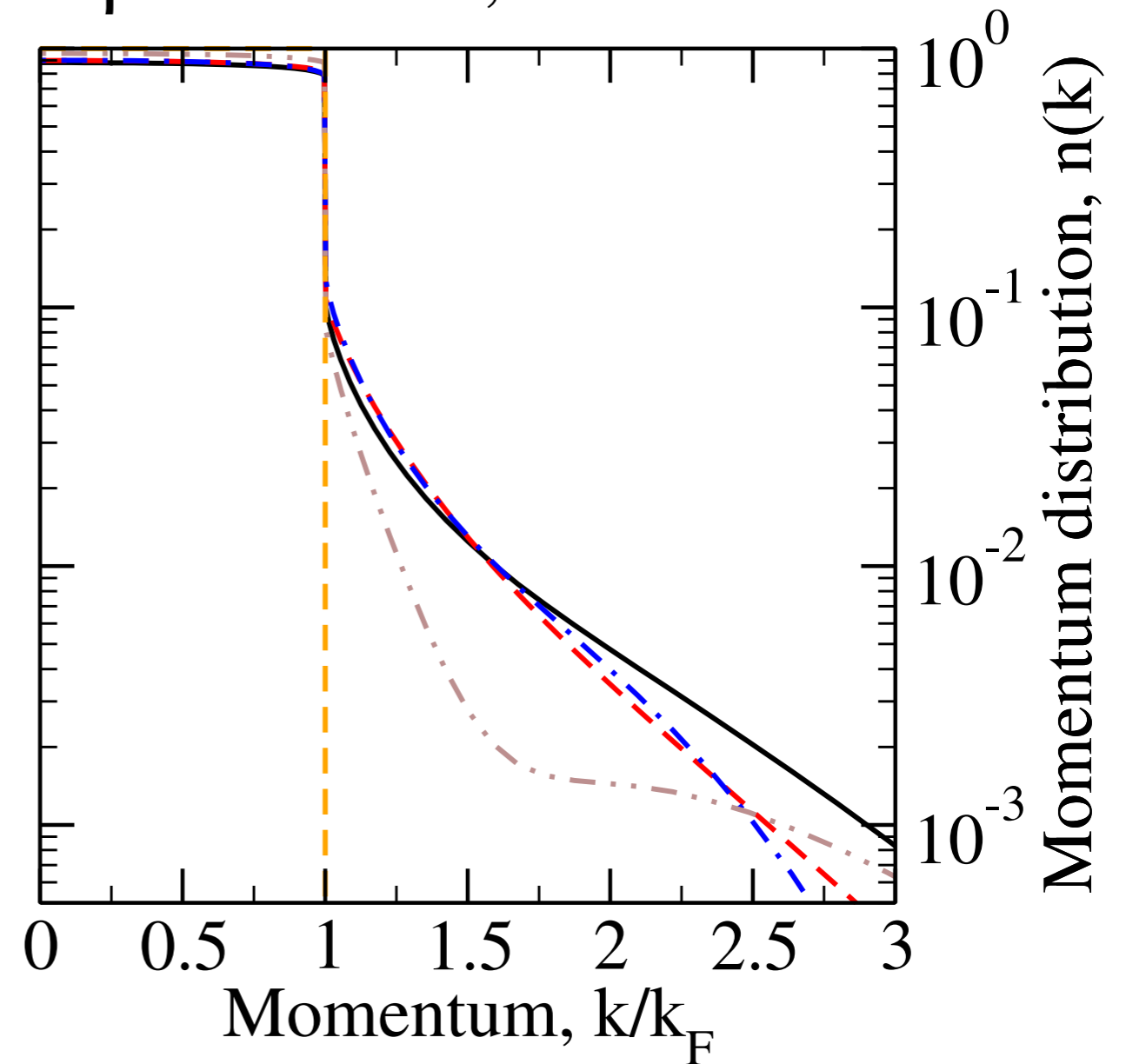
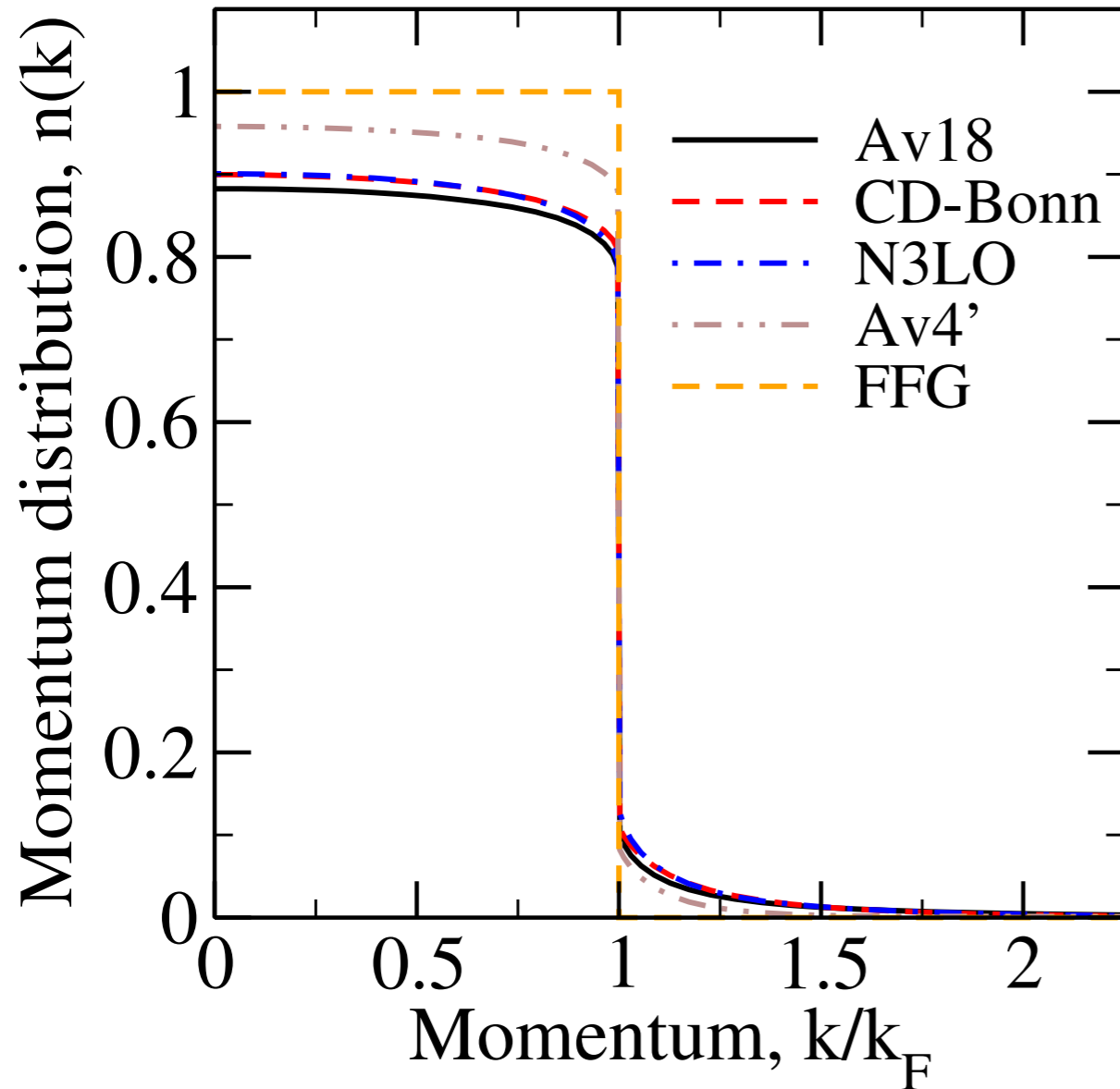
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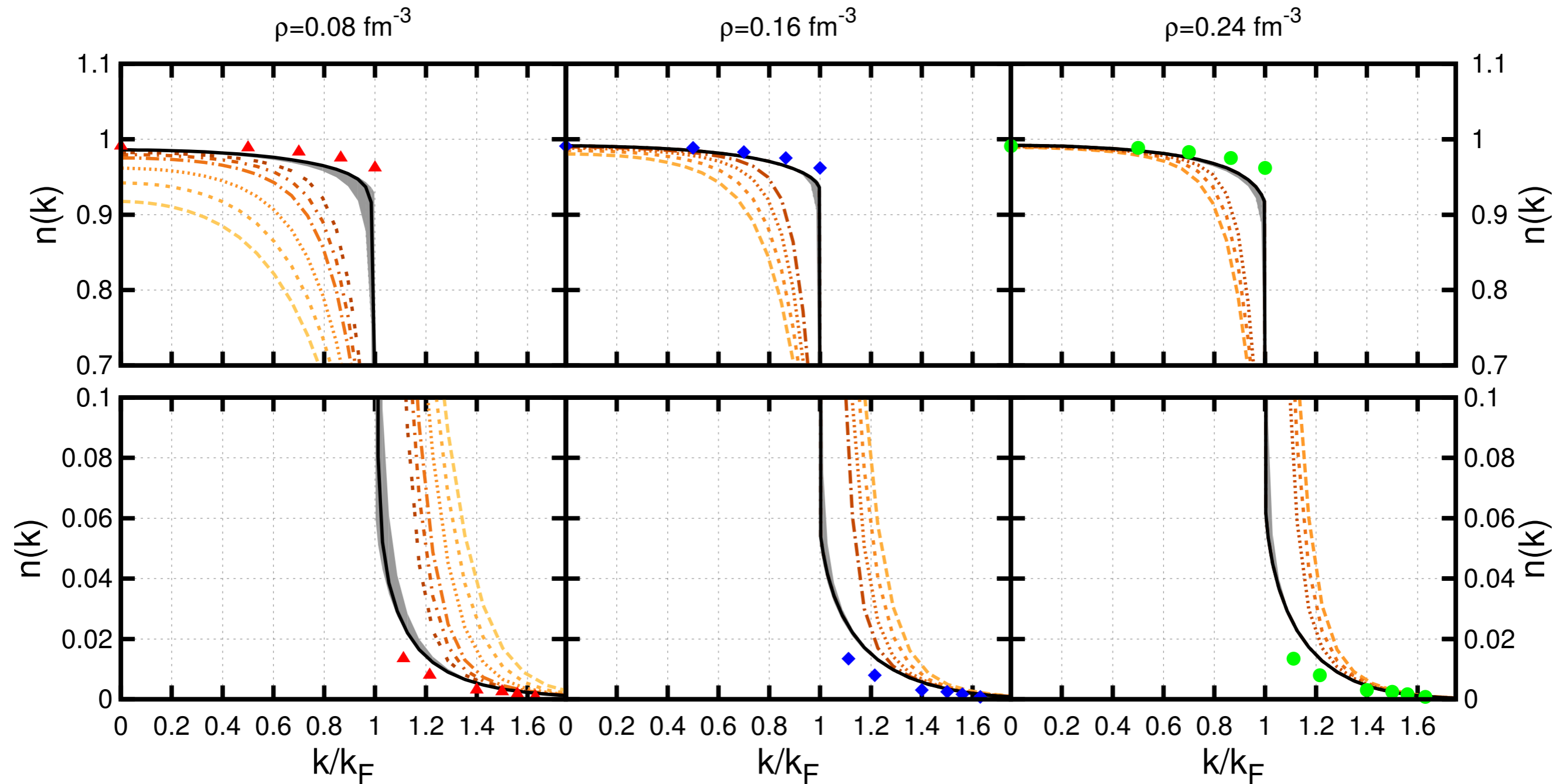
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- 11-13% depletion at low k , population at high k
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- $T=0$ extrapolated from finite T

CIMC vs SCGF

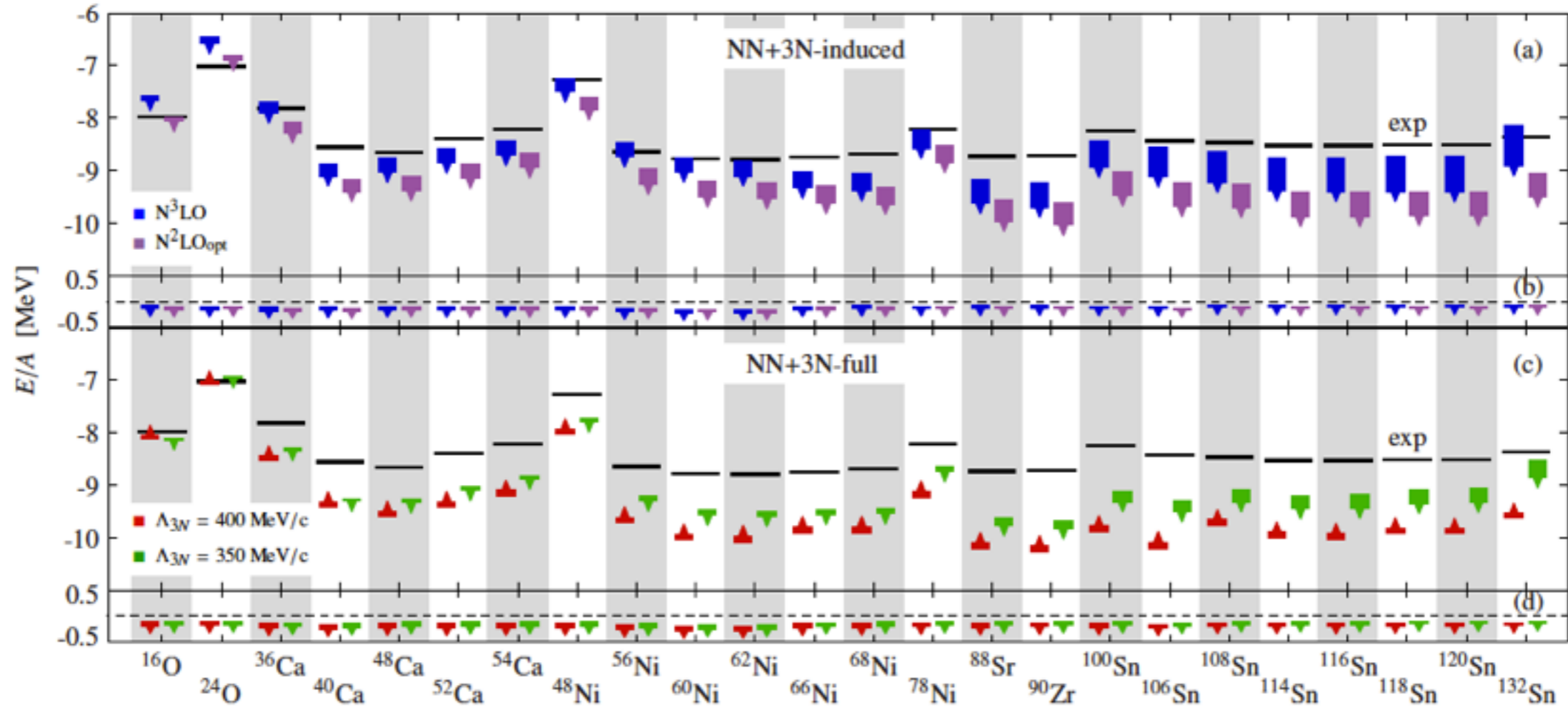


- Good agreement far from k_F
- Almost uncorrelated system at low density
- NNLO_{opt} potential

➔ *Motivation*

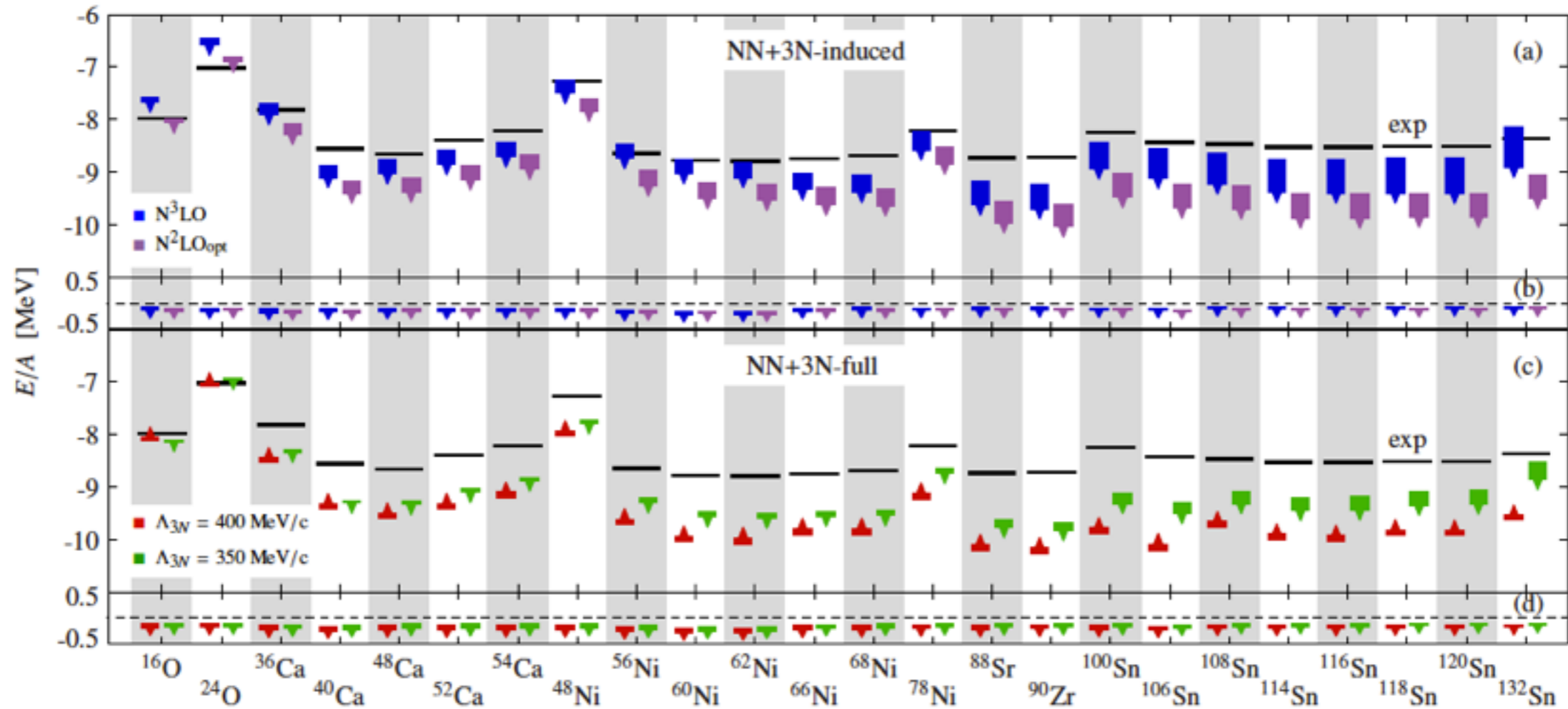
➔ **Nuclear matter: Equation of state with 3NFs**

➔ *Neutron matter: beyond-BCS pairing*

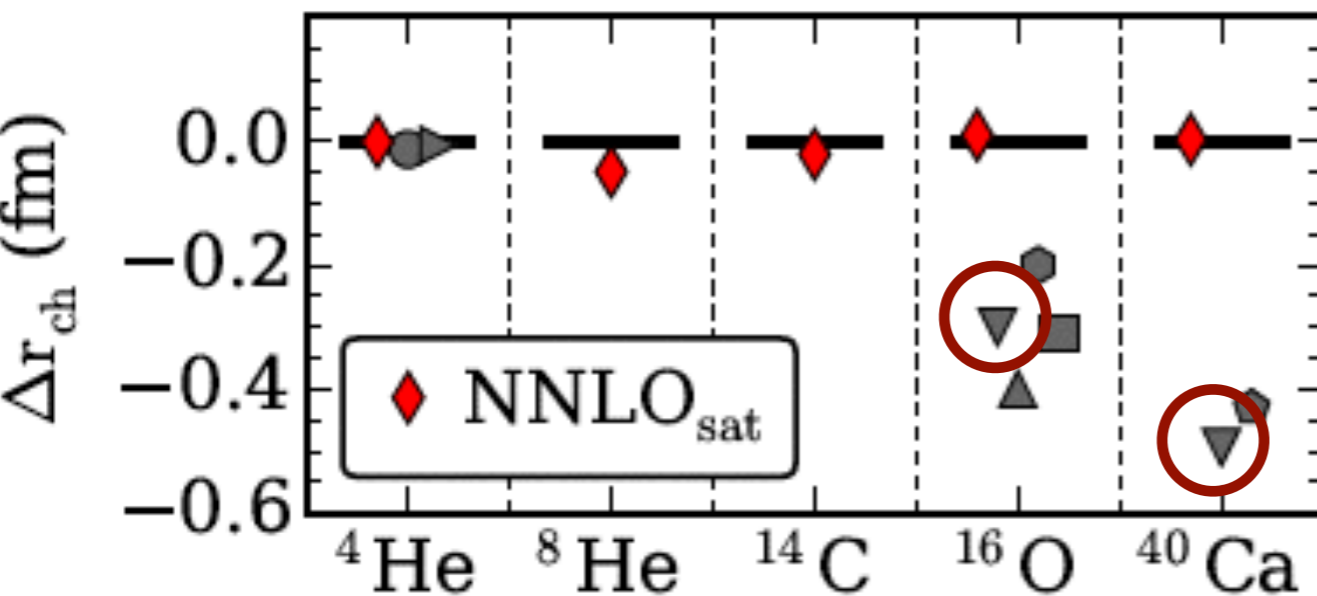


Binder, Langhammer, Calci & Roth, *PLB* **736** 119 (2014)

- Consistent calculations up to $A=132$
- Many-body (CC) errors under control
- Overbinding even when 3NF included
- Radii are **too small**



Binder, Langhammer, Calci & Roth, *PLB* **736** 119 (2014)



Ekstrom, et al. arXiv:1502.04682

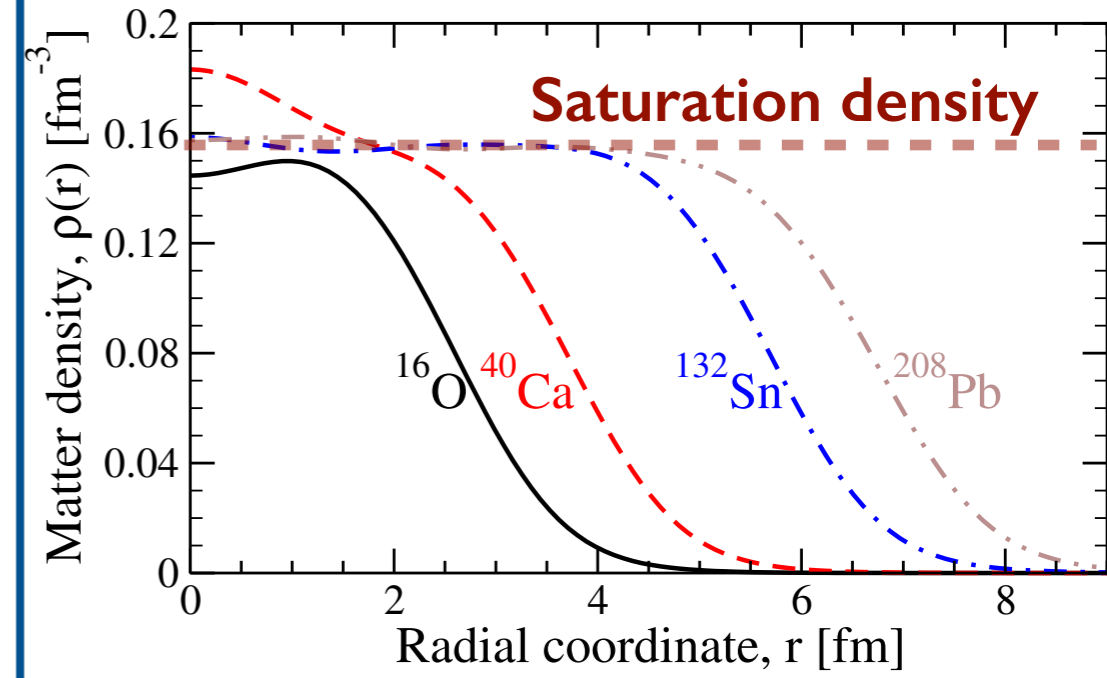
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Bethe-Weizsacker Formula

$$\frac{B}{A} = a_V + a_s A^{-1/3} + \mathcal{O}(A^{-1})$$

+

Nuclear density systematics



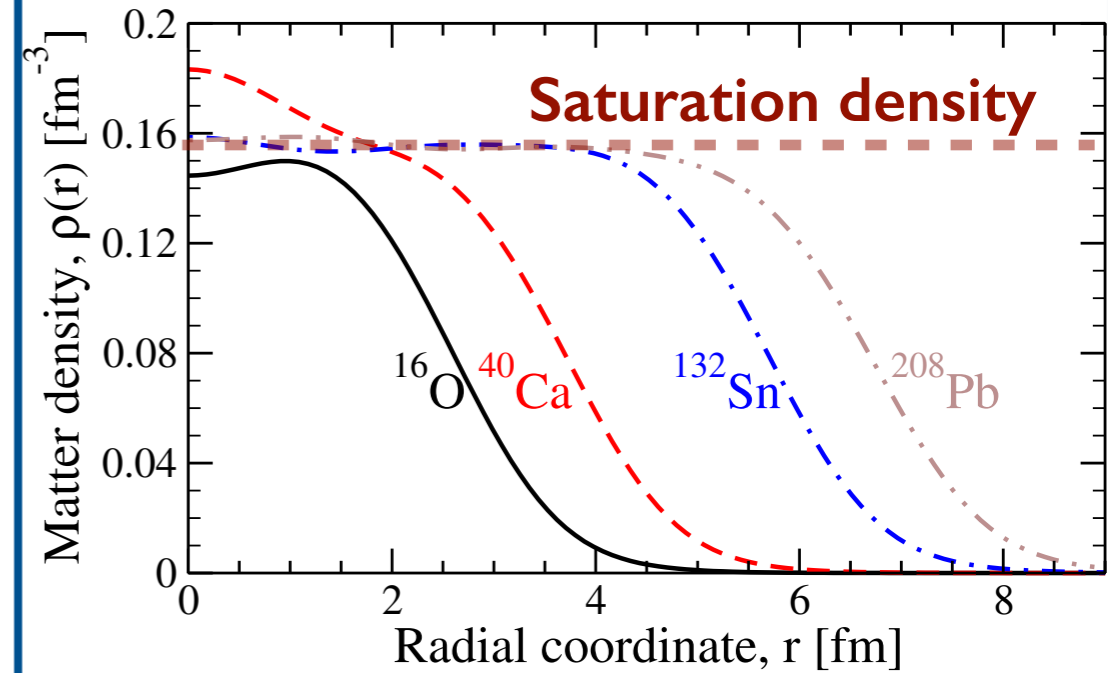
$$R \approx 1.2 \left(\frac{0.16}{\rho_0} \right)^{1/3} A^{-1/3} \text{ fm}$$

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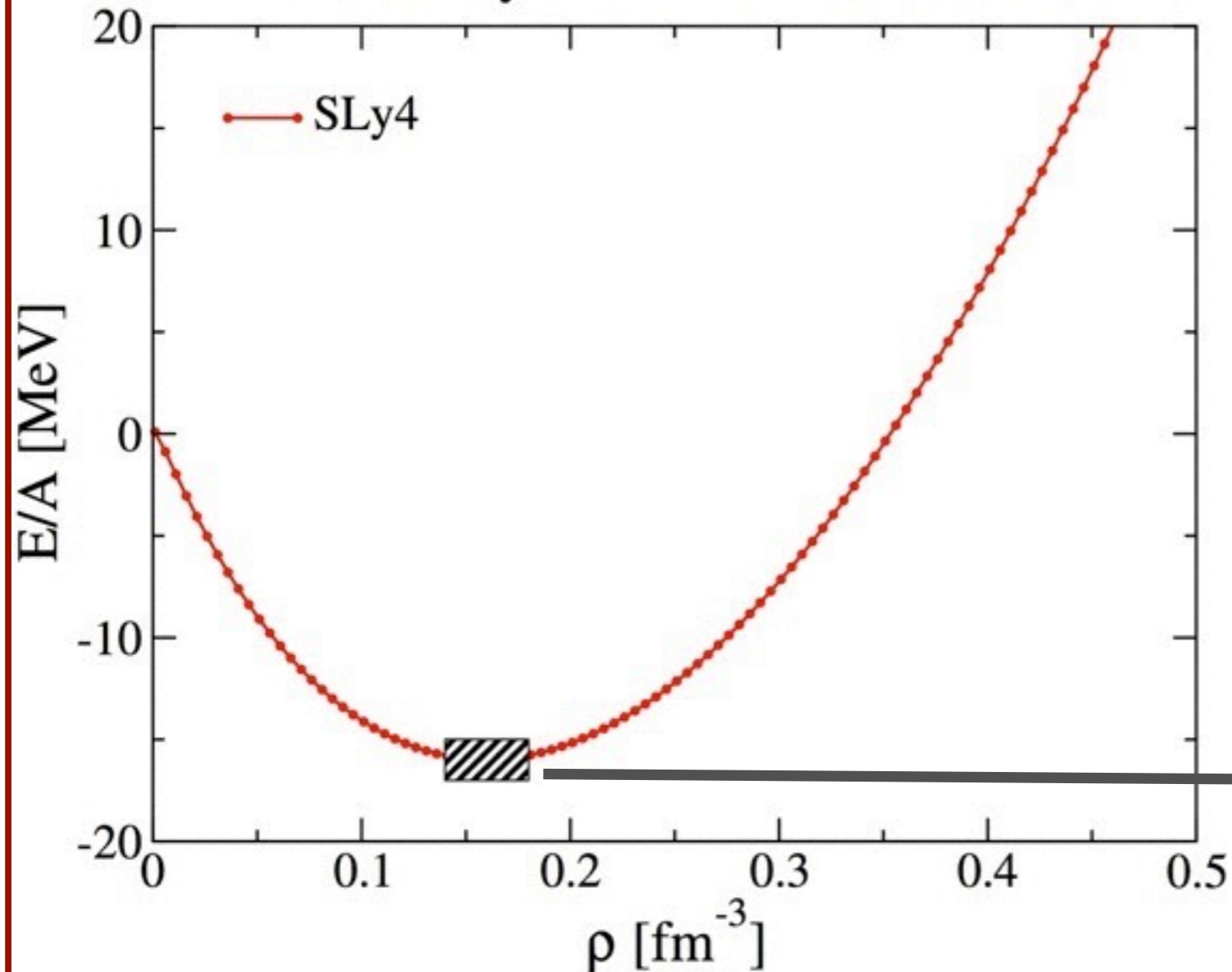
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Nuclear density systematics



$$R \approx 1.2 \left(\frac{0.16}{\rho_0} \right)^{1/3} A^{-1/3} \text{ fm}$$

Saturation of symmetric nuclear matter, T=0



Saturation point

$$\rho_0 = 0.16 \text{ fm}^{-3}$$

$$\frac{E}{A} = -16 \text{ MeV}$$

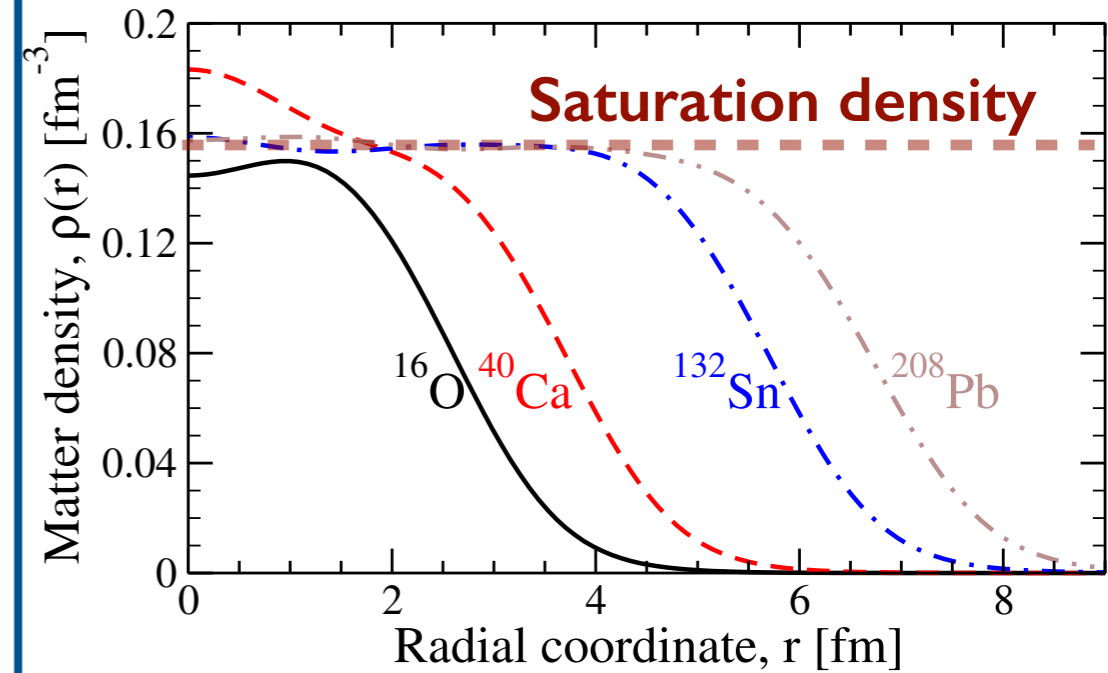
$$K = 210 \text{ MeV}$$

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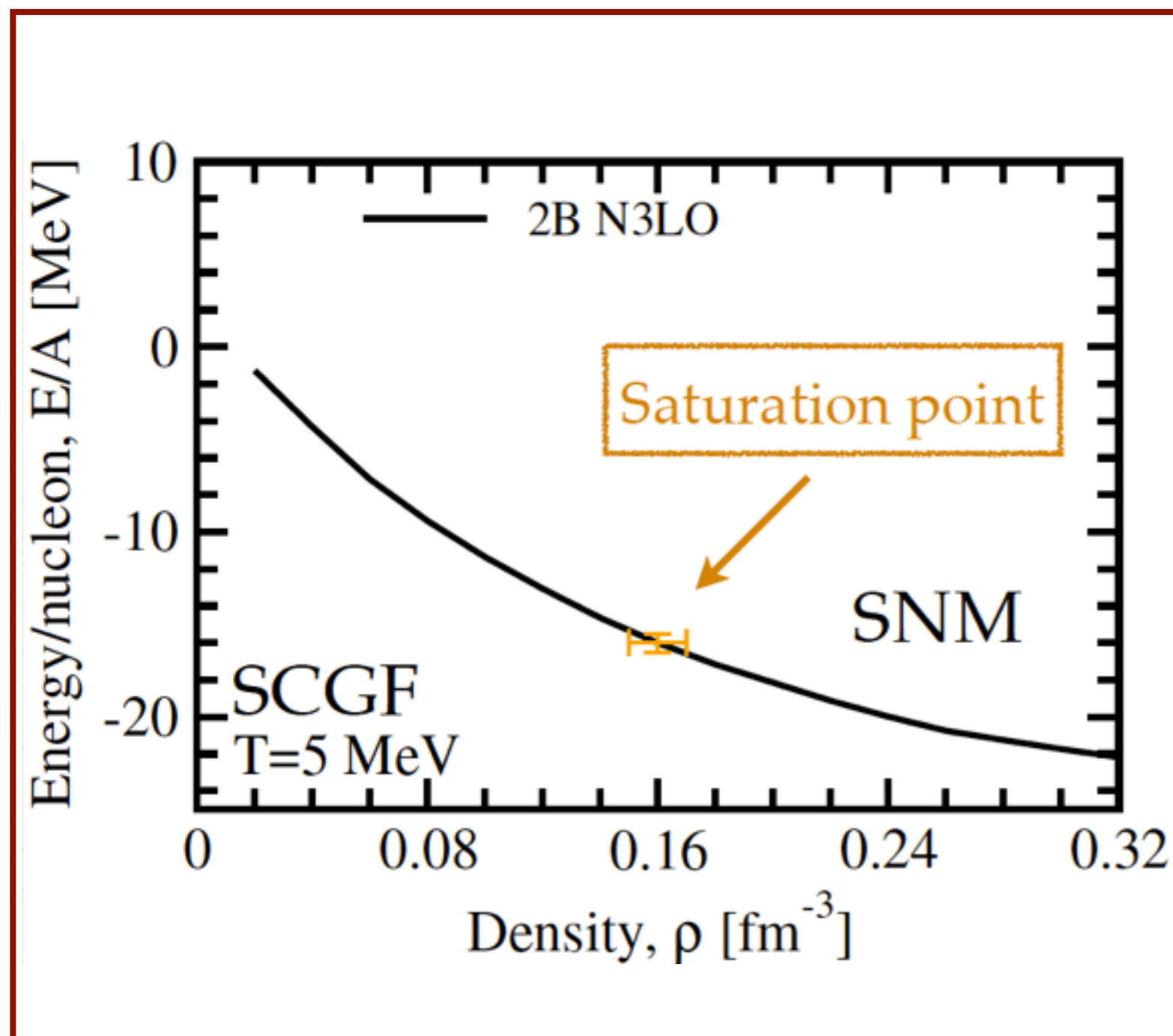
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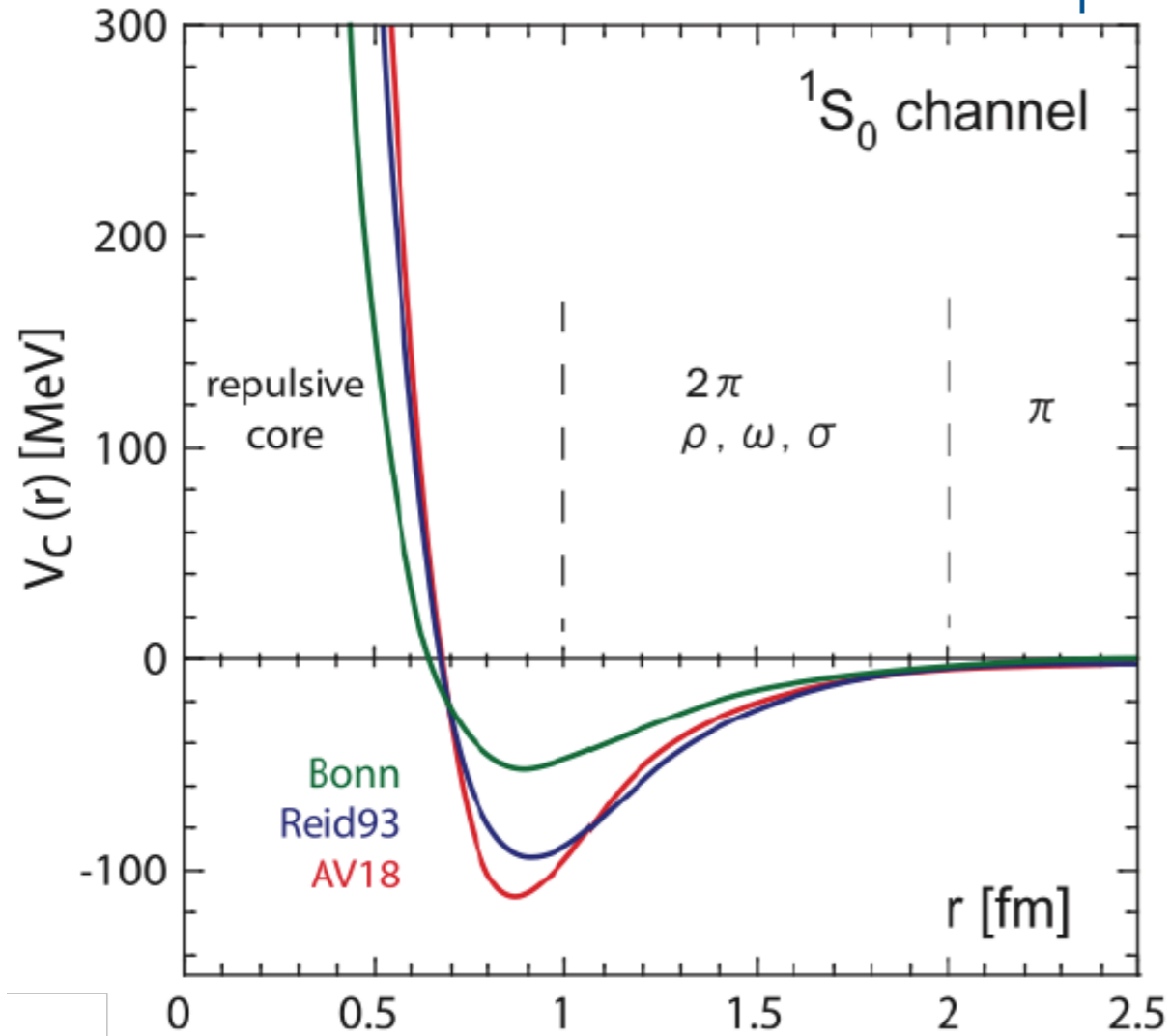
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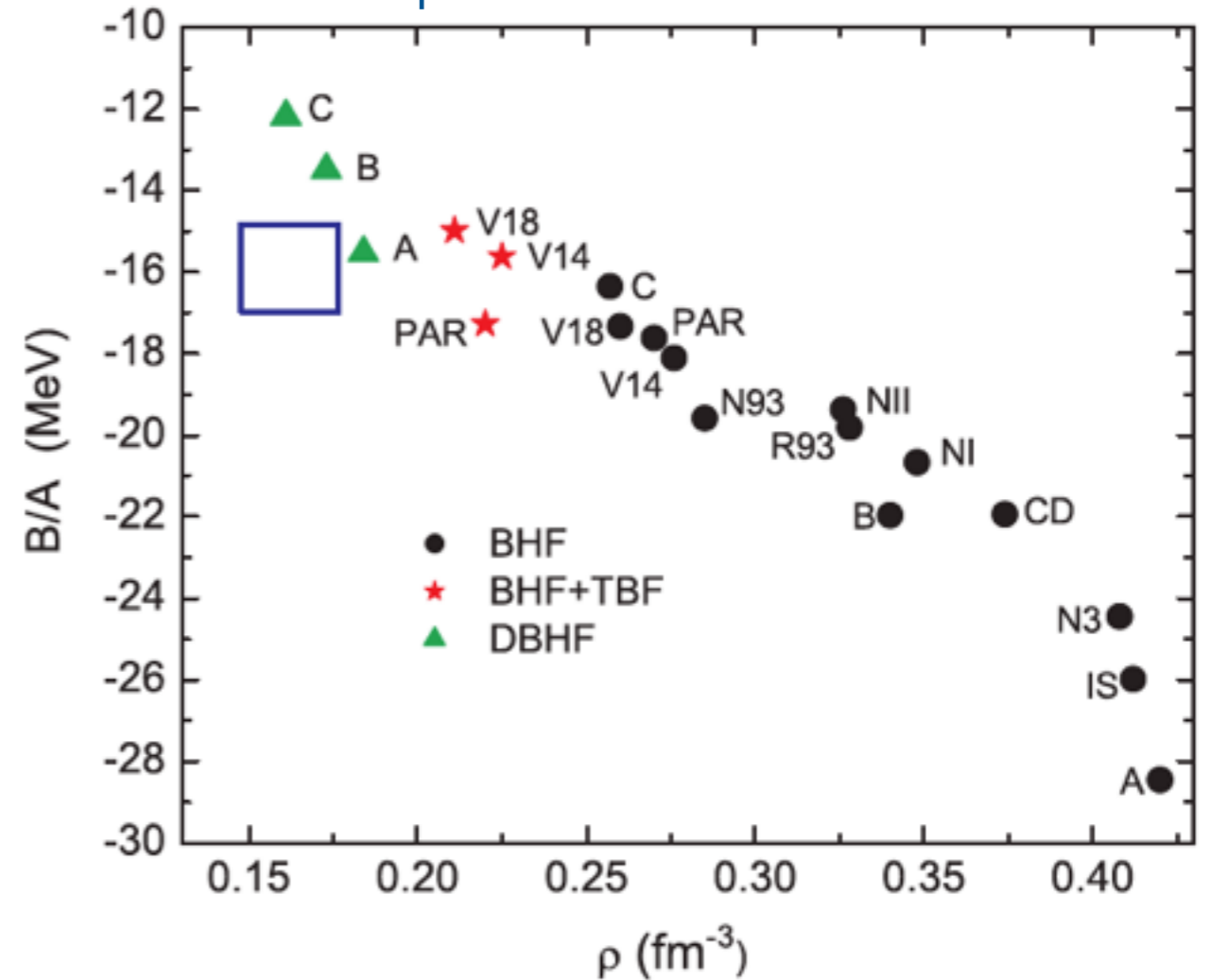


NN interaction is not unique



S. Aoki, et al. *Comput. Sci. Dis.* **1** 015009 (2008)

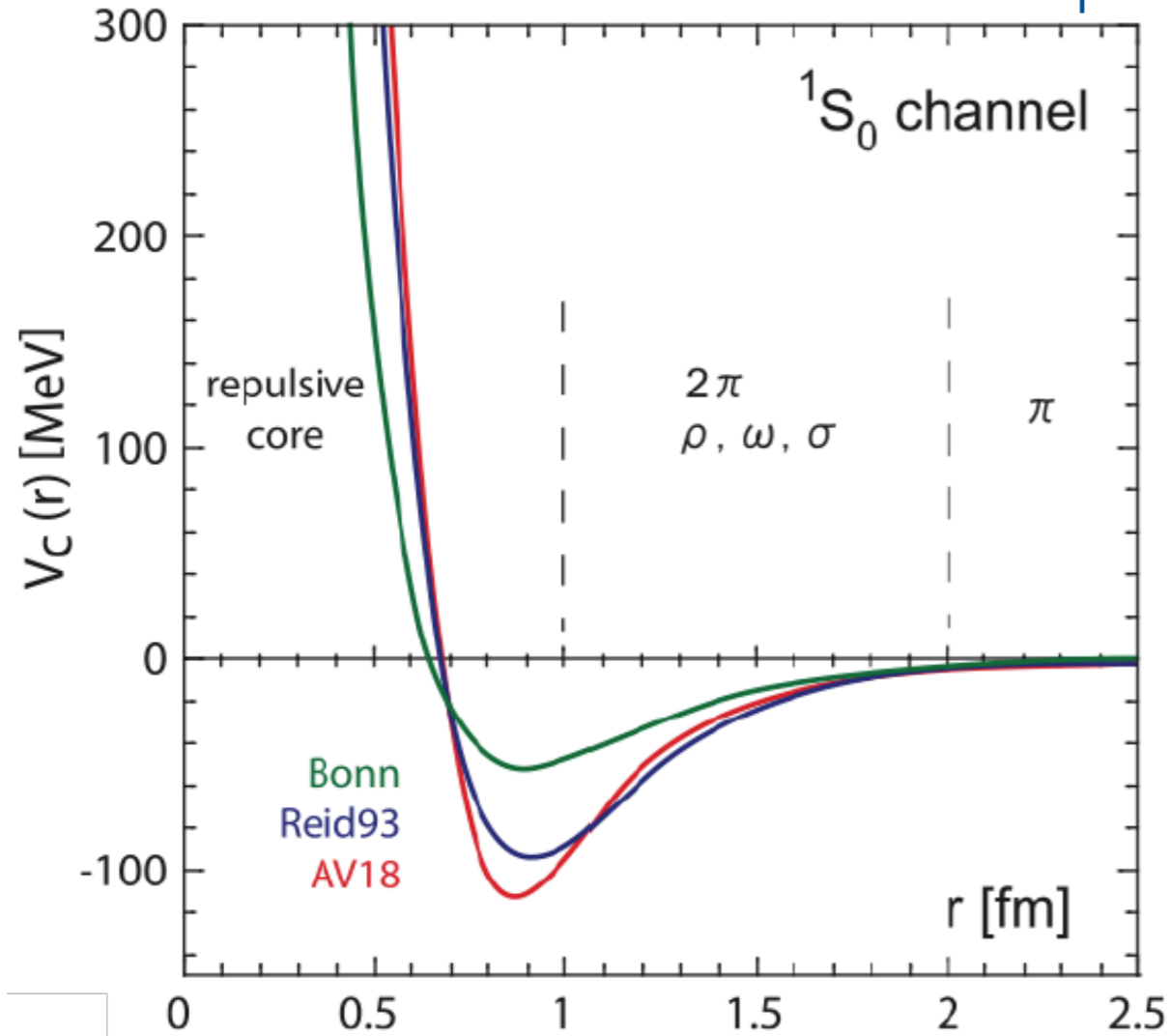
Saturation point of nuclear matter



Li, Lombardo, Schulze et al. *PRC* **74** 047304 (2006)

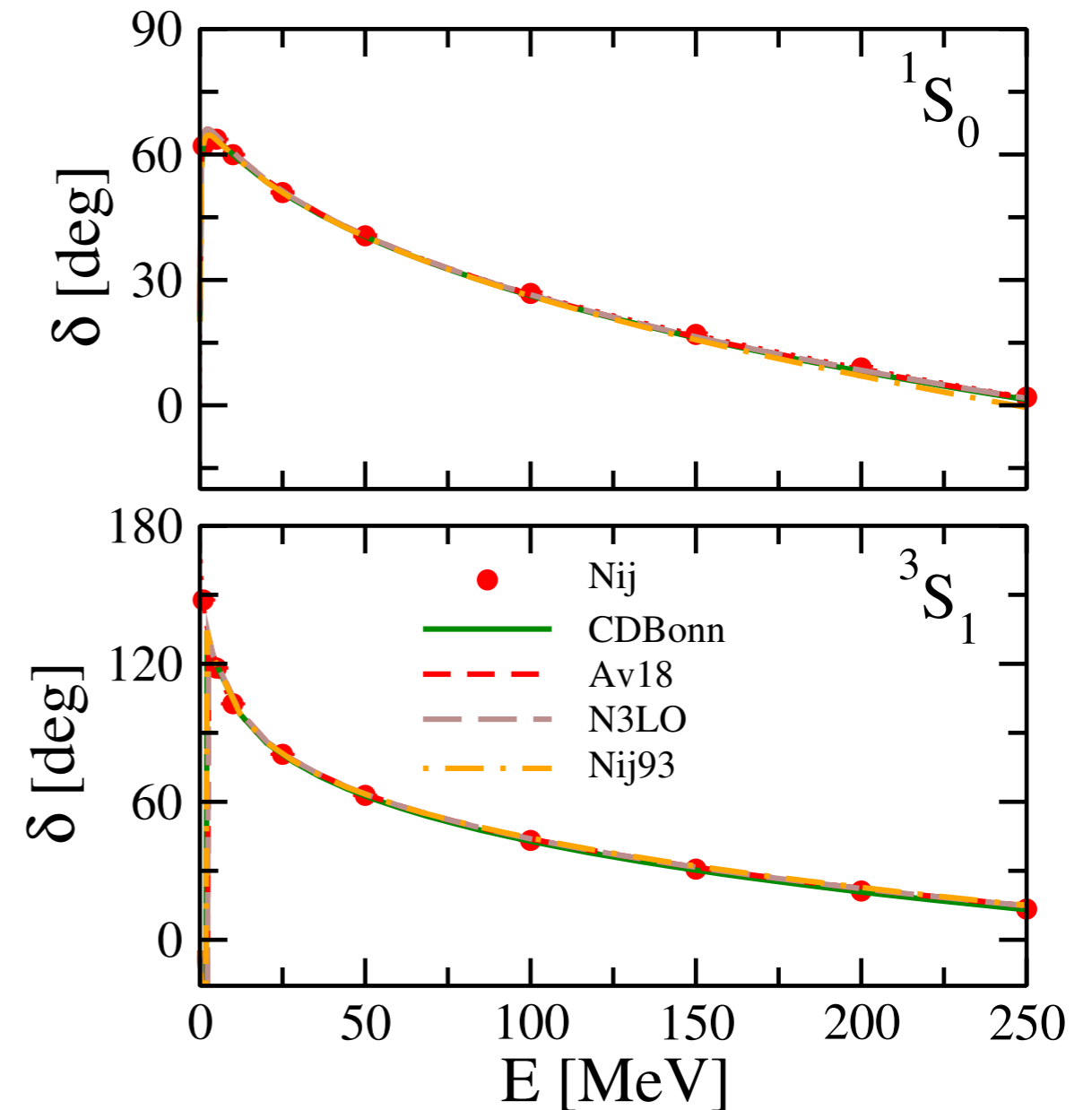
- Three-body forces needed for saturation ✗

NN interaction is not unique



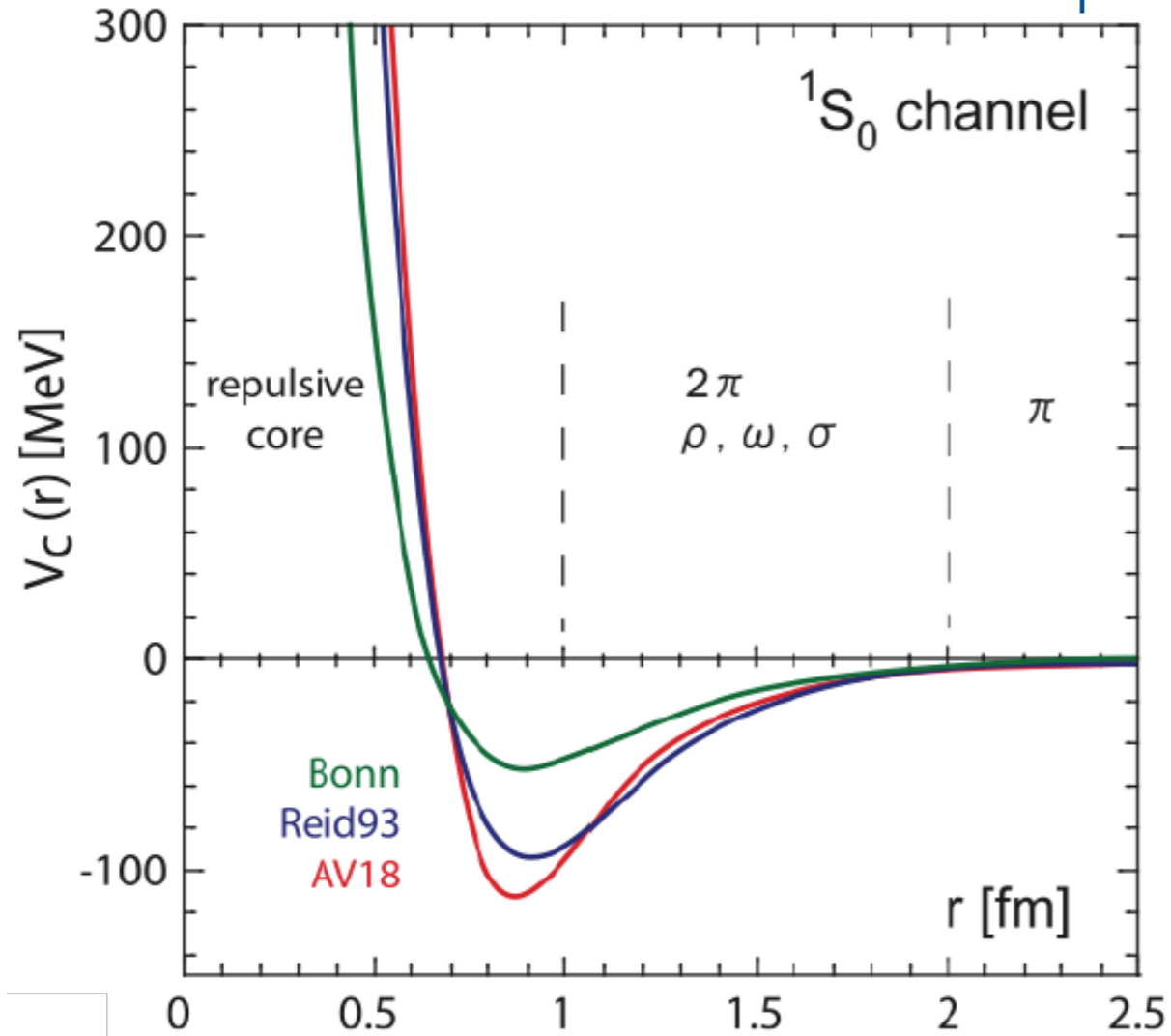
S. Aoki, et al. *Comput. Sci. Dis.* **1** 015009 (2008)

...but phase-shift equivalent!



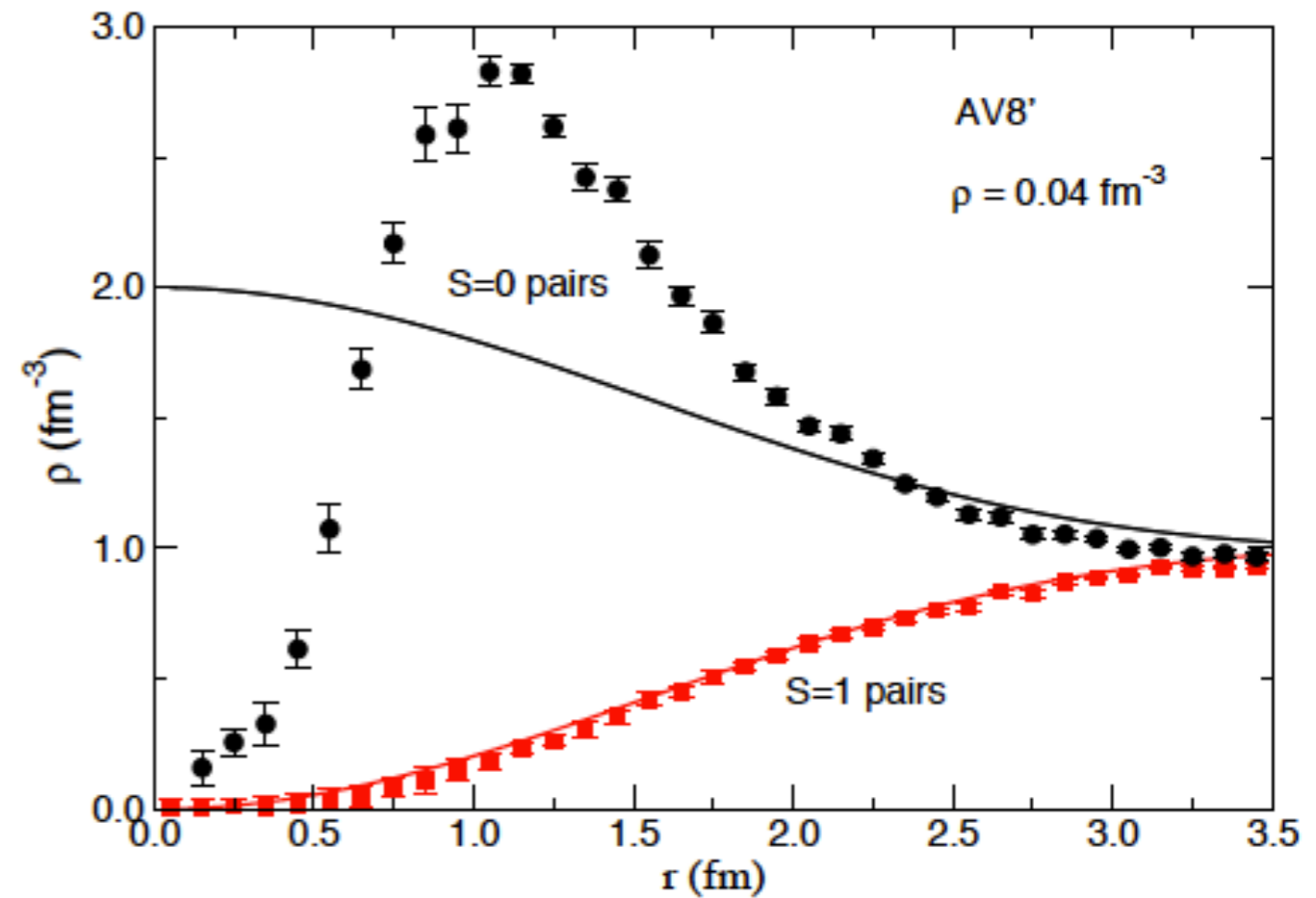
- Three-body forces needed for saturation ✗
- Non-uniqueness of nucleon forces ✗

NN interaction is not unique



S. Aoki, et al. *Comput. Sci. Dis.* **1** 015009 (2008)

Strong short-range correlations



Carlson et al., *Phys. Rev. C* **68** 025802 (2003)

- Three-body forces needed for saturation ✗
- Non-uniqueness of nucleon forces ✗
- Short-range core needs many-body treatment ✗

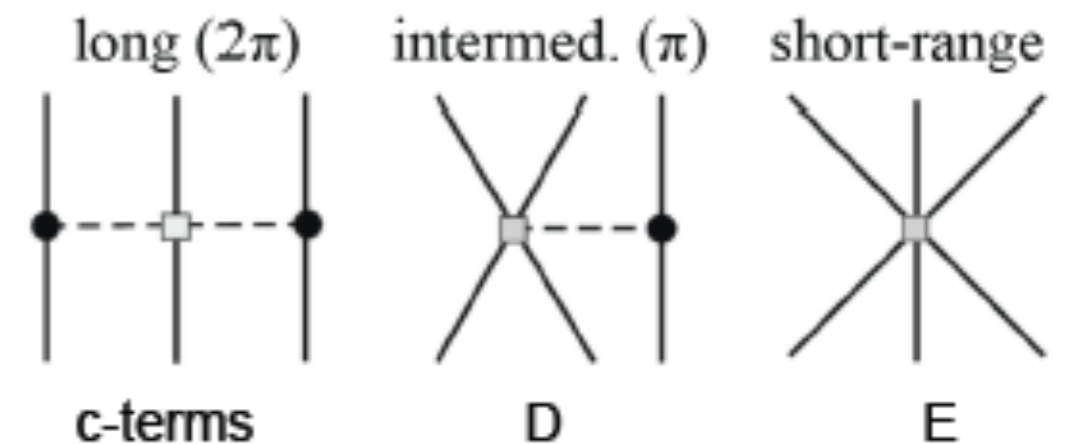
	NN	3N	4N
LO		—	—
NLO		—	—
N ² LO			—
N ³ LO			

$$\mathcal{O}\left(\frac{Q}{\Lambda}\right)$$

$$\Lambda \sim 1 \text{ GeV}$$

Chiral perturbation theory

- π and N as dof
- Systematic expansion
- 2N at N³LO - LECs from π N, NN
- 3N at N²LO - 2 more LECs
- (Often further renormalized)



Weinberg, *Phys. Lett. B* **251** 288 (1990), *Nucl. Phys. B* **363** 3 (1991)

Entem & Machleidt, *Phys. Rev. C* **68**, 041001(R) (2003)

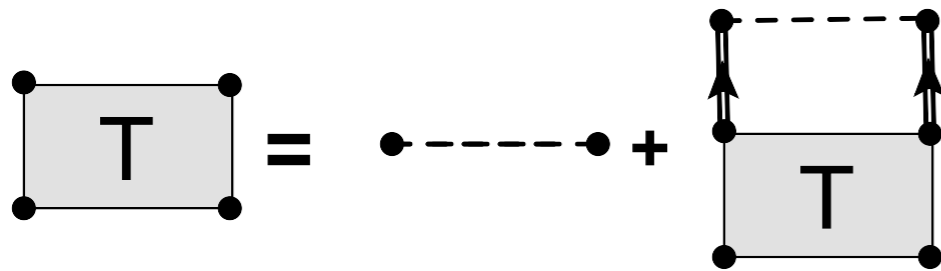
Tews, Schwenk et al., *Phys. Rev. Lett.* **110**, 032504 (2013) 13

Ladder approximation with 3BF

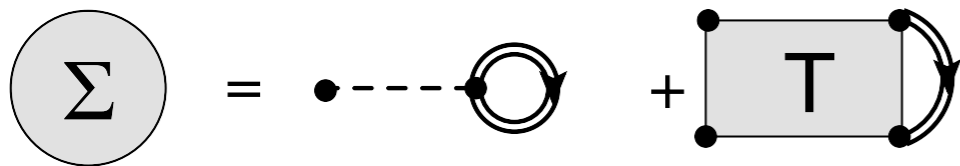
Two-body interaction



In-medium T-matrix

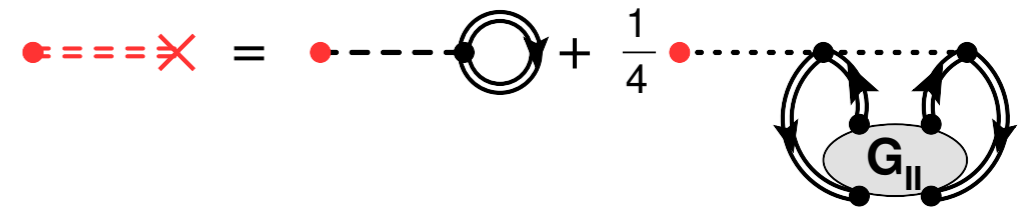


Self-energy

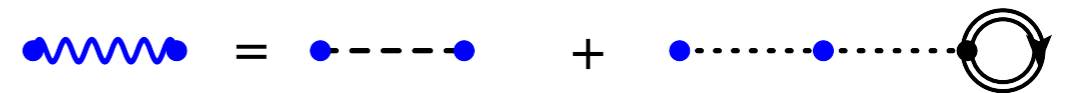


Effective interactions

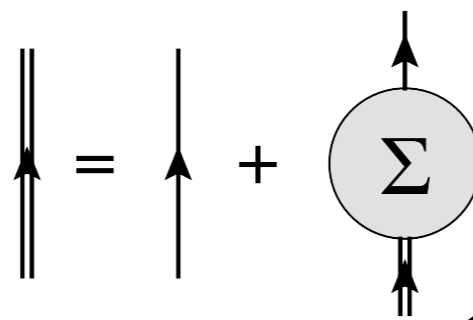
Effective one-body force



Effective two-body force



Dyson equation

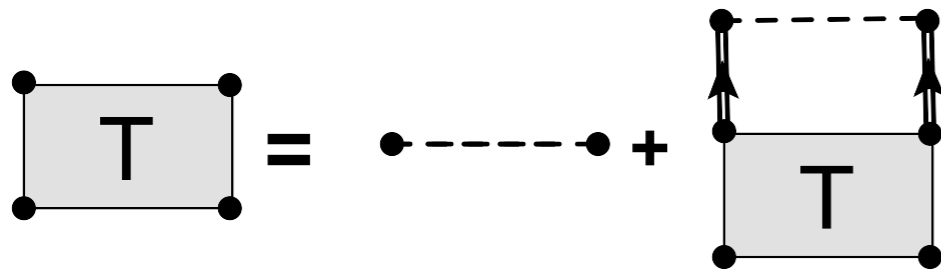


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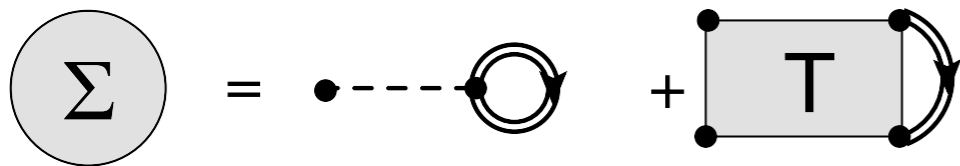
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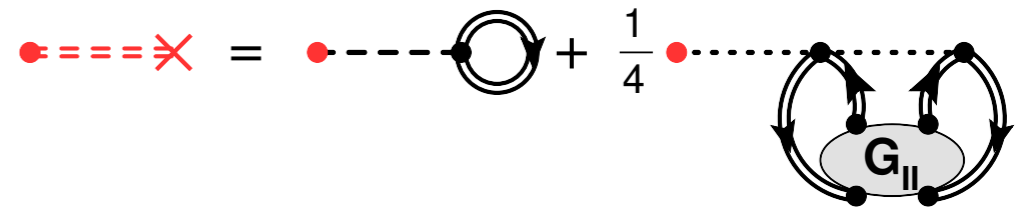


Self-energy

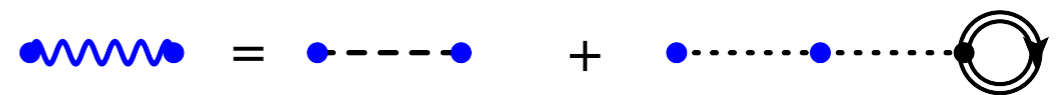


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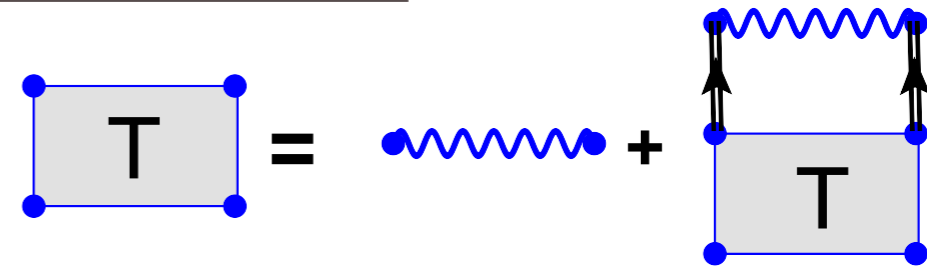
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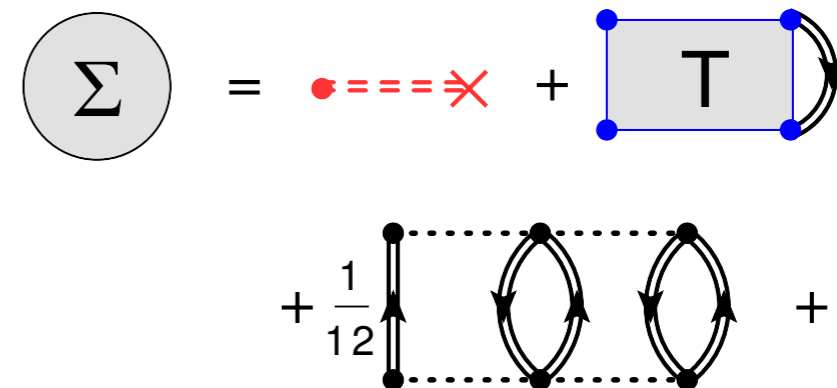
Effective two-body force



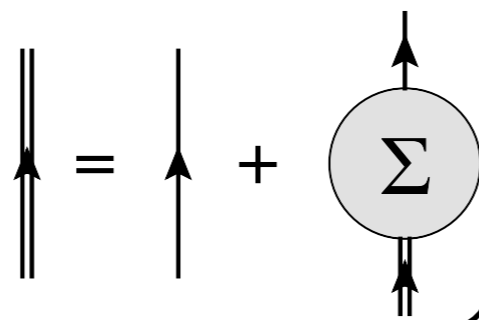
In-medium T-matrix



Self-energy



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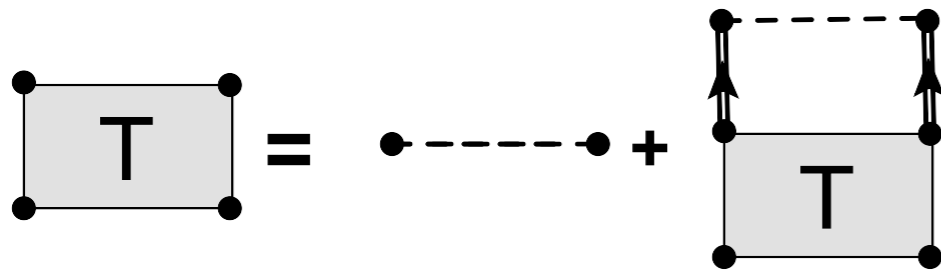


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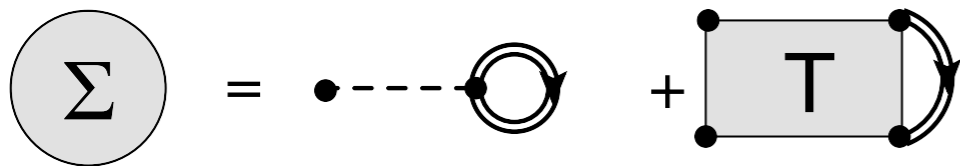
Two-body interaction



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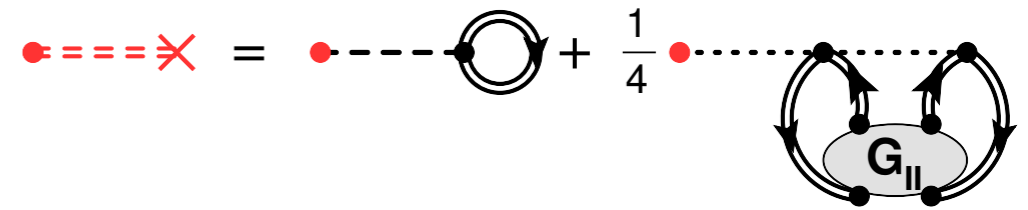


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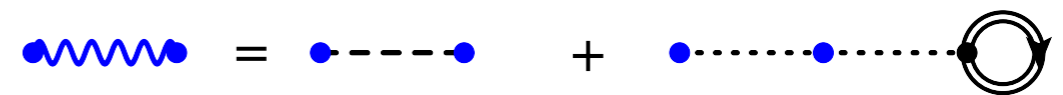


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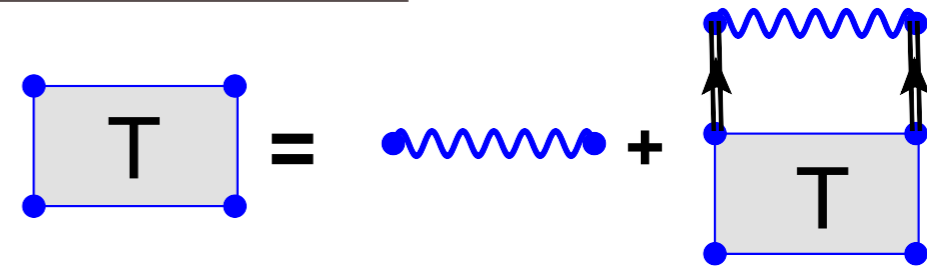
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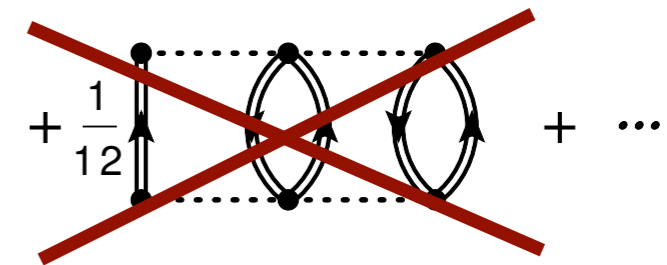
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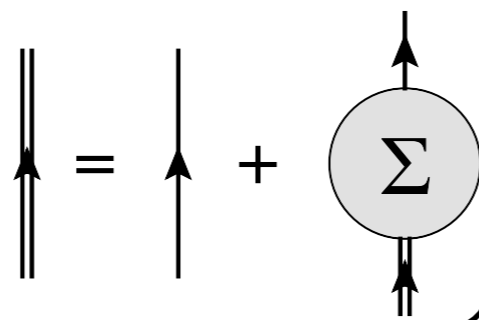
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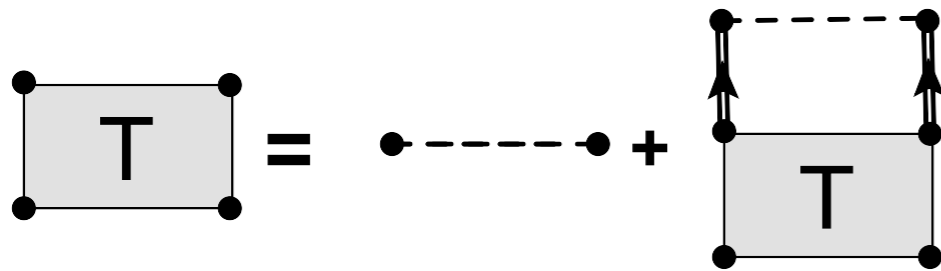


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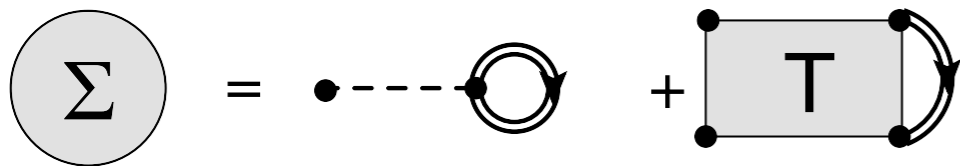
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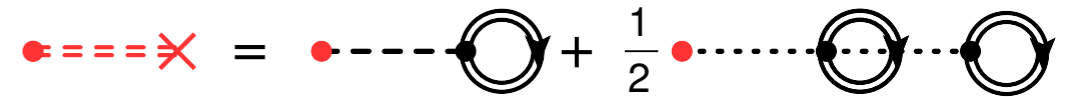


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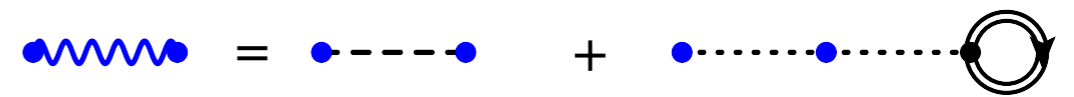


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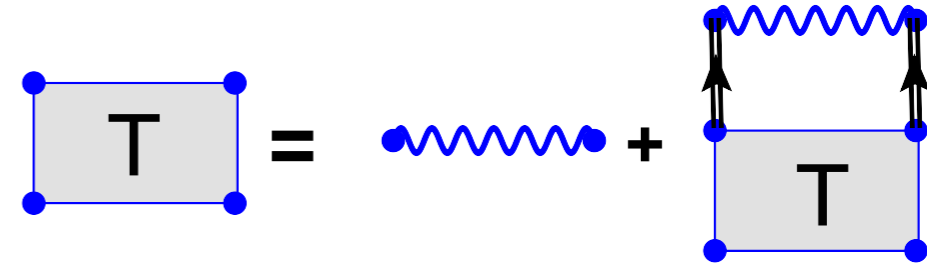
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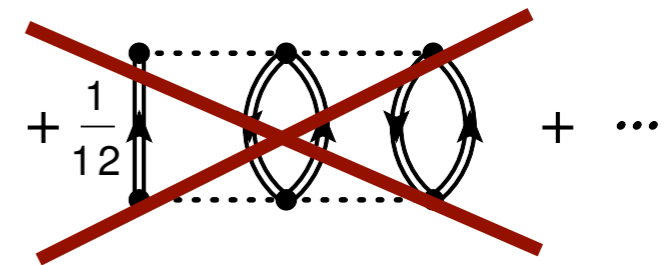
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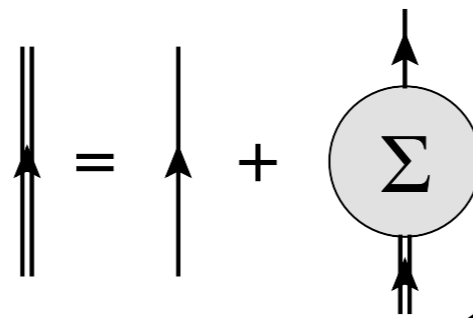
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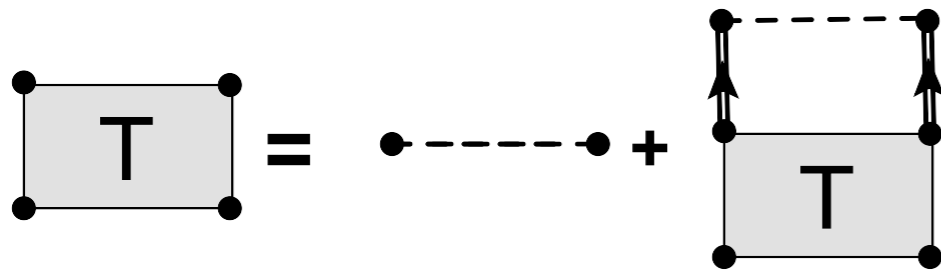


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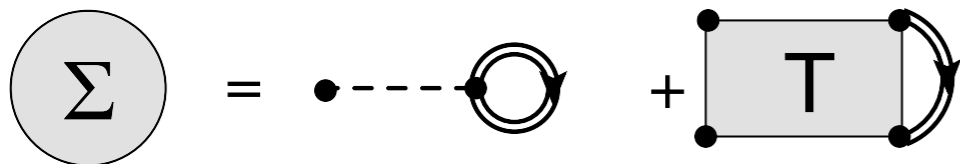
Two-body interaction



In-medium T-matrix

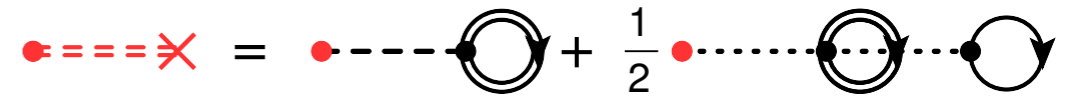


Self-energy

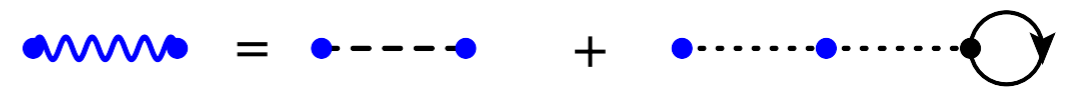


Effective interactions

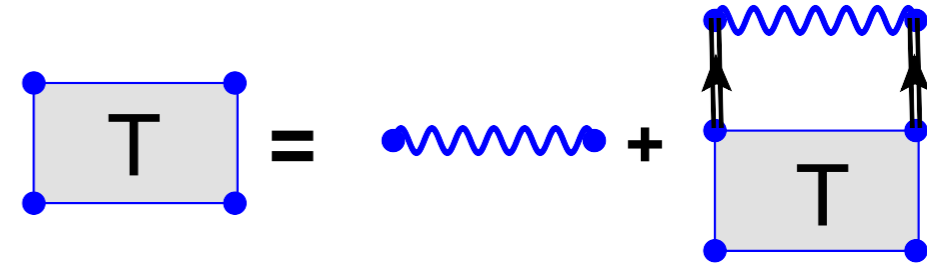
Effective one-body force



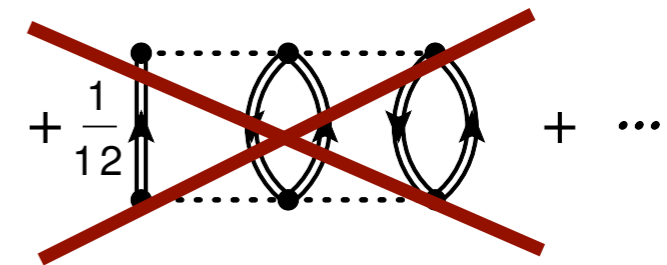
Effective two-body force



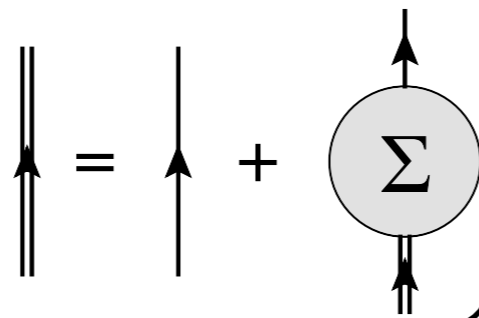
In-medium T-matrix



Self-energy



Dyson equation

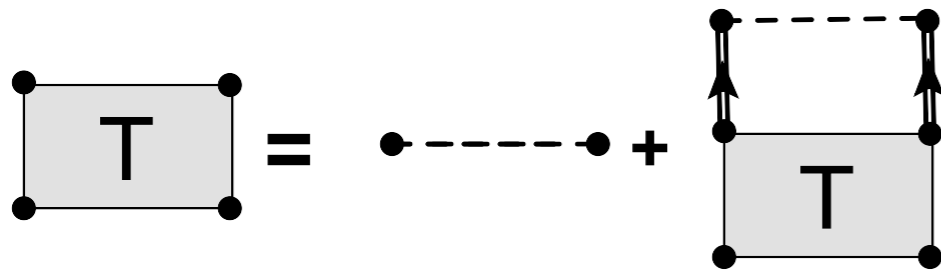


Ladder approximation with 3BF

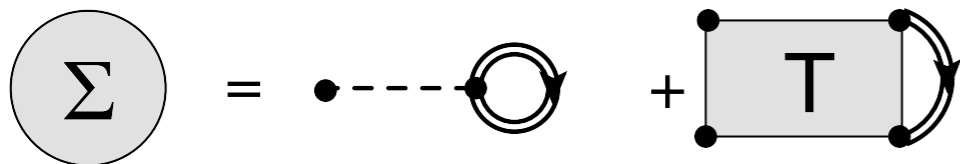
Two-body interaction



In-medium T-matrix

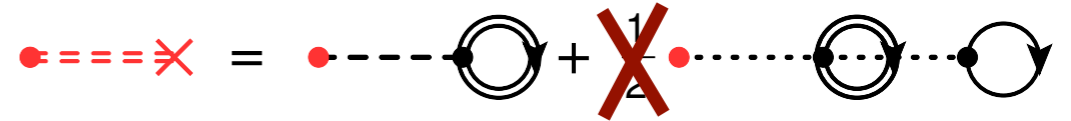


Self-energy

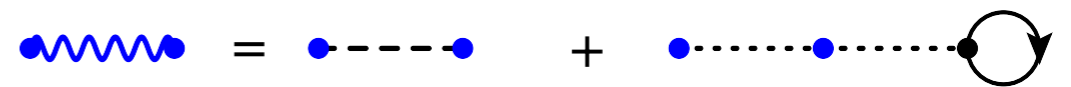


Effective interactions

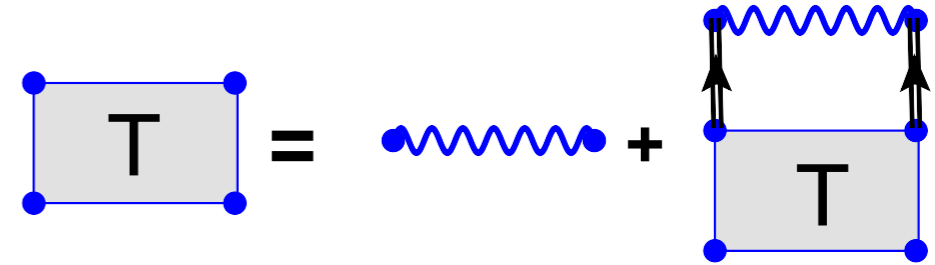
Effective one-body force



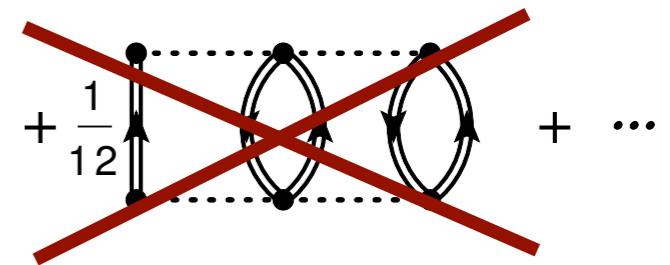
Effective two-body force



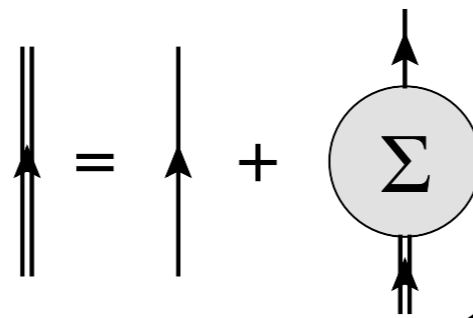
In-medium T-matrix



Self-energy



Dyson equation



Density-dependent interaction

Chiral NN effective 2B forces: symmetric matter

Two-body N3LO



Uncorrelated average¹



Correlated average²

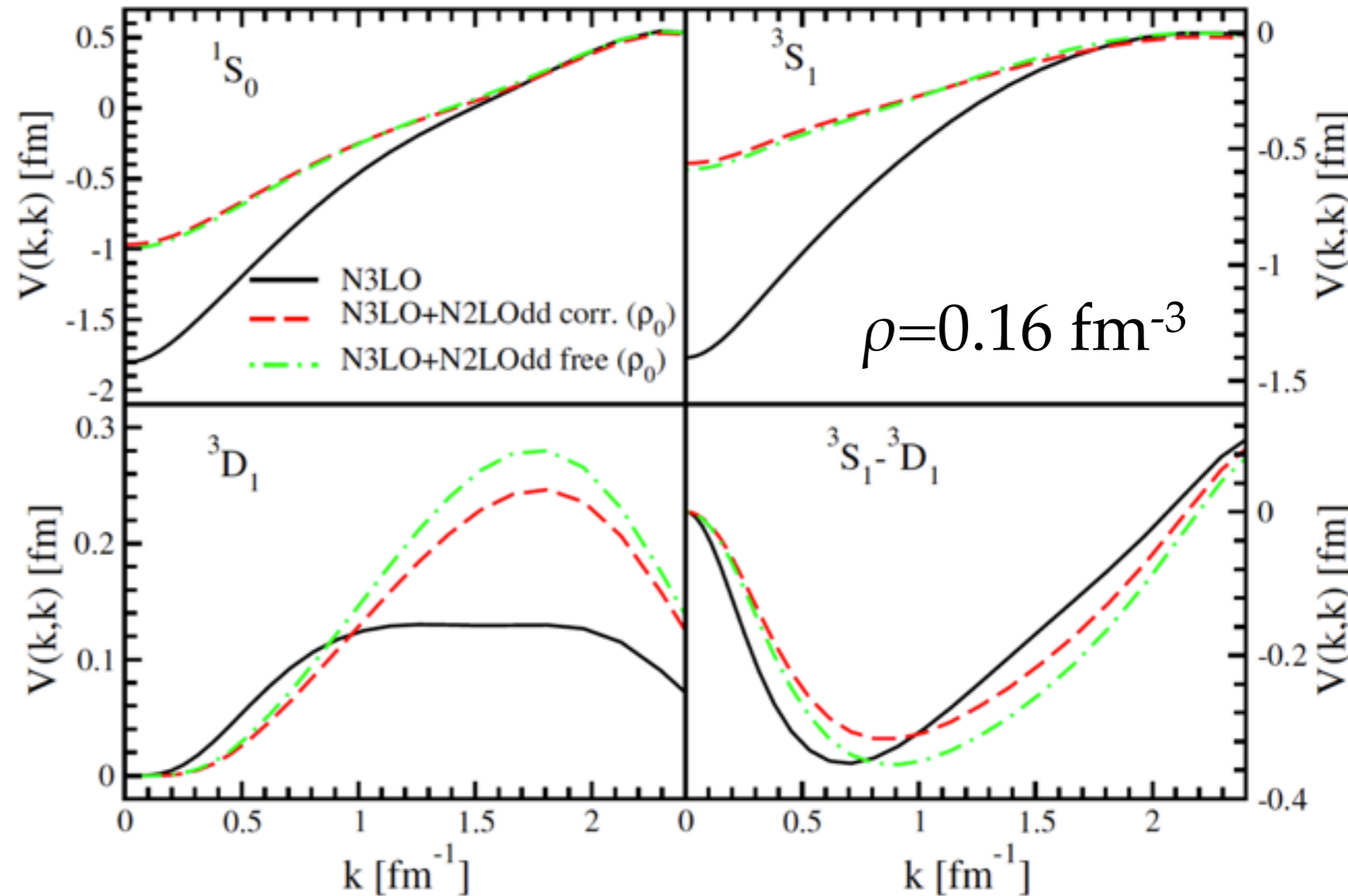


LECs

$$c_D = -1.11$$

$$c_E = -0.66$$

$$k \neq k' \Rightarrow \frac{1}{2}(k + k')$$



- 3NF bring **repulsion**: **correlated** & **uncorrelated** averages are similar
- **Correlated** average brings **small** corrections to 1/2 of terms
- Diagonal $k=k'$ matrix elements computed
- Off-diagonal extrapolated & **regulated**

¹Holt et al. Phys. Rev. C **81** 024002 (2010)

²Carbone, Polls & Rios, PRC **90**, 054322 (2014); A. Carbone, PhD thesis 15

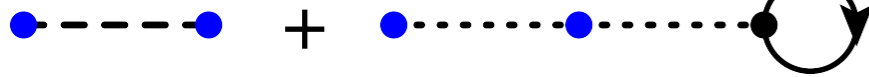
Symmetric matter

Theoretical uncertainties: average procedure

Two-body N³LO



Uncorrelated average



Correlated average



LECs

$$c_D = -1.11$$

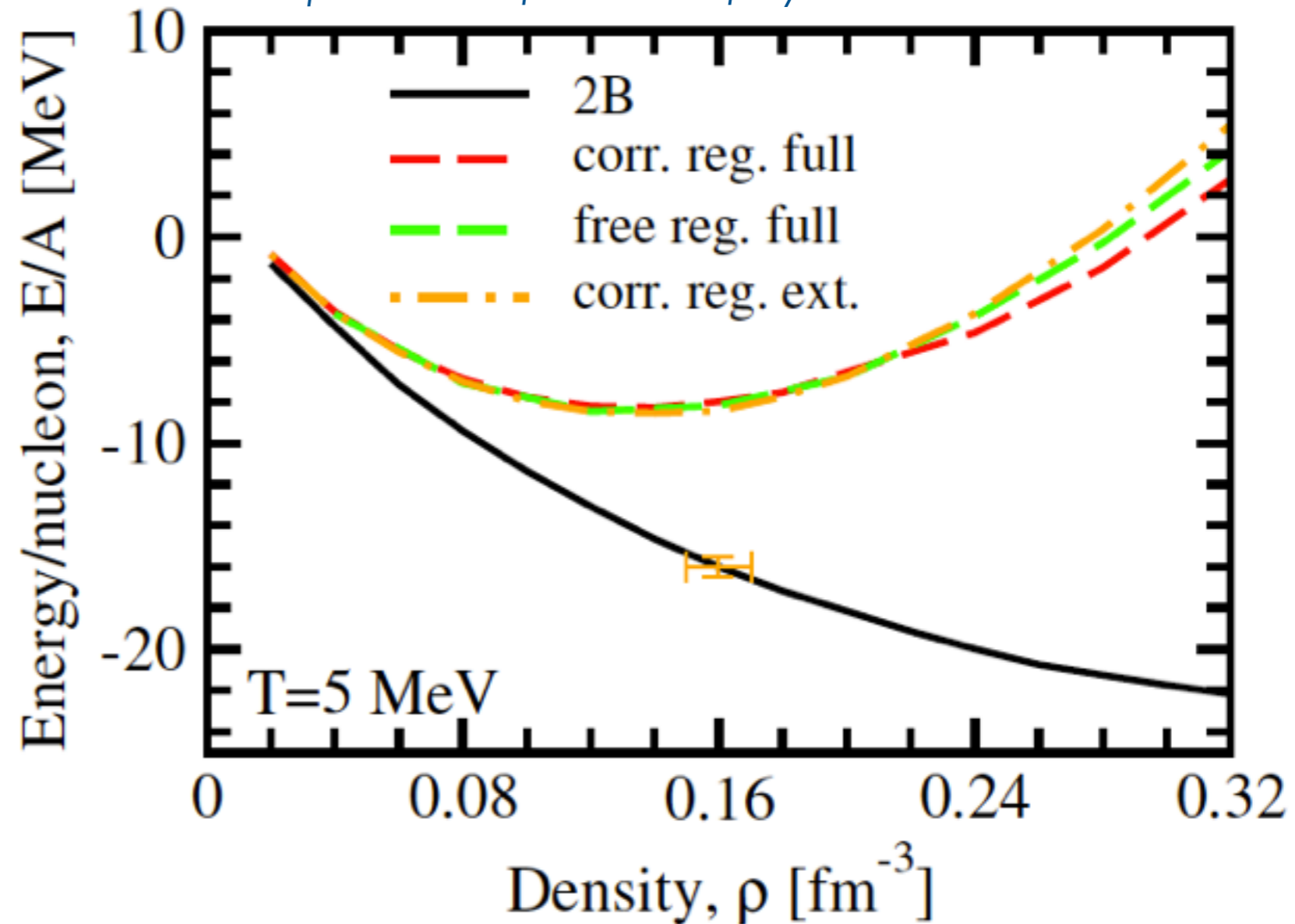
$$c_E = -0.66$$

$$K_0 \sim 60 \text{ MeV}$$

Energy from GMK sum-rule

... with 3N corrections¹ ...

Equation of state of symmetric matter



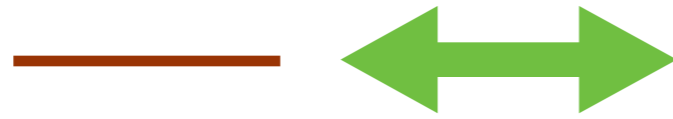
- Cor. reg. full=best we can do now is underbound
- Previous work with uncorrelated averages is validated
- Regulation at high momentum is irrelevant

¹Carbone, Polls & Rios, *PRC* **88** 044302 (2013)

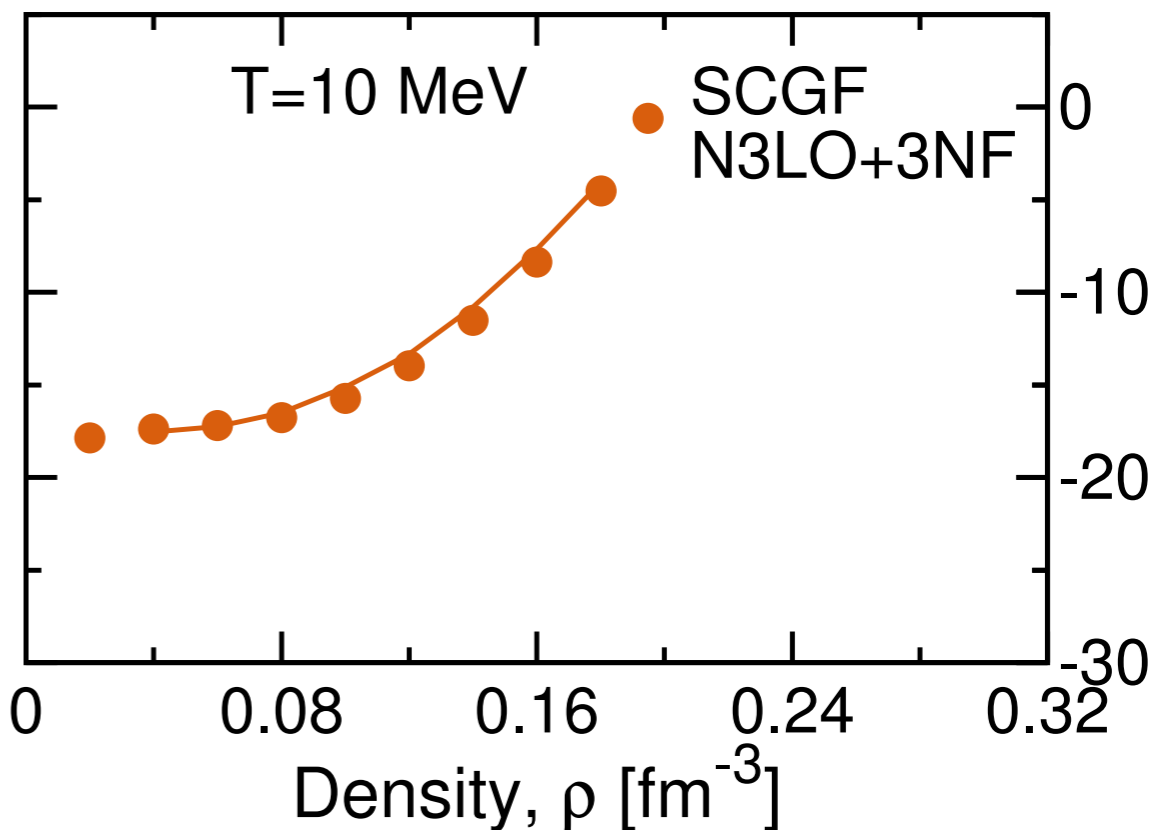
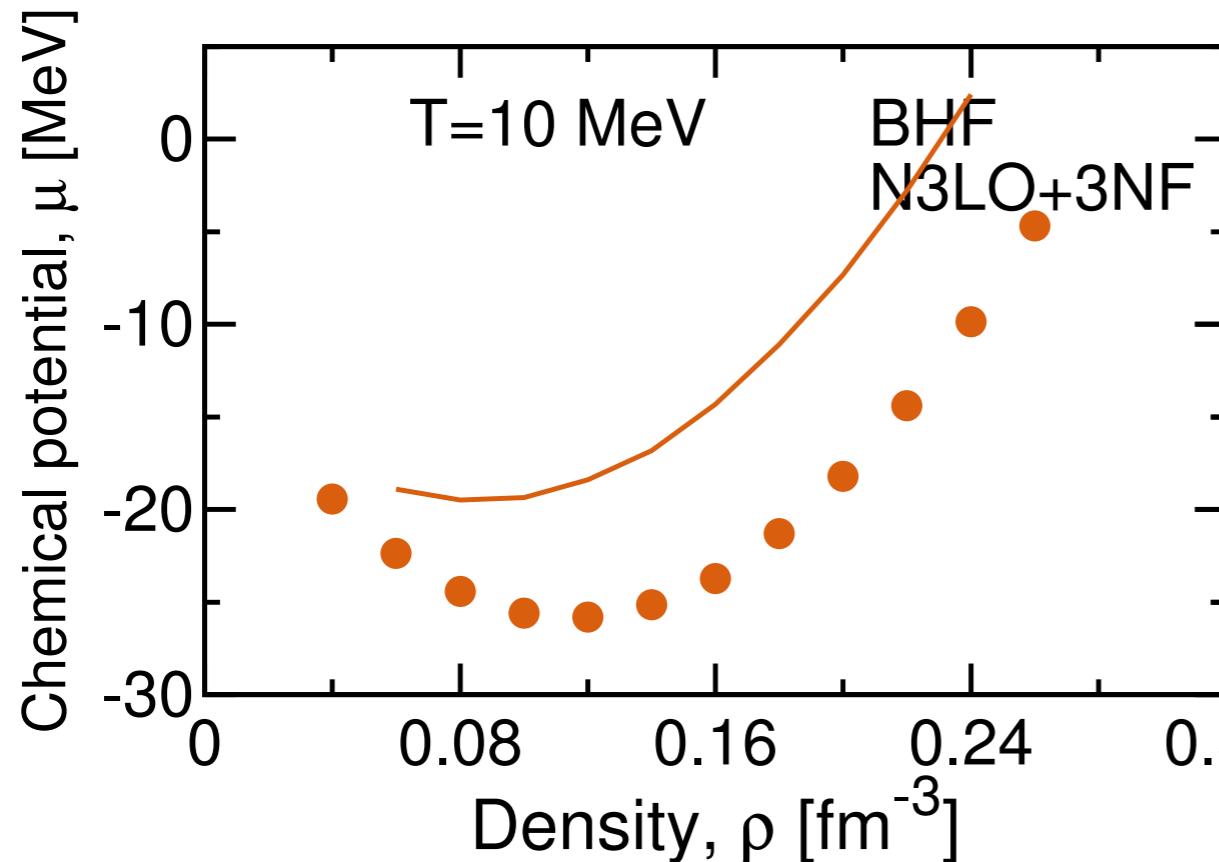
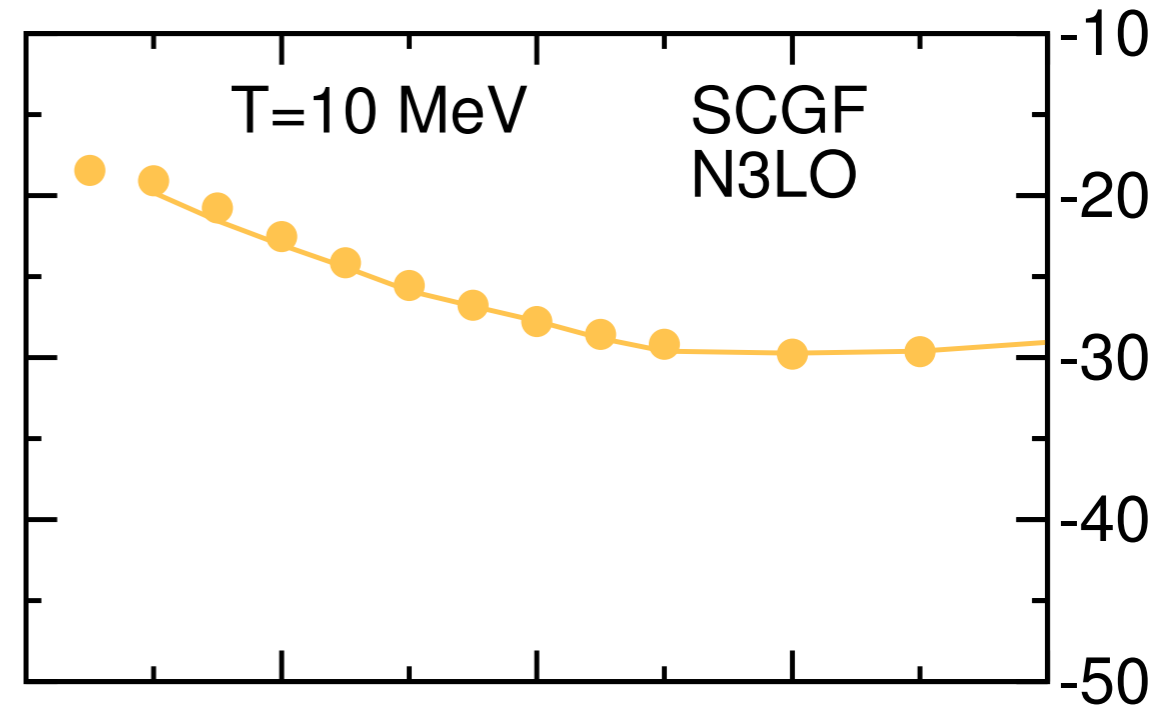
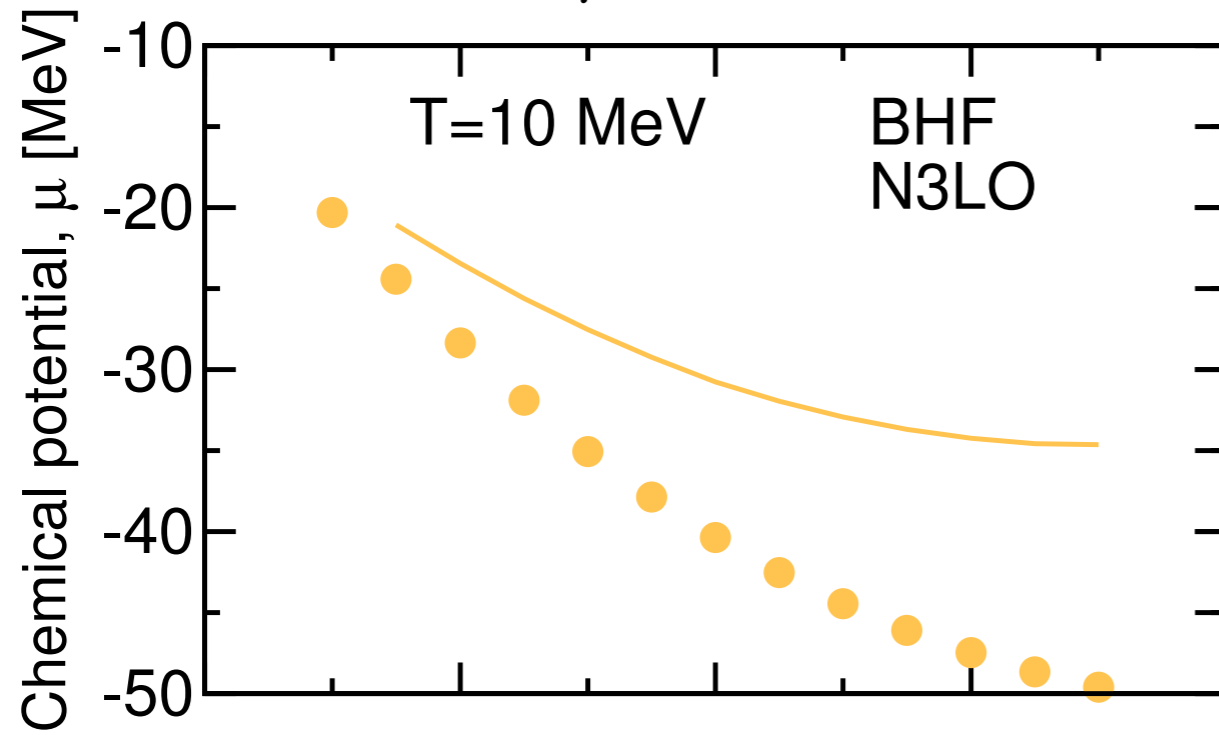
²Carbone, Polls & Rios, *PRC* **90** 054322 (2014); A. Carbone, PhD thesis 16

Thermodynamical consistency

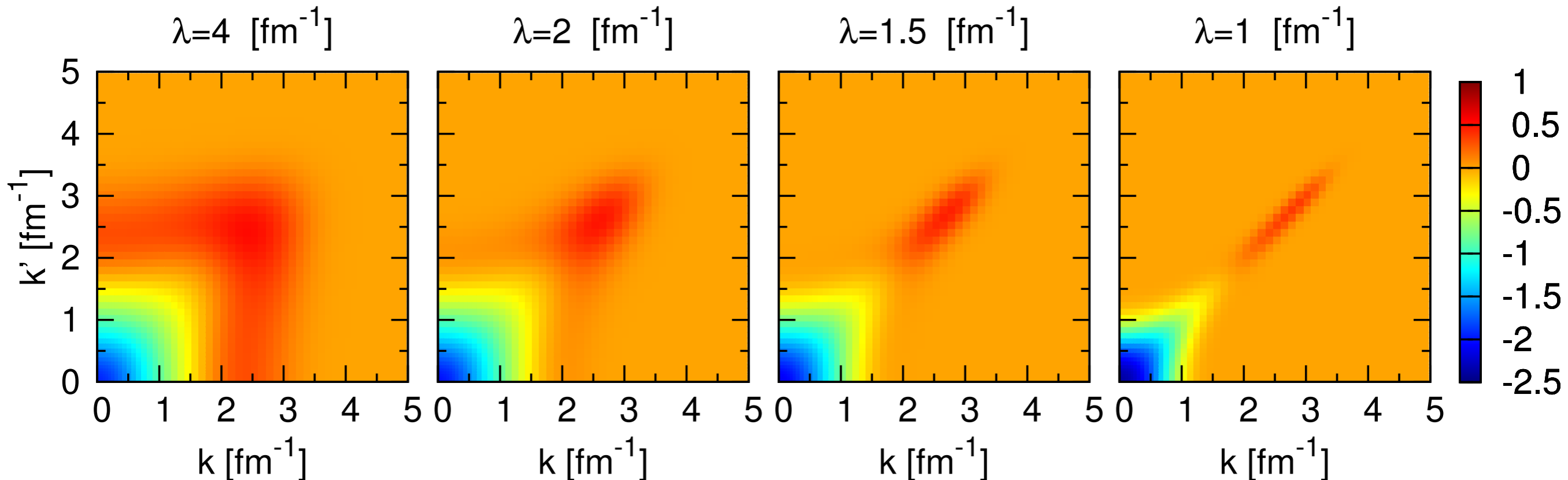
$$\mu = \frac{\partial E/\Omega}{\partial \rho}$$



$$\rho = \sum_k n(k; \tilde{\mu})$$



1S_0 NN matrix elements from N3LO

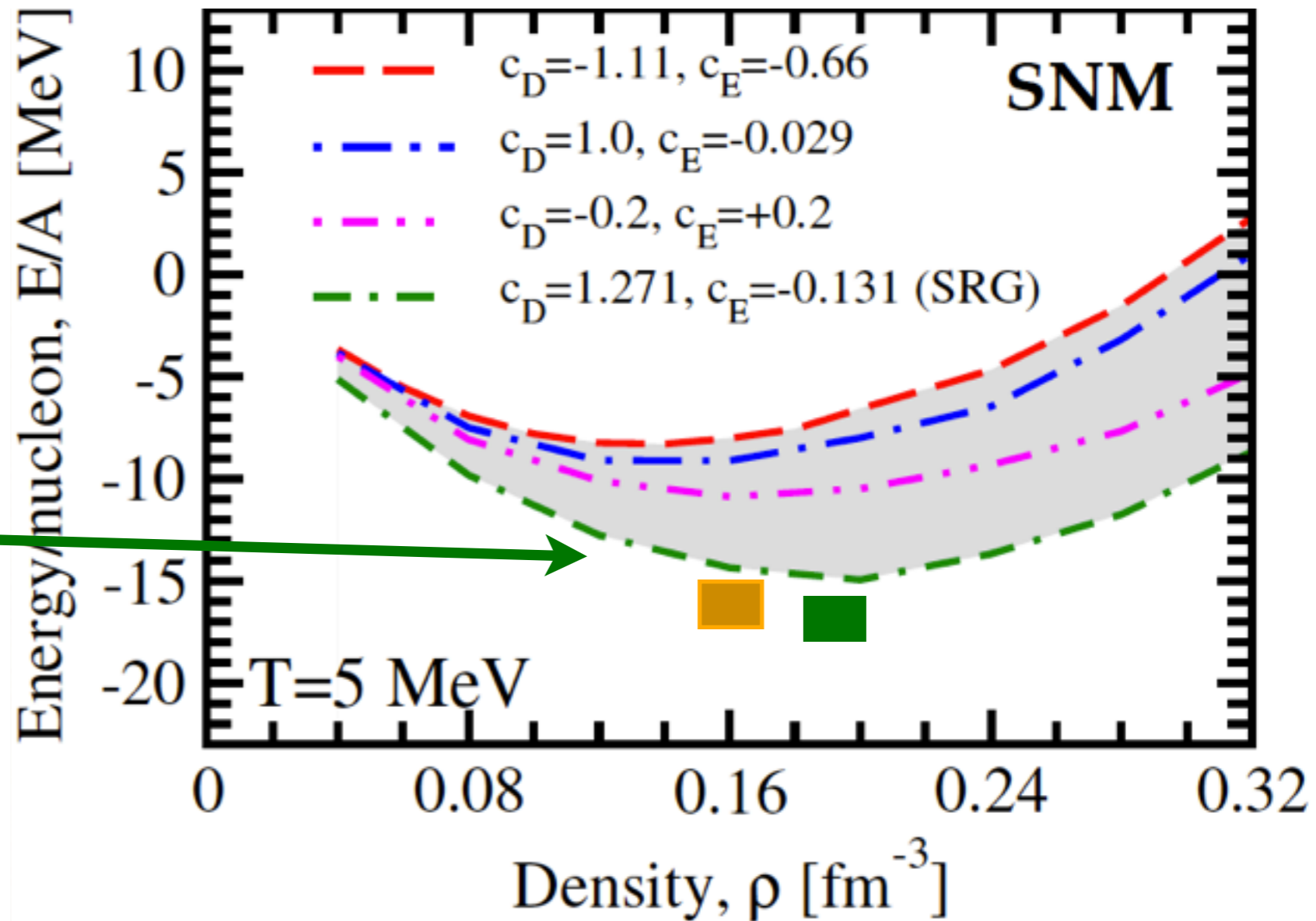


$$\frac{dH_s}{ds} = [[T_{rel}, H_s], H_s] \Leftrightarrow \lambda = s^{-1/4}$$

- Series of **unitary** transformation
- Observables **unaltered**, but force becomes **perturbative**
- **Induces** 3-, 4- and up to A-body forces...
- **If** these can be treated perturbatively, calculation is **easier**

Equation of state: LEC dependence

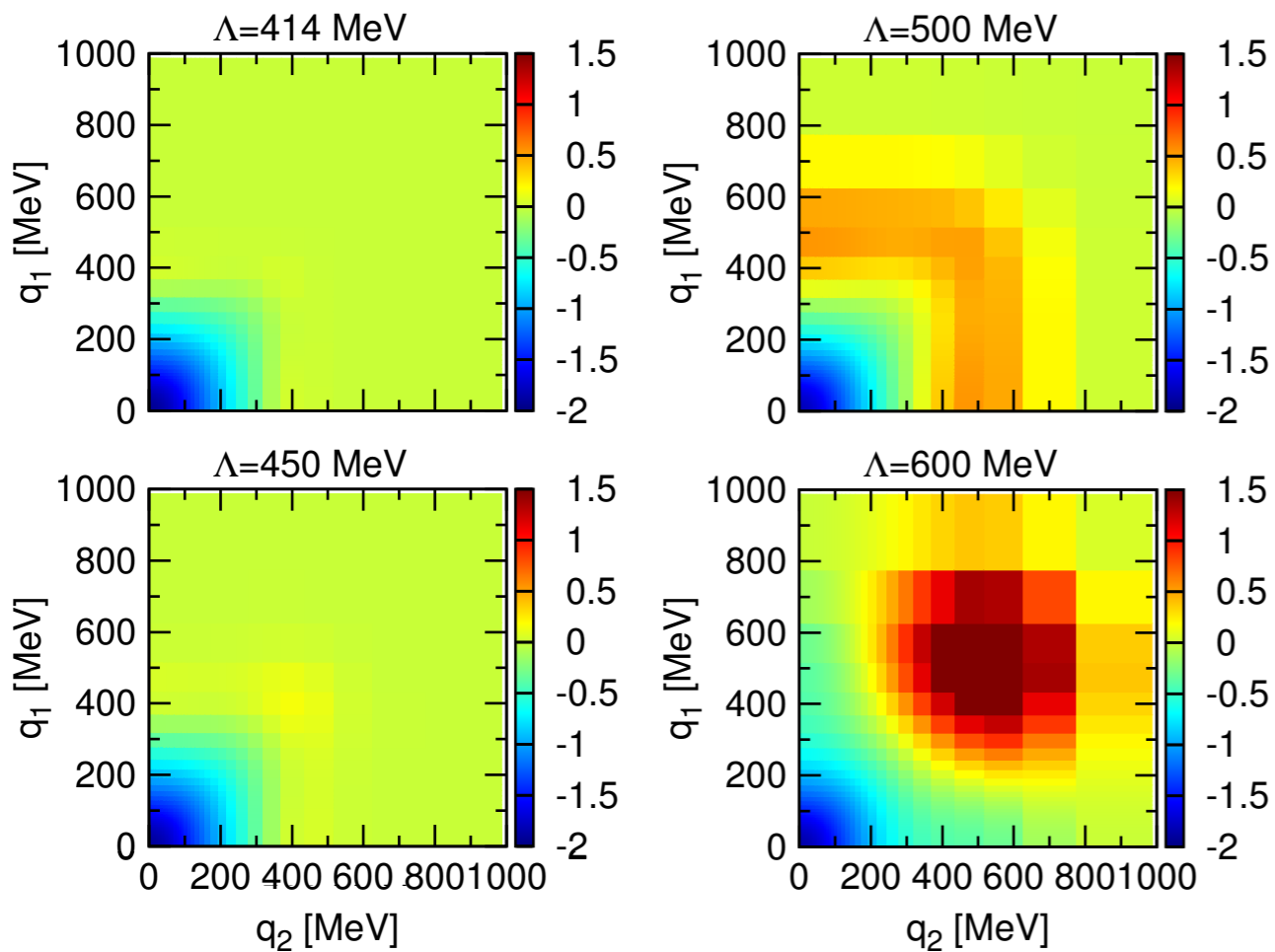
Uncorrelated average



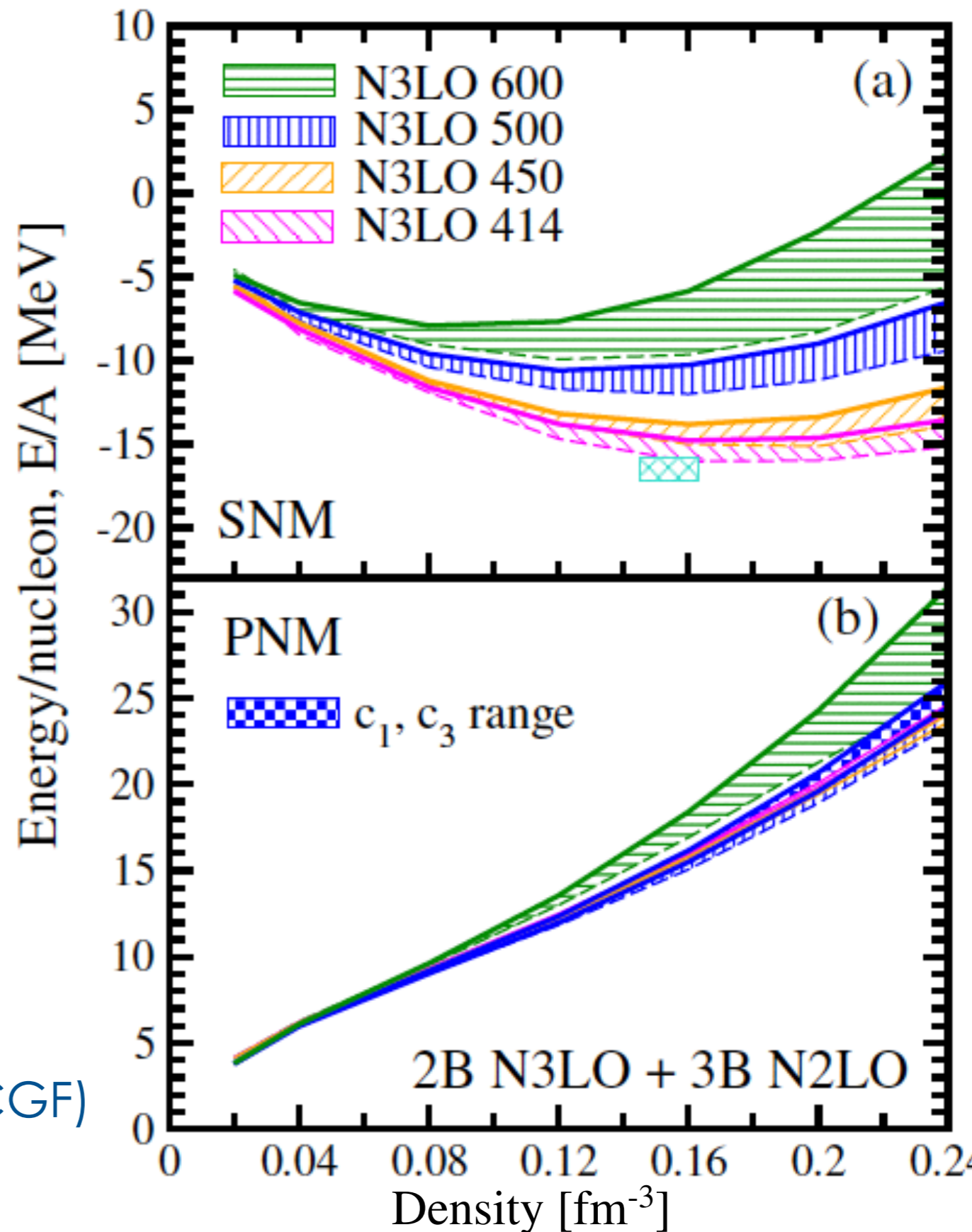
$\Lambda=2 \text{ fm}^{-1}$ SRG
well below c_D, c_E uncertainty

- LECs dependence is strong
- Renormalization via SRG: nuclear structure calculations?
- Small 3NF effects with larger saturation densities \Rightarrow smaller radii

Idaho chiral NN forces



- N3LO forces with **different** cut-offs
- (c_D, c_E) fitted to few-body, no SRG
- Explore **error** in many-body (BHF vs SCGF)
- Neutron matter is **more perturbative**



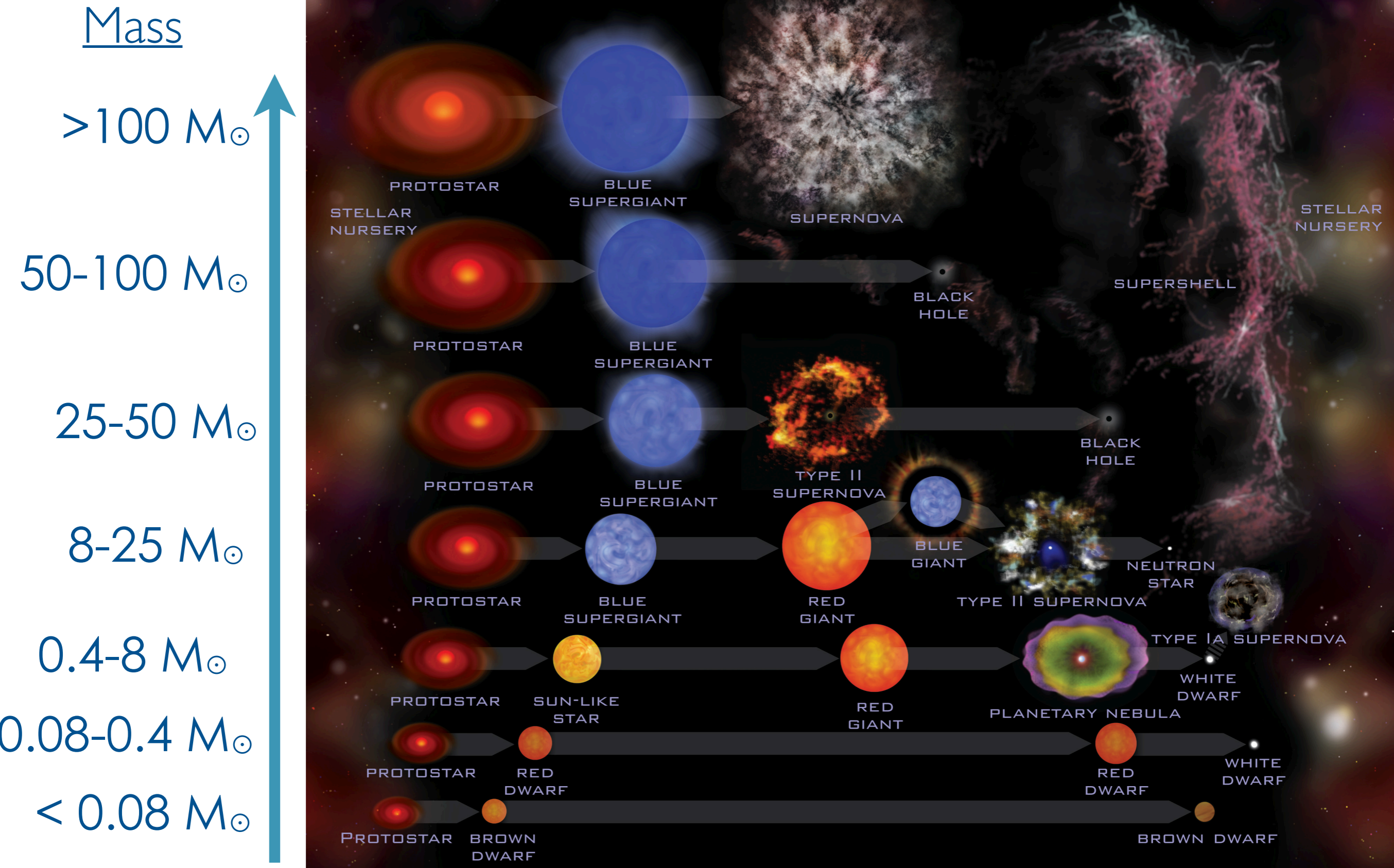
➡ *Motivation*

➡ *Nuclear matter: Equation of state with 3NFs*

➡ **Neutron matter: beyond-BCS pairing**

What is a neutron star?

Why care?



Neutron star 101

Mass

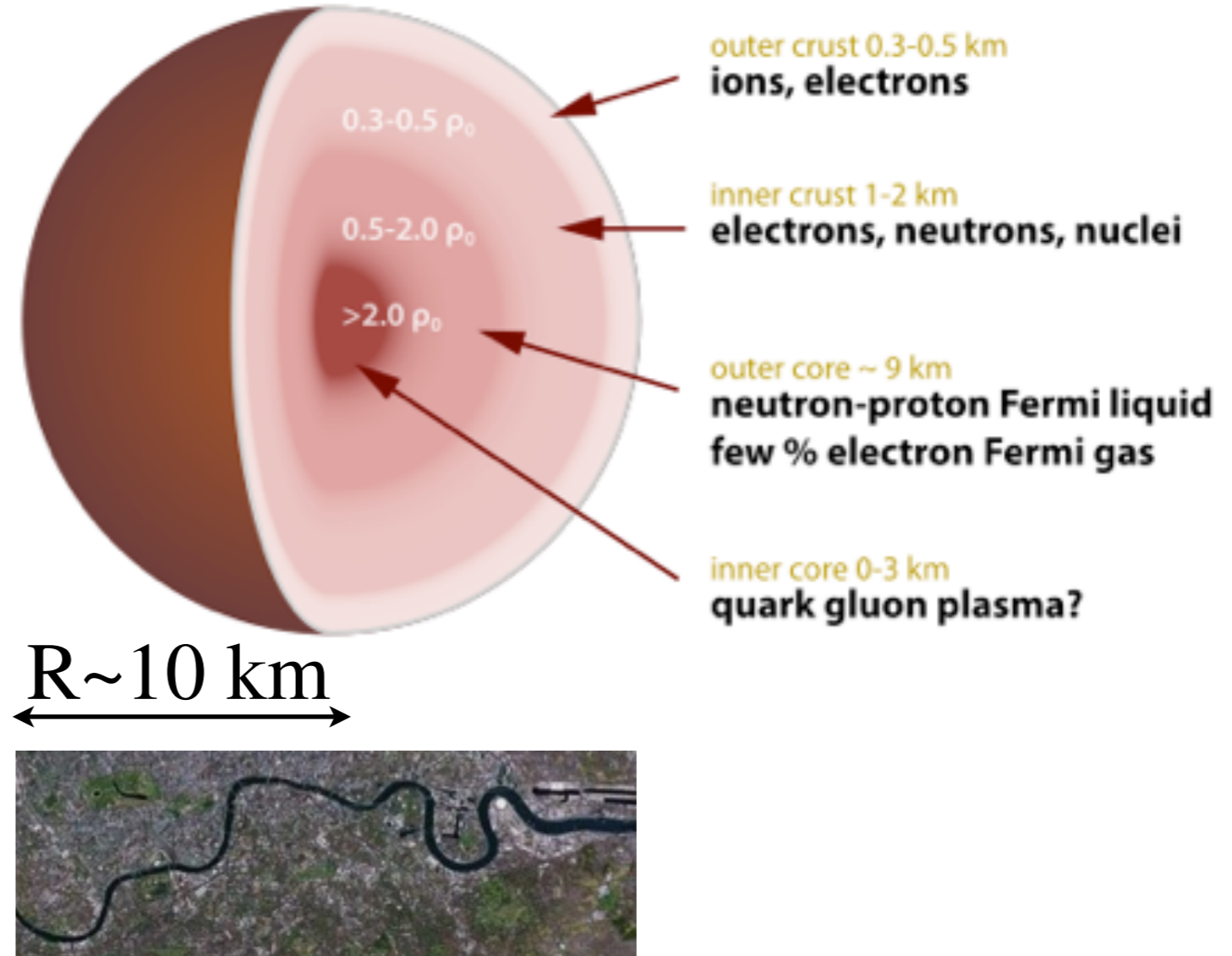
$$M \sim 1.5M_{\odot} = 3 \times 10^{30} \text{ kg}$$

Radius

$$R \approx 10 \text{ km}$$

Mass density

$$\rho = \frac{M}{V}$$



Neutron star 101

Mass

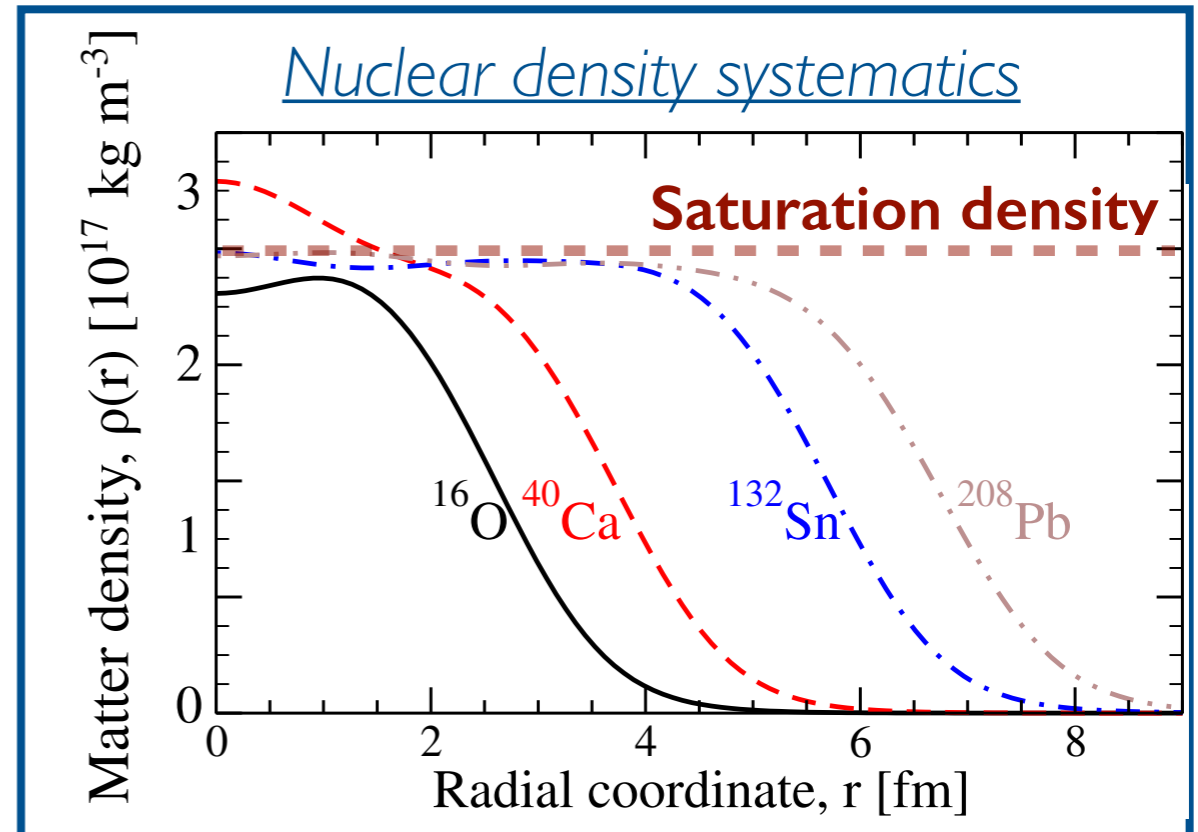
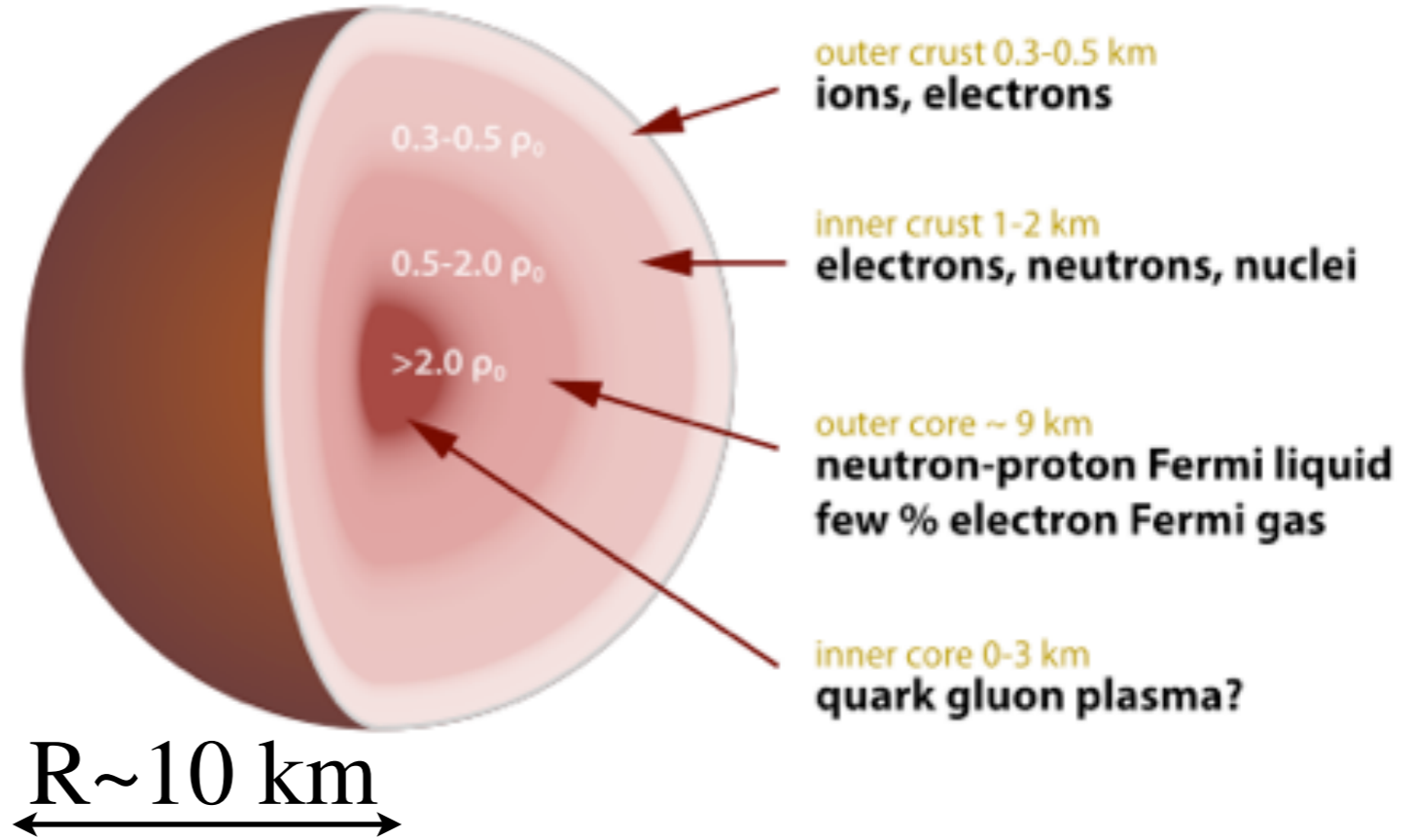
$$M \sim 1.5M_{\odot} = 3 \times 10^{30} \text{ kg}$$

Radius

$$R \approx 10 \text{ km}$$

Mass density

$$\rho = \frac{M}{V} \approx 7.5 \times 10^{17} \text{ kg m}^{-3}$$



Neutron star 101

Mass

$$M \sim 1.5M_{\odot} = 3 \times 10^{30} \text{ kg}$$

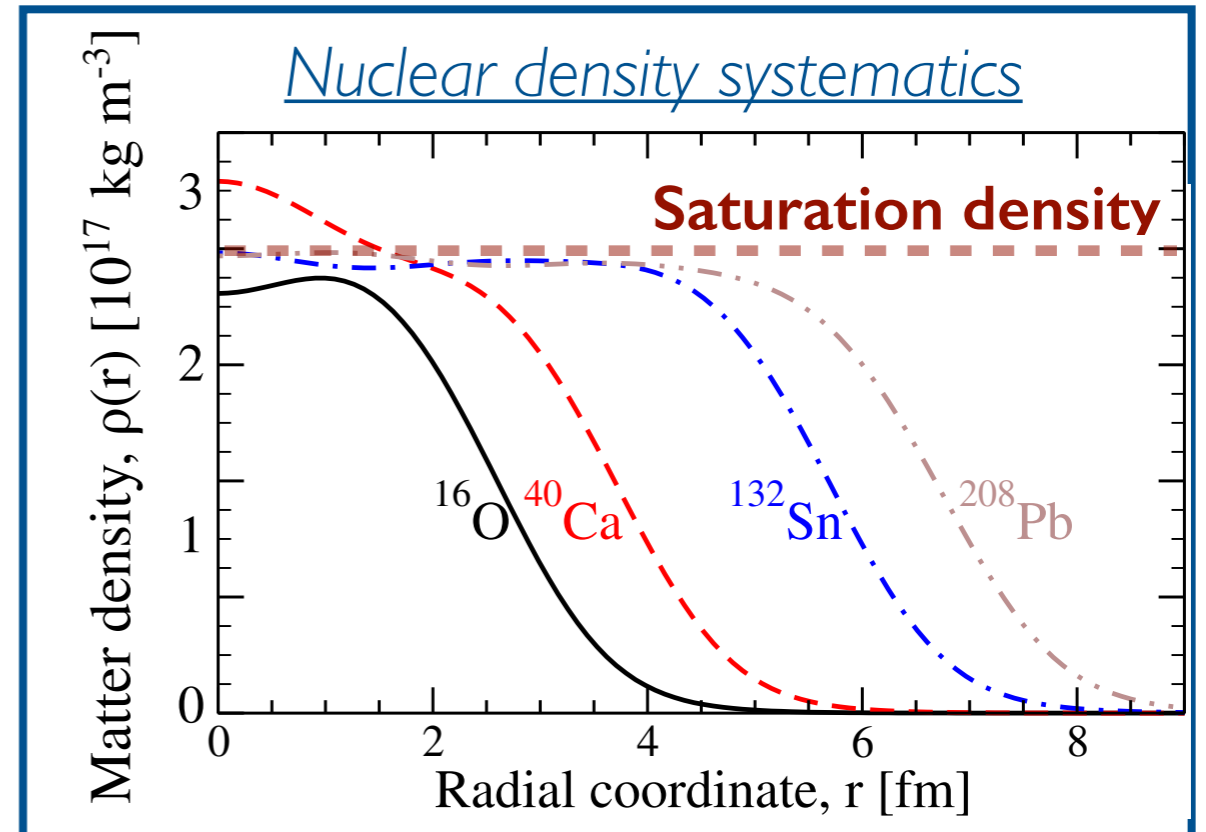
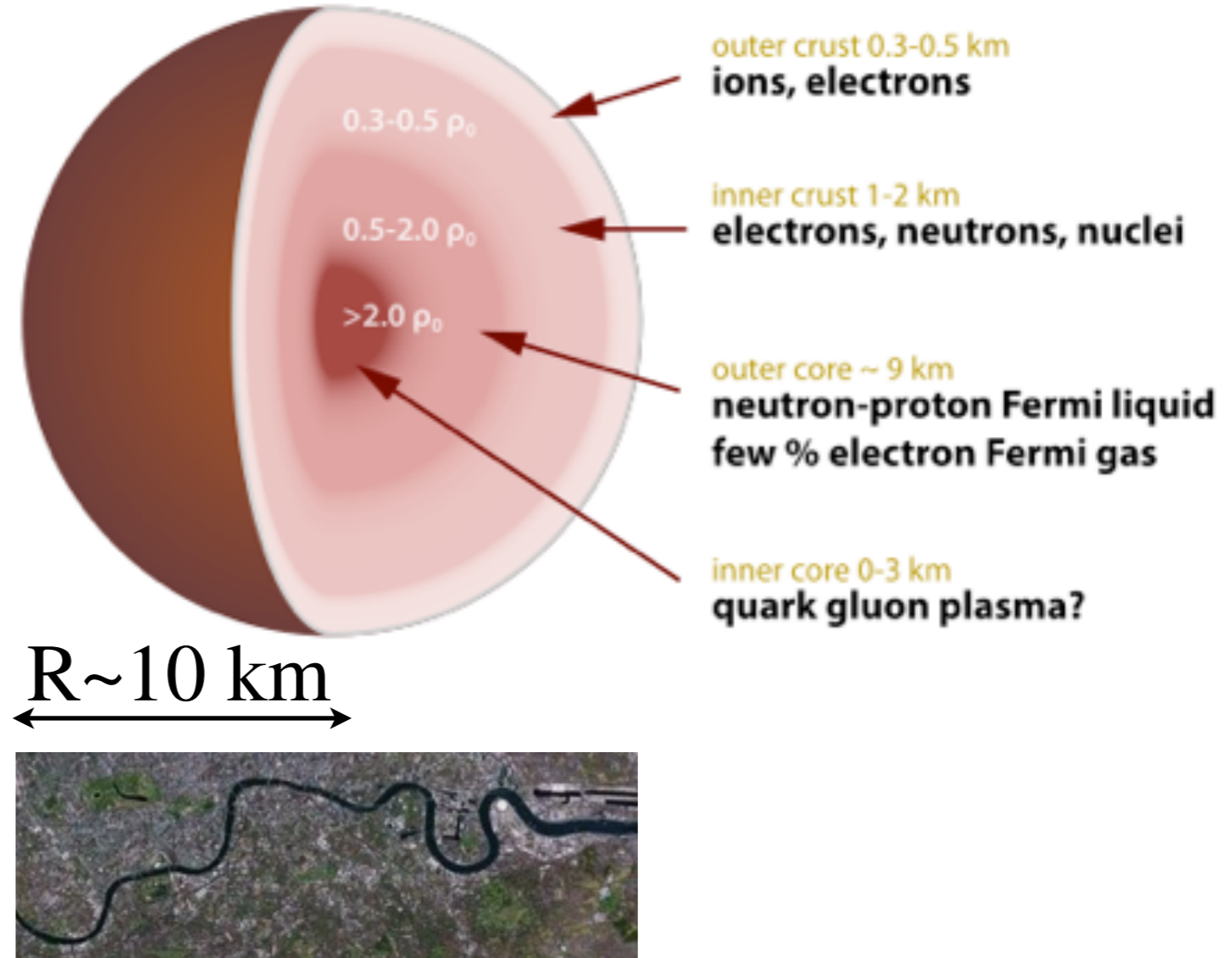
Radius

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Mass density

$$\rho = \frac{M}{V} \approx 7.5 \times 10^{17} \text{ kg m}^{-3}$$

A teaspoon weights a billions of tons!



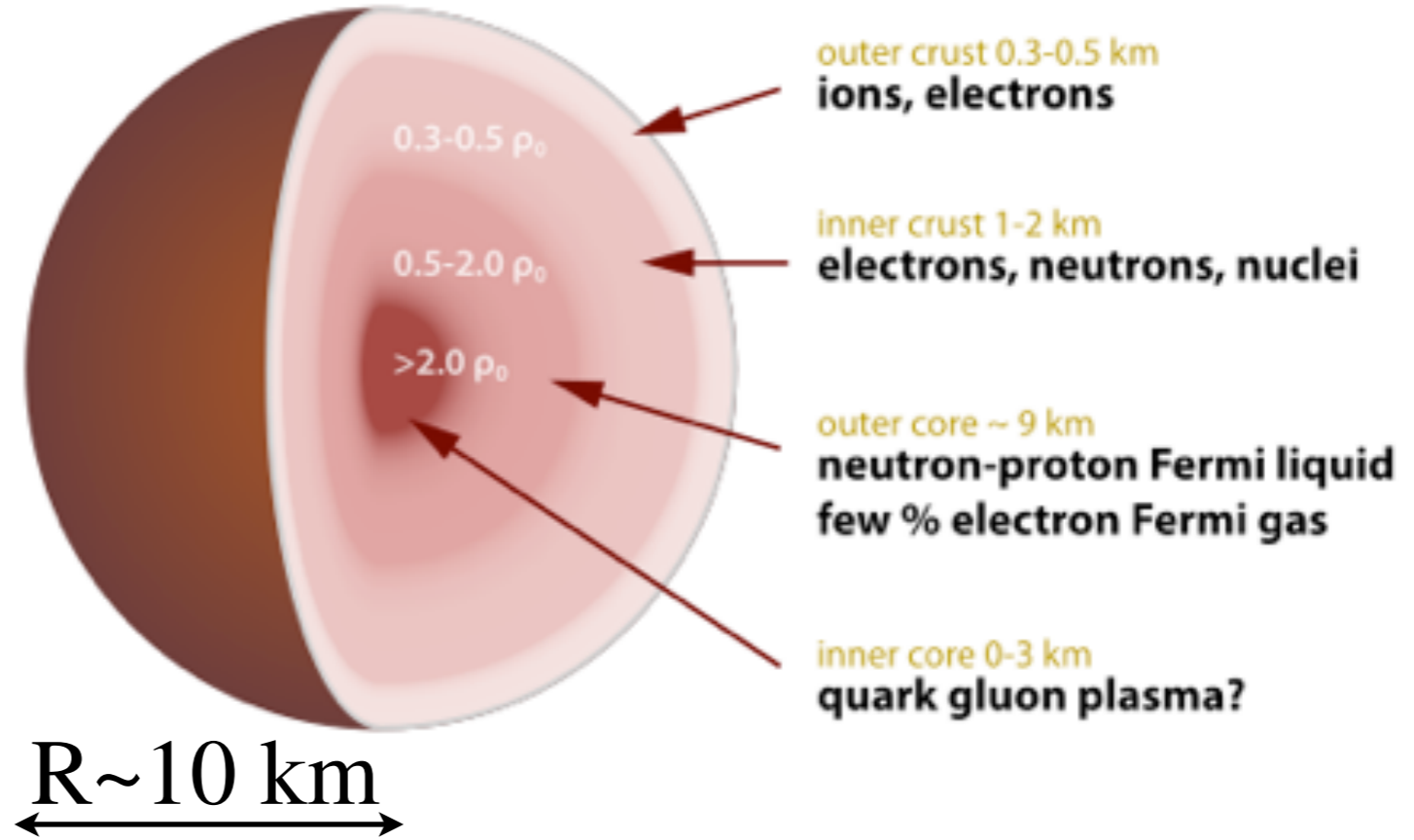
Neutron star 101

Mass

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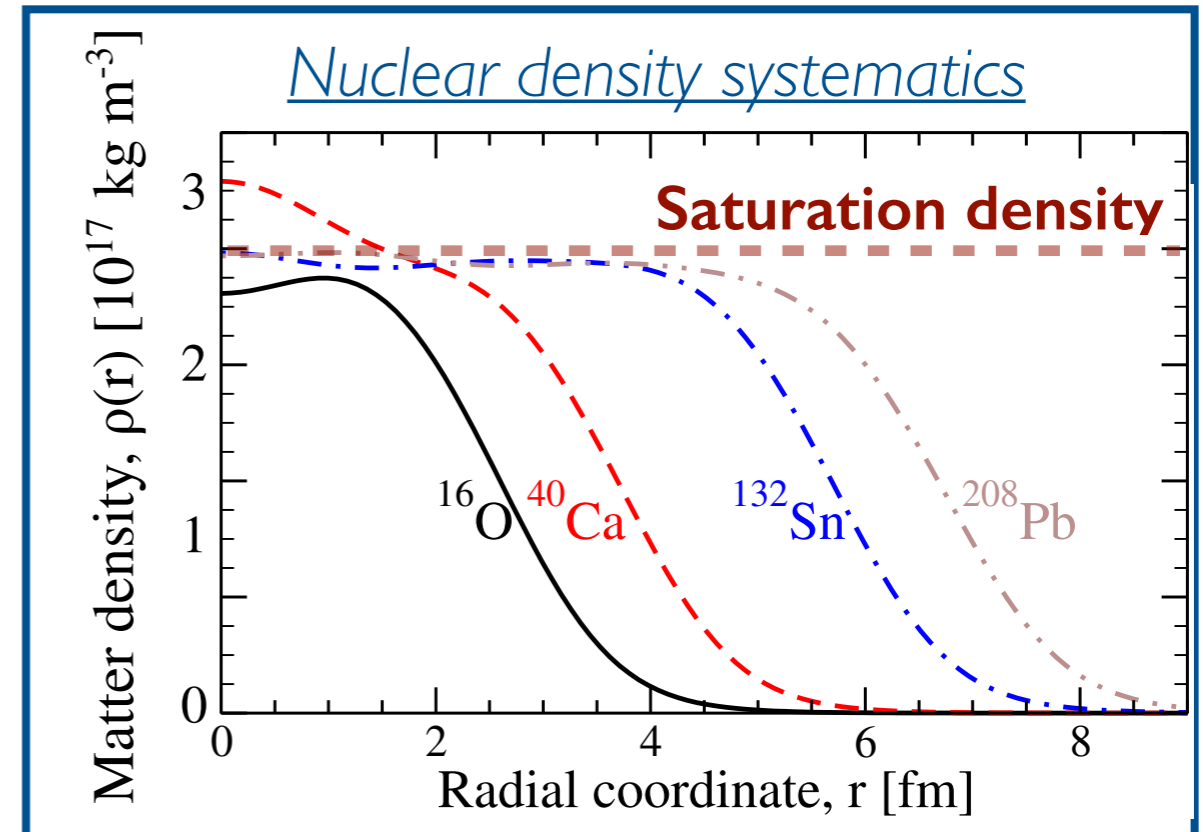
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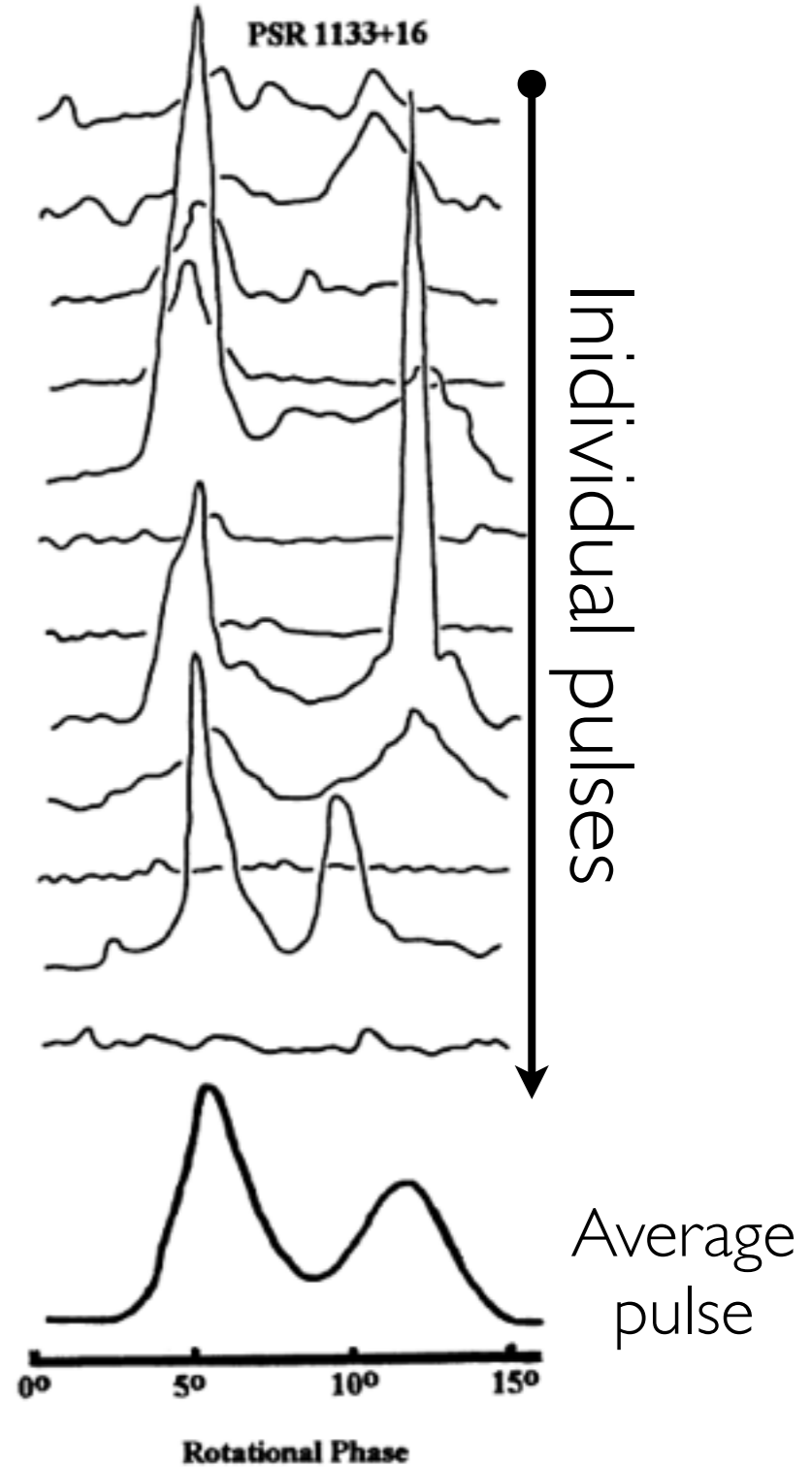
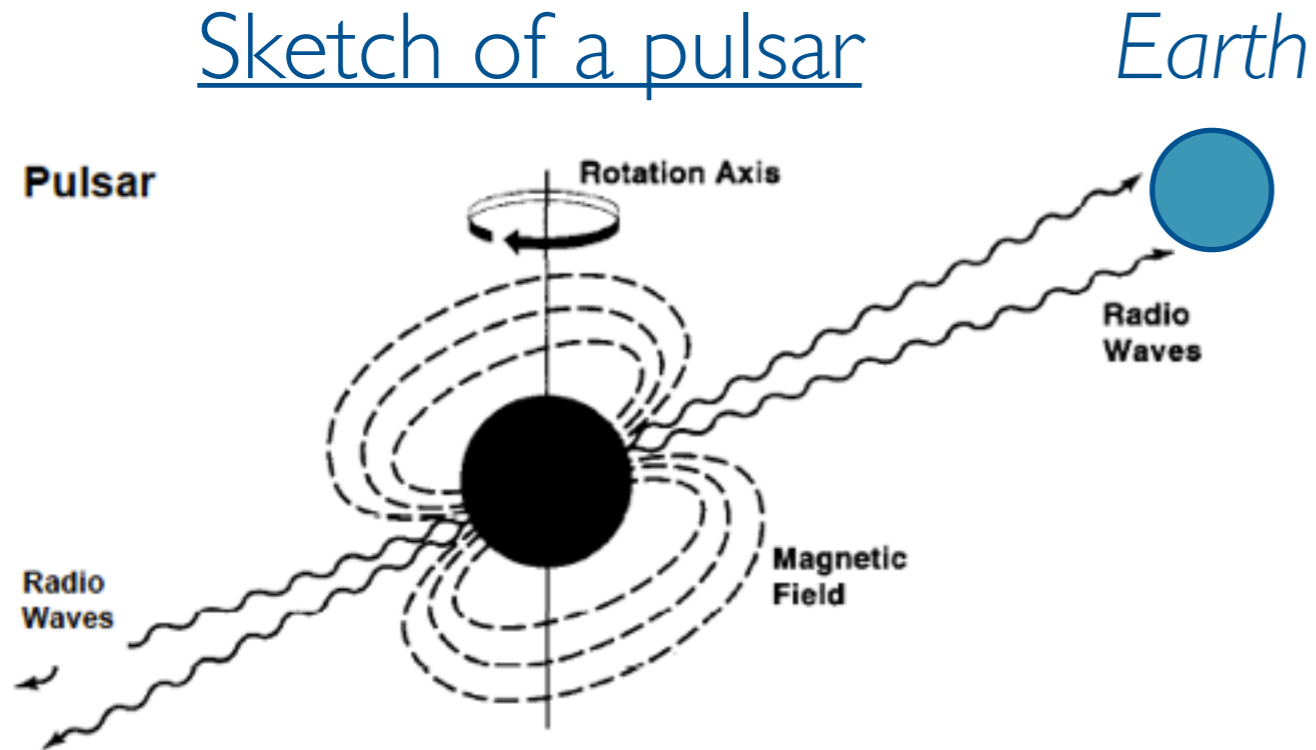
Escape velocity

$$v = \sqrt{\frac{2GM}{R}} \sim 0.5c$$



How do we find them?

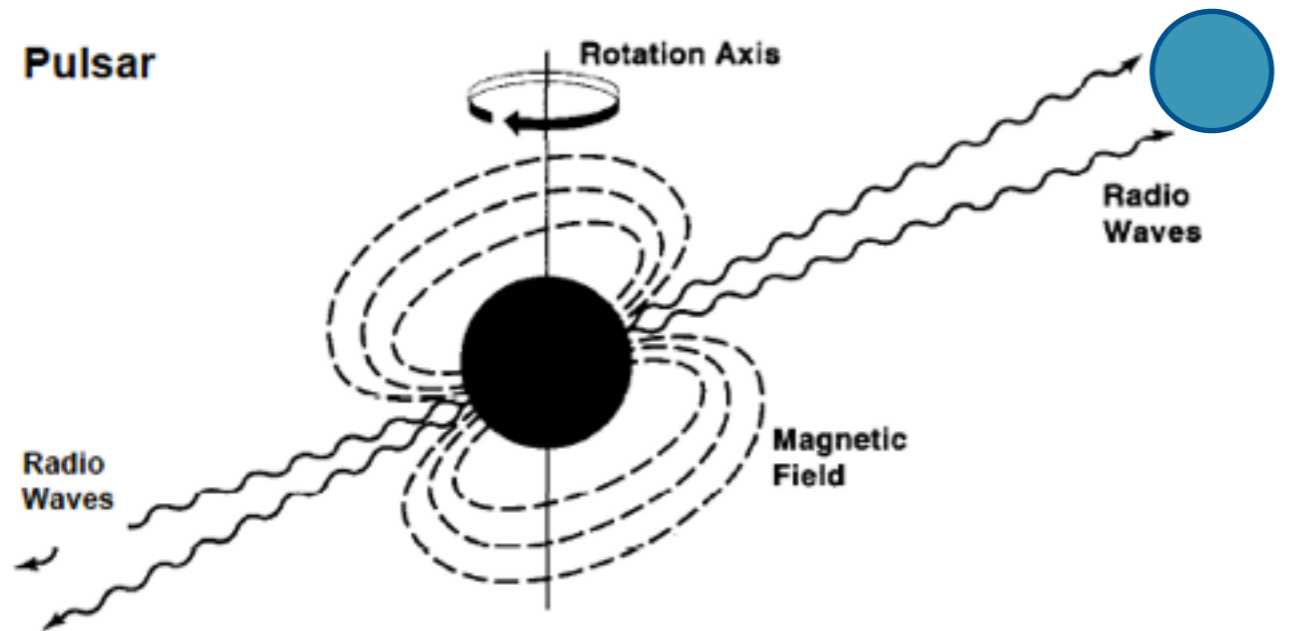
Sketch of a pulsar



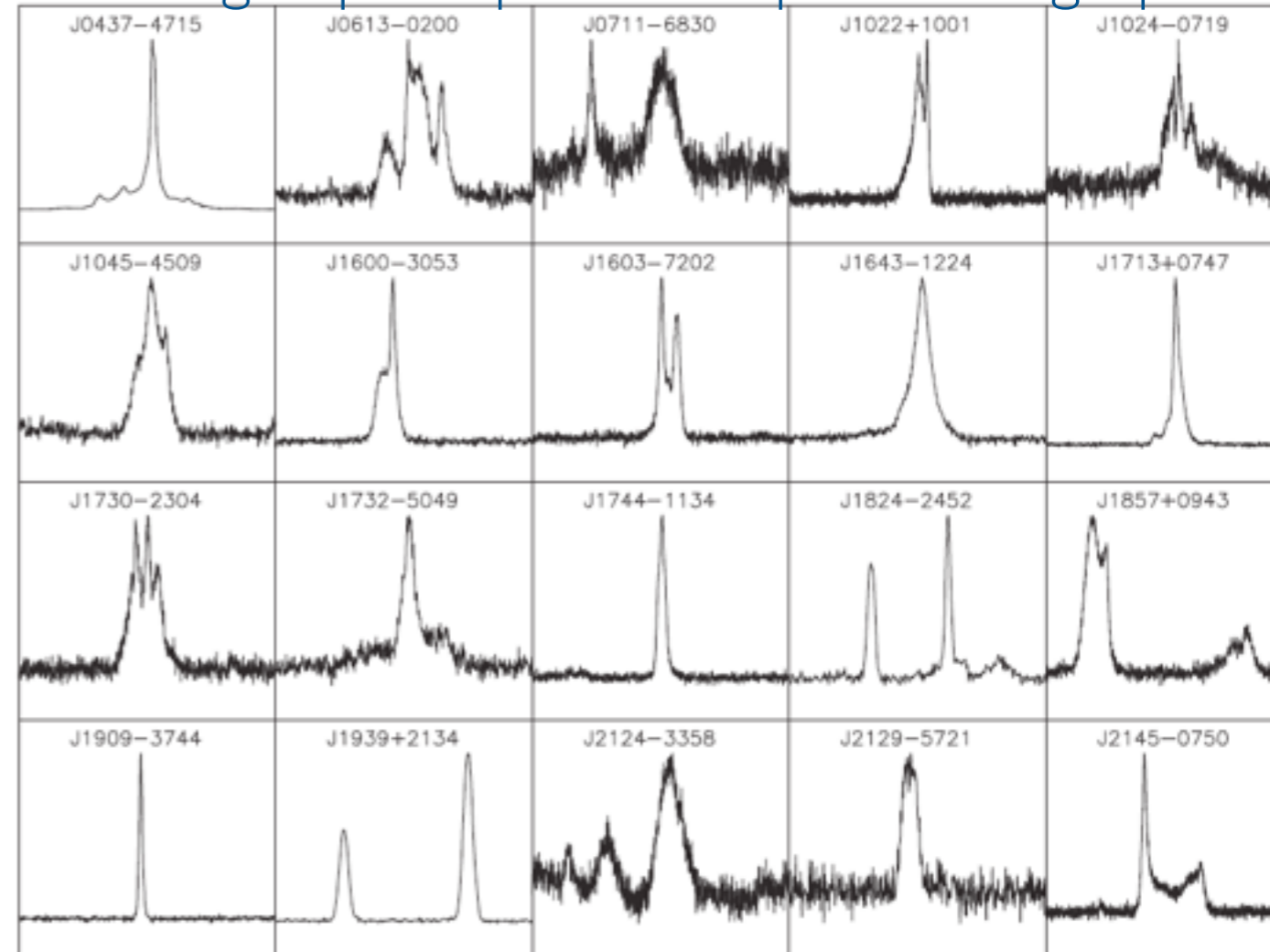
- **Highest** magnetic fields in nature: 10^{16} G
- Particles **accelerated** along field lines
- Rotation is **fast** and **accurate**, unique of each pulsar
- Radio & X-ray bursts point our way

How do we find them?

Sketch of a pulsar



Averaged pulse profile - a pulsar's fingerprint

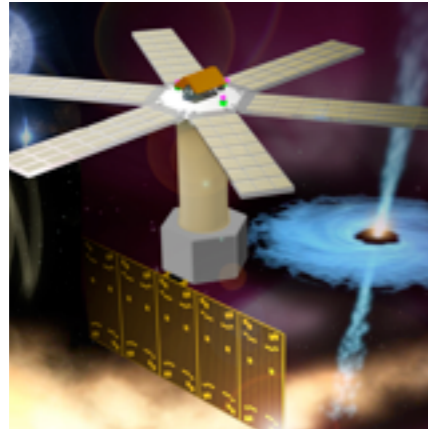


Hobbs et al., Pub. Astr. Soc. Aust. 202, 28 (2011)

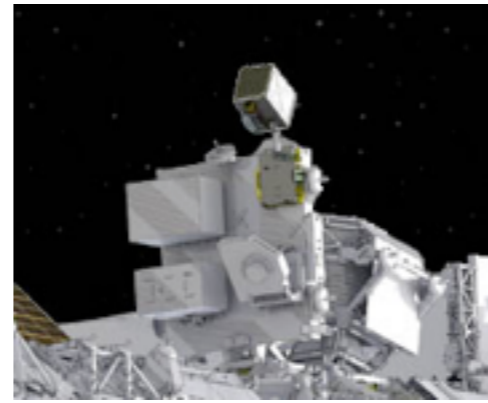
- **Highest** magnetic fields in nature: 10^{16} G
- Particles **accelerated** along field lines
- Rotation is **fast** and **accurate**, unique of each pulsar
- Radio & X-ray bursts point our way

EoS from future observations

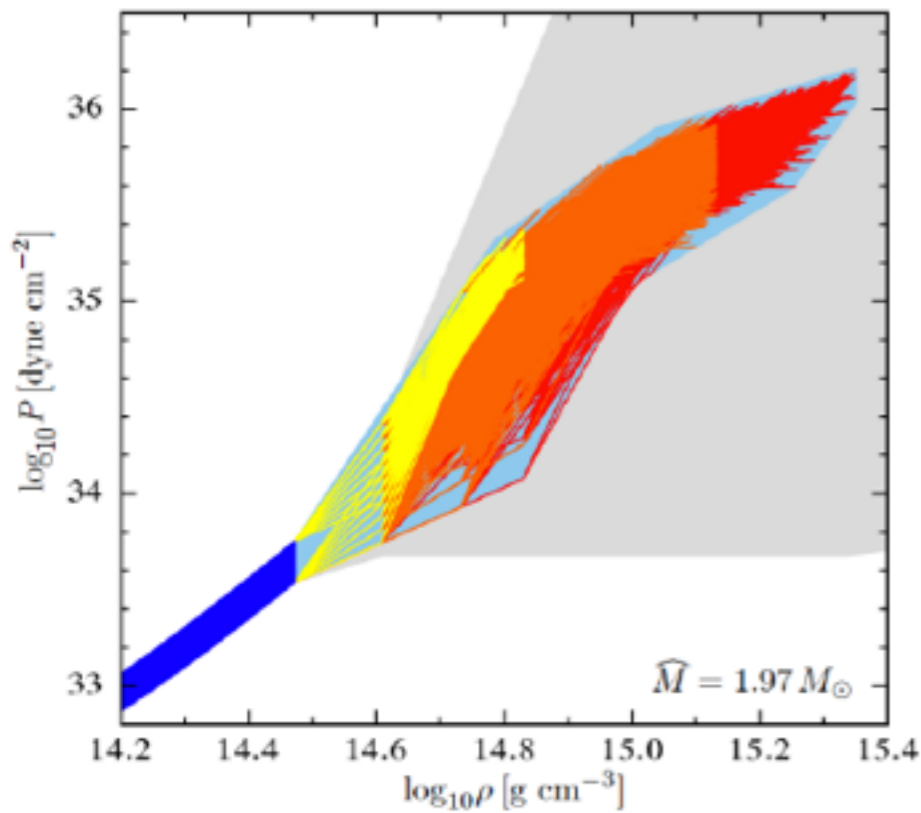
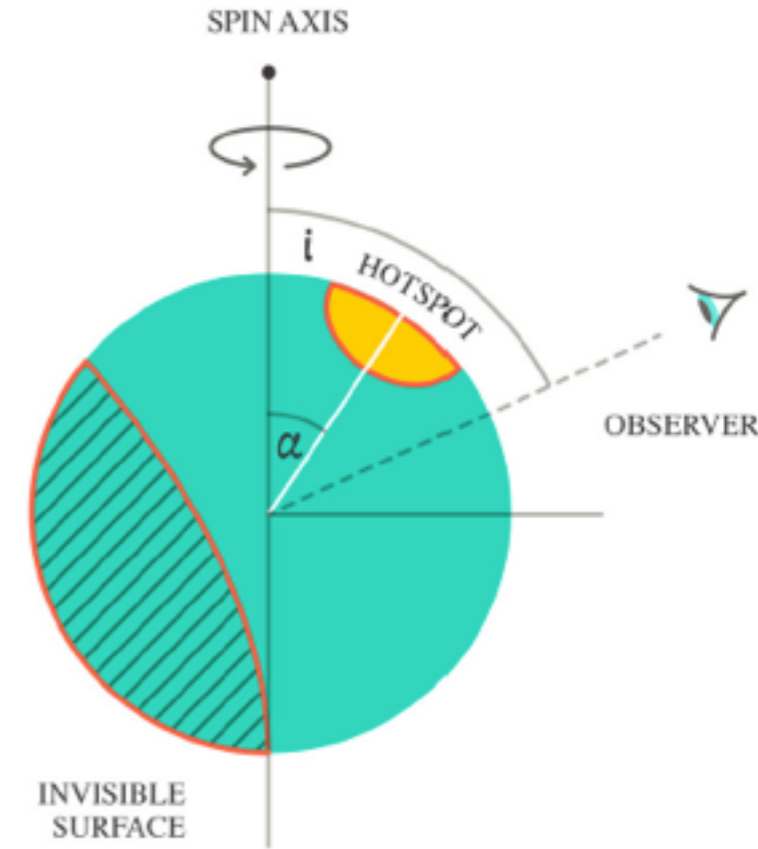
X-rays



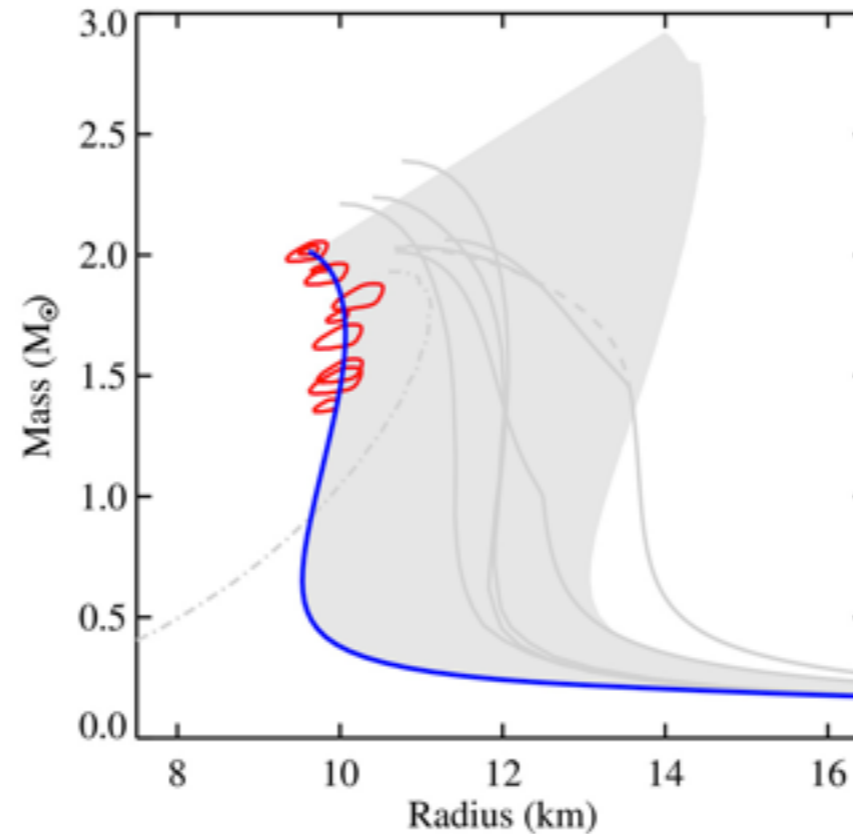
LOFT



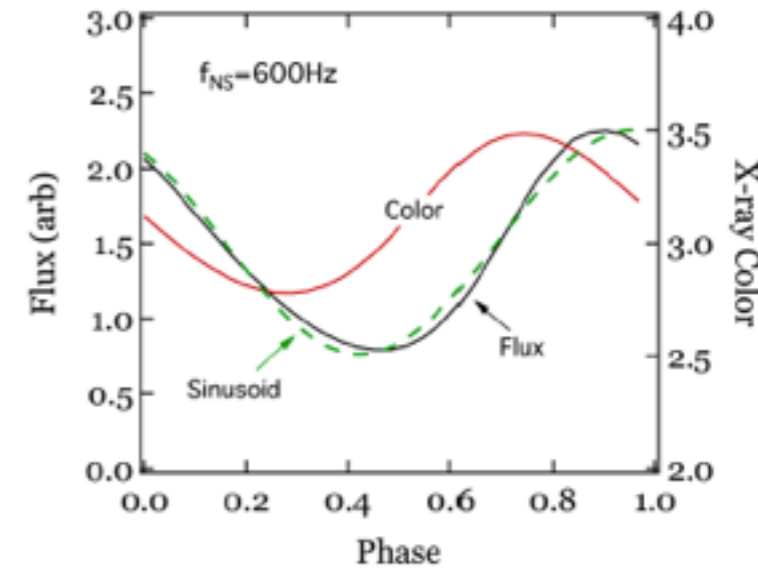
NICER



Hebeler et al. 2013



LOFT Yellow Book



- ESA X-ray observatory with dense matter program
- Burst oscillation of known sources will yield M-R to few % level
- Did not make it as M3... Might be put forward as M4?

Neutron matter mass and radius

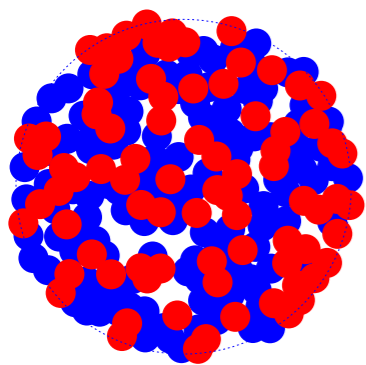
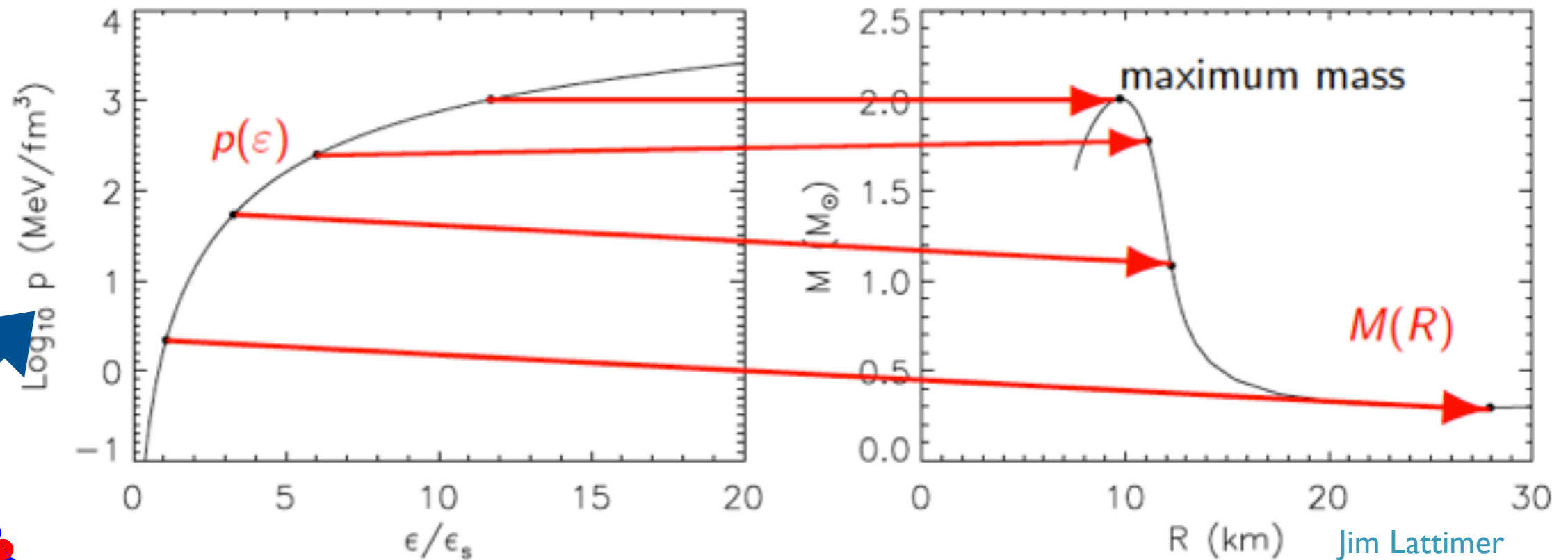
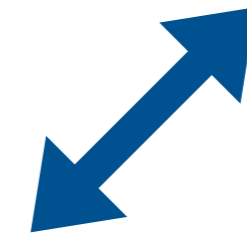
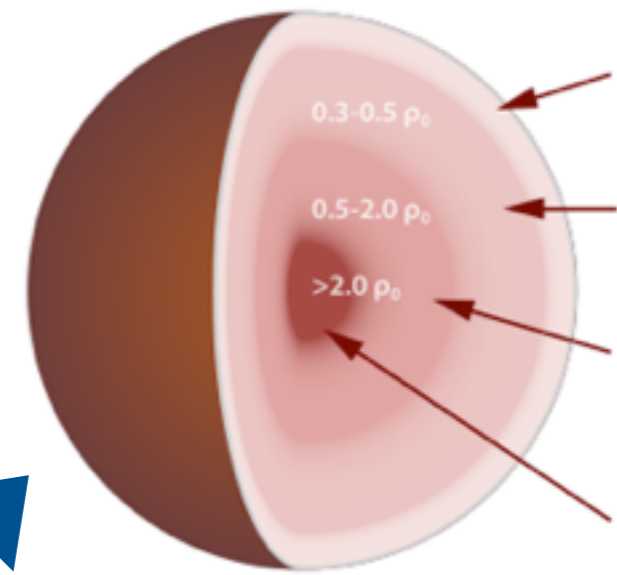
Tolman-Oppenheimer-Volkov equations

$$\frac{dp}{dr} = -\frac{G}{c^2} \frac{(m + 4\pi r^3)(\epsilon + p)}{r(r - 2Gm/c^2)}$$

$$\frac{dm}{dr} = \frac{4\pi}{c^2} \epsilon r^2$$

Equation of State

Mass-Radius relation



Jim Lattimer

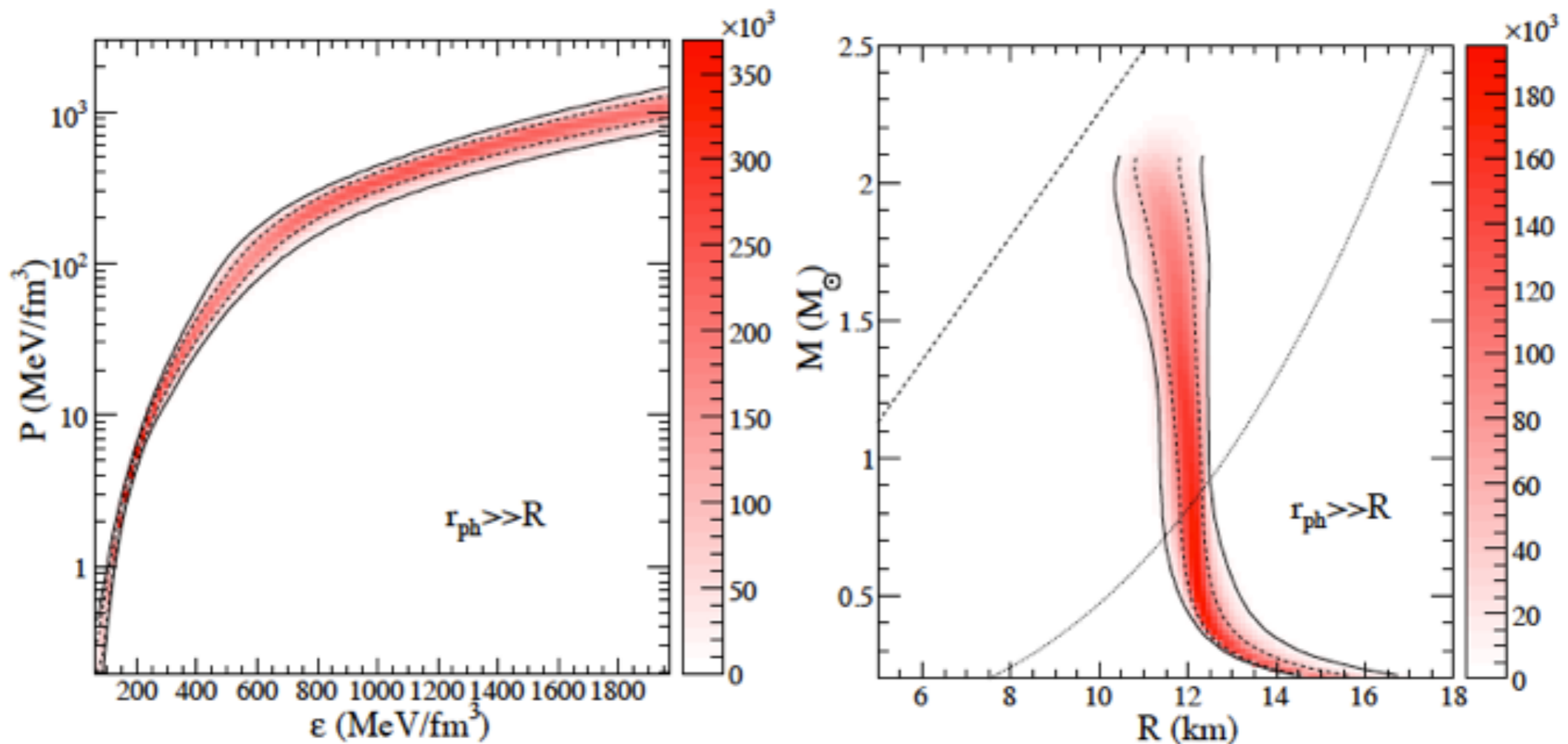
Neutron matter mass and radius

Tolman-Oppenheimer-Volkov equations

$$\frac{dp}{dr} = -\frac{G}{c^2} \frac{(m + 4\pi p r^3)(\epsilon + p)}{r(r - 2Gm/c^2)}$$

$$\frac{dm}{dr} = \frac{4\pi}{c^2} \epsilon r^2$$

Inferred EoS and M-R relation from observations



Steiner, Lattimer & Brown, ApJ **722**, 33 (2010)

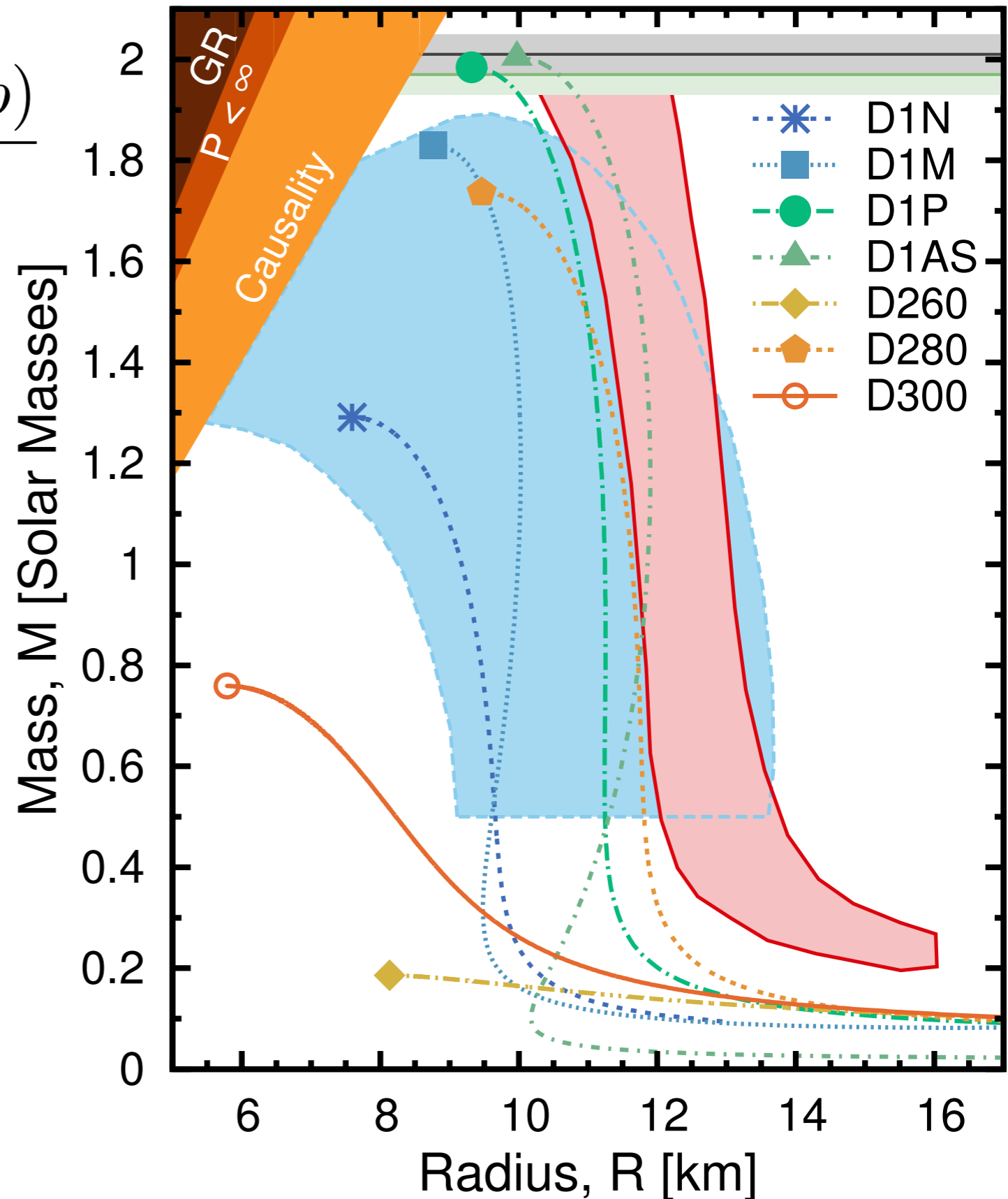
3 X-ray bursts, 3 X-ray binaries & 1 isolated NS

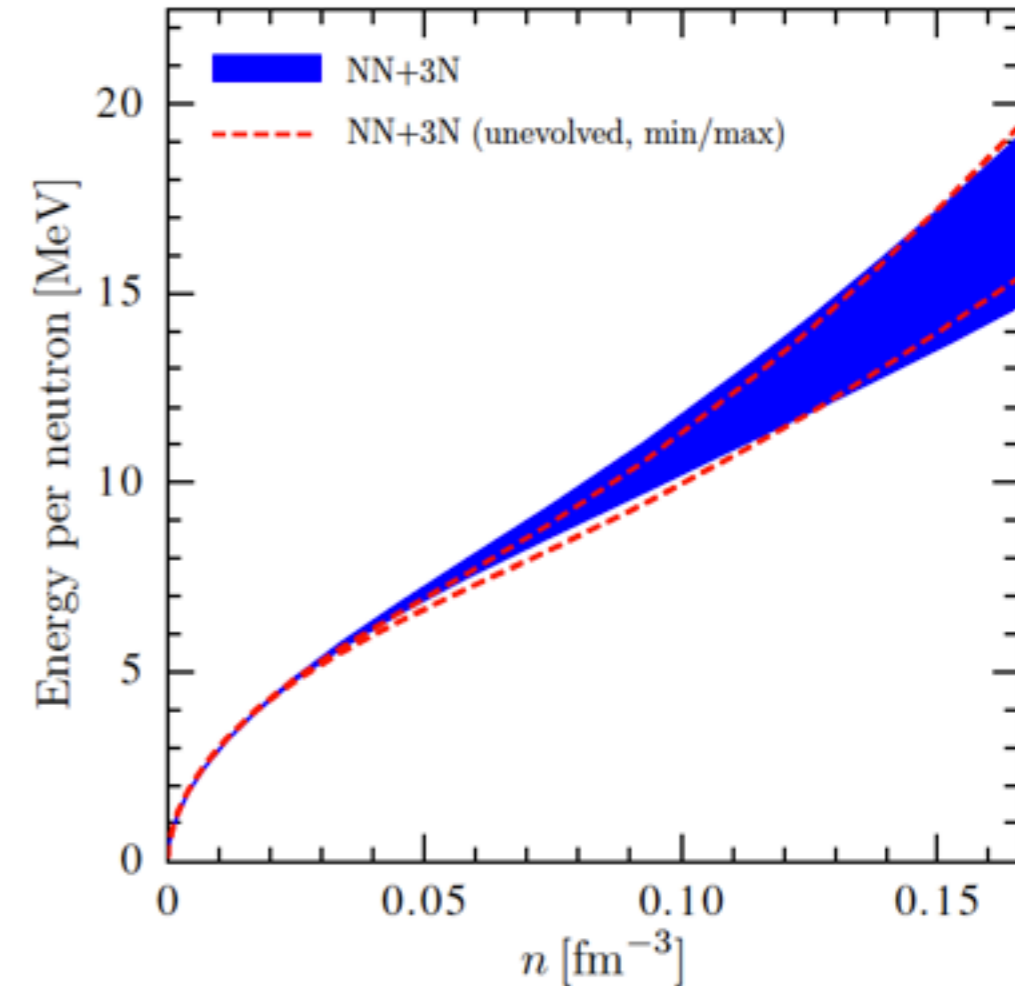
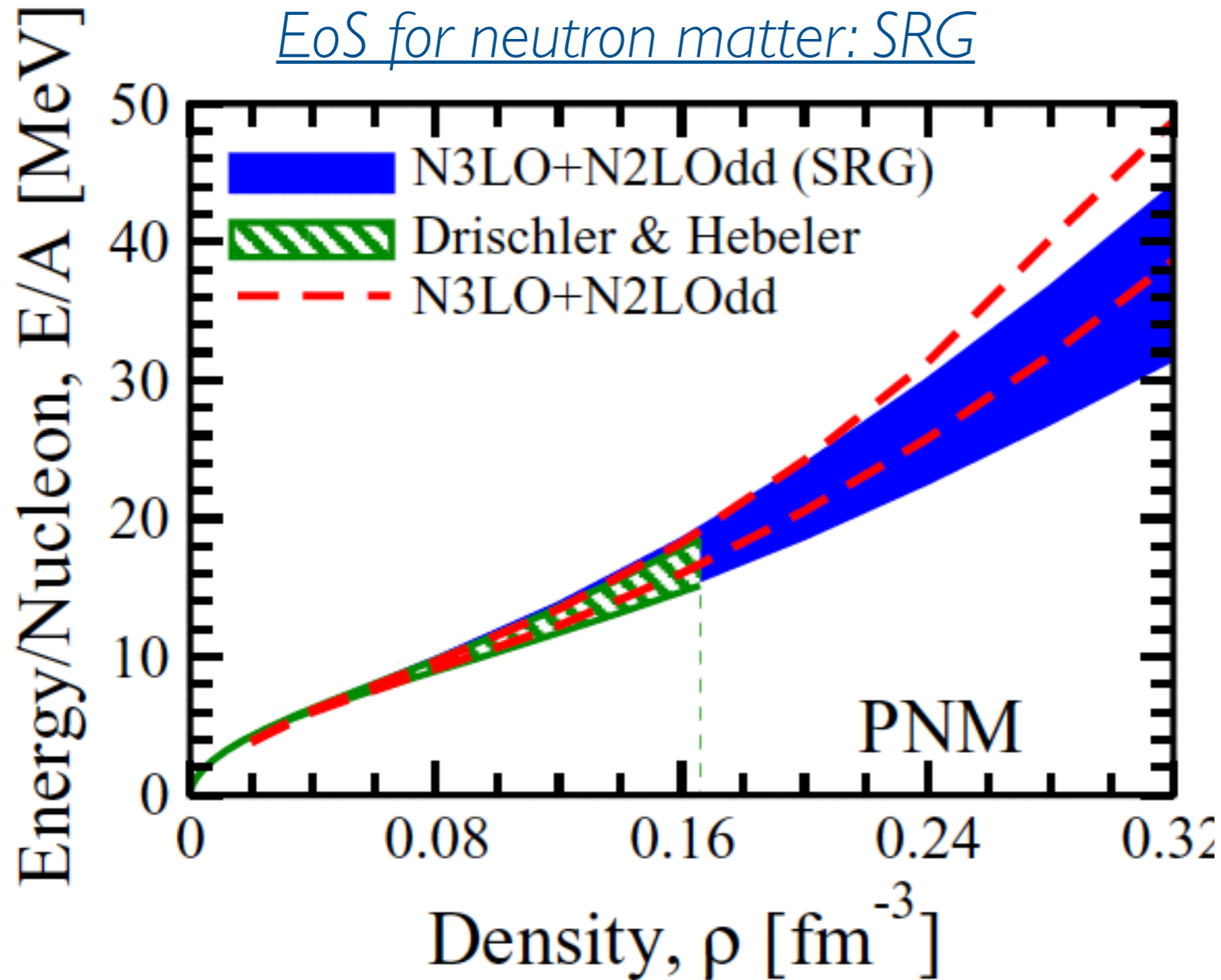
Neutron matter mass and radius

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$$\frac{dp}{dr} = -\frac{G}{c^2} \frac{(m + 4\pi pr^3)(\epsilon + p)}{r(r - 2Gm/c^2)}$$

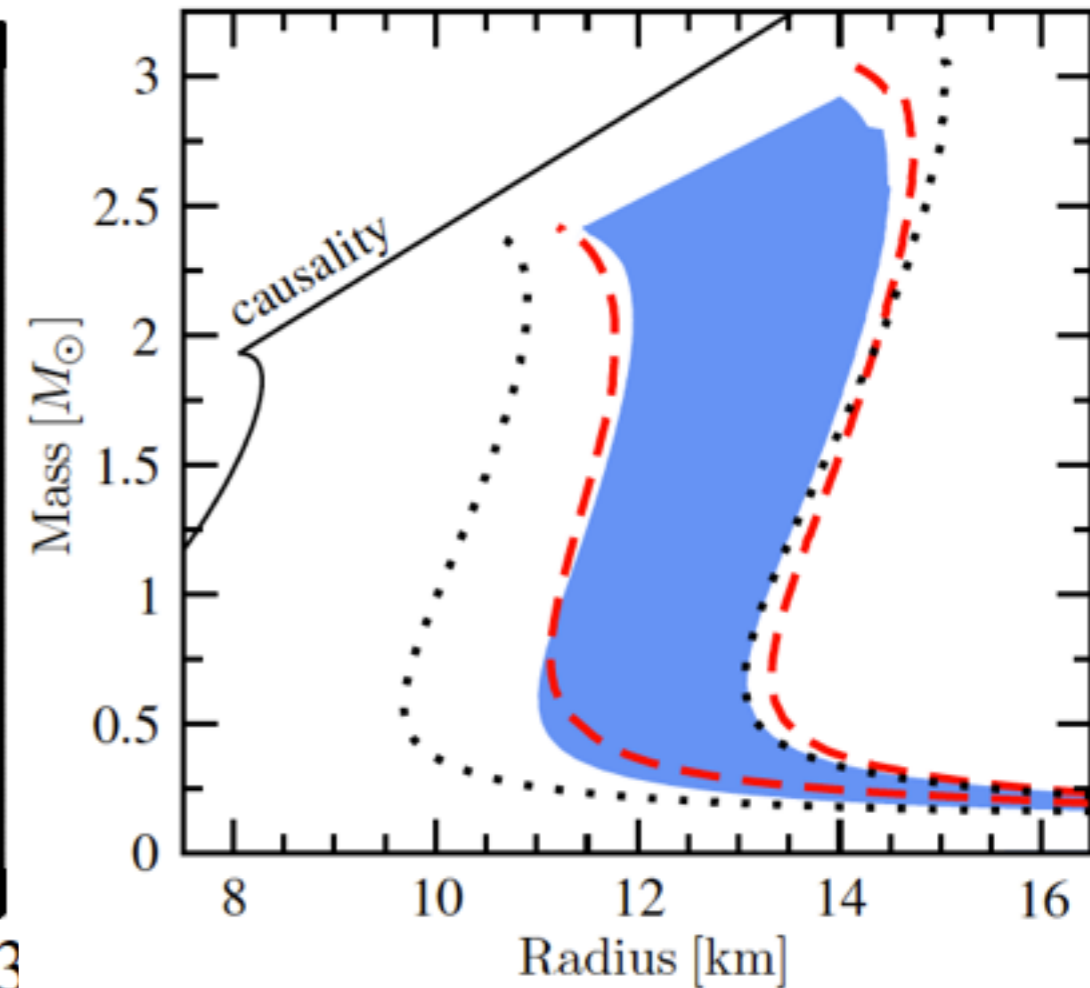
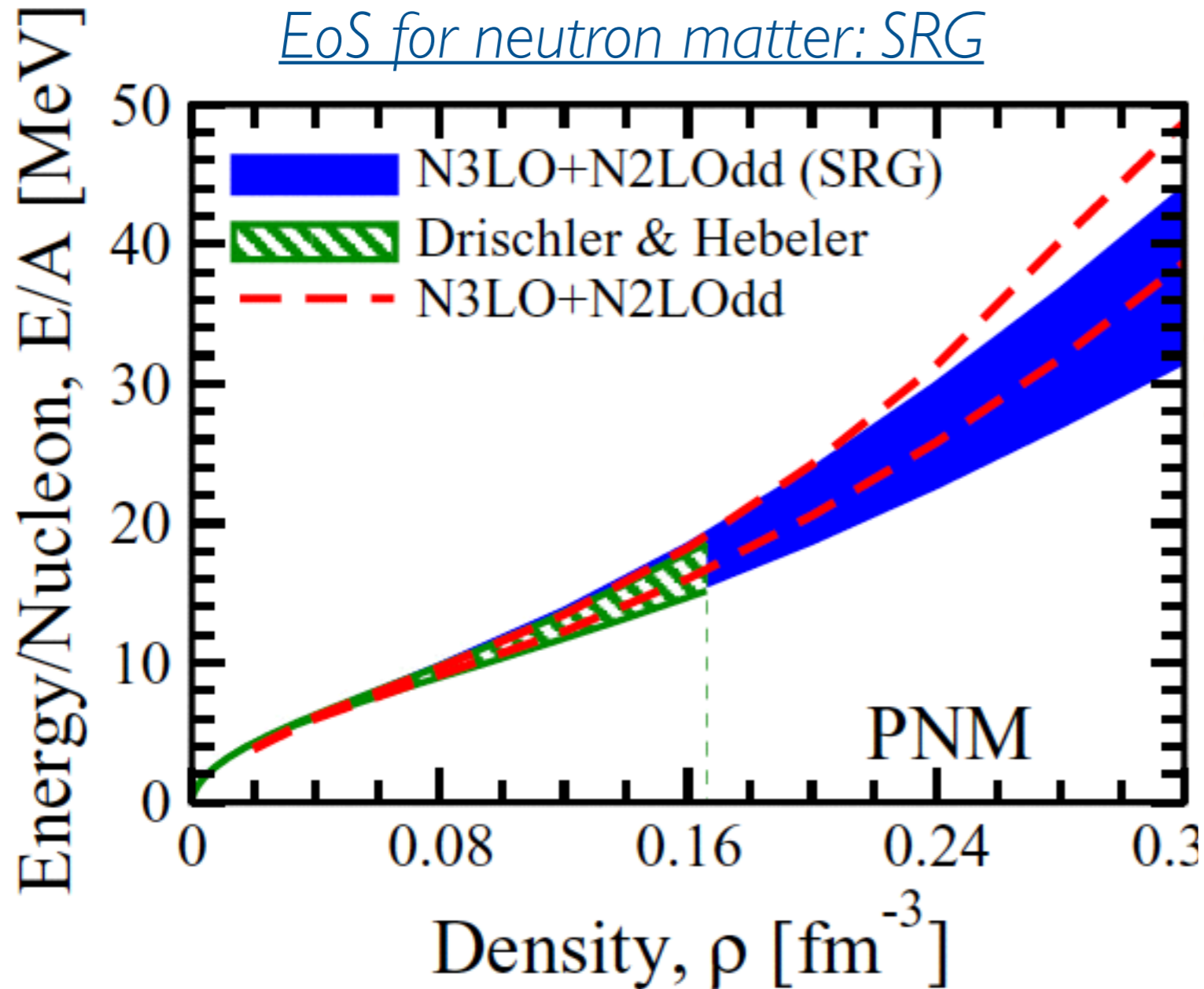
$$\frac{dm}{dr} = \frac{4\pi}{c^2} \epsilon r^2$$





Hebelner, Lattimer, Pethick, Schwenk
ApJ **773** 11 (2013)

- Error band from unknown ChPT c_1, c_3 parameters
- Finite temperature & higher densities available



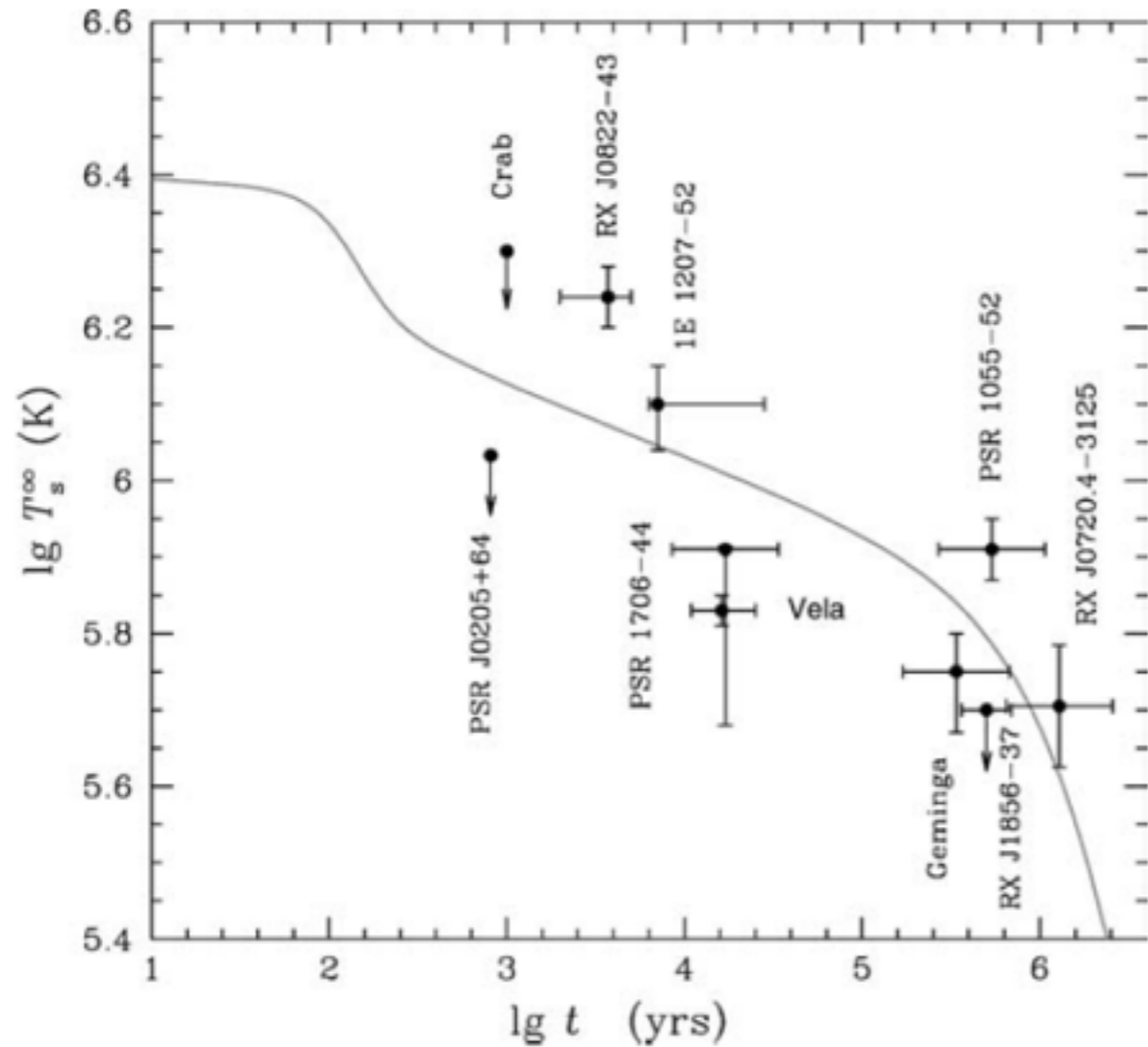
Hebeler, Lattimer, Pethick, Schwenk
ApJ **773** 11 (2013)

- Error band from unknown ChPT c_1, c_3 parameters
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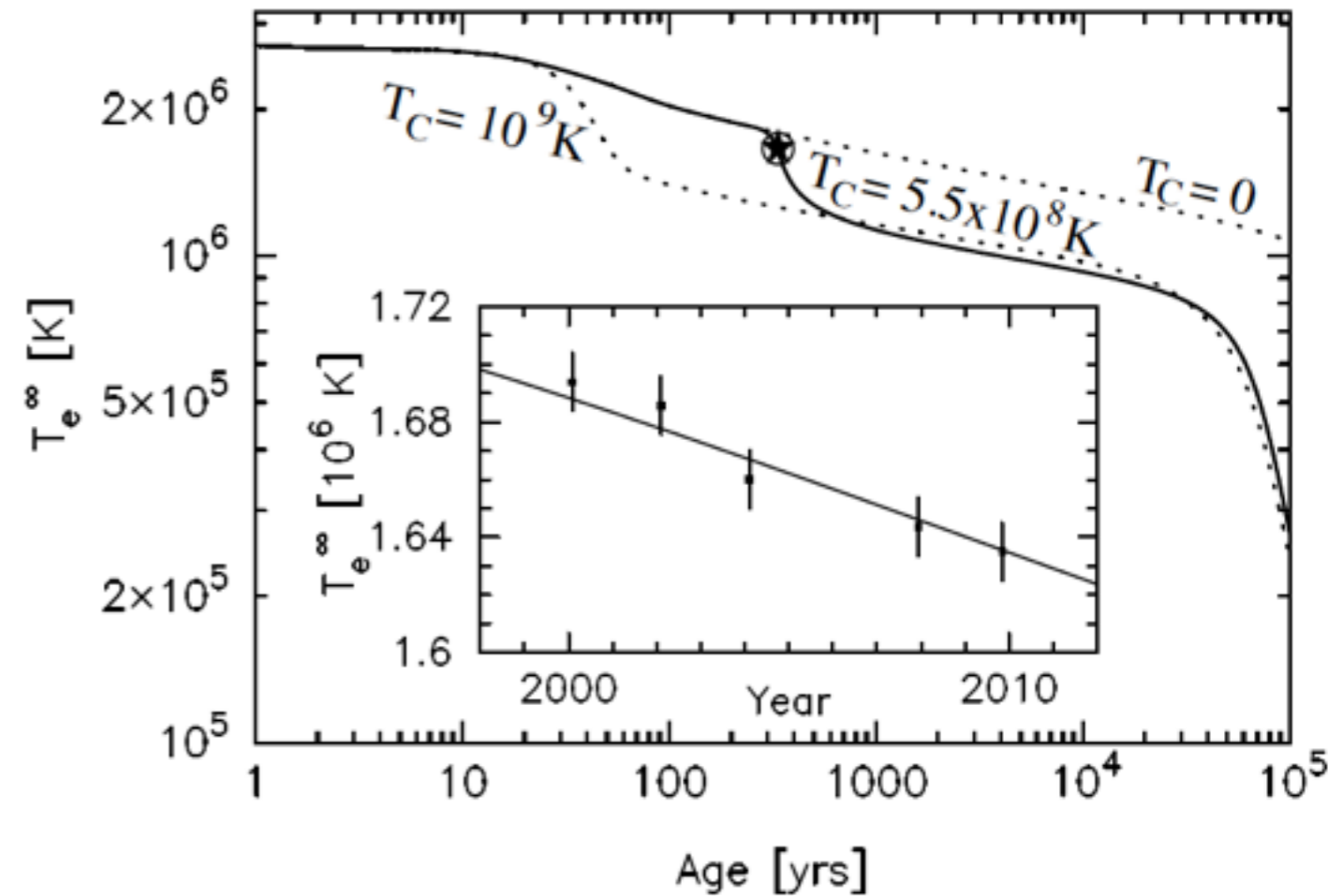
Cooling of neutron stars

NS cooling curves ($M=1.3M_{\odot}$)

3PF_2 pairing & CasA



Yakovlev & Pethick, ARAA **42** 169 (2004)

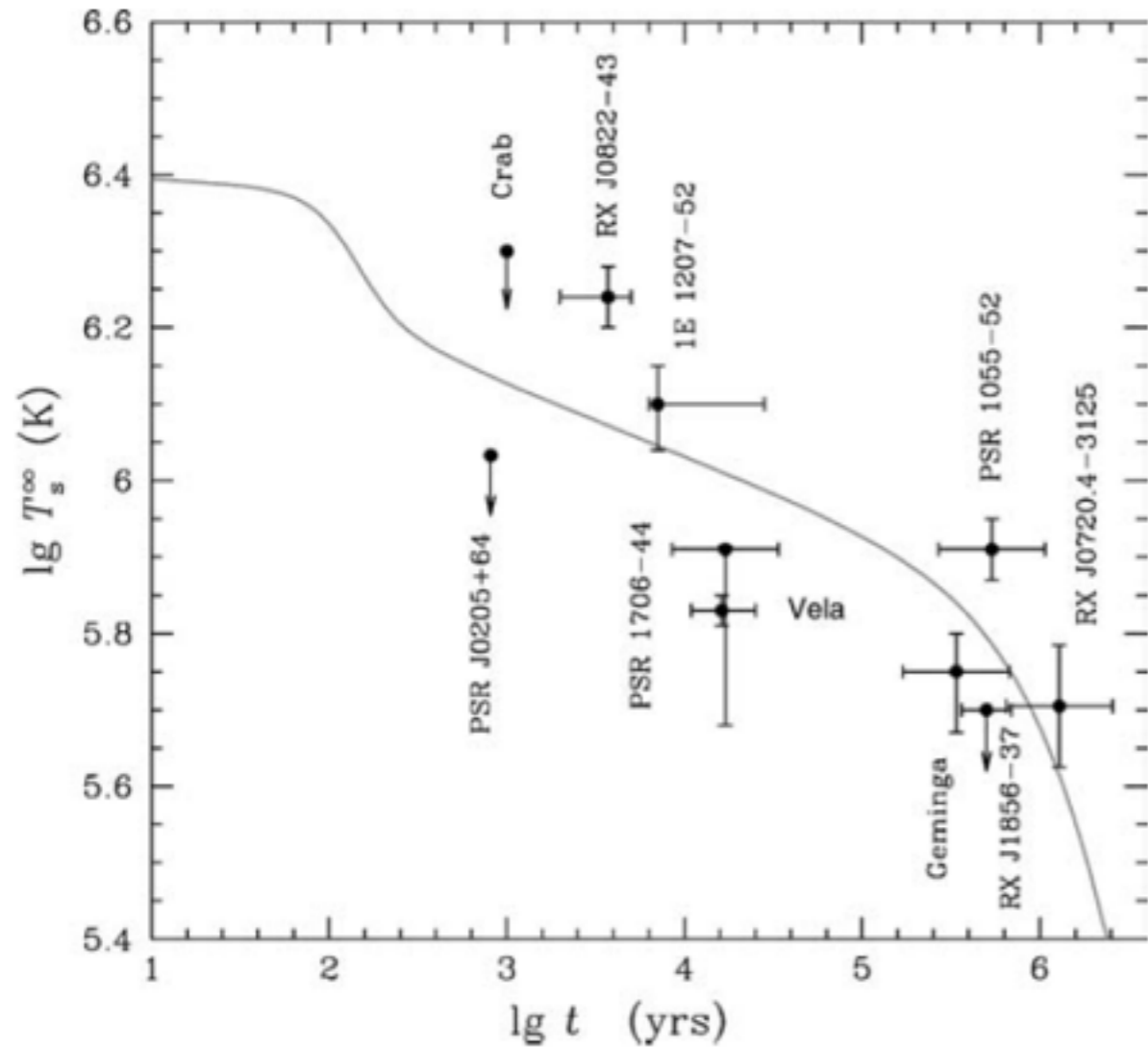


Page et al., PRL **106** 081101 (2011)

- Observational data available for a handful of systems
- Sensitive to **interior** physics!

Cooling of neutron stars

NS cooling curves ($M=1.3M_{\odot}$)



Yakovlev & Pethick, ARAA **42** 169 (2004)

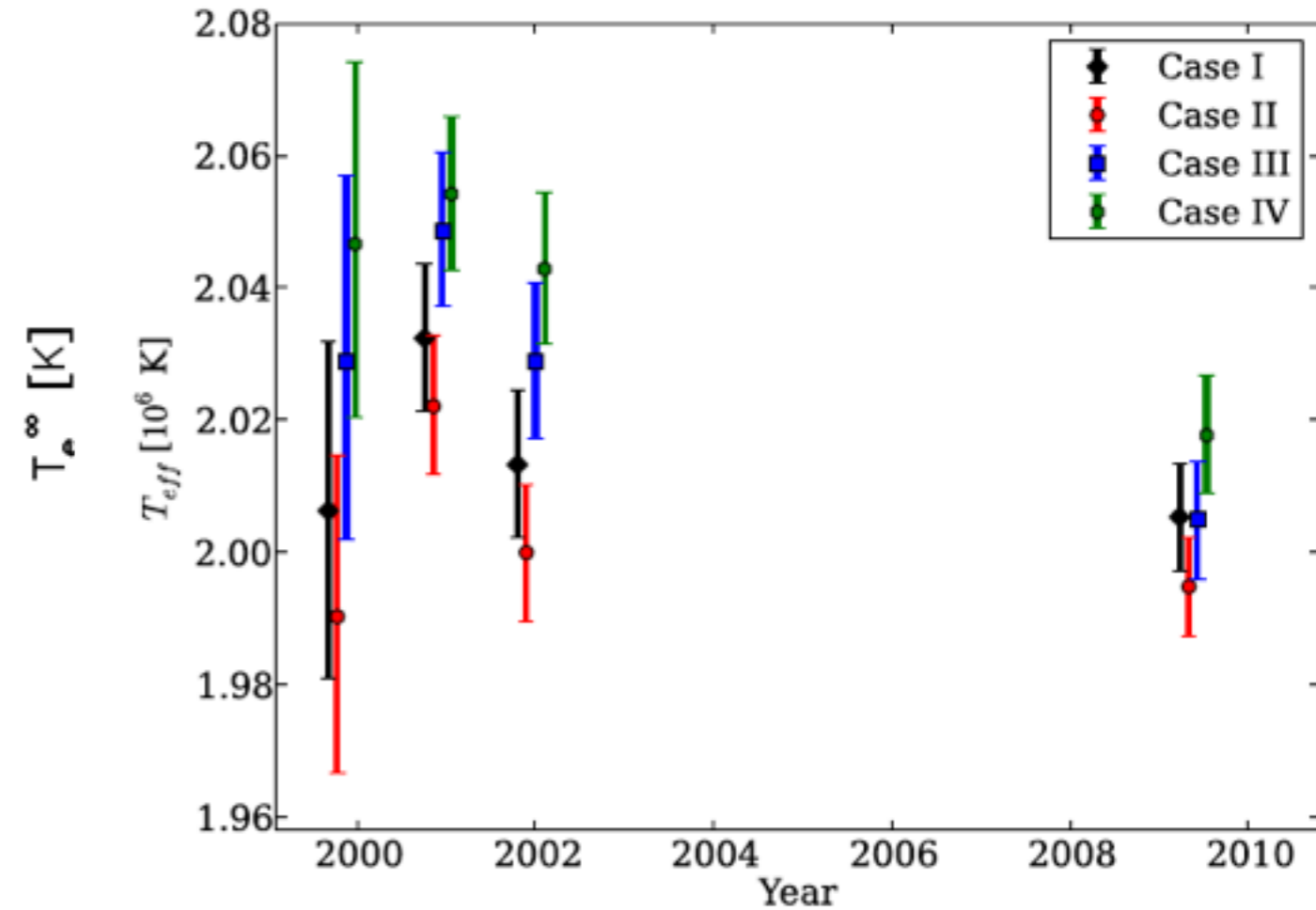
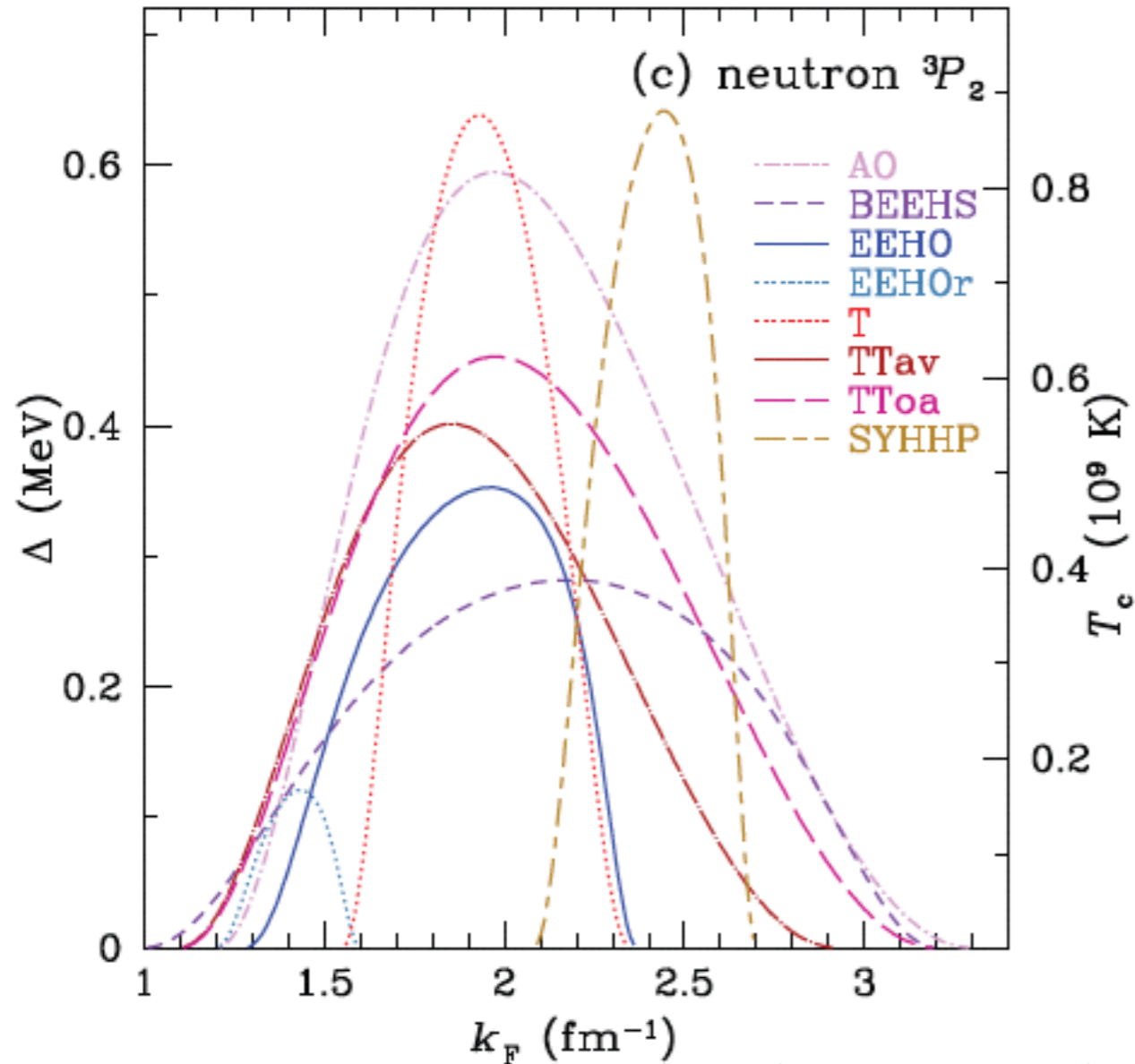


Figure 2. Inferred temperatures from HRC-S count rates for the NS in Cas A with different cases of source and background extraction regions (see Table 1 for case definitions). Cases II–IV have been shifted by a small offset in time (+0.1, +0.2, +0.3, respectively) to make them easier to distinguish. The temperature decline over 10 yr, for different cases, ranges from $0.9\% \pm 0.6\%$ ($\chi^2 = 2.7$ for 2 dof) to $2\% \pm 0.7\%$ ($\chi^2 = 1.3$ for 2 dof). Our preferred value for comparison with other detectors, Case I, exhibits a temperature decline of $1.0\% \pm 0.7\%$ ($\chi^2 = 1.8$ for 2 dof).

- Observational data available for a handful of systems
- Sensitive to **interior** physics!

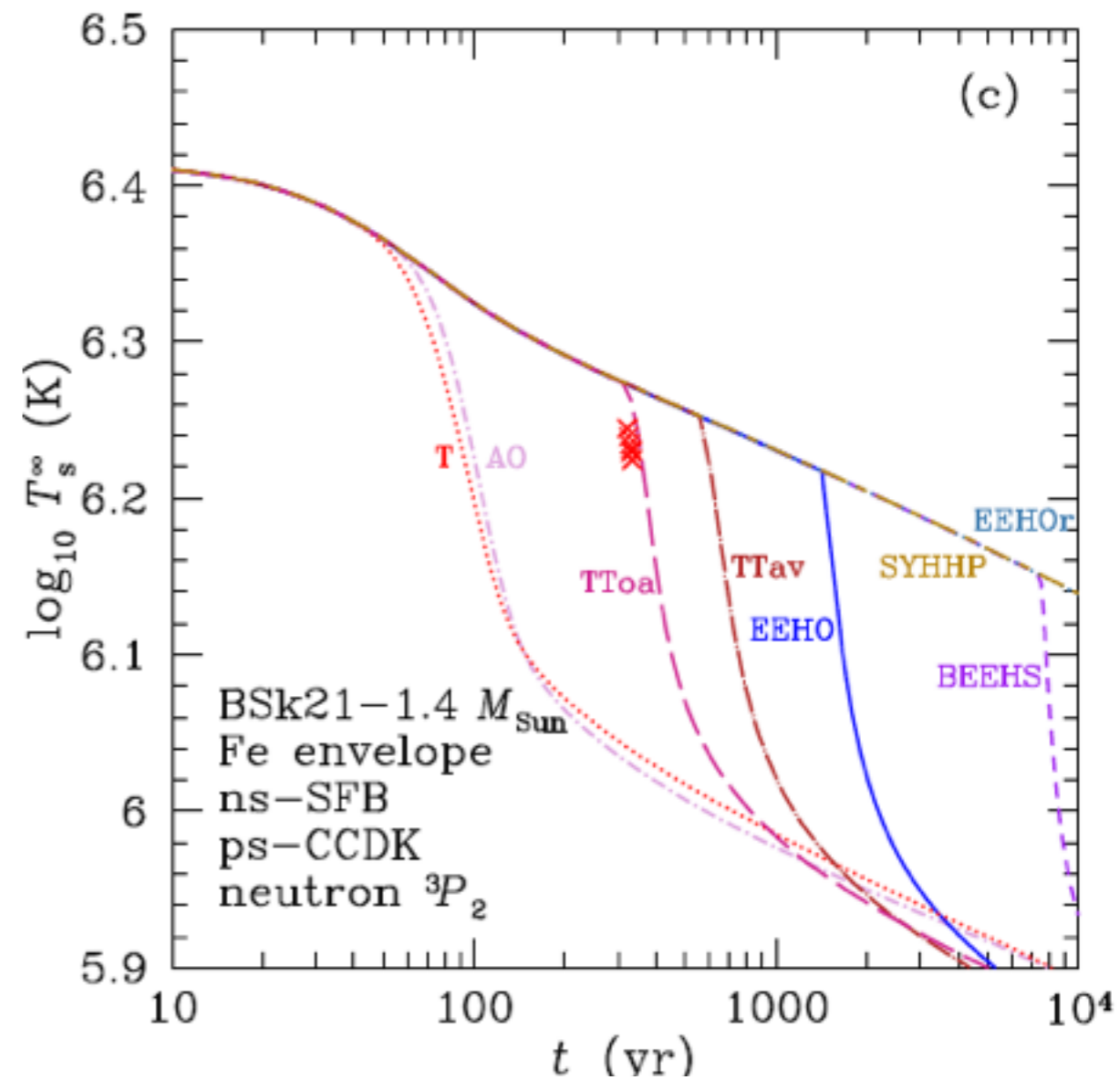
Cooling of Cassiopea A



Ho, Elshamouty, Heinke, Potekhin
PRC **91** 015806 (2015)

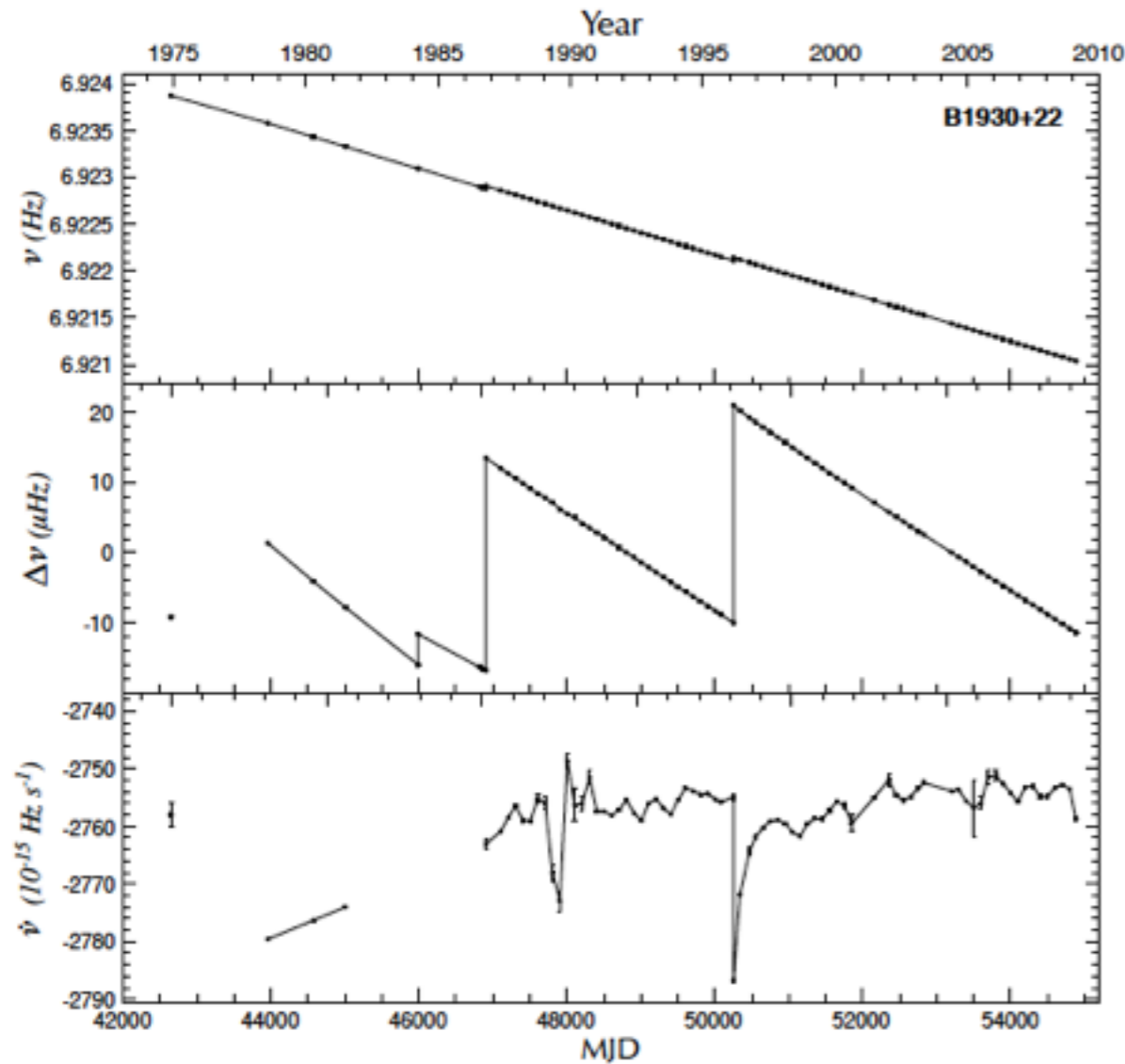
Ingredients

- (a) Mass of pulsar
- (b) EoS (determines radius)
- (c) Internal composition
- (d) Pairing gaps (1S_0 & 3P_2 channels)
- (e) Atmosphere composition

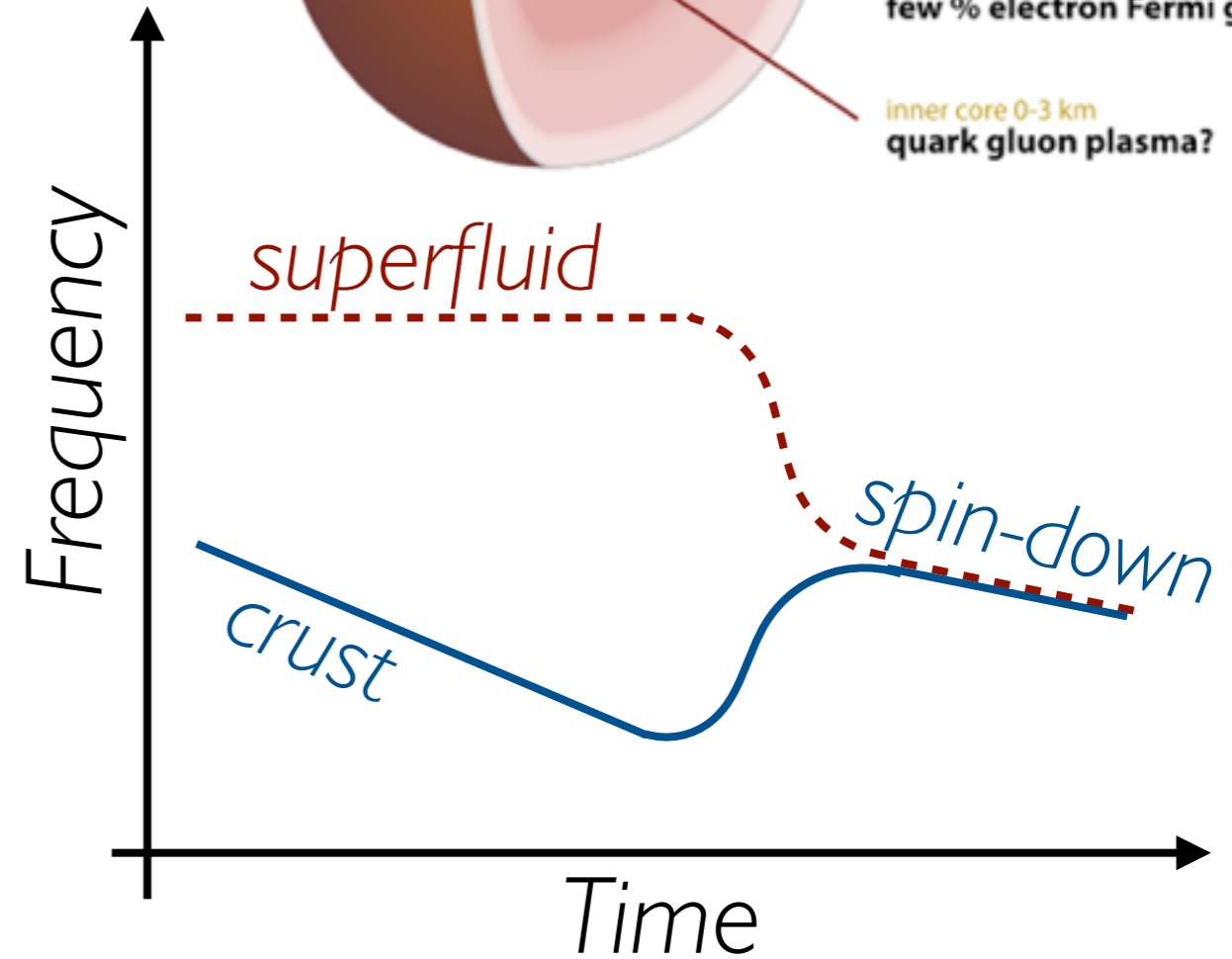
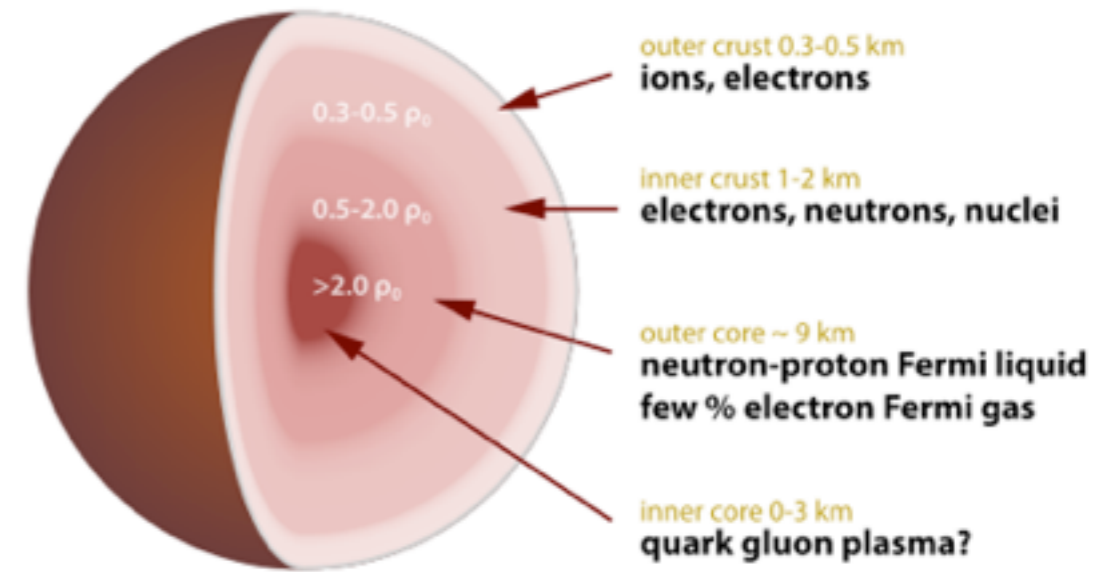


Name	Process	Emissivity ($\text{erg cm}^{-3} \text{s}^{-1}$)
Modified Urca (neutron branch)	$n + n \rightarrow n + p + e^- + \bar{\nu}_e$	$\sim 2 \times 10^{21} R T_9^8$
	$n + p + e^- \rightarrow n + n + \nu_e$	
Modified Urca (proton branch)	$p + n \rightarrow p + p + e^- + \bar{\nu}_e$	$\sim 10^{21} R T_9^8$
	$p + p + e^- \rightarrow p + n + \nu_e$	
Bremsstrahlungs	$n + n \rightarrow n + n + \nu + \bar{\nu}$	$\sim 10^{19} R T_9^8$
	$n + p \rightarrow n + p + \nu + \bar{\nu}$	
	$p + p \rightarrow p + p + \nu + \bar{\nu}$	
Cooper pair	$n + n \rightarrow [nn] + \nu + \bar{\nu}$	$\sim 5 \times 10^{21} R T_9^7$
	$p + p \rightarrow [pp] + \nu + \bar{\nu}$	$\sim 5 \times 10^{19} R T_9^7$
Direct Urca (nucleons)	$n \rightarrow p + e^- + \bar{\nu}_e$	$\sim 10^{27} R T_9^6$
	$p + e^- \rightarrow n + \nu_e$	

Pulsar glitches

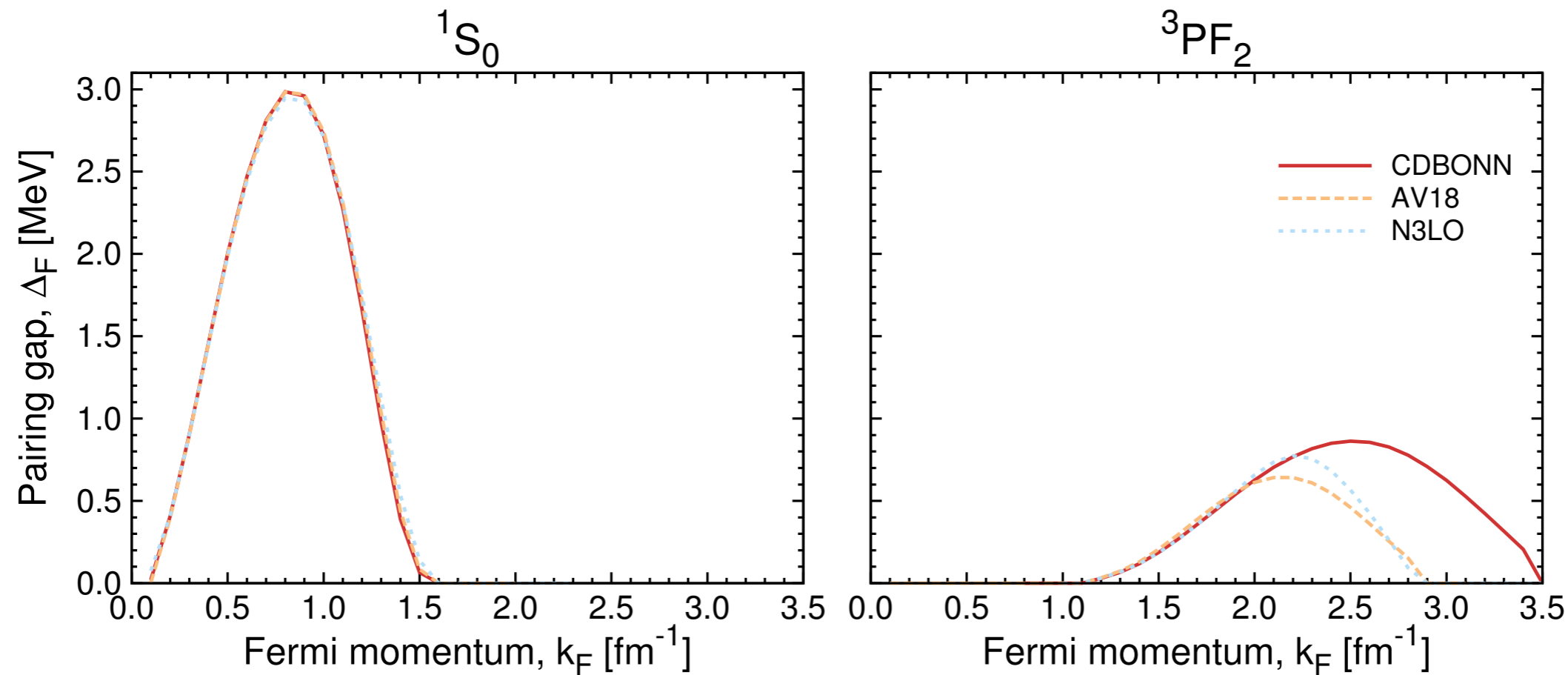


Espinoza, Lyne, Stappers & Kramer
MNRAS **414** 1679 (2011)



- **Crystalline** crust + dripped neutron **superfluid**
- Crust **slows down** due to magnetic braking
- **Superfluid** can only spin if **vortices** move out
- If vortices are **pinned** to nuclear lattice, they experience a time lag
- At some critical **pile-up**, a lot of vortices are **released** and crust spins up

Neutron matter BCS gaps



Dean & Hjorth-Jensen, *Rev. Mod. Phys.* **75** 607 (2003)

BCS equation

$$\Delta_k^L = - \sum_{L'} \int_{k'} \frac{\langle k | V_{nn}^{LL'} | k' \rangle}{2\chi_{k'}} \Delta_{k'}^{L'} + \chi_k = \sqrt{\varepsilon_k^2 + |\Delta_k|^2}$$

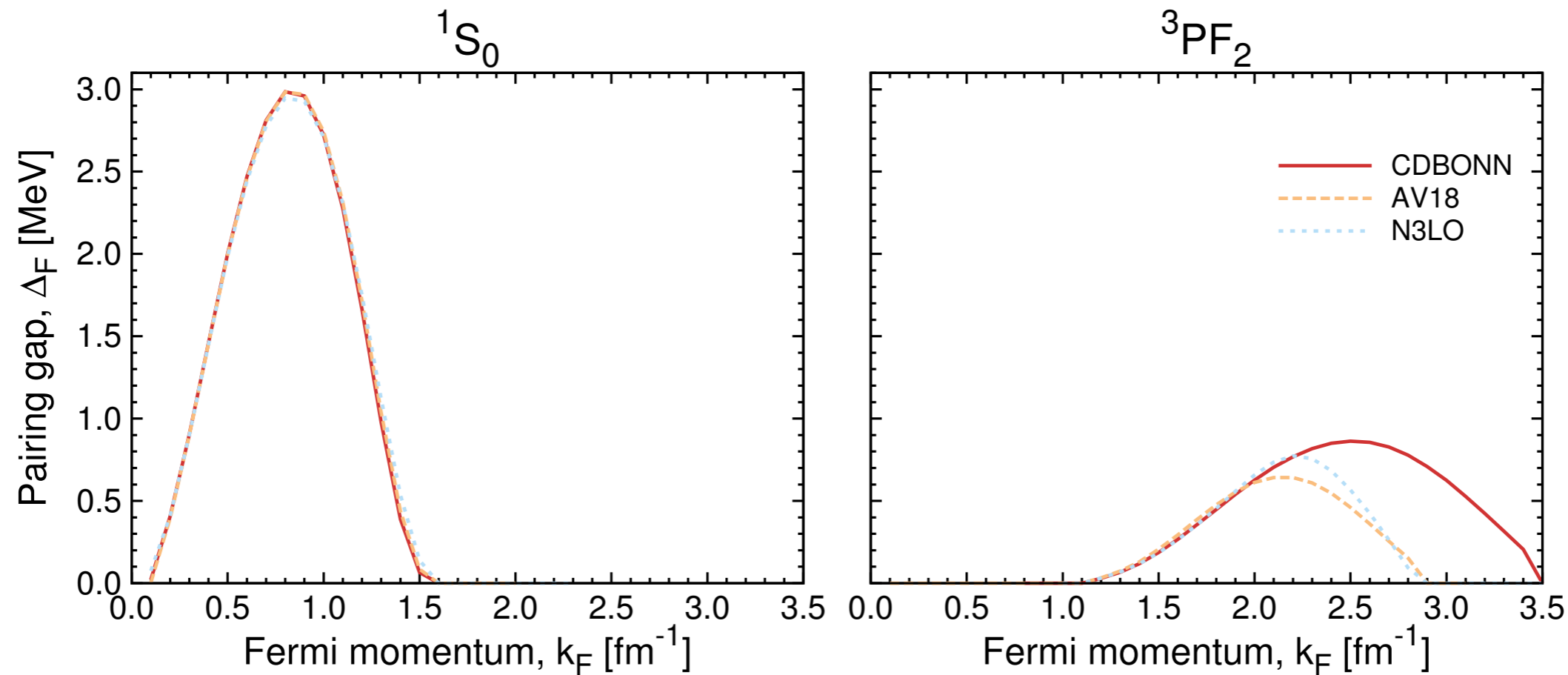
- Single-particle spectrum choice:

$$\varepsilon_k = \frac{k^2}{2m} + U(k) - \mu$$

- Angular gap dependence:

$$|\Delta_k|^2 = \sum_L |\Delta_k^L|^2$$

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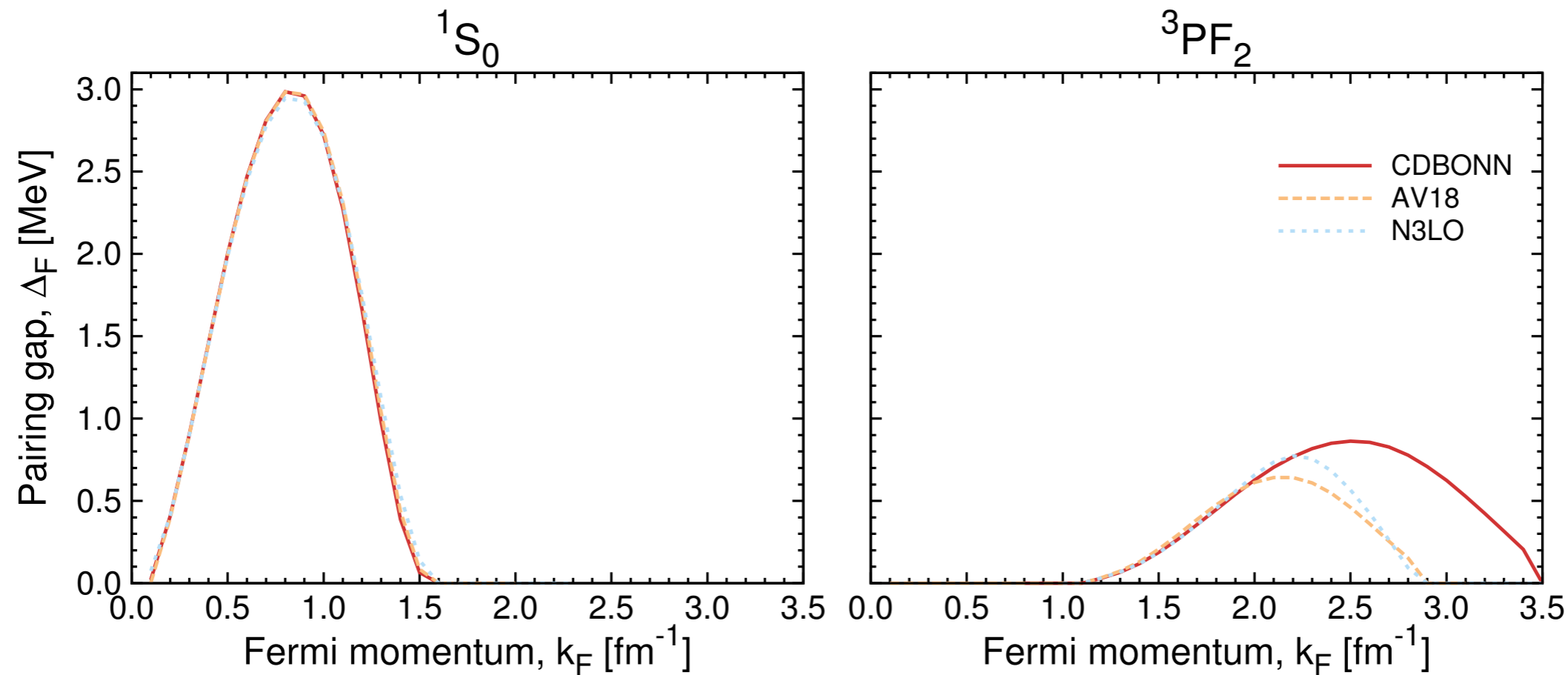
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Neutron matter BCS gaps



Dean & Hjorth-Jensen, *Rev. Mod. Phys.* **75** 607 (2003)

BCS equation

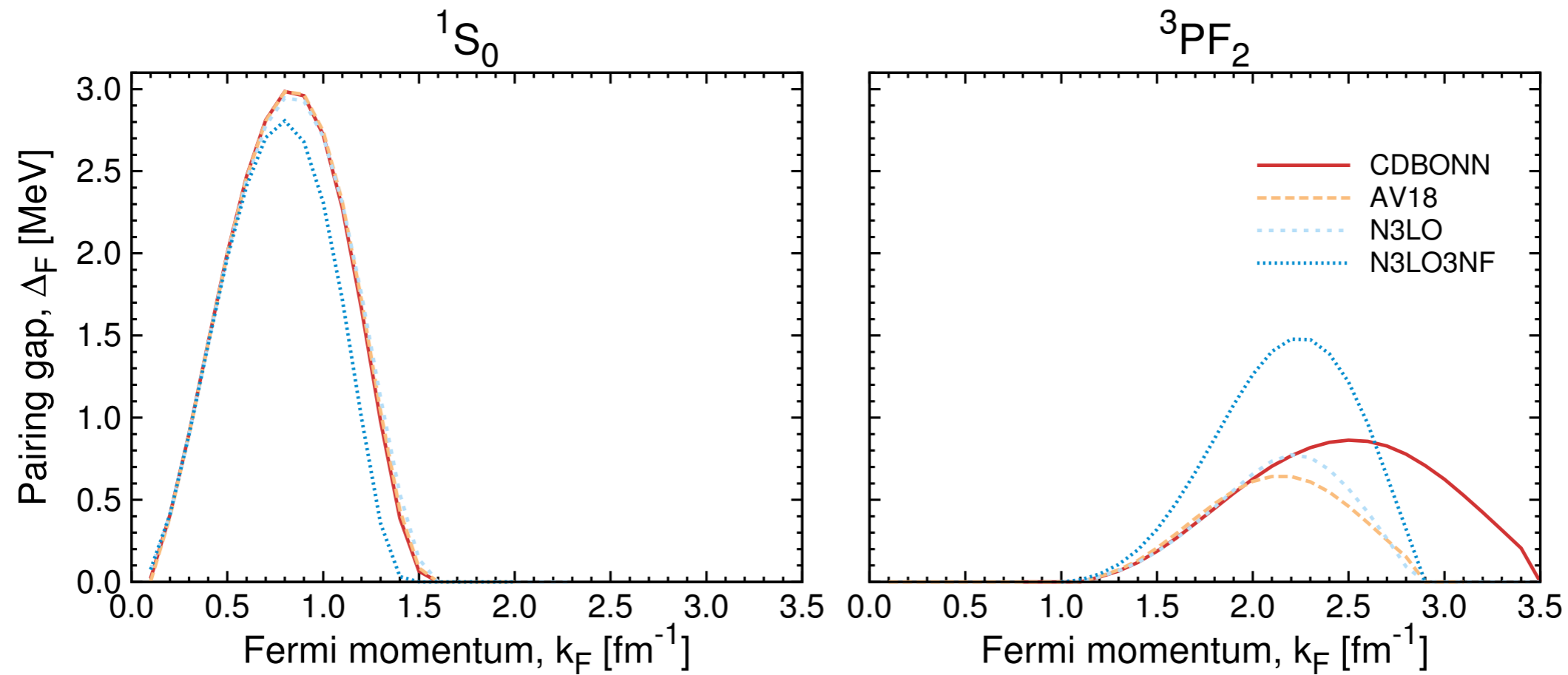
$$\Delta_k^L = - \sum_{L'} \int_{k'} \frac{\langle k | V_{nn}^{LL'} | k' \rangle}{2\chi_{k'}} \Delta_{k'}^{L'} + \chi_k = \sqrt{\varepsilon_k^2 + |\Delta_k|^2}$$

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- Angular gap dependence:

$$\varepsilon_k = \frac{k^2}{2m} + \cancel{U(k)} - \mu$$

$$|\Delta_k|^2 = \sum_L |\Delta_k^L|^2 \approx |\Delta_k^L|^2$$

Neutron matter BCS gaps



Dean & Hjorth-Jensen, *Rev. Mod. Phys.* **75** 607 (2003)

BCS equation

$$\Delta_k^L = - \sum_{L'} \int_{k'} \frac{\langle k | V_{nn}^{LL'} | k' \rangle}{2\chi_{k'}} \Delta_{k'}^{L'} + \chi_k = \sqrt{\varepsilon_k^2 + |\Delta_k|^2}$$

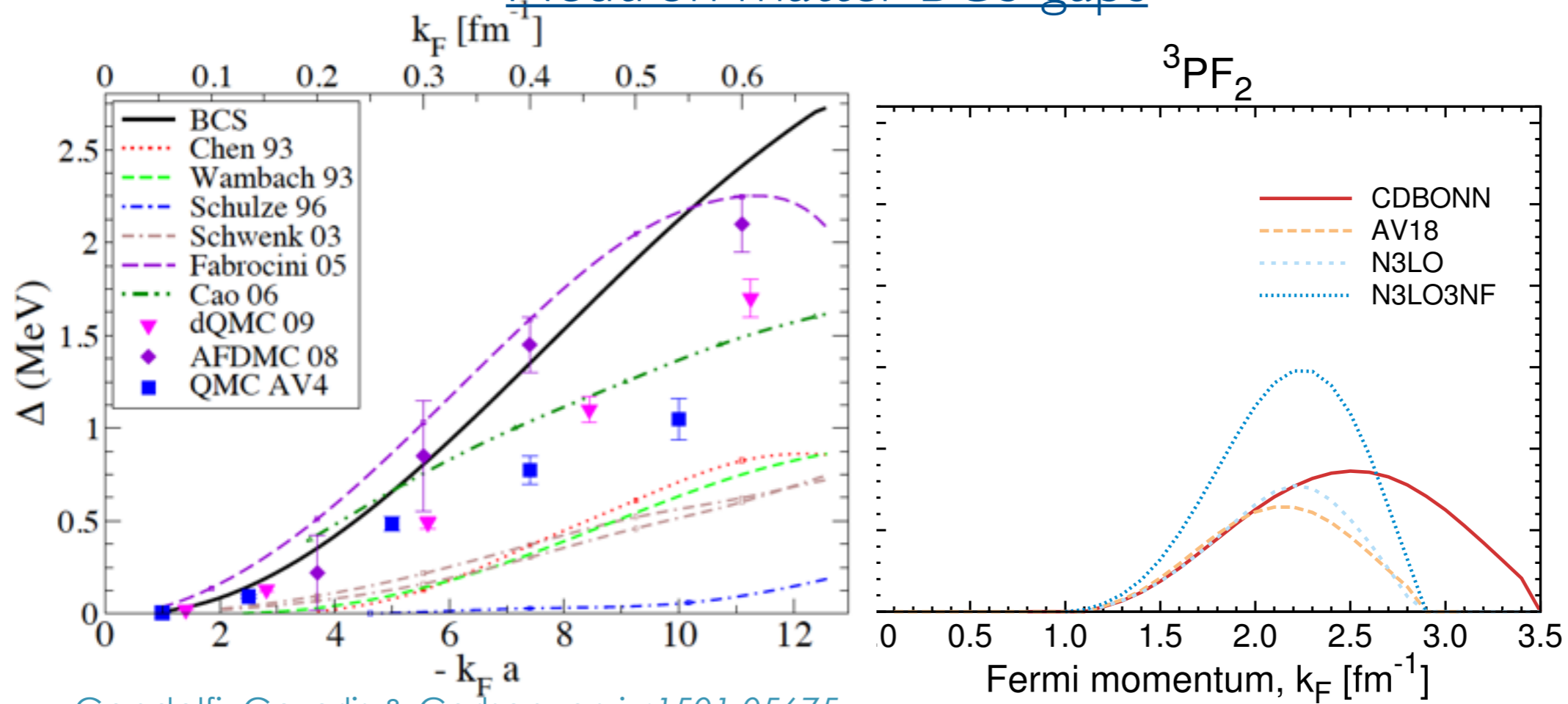
- Single-particle spectrum choice:

$$\varepsilon_k = \frac{k^2}{2m} + \cancel{U(k)} - \mu$$

- Angular gap dependence:

$$|\Delta_k|^2 = \sum_L |\Delta_k^L|^2 \approx |\Delta_k^L|^2$$

Neutron matter BCS gaps



Gandolfi, Gezerlis & Carlson, arxiv:1501.05675,

BCS equation Annu. Rev. Nucl. Part.Sci.

Dean & Hjorth-Jensen, Rev. Mod. Phys. **75** 607 (2003)

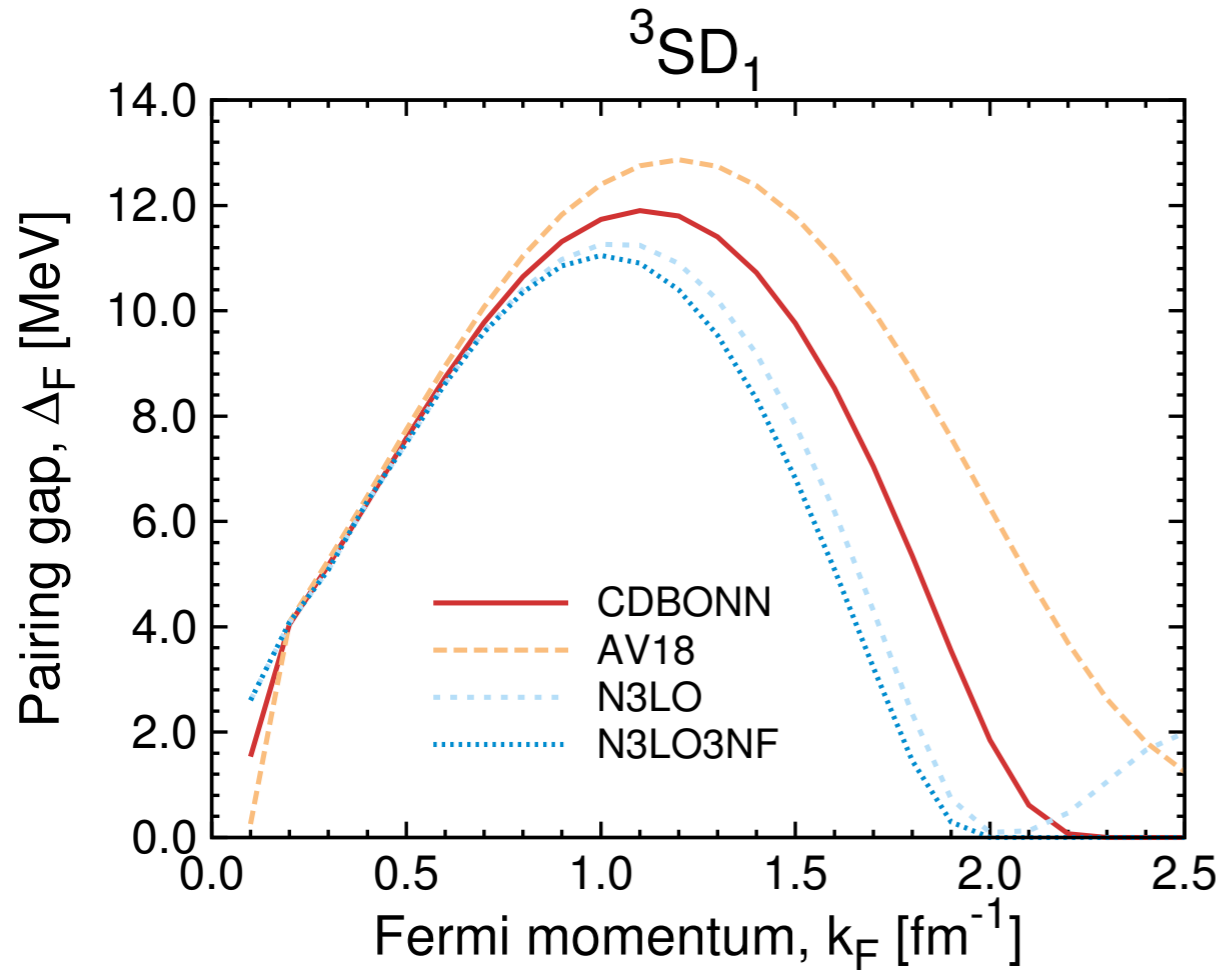
$$\Delta_k^L = - \sum_{L'} \int_{k'} \frac{\langle k | V_{nn}^{LL'} | k' \rangle}{2\chi_{k'}} \Delta_{k'}^{L'} + \chi_k = \sqrt{\varepsilon_k^2 + |\Delta_k|^2}$$

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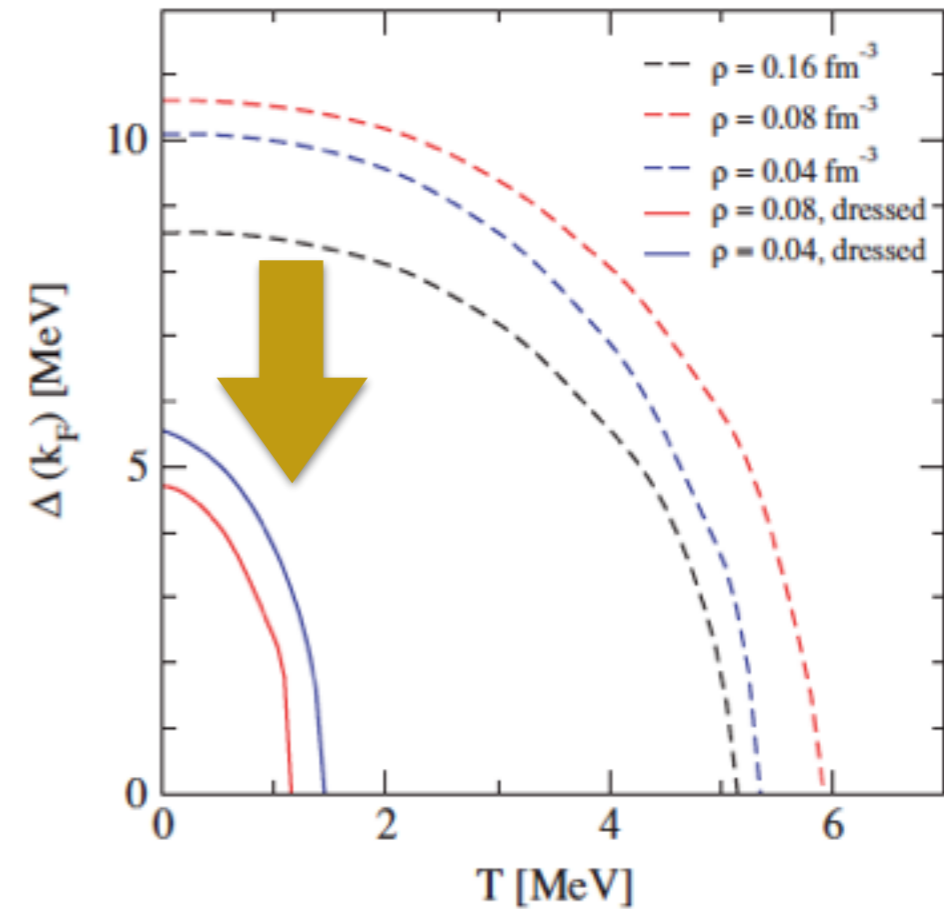
$$|\Delta_k|^2 = \sum_L |\Delta_k^L|^2 \approx |\Delta_k^L|^2$$

3SD_1 nuclear matter BCS gaps



Maurizio, Holt & Finelli, *PRC* **90**, 044003 (2014)

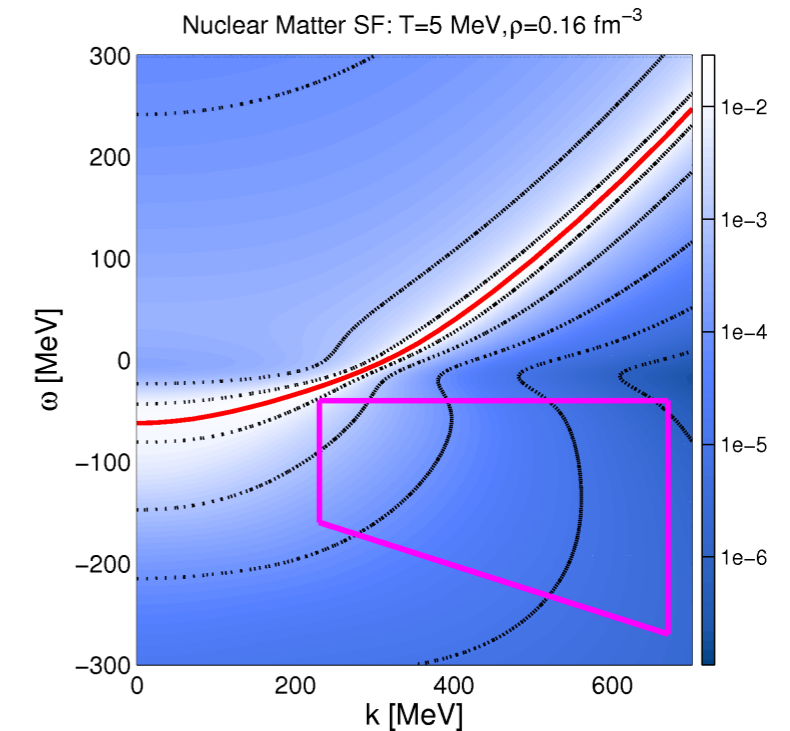
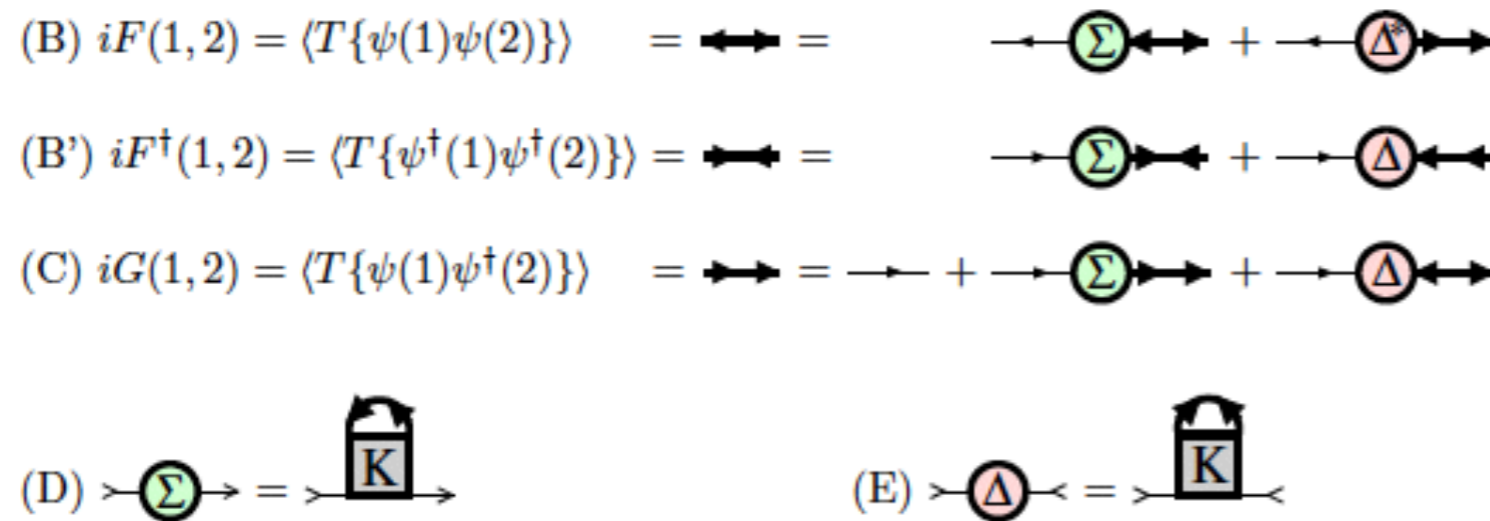
SRC-depleted 3SD_1 gaps



Muether & Dickhoff, *PRC* **72** 054313 (2005)

- **Massive** gaps 3SD_1 channel but...
- **No evidence** of strong np nuclear pairing
- Short-range correlations **deplete** gap
- 3BF effect? Short-range effects?

Beyond BCS 101: SRC

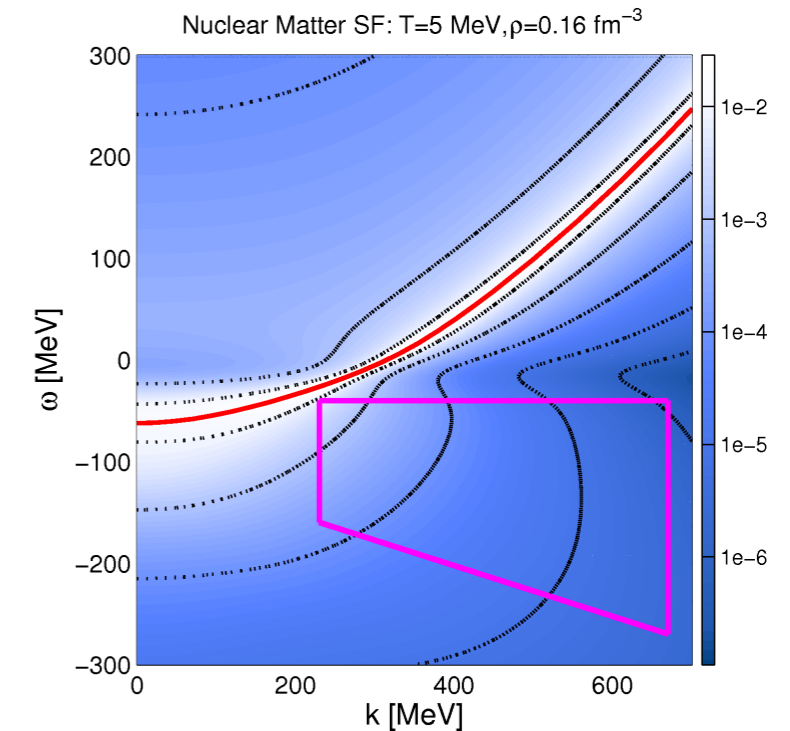
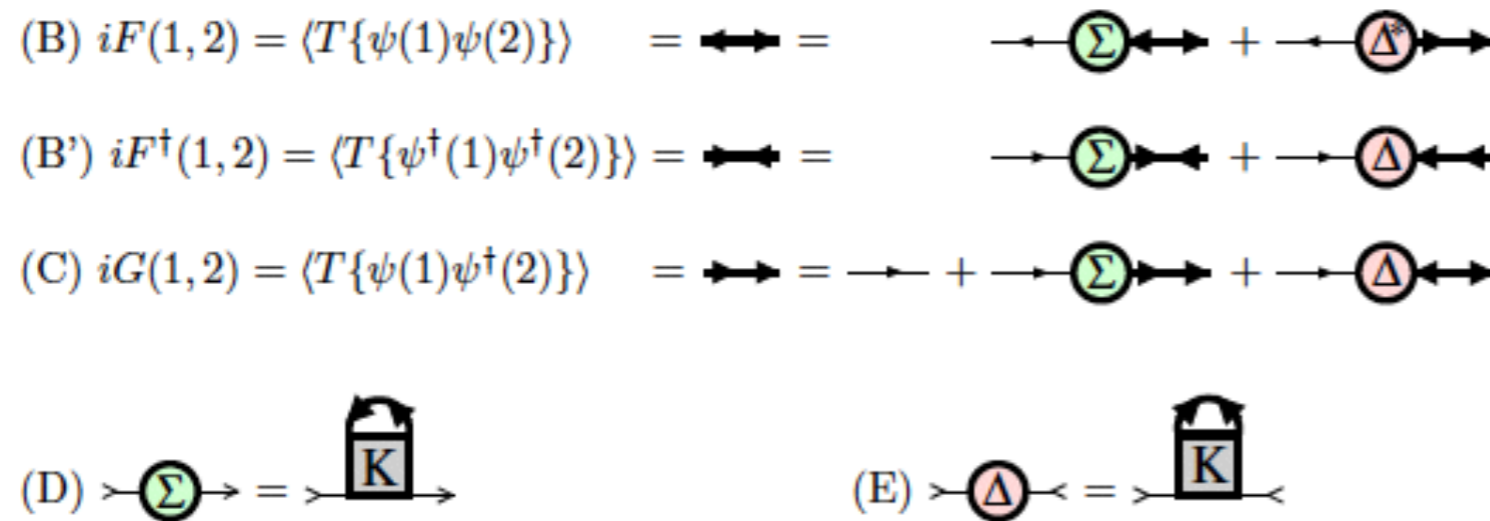


BCS+SRC equation

$$\Delta_k^L = - \sum_{L'} \int_{k'} \frac{\langle k | V_{nn}^{LL'} | k' \rangle}{2\chi_{k'}} \Delta_{k'}^{L'} + \frac{1}{2\chi_k} = \int_{\omega} \int_{\omega'} \frac{1 - f(\omega) - f(\omega')}{\omega + \omega'} A(k, \omega) A_s(k, \omega')$$

- **BCS** is lowest order in Gorkov Green's function expansion
- How does **removal** of **strength** affect pairing?
- Effective replacement of denominator

Beyond BCS 101: SRC

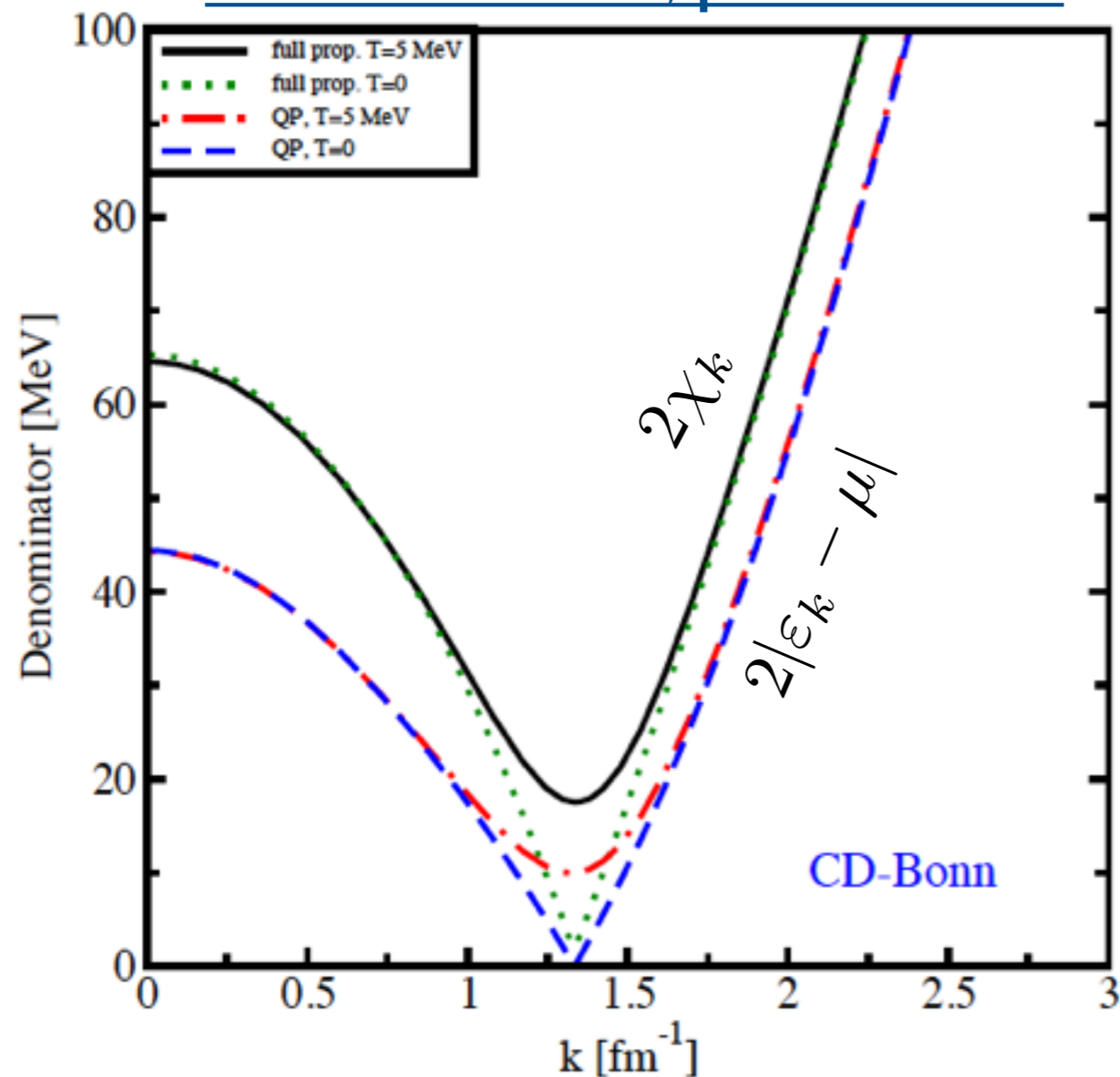


BCS+SRC equation

$$\Delta_k^L = - \sum_{L'} \int_{k'} \frac{\langle k | V_{nn}^{LL'} | k' \rangle}{2\chi_{k'}} \Delta_{k'}^{L'} + \frac{1}{2\chi_k} = \int_{\omega} \int_{\omega'} \frac{1 - f(\omega) - f(\omega')}{\omega + \omega'} A(k, \omega) A_{\text{S}}(k, \omega')$$

- **BCS** is lowest order in Gorkov Green's function expansion
- How does **removal** of **strength** affect pairing?
- Effective replacement of denominator

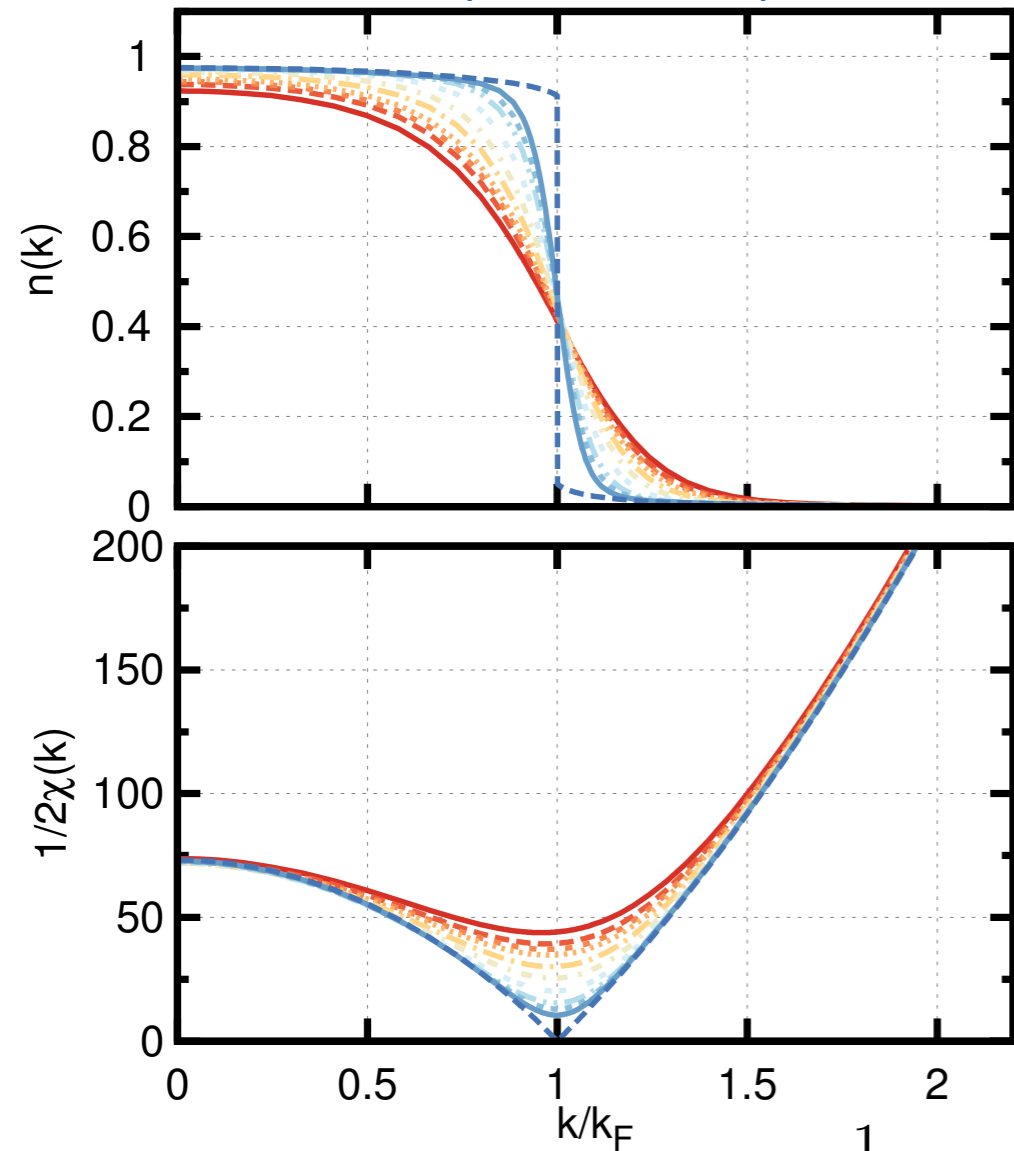
Nuclear matter, $\rho = 0.16 \text{ fm}^{-3}$



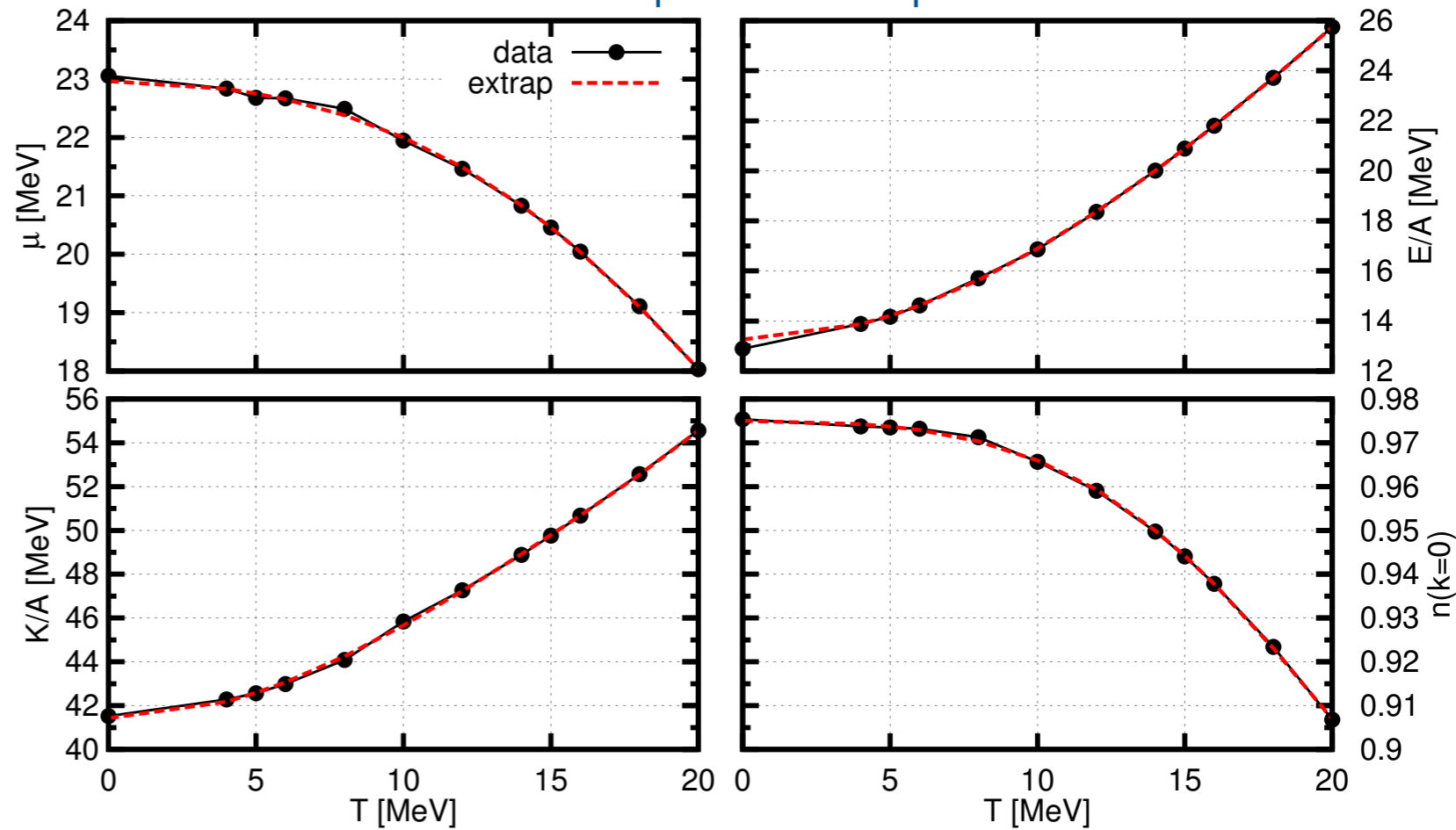
$$\frac{1}{2\chi_k} = \int_{\omega} \int_{\omega'} \frac{1 - f(\omega) - f(\omega')}{\omega + \omega'} A(k, \omega) A(k, \omega')$$

- **BCS** is lowest order in Gorkov Green's function expansion
- How does **removal** of **strength** affect pairing?
- Effective replacement of denominator

Microscopic extrapolation

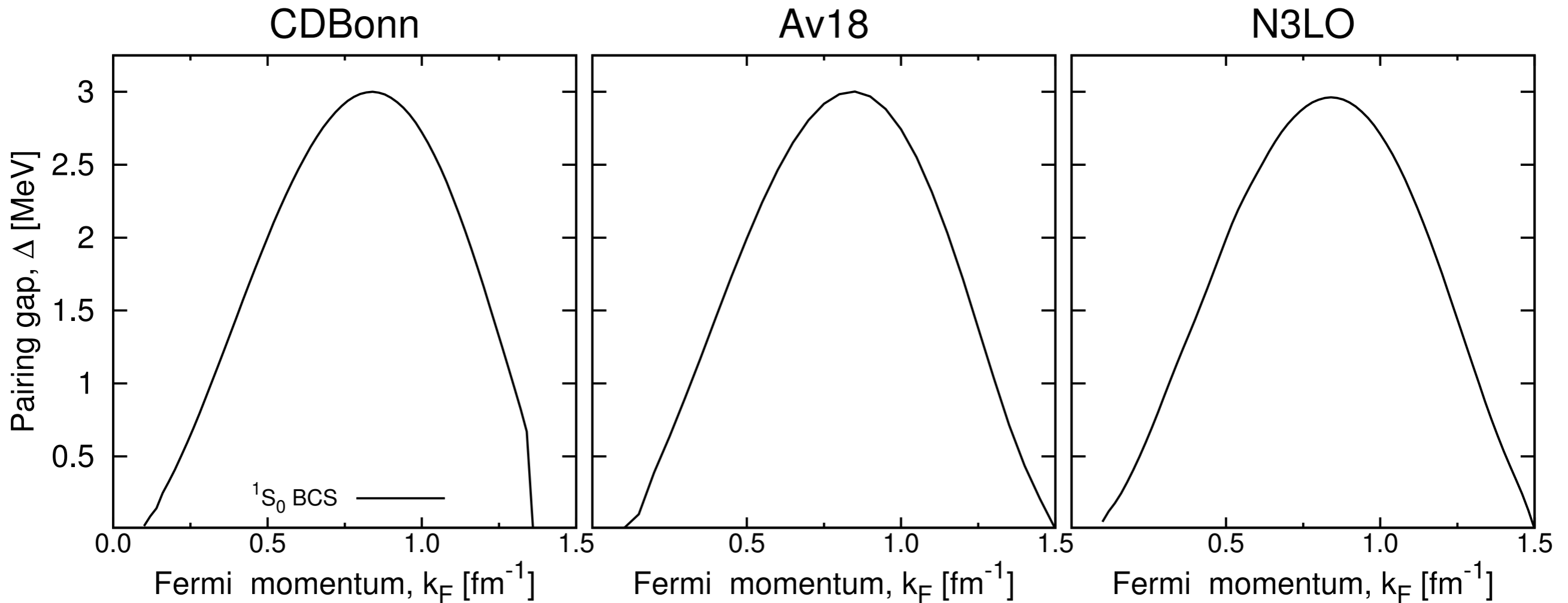


Macroscopic extrapolation

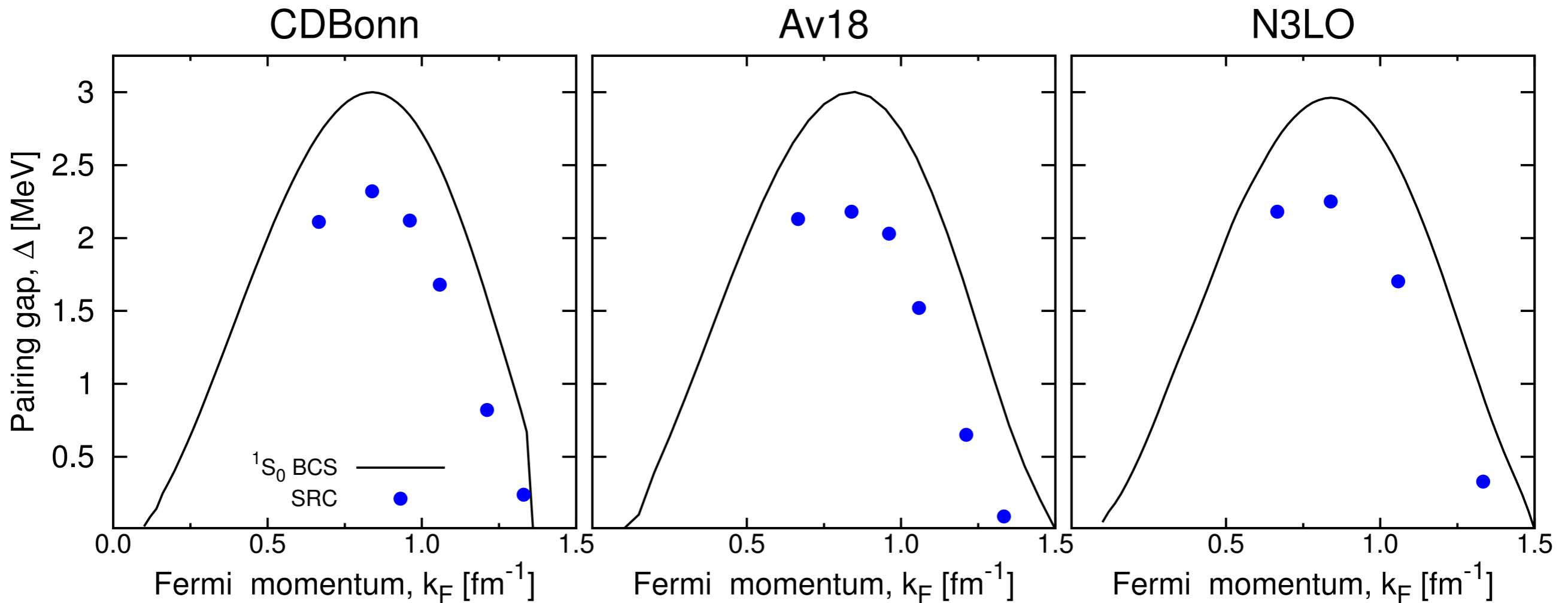


$$\frac{1}{2\chi_k} = \int_{\omega} \int_{\omega'} \frac{1 - f(\omega) - f(\omega')}{\omega + \omega'} A(k, \omega) A(k, \omega')$$

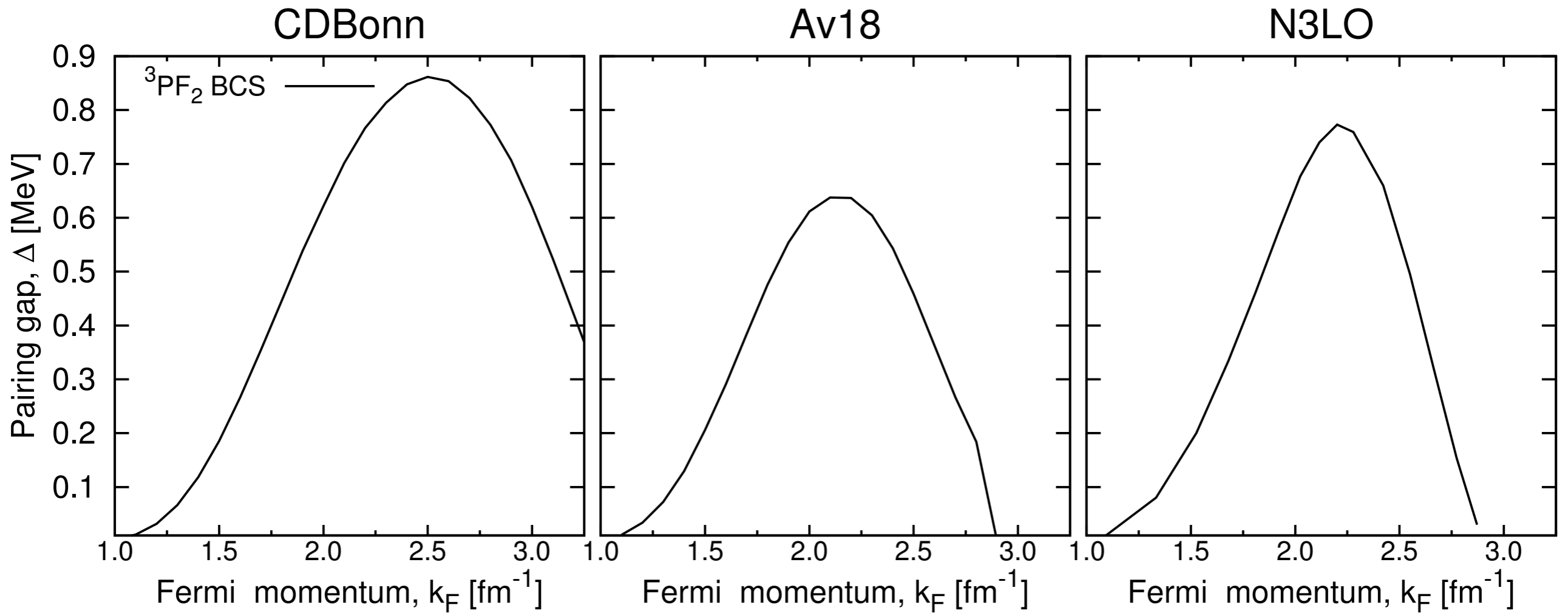
- Similar to G_2 component at finite T ... but need $T=0$ **data!**
- **Extrapolate** $\Sigma(k, \omega; T)$ (13 Gb worth data)
- **Constraining** with **macroscopic** properties



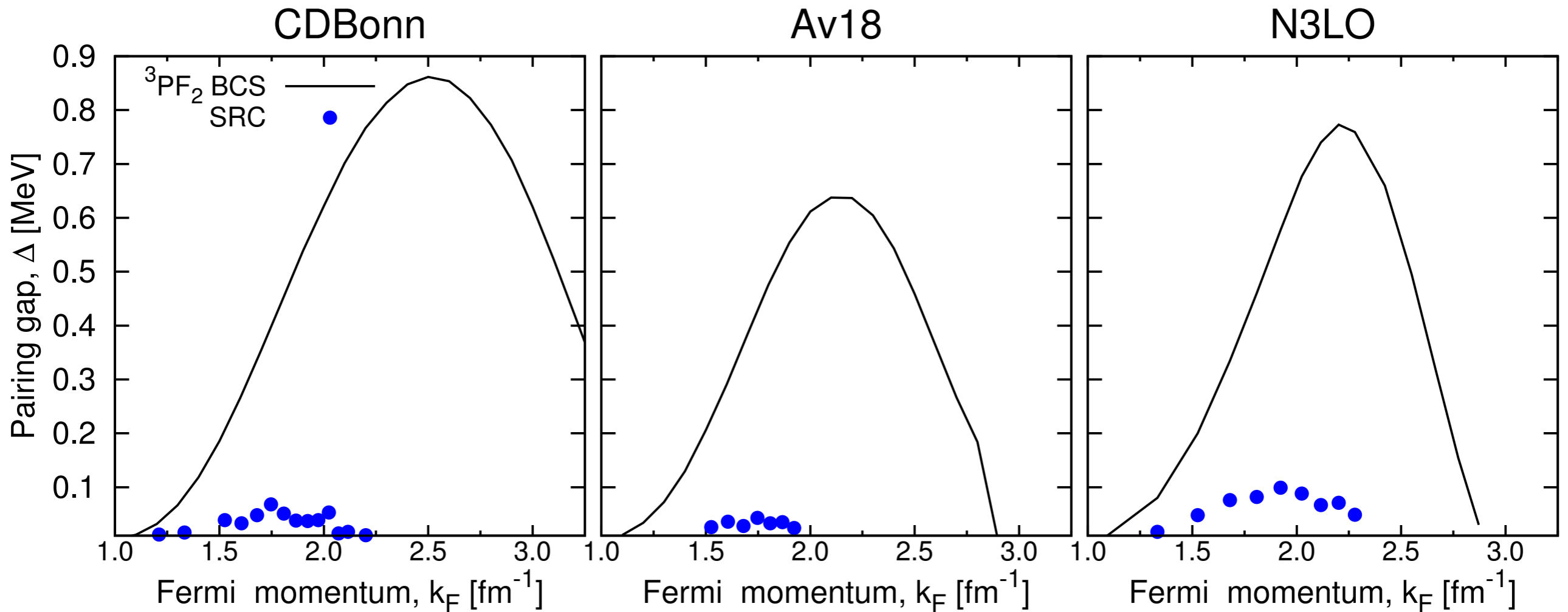
- **Moderate** screening due to SRC: $\Delta_{\max} \approx 2.5$ MeV
- Effect is **robust**: independent of NN potential
- **Similar** gap closure



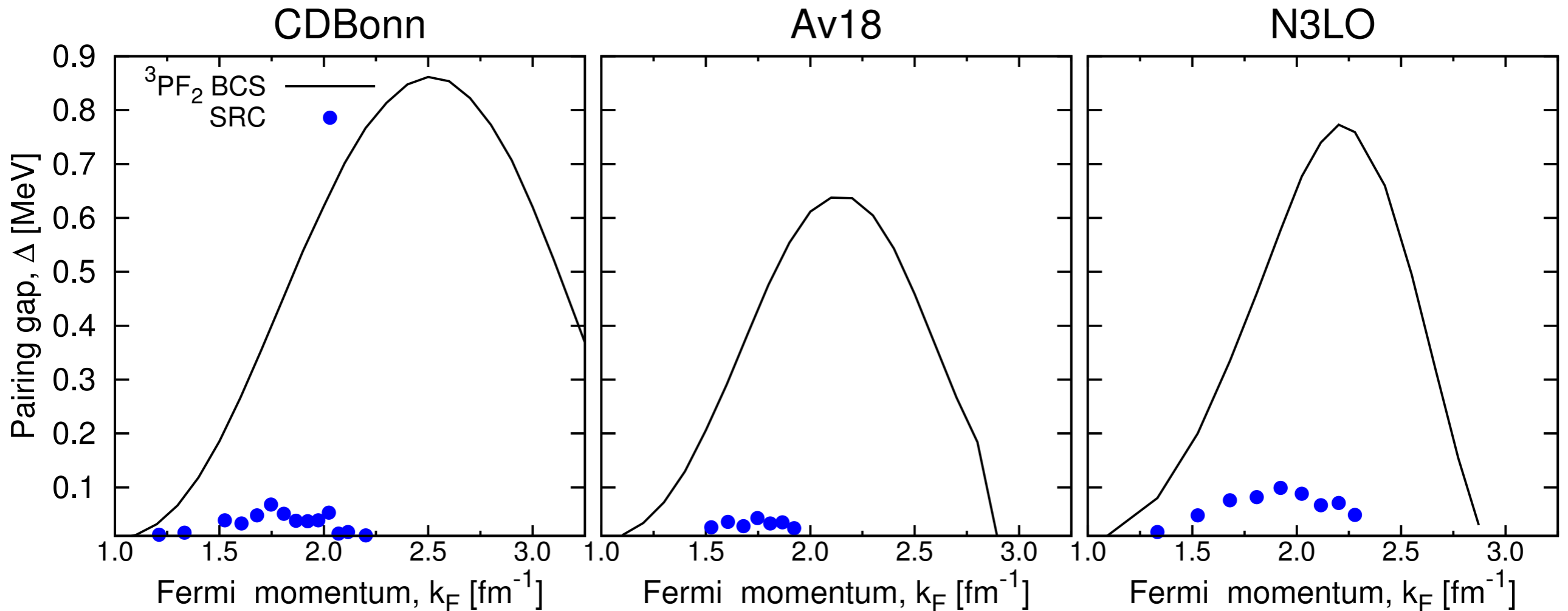
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- **Large** screening due to **SRC**: $\Delta \approx 0.1-0.05$ MeV
- **Lower** gap **closure**



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$$\Delta_k^L = - \sum_{L'} \int_{k'} \frac{\langle k | V_{nn}^{LL'} | k' \rangle}{2\chi_{k'}} \Delta_{k'}^{L'}$$

- Bare NN potential only is not the only possible interaction

$$\Delta_k^L = - \sum_{L'} \int_{k'} \frac{\langle k | V_{nn}^{LL'} | k' \rangle}{2\chi_{k'}} \Delta_{k'}^{L'}$$

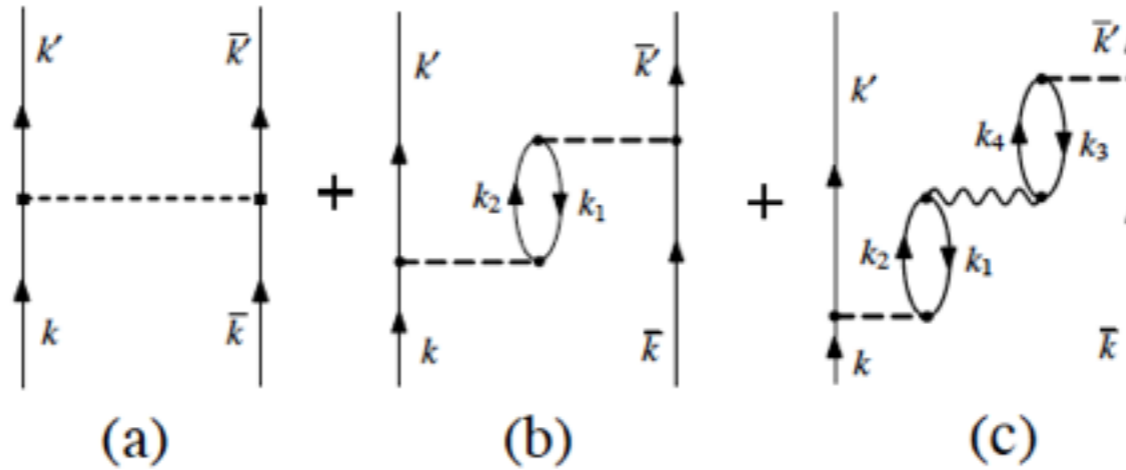
The equation features a green box around the numerator $\langle k | V_{nn}^{LL'} | k' \rangle$ with a green question mark above it, and a red box around the denominator $2\chi_{k'}$ with a red checkmark to its right.

- Bare NN potential only is not the only possible interaction

$$\Delta_k^L = - \sum_{L'} \int_{k'} \frac{\langle k | V_{nn}^{LL'} | k' \rangle}{2\chi_{k'}} \Delta_{k'}^{L'}$$

$\langle k | V_{nn}^{LL'} | k' \rangle$?
 $2\chi_{k'}$ ✓

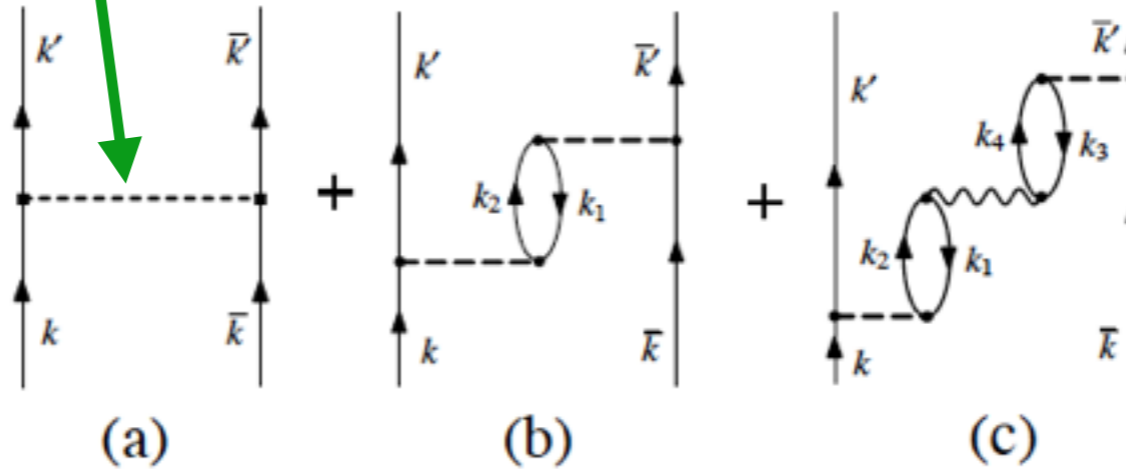
$\mathcal{V}_{\text{pair}} =$



- Bare NN potential only is not the only possible interaction

$$\Delta_k^L = - \sum_{L'} \int_{k'} \frac{\langle k | V_{nn}^{LL'} | k' \rangle}{2\chi_{k'}} \Delta_{k'}^{L'}$$

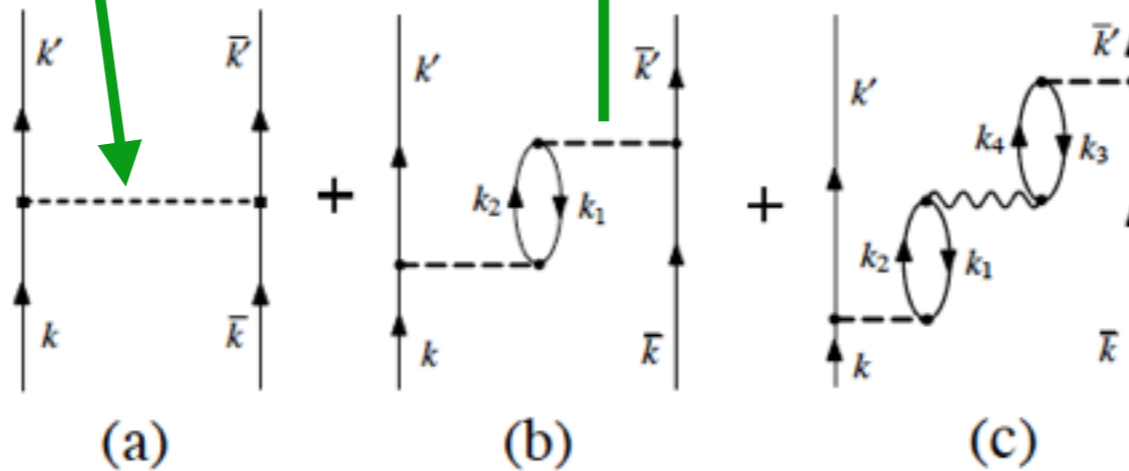
$\mathcal{V}_{\text{pair}} =$



- Bare NN potential only is not the only possible interaction
- Diagram (a): nuclear interaction

$$\Delta_k^L = - \sum_{L'} \int_{k'} \frac{\langle k | V_{nn}^{LL'} | k' \rangle}{2\chi_{k'}} \Delta_{k'}^{L'}$$

$\mathcal{V}_{\text{pair}} =$



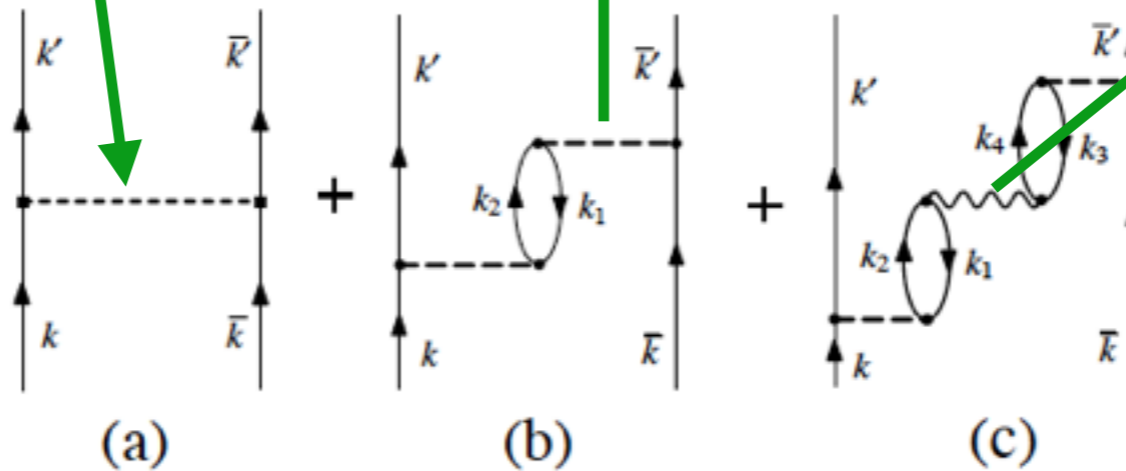
$$\langle 1\bar{1} | \mathcal{V} | 1\bar{1} \rangle = \frac{1}{4} \sum_{2,2'} \sum_{S,T} (-)^S (2S+1) \langle 12 | G_{ST}^{\text{ph}} | 1'2' \rangle_A \langle 2'\bar{1} | G_{ST}^{\text{ph}} | 2\bar{1}' \rangle_A \Lambda^0(22')$$

- Bare NN potential only is not the only possible interaction
- Diagram (a): nuclear interaction
- Diagram (b): in-medium interaction, density and spin fluctuations

$$\Delta_k^L = - \sum_{L'} \int_{k'} \langle k | V_{nn}^{LL'} | k' \rangle \Delta_{k'}^{L'}$$

$\langle k | V_{nn}^{LL'} | k' \rangle$ $2\chi_{k'}$

$\mathcal{V}_{\text{pair}} =$



$$\langle 1\bar{1} | \mathcal{V} | 1\bar{1} \rangle = \frac{1}{4} \sum_{2,2'} \sum_{S,T} (-)^S (2S+1) \langle 12 | G_{ST}^{\text{ph}} | 1'2' \rangle_A \langle 2'\bar{1} | G_{ST}^{\text{ph}} | 2\bar{1}' \rangle_A \Lambda^0(22')$$

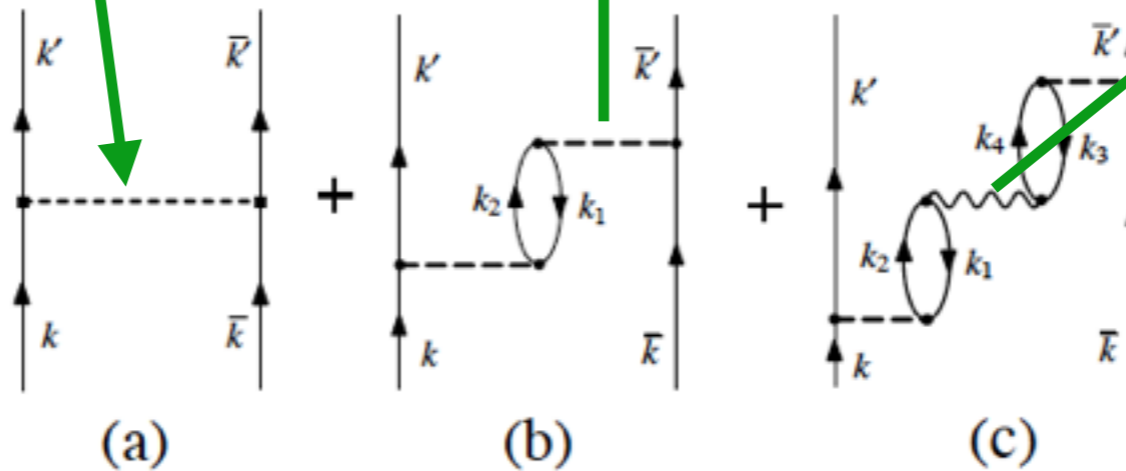
$$\Lambda_{ST}(q) = \frac{\Lambda_{ST}^0(q)}{1 - \Lambda_{ST}^0(q) \times F_{ST}}$$

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$$\Delta_k^L = - \sum_{L'} \int_{k'} \langle k | V_{nn}^{LL'} | k' \rangle \Delta_{k'}^{L'}$$

$\langle k | V_{nn}^{LL'} | k' \rangle$ $2\chi_{k'}$

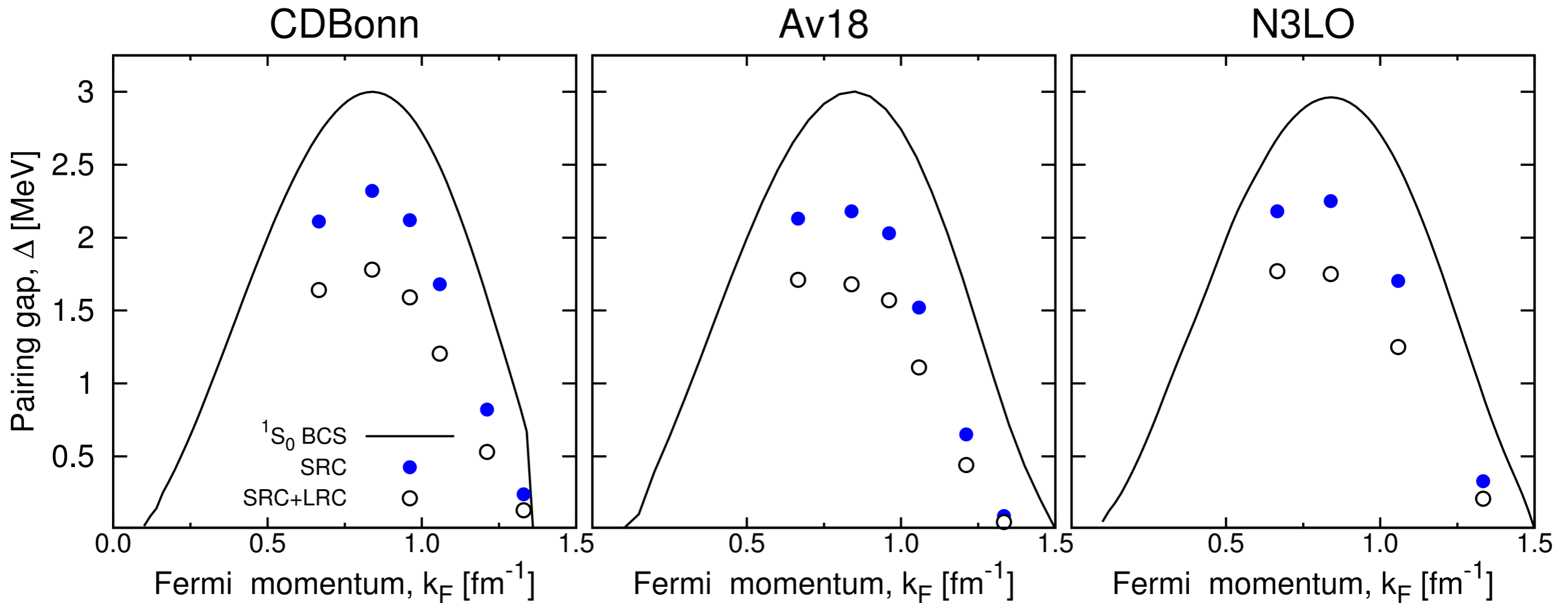
$\mathcal{V}_{\text{pair}} =$



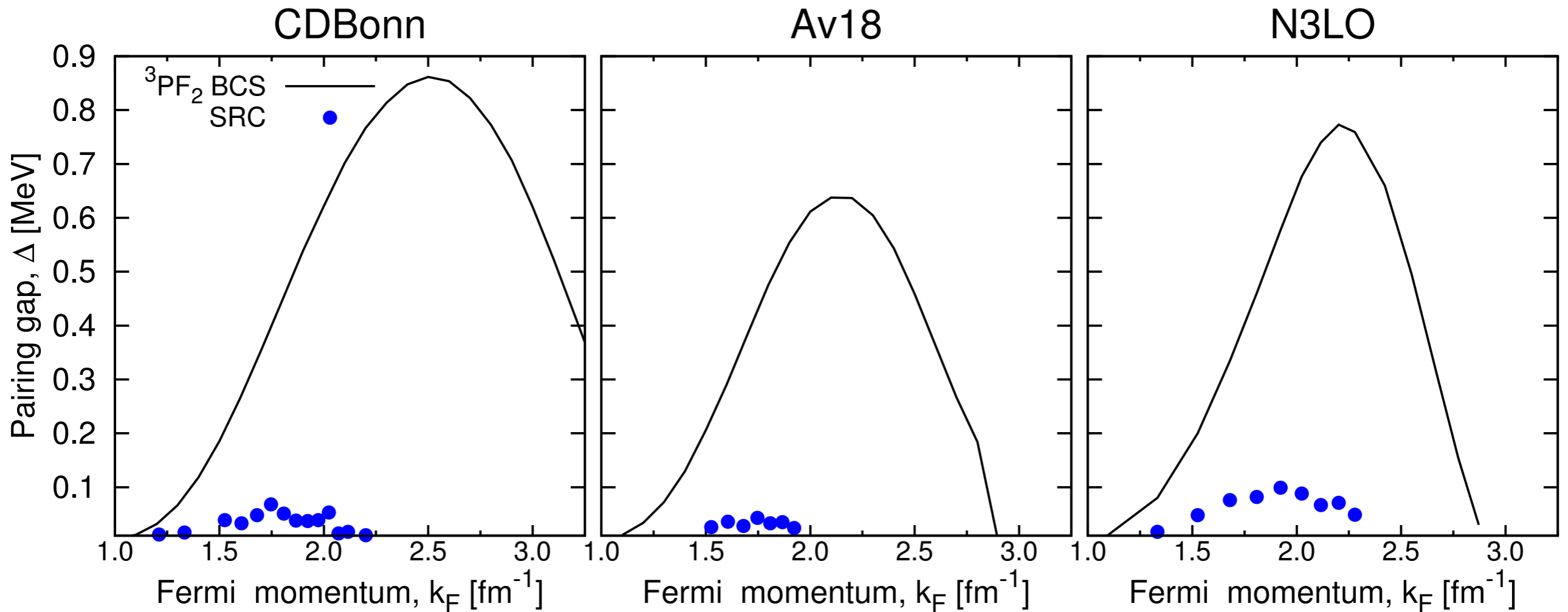
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$$\Lambda_{ST}(q) = \frac{\Lambda_{ST}^0(q)}{1 - \Lambda_{ST}^0(q) \times F_{ST}}$$

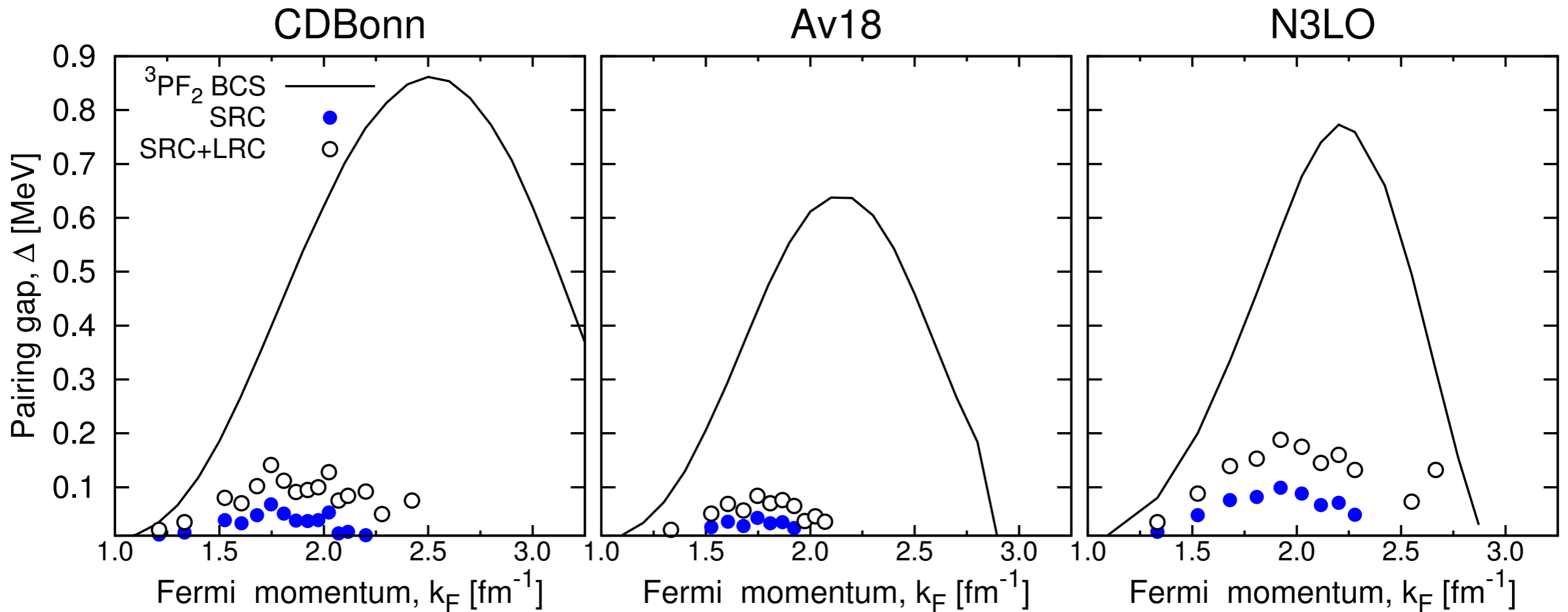
- Bare NN potential only is not the only possible interaction
- Diagram (a): nuclear interaction
- Diagram (b): in-medium interaction, density and spin fluctuations
- Diagram (c): included by Landau parameters



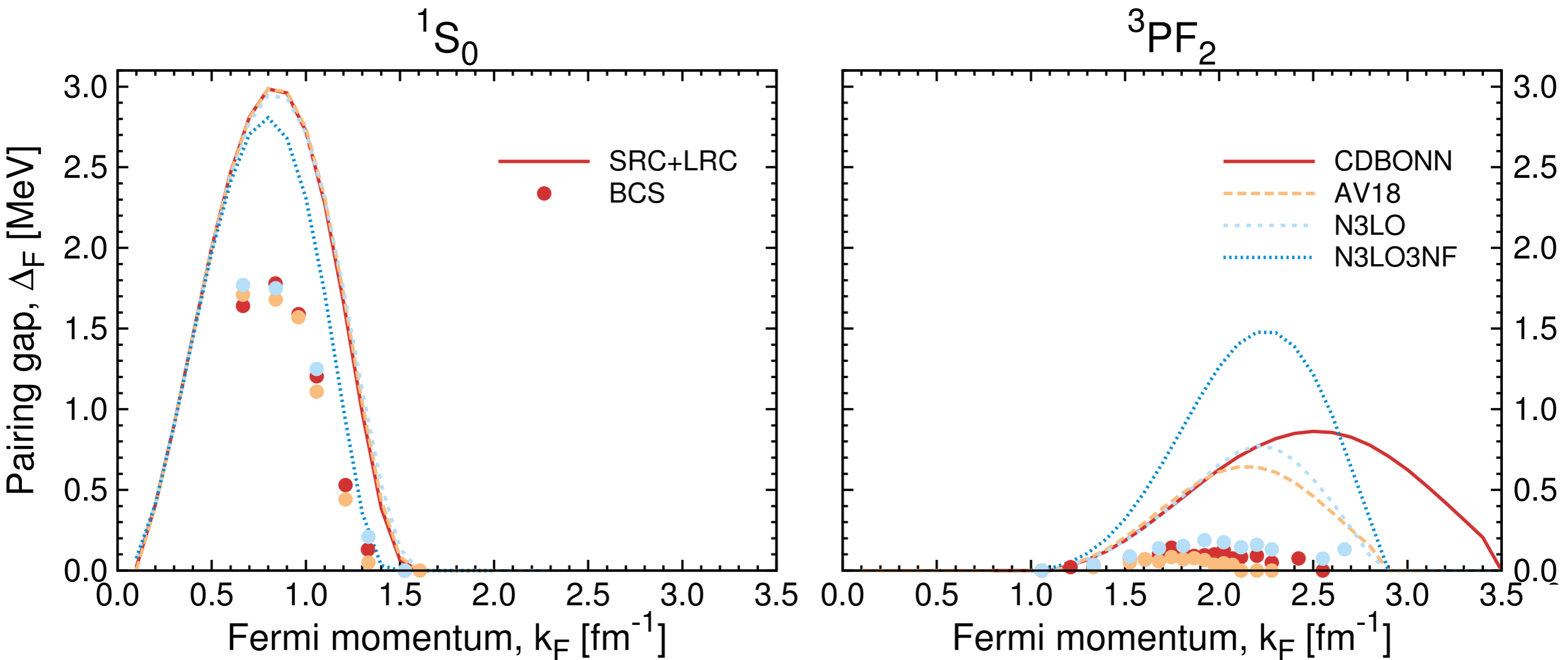
- **Moderate** screening due to SRC: $\Delta_{\max} \approx 2.5$ MeV
- **Additional** screening due to LRC: $\Delta_{\max} \approx 2$ MeV
- Effect is **robust**: independent of NN potential
- **Similar** gap closure



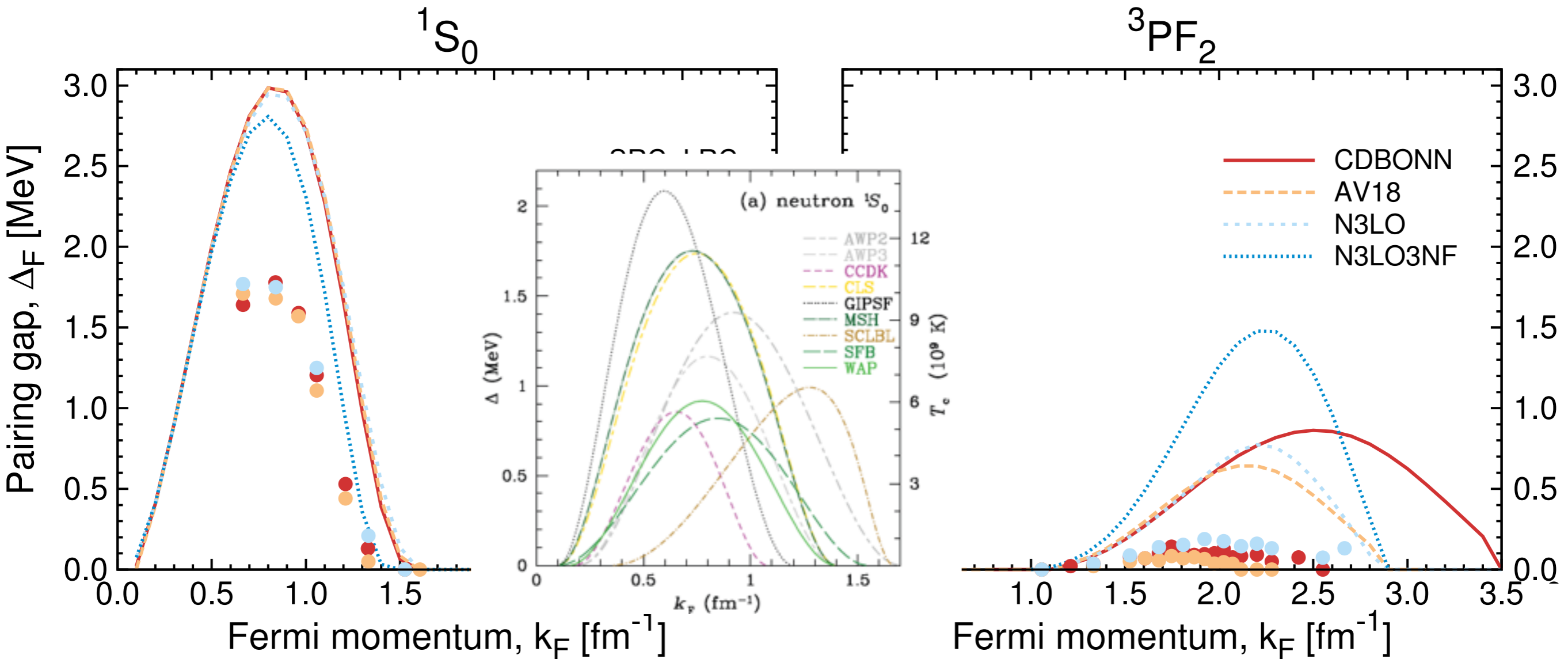
- **Anti**-screening induced by **SRC**: $\Delta \approx 0.15$ - 0.20 MeV
- **Lower** gap **closure**
- Effect is **robust**: independent of NN potential



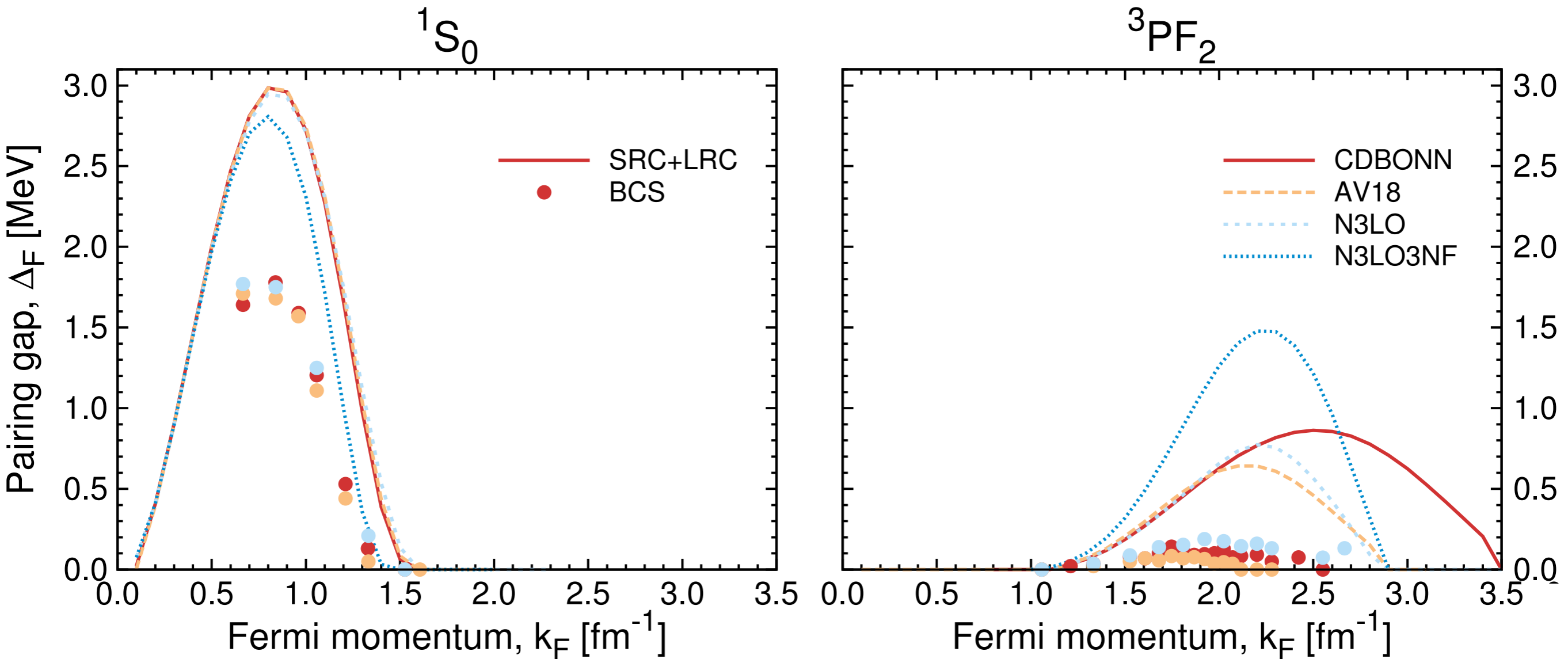
- **Anti**-screening induced by **SRC**: $\Delta \approx 0.15-0.20$ MeV
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- Effect is **robust**: independent of NN potential
- 3NF effect **not** included in SRC yet
- BCS indicates effects is **smaller** than correlations

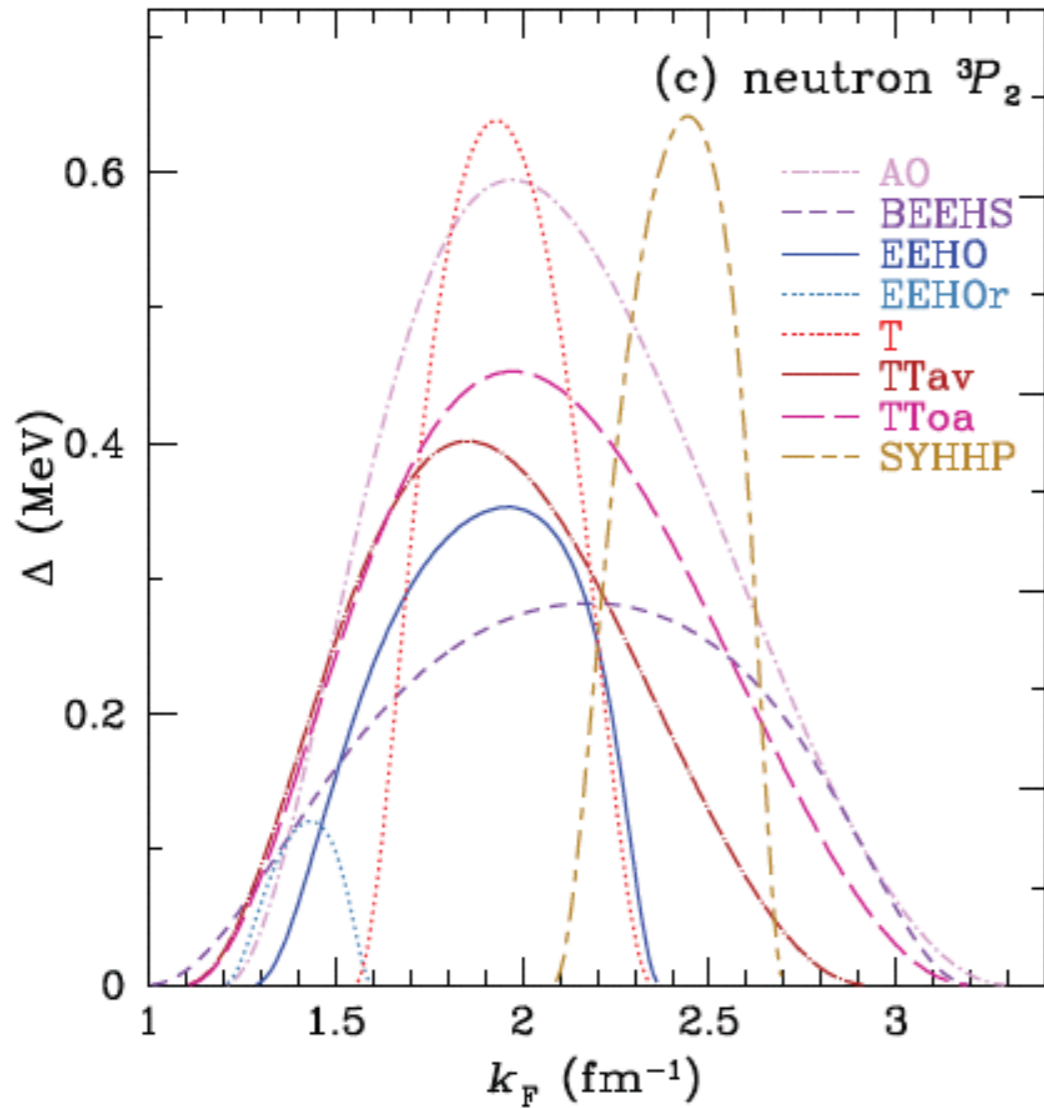


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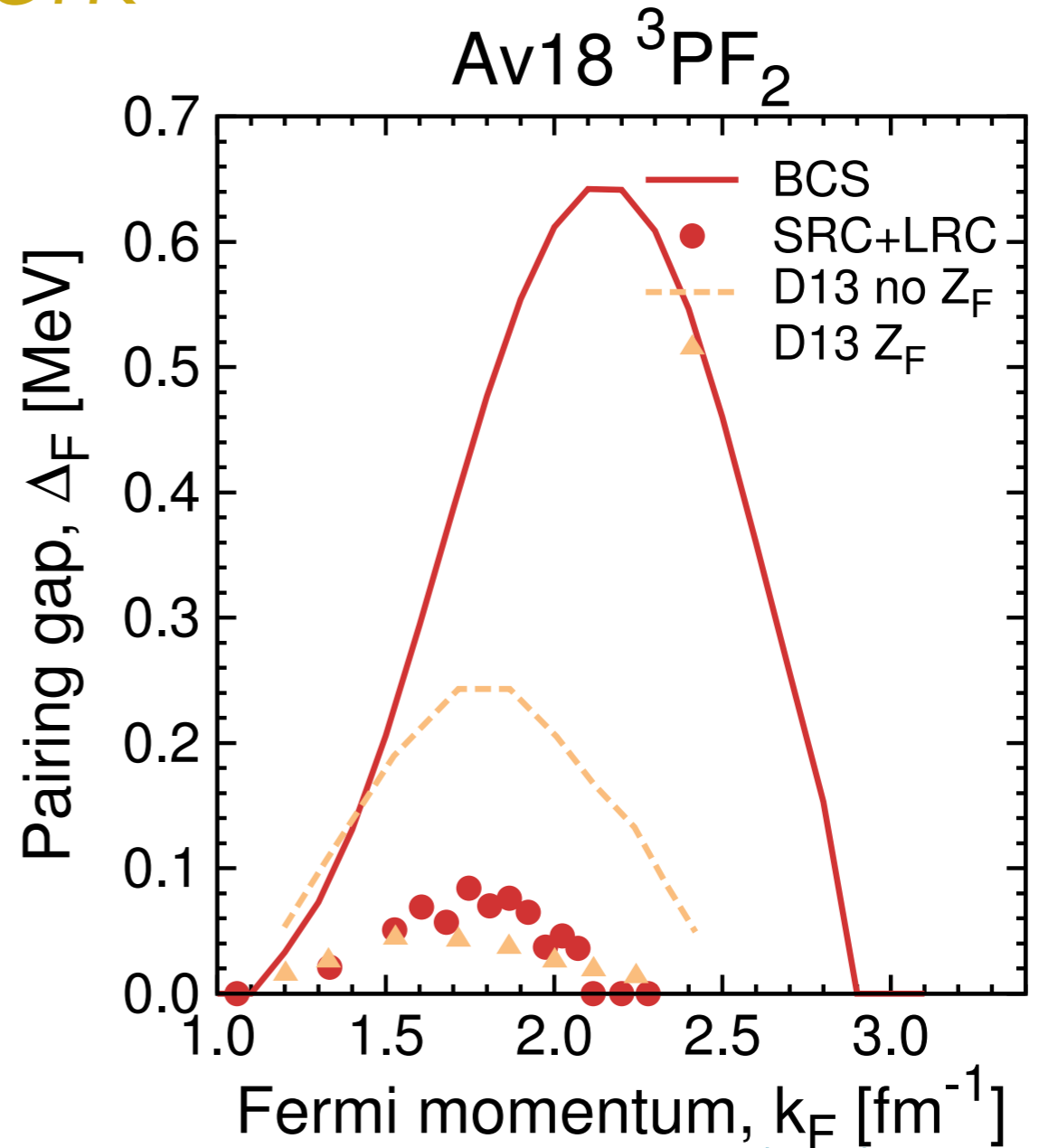
- Effect is **robust**: independent of NN potential
- 3NF effect **not** included in SRC yet
- BCS indicates effects is **smaller** than correlations

Comparison to other work



Ho, Elshamouty, Heinke, Potekhin,
PRC **91** 015806 (2015)

BCS+Z-factor equation



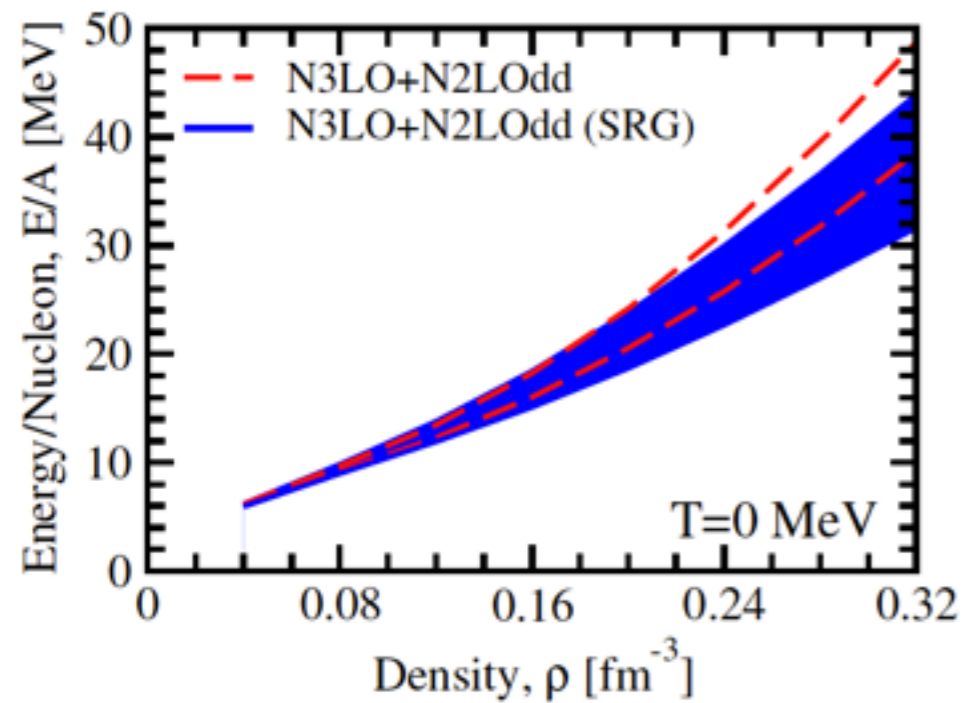
Dong, Lombardo, Zuo
PRC **87** 062801(R) (2013)

$$\Delta_k^L = Z_k \sum_{L'} \int_{k'} Z_{k'} \frac{\langle k | V_{nn}^{LL'} | k' \rangle}{2\chi_{k'}} \Delta_{k'}^{L'} + \chi_k = \sqrt{\varepsilon_k^2 + |\Delta_k|^2}$$

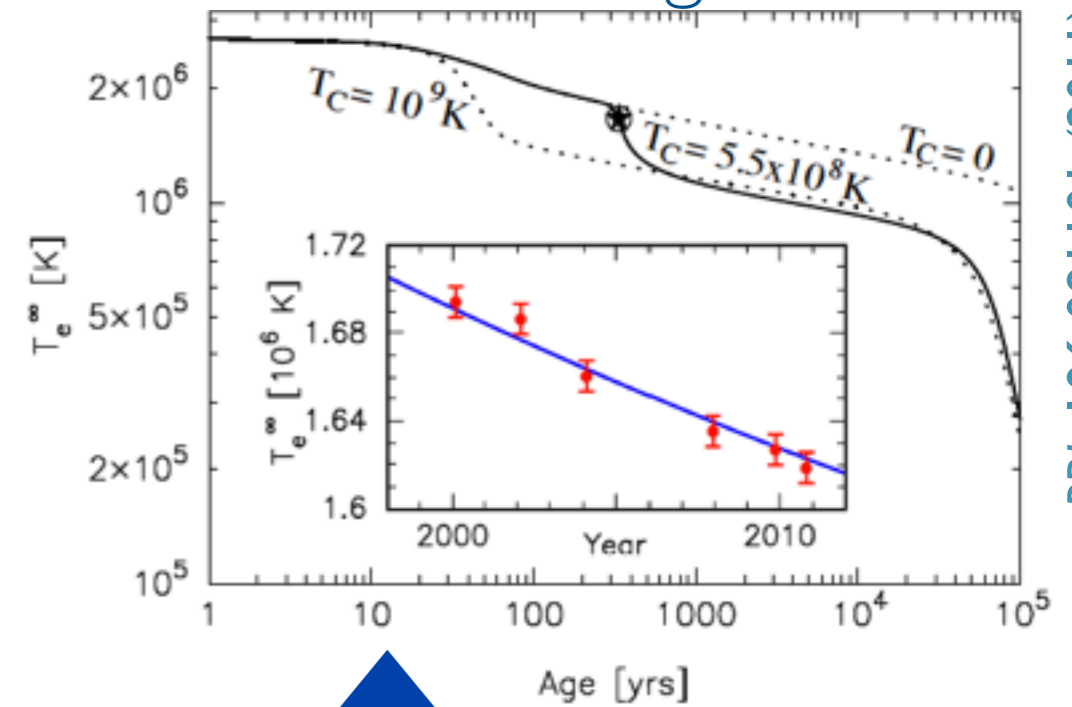
$$+ \varepsilon_k = \frac{k^2}{2m} + U(k) - \mu$$

Chiral EoS & pairing gaps

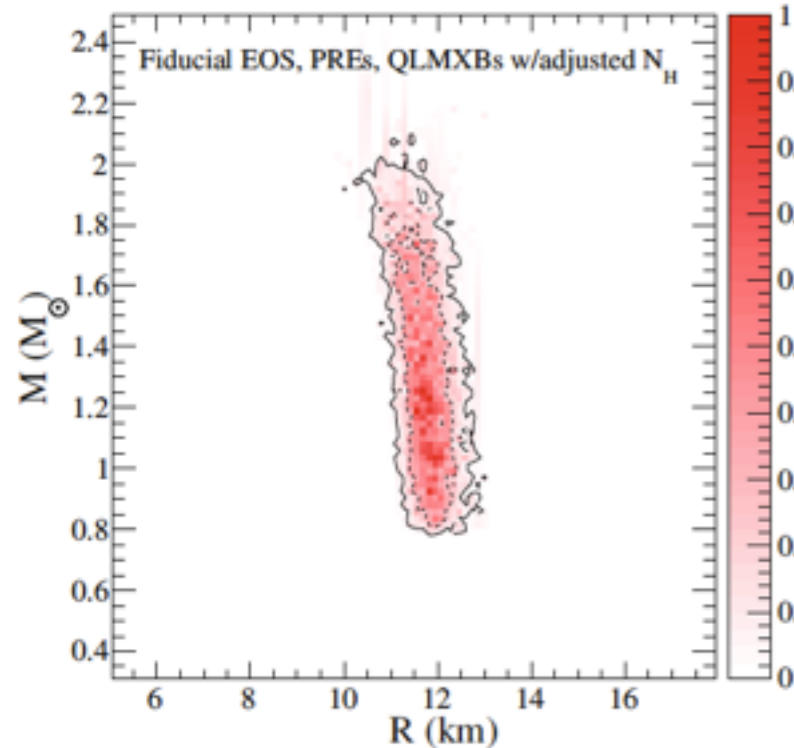
N^3 LO EoS with 3BFs



CasA cooling curve

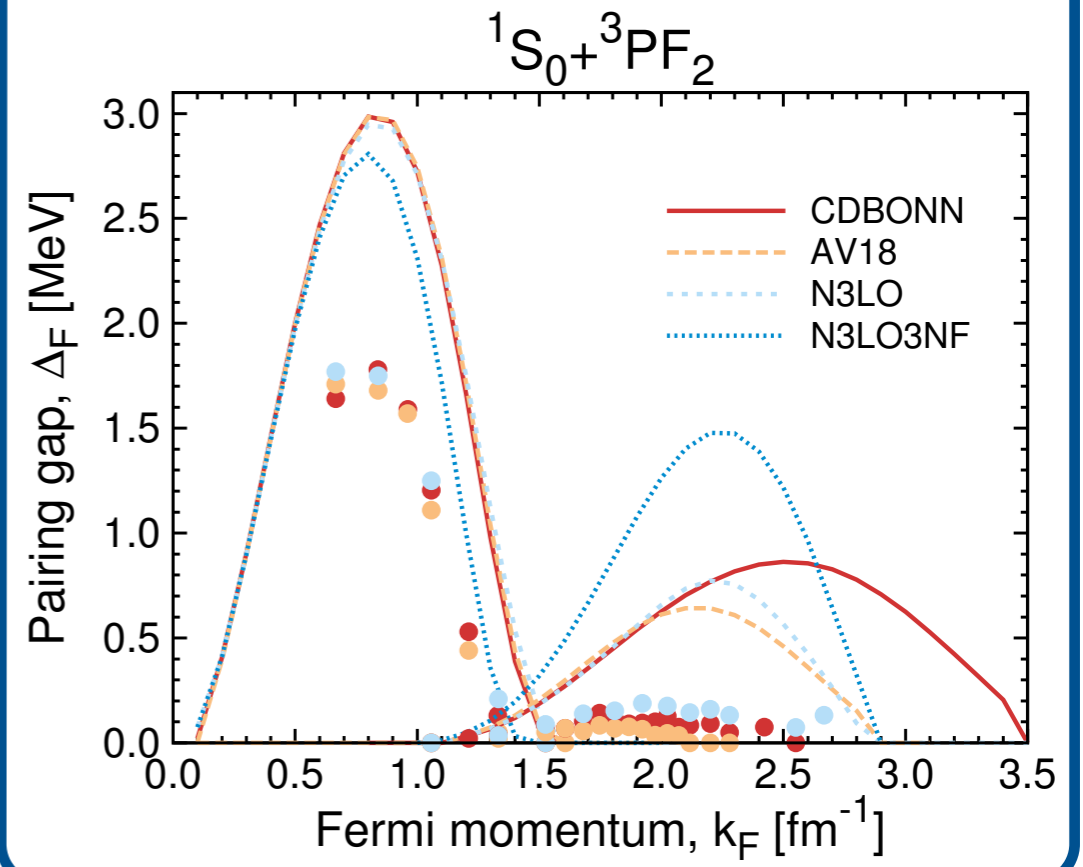


PRL 106 081101 (2011)



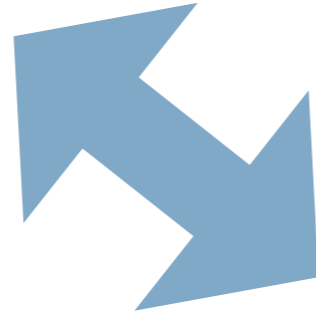
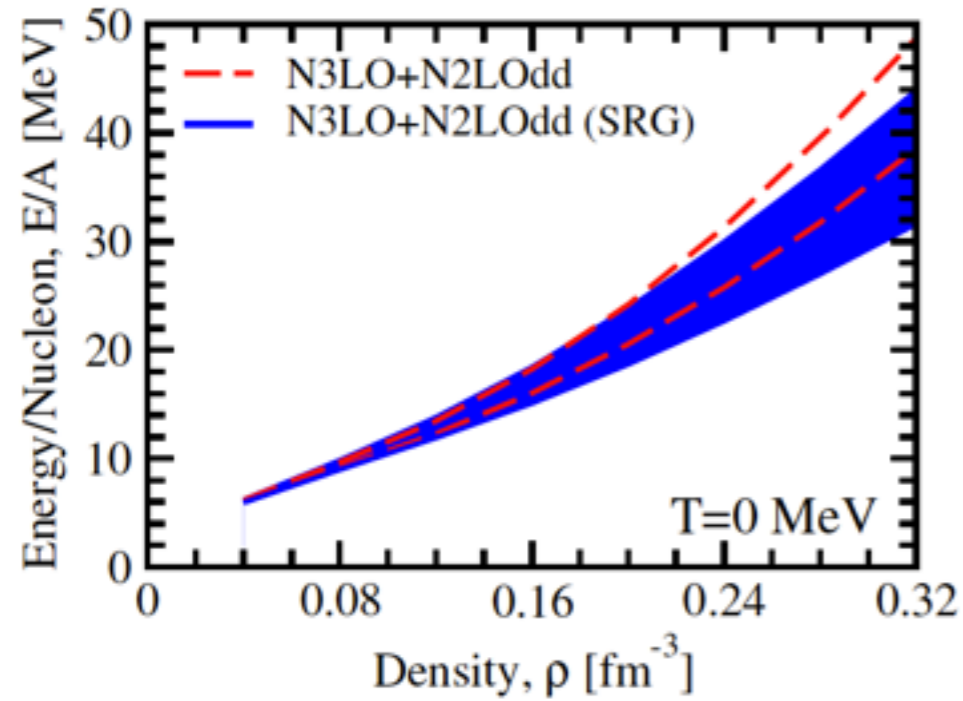
Steiner & Lattimer, EPJA 50 49 (2014)

Pairing gap beyond BCS

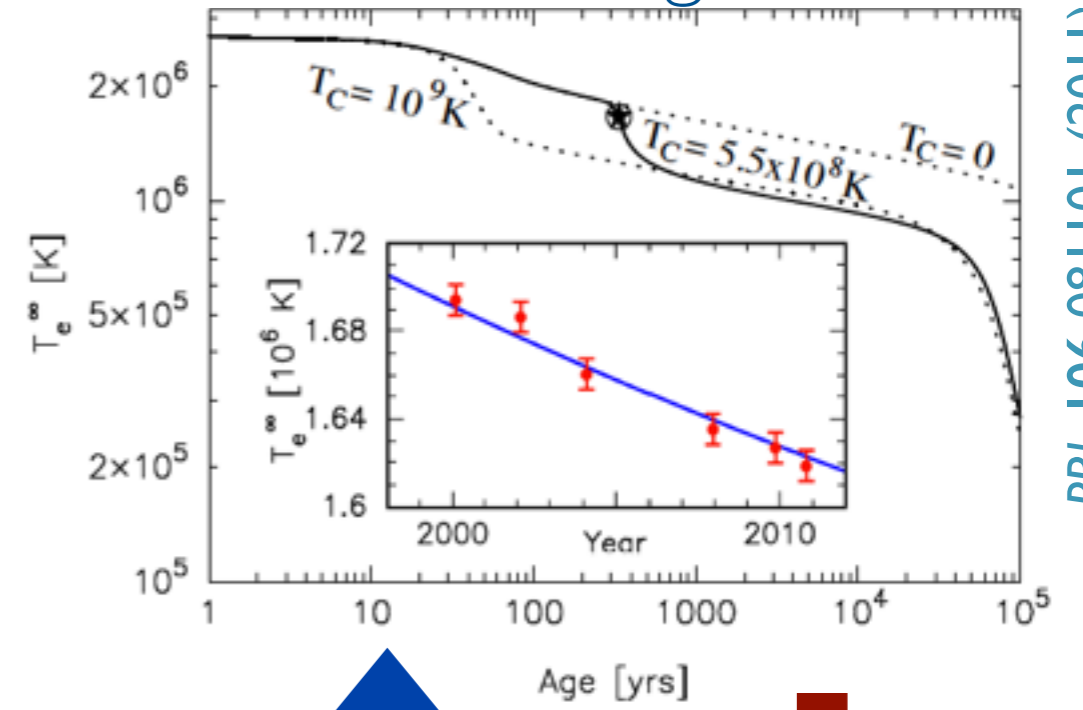


Chiral EoS & pairing gaps

N^3 LO EoS with 3BFs

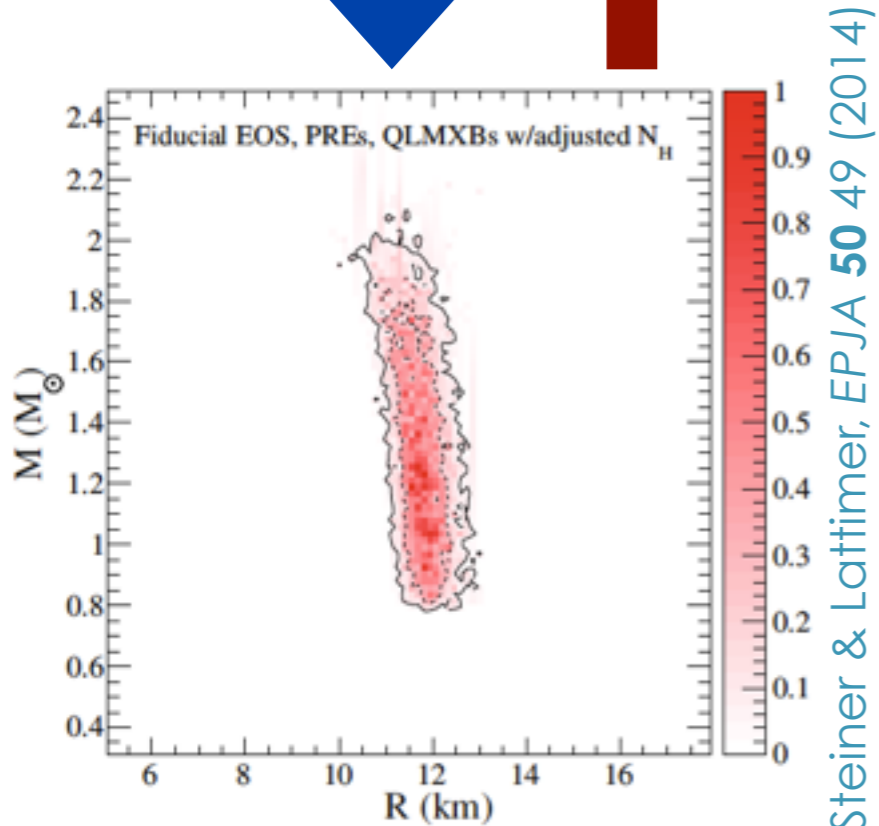
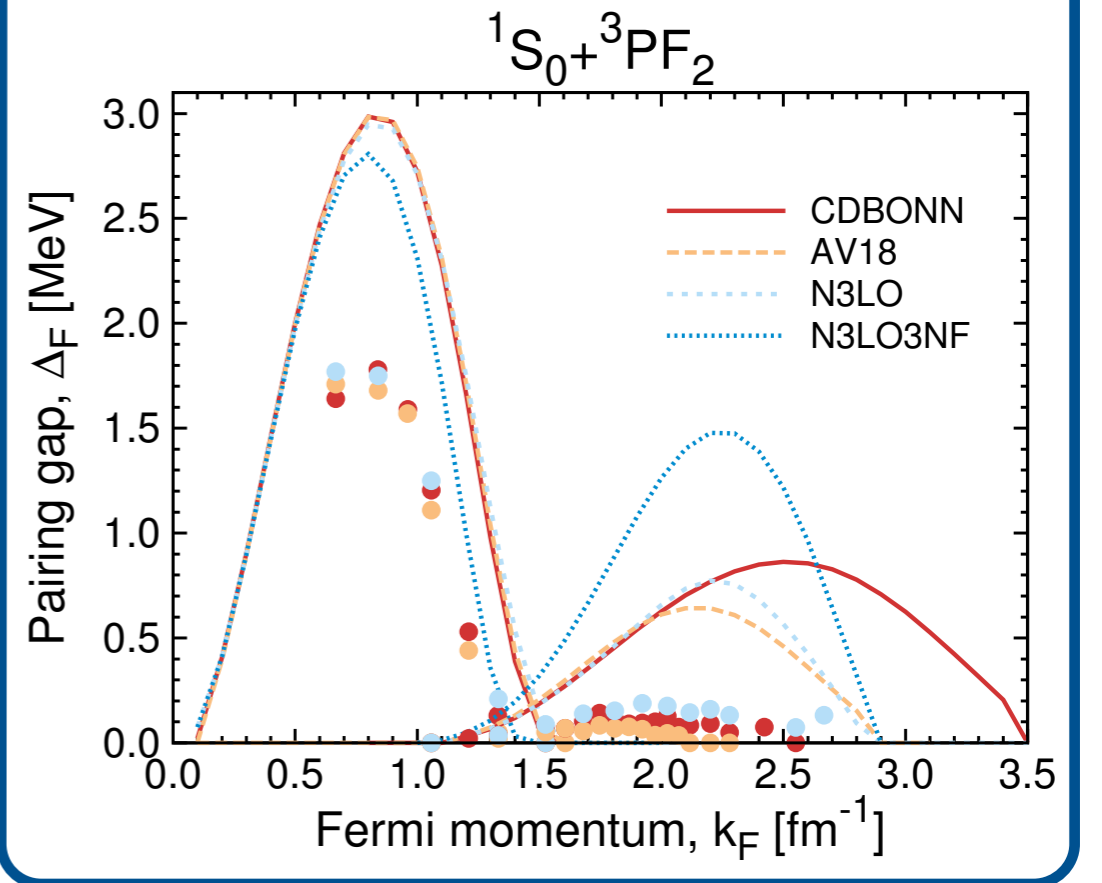


CasA cooling curve



PRL 106 081101 (2011)

Pairing gap beyond BCS



- Ab initio nuclear **theory** to treat **pairing** systems
- **Different** NN forces give robust predictions
- Approximations introduced in a **meaningful** way
- Challenges ahead:
 - **pp** pairing?
 - isospin **asymmetric** matter?
 - **Cooling?**



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