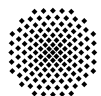




Internally contracted multireference coupled-cluster theory

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Universität Stuttgart

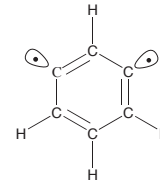
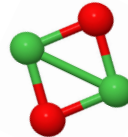


Deutsche
Forschungsgemeinschaft

Target now: Molecules

Actually: Electrons in the field of (point-like) nuclei fixed in space

(Keywords: Born-Oppenheimer approximation, clamped nuclei approximation)



We solve for (mostly) stationary states of the electronic (non-relativistic) Hamiltonian

$$h_p^q = \langle p | \left(-\frac{1}{2} \Delta - \sum_I \frac{Z_I}{|\mathbf{r}_1 - \mathbf{R}_I|} \right) | q \rangle$$

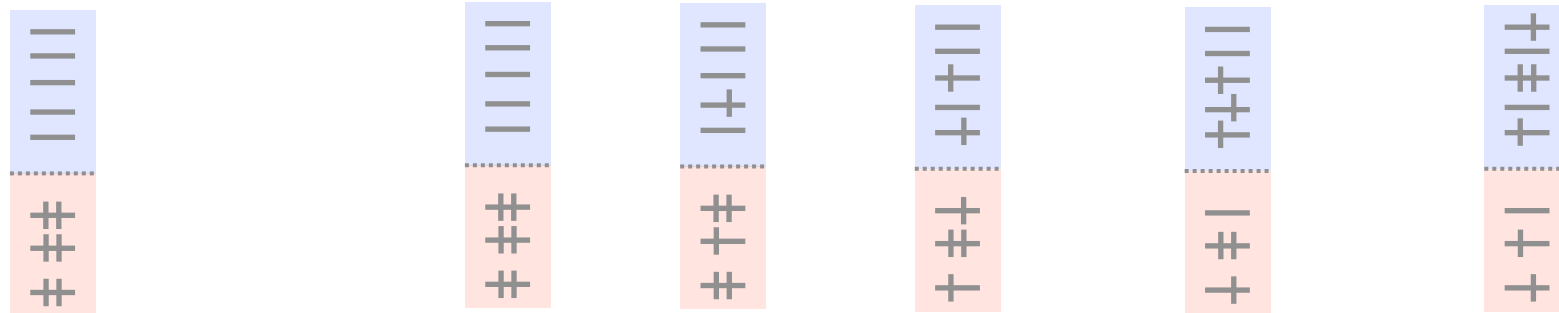
$$g_{pr}^{qs} = \langle pr | \frac{1}{|\mathbf{r}_1 - \mathbf{r}_2|} | qs \rangle$$

$$\hat{H} = \sum_{pq} h_p^q a_p^\dagger a_q + \frac{1}{4} \sum_{pqrs} g_{pr}^{qs} a_p^\dagger a_r^\dagger a_s a_q = h_p^q a_q^p + \frac{1}{4} g_{pr}^{qs} a_{qs}^{pr}$$

creation/annihilation of spin-orbitals

short-hand notation

Many-electron states in molecules



$$\begin{aligned}
 |\Phi_0\rangle \quad \Rightarrow \quad |\Psi\rangle &= c_0|\Phi_0\rangle + c_a^i|\Phi_i^a\rangle + c_{ab}^{ij}|\Phi_{ij}^{ab}\rangle + c_{abc}^{ijk}|\Phi_{ijk}^{abc}\rangle + c_{abcd}^{ijkl}|\Phi_{ijkl}^{abcd}\rangle + \dots \\
 &\propto |\Phi_0\rangle + |\Psi_{\text{korr}}\rangle
 \end{aligned}$$

Coupled-Cluster (CC) ansatz:

$$|\Psi_{\text{CC}}\rangle = e^{\hat{T}}|\Phi_0\rangle = |\Phi_0\rangle + \hat{T}|\Phi_0\rangle + \frac{1}{2}\hat{T}^2|\Phi_0\rangle + \dots$$

Coester, Kümmel (1958)
 Čížek, Paldus (1966)
 Bartlett (1978)

Coupled-cluster expansion

Ansatz:

$$|\Psi_{CC}\rangle = e^{\hat{T}} |\Phi_0\rangle \quad \leftarrow \text{single determinant}$$

$$\hat{T} = \hat{T}_1 + \hat{T}_2 + \hat{T}_3 + \dots$$

Features:

separability

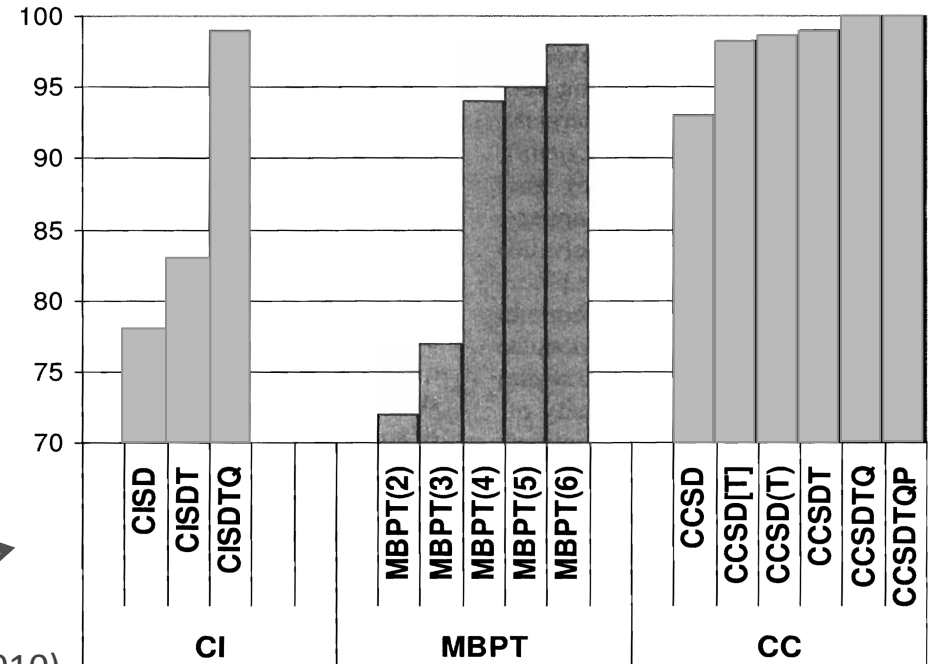
$$\begin{aligned} |\Psi_{CC}^{AB}\rangle &= e^{\hat{T}_A + \hat{T}_B} |\Phi_0^{AB}\rangle \\ &= \hat{A} \left(e^{\hat{T}_1} |\Phi_0^A\rangle \right) \left(e^{\hat{T}_2} |\Phi_0^B\rangle \right) \\ &= \hat{A} |\Psi_{CC}^A\rangle |\Psi_{CC}^B\rangle \end{aligned}$$

$$E^{AB} = E^A + E^B$$

fast convergence with cluster rank

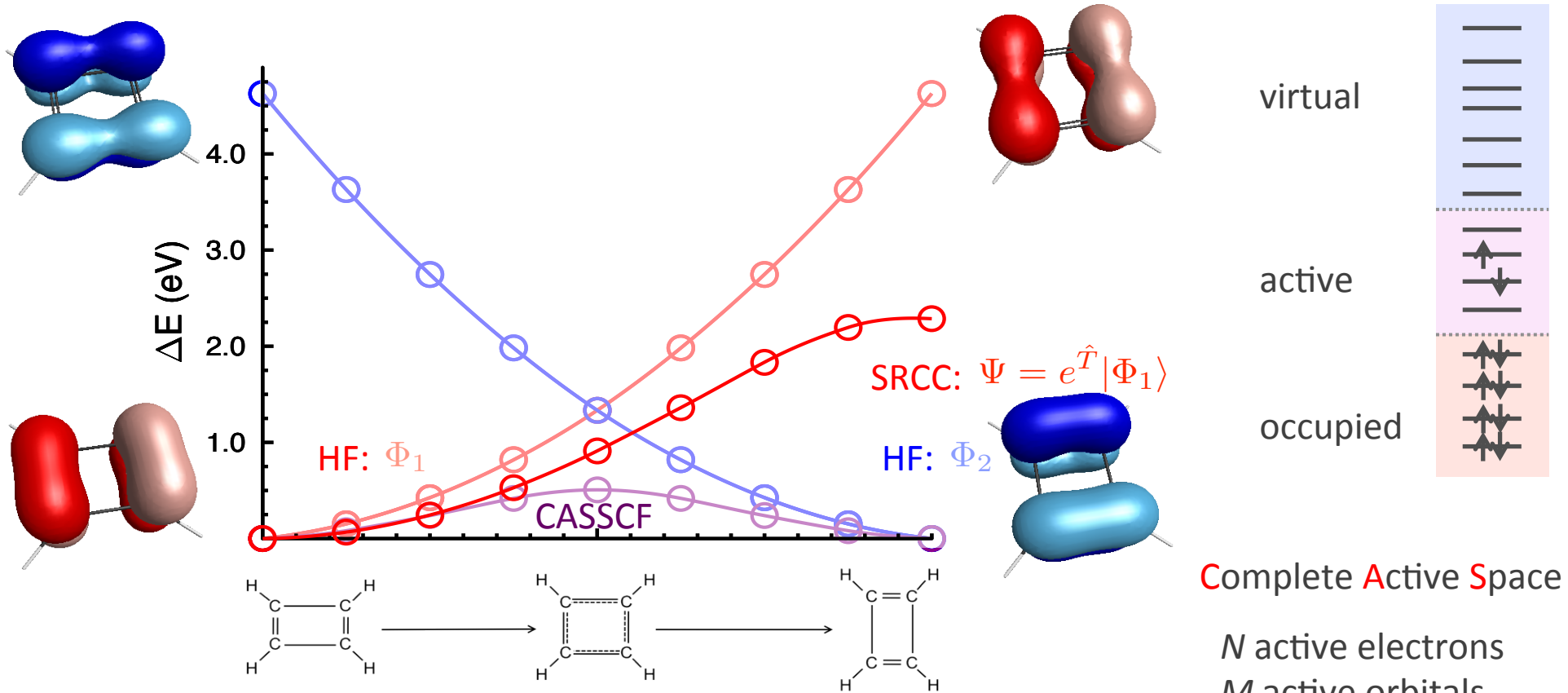


Percentage of E_{corr} recovered for set of small molecules
 Bartlett, Mol Phys 108, 2905 (2010)



Multireference cases

Example: Automerization of cyclobutadiene



Overview: Different approaches

$$|\Psi_{CC}\rangle = e^{\hat{T}} |\Phi_0\rangle \quad \longrightarrow \quad |\Psi_{\text{ext-SRCC}}\rangle = e^{\hat{T}_{\text{ex}}} e^{\hat{T}_{\text{int}}} |\Phi_0\rangle$$

Adamowicz (1992), Piecuch, Kallay, Olsen, Paldus, Bartlett, ...

$$|\Psi_{JM}\rangle = \sum_{\mu} e^{\hat{T}(\mu)} |\Phi_{\mu}\rangle c_{\mu}$$

Jeziorski, Monkhorst (1981)

$$|\Psi_{ic}\rangle = e^{\hat{T}} \sum_{\mu} |\Phi_{\mu}\rangle c_{\mu}$$

[Cicek, 1967], Mukherjee (1975)

**Valence universal:
Fock-Space MRCC**

Haque, Mukherjee (1984)
Stolarczyk, Monkhorst (1985)
Kaldor (1986)
recent work by: Musiał, Bartlett

State-specific: ic-MRCC

Banerjee, Simons (1981)
Nooijen et al. (2001,2011)
Olsen (2005)
Evangelista, Gauss (2011)
Hanauer, AK (2011)
compare: CT theory
Chan, Yanai (2006)

State-universal
Piecuch, Jeziorski; Kucharski, Bartlett; Li, Paldus

State-specific

Mášik, Hubač (1998)
Pittner et al. (2006)
Mahapatra, .
Datta, Mukherjee (1998)
Evangelista, Allen, Gauss et al. (2007)
Das, Kallay, Mukherjee (2009)
Hanrath (2005)

Reviews about MRCC: For example ...

Lyakh, Musiał, Lotrich, Bartlett, Chem. Rev. 112, 182 (2012)
Köhn, Hanauer, Mück, Jagau, Gauss, WIREs CMS 3, 176 (2013)

Design of a multireference CC theory

Single reference:

$$|\Phi_0\rangle \quad \longrightarrow \quad |\Psi\rangle = e^{\hat{T}} |\Phi_0\rangle$$

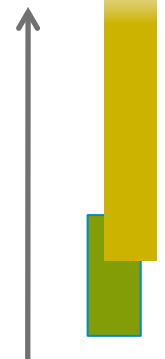
State-universal ansatz: (Jeziorski, Monkhorst)

$$\{|\Phi_\mu\rangle\} \quad \longrightarrow \quad |\Psi_{\text{JM}}\rangle = \sum_{\mu} e^{\hat{T}(\mu)} |\Phi_\mu\rangle c_\mu$$

Intruder state problems: State-specific versions

State-specific internally contracted ansatz

$$|\Psi_0\rangle = \sum_{\mu} |\Phi_\mu\rangle c_\mu \quad \longrightarrow \quad |\Psi_{\text{ic}}\rangle = e^{\hat{T}} \sum_{\mu} |\Phi_\mu\rangle c_\mu$$



+ Simpler, better control of intruders

- Questionable at state crossings (!)

Internally contracted MRCC: Why?

$$|\Psi_{ic}\rangle = e^{\hat{T}} \sum_{\mu} |\Phi_{\mu}\rangle c_{\mu}$$

“internally contracted”: originally used in CI context by Siegbahn, Mayer, Werner (1980s)

Promising features:

- Genuine multireference scheme (no bias within model space)
- Compact amplitude manifold (no scaling with model space size)
- Spin adaption or spin-free formulation is easy
- Orbital invariance

Evangelista, Gauss, JCP 134, 114102 (2011)
Hanauer, Köhn, JCP 134, 204111 (2011)

What is scary:

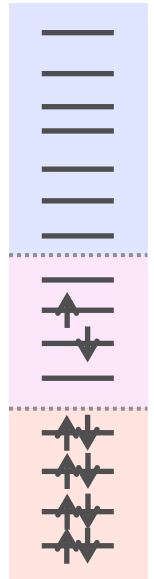
- Non-commuting operator manifold
- Linear dependencies

Simplified approaches, see e.g.

Banerjee, Simons, IJQC 19, 207 (1981)

Laidig, Saxe, Bartlett, JCP 86, 887 (1987)

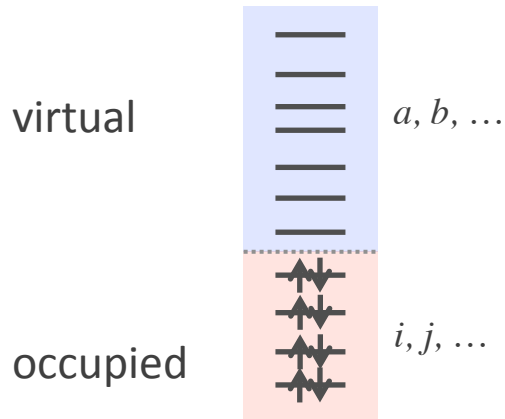
Fink, Staemmler, TCA 87, 129 (1993)



ic-MRCC vis-à-vis single-reference CC

Orbital space:

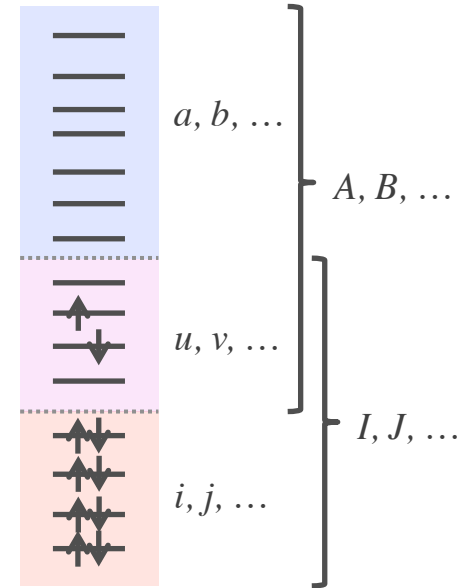
Hartree-Fock



virtual

active

occupied



CASSCF

Reference wave function:

$$|\Phi_0\rangle$$

$$|\Psi_0\rangle = \sum_{\mu} |\Phi_{\mu}\rangle c_{\mu}$$

Correlated wave function:

$$|\Psi\rangle = e^{\hat{T}} |\Phi_0\rangle$$

$$|\Psi_{ic}\rangle = e^{\hat{T}} |\Psi_0\rangle$$

Cluster operator:

$$\hat{T} = \sum_{ia} t_a^i a_i^a + \frac{1}{4} \sum_{ijab} t_{ab}^{ij} a_{ij}^{ab} + \dots$$

$$\hat{T} = \sum_{IA} t_A^I a_I^A + \frac{1}{4} \sum_{IJAB} t_{AB}^{IJ} a_{IJ}^{AB} + \dots$$

ic-MRCC vis-à-vis single-reference CC

Correlated wave function:

$$|\Psi\rangle = e^{\hat{T}}|\Phi_0\rangle$$

$$|\Psi_{ic}\rangle = e^{\hat{T}}|\Psi_0\rangle$$

Cluster operator:

$$\begin{aligned} \hat{T} &= \sum_{ia} t_a^i a_i^a \\ &+ \frac{1}{4} \sum_{ijab} t_{ab}^{ij} a_{ij}^{ab} \\ &+ \dots = \sum_{\rho} t_{\rho} \hat{\mathcal{T}}_{\rho} \end{aligned}$$

$$\begin{aligned} \hat{T} &= \sum_{IA} t_A^I a_I^A \\ &+ \frac{1}{4} \sum_{IJAB} t_{AB}^{IJ} a_{IJ}^{AB} \\ &+ \dots = \sum_{\rho} t_{\rho} \hat{\mathcal{T}}_{\rho} \end{aligned}$$

Excitation types (singles&doubles):

$$\hat{\mathcal{T}}_{\rho} = \begin{cases} \hat{a}_i^a \\ \hat{a}_{ij}^{ab} \end{cases}$$

$$\hat{\mathcal{T}}_{\rho} = \begin{cases} \hat{a}_i^a & \hat{a}_u^a \\ \hat{a}_i^u & \hat{a}_i^u \\ \hat{a}_{ij}^{ab} & \hat{a}_{uj}^{ab} & \hat{a}_{uv}^{ab} \\ \hat{a}_{ij}^{ub} & \hat{a}_{vj}^{ub} & \hat{a}_{vx}^{ub} \\ \hat{a}_{ij}^{uw} & \hat{a}_{vj}^{uw} \end{cases}$$

virtual

 a, b, \dots
 A, B, \dots

active

 u, v, \dots
 I, J, \dots

occupied

 i, j, \dots

avoid active-active excitations, instead: optimize

$$|\Psi_0\rangle = \sum_{\mu} |\Phi^{\mu}\rangle c_{\mu}$$

ic-MRCC vis-à-vis single-reference CC

Correlated wave function:

$$|\Psi\rangle = e^{\hat{T}}|\Phi_0\rangle$$

$$|\Psi_{\text{ic}}\rangle = e^{\hat{T}}|\Psi_0\rangle$$

Energy

$$E = \langle\Phi_0|e^{-\hat{T}}\hat{H}e^{\hat{T}}|\Phi_0\rangle$$

$$E = \langle\Psi_0|e^{-\hat{T}}\hat{H}e^{\hat{T}}|\Psi_0\rangle$$

Equations for reference function

$$|\Psi_0\rangle = \sum_{\mu} |\Phi^{\mu}\rangle c_{\mu}$$

$$\sum_{\nu} \langle\Phi_{\mu}|e^{-\hat{T}}\hat{H}e^{\hat{T}}|\Phi^{\nu}\rangle c_{\nu} = E c_{\mu} \quad \equiv \quad \hat{\mathbf{H}}^{\text{eff}} \mathbf{c} = E \mathbf{c}$$

Amplitude equations

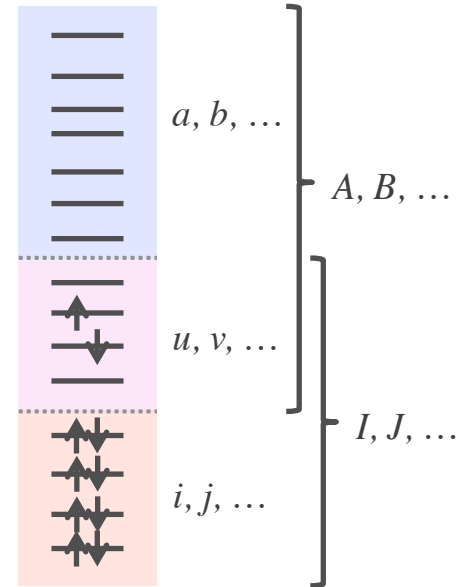
$$0 = \langle\Phi_0|(\hat{\tau}_{\rho})^{\dagger} e^{-\hat{T}}\hat{H}e^{\hat{T}}|\Phi_0\rangle$$

$$0 = \langle\Psi_0|(\hat{\tau}_{\rho})^{\dagger} e^{-\hat{T}}\hat{H}e^{\hat{T}}|\Psi_0\rangle$$

virtual

active

occupied



Here it becomes more difficult!

Linear dependencies

$$0 = \langle \Psi_0 | (\hat{\tau}_\rho)^\dagger e^{-\hat{T}} \hat{H} e^{\hat{T}} | \Psi_0 \rangle$$

non-orthogonal

$$S_{\rho\sigma} = \langle \Psi_0 | (\hat{\tau}_\rho)^\dagger \hat{\tau}_\sigma | \Psi_0 \rangle \neq \delta_{\rho\sigma}$$

linear dependencies, e.g.

$$\hat{a}_{1\bar{1}}^{ab} (c_1 |1\bar{1}\rangle + c_2 |2\bar{2}\rangle) = c_1 |ab\rangle$$

$$\hat{a}_{2\bar{2}}^{ab} (c_1 |1\bar{1}\rangle + c_2 |2\bar{2}\rangle) = c_2 |ab\rangle$$

$$\sum_u \hat{a}_{ui}^{ua} | \Psi_0 \rangle = n_a \hat{a}_i^a | \Psi_0 \rangle$$

Consequences:

covariant and contravariant quantities

$$\hat{A} = \sum_\rho A_\rho \hat{\tau}_\rho$$

$$\bar{A}_\rho = \langle \Psi_0 | (\hat{\tau}_\rho)^\dagger \hat{A} | \Psi_0 \rangle = \sum_\sigma S_{\rho\sigma} A_\sigma$$

projection $\langle \Psi_0 | (\hat{\tau}_\rho)^\dagger$ has less degrees of freedom than operator basis $\sum_\rho t_\rho \hat{\tau}_\rho$

particle ranks are not orthogonal, e.g.

$$\langle \Psi_0 | a_{vw}^{ui} a_i^w | \Psi_0 \rangle \neq 0$$

Linear dependencies

project to orthogonalized, non-redundant basis:

technically identical to ic-MRCI, CASPT2 etc. $S_{\rho\sigma} = \langle \Psi_0 | (\hat{\tau}_\rho)^\dagger \hat{\tau}_\sigma | \Psi_0 \rangle = S_{\rho_a \sigma_a}^{\text{act.}} \otimes \mathbf{1}_{\rho_i \sigma_i}^{\text{inact.}}$

icMRCCSD: up to
3-particle density

$$\Rightarrow \hat{\tau}'_{\tilde{\rho}} = X_{\tilde{\rho}}^\sigma \tau_\sigma \quad \text{where} \quad \hat{\tau}' = \hat{\tau}^{\parallel} + \hat{\tau}^{\perp} \quad \text{with} \quad \hat{\tau}'_{\tilde{\rho}} | \Psi_0 \rangle = 0$$

$$\Longrightarrow \hat{T} = \hat{T}^{\parallel} + \hat{T}^{\perp}$$

we have the *freedom* to set $\hat{T}^{\perp} = 0$ (will still converge to full CI)*

*Evangelista, Hanauer, Köhn, Gauss, JCP 136, 204108 (2012)

Modified amplitude equations

$$0 = \langle \Psi_0 | (\hat{\tau}'_{\tilde{\rho}})^\dagger e^{-\hat{T}^{\parallel}} \hat{H} e^{\hat{T}^{\parallel}} | \Psi_0 \rangle$$

! A nasty problem: near-linear dependencies !

Technical note: $X_{\tilde{\rho}}^\sigma = (S_\eta^{-\frac{1}{2}})_{\tilde{\rho}}^\sigma$ + Schmidt-orthog. betw. particle ranks

$$\Omega_{\tilde{\rho}}^{\parallel} = X_{\tilde{\rho}}^\sigma \langle \Psi_0 | (\tau_\sigma)^\dagger \bar{H} | \Psi_0 \rangle$$

$$\Delta t_{\tilde{\rho}}^{\parallel} = (X^\dagger)_{\tilde{\rho}}^{\tilde{\sigma}} \left(\Omega_{\tilde{\sigma}}^{\parallel} / \Delta \epsilon_{\tilde{\sigma}} \right)$$

$$\longrightarrow \langle \Psi_0 | \hat{T}_1^\dagger \hat{T}_2 | \Psi_0 \rangle \stackrel{!}{=} 0$$

proper choice important for extensivity!

Hanauer, Köhn, JCP 134, 204111 (2011)

Hanauer, Köhn, JCP 137, 131103 (2012)

Numerical evaluation

$$\sum_{\nu} \langle \Phi_{\mu} | e^{-\hat{T}} \hat{H} e^{\hat{T}} | \Phi^{\nu} \rangle c_{\nu} = E c_{\mu} \quad 0 = \langle \Psi_0 | (\hat{\tau}_{\tilde{\rho}}^{\parallel})^{\dagger} e^{-\hat{T}^{\parallel}} \hat{H} e^{\hat{T}^{\parallel}} | \Psi_0 \rangle$$

$$e^{-\hat{T}} \hat{H} e^{\hat{T}} = \hat{H} + [\hat{H}, \hat{T}] + \frac{1}{2} [[\hat{H}, \hat{T}], \hat{T}] + \dots$$

$[\hat{\tau}_{\rho}, \hat{\tau}_{\sigma}] \neq 0$

$$[\hat{H}, \hat{T}] = \overbrace{\hat{H}\hat{T}} - \overbrace{\hat{T}\hat{H}}$$

icMRCCSD Energy: up to 4-fold comm.

Equations: up to 8-fold comm.

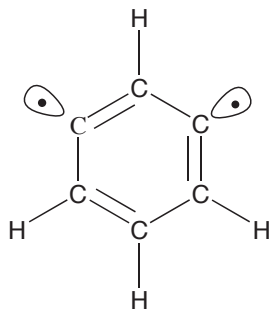
$e^{-\hat{T}} \hat{H} e^{\hat{T}}$	CCSD	ic-MRCCSD / CAS(6,6)
$= \hat{H}$	2	27
$+ [\hat{H}, \hat{T}]$	13	690
$+ \frac{1}{2} [[\hat{H}, \hat{T}], \hat{T}]$	21	8 766
$+ \frac{1}{6} [[[\hat{H}, \hat{T}], \hat{T}], \hat{T}]$	8	70 997
$+ \frac{1}{24} [[[[\hat{H}, \hat{T}], \hat{T}], \hat{T}], \hat{T}]$	1	399 839
$+ \dots$	0	...

Commutator approximation

Singlet-Triplet separation of m-Benzyne:

icMRCCSD, cc-pVDZ basis, CAS(2,2)

frozen core



$$\langle \Psi_0 | (\tau'_\rho)^\dagger \hat{H} | \Psi_0 \rangle = 0$$

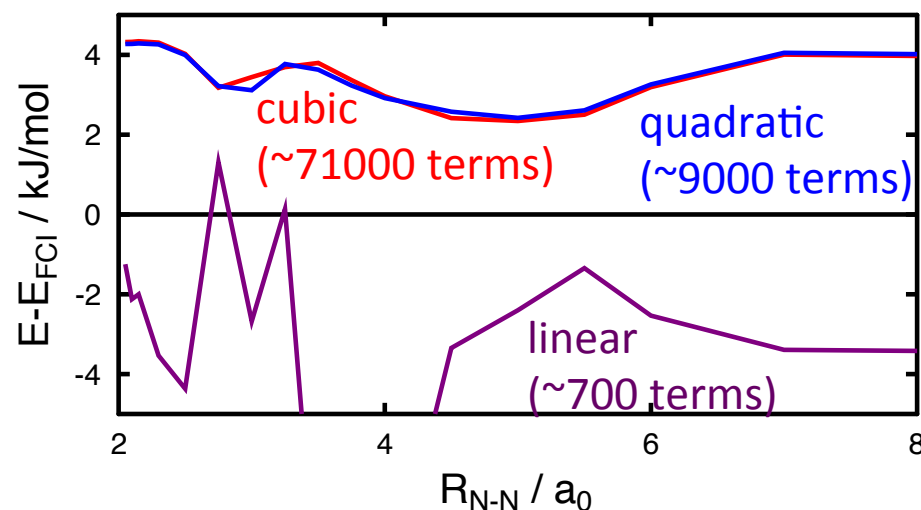
$$\sum_\nu \langle \Phi_\mu | \hat{H} | \Phi^\nu \rangle c_\nu = E c_\mu$$

$N_{\text{com}}(\tau)$	$N_{\text{com}}(E)$	$\Delta E_{\text{ST}} / \text{kcal mol}^{-1}$
4	4	17.289
3	3	17.289
2	2	17.415
2	4	17.415

N₂ stretching:

N₂, cc-pVDZ, CAS(6,6)

ic-MRCCSD vs. FCI



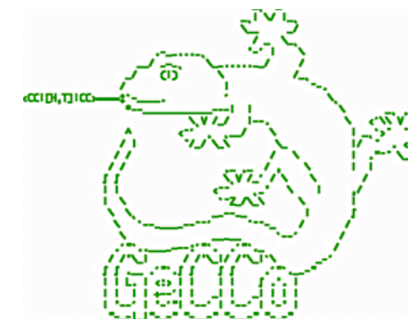
Numerical evaluation

Dealing with complex equations:

- Derivation by **automated techniques**

using GeCCo program

(Integrals, CASSCF: DALTON, GAMESS, Molpro)



- **Approximations:** Truncate after double commutator

tested: Evangelista, Gauss, JCP 134, 114102(2011)

Hanauer, Köhn, J. Chem. Phys. 134, 204111 (2011)

- Present implementation (GeCCo):

Certainly not perfect ...

... BUT: **correct N^6 scaling** (for ic-MRCCSD and fixed active space)



GeCCo

Very new:

PYTHON powered
input language for F90 back-end

define operators

symbolic algebra with operators

general order contraction kernel for numerical evaluation

```
#####
# operators associated with T
#####
new_target('T-Operators')

PRINT({STRING: '=== Cluster Operators ==='})

DEF_OP_FROM_OCC({LABEL: 'Tex', DESCR: ' [PV],H|P,V|P[PV], [HV] [HV] |VV,H[HV]
DEF_ME_LIST({LIST: 'TLST', OPERATOR: 'Tex', IRREP: 1, '2MS': 0, AB_SYM: +1})
PRINT_MEL_INFO({LIST: 'TLST'})

CLONE_OPERATOR({LABEL: 'Tp', TEMPLATE: 'Tex'})

# Formula for overlap matrix
EXPAND_OP_PRODUCT({LABEL: 'F_SMAT', NEW: True, OP_RES: 'SSCAL',
  OPERATORS: ['C0^+', 'Tv^+', 'Tv', 'C0'],
  IDX_SV: [1, 2, 3, 4]})
FACTOR_OUT({LABEL_RES: 'F_SMAT', LABEL_IN: 'F_SMAT', INTERM: 'F_GAM0'})
PRINT_FORMULA({LABEL: 'F_SMAT'})

OPTIMIZE({LABEL_OPT: 'FOPT_H_C0', LABELS_IN: 'F_H_C0'})

PRINT({STRING: 'Prepare diagonal ...'})
PRECONDITIONER({LIST_PRC: 'D0LST', LIST_INP: 'H0', MODE: 'dia-H'})

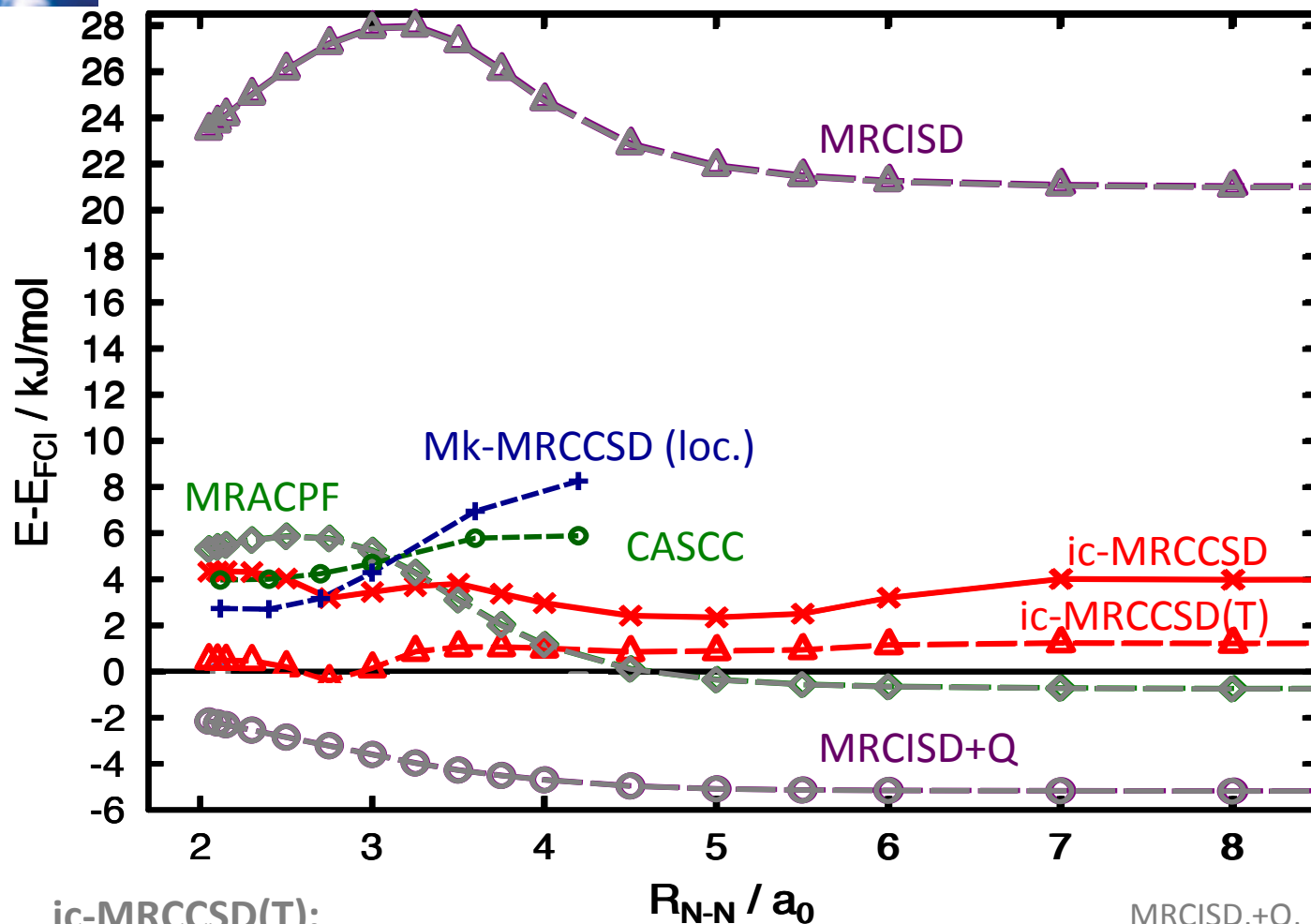
PRINT({STRING: 'Solving equations ...'})

SOLVE_EVP({LIST_OPT: 'C0LST', LIST_PRC: 'D0LST', OP_MVP: 'H_C0', OP_SVP: 'C0',
  FORM: 'FOPT_H_C0', MODE: 'DIA', N_ROOTS: 1})
```

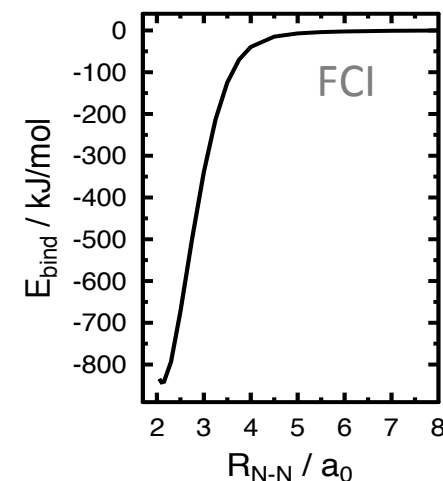


A few examples

N₂ dissociation



cc-pVDZ basis set



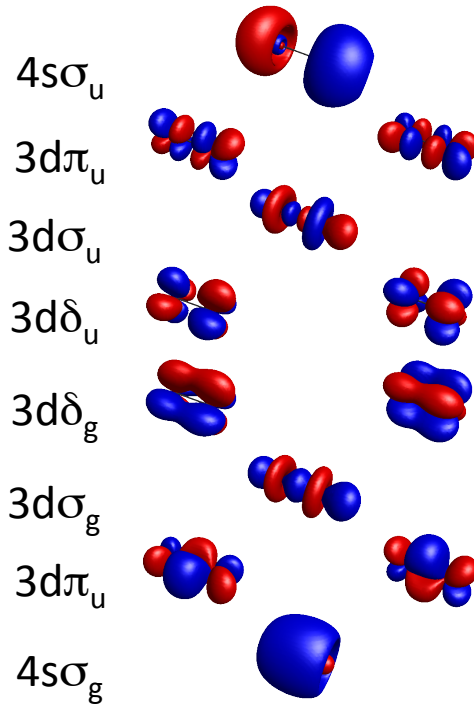
all MR methods use
CAS(6,6)

ic-MRCCSD(T):
Hanauer, Köhn, JCP 136, 204107 (2012)

MRCISD,+Q, MRACPF: Werner, private comm.
MRAQCC: Yanai, Chan, JCP 127, 104107
CASCC: Lyakh et al., Mol. Phys. 105, 1335 (2007)
Mk-MRCCSD: Das, Kallay, Mukherjee, CP 392, 83 (2012)

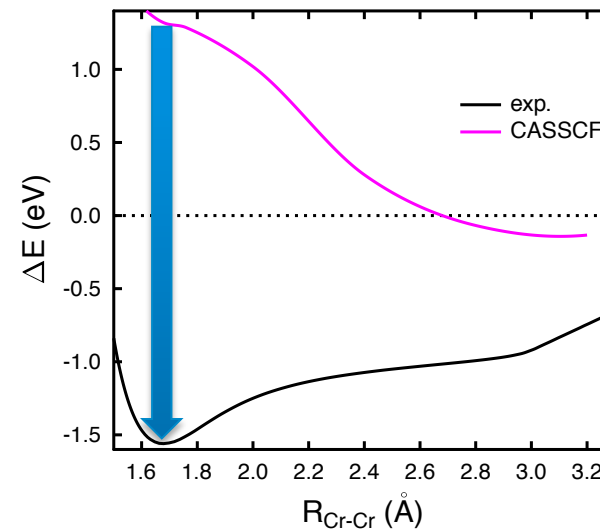
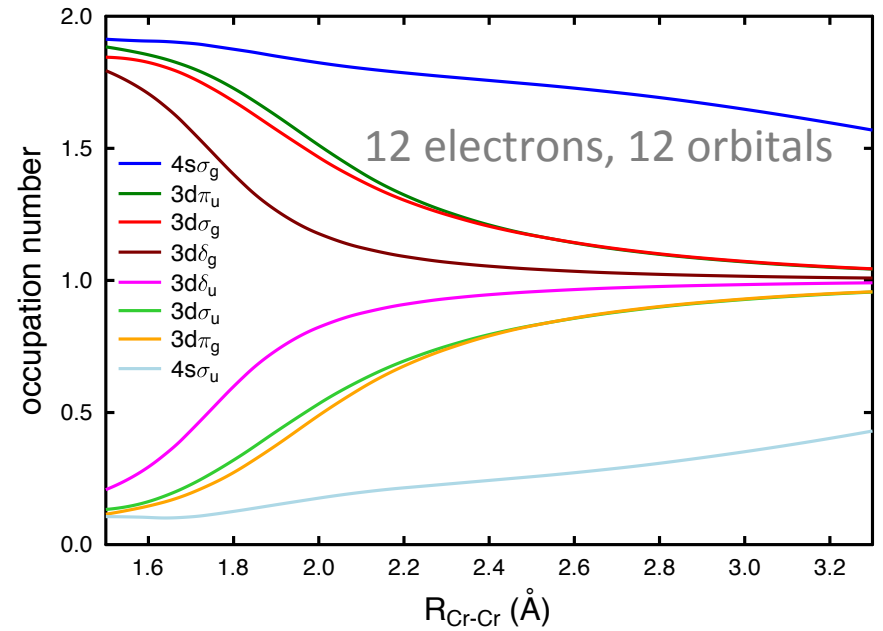
Chromium dimer

Strong
multireference
character

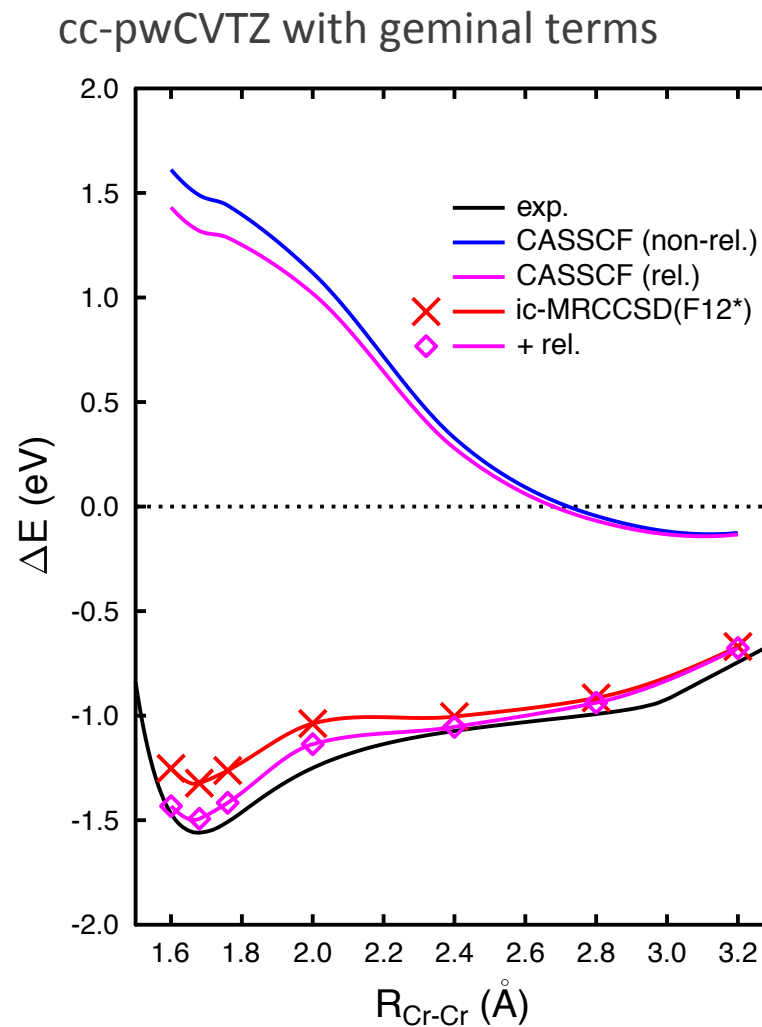
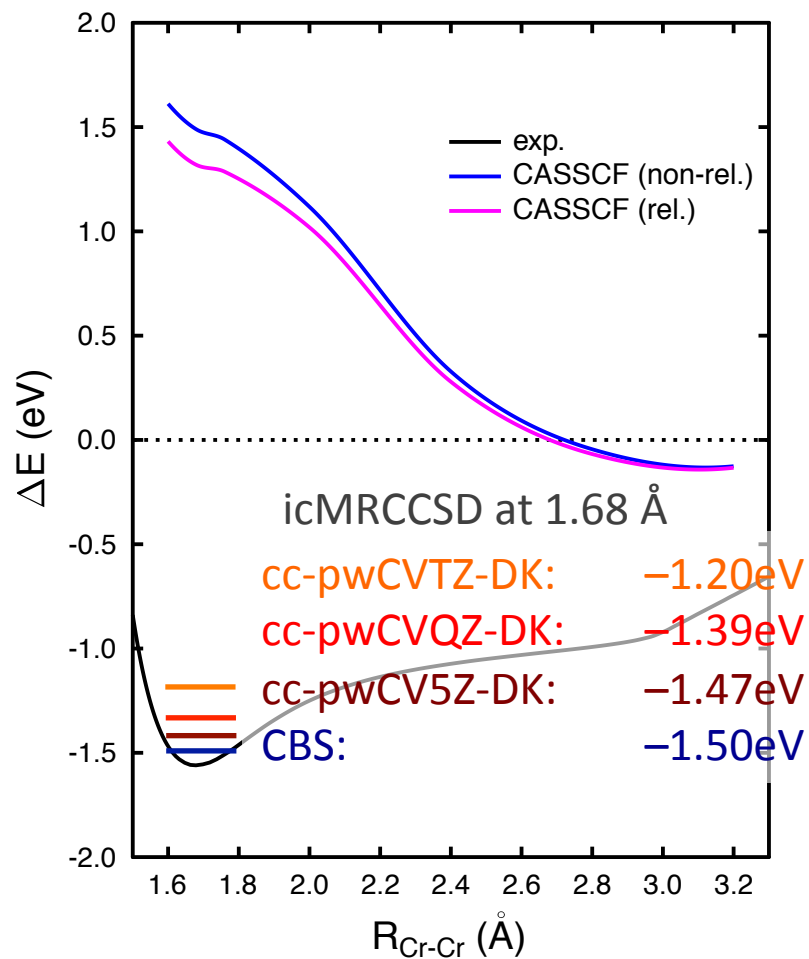


Huge contribution from
dynamic correlation

3s3p-correlation important!!



Some preliminary results



2nd-order DKH transf. for scalar rel. effects

F12 methods

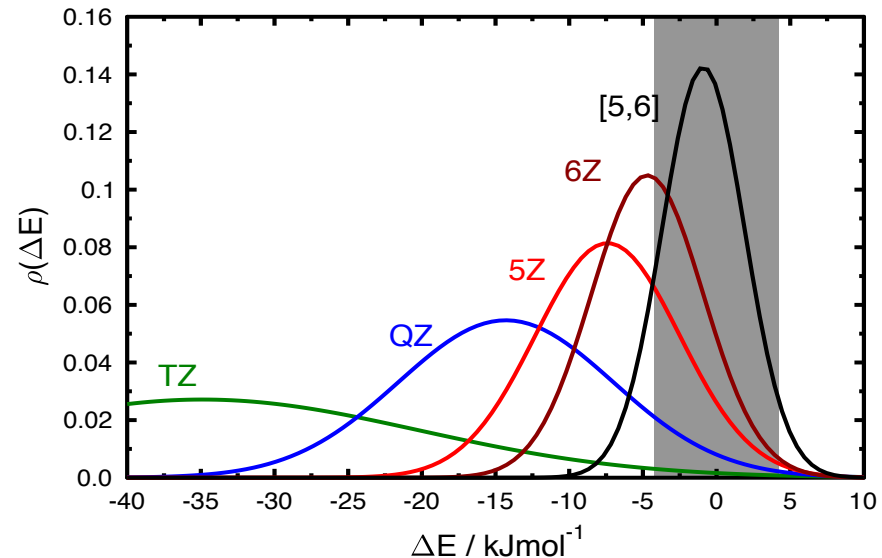
Molecular CC calculations need large basis sets!

Example:

CCSD(T) atomization energies
(20 small molecules)

Error in electronic contribution
compared to reference values
(derived from experiment)

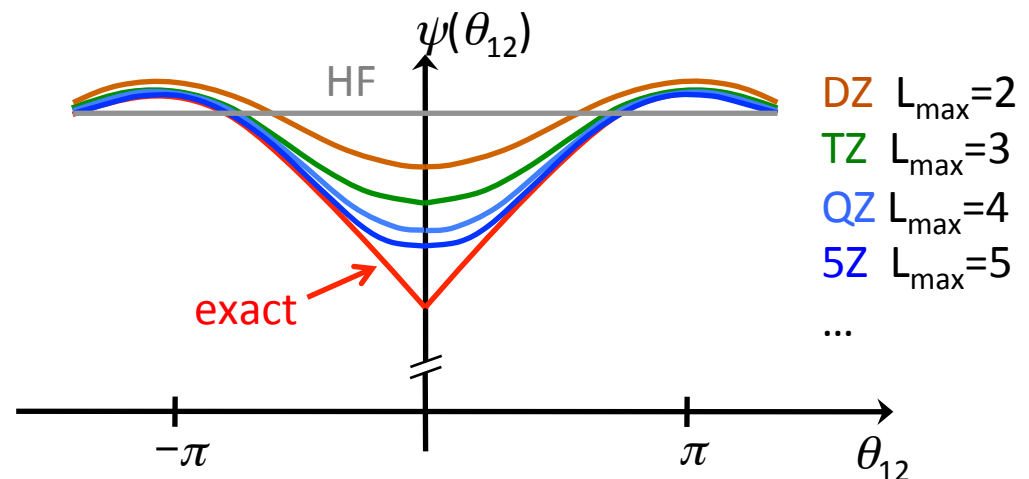
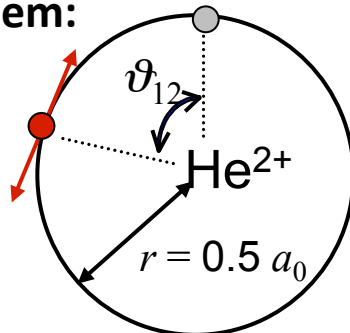
Values from: Helgaker, Jørgensen, Olsen,
"Electronic structure theory", Wiley (2000)



Root of the problem:

Behavior of ψ
for $r_1 \rightarrow r_2$ due to

$$\hat{V}_{ee} \propto \frac{1}{r_{12}}$$



- DZ $L_{\max}=2$
- TZ $L_{\max}=3$
- QZ $L_{\max}=4$
- 5Z $L_{\max}=5$
- ...

F12 theory: concepts

Ansatz for pair-function: $|u_{ij}\rangle = \sum_{a>b} t_{ab}^{ij} |ab\rangle + \sum_{k>l} c_{kl}^{ij} \hat{Q}_{12} f(r_{12}) |kl\rangle$

Correlation factor:

F12-theory (Ten-no):

$$f(r_{12}) = \frac{1}{\gamma} (1 - e^{-\gamma r_{12}})$$

$\gamma \approx 1 \text{ a}_0^{-1}$ is a fixed basis set parameter

Combination with CC theory (Kutzelnigg, Klopper, Noga)

$$T_{2'} = \frac{1}{8} c_{kl}^{ij} \langle \alpha\beta | \hat{Q}_{12} f(r_{12}) | kl \rangle a_{ij}^{\alpha\beta}$$

Efficient combination with CC theory possible:

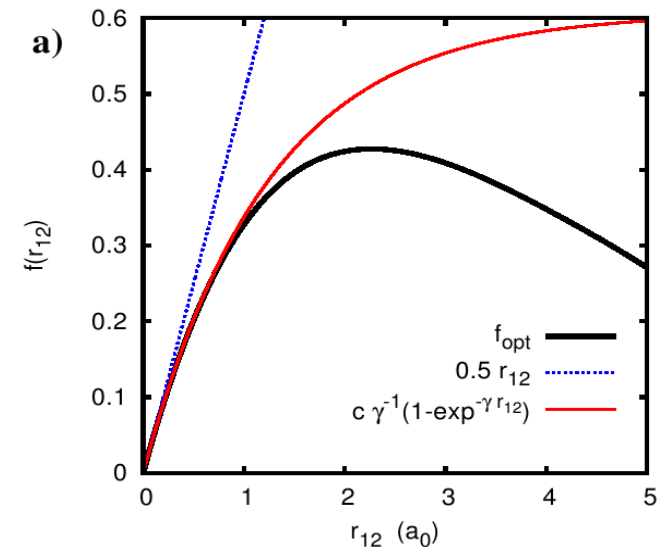
CCSD-F12a/b (Werner, Adler, Knizia),
 CCSD(F12*) (Hättig, Köhn, Tew)

Analogous development: icMRCSD(F12*)

Liu, Hanauer, Köhn, Chem. Phys. Lett. 565, 122 (2013)

$$T_{2'} = \frac{1}{8} c_{KL}^{IJ} \langle \alpha\beta | \hat{Q}_{12} f(r_{12}) | KL \rangle a_{IJ}^{\alpha\beta}$$

correlation factor



From: Tew, Hättig, Klopper, Bachorz. In: "Recent Progress in Coupled Cluster Methods ...", Csarsky, Pittner, Paldus (Eds.), Springer

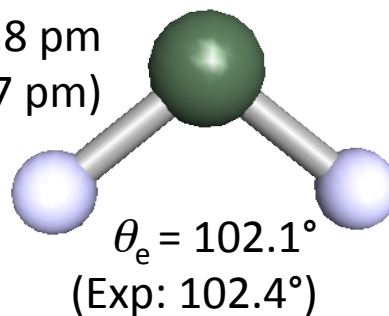
Kutzelnigg (1985), Kutzelnigg, Klopper (1987, 1991),
 Ten-no (2004), May, Manby (2004), Tew, Klopper (2005)

icMRCCSD(F12*) theory

CH₂

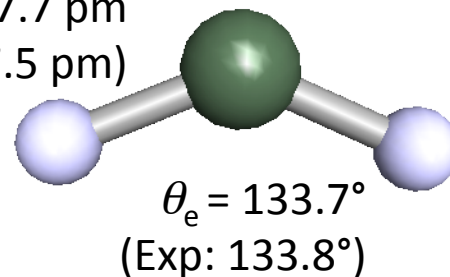
Singlet

$r_e = 110.8$ pm
 (Exp: 110.7 pm)



Triplet

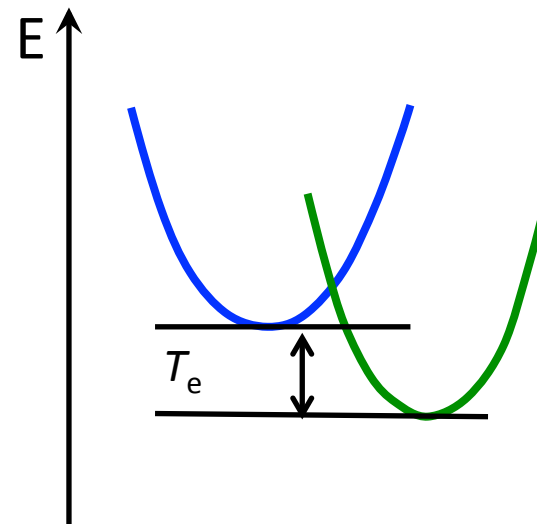
$r_e = 107.7$ pm
 (Exp: 107.5 pm)



CAS(2,2)

ic-MRCCSD(F12*)+(T)/cc-pVQZ-F12

(frozen core, numerical gradients)



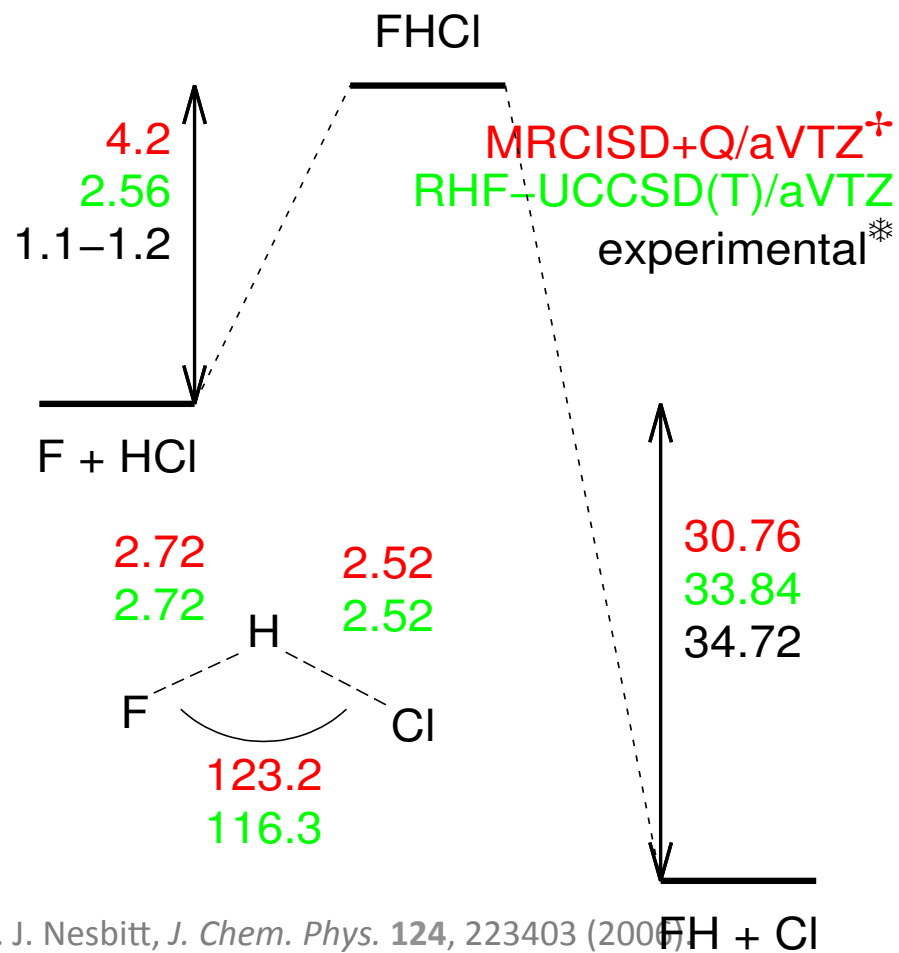
$T_e = 37.2$ kJ/mol
 (Exp: 37.7 kJ/mol)

Exp. values: Petek et al, JCP 91, 6566 (1989)
 Bunker, Sears, JCP 83, 4866 (1985)
 Bunker, Jensen, JCP 79, 1224 (1983)

The F + HCl → Cl + HF reaction

Traditional theoretical calculations overestimate experimental barrier height! †★

Experimental ro-vibrational distribution not reproduced by quantum dynamics.★



[†]M. P. Deskevich, M. Y. Hayes, K. Takahashi, R. T. Skodje and D. J. Nesbitt, *J. Chem. Phys.* **124**, 223403 (2006)

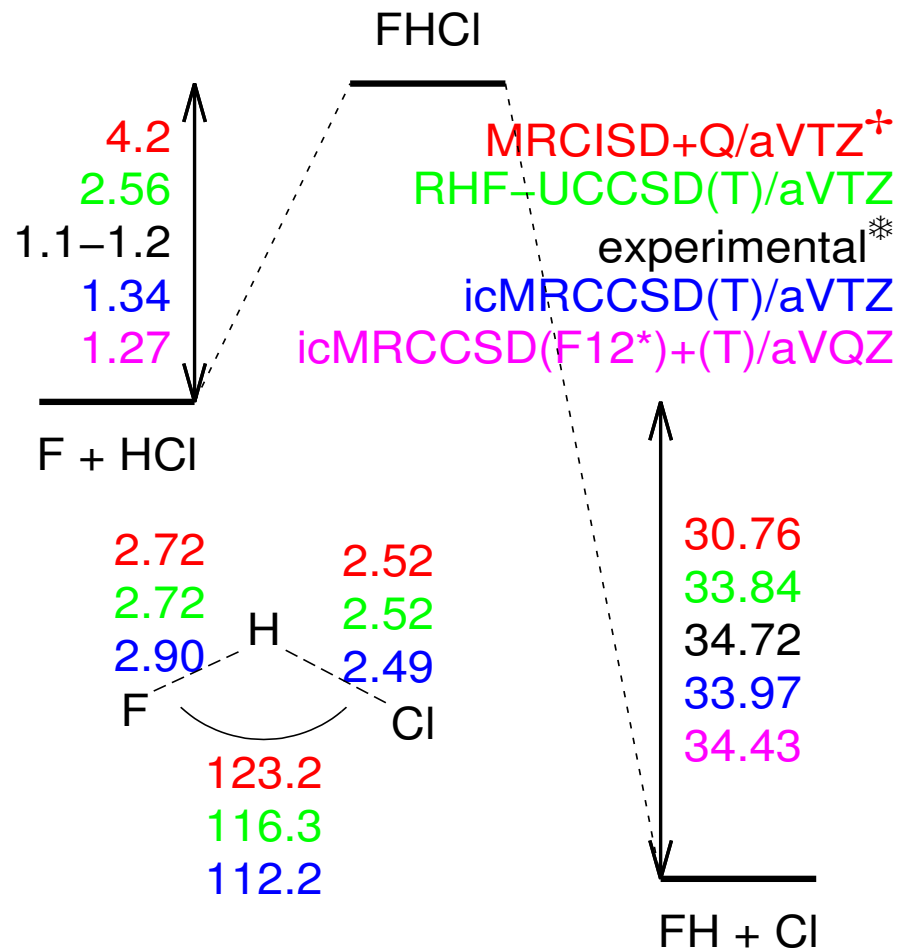
★A. Li, H. Guo, Z. Sun, J. Klos and M. H. Alexander, *Phys. Chem. Chem. Phys.* **15**, 15347 (2013).

*N. Jonathan *et al.*, *Chem. Phys. Lett.* **7**, 257 (1970); A. M. G. Ding *et al.*, *Faraday Discuss. Chem. Soc.* **55**, 252 (1973).

The F + HCl → Cl + HF reaction

Barrier height and exothermicity obtained with icMRCCSD(T) theory are very close to the experimental values!

CAS(3,3) sufficient!
 [MRCI uses CAS(15,9)]

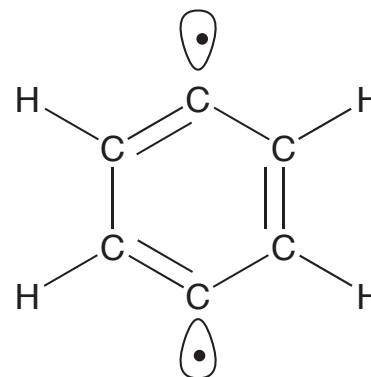


[†]M. P. Deskevich, M. Y. Hayes, K. Takahashi, R. T. Skodje and D. J. Nesbitt, *J. Chem. Phys.* **124**, 223403 (2006).

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Active-space size

Singlet-Triplet splitting of p-benzyne



keep the active space as
small as physically
necessary

$\Delta E_{S-T}/\text{kcal mol}^{-1}$	icMRCCSD	icMRCCSD(T)
CAS(2,2)	3.7	4.8
CAS(8,8)	4.7	4.9

cc-pVDZ basis

Structures: RHF/UHF-CCSD(T)/cc-pVDZ

from: Evangelista et al., JCP 127, 024102 (2007)

Hanauer, Köhn, JCP 136, 204107 (2012)

see also: Evangelista, Hanauer, Köhn, Gauss, JCP 136, 204108 (2012)

Additional aspects (now and future)

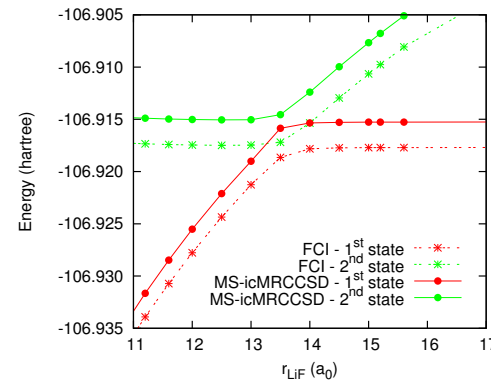
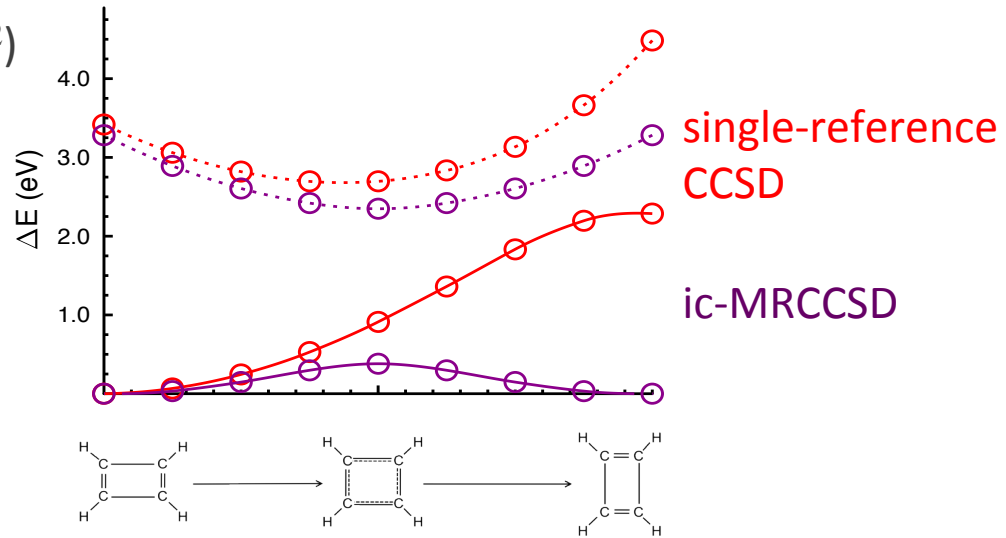
Full spin adaptation (eigenfunctions of S^2)

Excited states via linear-response and/or EOM frameworks

Multi-state extension

Response properties

Lower-order approximations (perturbation theory, CEPA-like theories)



Success and criticism

Full analogy to single-reference CC

Proper formal scaling, compact excitation manifold

Fully extensive energy, orbital invariance

Very accurate results

$$|\Psi_{ic}\rangle = e^{\hat{T}} \sum_{\mu} |\Phi_{\mu}\rangle c_{\mu}$$

Problems, criticism:

Threshold for near-linear dependencies

Commutator approximation

Some consistency problems remain (for dissociating singlets)

Efficiency

