

3D angular momentum and particle number restored calculations with the Gogny EDF

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Outline

- 1. Theoretical framework**
- 2. Applications**
- 3. Conclusions and outlook**

Solving the nuclear many body problem

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1. Theoretical framework

2. Applications

3. Conclusions and outlook

In order to give a proper description of the nuclear system we need:

- ➔ A good **interaction** or energy density functional (EDF) that describes the dynamics of the constituent nucleons.
- ➔ A good **method** -adapted to the corresponding interaction/EDF- for solving the quantum many body problem.

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- Density matrices (normal and abnormal) and spatial densities are written in terms of these auxiliary wave functions.

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- All the observables and variational procedures in EDF methods are expressed in terms of density matrices and spatial densities.

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- Density matrices (normal and abnormal) and spatial densities are written in terms of these auxiliary wave functions.
- All the observables and variational procedures in EDF methods are expressed in terms of density matrices and spatial densities.
- The usual EDF methods consist in determining first the “best” auxiliary many body wave functions and, afterwards, in including more correlations.

Theoretical Framework

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1. Theoretical framework

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- Initial intrinsic states: PN-VAP $\delta E^{N,Z} [\bar{\Phi}(\beta, \gamma)] \Big|_{\bar{\Phi}=\Phi} = 0$

$$E^{N,Z}[\Phi] = \frac{\langle \Phi | \hat{H}_{2b} \hat{P}^N \hat{P}^Z | \Phi \rangle}{\langle \Phi | \hat{P}^N \hat{P}^Z | \Phi \rangle} + \varepsilon_{DD}^{N,Z}(\Phi) - \lambda_{q_{20}} \langle \Phi | \hat{Q}_{20} | \Phi \rangle - \lambda_{q_{22}} \langle \Phi | \hat{Q}_{22} | \Phi \rangle$$

- Intermediate Particle Number and Angular Momentum Projected states

$$|IMK; NZ; \beta\gamma\rangle = \frac{2I+1}{8\pi^2} \int \mathcal{D}_{MK}^{I*}(\Omega) \hat{R}(\Omega) \hat{P}^N \hat{P}^Z |\Phi(\beta, \gamma)\rangle d\Omega$$

- Final GCM states $|IM; NZ\sigma\rangle = \sum_{K\beta\gamma} f_{K\beta\gamma}^{I;NZ,\sigma} |IMK; NZ; \beta\gamma\rangle$

$$\sum_{K'\beta'\gamma'} \left(\mathcal{H}_{K\beta\gamma K'\beta'\gamma'}^{I;NZ} - E^{I;NZ;\sigma} \mathcal{N}_{K\beta\gamma K'\beta'\gamma'}^{I;NZ} \right) f_{K'\beta'\gamma'}^{I;NZ;\sigma} = 0$$

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Recent 3D AMP implementations

- Int Skyrme: M. Bender, P.-H. Heenen, Phys. Rev. C 78, 024309 (2008)
 - states - Particle number and angular momentum restoration of intrinsic LN states.
 - Relativistic: J.M. Yao et al., Phys. Rev. C 81, 044311 (2010)
 - Angular momentum restoration of intrinsic HFB states.
- Fin Gogny: T.R.R., J.L. Egido, Phys. Rev C 81, 064323 (2010)
 - Particle number and angular momentum restoration of PN-VAP states.

$$\sum_{K'\beta'\gamma'} \left(\mathcal{H}_{K\beta\gamma K'\beta'\gamma'}^{I;NZ} - E^{I;NZ;\sigma} \mathcal{N}_{K\beta\gamma K'\beta'\gamma'}^{I;NZ} \right) f_{K'\beta'\gamma'}^{I;NZ;\sigma} = 0$$

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Effective nucleon-nucleon interaction: Gogny force (DIS) that is able to describe properly many phenomena along the whole nuclear chart.

$$\begin{aligned}
 V(1, 2) = & \sum_{i=1}^2 e^{-(\vec{r}_1 - \vec{r}_2)^2 / \mu_i^2} (W_i + B_i P^\sigma - H_i P^\tau - M_i P^\sigma P^\tau) \\
 & + iW_0 (\sigma_1 + \sigma_2) \vec{k} \times \delta(\vec{r}_1 - \vec{r}_2) \vec{k} + V_{\text{Coulomb}}(\vec{r}_1, \vec{r}_2) \\
 & + t_3 (1 + x_0 P^\sigma) \delta(\vec{r}_1 - \vec{r}_2) \rho^\alpha ((\vec{r}_1 + \vec{r}_2) / 2)
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2-body potential

$$+ t_3 (1 + x_0 P^\sigma) \delta(\vec{r}_1 - \vec{r}_2) \rho^\alpha ((\vec{r}_1 + \vec{r}_2) / 2)$$

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Density dependent term

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Density dependent term


Prescriptions:

$$\varepsilon_{DD}^{N,Z}(\Phi)$$

$$\rho_H^{NZ}(\vec{r}) \equiv \frac{\langle \Phi | \hat{\rho}(\vec{r}) P^N P^Z | \Phi \rangle}{\langle \Phi | P^N P^Z | \Phi \rangle}$$

$$\varepsilon_{DD}^{IKK';NZ}[\Phi(\beta, \gamma), \Phi'(\beta', \gamma')]$$

$$\rho_H^{NZ}(\Omega, \vec{r}) \equiv \frac{\langle \Phi | \hat{\rho}(\vec{r}) \hat{R}(\Omega) P^N P^Z | \Phi' \rangle}{\langle \Phi | \hat{R}(\Omega) P^N P^Z | \Phi' \rangle}$$

Theoretical Framework

Axial calculations ^{24}Mg

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First step: Particle Number Projection
 (before the variation) of HFB-type wave functions.

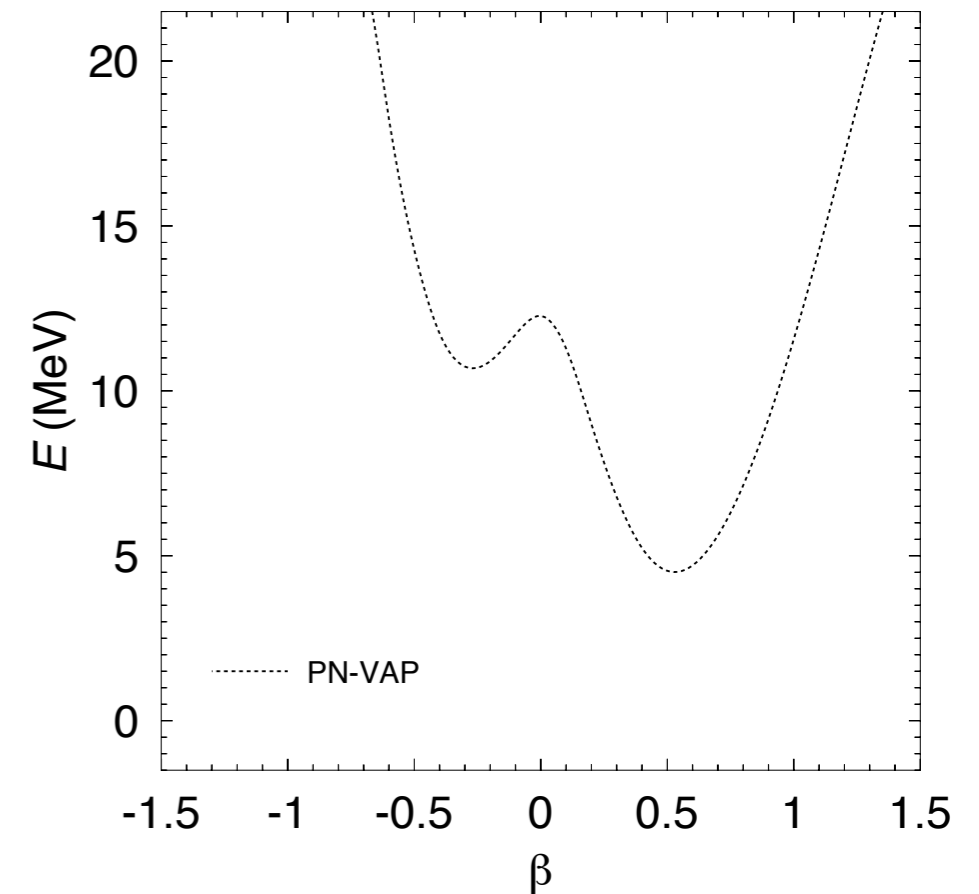
$|\Phi(\beta)\rangle$ product-type many body wave function



$$|\Phi^{N,Z}(\beta)\rangle = \hat{P}^N \hat{P}^Z |\Phi(\beta)\rangle$$



$$E^{N,Z}(\beta) \text{ Potential Energy Surface}$$



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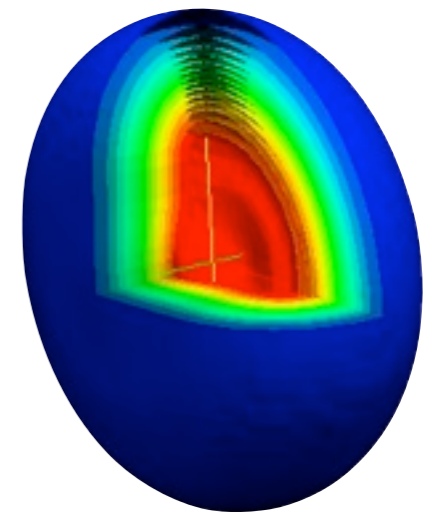
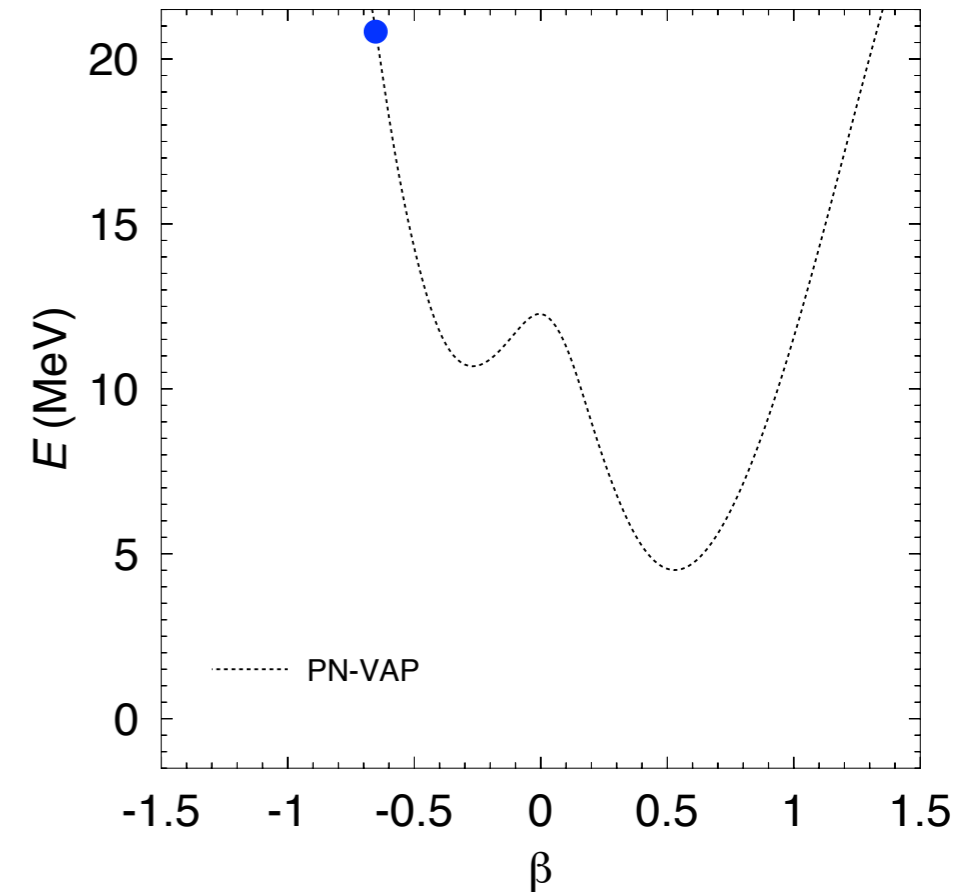
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Second step: Simultaneous Particle Number and Angular Momentum Projection

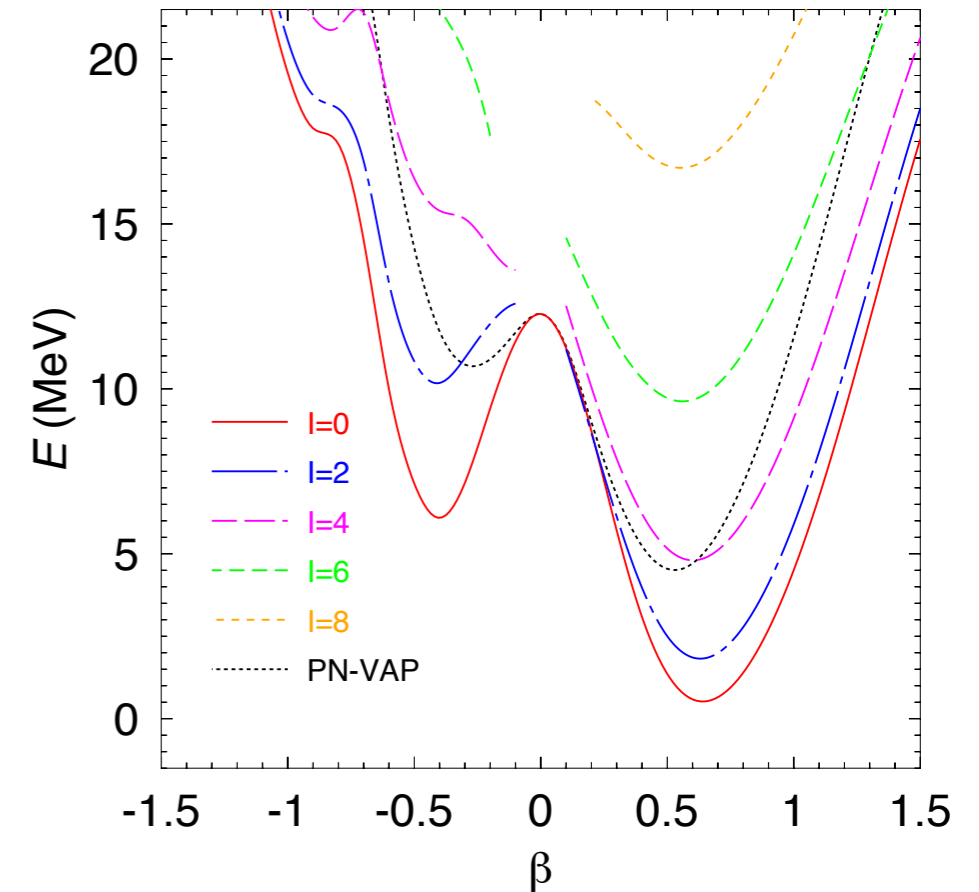
$|\Phi(\beta)\rangle$ product-type many body wave function



$$|\Phi^{I,M;N,Z}(\beta)\rangle = \hat{P}_{00}^I \hat{P}^N \hat{P}^Z |\Phi(\beta)\rangle$$



$E^{I,N,Z}(\beta)$ Potential Energy Surface



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Axial calculations ^{24}Mg

Third step: Configuration mixing within the framework of the **Generator Coordinate Method (GCM)**.

$$|\Psi^{I;N,Z;\sigma}\rangle = \sum_{\beta} f^{I;N,Z;\sigma}(\beta) \hat{P}_{00}^I \hat{P}^N \hat{P}^Z |\Phi(\beta)\rangle$$

Configuration mixing

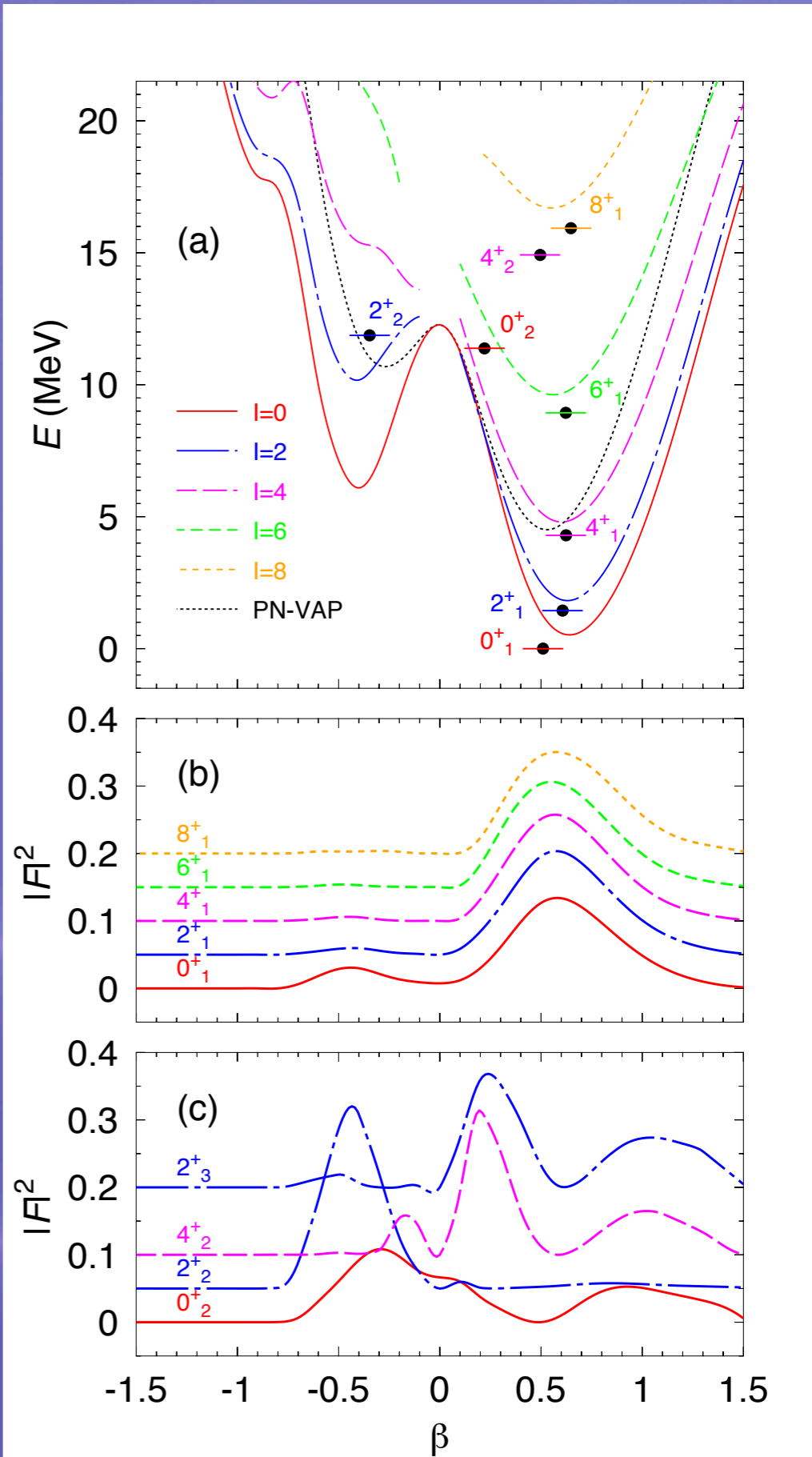
$$\sum_{\beta'} \left(\mathcal{H}_{\beta\beta'}^{I;NZ} - E^{I;NZ;\sigma} \mathcal{N}_{\beta\beta'}^{I;NZ} \right) f_{\beta'}^{I;NZ;\sigma} = 0$$

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Hill-Wheeler-Griffin equations

- Energy spectrum
- Observables (mass, radius, B(E2), etc.)
- "Collective w.f."



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Theoretical Framework

Triaxiality in nuclei

- Gamma bands and gamma softness in the low energy spectra
- Shape coexistence and/or shape transitions in transitional regions
- Impact of triaxiality in the mass
- Impact of triaxiality in the fission barriers
- Triaxiality at high spin
- ...

Theoretical Framework

First step: Particle Number Projection (before the variation) of HFB-type wave functions.

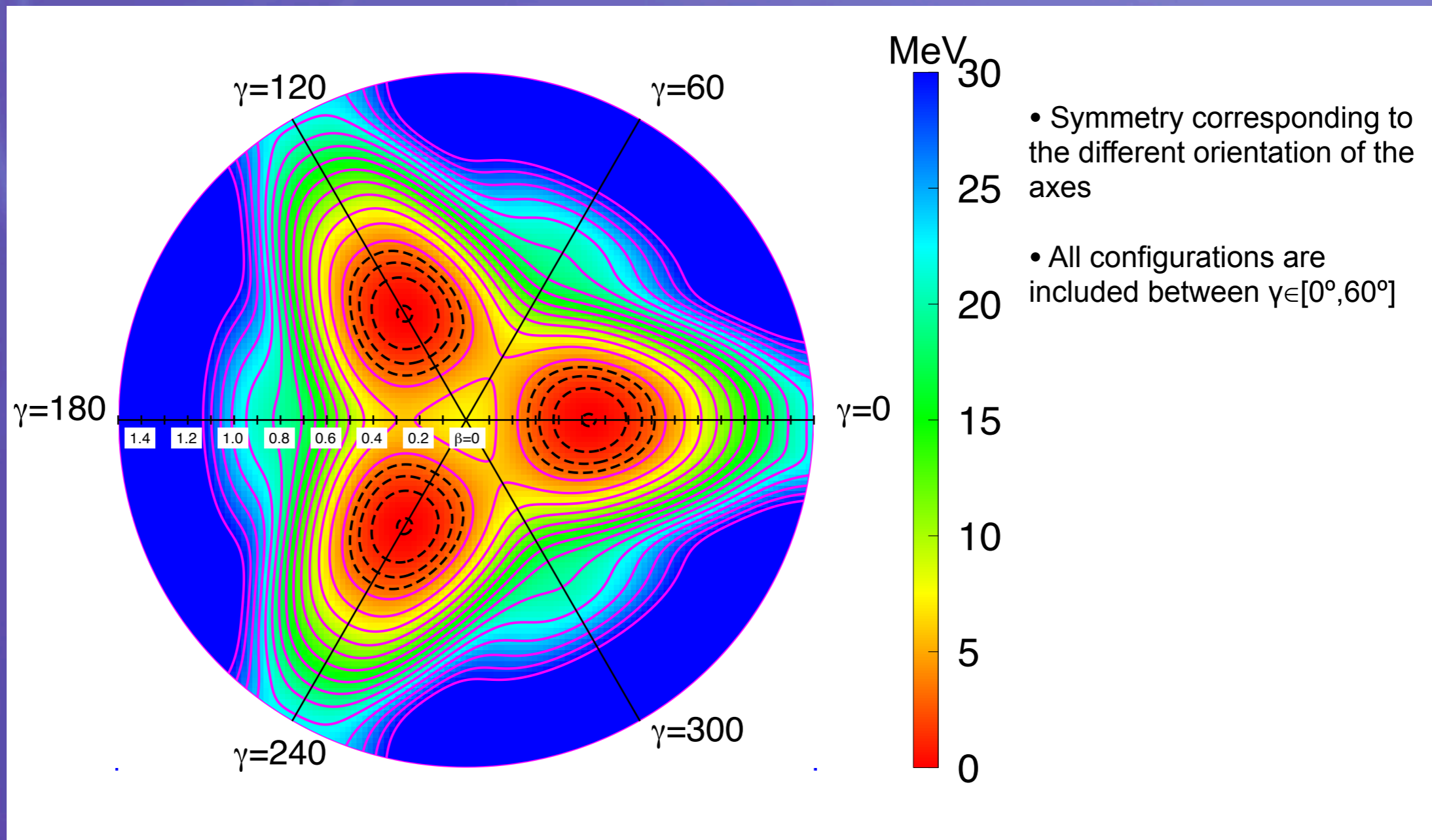
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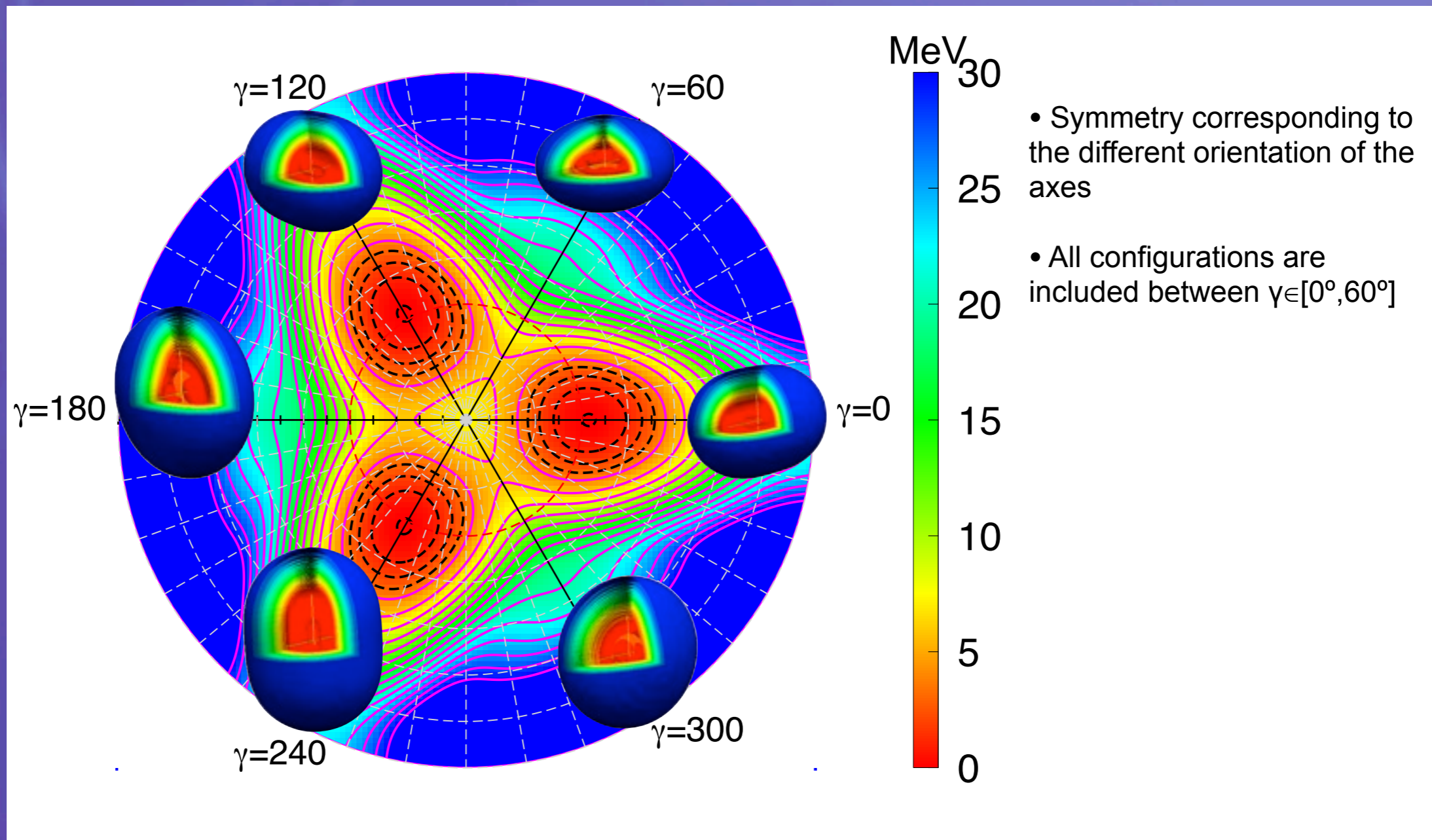
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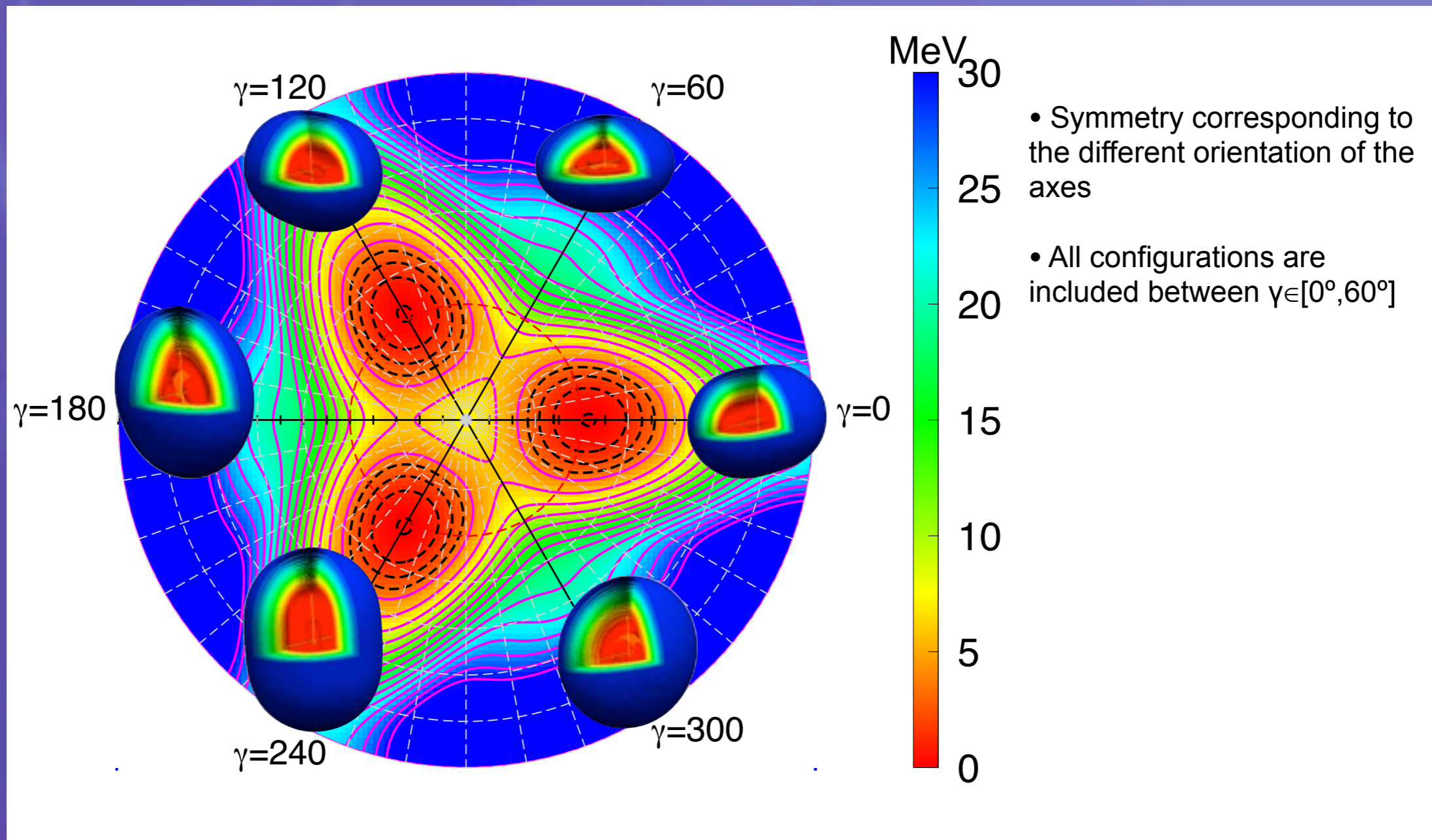
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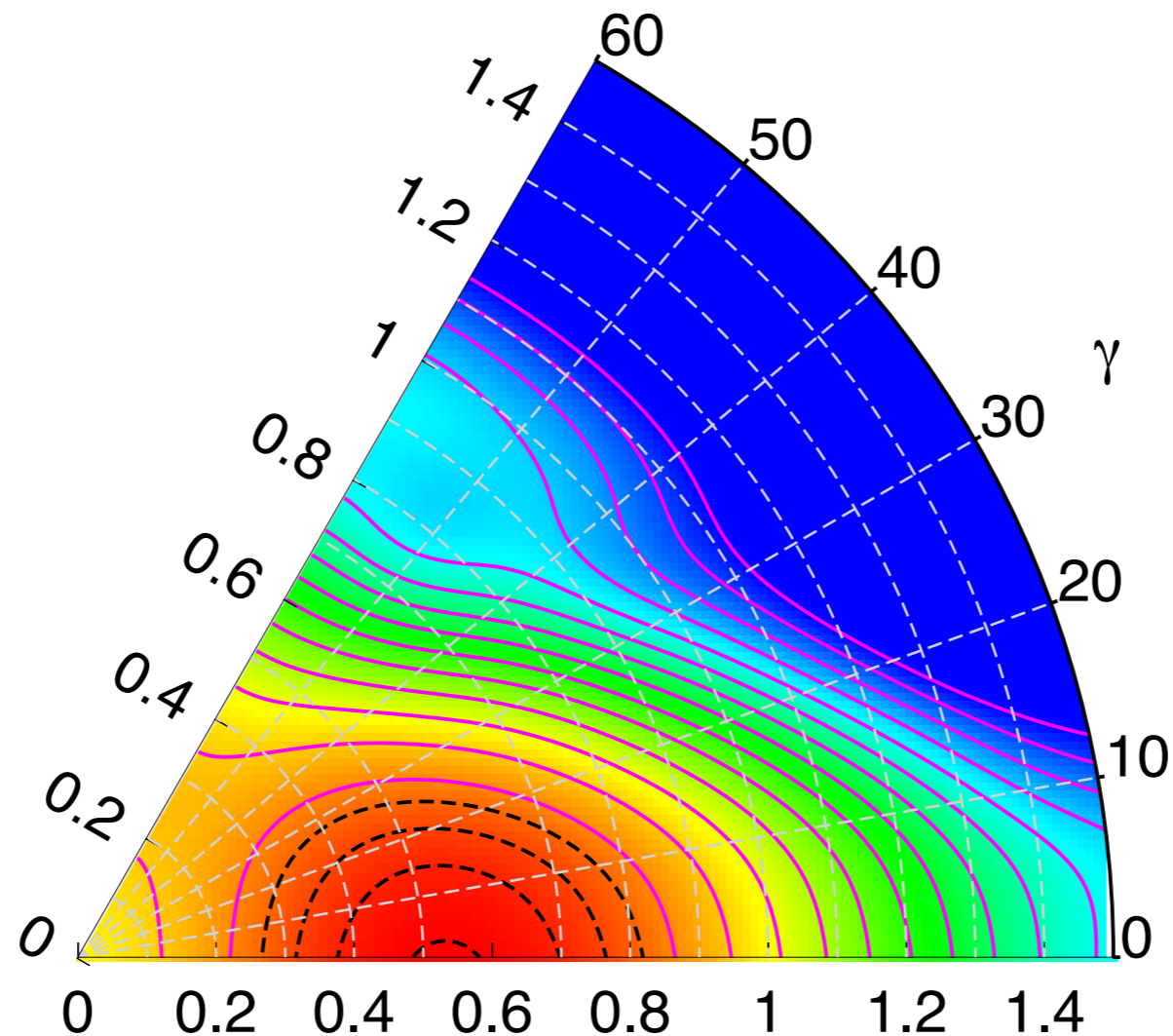
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- Symmetry corresponding to the different orientation of the axes
- All configurations are included between $\gamma \in [0^\circ, 60^\circ]$

Theoretical Framework

Second step: Simultaneous Particle Number and Angular Momentum Projection

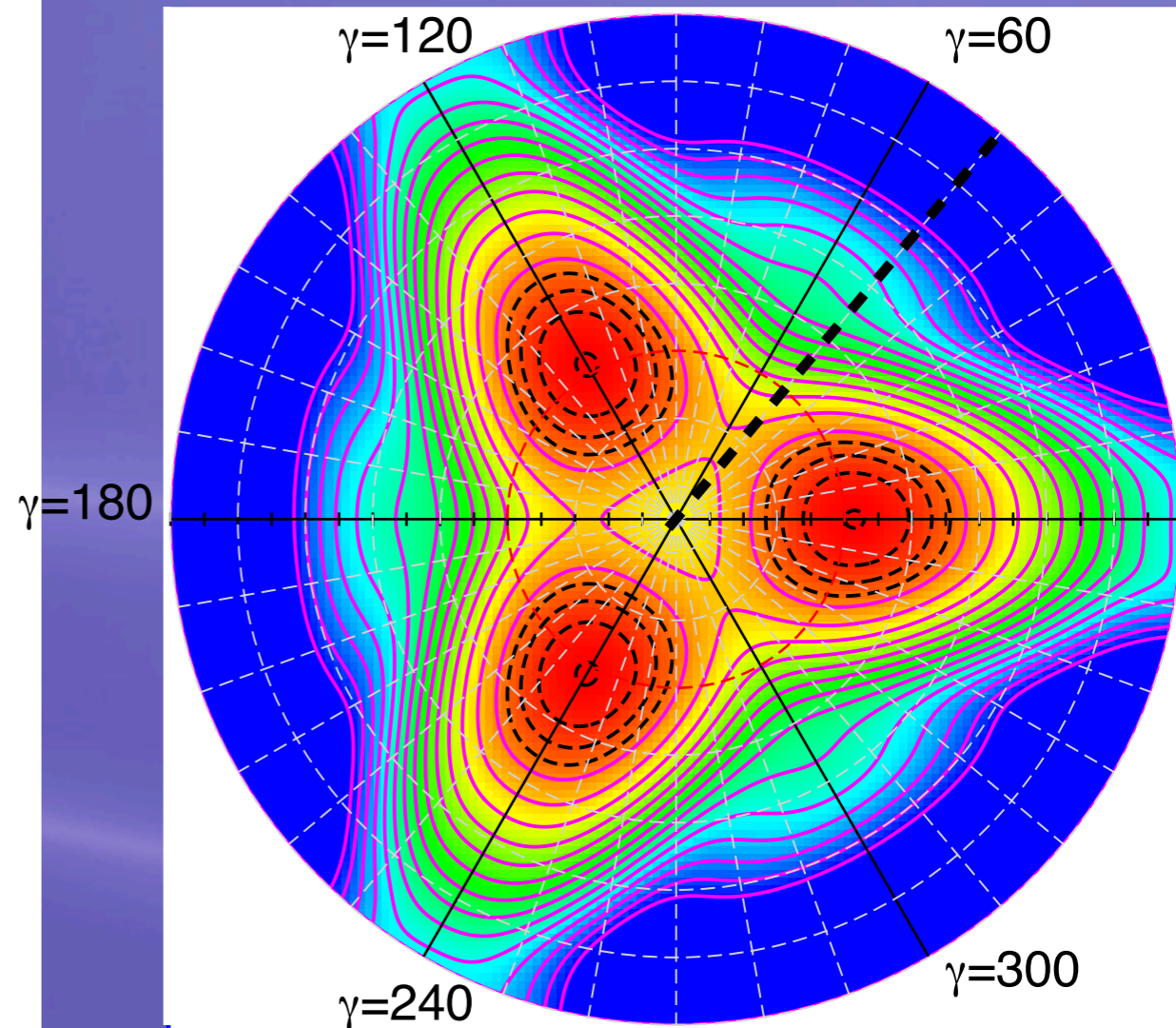
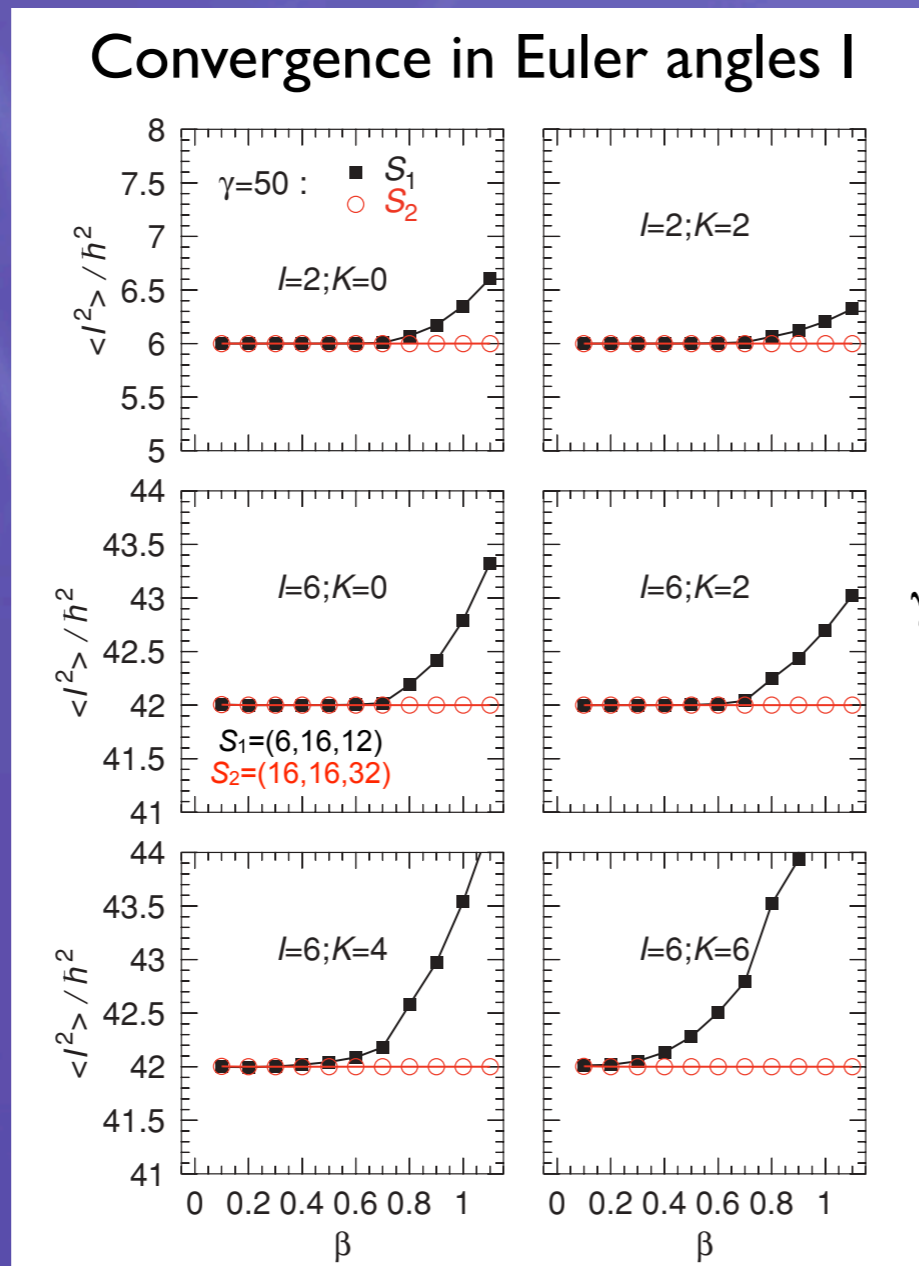
$$|IMK; NZ; \beta\gamma\rangle = \frac{2I+1}{8\pi^2} \int \mathcal{D}_{MK}^{I*}(\Omega) \hat{R}(\Omega) \hat{P}^N \hat{P}^Z |\Phi(\beta, \gamma)\rangle d\Omega$$

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Theoretical Framework

Second step: Simultaneous Particle Number and Angular Momentum Projection

$$|IMK; NZ; \beta\gamma\rangle = \frac{2I+1}{8\pi^2} \int \mathcal{D}_{MK}^{I*}(\Omega) \hat{R}(\Omega) \hat{P}^N \hat{P}^Z |\Phi(\beta, \gamma)\rangle d\Omega$$

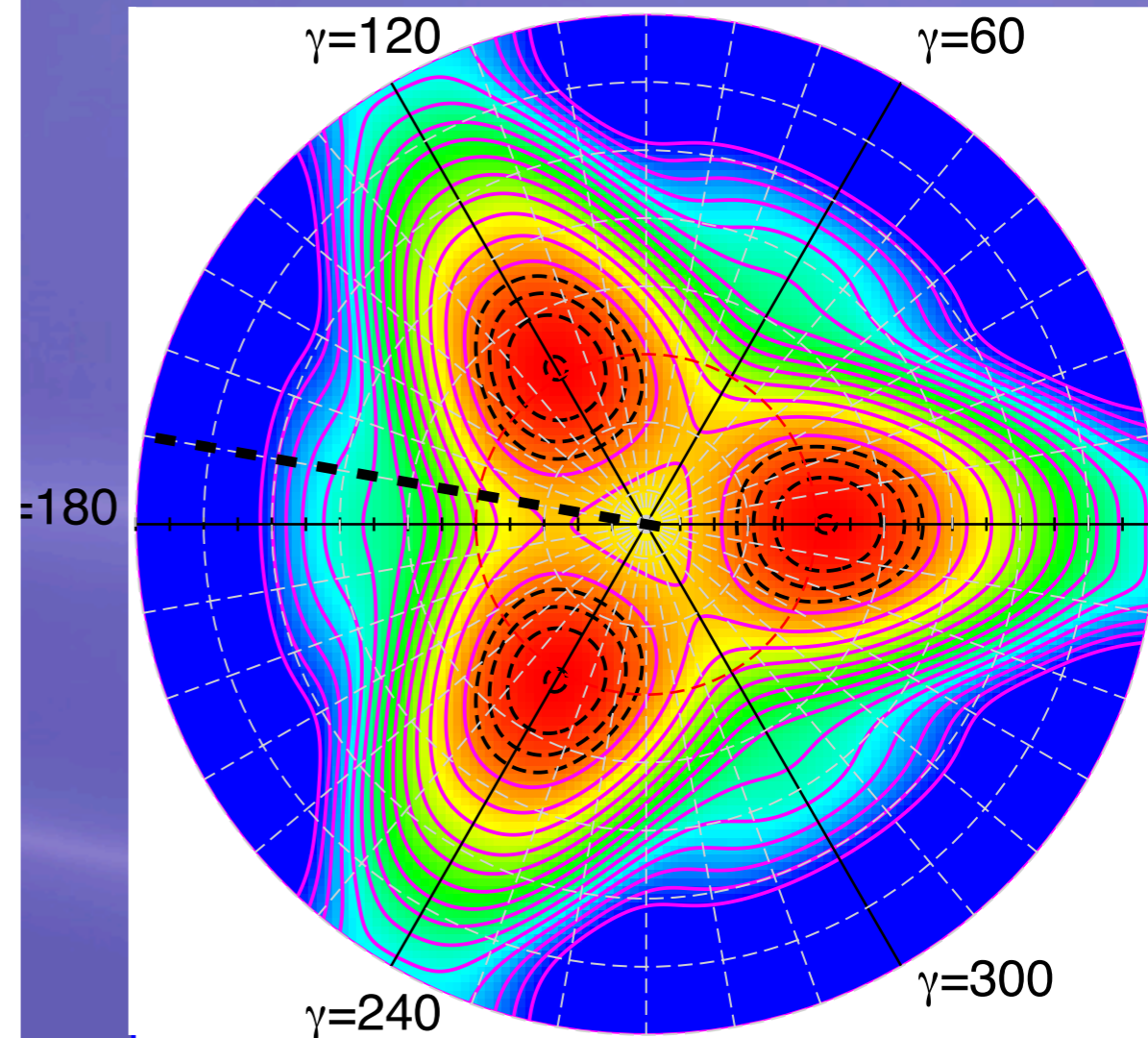
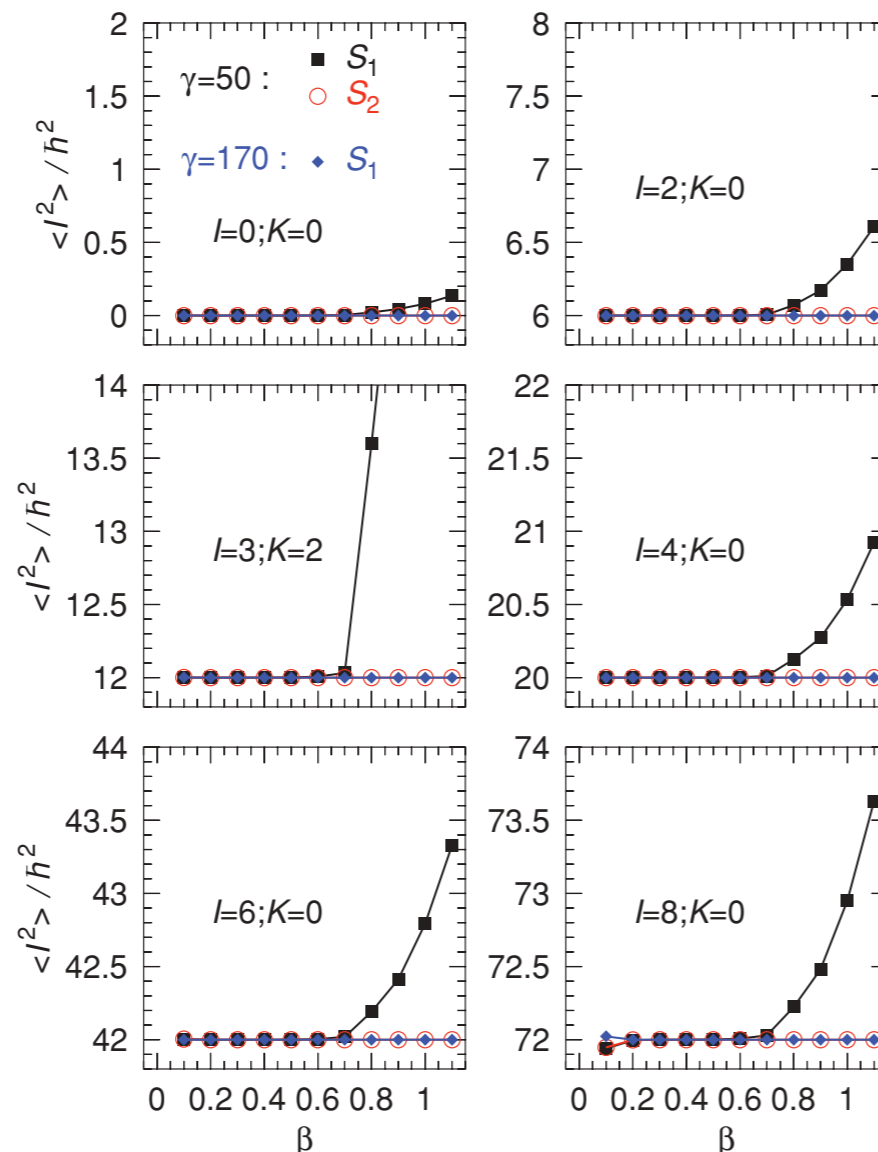
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Convergence in Euler angles II



Theoretical Framework

Second step: Simultaneous Particle Number and Angular Momentum Projection

$$|IMK; NZ; \beta\gamma\rangle = \frac{2I+1}{8\pi^2} \int \mathcal{D}_{MK}^{I*}(\Omega) \hat{R}(\Omega) \hat{P}^N \hat{P}^Z |\Phi(\beta, \gamma)\rangle d\Omega$$

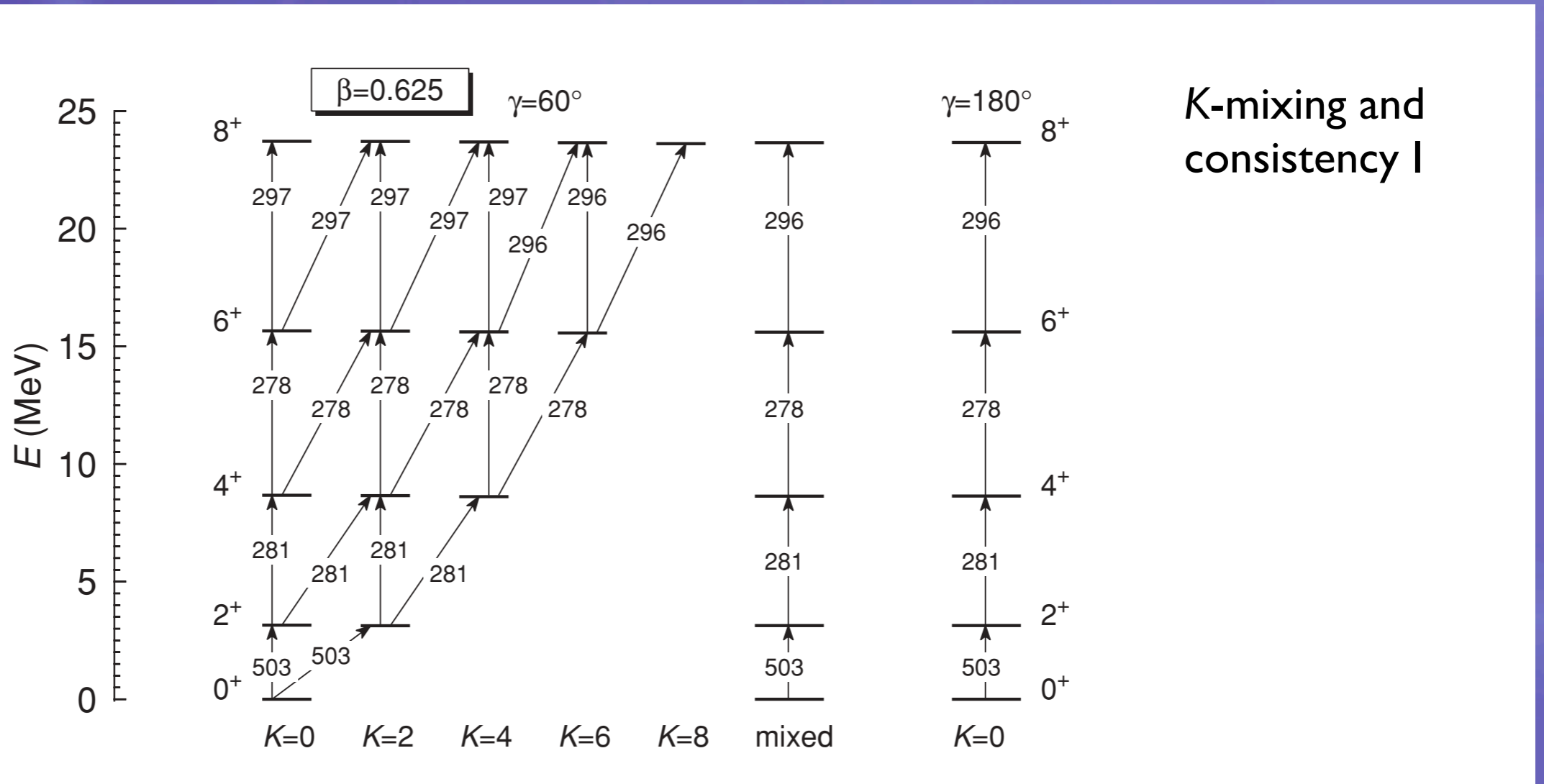
$$|IM; NZ; \beta\gamma\rangle = \sum_K g_K^{IM; NZ; \beta\gamma} |IMK; NZ; \beta\gamma\rangle$$

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1. Theoretical framework

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K-mixing and consistency I

Theoretical Framework

Second step: Simultaneous Particle Number and Angular Momentum Projection

$$|IMK; NZ; \beta\gamma\rangle = \frac{2I+1}{8\pi^2} \int \mathcal{D}_{MK}^{I*}(\Omega) \hat{R}(\Omega) \hat{P}^N \hat{P}^Z |\Phi(\beta, \gamma)\rangle d\Omega$$

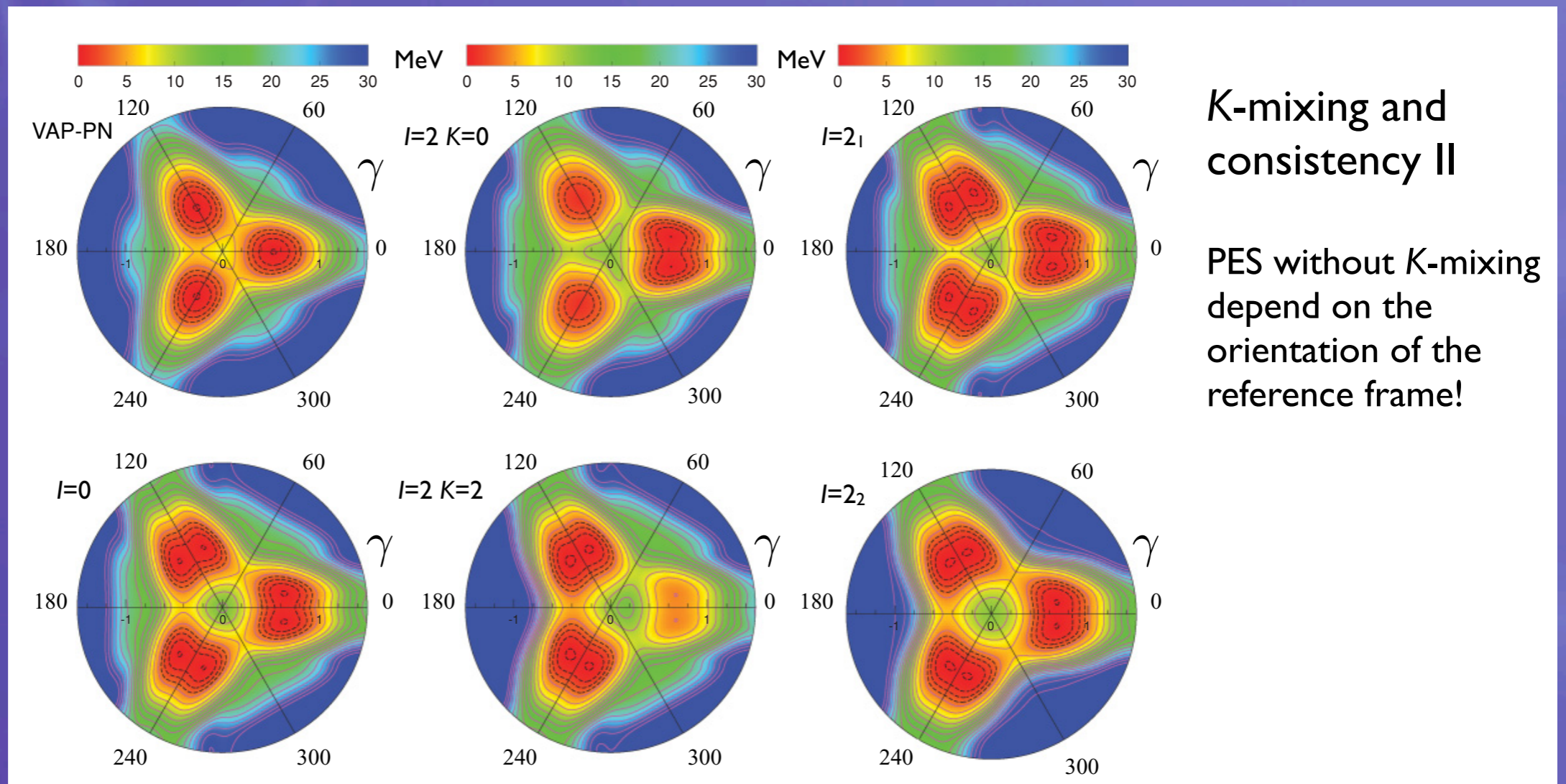
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Theoretical Framework

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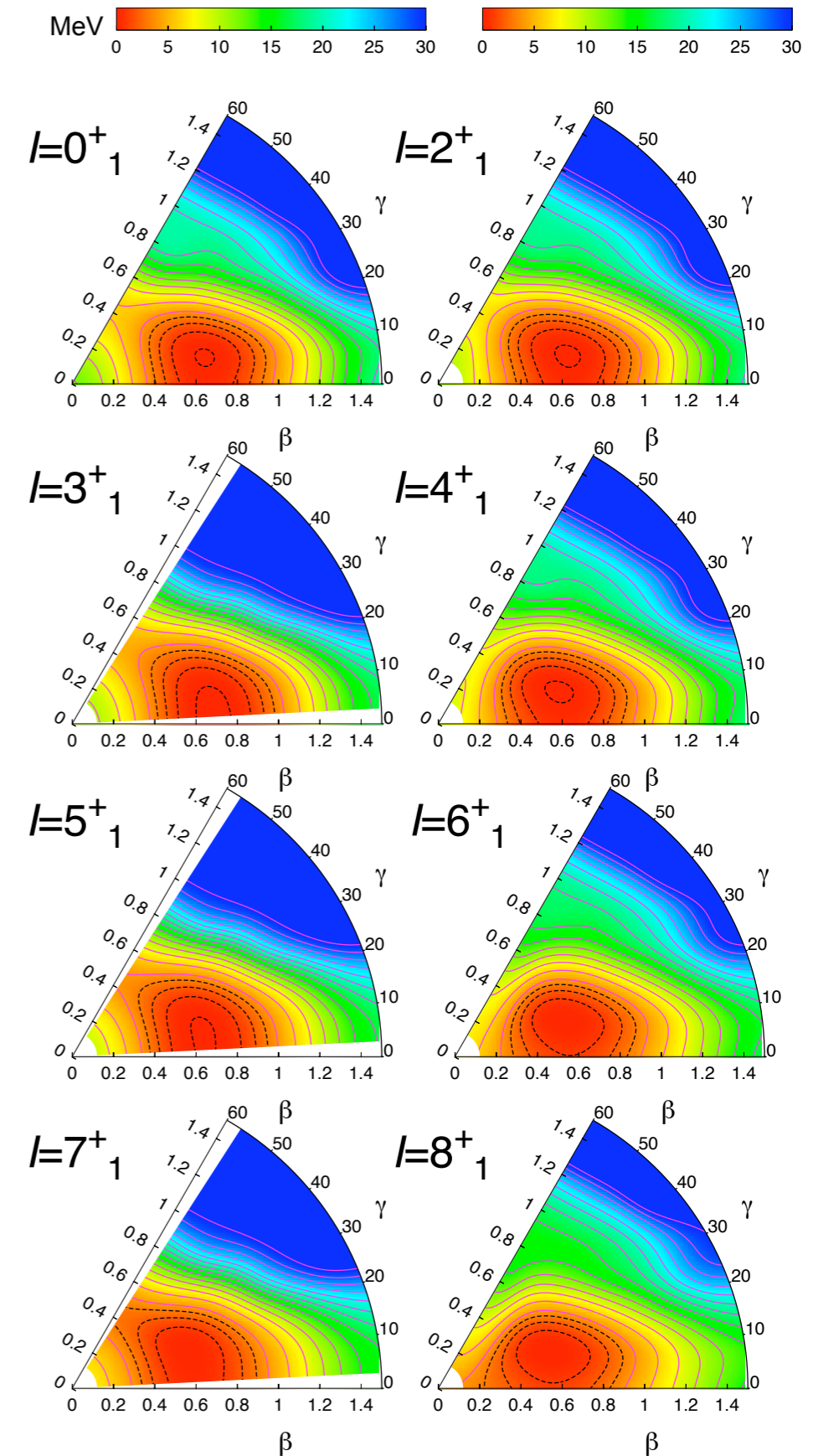
Triaxial calculations ^{24}Mg

Second step: Simultaneous Particle Number and Angular Momentum Projection

$$|IMK; NZ; \beta\gamma\rangle = \frac{2I+1}{8\pi^2} \int \mathcal{D}_{MK}^{I*}(\Omega) \hat{R}(\Omega) \hat{P}^N \hat{P}^Z |\Phi(\beta, \gamma)\rangle d\Omega$$

$$|IM; NZ; \beta\gamma\rangle = \sum_K g_K^{IM; NZ; \beta\gamma} |IMK; NZ; \beta\gamma\rangle$$

- Minimum displaced to triaxial shapes.
- Projection onto odd l angular momentum
- Softening of PES with increasing l .
- Difference between triaxial minimum and axial saddle point of ~ 0.7 MeV (0^+)



Theoretical Framework

Third step: Configuration mixing within the framework of the **Generator Coordinate Method (GCM)**. *K* and deformation mixing

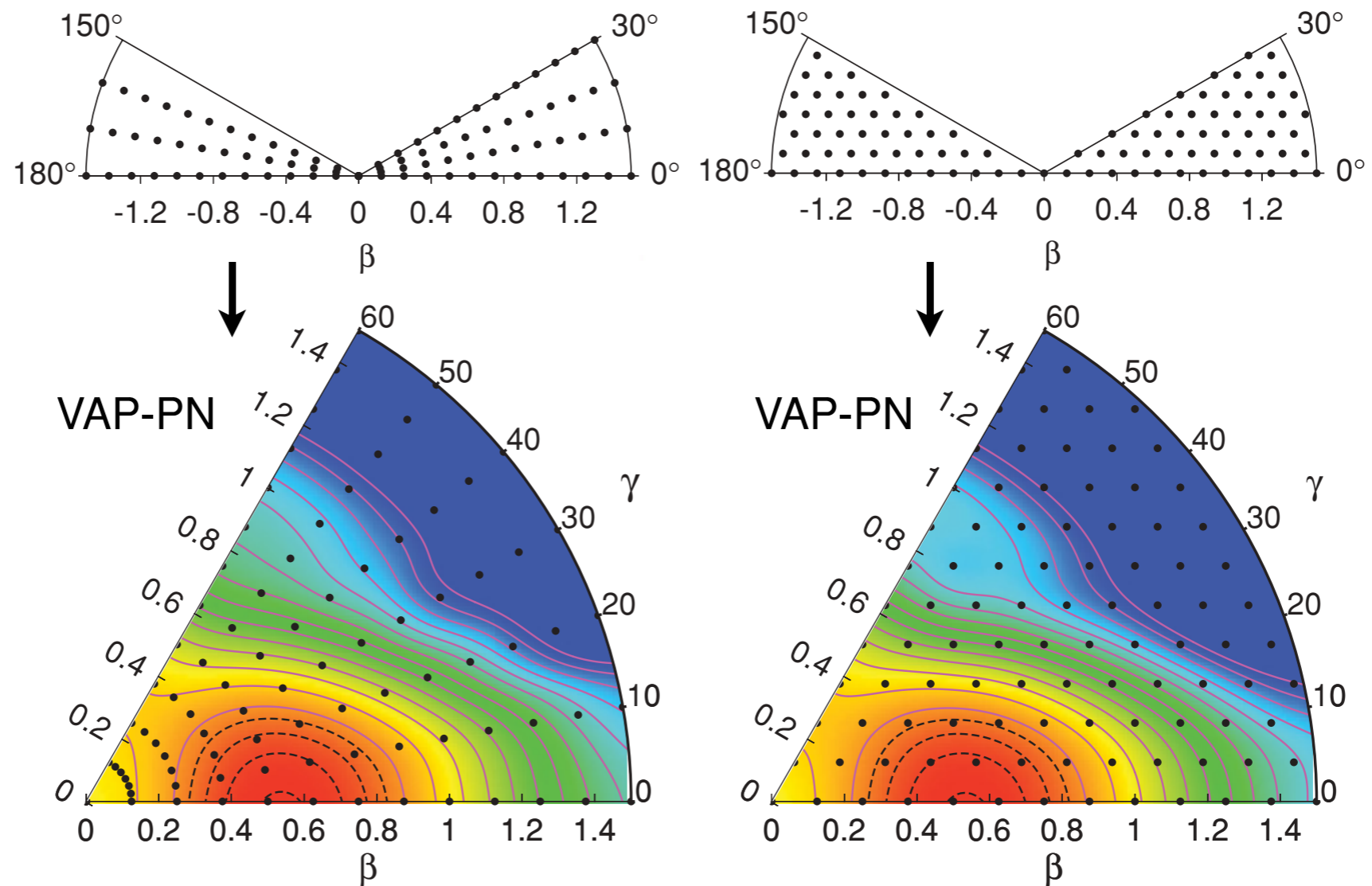
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Selection of the mesh. Resolution.



Theoretical Framework

Third step: Configuration mixing within the framework of the **Generator Coordinate Method (GCM)**. *K* and deformation mixing

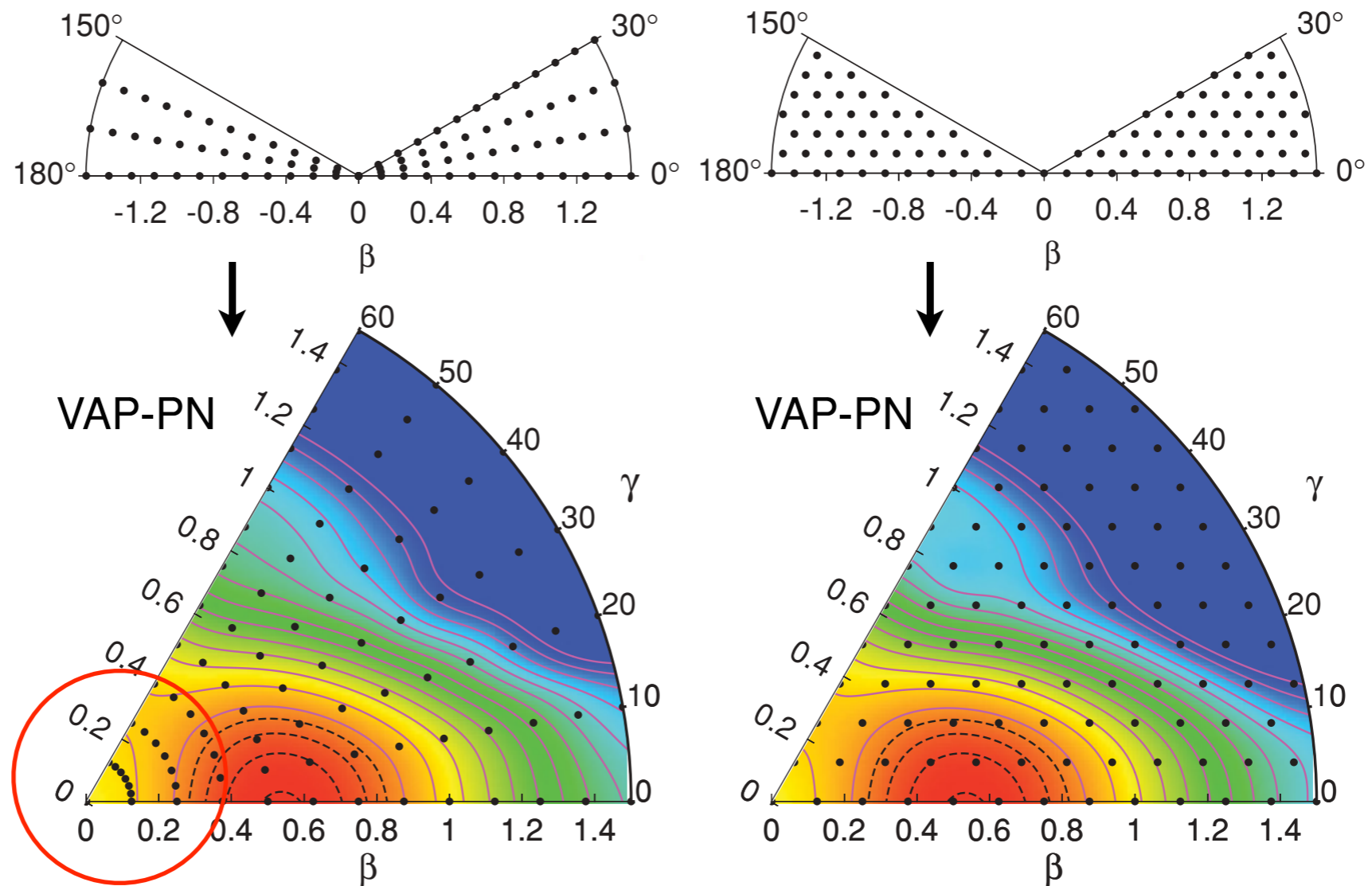
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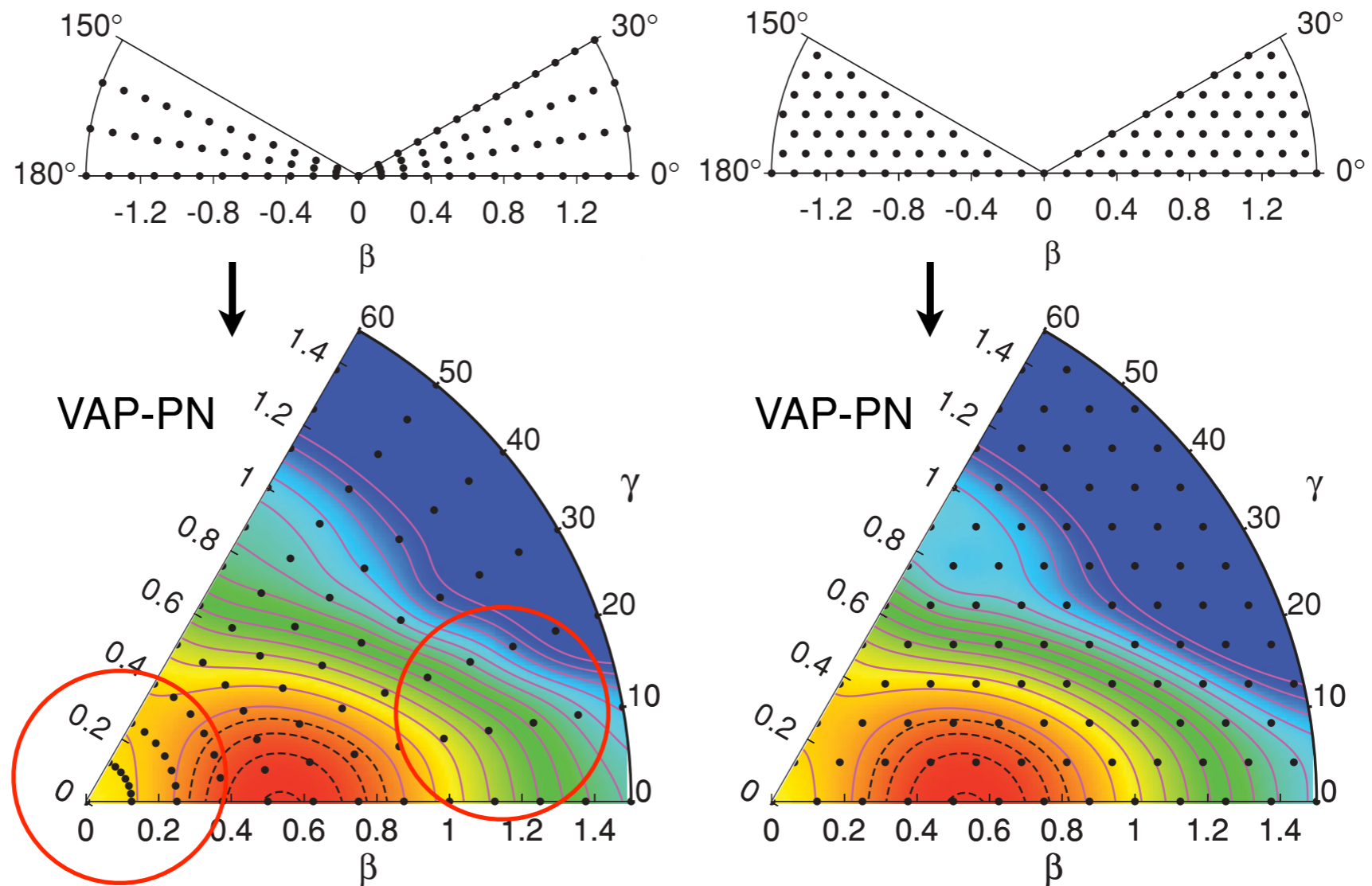
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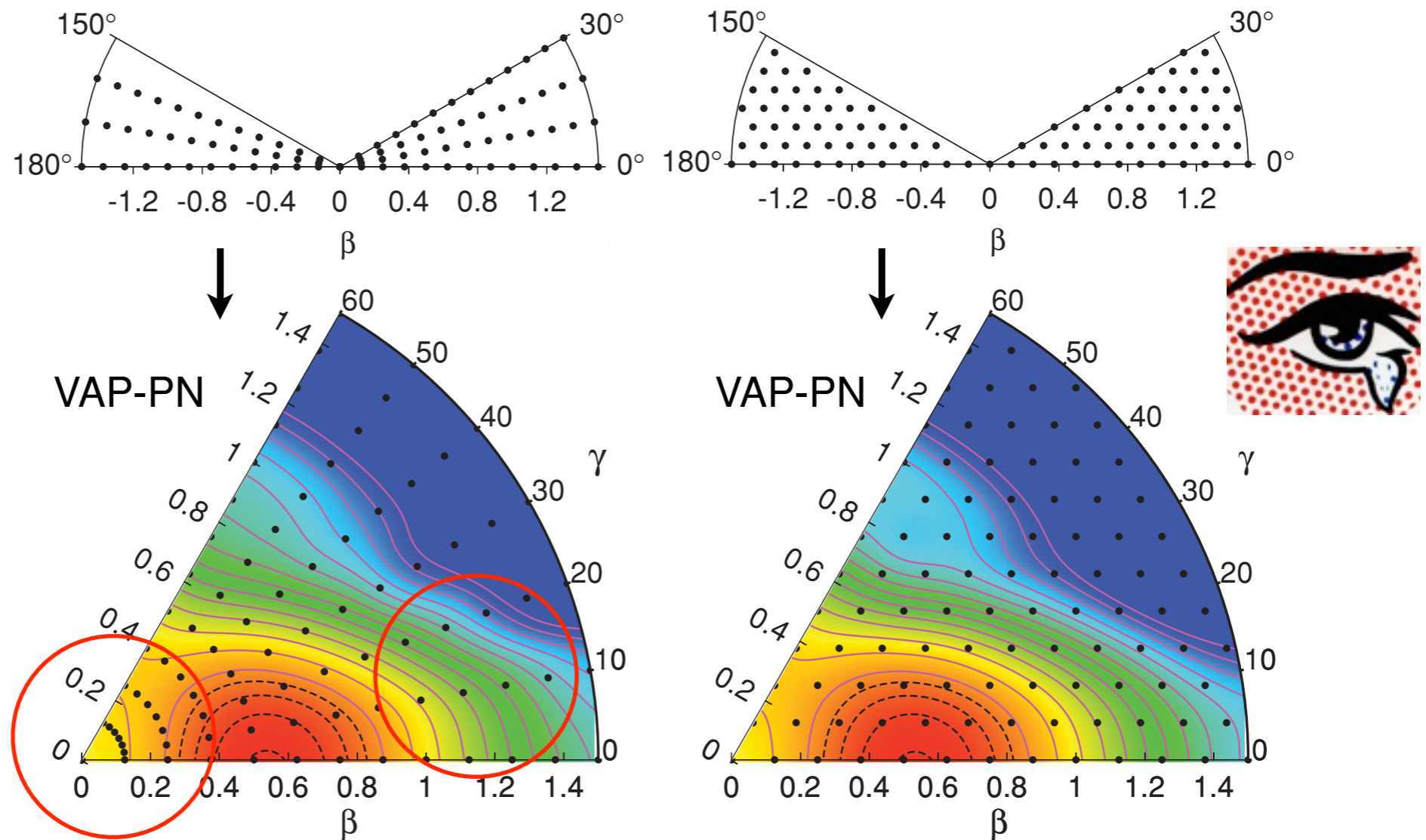
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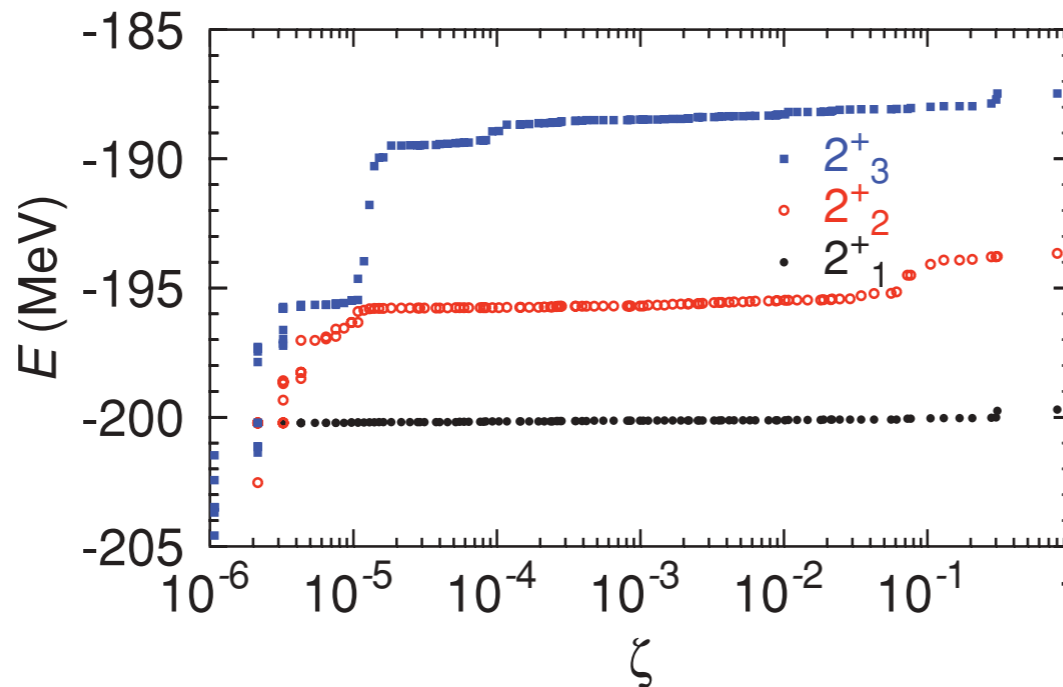
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Third step: Configuration mixing within the framework of the **Generator Coordinate Method (GCM). K and deformation mixing**

$$|IM; NZ\sigma\rangle = \sum_{K\beta\gamma} f_{K\beta\gamma}^{I;NZ,\sigma} |IMK; NZ; \beta\gamma\rangle$$

$$\sum_{K'\beta'\gamma'} \left(\mathcal{H}_{K\beta\gamma K'\beta'\gamma'}^{I;NZ} - E^{I;NZ;\sigma} \mathcal{N}_{K\beta\gamma K'\beta'\gamma'}^{I;NZ} \right) f_{K'\beta'\gamma'}^{I;NZ;\sigma} = 0$$

Convergence of the GCM states



- Plateau condition as a function of natural states.
- Orthogonalization requirements.

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1. Theoretical framework

2. Applications

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Theoretical Framework

Third step: Configuration mixing within the framework of the **Generator Coordinate Method (GCM)**. *K* and deformation mixing

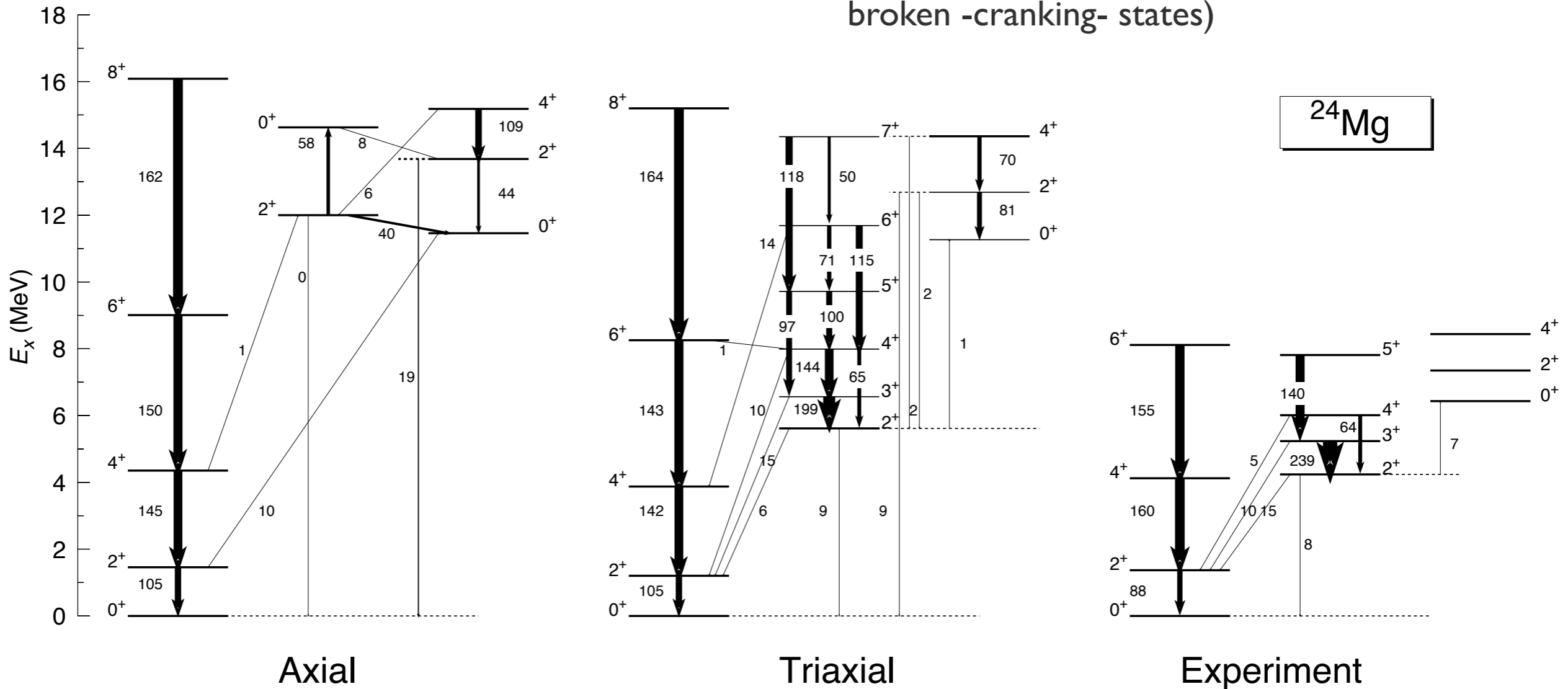
- Axial ground state rotational band well described with axial calculations in this nucleus
- Second band associated to a gamma band
- Overall qualitative agreement between experimental data and triaxial calculations (energies and B(E2))
- Too high energies for the second and third band heads (lack of time reversal symmetry broken -cranking- states)

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1. Theoretical framework

2. Applications

3. Conclusions



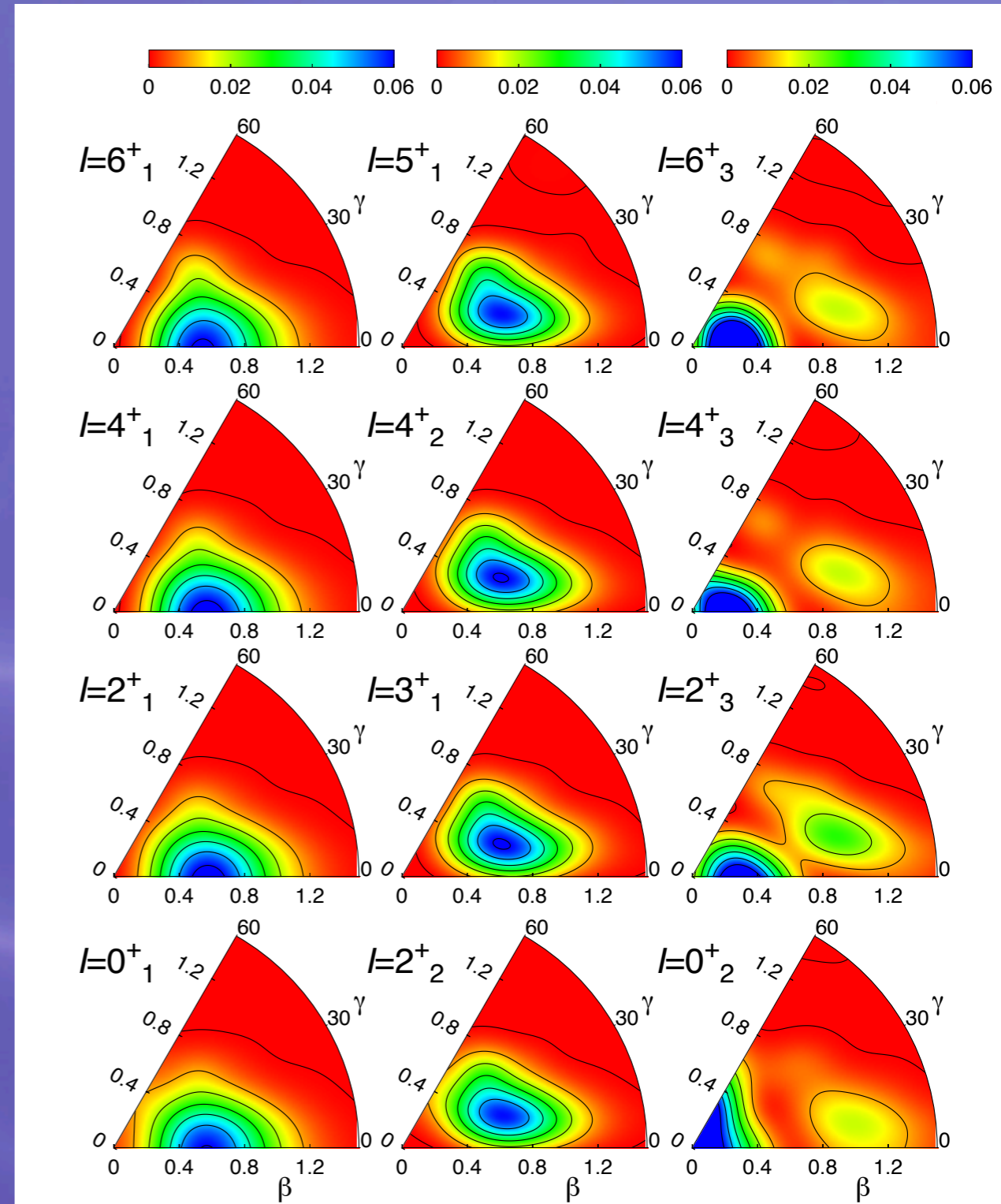
Theoretical Framework

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- Axial ground state rotational band
- Second band associated to a gamma band
- Third band with shape mixing



Theoretical Framework

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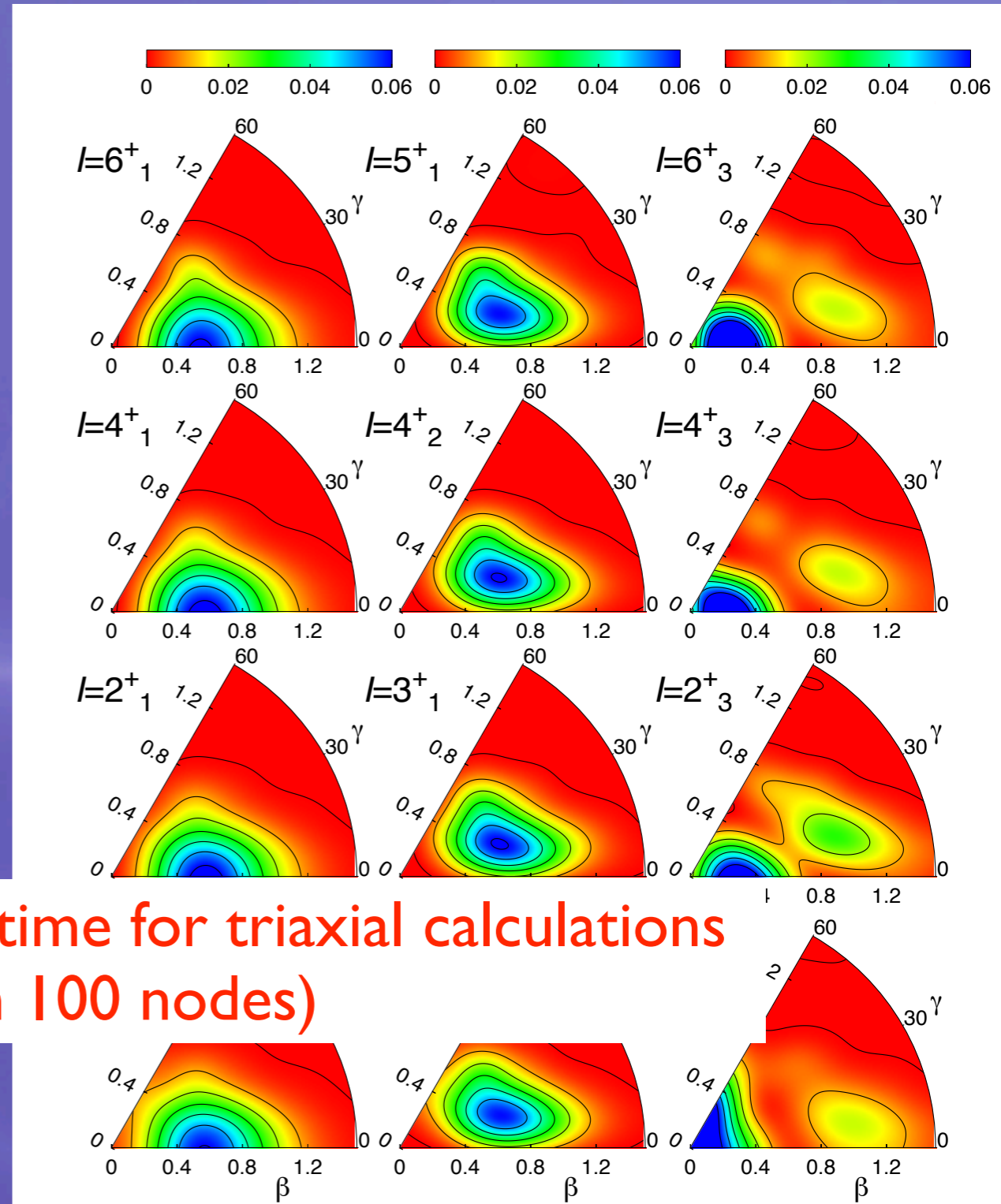
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- Axial ground state rotational band

Problem: Computational time for triaxial calculations
 (~1 month per nucleus in 100 nodes)

- Third band with shape mixing



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1. Theoretical framework

2. Applications

3. Conclusions and outlook

Theoretical Framework

Recipe for cooking triaxial calculations

- Check the interval of deformations and the number of oscillator shells with axial calculations.
- Use a high resolution mesh in the triaxial plane.
- Check the number of integration points in the Euler angles studying the corresponding known mean values. Use rotated states to improve the convergence.
- Check the convergence of the final GCM calculations (plateau condition and orthonormalization requirements).

Applications

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1. Theoretical framework

2. Applications

3. Conclusions and outlook

Skyrme: M. Bender, P.-H. Heenen, Phys. Rev. C 78, 024309 (2008)
- Particle number and angular momentum restoration of intrinsic LN states.

Relativistic: J.M. Yao et al., Phys. Rev. C 81, 044311 (2010)
- Angular momentum restoration of intrinsic HFB states.

Gogny: T.R.R., J.L. Egido, Phys. Rev. C 81, 064323 (2010)
- Particle number and angular momentum restoration of PN-VAP states.

^{24}Mg is a bad choice to see triaxial effects!!

Relativistic: J.M. Yao et al., Phys. Rev. C 83, 014308 (2011), Phys. Rev. C 84, 024306 (2011)
Gogny: T.R.R. and J.L.E., Journal of Physics: Conference Series INPC (2010),
Phys. Lett. B submitted, Phys. Rev. C submitted.

The need of triaxiality: ^{126}Xe as an example

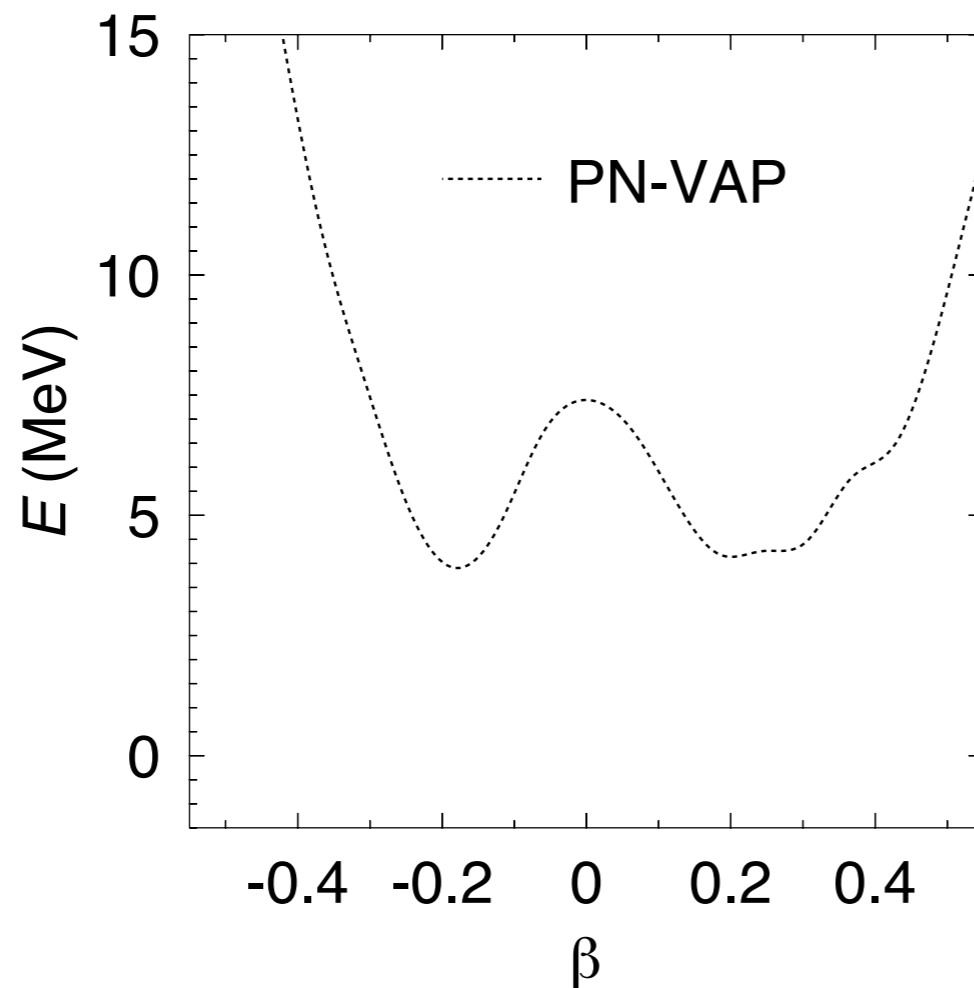
Axial calculations ^{126}Xe

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1. Theoretical framework

2. Applications

3. Conclusions and outlook



✓ AXIAL calculations

✓ Two minima almost degenerated in the potential energy surface

✓ The collective wave function of the ground state is distributed in these two minima (shape coexistence)

✓ TRIAXIAL calculations?

The need of triaxiality: ^{126}Xe as an example

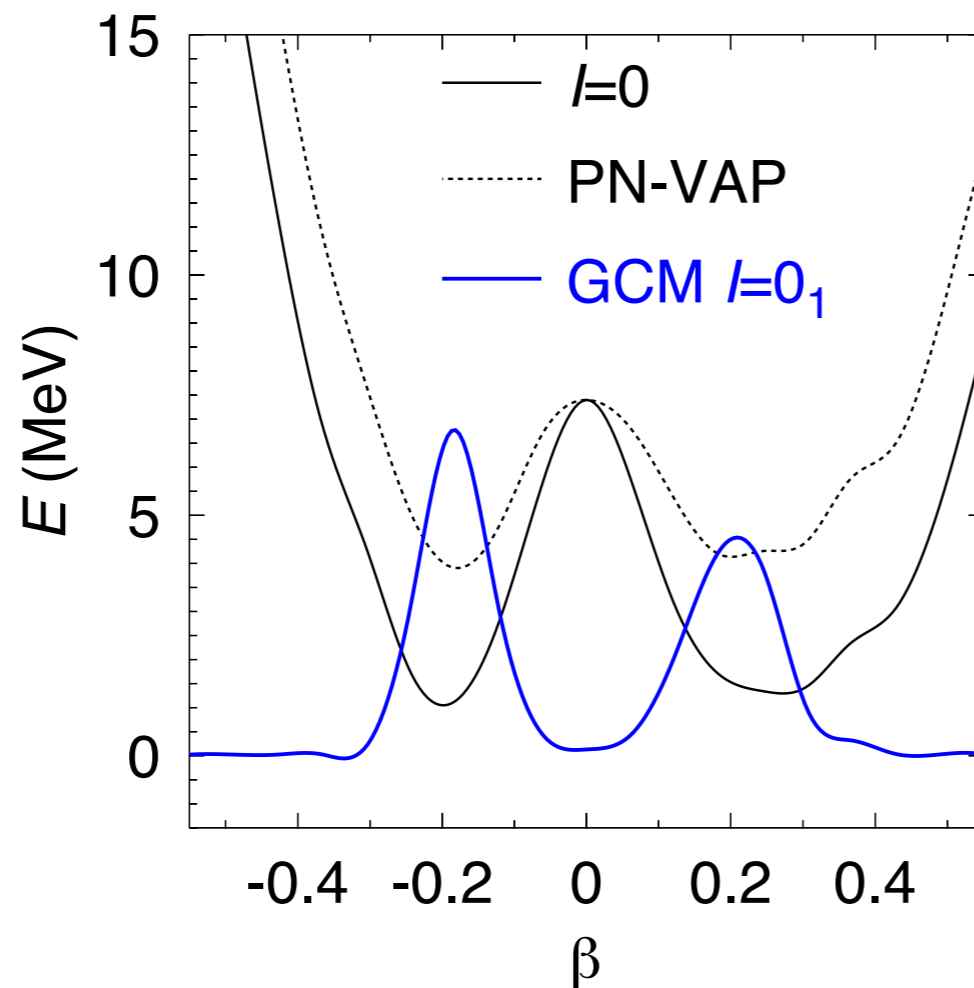
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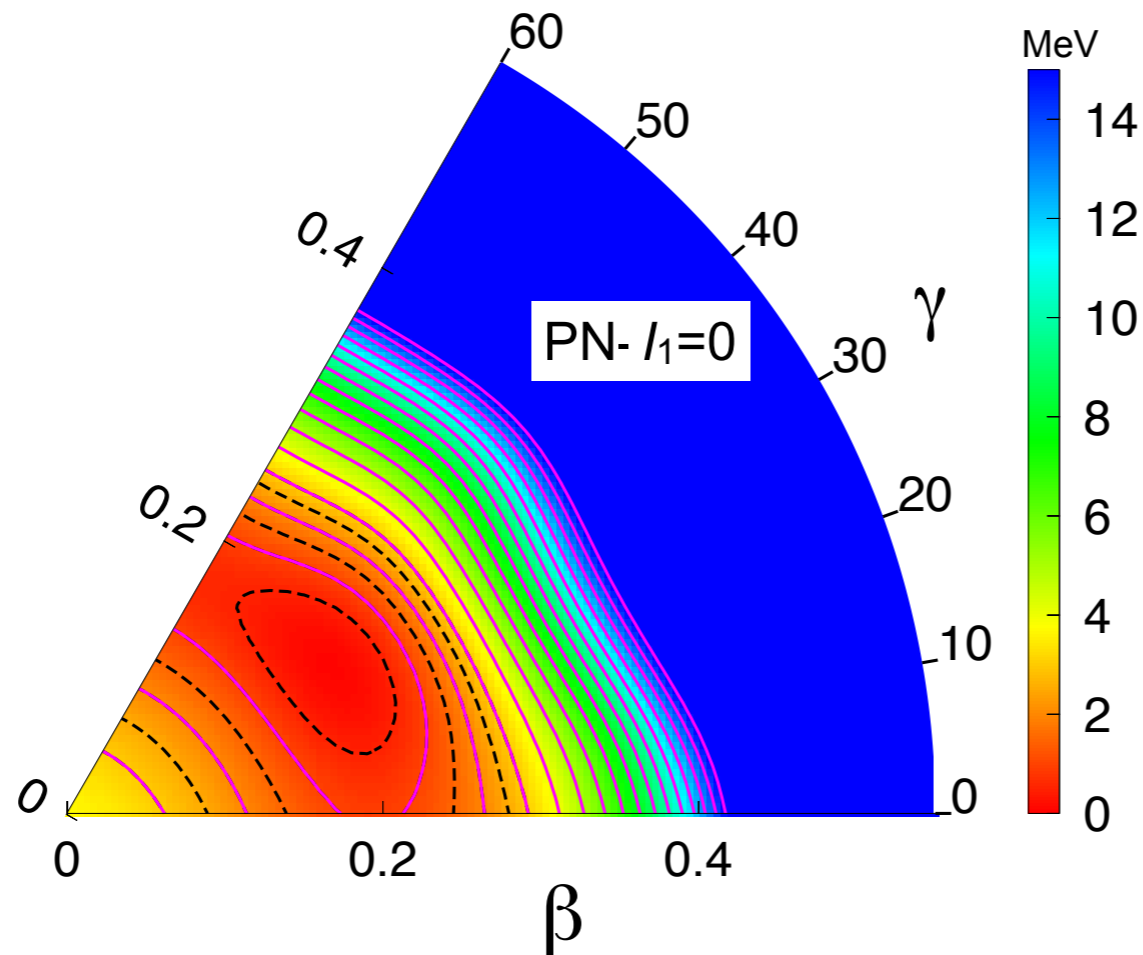
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Triaxial calculations ^{126}Xe in a reduced configuration space (seven shells)

T.R.R and J.L. Egido, Journal of Physics: Conference Series (2011)



✓ TRIAXIAL calculations

✓ One single minimum in $\gamma=30^\circ$ and saddle points in the axial configurations

✓ PES very soft in the γ degree of freedom

✓ After GCM, there is not coexistence of prolate and oblate configurations for the ground state, just a triaxial state.

The need of triaxiality: ^{126}Xe as an example

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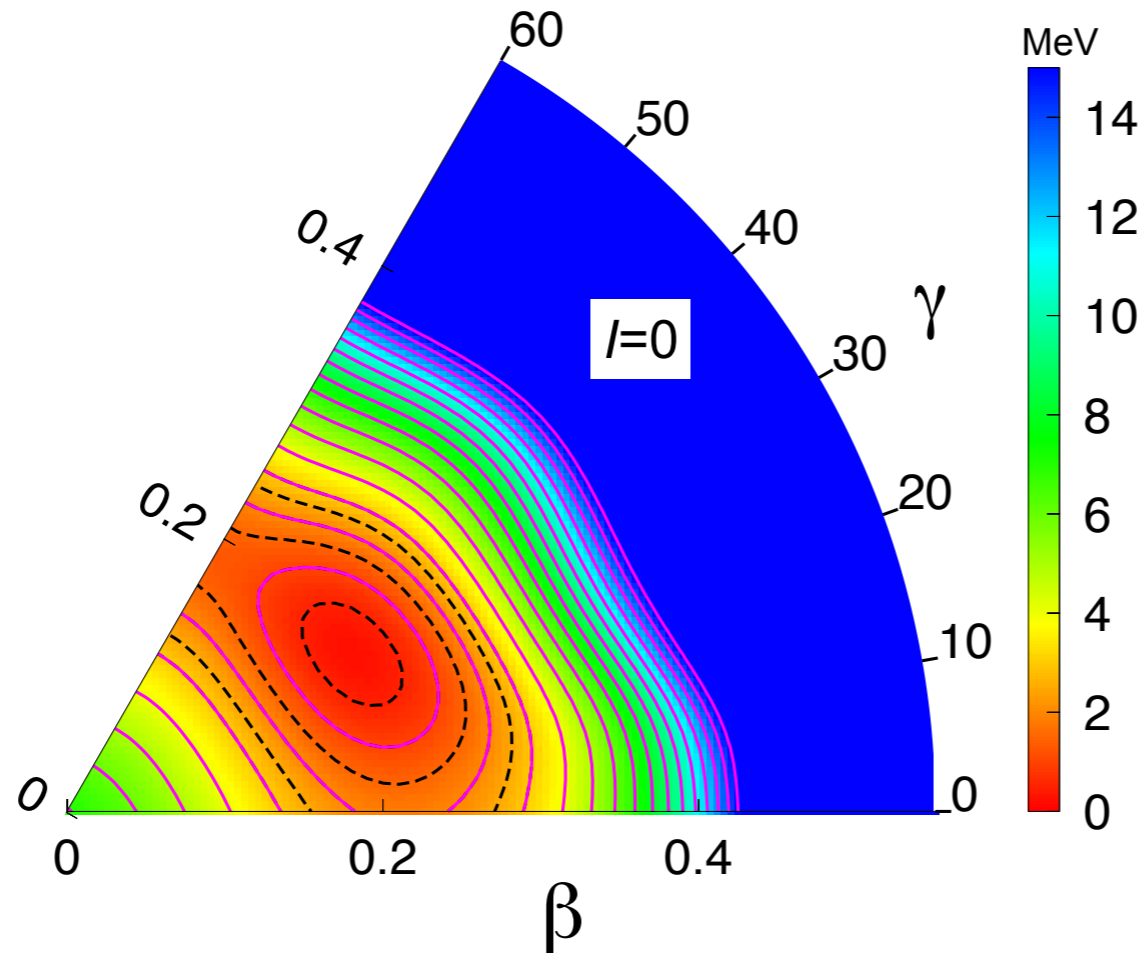
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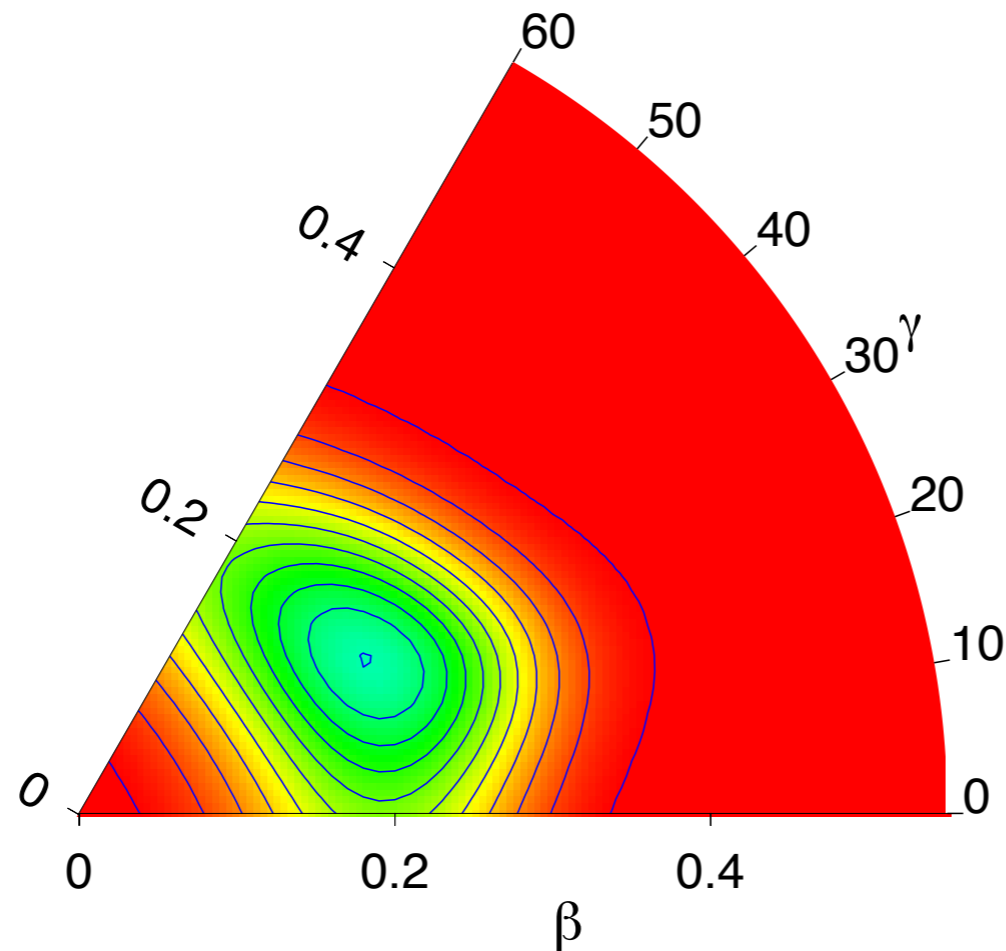
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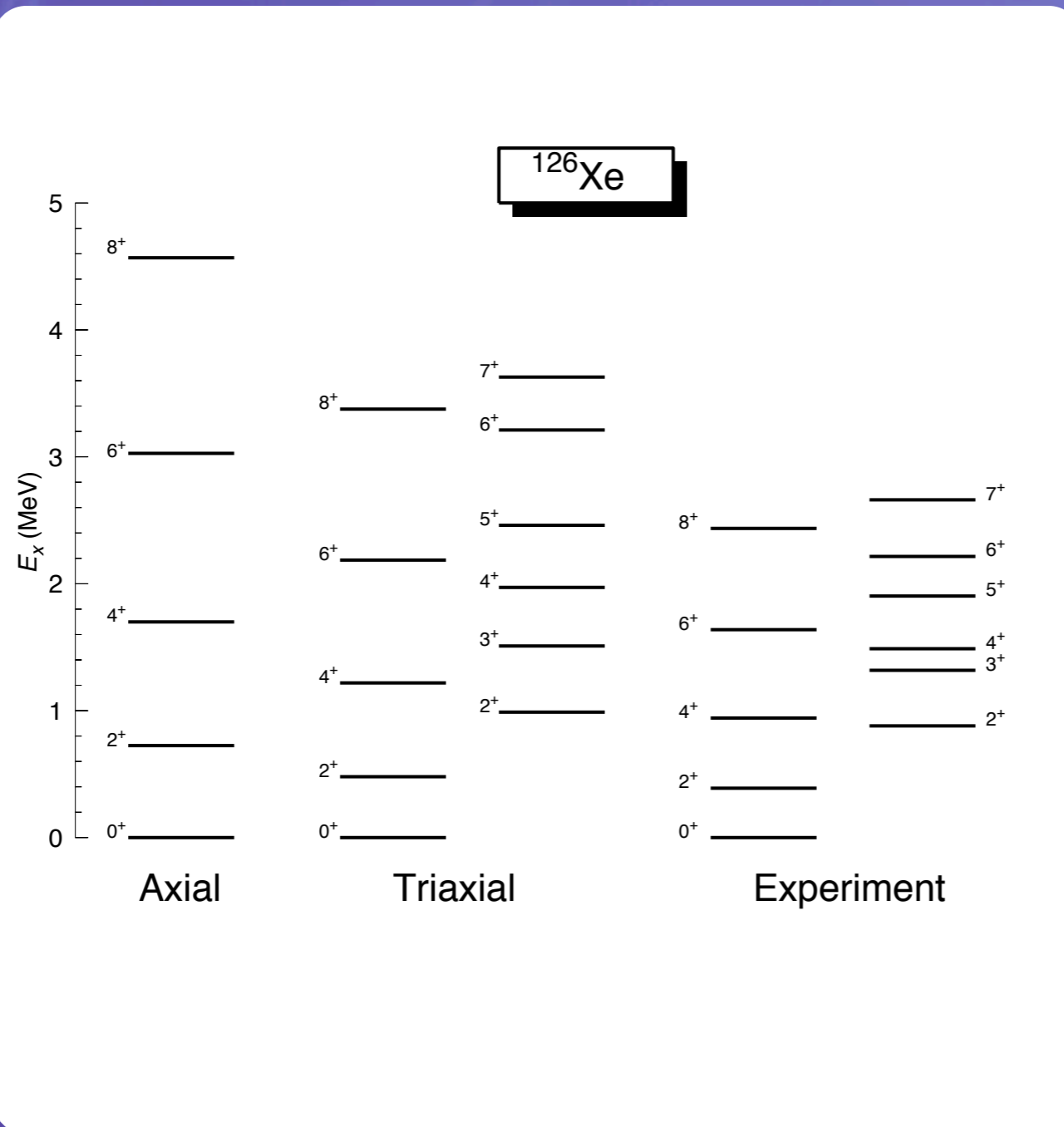
Triaxial calculations ^{126}Xe

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2. Applications

3. Conclusions and outlook



✓ TRIAXIAL calculations

✓ Triaxial calculations are able to describe qualitatively the experimental data

✓ Branching ratios for the B(E2) nicely reproduced.

$I_i \rightarrow I_f$	Exp.	Theory
$2_2^+ \rightarrow 2_1^+$	100.	100.
$2_2^+ \rightarrow 0_1^+$	1.5 ± 0.4	0.001
$3_1^+ \rightarrow 4_1^+$	$35._{-34}^{+10}$	40.48
$3_1^+ \rightarrow 2_2^+$	100.	100.
$3_1^+ \rightarrow 2_1^+$	$2.0_{-1.7}^{+0.6}$	0.000
$4_2^+ \rightarrow 4_1^+$	$76. \pm 22$	80.6
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$5_1^+ \rightarrow 4_1^+$	2.9 ± 0.8	0.02
$6_2^+ \rightarrow 6_1^+$	$34. \pm 15$	27.1
$6_2^+ \rightarrow 4_2^+$	100.	100.
$6_2^+ \rightarrow 4_1^+$	0.49 ± 0.15	0.003
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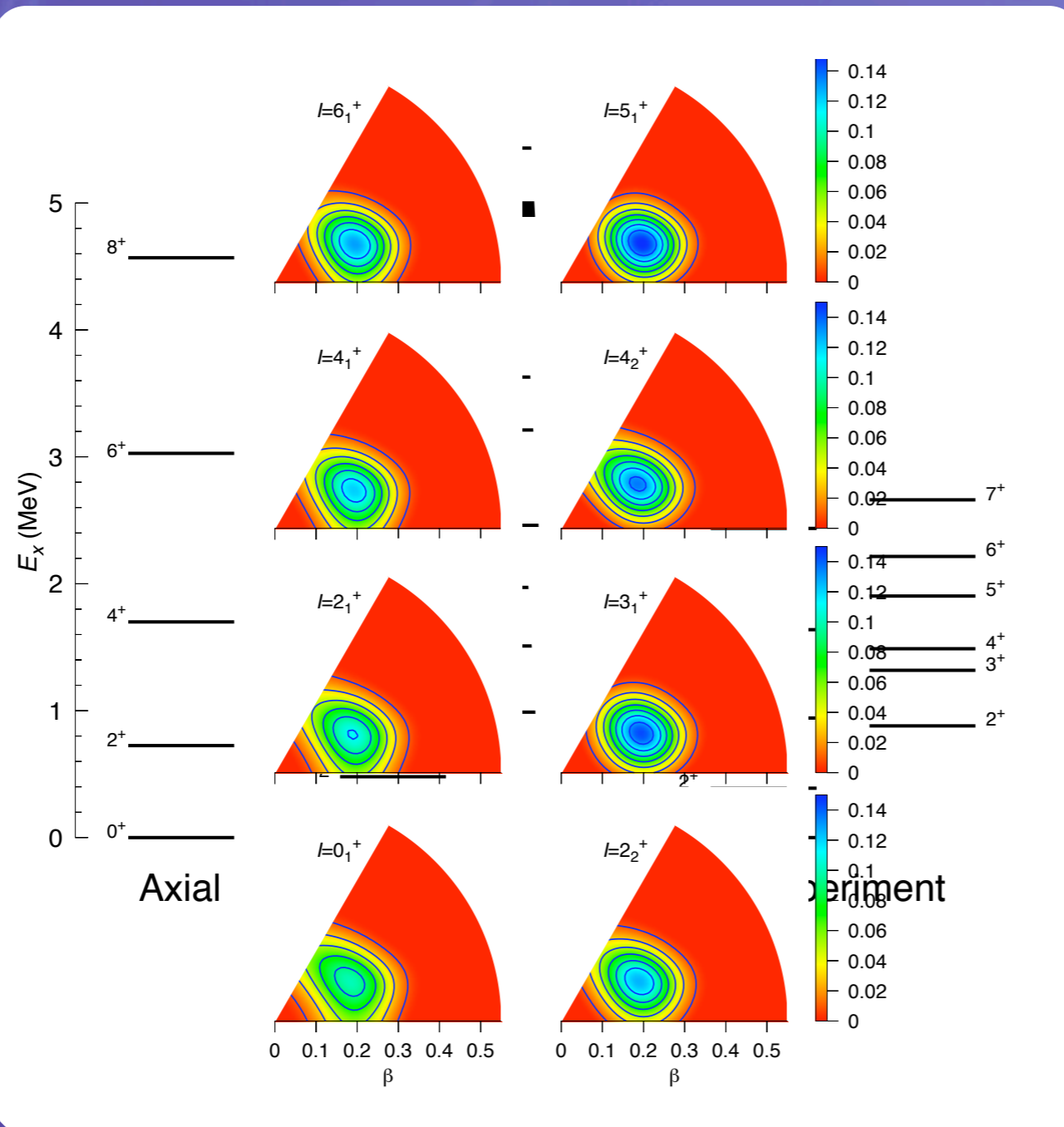
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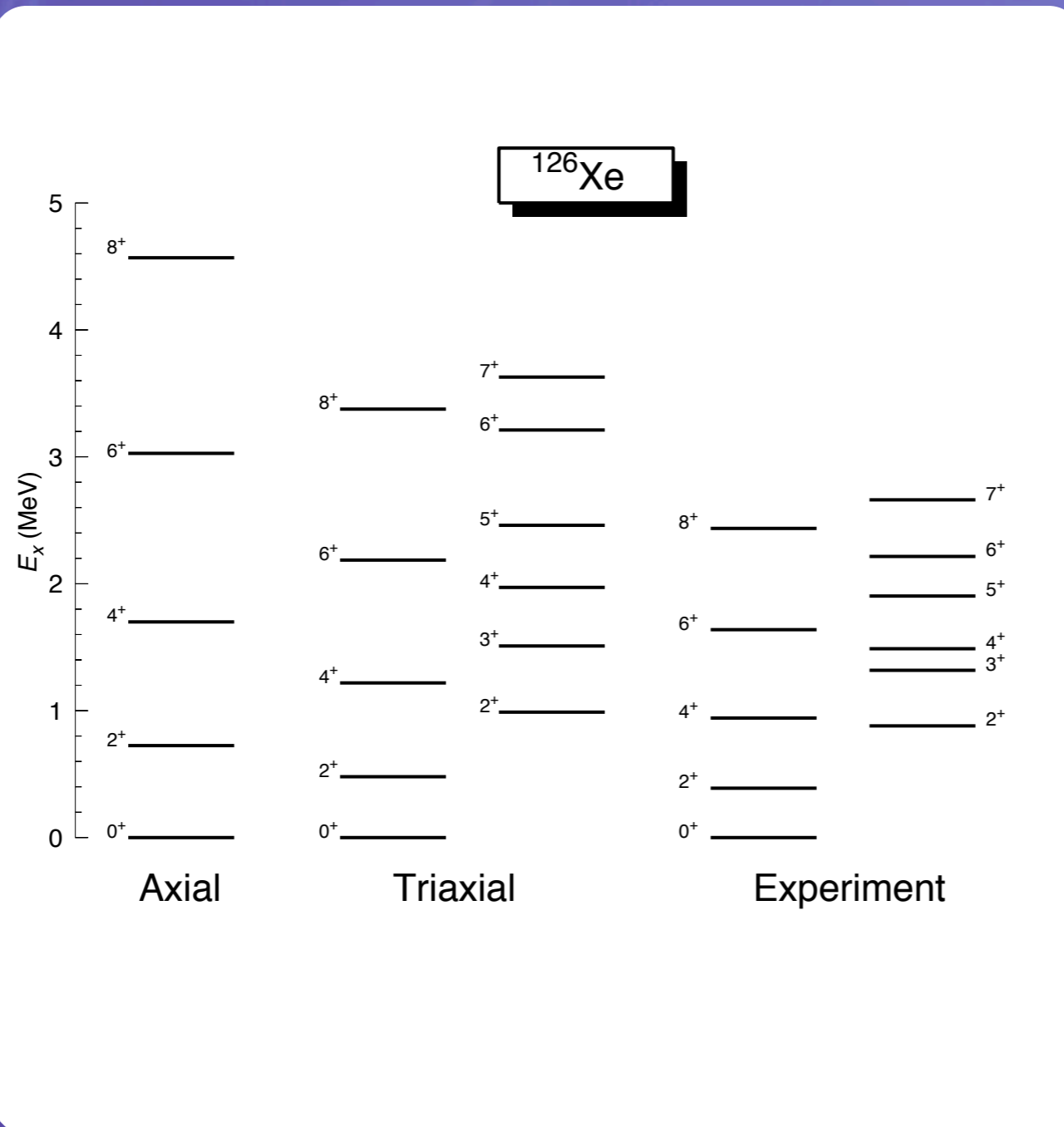
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Triaxial calculations ^{80}Zr

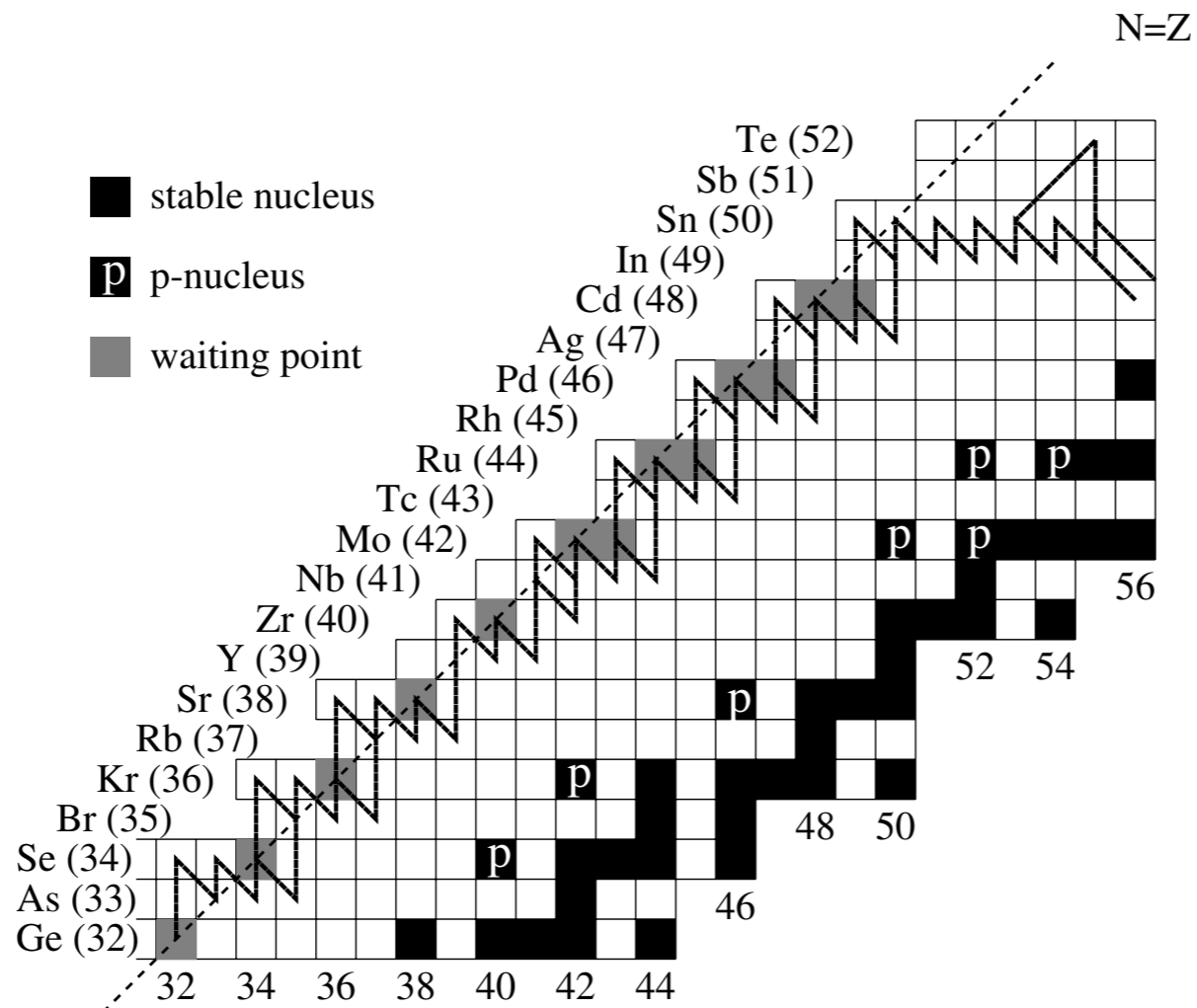
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1. Theoretical framework

2. Applications

3. Conclusions and outlook



H. Schatz *et al.*, Phys. Rep. 294, 167 (1998)

✓ $N=Z$ neutron deficient nuclei are the waiting points for rp-process nucleosynthesis

✓ Small proton capture cross sections and large beta half-lives

✓ Nuclear structure of these nuclei determines the parameters relevant to this process (masses, beta decay half-lives and branching ratios, ...)

Triaxial calculations ^{80}Zr

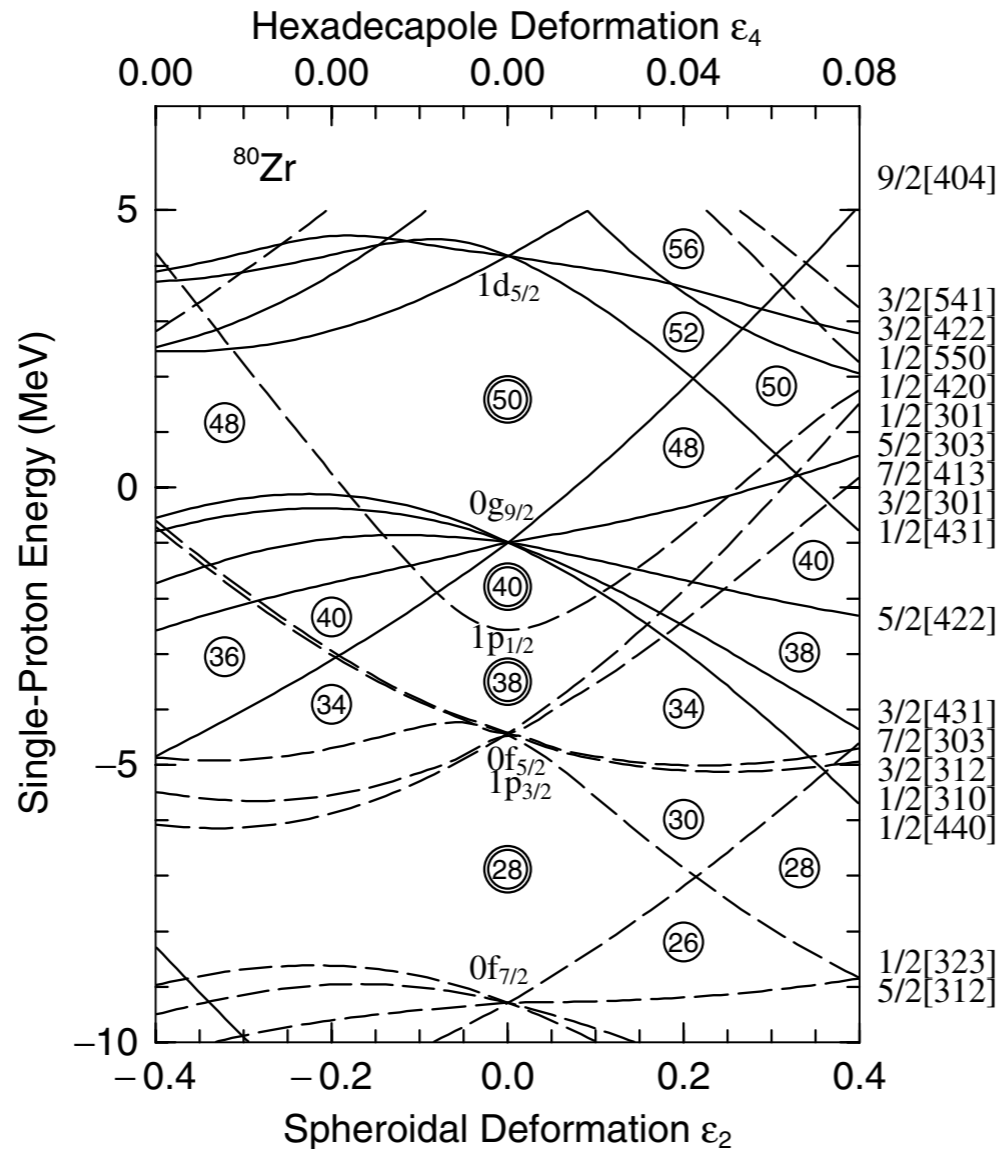
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1. Theoretical framework

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3. Conclusions and outlook



P. Möller *et al.*, At. Data Nucl. Data Tables 66,131 (1997)

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✓ Shape coexistence appears in most of these nuclei

✓ To which extent the appearance of “shape-isomeric” states affects the decay of these nuclei?

Triaxial calculations ^{80}Zr

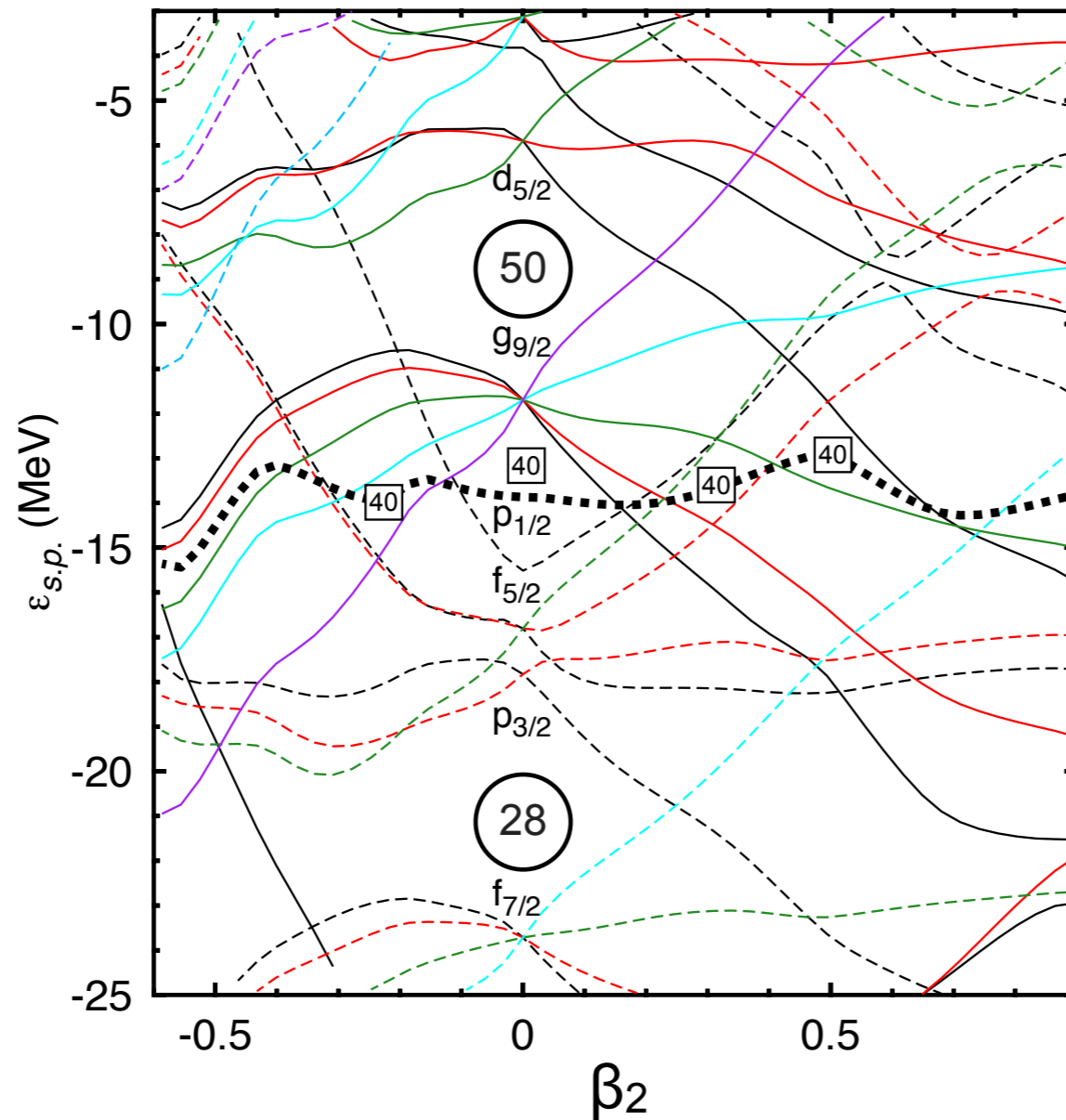
Motivation

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3. Conclusions and outlook



T.R.R. and J.L. Egido, submitted to PLB

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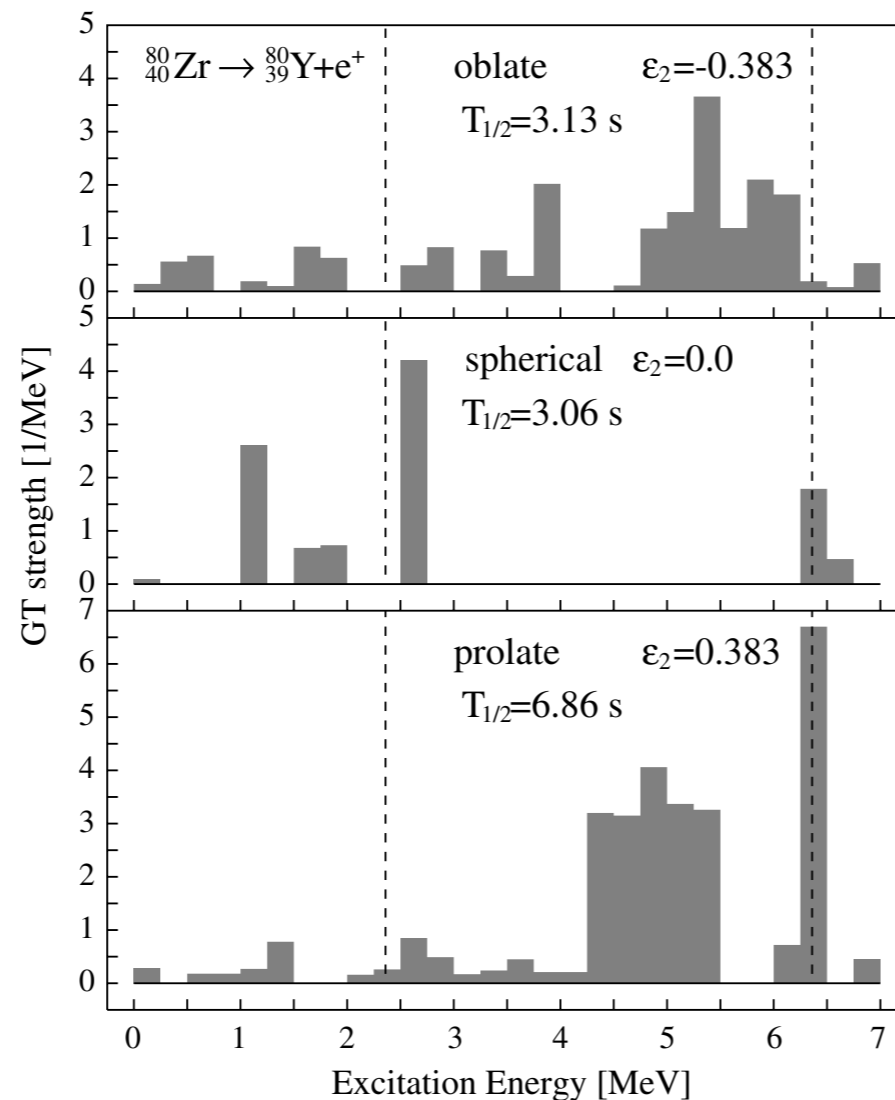
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✓ Shape coexistence appears in most of these nuclei

✓ Different GT strength distribution depending on the shape of the mother nucleus.

Triaxial calculations ^{80}Zr

Motivation

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1. Theoretical framework

2. Applications

3. Conclusions and outlook

There are very few experimental data information of this nucleus:

- Ground state band energies (almost rotational)

(C.J. Lister et al, Phys. Rev. Lett. 59, 1270 (1987) and S. M. Fischer et al., Phys. Rev. Lett. 87, 132501 (2001))

- β^+ half-life (4.1 s)

(J.J. Ressler et al., Phys. Rev. Lett. 84, 2104 (2000))

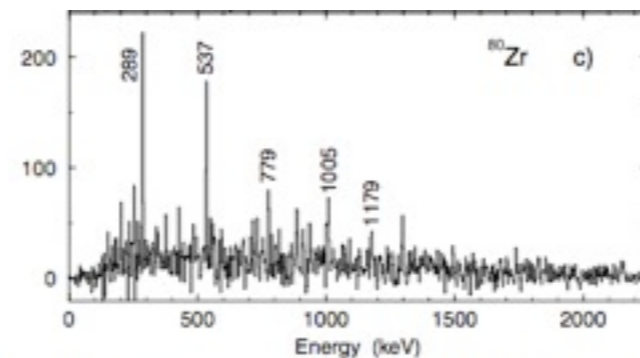


FIG. 1. Gamma-ray spectra of the ground state cascades in $N = Z$ nuclei: (a) ^{72}Kr (2α -gated $\gamma\gamma\gamma$ data, sum of $\gamma\gamma$ gates with one transition below and one transition above $J = 14$, production cross section $\sim 120 \mu\text{b}$); (b) ^{76}Sr ($A/Q = 76/24$ and dE/dx gated $\gamma\gamma$ data, sum of γ gates for transitions below $J = 10$, $\sim 20 \mu\text{b}$); (c) ^{80}Zr ($A/Q = 80/25$ and dE/dx gated γ singles, $\sim 10 \mu\text{b}$).

S. M. Fischer et al., Phys. Rev. Lett. 87, 132501 (2001)

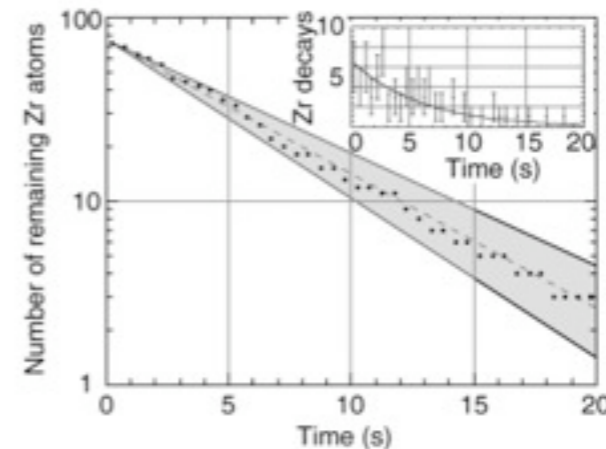


FIG. 4. Number of remaining Zr atoms as a function of time. The straight line represents a 4.1-s decay line, while the shaded areas show the range of error—4.9 s as an upper limit and 3.5 s as a lower. The inset shows the background-corrected counts per 0.5 s.

J. J. Ressler et al., Phys. Rev. Lett. 84, 2104 (2000)

Triaxial calculations ^{80}Zr

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3. Conclusions and outlook

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 (J.J. Ressler et al., Phys. Rev. Lett. 84, 2104 (2000))

^{80}Zr :



FIG. 1. Gamma-ray spectra of the ground state cascades in $N = Z$ nuclei: (a) ^{78}Kr (2α -gated $\gamma\gamma\gamma$ data, sum of $\gamma\gamma$ gates with one transition below and one transition above $J = 14$, production cross section $\sim 120 \mu\text{b}$); (b) ^{78}Sr ($A/Q = 76/24$ and dE/dx gated $\gamma\gamma$ data, sum of γ gates for transitions below $J = 10$, $\sim 20 \mu\text{b}$); (c) ^{80}Zr ($A/Q = 80/25$ and dE/dx gated γ singles, $\sim 10 \mu\text{b}$).

S. M. Fischer et al., Phys. Rev. Lett. 87, 132501 (2001)

✓ RELEVANT IN NUCLEOSYNTHESIS
✓ RELEVANT IN NUCLEAR STRUCTURE

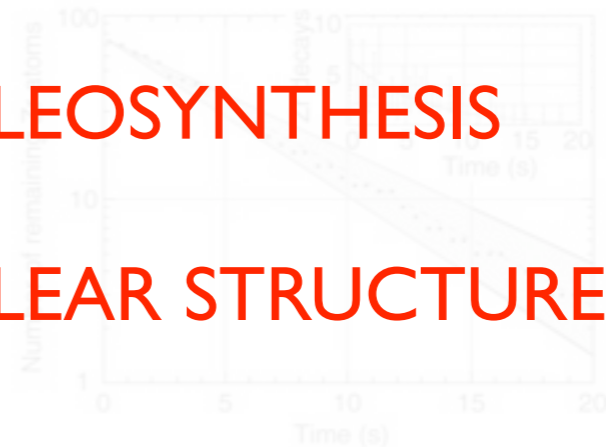


FIG. 4. Number of remaining Zr atoms as a function of time. The straight line represents a 4.1-s decay line, while the shaded areas show the range of error — 4.9 s as an upper limit and 3.5 s as a lower. The inset shows the background-corrected counts per 0.5 s.

J. J. Ressler et al., Phys. Rev. Lett. 84, 2104 (2000)

Triaxial calculations ^{80}Zr

First step: Particle Number Projection (before the variation) of HFB-type wave functions.

$$\delta E^{N,Z} [\bar{\Phi}(\beta, \gamma)] \Big|_{\bar{\Phi}=\Phi} = 0 \quad E^{N,Z}[\Phi] = \frac{\langle \Phi | \hat{H}_{2b} \hat{P}^N \hat{P}^Z | \Phi \rangle}{\langle \Phi | \hat{P}^N \hat{P}^Z | \Phi \rangle} + \varepsilon_{DD}^{N,Z}(\Phi) - \lambda_{q_{20}} \langle \Phi | \hat{Q}_{20} | \Phi \rangle - \lambda_{q_{22}} \langle \Phi | \hat{Q}_{22} | \Phi \rangle$$

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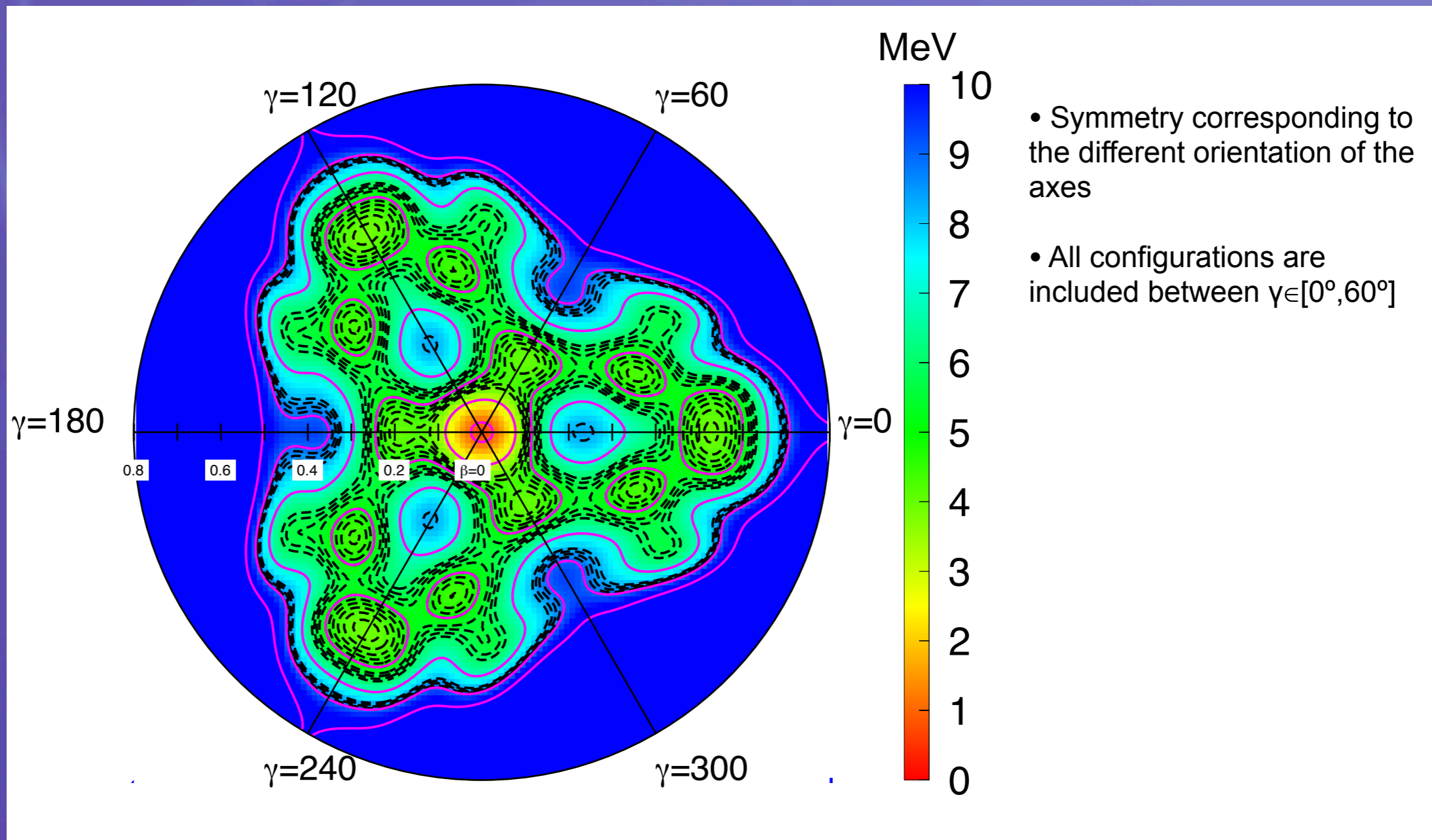
$$\delta E^{N,Z} [\bar{\Phi}(\beta, \gamma)] \Big|_{\bar{\Phi}=\Phi} = 0 \quad E^{N,Z}[\Phi] = \frac{\langle \Phi | \hat{H}_{2b} \hat{P}^N \hat{P}^Z | \Phi \rangle}{\langle \Phi | \hat{P}^N \hat{P}^Z | \Phi \rangle} + \varepsilon_{DD}^{N,Z}(\Phi) - \lambda_{q_{20}} \langle \Phi | \hat{Q}_{20} | \Phi \rangle - \lambda_{q_{22}} \langle \Phi | \hat{Q}_{22} | \Phi \rangle$$

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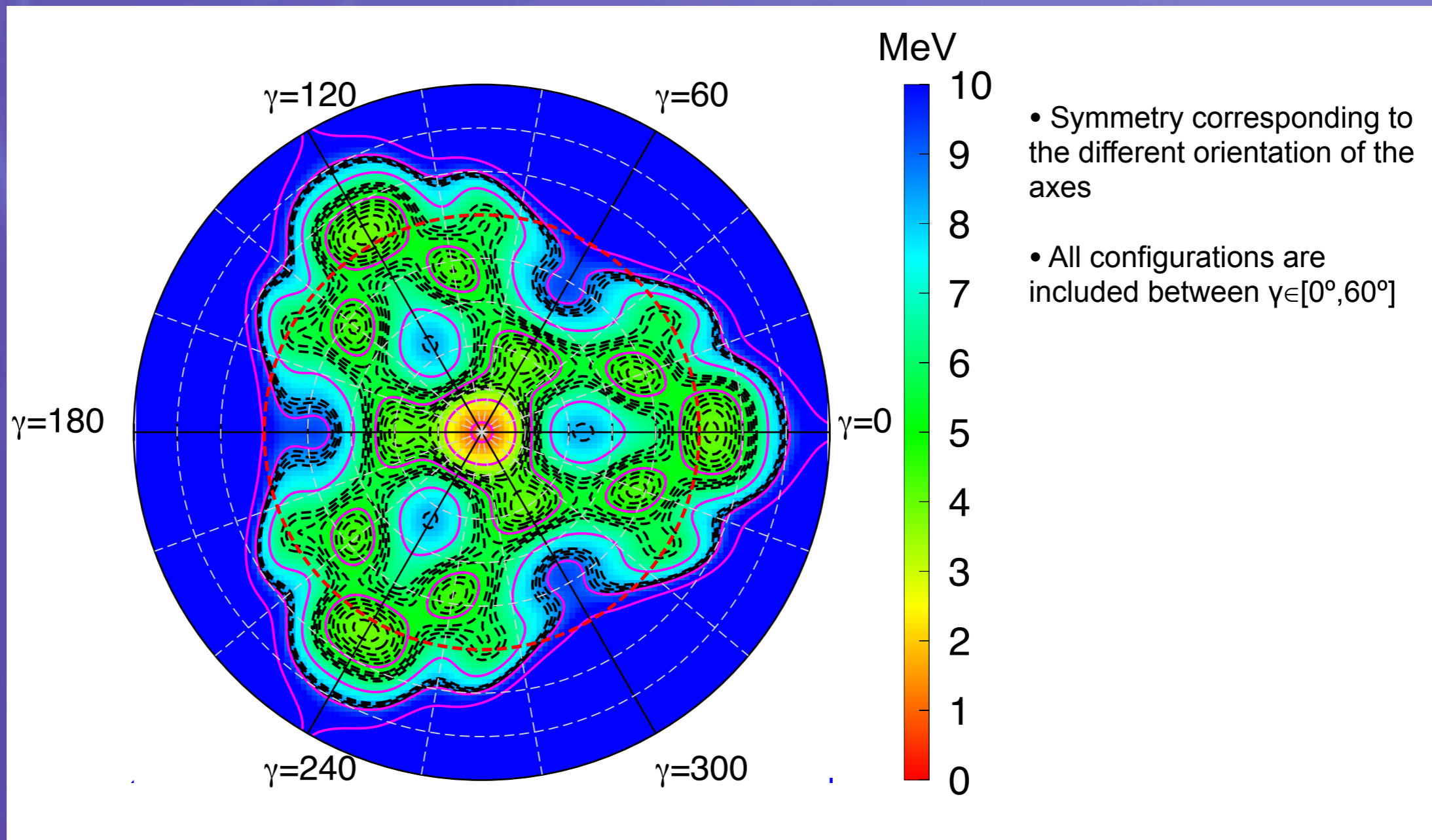
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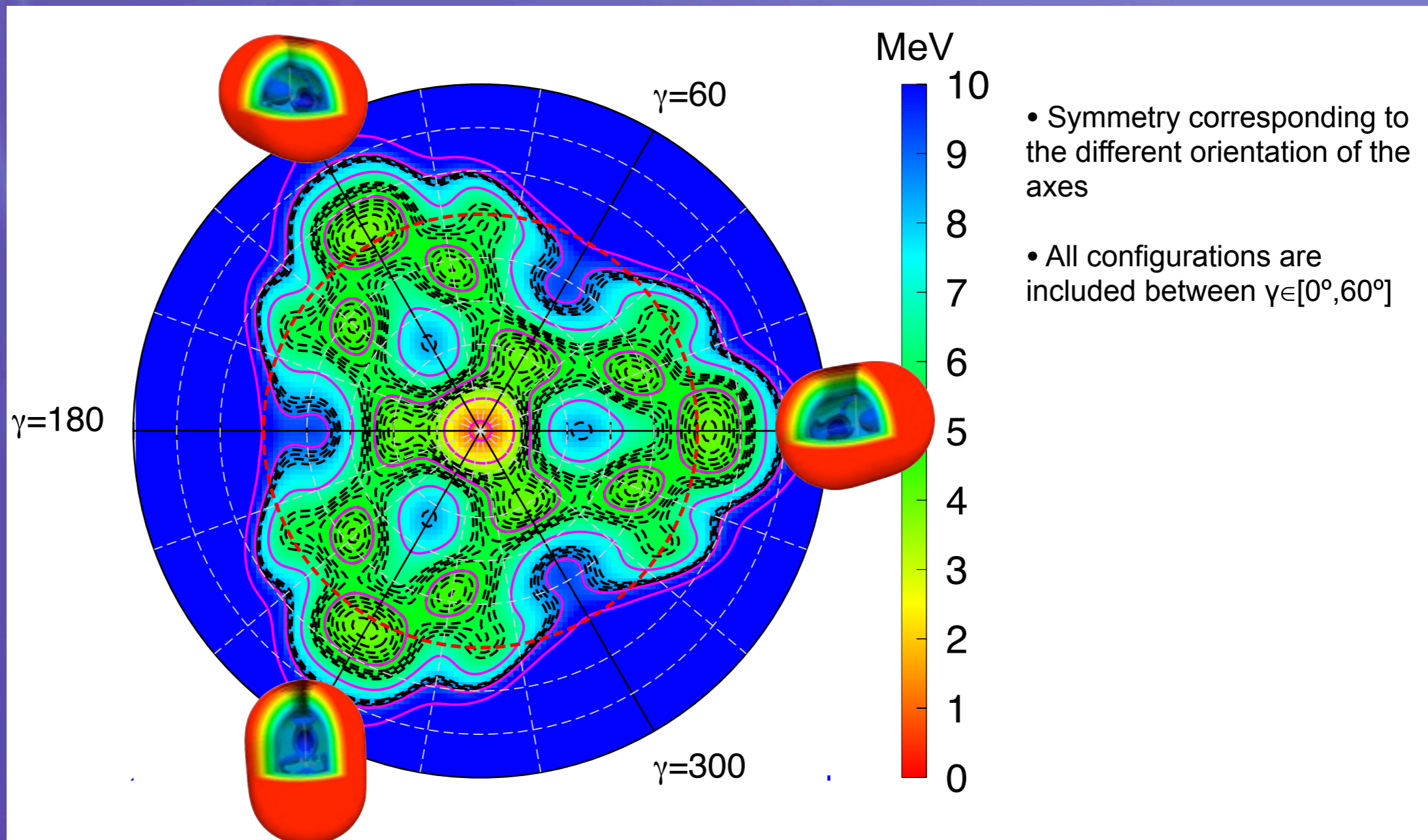
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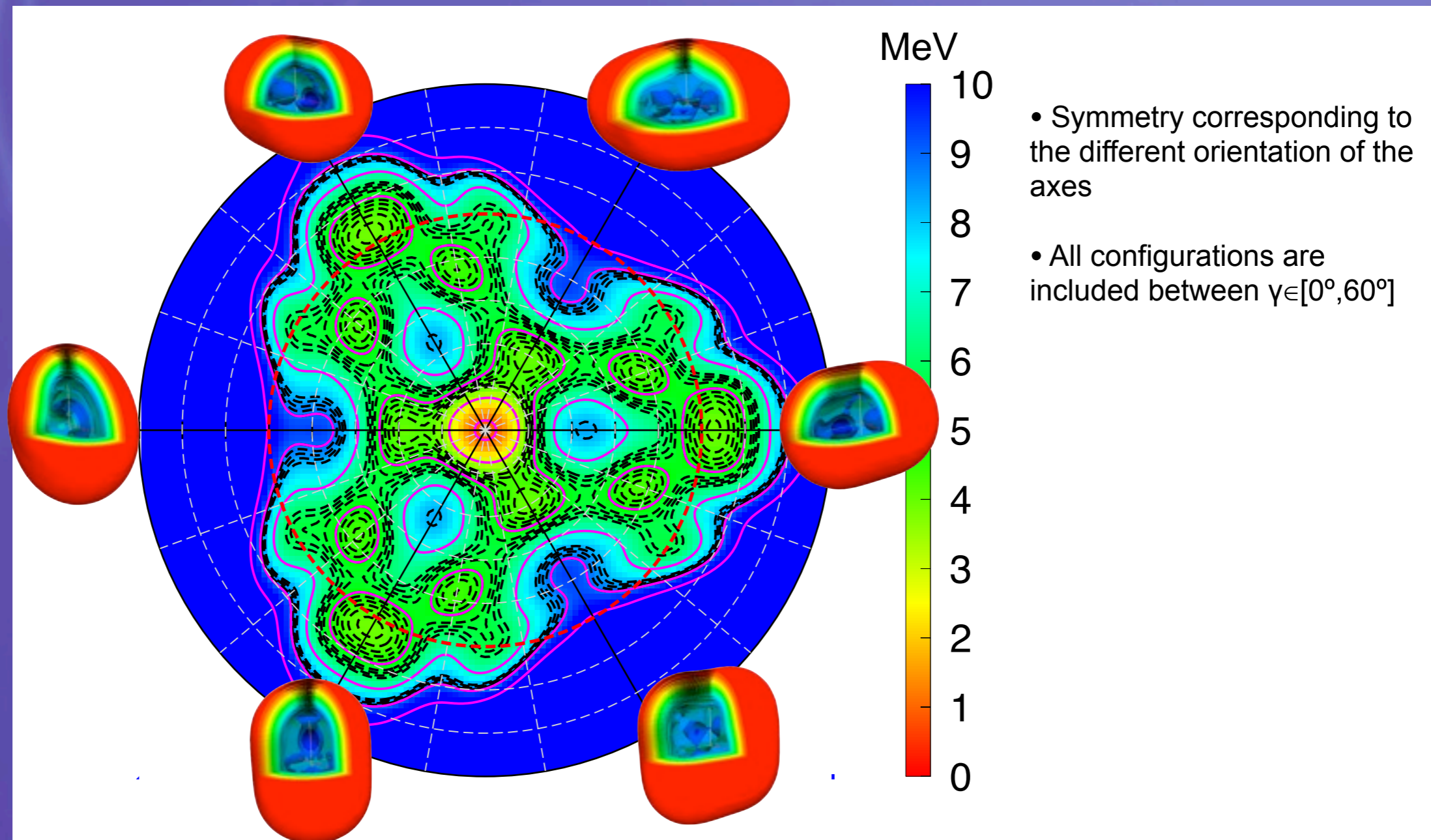
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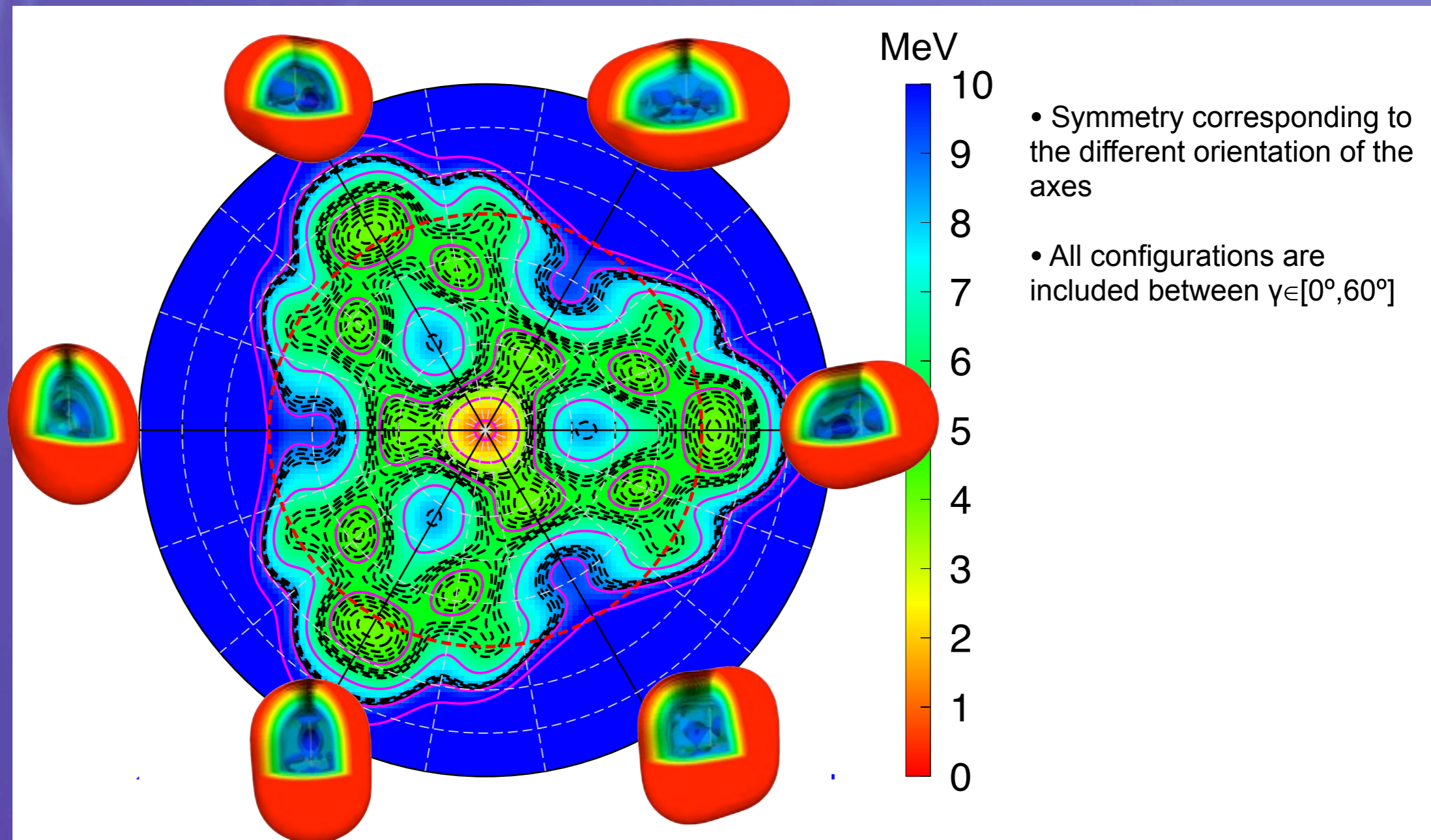
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Triaxial calculations ^{80}Zr (83 states, 9 shells)

First step: Particle Number Projection (before the variation) of HFB-type wave functions.

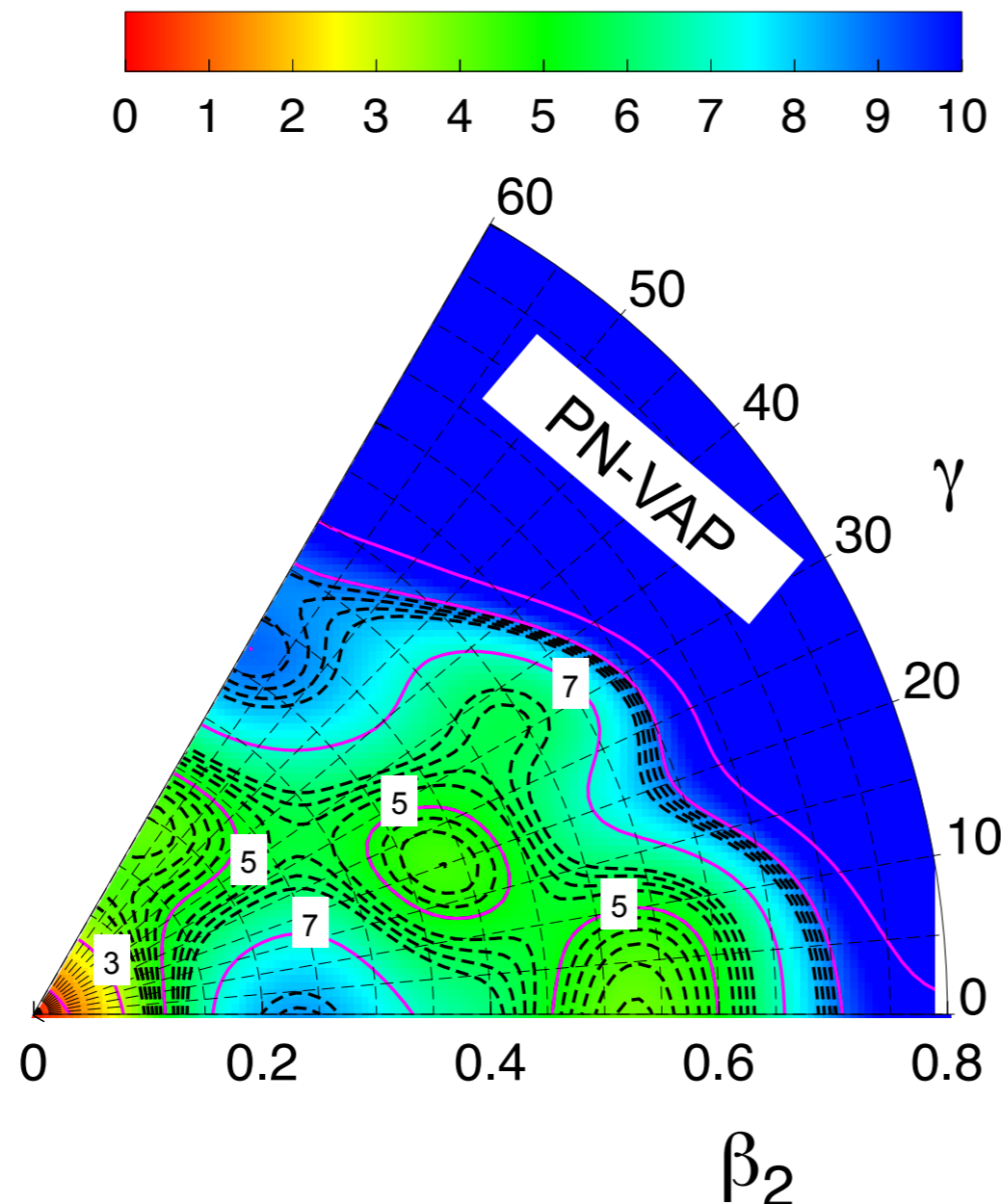
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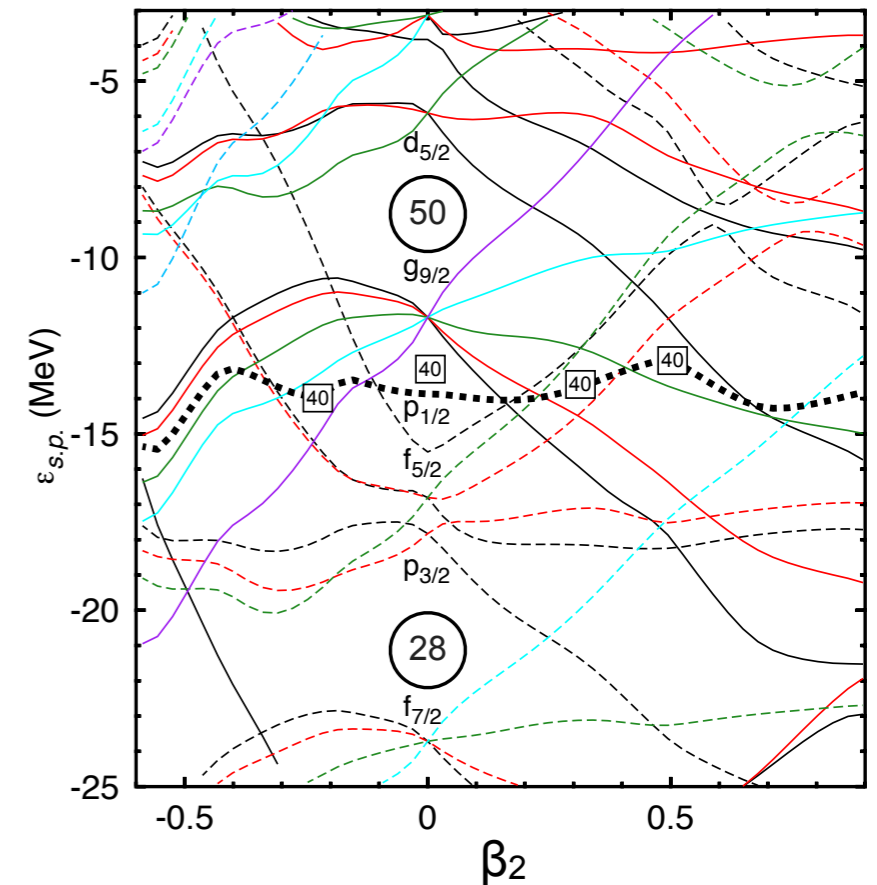
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- Up to five minima in the potential energy surface.
- Absolute minimum corresponds to spherical configuration ($N=40$ spherical gap)
- Other minima related to the filling in and out of $g_{9/2}$, $p_{1/2}$, $f_{5/2}$ and $d_{5/2}$ orbits.



Triaxial calculations ^{80}Zr (83 states, 9 shells)

Second step: Simultaneous Particle Number and Angular Momentum Projection

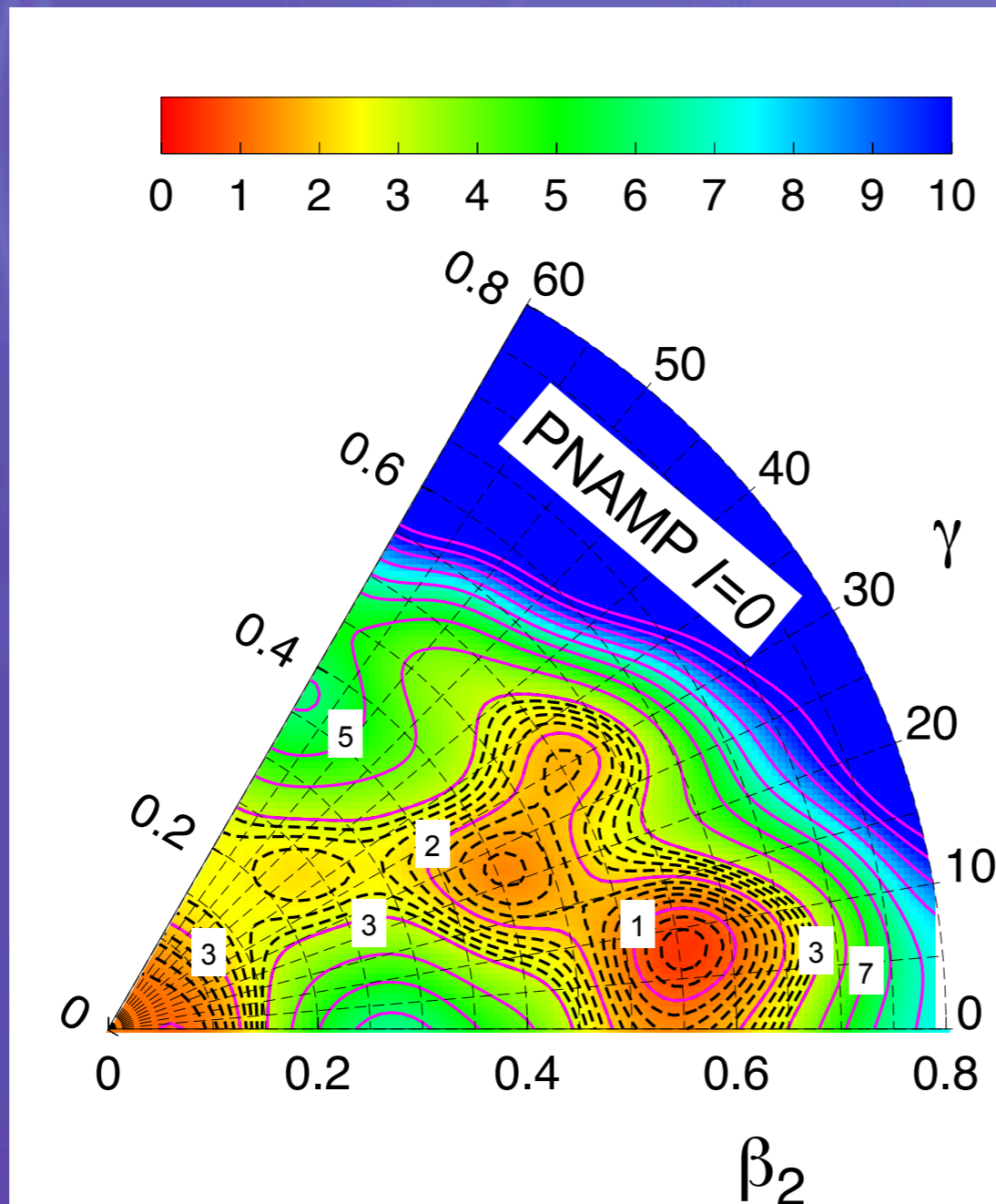
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$$|IMK; NZ; \beta\gamma\rangle = \frac{2I+1}{8\pi^2} \int \mathcal{D}_{MK}^{I*}(\Omega) \hat{R}(\Omega) \hat{P}^N \hat{P}^Z |\Phi(\beta, \gamma)\rangle d\Omega \rightarrow |IM; NZ; \beta\gamma\rangle = \sum_K g_K^{IM; NZ; \beta\gamma} |IMK; NZ; \beta\gamma\rangle$$



- Five minima are closer in energy whenever the rotational invariance is restored.
- Absolute minima corresponds to deformed configuration $\beta \sim 0.55$
- Barriers between the minima are less than 1 MeV. Mixing?

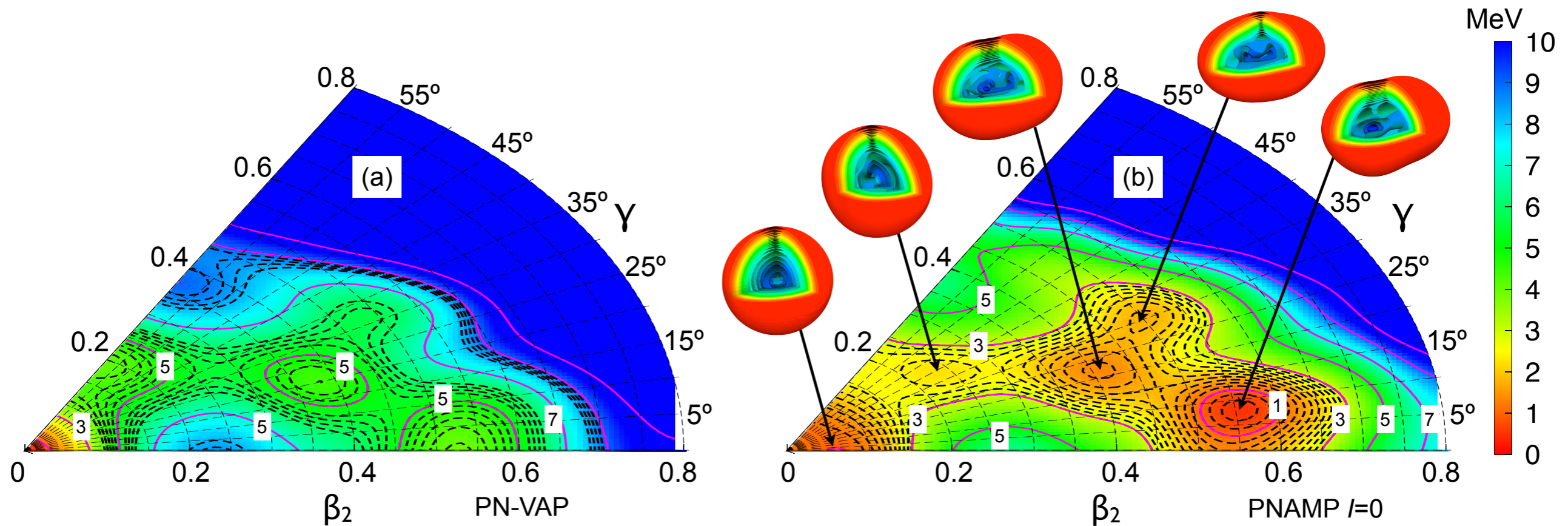
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Relevance of angular momentum projection

(Similar feature as in ^{32}Mg , see R. Rodríguez-Guzmán et al., Nucl. Phys. A 709, 201 (2002))



Triaxial calculations ^{80}Zr (83 states, 9 shells)

Final step: Configuration mixing within the framework of the **Generator Coordinate Method (GCM)**. K and deformation mixing

$$|IM; NZ\sigma\rangle = \sum_{K\beta\gamma} f_{K\beta\gamma}^{I;NZ,\sigma} |IMK; NZ; \beta\gamma\rangle$$

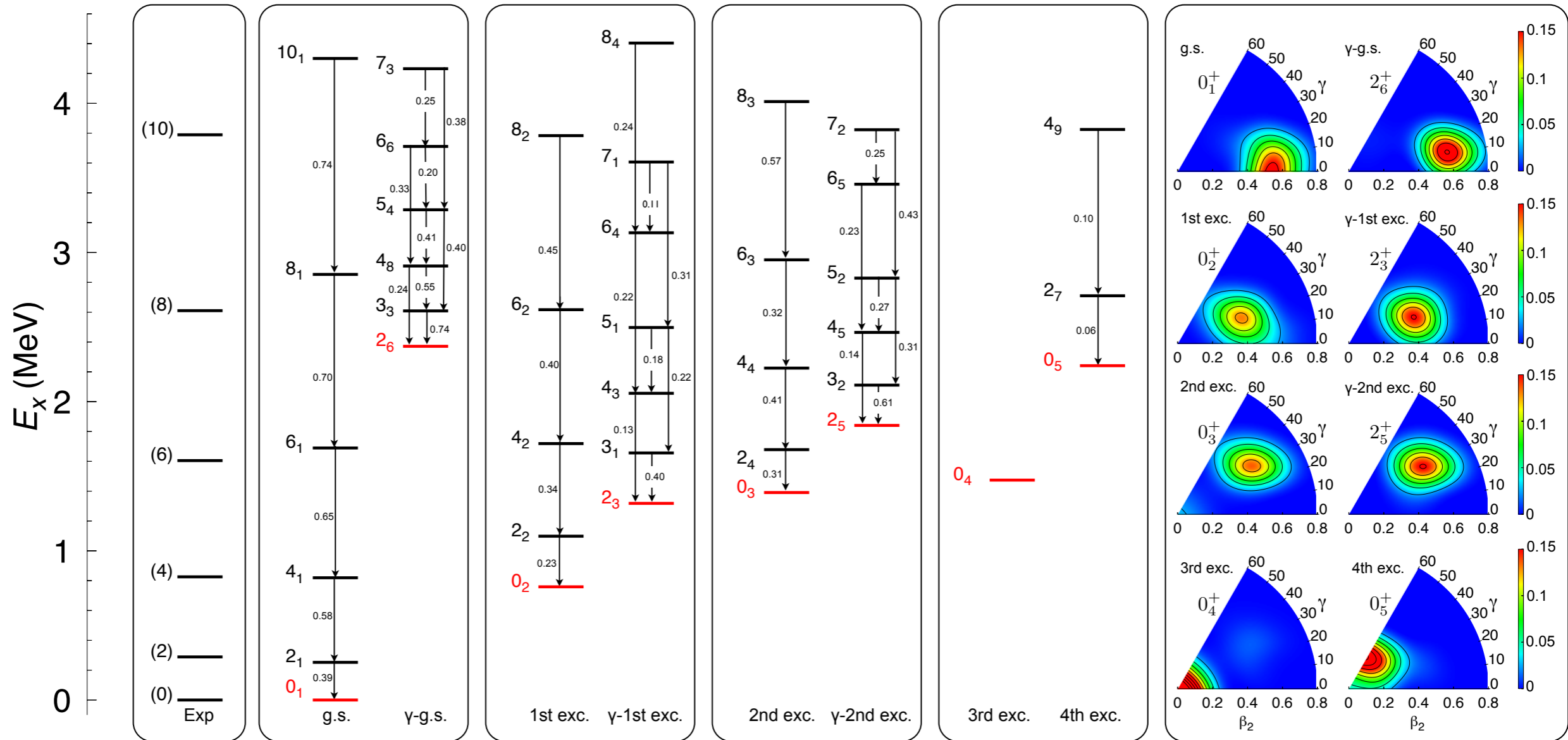
T.R.R and J.L. Egido, Phys. Lett. B submitted.

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Triaxial calculations ^{80}Zr (83 states, 9 shells)

Final step: Configuration mixing within the framework of the **Generator Coordinate Method (GCM)**. **K** and deformation mixing

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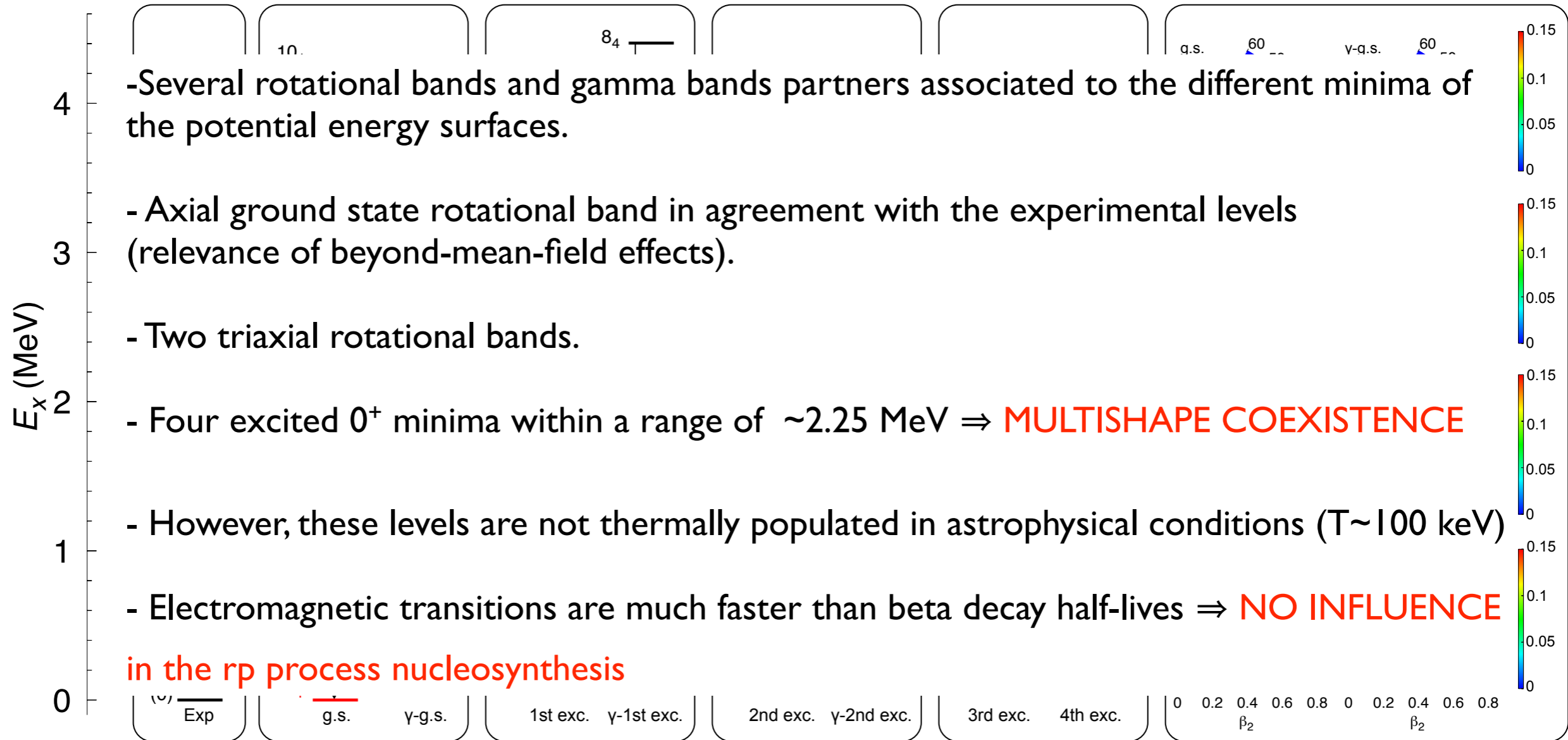
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Conclusions and outlook

- ✓ Current phenomenological Energy Density Functional methods including triaxial shapes provide a very good description and physical insight of many phenomena in nuclei along the whole nuclear chart.
- ✓ It is a competitive alternative and/or complement to shell model calculations.
- ✓ Computational time is still a problem.
- ✓ Some improvements have to be performed yet:
 - Projection of cranking states (time reversal symmetry breaking).
 - Include quasiparticle states (blocking) to describe single particle excitations, odd nuclei and β transitions.
 - Include parity and isospin breaking.
 - Fit the interaction with beyond mean field calculations (relevant to mass table calculations) and/or develop new non-empirical functionals based on QCD (relevant to mass table calculations).

Acknowledgments

J. L. Egido (Universidad Autónoma de Madrid)

G. Martínez-Pinedo (GSI-Darmstadt)

Andrea Jungclaus (IEMM-CSIC Madrid)

C. Domingo-Pardo (IFIC-CSIC Valencia)

Nuria López Vaguero (Universidad Autónoma de Madrid)

THANK YOU FOR YOUR ATTENTION!!!